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Graph-convolutional neural networks for (QM)ML/MM molecular dynamics simulations†

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To accurately study the chemical reactions in the condensed phase or within enzymes, both quantum-mechanical description and sufficient configurational sampling are required to reach converged estimates. Here, quantum mechanics/molecular mechanics (QM/MM) molecular dynamics (MD) simulations play an important role, providing QM accuracy for the region of interest at a decreased computational cost. However, QM/MM simulations are still too expensive to study large systems on longer time scales. Recently, machine learning (ML) models have been proposed to replace the QM description. The main limitation of these models lies in the accurate description of long-range interactions present in condensed-phase systems. To overcome this issue, a recent workflow has been introduced combining a semi-empirical method (*i.e.* density functional tight binding (DFTB)) and a high-dimensional neural network potential (HDNNP) in a Δ -learning scheme. This approach has been shown to be capable of correctly incorporating long-range interactions within a cutoff of 1.4 nm. One of the promising alternative approaches to efficiently take long-range effects into account is the development of graph-convolutional neural networks (GCNNs) for the prediction of the potential-energy surface. In this work, we investigate the use of GCNN models – with and without a Δ -learning scheme – for (QM)ML/MM MD simulations. We show that the Δ -learning approach using a GCNN and DFTB as a baseline achieves competitive performance on our benchmarking set of solutes and chemical reactions in water. This method is additionally validated by performing prospective (QM)ML/MM MD simulations of retinoic acid in water and S-adenosylmethionine interacting with cytosine in water. The results indicate that the Δ -learning GCNN model is a valuable alternative for the (QM)ML/MM MD simulations of condensed-phase systems.

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1 Introduction

A key goal of computational chemistry is the molecular level understanding of chemical reactions in solution and enzymes. For this, the free-energy change (rather than the change in potential energy) during a reaction process is the central property. To calculate free-energy differences, molecular dynamics (MD) simulations from tens of picoseconds to hundreds of nanoseconds are typically required to obtain sufficiently converged results. While classical force fields can be used for MD simulations of condensed-phase systems over long time scales,^{1,2} higher level quantum-mechanical (QM) methods are required for an accurate description of molecular interactions and chemical reactions. Unfortunately, QM calculations are much more computationally intensive, limiting the accessible time and spatial scales.

In order to overcome this bottleneck, the combined QM and molecular mechanical (QM/MM) approach provides a QM description of the region of interest (QM zone) coupled with a realistic modelling of the long-range interactions of the surrounding condensed-phase system (MM zone).^{3–6} The interaction between the zones is then calculated based on either mechanical constraints (*i.e.*, a “mechanical embedding” scheme) or electronic perturbations (*i.e.*, an “electrostatic embedding” scheme). Generally, the electrostatic embedding scheme has been shown to be more accurate and it is currently the gold standard for QM/MM simulations.^{5,7–10} In the QM/MM scheme, the QM zone requires electronic structure calculations at each time step and is thus the computational bottleneck. While the computational costs of QM/MM MD simulations are reduced compared to full *ab initio* simulations, the accessible time and spatial scales are still not sufficient for most free-energy calculations. This issue can be partially circumvented by using semi-empirical methods to describe the QM zone.^{9,11,12} However, this reduces not only the computational cost but also the achievable accuracy. An alternative is to use machine-learned (ML) potentials to describe the QM zone.

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In recent years, there have been major advances in the development of ML models trained to reproduce the potential-energy surface (PES) of chemical systems.^{13–28} For small-to-medium sized compounds in the gas phase or in periodic materials, state-of-the-art ML models have been shown to achieve chemical accuracy, for both the predicted energies and forces.^{20–24,29} However, condensed-phase systems pose an additional challenge to ML models due to the typically large number of involved atoms, element types, and rotatable bonds without any exploitable symmetries, and important long-range interactions and non-local charge transfer. In particular, the long-range interactions and the non-local charge transfer currently limit the established ML methods, as these often rely on local descriptors and are thus unable to take global changes in the electronic structure into account.^{26,30,31} There are two main approaches to overcome this locality dependence of the descriptors used in the current ML models. Either non-local information transfer is incorporated directly in the ML models^{26,30–32} or ML methods are combined with a fast or semi-empirical QM method with an explicit treatment of long-range interactions in a Δ -learning scheme.^{33–36} The Δ -learning approach has already been shown to be promising for the (QM)ML/MM MD simulations of condensed-phase systems.^{34–36}

Next to the descriptor-based approaches, there has also been a lot of development on message passing approaches trained to reproduce the PES of QM systems.^{22,23,37–39} These graph networks typically use dense layers of neural networks as non-linear functions for the message passing convolutions and are thus known as graph-convolutional neural networks (GCNNs). One of the main advantages of GCNNs compared to descriptor-based ML models is that no specialized descriptors have to be developed for chemical systems. Instead, this is achieved directly through the graph-convolutional layers. A further advantage is that through iterative message passing operations more distant information is taken into account and thus the local dependence of descriptor based models can be (partially) avoided. However, each consecutive convolution substantially increases the model size. Thus, in practice, long-range contributions are terminated at a certain cutoff and global changes in the electronic structure are not considered. To address this issue, the total energy can be separated into a short-range electrostatic term and a long-range electrostatic term, for which a GCNN is used to predict atomic charges.^{23,31} Furthermore, long-range charge transfer and global changes in the electronic structure may be taken into account by including a global state or a self-attention mechanism in the message passing operation.^{23,31,32,40}

In this work, we assess the applicability of GCNNs to reproduce the PES of condensed-phase systems and their use in (QM)ML/MM MD simulations. For this, we validate a GCNN model with and without a Δ -learning scheme on different molecular systems in water and compare it with the previously developed high-dimensional neural network potentials (HDNNPs),^{34,41} which use the same Δ -learning scheme with DFTB^{42,43} as the baseline method. In the Theory section, we briefly describe the relevant concepts behind the QM/MM approach as well as the GCNNs. In the Methods section, we

describe the setup used for the QM/MM simulations, the implemented GCNN architecture as well as the training setup. In the Results section, we evaluate different GCNN models, including different global information transfer schemes, with different (QM)ML/MM and training setups. Finally, we compare the resulting GCNN model with the previous HDNNP model.³⁴

2 Theory

2.1 QM/MM scheme

QM/MM is a multi-scale approach, which incorporates a QM zone within a larger MM zone describing the condensed-phase system.^{3–6} This enables the accurate calculation of the region of interest coupled with a realistic modelling of long-range interactions with the surrounding environment. The main challenge here is how to describe the interaction between these two different zones. In order to calculate the total energy ($E_{\text{QM/MM}}(\vec{R})$) of the combined QM/MM system, an additive or subtractive scheme can be chosen. In the more prominent additive scheme, the total energy ($E_{\text{QM/MM}}(\vec{R})$) is described as the sum of the energy of the QM subsystem ($E_{\text{QM}}(\vec{R}_{\text{QM}})$) and the MM subsystem ($E_{\text{MM}}(\vec{R}_{\text{MM}})$) plus the electrostatic ($E_{\text{QM-MM}}^{\text{el}}(\vec{R})$) and short-range van der Waals interactions ($E_{\text{QM-MM}}^{\text{vdW}}(\vec{R})$) between the two subsystems.

$$E_{\text{QM/MM}}(\vec{R}) = E_{\text{QM}}(\vec{R}_{\text{QM}}) + E_{\text{MM}}(\vec{R}_{\text{MM}}) + E_{\text{QM-MM}}^{\text{el}}(\vec{R}) + E_{\text{QM-MM}}^{\text{vdW}}(\vec{R}) \quad (1)$$

