

Cite this: *Catal. Sci. Technol.*, 2026, 16, 3137Received 14th February 2026,  
Accepted 21st March 2026

DOI: 10.1039/d6cy00192k

rsc.li/catalysis

Activation of  $N_2$  with metallic lithium is a key step of the lithium redox-mediated nitrogen reduction reaction (Li-NRR) – a process that can enable practical ammonia electrosynthesis under mild conditions. Our present DFT study revealed that  $N_2$  capture and  $N\equiv N$  bond cleavage is substantially enhanced on and within  $Li^0$  nanoparticles as compared to bulk slabs. This behaviour is not limited to single-molecule adsorption but is also observed during the successive capture of additional  $N_2$  molecules. The results presented herein might assist in the design of higher performance Li-NRR systems.

The nitrogen reduction reaction (NRR) is one of the most extensively studied processes of the last decade in the field of applied electrocatalysis, owing to the importance of ammonia as a fertiliser feedstock and a future energy vector.<sup>1</sup> However, research on the electrochemical  $N_2$  activation under mild conditions has demonstrated that the direct electrocatalytic reduction of  $N_2$  is challenging to achieve at production rates that are viable on a practical scale.<sup>2</sup> In contrast, the lithium redox-mediated NRR (Li-NRR) has been unambiguously proven to provide practical levels of genuine conversion of  $N_2$  into  $NH_3$  at ambient temperature and mild pressure  $<20$  bar.<sup>3–7</sup>

In the Li-NRR, the electrochemically generated metallic lithium mediates  $N_2$  cleavage with the formation of lithium nitride *via* a dissociative mechanism, which demonstrates a range of distinctive features<sup>8–10</sup> as compared to the mechanisms of  $N_2$  activation proposed for the direct electrocatalytic NRR.<sup>11–18</sup> Next, the nitrogen atoms incorporated into the  $Li_3N$  lattice are converted into  $NH_3$  through successive protonation by a suitable proton carrier.<sup>19,20</sup> A similar mechanism has recently been reported for calcium<sup>21</sup> and magnesium.<sup>22</sup>

## Enhanced $N_2$ capture and cleavage promoted by lithium nanoparticles

José Miguel Doña-Rodríguez,<sup>a</sup> Raúl Quesada-Cabrera,<sup>a</sup>  
Douglas R. MacFarlane,<sup>b</sup> Alexandr N. Simonov<sup>\*b</sup> and Luis Miguel Azofra<sup>\*a</sup>

At a theoretical level, this behaviour has been modelled in metal slabs, in the first instance for the  $N_2$  cleavage mechanism (Fig. 1),<sup>8–10</sup> but also in subsequent protonation steps.<sup>23,24</sup> These previous studies established that: (i) end-on adsorption of  $N_2$  on  $Li^0$  is non-spontaneous; (ii) side-on adsorption, by contrast, is spontaneous, overcoming the physisorption regime, as  $N_2$  becomes inserted into the surface structure of the  $Li^0$  slabs, with binding free energies characteristic of chemisorption or capture events; (iii)  $N_2$  can diffuse into the inner layers of the  $Li^0$  bulk; and (iv) the activation free-energy barriers for  $N_2$  cleavage on the  $Li^0$  surface or within the  $Li^0$  bulk are accessible at room temperature.

However, metal slabs represent restricted structures, in contrast to nanoparticles which exhibit diverse structural features, like different facets as well as edges and corners. These features are likely relevant to the real state of the electroplated  $Li^0$  intermediate during the Li-NRR, since electrodeposition of lithium metal typically produces highly heterogeneous nanostructured aggregates rather than well-

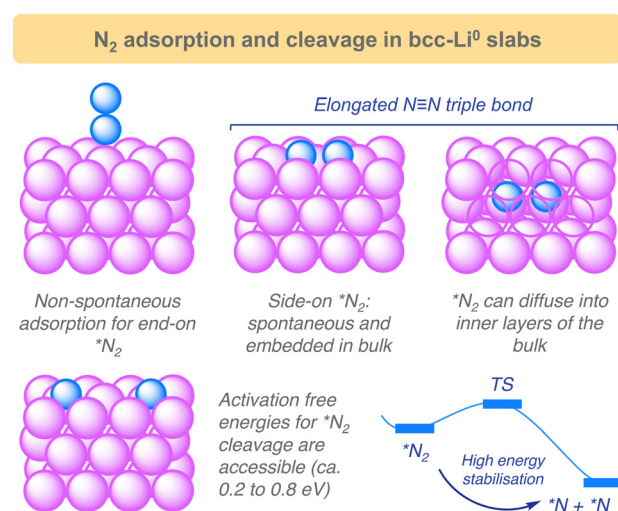


Fig. 1 Summary of the main chemical events observed during the  $N_2$  capture and cleavage on and within the  $bcc-Li^0$  slabs.<sup>8,9</sup>

<sup>a</sup> Instituto de Estudios Ambientales y Recursos Naturales (i-UNAT), Universidad de Las Palmas de Gran Canaria (ULPGC), Campus de Tafira, 35017 Las Palmas de Gran Canaria, Spain. E-mail: luismiguel.azofra@ulpgc.es

<sup>b</sup> School of Chemistry, Monash University, Clayton, Victoria 3800, Australia. E-mail: alexandr.simonov@monash.edu

defined crystalline surfaces.<sup>25–27</sup> To investigate this, we conducted a density functional theory (DFT) study (pseudopotential plane-wave method, RPBE<sup>28</sup> + D3<sup>29,30</sup> level of theory, VASP,<sup>31–34</sup> see SI, for computational details) of a series of structures that were modelled by spherically projecting the body-centred cubic (bcc) Li<sup>0</sup> unit cell with radii ranging from 1 to 6 Å. This resulted in a series of initial models: the Li<sub>1</sub> atomic dot (1 Å radius), the Li<sub>2</sub> (2 Å), Li<sub>6</sub> (3 Å), Li<sub>15</sub> (4 Å) and Li<sub>27</sub> (5 Å) sub-nanoclusters, and the Li<sub>51</sub> (6 Å) nanoparticle (Fig. 2).