Note the distinction between \vec{R} referring to all nuclei in the system, and \vec{R}_{QM} and \vec{R}_{MM} referring to the nuclei of the QM and MM zones, respectively. In the additive scheme, the interaction terms between the QM and MM zones are described either *via* a mechanical or an electrostatic embedding scheme,⁷ where the latter scheme has been shown to be more accurate.⁵ For this, two Hamiltonians are introduced in the QM calculation. In atomic units, $\hat{H}_{\text{QM-MM}}^{\text{el}}$ is given as

$$\hat{H}_{\text{QM-MM}}^{\text{el}} = - \sum_i^{N_{\text{MM}}} \sum_j^{N_{\text{el}}} \frac{q_i}{|\vec{R}_{\text{MM},i} - \vec{r}_j|} + \sum_i^{N_{\text{QM}}} \sum_j^{N_{\text{MM}}} \frac{Z_i q_j}{|\vec{R}_{\text{QM},i} - \vec{R}_{\text{MM},j}|} \quad (2)$$

where $\hat{H}_{\text{QM-MM}}^{\text{vdW}}$ is treated classically and is given as

$$\hat{H}_{\text{QM-MM}}^{\text{vdW}} = E_{\text{QM-MM}}^{\text{vdW}}(\vec{R}) = \sum_i^{N_{\text{QM}}} \sum_j^{N_{\text{MM}}} 4\epsilon_{ij} \left(\left(\frac{\sigma_{ij}}{|\vec{R}_i - \vec{R}_j|} \right)^{12} - \left(\frac{\sigma_{ij}}{|\vec{R}_i - \vec{R}_j|} \right)^6 \right), \quad (3)$$

where q_i is the partial charge of the MM atom i , and ϵ_{ij} and σ_{ij} are fitted parameters. This means that the QM subsystem is directly influenced by the MM partial charges, while the MM subsystem “feels” a force from the perturbed QM subsystem.



Therefore, eqn (1) in the electrostatic embedding scheme becomes

$$E_{\text{QM/MM}}(\vec{R}) = \frac{\langle \psi(\vec{r}) | (\hat{H}_{\text{QM}} + \hat{H}_{\text{QM-MM}}^{\text{el}}) \psi(\vec{r}) \rangle}{\langle \psi(\vec{r}) | \psi(\vec{r}) \rangle} + E_{\text{MM}}(\vec{R}_{\text{MM}}) + E_{\text{QM-MM}}^{\text{vdw}}(\vec{R}) \quad (4)$$

Note that only MM particles within a given cutoff radius R_c of the QM zone are included in the summations in eqn (2) and (3). Typically, a relatively large cutoff radius (R_c) of around 1.4 nm is required to achieve sufficiently converged results.^{8,34} The arising issue due to the non-continuous PES at the cutoff can be partially resolved using adaptive resolution schemes.⁴⁴

2.2 Graph-convolutional neural networks

In our description of graph-convolutional neural networks (GCNNs) or message passing neural networks, we follow the notations used by ref. 39 and 45. Here, GCNNs are permutation-invariant ML models that operate on the graph structured data. In the present work, atoms are represented as nodes ν and interactions between atoms are represented as edges e , whereby only interactions up to a certain cutoff R_{edge} are considered as edges. Note that we will refer to two different cutoffs throughout the manuscript. R_c is the cutoff used in the QM/MM scheme, and the summations in eqn (2) and (3) run over all partial charges within R_c . R_{edge} , on the other hand, is the cutoff used for the edge definition in the GCNNs. The GCNNs contain consecutive graph-convolutional layers, where each layer consists of an edge or message update operation (eqn (5)), an aggregation operation (eqn (6)), and a node update operation (eqn (7)). Considering a graph $G = (V, E)$ with nodes $\nu_i \in V$ and edges $e_{ij} \in E$, message passing can be defined as follows:

$$m_{ij} = \phi_e(h_i^l, h_j^l, a_{ij}) \quad (5)$$

$$m_i = \sum_{j \in N(i)} m_{ij} \quad (6)$$

$$h_i^{l+1} = \phi_h(h_i^l, m_i). \quad (7)$$

Here, the superscript l denotes the current layer, $h_i^l \in \mathbf{R}^n$ describes the hidden-feature vector of a node ν_i at a layer l , $a_{ij} \in \mathbf{R}^n$ describes the edge feature of the edge e_{ij} between nodes i and j , $N(i)$ denotes the set of neighbors of the node i , and ϕ_e and ϕ_h describe update functions. Different GCNN models commonly differ by their used features, update functions (ϕ_e and ϕ_h), and aggregation functions.^{39,45} The update functions are most commonly approximated by multilayer perceptrons. Note that we use a summation as an aggregation function (eqn (6)), but other aggregation functions such as min or max functions have also been investigated.⁴⁶

3 Methods

3.1 Systems

To allow a direct comparison, the same systems as in ref. 34 were investigated. For the validation of different GCNN models

and training setups, we used two single-solute systems with different MM cutoff radii (R_c): (i) benzene in water, and (ii) uracil in water. For the comparison to the previous HDNNP model,³⁴ we used the largest single-solute system (retionic acid in water), and two chemical reactions in water (constrained close to the transition state): (i) the second-order nucleophilic substitution (S_N2) reaction of CH_3Cl with Cl^- , and (ii) the reaction of *S*-adenosylmethionine (SAM) with cytosine. We investigated the accuracy of the different models in training/validation/test splits and the performance in a prospective (QM)ML/MM MD simulation. The data sets are freely available on <https://www.research-collection.ethz.ch/handle/20.500.11850/512374>.

3.2 General computational details

All QM/MM and (QM)ML/MM MD simulations were performed using the GROMOS software package^{47,48} interfaced to DFTB+/19.2^{42,43} and ORCA/4.2.0.⁴⁹ All structures used in the training, validation and test sets were taken from ref. 34 (available at <https://www.research-collection.ethz.ch/handle/20.500.11850/512374>). These come from QM/MM MD trajectories, where the first 70% of the 10 000 frames were considered to be the training set, the following 20% constitute the validation set, and the last 10% are taken as the test set. The computational details of the (QM)ML/MM simulations are provided in ref. 34. Note that a MM cutoff radius of $R_c = 0.6$ nm was used for benzene in water, while $R_c = 1.4$ nm was used for all other systems. The GCNNs were implemented using Tensorflow/keras.⁵⁰ The models were trained in Python and then exported to the C++ GROMOS code as described in ref. 34 for the HDNNPs.

For the validation of the different GCNN models, we compare both the full-QM learning task and the Δ -learning scheme. For the full-QM learning task, we use the GCNN models to directly predict the DFT⁴⁹ energies and forces, whereas DFTB^{42,43} is used as the baseline method in the Δ -learning scheme and the GCNN models predicts the difference between the DFT and DFTB properties. After the initial validation and comparison of the two learning tasks, we continue only with the Δ -learning approach throughout the remainder of the study.

3.3 (QM)ML/MM MD simulations

For retionic acid in water and the SAM/cytosine transition state in water, we performed (QM)ML/MM MD simulations using the trained GCNN models with the Δ -learning setup. To ensure comparability, we used the same MD and DFTB settings as in ref. 34. The time step was set to 0.5 fs, the temperature was set to $T = 298$ K, and the pressure was set to 1 bar. For the SAM/cytosine system, we set the force constant for the position restraints to $2000 \text{ kJ mol}^{-1} \text{ nm}^{-2}$. All point charges within the cutoff radius $R_c = 1.4$ nm were included in an electrostatic embedding scheme in the QM(ML) computation. The selected solvent atoms beyond the cutoff were included to avoid the bond-breaking and creation of artificial charges. Long-range electrostatic interactions beyond R_c were included using a reaction-field method.⁵¹ Note that the reaction field acts only on the MM particles. As starting coordinates for the prospective



simulations, we used the last snapshot from the test set (which originated from the initial QM/MM MD trajectory). The simulations were performed for 170 000 steps for retionic acid in water and 110 000 steps for the SAM/cytosine transition state in water.