In our previous study on the interaction of N<sub>2</sub> on and within the bcc-Li<sup>0</sup> slabs, we computed binding free energies for the (001) and (011) flat surfaces, amongst others.<sup>8,9</sup> We found that end-on N<sub>2</sub> adsorption exhibited non-spontaneous values of 0.22 eV in both cases, whereas side-on adsorption was spontaneous, with values of –0.46 and –0.66 eV, respectively. The computational protocol used in that study differs from the one employed herein in the calculation of vibrational frequencies only, where atomic positions of the metal were frozen in the previously published work.<sup>8,9</sup> While slabs can be considered constrained structures due to their periodicity, this assumption does not hold for nanostructures. Consequently, when recalculating the side-on N<sub>2</sub> binding on the (001) and (011) facets without structural constraints in the frequency calculation, the values changed slightly to –0.35 and –0.55 eV,

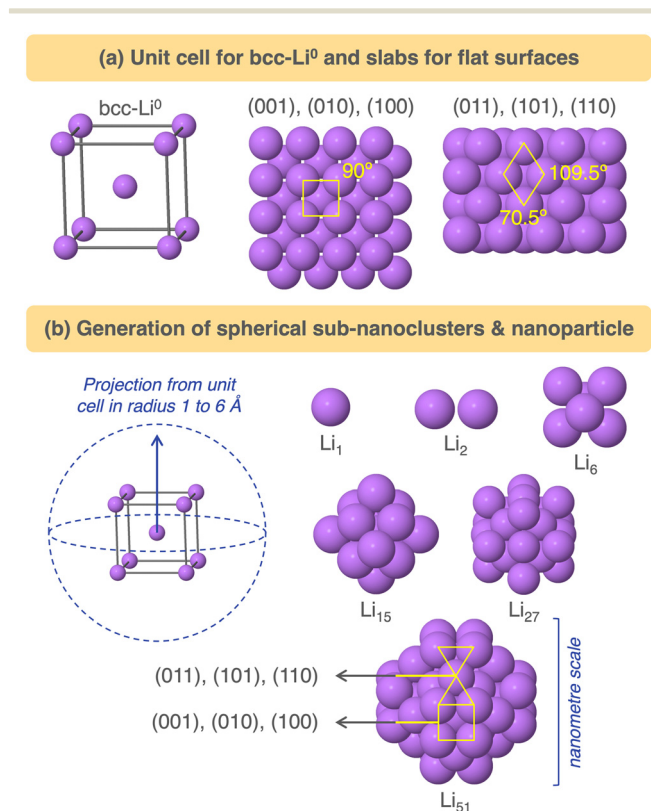
respectively. By applying a consistent methodology, comparison of the N<sub>2</sub>-adsorption results computed here for our sub-nanoclusters and nanoparticle models with those for metal slabs was therefore straightforward.

Interestingly, both Li<sub>1</sub> and Li<sub>2</sub> – the simplest models examined herein – exhibited highly non-spontaneous N<sub>2</sub> binding free energies ( $\Delta G_b$ ). For instance, the linear end-on minima (labelled [1] in both cases in Fig. 3) showed  $\Delta G_b$  values of 0.36 and 0.34 eV, respectively, along with minimal N≡N bond elongations of 0.03 and 0.01 Å, compared with the nitrogen–nitrogen distance of 1.12 Å for N<sub>2</sub>(g) obtained using this computational protocol. These energy values were similar or even less favourable than those others considered herein, specifically minimum [2] for Li<sub>1</sub> of side-on type (0.54 eV), and minima [2] and [3] for Li<sub>2</sub> of angular end-on (0.31 eV) and side-on (0.55 eV) types.

In principle, these results qualitatively follow the trends established before for the Li<sup>0</sup> slabs, where the side-on adsorption was energetically more favourable than end-on.<sup>8,9</sup> However, it was somewhat surprising to find that the interactions between N<sub>2</sub> and these Li<sup>0</sup> structures were so endergonic, given that both Li<sub>1</sub> and Li<sub>2</sub> represent single-site models. Indeed, single-site models have demonstrated superior substrate-interaction capabilities for many other systems. For example, this was shown for the end-on interaction of N<sub>2</sub> with iron slabs<sup>35</sup> and with a single iron atom,<sup>36</sup> where the binding energy in the latter case increases by nearly a factor of seven, although we acknowledge that, in this model, the single-site iron atom was not isolated but supported on a MoS<sub>2</sub> sheet. However, in the case of Li<sup>0</sup>, energetically favourable N<sub>2</sub>–metal interactions seem to demonstrate an opposite trend.

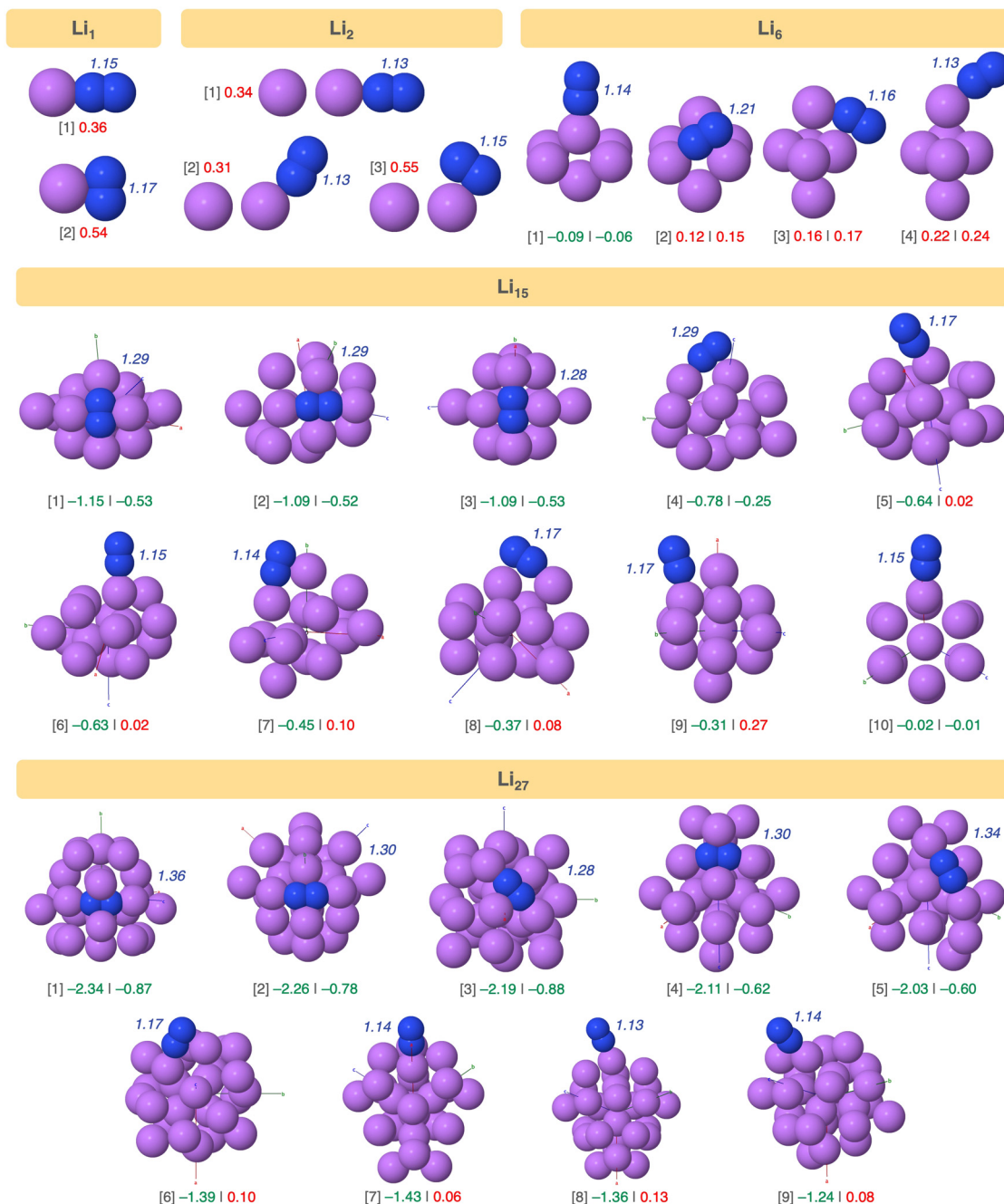
Li<sub>6</sub> formally represents a cluster-type structure of a square-based bipyramid. In this case, we observe binding events that deviate from the general trend in N<sub>2</sub>-philicity exhibited by Li<sup>0</sup> slabs. Four minima were identified for this model: the linear end-on minimum, labelled [1], was slightly physisorbed with a spontaneous  $\Delta G_b$  (ca. –0.1 eV), whereas  $\Delta G_b$  values were endergonic for minima [2]–[4], despite [2] and [3] being side-on (the former interacting with three Li centres and the latter with two) (Fig. 3).

Note that two binding free energy values are reported for all lithium models, except Li<sub>1</sub> and Li<sub>2</sub>, considered herein because, unlike metallic slabs, they undergo distortions relative to their original structures. The first  $\Delta G_b$  value refers to the free energy difference between the minimum in question and the perfectly symmetric cluster. The second  $\Delta G_b$  was calculated by re-optimising the cluster structure after removal of the adsorbed N<sub>2</sub> molecule. The difference between the two values provides a measure of the energy associated with cluster structure deformation. It should be noted that the deformation energy discussed herein does not necessarily reflect stability of Li<sup>0</sup> nanostructures during Li-NRR. A proper assessment of this aspect would require a separate investigation involving additional species participating in the reaction.



**Fig. 2** (a) Most common flat surfaces (top view) of bcc-Li<sup>0</sup>: (001), equivalent to (010) and (100), and (011), equivalent to (101) and (110). (b) Generation of sub-nanoclusters and nanoparticle models via spherical projection of the bcc-Li<sup>0</sup> unit cell, resulting in Li<sub>1</sub> (radius,  $r = 1$  Å), Li<sub>2</sub> ( $r = 2$  Å), Li<sub>6</sub> ( $r = 3$  Å), Li<sub>15</sub> ( $r = 4$  Å), Li<sub>27</sub> ( $r = 5$  Å), and Li<sub>51</sub> ( $r = 6$  Å).