3.4 GCNN architectures

The basic building blocks of the GCNN are so-called dense layers. They take an input vector $x \in R^{n_{in}}$ and return an output vector $y \in R^{n_{out}}$ according to the transformation

$$y = Wx + b. \quad (8)$$

Here, $W \in R^{n_{in} \times n_{out}}$ and $b \in R^{n_{out}}$ are learnable parameters. In order to model arbitrary non-linear relationships, at least two dense layers need to be stacked and combined with a (non-linear) activation function σ . Here, we use a generalized SiLU (sigmoid linear unit) activation known as the Swish activation function,⁵² which is given as $\sigma(x) = x \cdot \text{sigmoid}(x)$ and has been used successfully in the recently published GCNNs for molecular systems.^{31,38,39} The inputs to the GCNN (v_i) are the nuclear charges $Z_i \in \mathbb{N}$ and positions $\vec{r}_i \in R^3$. We evaluate four different GCNN architectures with different global information transfer schemes. An overview of the basic network architecture is given in Fig. 1A. In the following sections, this architecture is referred to as the GCNN model.

Embedding block. An embedding is a mapping from a discrete object to a vector of real numbers. Here, the atomic numbers are mapped to embeddings $e_z \in R^F$, where the entries of e_z are learnable parameters and n_f denotes the number of features. Note that the number of features is kept constant

throughout the network. The embedding vector is then used to initialize the atomic feature vector h_i^0 .

Edge embedding block. A continuous filter convolution block is used to generate the edge representations a_{ij} (Fig. 1B). First, all edges are expressed as Euclidean distances. These are subsequently transformed to linear combinations of rationally invariant filters (RBFs) of radial basis functions $\sin\left(\frac{n\pi}{R_{\text{edge}}}\right) \frac{\|\vec{r}_{ij}\|}{\|\vec{r}_{ij}\|}$ as proposed by Klicpera *et al.*³⁸ Additionally, we apply a cosine cutoff to the filters,⁴¹ which ensures continuous behavior when an atom enters or leaves the cutoff sphere.

Interaction block. The interaction blocks calculate the message passing operation as defined in eqn (5)–(7) (Fig. 1E), generating the atomic feature vectors h_i^l at layers l .

Output block. The atomic feature vectors h_i^l at each layer l are passed through an output block, consisting of two stacked dense layers combined with an activation function (σ) (Fig. 1C). The output vectors s_i^l of each layer l are then summed and passed through a post-processing block, consisting of two dense layers, each combined with an activation function (Fig. 1D). The outputs of the post-processing block are the atomic energies. Finally, the total energy is calculated as the sum over all atomic energies. The forces are subsequently obtained as the derivatives of the total energy with respect to the Cartesian coordinates of the atoms. For this, we use the reverse mode automatic differentiation implemented using Tensorflow.

Unke *et al.*³¹ introduced a GCNN model, which takes into account non-local effects and charge transfer. They achieved this by introducing non-local interactions using a self-attention layer.^{40,53} An attention layer maps a matrix $X \in R^{T \times n_{in}}$ of

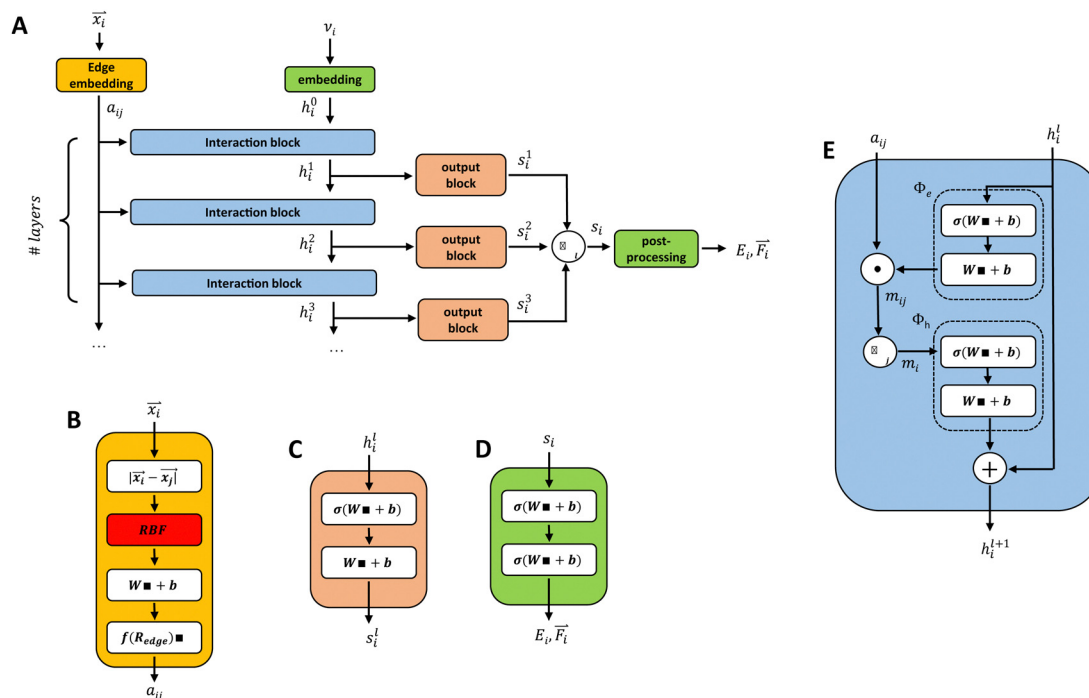


Fig. 1 Overview of the employed GCNN models with the full architecture (A), the cfconv block (B), the output block (C), the post-processing block (D), and the interaction block (E). For all linear layers $Wx + b$, we use n_f . For all activation functions $\sigma(\cdot)$, we use a Swish activation function.



query (Q) tokens T to n_{out} dimensions using a matrix $Y \in R^{T \times n_{\text{in}}}$ of key (K) tokens T' as follows:

$$\text{Attention}(Q, K, V) = \sigma\left(\frac{QK^T}{\sqrt{(d_k)}}\right)V \quad (9)$$

$$Q = XW_Q, K = YW_K, V = YW_V. \quad (10)$$

The layer is parametrized by a query matrix W_Q in $R^{n_{\text{in}} \times n_k}$, a key matrix W_K in $R^{n_{\text{in}} \times n_k}$, and a value matrix W_V in $R^{n_{\text{in}} \times n_{\text{out}}}$. A self-attention layer uses attention on the same sequence ($X = Y$). In this work, we use a multi-head self-attention mechanism as described in ref. 53. Here, the attention is calculated for multiple heads using a reduced attention $n_{\text{out}} = n_k < n_{\text{in}}$ dimensionality. The multiple attentions are then concatenated and transformed as

$$\text{MultiAttention} = \{\text{Attention}_0 | \text{Attention}_1 | \dots | \text{Attention}_{N_{\text{heads}}}\} W_{\text{multi}}, \quad (11)$$

where $W_{\text{multi}} \in R^{N_{\text{heads}} n_{\text{out}} \times n_{\text{multi}}}$ is used to parametrize the layer.

We introduce the multi-head self-attention layer at two different stages in our basic GCNN architecture: (i) within the interaction block (Fig. 2A) as described in ref. 31, and (ii) after

at the end of the model (Fig. 2B). This way we can investigate the effect of repeating the global information transfer within each message passing update as compared to a single transfer step using the final atom features. In the following sections, these models are referred to as the ‘‘interaction GCNN’’ and ‘‘attention GCNN’’.