**Fig. 3** \*N<sub>2</sub> minima obtained for the Li<sub>1</sub> atomic dot and the Li<sub>2</sub>, Li<sub>6</sub>, Li<sub>15</sub>, and Li<sub>27</sub> sub-nanoclusters. N≡N distances (in Å) are shown in *blue*. Binding free energies (in eV) are indicated in green for spontaneous and red for non-spontaneous values. For the Li<sub>15</sub> and Li<sub>27</sub> models, Cartesian axes (*a*, *b*, *c*, corresponding to OX, OY, and OZ, respectively) are also shown as a reference for orientation, facilitating comparison with the symmetric structures presented in Fig. 2.

For the Li<sub>6</sub> cluster, this difference was negligible, meaning almost no deformation with respect to the perfectly symmetrical cluster of this type (Fig. 3). In contrast, the change was pronounced for the Li<sub>15</sub> model, where the cluster deformation energies were, in most cases, around -0.6 eV; that is, the distorted structures were more stable than the initial perfectly symmetric cluster. From the perspective of formal analysis of N<sub>2</sub> stabilisation, it is the second  $\Delta G_b$  value that is of major interest.

For Li<sub>15</sub>, the binding free energies of N<sub>2</sub> were similar to those computed for the slabs, with values of -0.53, -0.52, and -0.53 eV for the first three minima, which: (i) are of side-on type; (ii) correspond to molecules embedded within the cluster surface; and (iii) show that superficial interactions lead to the decreased  $\Delta G_b$  values (Fig. 3). We also observed that the captured N<sub>2</sub> molecules in these more stable configurations were activated, with elongated N≡N bonds to approximately 1.3 Å. Although the  $\Delta G_b$  values for these three



most stable minima were similar, their electronic energies, *i.e.*, not including enthalpic and entropy contributions, differed by more than 0.1 eV ( $\approx 9.65$  kJ mol<sup>-1</sup>), confirming that they represent distinct structures on the potential energy surface (PES).

The binding free energy results for the Li<sub>27</sub> model further emphasised the intensification of the N<sub>2</sub>-Li<sup>0</sup> interactions within atom clusters as compared to the slab models. We identified a series of side-on and embedded N<sub>2</sub> minima, labelled [1]-[5], whereas minima [6]-[9] correspond to mixed ([6]), angular ([7]), and linear end-on ([8]-[9]) types (Fig. 3). For the former (side-on),  $\Delta G_b$  was exergonic, while for the latter (mixed and end-on) it was not. This behaviour, although somewhat expected, would not be particularly noteworthy if the N<sub>2</sub> binding free energy values did not significantly exceed the reference values obtained for slabs (-0.55 eV). In these clusters, the interactions reached as much as -0.9 eV, indicating that N<sub>2</sub> capture was strongly enhanced, accompanied by substantial activation of the molecule. For instance, for side-on minima [1]-[5], N≡N bond lengths ( $r_{N\equiv N}$ ) elongated up to 1.28-1.36 Å (Fig. 3).

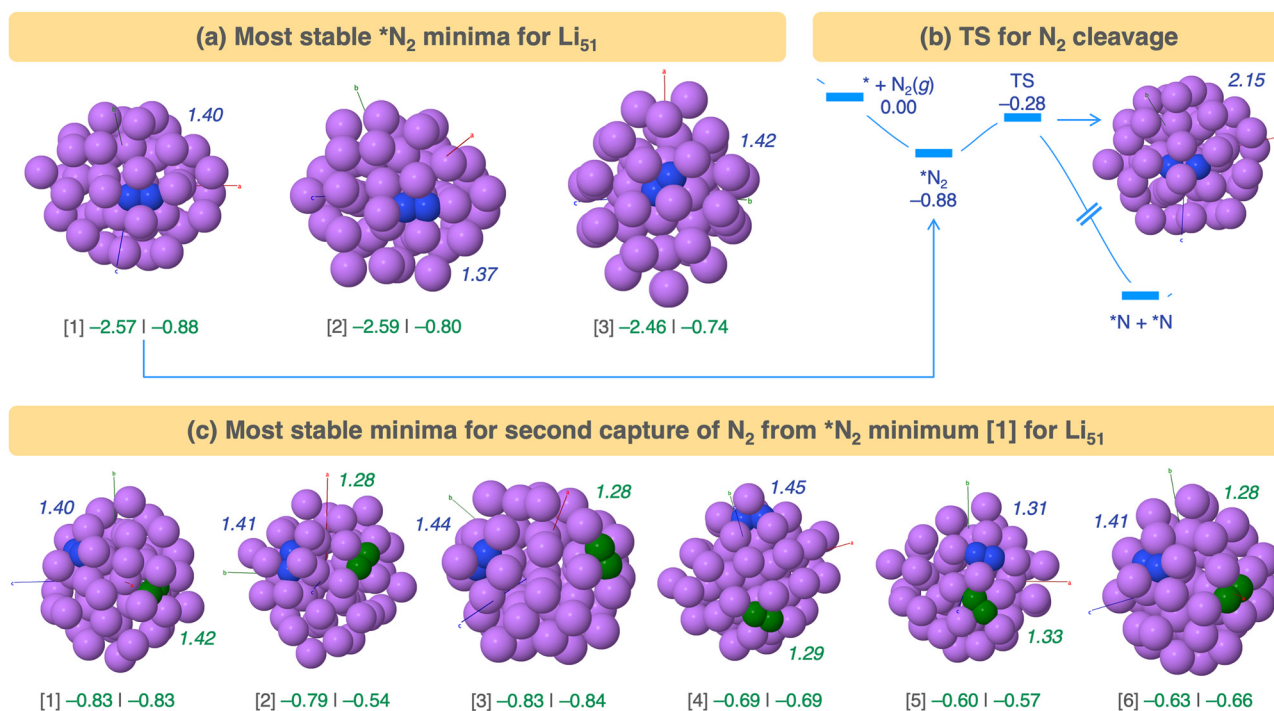
To explore if this improved N<sub>2</sub> fixation could be maintained, and to what extent, at the nano- rather than sub-nanometre scale discussed above, we further studied the Li<sub>51</sub> nanoparticle model. The intensity of N<sub>2</sub> capture in Li<sub>51</sub> followed a pattern similar to that observed in the Li<sub>27</sub> cluster: side-on and embedded N<sub>2</sub> structures on the nanoparticle surface exhibited  $\Delta G_b$  values ranging from -0.88 eV for the most stable to -0.60 eV for the least stable, whereas end-on

interactions remained endergonic. This is exemplified in Fig. 4a for the three most stable N<sub>2</sub> minima identified for the Li<sub>51</sub> model, with the remaining minima provided in Fig. S2. Notably, the N≡N bond elongation was slightly greater than that observed for the corresponding minima in the Li<sub>27</sub> model, and distortion energies were somewhat higher, indicative of enhanced stability in the non-symmetric Li<sub>51</sub> nanoparticle compared with the symmetric structure (*cf.* data in Fig. 3 and 4a). A correlation between  $\Delta G_b$  and the number of lithium atoms interacting with N<sub>2</sub> was noted for \*N<sub>2</sub> minima in both the Li<sub>27</sub> and Li<sub>51</sub> models (Fig. S5), consistent with our observation that binding free energy became stronger as the number of the Li...N bonds increased.

In light of these results obtained for the larger nanostructures, Li<sup>0</sup> appears to exhibit a pronounced affinity towards N<sub>2</sub>. Nevertheless, further investigations including explicit solvent molecules will be required to evaluate possible competitive adsorption effects under more realistic conditions.

Further, we sought to explore fundamental aspects of the N<sub>2</sub> cleavage mediated by the lithium metal nanoparticle by addressing two questions: (i) what is the energy cost (activation barrier) associated with cleaving the \*N<sub>2</sub> molecule into N adatoms (\*N) during the formation of lithium nitride; and (ii) is capture of a second N<sub>2</sub> molecule also an energetically favourable process?

Addressing the first question, the computed activation energies for the lithium-mediated dinitrogen dissociation process were low to moderate, that is, with values less than 1



**Fig. 4** (a) Most stable \*N<sub>2</sub> minima (all side-on) obtained for the Li<sub>51</sub> nanoparticle. (b) Potential energy surface (PES) for N<sub>2</sub> capture and cleavage of minimum [1] in the Li<sub>51</sub> model, including the optimised structure of the N≡N cleavage transition state. (c) Most stable minima for a second N<sub>2</sub> capture, where the original \*N<sub>2</sub> from minimum [1] is shown in blue and the second \*N<sub>2</sub> in green. N≡N distances (in Å) are shown in *italics*, following the same colour code.



eV. Specifically, we performed calculations to locate the transition states (TSs) for N<sub>2</sub> cleavage of the six most stable side-on N<sub>2</sub> minima in the Li<sub>51</sub> model, obtaining values ranging from 0.46 eV for the lowest barrier to 0.64 eV for the highest (Table S3 and Fig. S3). Remarkably, all activation barriers were smaller, in absolute terms, than the energy released upon N<sub>2</sub> chemisorption. This is exemplified in Fig. 4b for the minimum [1] of the Li<sub>51</sub> model. In this case, the \* + N<sub>2</sub>(g) → \*N<sub>2</sub> physical process exhibited a free energy drop of -0.88 eV, as discussed previously, and an activation barrier of 0.61 eV relative to this state. Additionally, the chemisorbed N<sub>2</sub> molecule showed a bond elongation of 0.28 Å ( $r_{\text{N}=\text{N}} = 1.40$  Å), while the nitrogen–nitrogen distance in the TS was 2.15 Å.

To assess the kinetic accessibility of the embedded N<sub>2</sub>-in-Li<sup>0</sup> configurations, a constrained approach scan was performed for the Li<sub>51</sub> model in which the distance between the central Li atom and N<sub>2</sub> was progressively reduced while relaxing the remaining degrees of freedom (see Fig. S6). The resulting profile reveals a very small electronic energy barrier of 0.14 eV for N<sub>2</sub> penetration into the subsurface region, followed by stabilisation of the embedded state. This indicates that subsurface incorporation is kinetically accessible within the capabilities of the static DFT approach employed herein.

To explore the capability of a lithium nanoparticle to activate multiple dinitrogen molecules, a second N<sub>2</sub> was randomly placed at various positions onto the \*N<sub>2</sub> minimum [1] of the Li<sub>51</sub> model. The capture of an additional N<sub>2</sub> molecule was found to be spontaneous, with structural distortion energies being practically negligible. In more detail, Fig. 4c shows six of the nineteen minima identified for this case, while the remaining structures are provided in Fig. S4. Notably, for all these minima, the second N<sub>2</sub> molecule adopted a side-on configuration in the most stable structures (shown in green in Fig. 4c), with spontaneous  $\Delta G_b$  values generally comparable to those computed for the first captured N<sub>2</sub> molecule (blue in Fig. 4c). Although the N≡N bond in the second N<sub>2</sub> molecule was also elongated, the magnitude of this effect was less pronounced. By contrast, for all these minima, the bond elongation of the first captured N<sub>2</sub> molecule was maintained or even enhanced, with N≡N distances exceeding 1.4 Å in some cases (Fig. S4).