A different GCNN architecture, which considers global information transfer, has been reported by Chen *et al.*³² Here, the authors introduce a global state u^l into the message passing operation. We adapted this idea by changing the message update step (eqn (7)) to

$$u^{l+1} = u^l + \phi_u\left(\sum_i h_i^l, \sum_i m_i, \sum_i h_i^l m_i\right) \quad (12)$$

$$h^{l+1} = h^l + \phi_h(h_i^l, m_i, u^{l+1}), \quad (13)$$

where $u^0 \in R^F$ is initialized as zero. The updated interaction block is shown in Fig. 2C. In the following sections, we refer to this architecture as the ‘‘global GCNN’’.

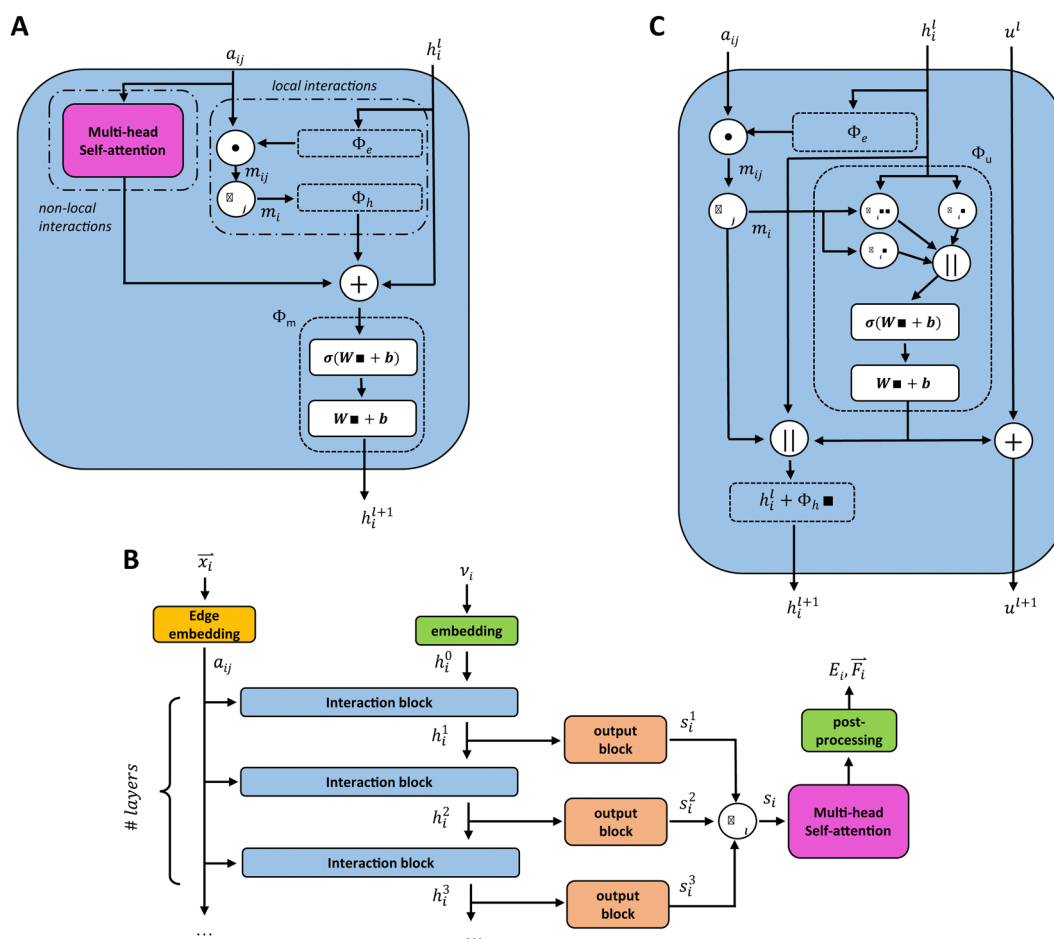


Fig. 2 Changes in the architecture of the GCNN models with different global information transfer schemes. (A) Interaction block of the GCNN model with multi-head self-attention layers in the interaction block. (B) Full architecture of the GCNN model with multi-head self-attention layers prior to the post-processing block. (C) Interaction block of the GCNN model with a global state u^l .



3.5 Training setup

We implemented the four different GCNN model architectures using Tensorflow/keras. The models were trained using the Adam optimizer⁵⁴ with an exponentially decaying learning rate ([initial learning-rate, decay steps, and decay rate] = [$1e^{-3}$, $5e^3$, and 0.96]). The Adam optimizer is one of the most well-established methods for the training of neural network models. It is based on a stochastic gradient descent optimizer that uses an adaptive estimation of the first and second-order moments. The models were trained for up to 2000 epochs (or until convergence) using a batch sampling of 2–32, depending on the memory requirements of the models and systems. Note that the models were trained using Tensorflow 2.7.0, but the models were saved using Tensorflow 1.15 for the integration with the GROMOS C++ code (as described in ref. 34 for the HDNNPs).

As a loss function, we used the weighted mean-squared-error (MSE) of the energies and forces

$$L = \frac{1}{N} \sum_i^N (E_i - \tilde{E}_i)^2 + \frac{\omega_0}{3N_{\text{QM}}} \sum_i^{N_{\text{QM}}} \sum_x^3 (F_{ix} - \tilde{F}_{ix})^2 + \frac{\omega_1}{3N_{\text{MM}}} \sum_i^{N_{\text{MM}}} \sum_x^3 (F_{ix} - \tilde{F}_{ix})^2 \quad (14)$$

where N_{QM} is the number of QM particles, N_{MM} is the number of MM particles, and ω_0 and ω_1 are the weight parameters of the gradient contributions. We monitored the loss during the training process and recovered the model with the lowest loss on the validation set after training.

4 Results and discussion

In order to evaluate the different model architectures and parametrizations as well as the training setup, the two simplest test systems from ref. 34, *i.e.*, benzene in water and uracil in water, were used with the same training/validation/test split as in the original publication. The use case here is that the training set is generated from a short initial MD simulation, from which the energies and forces of the subsequent MD steps can be predicted. In the first step, we evaluated the full-QM learning task and a Δ -learning scheme with DFTB^{42,43} as the baseline for the different GCNN architectures. The Δ -learning

scheme simplifies the learning task, and we adopted the same approach as used in ref. 34 for the HDNNPs. For the basic setup of our GCNN model, 128 features per dense layer (n_f), five interaction blocks, and $R_{\text{edge}} = 0.5$ nm were used. For the multi-head self-attention models, we used four heads with $n_k = 32$. Finally, we compare the best GCNN model with the previous HDNNP model in ref. 34 for all five test systems: (i) benzene in water, (ii) uracil in water, (iii) retionic acid in water, (iv) (close to) the transition state of the $S_{\text{N}}2$ reaction of CH_3Cl with Cl^- in water, and (v) (close to) the transition state of SAM with cytosine in water. Note that, for the parametrization and evaluation of the GCNN models and training procedure, we solely used the training and validation sets. The test sets were only taken for the comparison between the final GCNN model with the previous HDNNP model.

4.1 Model architectures

In the first step, we compared the different model architectures in order to evaluate how the global transfer schemes influence the inclusion of the long-range information directly into the model. For this, four different GCNN models were trained: (a) a basic GCNN model, (b) a GCNN model including multi-head self-attention at the post-processing stage (labeled as the “attention GCNN”), (c) a GCNN model including multi-head self-attention in the interaction layers (labeled as the “interaction GCNN”), and (d) a GCNN model including a global state (labeled as the “global GCNN”). The two test systems consist of benzene (apolar molecule) in water with a short MM cutoff radius of $R_c = 0.6$ nm, and uracil (polar molecule) in water with a longer $R_c = 1.4$ nm.