Based on these results, Li<sup>0</sup> nanoparticles can be considered exceptionally effective N<sub>2</sub> capturers, even more effective than the slab models. To further support this conclusion, we performed NVT [substance (N), volume (V), temperature (T); canonical ensemble] molecular dynamics (MD) simulations starting from an initial configuration in which thirty-six N<sub>2</sub> molecules were positioned in an end-on orientation, pointing towards each of the surface lithium atoms of the Li<sub>51</sub> nanoparticle and placed at a distance of 3 Å. This setup established unfavourable initial conditions, namely, N<sub>2</sub> molecules at distances corresponding to non-adsorptive interactions and in an end-on orientation, which, as confirmed by our calculations, were non-spontaneous. If the Li<sub>51</sub> nanoparticle indeed behaves as a highly potent N<sub>2</sub>

capturer, the MD simulation could be expected to evolve towards spontaneous adsorption and activation.

The results in this regard were highly instructive. On the picosecond timescale, a progressive insertion of numerous N<sub>2</sub> molecules onto and within the nanoparticle was observed (Fig. 5). As the MD simulation progressed, convergence of the centres of mass of the N<sub>2</sub> molecules with those of the Li<sup>0</sup> atoms was found, indicating the point of interaction between dinitrogen and the lithium nanoparticle surface. Specifically, a substantial number of N<sub>2</sub> molecules were captured within the inner layers of Li<sub>51</sub> after approximately 1.5 ps of the total 4 ps simulation. Throughout this evolution, up to nearly 4 ps, N<sub>2</sub> capture by the Li<sup>0</sup> nanoparticle was accompanied by a decrease in the electronic energy (shown in dark red in Fig. 5), indicating that these adsorption events were

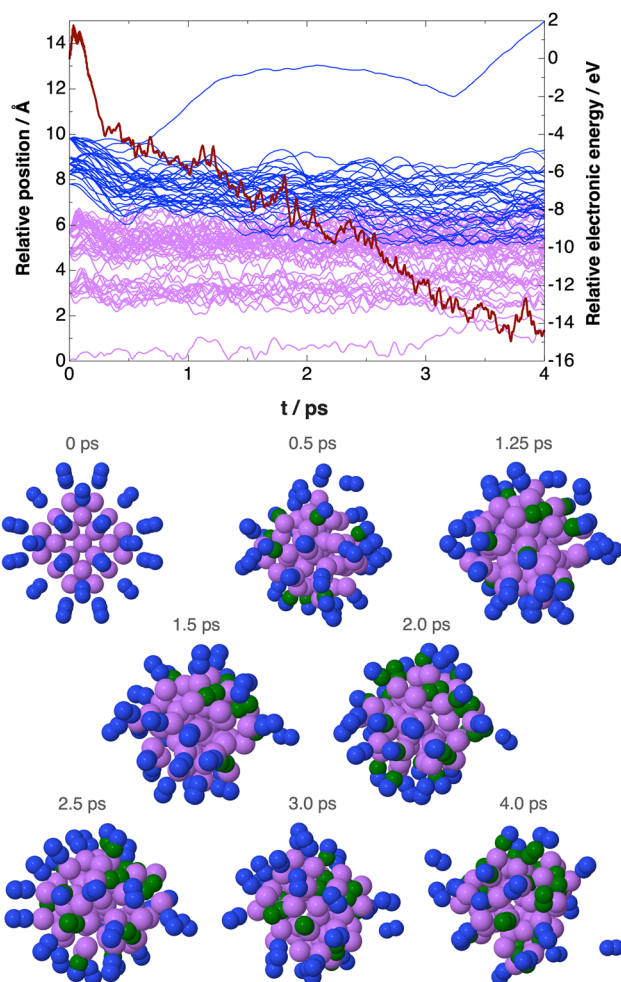


Fig. 5 NVT molecular dynamics for the Li<sub>51</sub> nanoparticle model surrounded by thirty-six N<sub>2</sub> molecules, each initially placed 3 Å away from the closest N atom interacting with Li. Relative positions (in Å) correspond to the distance of each Li atom (purple) and the centre of mass of each N<sub>2</sub> molecule (blue) relative to the nanoparticle's centre of mass. Evolution of electronic energy (dark red, in eV) is also shown, together with selected snapshots along the NVT MD trajectory. In these structures, N atoms of the N<sub>2</sub> molecules that are closer to the nanoparticle's centre of mass than the outermost Li<sup>0</sup> atom are highlighted in green.



energetically stabilising. Fig. 5 also includes selected snapshots along the MD trajectory, visually illustrating the progressive N<sub>2</sub> capture behaviour of the lithium nanoparticle.

The findings discussed above provide new insights into the chemistry governing the first critical step of the Li-NRR, *viz.* activation of dinitrogen. In the future, it will be important to investigate how the cleavage of a sufficiently large number of captured N<sub>2</sub> molecules ultimately leads to nitridation of the Li<sup>0</sup> nanoparticles. Simultaneous capture of multiple N<sub>2</sub> molecules by the Li<sup>0</sup> nanoparticles may introduce electronic interactions between adsorbates, which could influence the activation of subsequent N<sub>2</sub> molecules. Of an even higher interest is the location of the nitride formation, *viz.* at the surface or in the bulk of the particles, which might result in very different scenarios. Indeed, formation of a lithium nitride surface layer may potentially passivate Li<sup>0</sup> and suppress subsequent N<sub>2</sub> capture. On the contrary, formation of the lithium nitride in the bulk, which might be possible owing to the N<sub>2</sub> diffusion through Li<sup>0</sup> demonstrated herein and previously,<sup>8–10</sup> might allow for the most effective utilisation of the electrochemically generated lithium metal. A detailed assessment of these effects would require a more extensive exploration of the system and is left for future investigation.

## Conclusions

In summary, our DFT investigation demonstrated that Li<sup>0</sup> sub-nanoclusters and especially nanoparticles exhibit a remarkable capability for N<sub>2</sub> capture and activation, surpassing that of the previously studied metal slabs. Our calculations revealed that side-on and embedded adsorption configurations lead to significant elongation of the N≡N bond and accessible activation barriers for cleavage, and that multiple N<sub>2</sub> molecules can be captured successively. NVT molecular dynamics simulations indicate that the highly energetically favourable capture events take place on a picosecond scale, with N<sub>2</sub> molecules progressively inserted into the nanoparticle's inner layers. These findings highlight the unique mechanistic role of Li<sup>0</sup> nanostructures in the lithium redox-mediated nitrogen reduction reaction and suggest that such nanoparticles may serve as highly efficient platforms for ammonia synthesis under mild electrochemical conditions. Overall, this study provides fundamental insights into the interplay between size, structure, and reactivity against N<sub>2</sub> of lithium metal, opening new avenues for the design of Li-NRR catalysts.

## Conflicts of interest

J. M. D.-R., R. Q.-C., and L. M. A. declare no conflicts of interest. D. R. M. and A. N. S. declare competing financial interests since they are shareholders in Jupiter Ionics Pty Ltd, a start-up company focused on commercialising ammonia production technology.

## Data availability

The material presented herein does not possess any characteristics that require additional guidance regarding its findable, accessible, interoperable, and reusable (FAIR) aspects. All supplementary materials have been prepared following best practices to ensure reproducibility, providing detailed instructions for any reader to replicate the results, subject to the necessary permissions to use the software employed in this study. Furthermore, upon publication of this article, a preprint version of its content will be deposited in my university's institutional repository, in accordance with standard institutional practice.

Supplementary information (SI): methodological aspects, full computational details. See DOI: <https://doi.org/10.1039/d6cy00192k>.

## Acknowledgements

J. M. D.-R., R. Q.-C., and L. M. A. are grateful to MCIN/AEI and NextGenerationEU/PRTR for the financial support through project ref. PID2022-143294OB-I00. D. R. M. and A. N. S. acknowledge the financial support from the Australian Research Council (DP250102783). L. M. A. is a Ramón y Cajal fellow (ref. RYC2021-030994-I). Gratitude is also due to the KAUST Supercomputer Laboratory (KSL), KSA, for providing the computational resources (Shaheen III). This paper is dedicated to Prof Steve Scheiner (Utah State University) in recognition of his outstanding contributions to the study of non-covalent interactions.

## Notes and references

- 1 D. R. MacFarlane, P. V. Cherepanov, J. Choi, B. H. R. Suryanto, R. Y. Hodgetts, J. M. Bakker, F. M. Ferrero Vallana and A. N. Simonov, *Joule*, 2020, **4**, 1186–1205.
- 2 H. Zhang, Z. Lu and Z. Wen, *Nat. Rev. Chem.*, 2025, **9**, 285–286.
- 3 J. M. McEnaney, A. R. Singh, J. A. Schwalbe, J. Kibsgaard, J. C. Lin, M. Cargnello, T. F. Jaramillo and J. K. Nørskov, *Energy Environ. Sci.*, 2017, **10**, 1621–1630.
- 4 S. Z. Andersen, M. J. Statt, V. J. Bukas, S. G. Shapel, J. B. Pedersen, K. Krempel, M. Saccoccio, D. Chakraborty, J. Kibsgaard, P. C. K. Vesborg, J. Nørskov and I. Chorkendorff, *Energy Environ. Sci.*, 2020, **13**, 4291–4300.
- 5 K. Li, S. Z. Andersen, M. J. Statt, M. Saccoccio, V. J. Bukas, K. Krempel, R. Sažinas, J. B. Pedersen, V. Shadravan, Y. Zhou, D. Chakraborty, J. Kibsgaard, P. C. K. Vesborg, J. K. Nørskov and I. Chorkendorff, *Science*, 2021, **374**, 1593–1597.
- 6 B. H. R. Suryanto, K. Matuszek, J. Choi, R. Y. Hodgetts, H.-L. Du, J. M. Bakker, C. S. M. Kang, P. V. Cherepanov, A. N. Simonov and D. R. MacFarlane, *Science*, 2021, **372**, 1187–1191.
- 7 H.-L. Du, M. Chatti, R. Y. Hodgetts, P. V. Cherepanov, C. K. Nguyen, K. Matuszek, D. R. MacFarlane and A. N. Simonov, *Nature*, 2022, **609**, 722–727.