Table 1 shows the mean absolute error (MAE) of the different model architectures for the validation sets of benzene in water and uracil in water. When the models were trained on the full QM energies and forces, we observe for benzene that all models achieve a similar accuracy on the energies, while the basic GCNN outperforms the other architectures for the forces (F_{QM} and F_{MM}). When using the Δ -learning scheme, the accuracy is generally improved, with similar performances of the different architectures (Table 1). Only the interaction GCNN model has a significantly higher MAE for both the energies and the forces. For uracil, we observe similar trends as for benzene

Table 1 Mean absolute error (MAE) on the validation set (2000 frames) of GCNN models containing global information for the test systems benzene in water and uracil in water. For each property, the model with the lowest MAE is marked in bold

GCNN model	Full QM			Δ -Learning		
	E (kJ mol ⁻¹)	F_{QM} (kJ mol ⁻¹ nm ⁻¹)	F_{MM} (kJ mol ⁻¹ nm ⁻¹)	E (kJ mol ⁻¹)	F_{QM} (kJ mol ⁻¹ nm ⁻¹)	F_{MM} (kJ mol ⁻¹ nm ⁻¹)
Benzene						
Basic	2.7	48.7	8.3	1.0	22.9	3.5
Global	2.4	54.1	8.9	1.0	22.7	3.8
Interaction	2.7	52.8	9.7	3.9	36.2	7.3
Attention	2.7	61.7	10.0	1.2	23.7	4.2
Uracil						
Basic	7.3	135.2	5.8	2.2	60.2	1.3
Global	8.1	129.7	6.4	2.3	59.8	1.7
Interaction	8.0	138.4	4.4	2.6	55.5	1.6
Attention	7.4	134.9	7.3	2.5	55.5	1.4



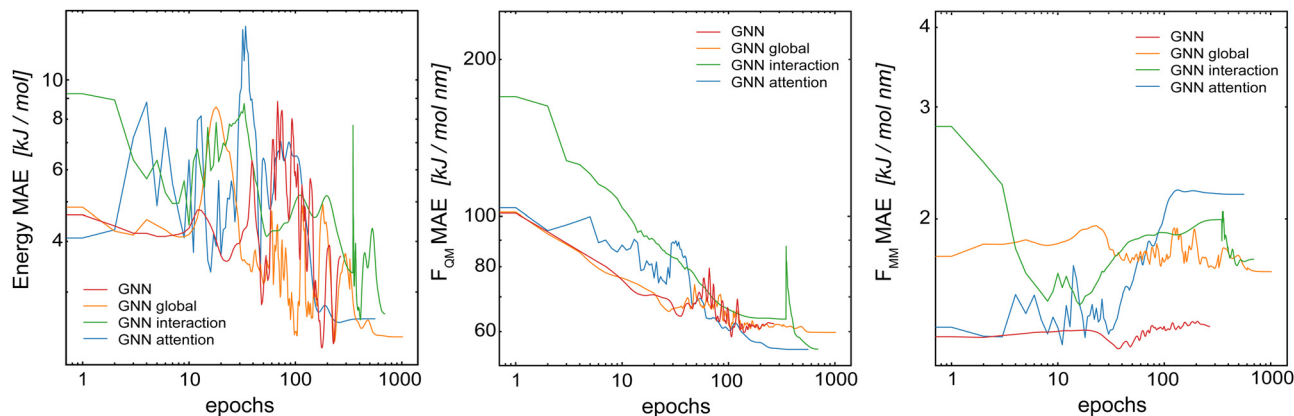


Fig. 3 Mean absolute error (MAE) on the validation set as a function of the training epochs (learning curves) of the Δ -learning GCNN models containing global information for the test system uracil in water. (left): MAE of the energies in the QM zone. (middle): MAE of the forces on the QM particles. (right): MAE of the forces on the MM particles from the QM zone.

(Table 1). Again, no large differences are observed between the models for both the “full QM” and the Δ -learning setup (Fig. 3). In contrast to benzene, also the interaction GCNN models achieve a similar accuracy compared to the other models.

As observed for HDNNPs in ref. 34, the use of a Δ -learning scheme with DFTB as the baseline method reduces the prediction error by about two folds. This is also the case for the GCNN models containing global information. This indicates that the Δ -learning approach is more suited to incorporate the long-range interactions. A reason for this could be that while the GCNNs containing global information can in theory learn long-range interactions, the large number of MM atoms present leads to an exponential increase in possible system configurations. This in turn requires a huge number of training data points in order to accurately capture the long-range interactions. If the size of the training set is limited to a practically useful number (as done here), the training set is evidently not large enough to achieve the desired accuracy.

In conclusion, the Δ -learning scheme leads to a clear and consistent performance improvement, whereas no large differences

are observed between the four different GCNN architectures for both test systems. While the improved GCNN models can reach a slightly lower MAE than the basic GCNN model for some setups, these improvements do not justify their increased model complexity and the subsequently higher computational requirements. For comparison, the basic GCNN model has around 710 000 tunable parameters, while the global GCNN model has around 1 200 000 parameters (870 000 in the interaction GCNN and 840 000 in the attention GCNN). In addition, the interaction GCNN and attention GCNN models require dot products of the query ($Q \in R^{n_{in} \times n_k}$) and key ($K \in R^{n_{in} \times n_k}$) matrices, which are computationally expensive (although the costs can be reduced through the use of multi-head self-attention and $n_{out} = n_k < n_{in}$). For these reasons, we decided to focus on the basic GCNN model and the Δ -learning scheme in the following.

4.2 Model parametrization

For the basic GCNN architecture, we explored different model parametrizations: an edge-cutoff value (R_{edge}) from 0.3 to

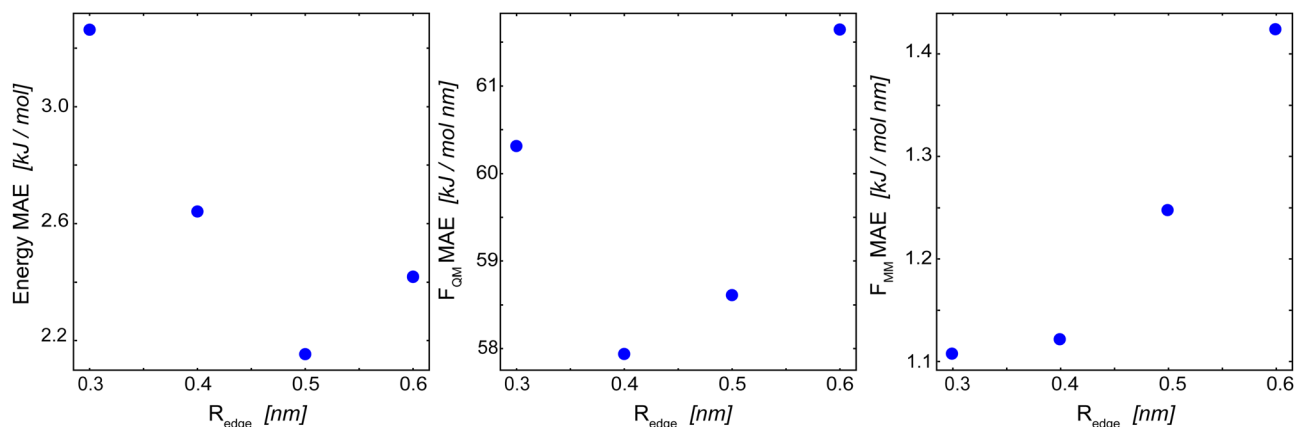


Fig. 4 Influence of R_{edge} on the mean absolute error (MAE) of the Δ -learning basic GCNN model for uracil in water. (left): MAE of the energies in the QM zone. (middle): MAE of the forces on the QM particles. (right): MAE of the forces on the MM particles from the QM zone. The numerical values are given in Table 2.