- 8 D. R. MacFarlane, A. N. Simonov, T. M. Vu, S. Johnston and L. M. Azofra, *Faraday Discuss.*, 2023, **243**, 557–570.
- 9 L. M. Azofra, J. M. Doña-Rodríguez, D. R. MacFarlane and A. N. Simonov, *J. Phys. Chem. C*, 2025, **129**, 1198–1205.
- 10 T. Ludwig, A. R. Singh and J. K. Nørskov, *J. Phys. Chem. C*, 2020, **124**, 26368–26378.
- 11 Y. Abghoui, A. L. Garden, J. G. Howalt, T. Vegge and E. Skúlason, *ACS Catal.*, 2016, **6**, 635–646.
- 12 D. Wang, L. M. Azofra, M. Harb, L. Cavallo, X. Zhang, B. H. R. Suryanto and D. R. MacFarlane, *ChemSusChem*, 2018, **11**, 3416–3422.
- 13 B. H. R. Suryanto, D. Wang, L. M. Azofra, M. Harb, L. Cavallo, R. Jalili, D. R. G. Mitchell, M. Chatti and D. R. MacFarlane, *ACS Energy Lett.*, 2019, **4**, 430–435.
- 14 S. Kamiguchi, K. Asakura, T. Shibayama, T. Yokaichiya, T. Ikeda, A. Nakayama, K. Shimizu and Z. Hou, *Chem. Sci.*, 2024, **15**, 2914–2922.
- 15 Y. Jiang, Z. Chen, T. Peng, L. Jiao, X. Pan, H.-L. Jiang and X. Bao, *Angew. Chem., Int. Ed.*, 2025, **64**, e202501190.
- 16 S. C. Jesudass, S. Surendran, J. Y. Kim, T.-Y. An, G. Janani, T.-H. Kim, J. K. Kim and U. Sim, *Electrochem. Energy Rev.*, 2023, **6**, 27.
- 17 W.-Q. Li, M. Xu, J.-S. Chen and T.-N. Ye, *Adv. Mater.*, 2024, **36**, 2408434.
- 18 T. Žibert, B. Likozar and M. Huš, *ChemSusChem*, 2024, **17**, e202301730.
- 19 X. Fu, A. Xu, J. B. Pedersen, S. Li, R. Sažinas, Y. Zhou, S. Z. Andersen, M. Saccoccio, N. H. Deissler, J. B. V. Mygind, J. Kibsgaard, P. C. K. Vesborg, J. K. Nørskov and I. Chorkendorff, *Nat. Commun.*, 2024, **15**, 2417.
- 20 H.-L. Du, K. Matuszek, R. Y. Hodgetts, K. Ngoc Dinh, P. V. Cherepanov, J. M. Bakker, D. R. MacFarlane and A. N. Simonov, *Energy Environ. Sci.*, 2023, **16**, 1082–1090.
- 21 X. Fu, V. A. Niemann, Y. Zhou, S. Li, K. Zhang, J. B. Pedersen, M. Saccoccio, S. Z. Andersen, K. Enemark-Rasmussen, P. Benedek, A. Xu, N. H. Deissler, J. B. V. Mygind, A. C. Nielander, J. Kibsgaard, P. C. K. Vesborg, J. K. Nørskov, T. F. Jaramillo and I. Chorkendorff, *Nat. Mater.*, 2024, **23**, 101–107.
- 22 M. Krebsz, R. Y. Hodgetts, S. Johnston, C. K. Nguyen, Y. Hora, D. R. MacFarlane and A. N. Simonov, *Energy Environ. Sci.*, 2024, **17**, 4481–4487.
- 23 V. Azumah, L. Kavalsky and V. Viswanathan, *J. Phys. Chem. C*, 2025, **129**, 2488–2501.
- 24 V. Azumah, L. Kavalsky and V. Viswanathan, *J. Catal.*, 2025, **450**, 116250.
- 25 K. N. Wood, M. Noked and N. P. Dasgupta, *ACS Energy Lett.*, 2017, **2**, 664–672.
- 26 K. B. Hatzell, X. C. Chen, C. L. Cobb, N. P. Dasgupta, M. B. Dixit, L. E. Marbella, M. T. McDowell, P. P. Mukherjee, A. Verma, V. Viswanathan, A. S. Westover and W. G. Zeier, *ACS Energy Lett.*, 2020, **5**, 922–934.
- 27 K. Dong, Y. Xu, J. Tan, M. Osenberg, F. Sun, Z. Kochovski, D. T. Pham, S. Mei, A. Hilger, E. Ryan, Y. Lu, J. Banhart and I. Manke, *ACS Energy Lett.*, 2021, **6**, 1719–1728.
- 28 B. Hammer, L. B. Hansen and J. K. Nørskov, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 7413–7421.
- 29 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, *J. Chem. Phys.*, 2010, **132**, 154104.
- 30 S. Grimme, S. Ehrlich and L. Goerigk, *J. Comput. Chem.*, 2011, **32**, 1456–1465.
- 31 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1993, **47**, 558–561.
- 32 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **49**, 14251–14269.
- 33 G. Kresse and J. Furthmüller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169–11186.
- 34 G. Kresse and J. Furthmüller, *Comput. Mater. Sci.*, 1996, **6**, 15–50.
- 35 B. H. R. Suryanto, C. S. M. Kang, D. Wang, C. Xiao, F. Zhou, L. M. Azofra, L. Cavallo, X. Zhang and D. R. MacFarlane, *ACS Energy Lett.*, 2018, **3**, 1219–1224.
- 36 L. M. Azofra, C. Sun, L. Cavallo and D. R. MacFarlane, *Chem. Eur. J.*, 2017, **23**, 8275–8279.