0.6 nm, a number of features per dense layer (n_f) from 32 to 256, and a model depth (*i.e.*, number of interaction blocks) from 2 to 8. The edge-cutoff value and the model depth directly influence the contributions from long-range interactions, as they both determine (directly and indirectly) the distance up to which atoms contribute to the message update function. The number of features per dense layer, on the other hand, determines the maximum possible complexity of each message update function.

Fig. 4 shows the influence of the edge-cutoff value (R_{edge}) on the model performance for the Δ -learning GCNN model of uracil in water. For the energy, the error is the lowest for a R_{edge} of 0.5 nm. For both force terms (F_{QM} and F_{MM}), however, smaller R_{edge} values of 0.4 nm and 0.3 nm, respectively, give the lowest errors. Note that, for both the energy and the force terms, the error increases again for larger R_{edge} values. The reason for this could be that larger cutoffs are connected with an increase in possible system configurations, requiring in turn more training data points. If the same training set is used, overfitting might occur. For benzene in water, a different trend is observed. Here, a large R_{edge} of 0.6 nm still leads to an improvement in the MAE of the forces (Table 2). This is especially interesting as a different MM cutoff radius is used in the reference QM/MM calculations of benzene and uracil ($R_c = 0.6$ nm and 1.4 nm, respectively). A reason could be that the reduced MM cutoff radius (R_c) leads to a smaller configuration space in the MM-region and thus a GCNN model with a larger R_{edge} is able to sufficiently learn the important contributions.

For the model depth, we observe a similar trend as for the edge-cutoff value. Namely that including the information about further distant nodes (atoms) is not always beneficial for the model performance (Table S1 in the ESI[†]). For the test system uracil in water, the MAE of both the energies and forces is lowest at around 2–4 interaction layers, while up to 6–8 interaction layers are required for the benzene in water test system. The results for the number of features per dense layer (n_f) are given in Table S2 in the ESI[†]. In general, a higher number of features per dense layer does not lead to a lower error, with

convergence at around 64 features per dense layer for uracil in water. For the test system of benzene in water, on the other hand, a more complex model leads again to a decrease in the errors.

4.3 Loss contribution of the forces

Böselt *et al.*³⁴ showed that the relative weighting of the different loss terms (energy loss, F_{QM} loss, and F_{MM} loss) can have a significant influence on the prediction accuracy. Thus, we systematically trained and evaluated GCNN models by varying the relative loss weightings (wE , wF_{QM} , and wF_{MM}). We decided to keep wE constant at 1.0, while changing F_{QM} from 0.001 to 100 and F_{MM} from 0.1 to 1000.

Fig. 5 shows the MAE on the training and validation set of the basic GCNN model with the Δ -learning scheme for the test system of uracil in water when varying wF_{QM} and wF_{MM} . For the training set, we observe the expected behaviour, *i.e.* the error on F_{QM} decreases when the weight wF_{QM} is increased and the same for F_{MM} with wF_{MM} . Interestingly, the results on the validation set are different. While the error on F_{MM} still decreases with increasing wF_{MM} as expected (the bottom right panel in Fig. 5), the MAE on F_{QM} is the smallest for low wF_{QM} and high wF_{MM} (the bottom middle panel in Fig. 5). This observation is further illustrated in Fig. 6, which shows the learning curves for F_{QM} in the training and validation sets (corresponding to the horizontal line at $wF_{\text{MM}} = 100$ and the vertical line at $wF_{\text{QM}} = 0.1$ in the bottom middle panel in Fig. 5).

For the training set, the learning curves show again the expected dependence with F_{QM} and F_{MM} , respectively (top panels in Fig. 6). However, for the validation set, the lowest error of F_{QM} is observed with $wF_{\text{QM}} = 0.1$ and not with $wF_{\text{QM}} = 100$. In general, a relative weighting of $\frac{F_{\text{MM}}}{F_{\text{QM}}}$ around 100–1000 results in the lowest MAE of F_{QM} . Higher relative wF_{QM} values seem to result in over-fitting of F_{QM} . Thus, for the final GCNN model, we use the following relative loss weights: $wE = 1$, $wF_{\text{QM}} = 0.1$, and $wF_{\text{MM}} = 10$. Note that the same observations were made for the Δ -learning GCNN model with benzene in water (Fig. S1 and S2 in the ESI[†]).

4.4 Neighborhood reduction

The main increase in computational requirements for the GCNN models (especially memory requirements when using GPUs during the batched training procedure) comes from the growing number of edges with an increasing size of the edge-cutoff (R_{edge}). For example, $R_{\text{edge}} = 0.4$ nm leads to around 21 000 edges for one of the uracil in water snapshots, while $R_{\text{edge}} = 0.5$ nm already leads to around 38 000 edges. Therefore, we investigated whether the number of edges can be reduced using different neighborhood schemes, without decreasing the model performance. To limit the number of edges around each atom, we explored the k -nearest-neighbour⁵⁵ method with eight (KNN-8) and twelve (KNN-12) neighbours, and the Voronoi–Dirichlet polyhedra (VD).^{56,57} With the same $R_{\text{edge}} = 0.5$ nm, KNN-8 leads to around 12 000 edges for a uracil in water snapshot, KNN-12 to around 18 000 edges, and VD to around

Table 2 Influence of R_{edge} on the mean absolute error (MAE) on the validation set of the basic GCNN models for the test systems benzene in water and uracil in water. For each property, the model with the lowest MAE is marked in bold

R_{edge} (nm)	Δ -Learning		
	E (kJ mol ⁻¹)	F_{QM} (kJ mol ⁻¹ nm ⁻¹)	F_{MM} (kJ mol ⁻¹ nm ⁻¹)
Benzene			
0.3	2.4	27.8	5.5
0.4	1.3	23.7	4.2
0.5	1.0	22.9	3.5
0.6	1.1	21.0	3.1
Uracil			
0.3	3.3	60.4	1.1
0.4	2.6	59.2	1.1
0.5	2.2	60.2	1.3
0.6	2.4	62.9	1.5





Fig. 5 Influence of relative weights (wF_{QM} and wF_{MM}) for the different loss terms (QM forces and MM forces) on the MAE of the energy (left), the forces on the QM particles (middle), and the forces on the MM particles from the QM zone (right). The basic GCNN with the Δ -learning scheme and the test system uracil in water was used. The weight of the energy loss (wE) is kept constant at 1.0. (top): MAE for the training set. (bottom): MAE for the validation set. The color map and the relative size of the points indicate the MAE.

11 000 edges. Thus, the memory requirement for the GCNN model is drastically reduced by all these approaches.

For the following results, we used the same settings for the GCNN models, *i.e.*, $R_{\text{edge}} = 0.5$ nm, $n_f = 128$ with five interaction layers, and loss weightings of $wE = 1.0$, $wF_{QM} = 0.1$, and $wF_{MM} = 10$. Fig. 7 shows the learning curves on the validation set of the Δ -learning GCNN model for the test system uracil in water. While the MAE on the energies is largely unaffected (within the fluctuations of the error) by the choice of the neighborhood scheme (KNN-12 being the closest to the complete model), the errors on the force terms (F_{QM} and F_{MM}) increase substantially with all neighborhood schemes compared to the complete model. The same trends are also observed for the Δ -learning model of benzene in water (Table S3 in the ESI†).

4.5 Data set ordering

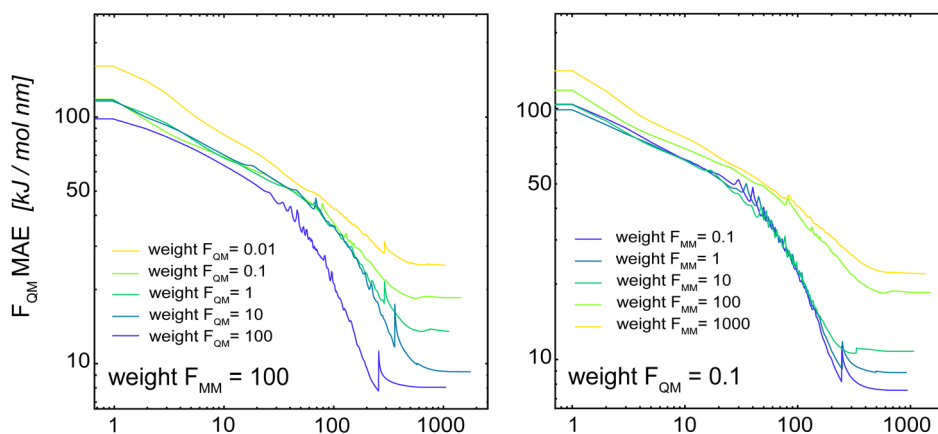
The split into training/validation/test sets was chosen with the future application in MD simulations in mind. The idea is that a short initial QM/MM MD run of the target system can be used

as the training set for subsequent longer (QM)ML/MM MD simulations. To mimic this, the first 70% of the frame from the initial QM/MM trajectory was taken as the training set, the following 20% as the validation set and the final 10% as the test set.³⁴ This leads to a time-based ordering of the frames within the training set. Here, we investigate if this correlated ordering has an effect on the model performance. For this, we compared five models trained with the same training set but using a different frame ordering within the set: (i) original time-based ordering from the MD simulation, (ii–iv) random-ordering using three different random number seeds, and (v) a farthest-point sampling (fps) as described in the ESI.† For all models, we used $R_{\text{edge}} = 0.5$ nm, $n_f = 128$ with five interaction layers, and loss weightings of $wE = 1.0$, $wF_{QM} = 0.1$, and $wF_{MM} = 10$.

Fig. 8 shows the learning curves for the training and validation sets of the basic GCNN model with the Δ -learning scheme trained with different orderings of the data points in the training set. First, the results clearly show that the order of



training-set



validation-set

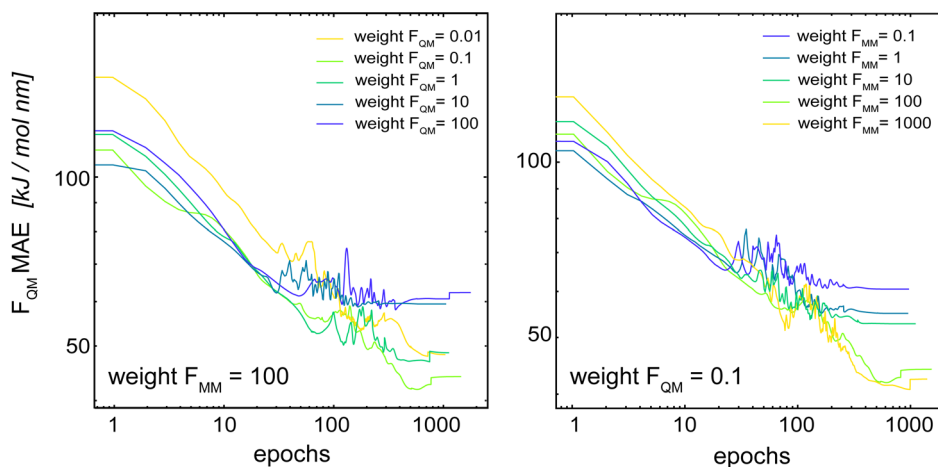


Fig. 6 Learning curves of F_{QM} in the training set (top) and validation set (bottom) when varying the relative loss weightings wF_{QM} (left) and wF_{MM} (right). The basic GCNN with the Δ -learning scheme and the test system uracil in water was used. The weight of the energy loss (wE) is kept constant at 1.0.

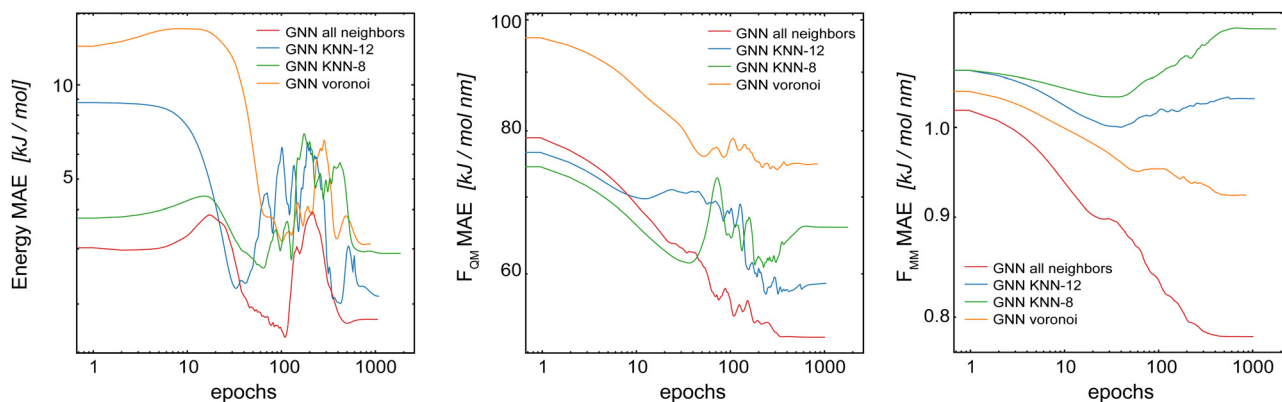


Fig. 7 Learning curves on the validation set of the Δ -learning GCNN model when using different neighborhood selection schemes. The test system uracil in water was used.

the training data points does indeed affect the model performance. For all three properties (energies and forces), the error for the training set is clearly smallest with the fps-ordering (orange line) or MD-ordering (purple line). Interestingly, using

an fps-ordering of the training data points results in a higher MAE for the validation set, indicating that the fps-ordering leads to an over-fitting of the model for this setup. The situation is different for the MD-ordering, which leads to



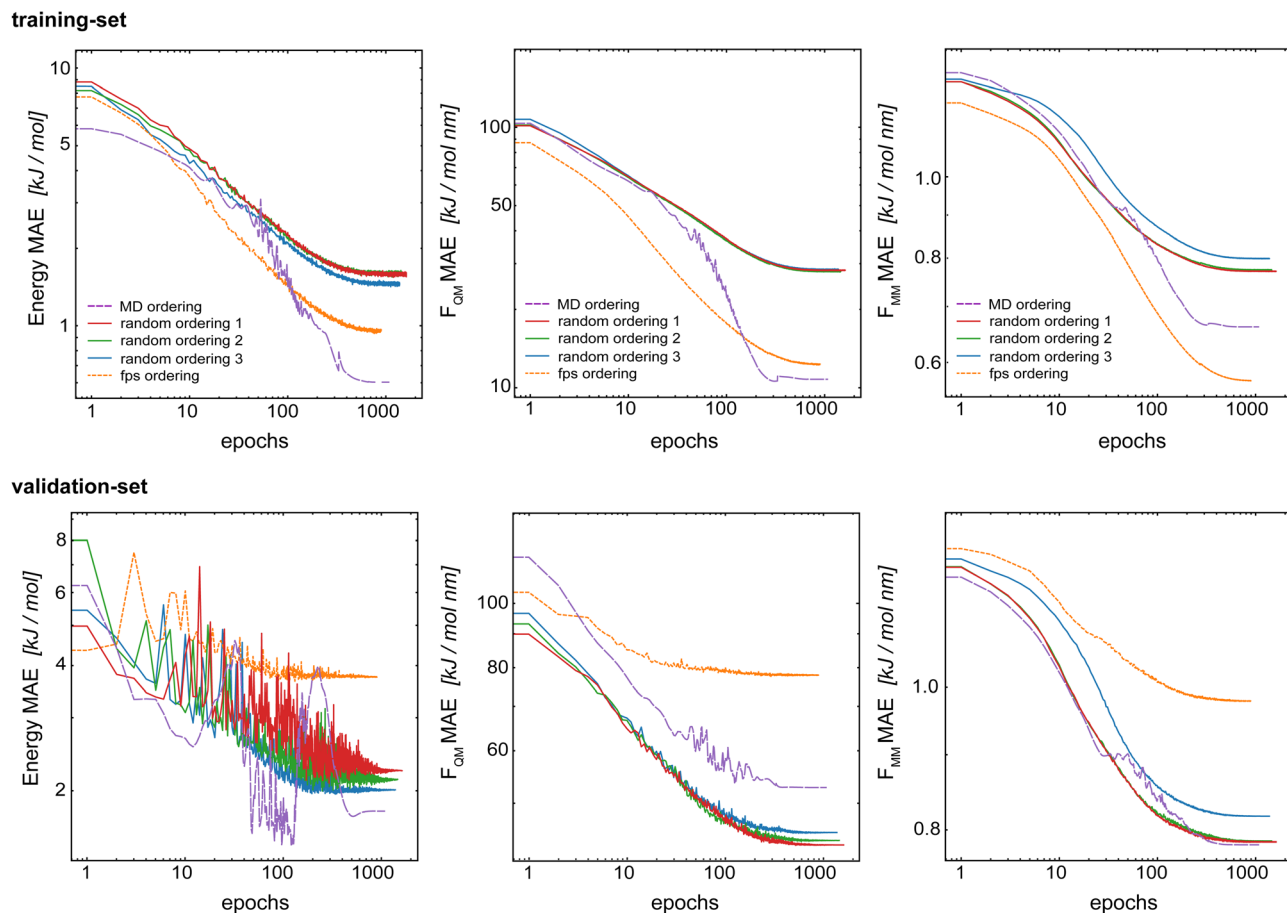


Fig. 8 Learning curves for the training set (top) and validation set (bottom) when varying the order of the data points in the training set. The Δ -learning GCNN for uracil in water is shown.

comparable results on the validation set as the random-ordering, except for F_{QM} . Here, the accuracy is between the random and the fps-ordering. For the test system of benzene in water (Fig. S3 in the ESI[†]), the same trends are observed for fps-ordering on the training and validation sets, while the MD-ordering shows again a comparable accuracy as the random-orderings for both the training and the validation sets. However, in all cases, the random-orderings converge earlier (at around 1000 epochs) compared to the MD-ordering (around 2000 epochs). Thus, using a random-ordering of the MD frames in the training set seems to be beneficial by reducing the training costs and potentially increasing robustness.

4.6 Comparison to HDNNPs

As a final evaluation, we compare the GCNN models with the previously developed HDNNPs,^{34,41} which use the same Δ -learning scheme with DFTB^{42,43} as the baseline, for all five test systems: (i) benzene in water, (ii) uracil in water, (iii) retionic acid in water, (iv) (close to) the transition state of the $S_{\text{N}}2$ reaction of CH₃Cl with Cl⁻ in water, and (v) (close to) the transition state of SAM with cytosine in water. The model performance is compared on the training set (7000 frames), validation set (2000 frames), and test set (1000 frames). Note that the models were solely developed using the training and

validation sets, and are only evaluated on the test set at this stage. Furthermore, hyperparameters were only tuned on the two test systems benzene and uracil in water. For all models, we used $R_{\text{edge}} = 0.5$ nm, $n_{\text{f}} = 128$ with four (systems i and ii) or five (systems iii–v) interaction layers, and loss weightings of $wE = 1.0$, $wF_{\text{QM}} = 0.1$, and $wF_{\text{MM}} = 10$.

Fig. 9 compares the learning curves and final MAE between the HDNNP and the GCNN models for benzene in (top panels) and uracil in water (bottom panels). For both systems, the learning curves for the GCNN validation and test sets show a very similar performance. This indicates that the models are not overfitted and are able to generalize to new data points within the same MD trajectory. For the HDNNP models, we observe a similar behaviour. However, the MAE values for the test set of benzene in water are consistently above the MAE for the validation set. For this system, we can directly compare the effect of the training set size between the HDNNP and GCNN models. While both models show a similar MAE for small training set sizes, the GCNN model outperforms the HDNNP model when all training data points are used. A possible reason for this observation is that the GCNN model has to learn the descriptor for the atomic environments and would thus profit from a larger training set. This is also indicated by the fact that the learning curves of the GCNN models are still continuously



Benzene in water



Uracil in water



Fig. 9 Learning curves as a function of the training-set size of the Δ -learning models for benzene in water (top) and uracil in water (bottom). The dark blue circles show the GCNN validation MAE, the dark red circles show the HDNNP validation MAE, the light blue triangles show the GCNN test MAE, and the light red triangles show the HDNNP test MAE. The numerical values are given in Table 3.

decreasing even at 7000 training structures. For the system of uracil in water, the final MAE values are comparable for all models. Here, we note that the benzene test system has a shorter MM-cutoff radius (R_c), which might simplify the learning task and thus a similar behaviour might be observed for uracil at a larger training set size.

Table 3 shows the MAE values on the training, validation, and test sets for all five test systems. Overall, the GCNN models perform similar or slightly better than the HDNNP models. A major exception to this is the SAM/cyt in water system. Here, the GCNN model reaches the same accuracy as the HDNNP

model on the training set, but the MAE on the validation and test set is two to five times higher than with the HDNNP model. This may be because the GCNN models were parametrized and benchmarked on the simpler systems (*i.e.*, benzene and uracil in water) and there is less transferability of the hyperparameters than that with the HDNNP.

Up to this point, we have evaluated the performance of the GCNN models in terms of the MAE on reasonable validation and test sets. Bösel *et al.*³⁴ have already emphasized that it is important to test a ML model for the intended application, which in our case are (QM)ML/MM MD simulations. While rare

Table 3 Mean absolute error (MAE) on the training set (7000 frames), validation set (2000 frames), and test set (1000 frames) of the Δ -learning HDNNP models and the Δ -learning GCNN models for all five test systems: (i) benzene in water, (ii) uracil in water, (iii) retionic acid in water, (iv) (close to) the transition state of the S_N2 reaction of CH_3Cl with Cl^- in water, and (v) (close to) the transition state of SAM with cytosine in water. For each test system, the model with the lowest MAE is marked in bold

System	HDNNP			GCNN		
	E (kJ mol ⁻¹)	F_{QM} (kJ mol ⁻¹ nm ⁻¹)	F_{MM} (kJ mol ⁻¹ nm ⁻¹)	E (kJ mol ⁻¹)	F_{QM} (kJ mol ⁻¹ nm ⁻¹)	F_{MM} (kJ mol ⁻¹ nm ⁻¹)
Benzene	1.8/1.6/3.2	23.7/22.8/29.4	3.5/3.0/4.7	0.4/0.7/0.7	8.6/14.7/15.3	1.3/1.9/2.0
Uracil	1.2/2.8/ 2.0	34.3/ 45.9/48.1	1.1/1.0/1.0	0.8/1.9/2.1	14.1/51.0/51.0	0.7/0.8/0.8
Chloroform	1.8/1.7/2.2	24.9/29.5/30.7	1.0/1.0/1.0	0.5/1.1/1.2	5.3/21.4/22.1	0.5/0.6/0.6
Retionic acid	3.9/ 4.4 /—	43.8/44.8/—	1.1/1.1/—	0.6/4.4/4.4	20.5/37.0/37.0	0.7/0.9/0.9
SAM/cyt	6.3/ 8.5 /—	74.8/74.6 /—	2.3/2.3 /—	5.6/21.1/28.8	79.7/130.4/130.4	11.9/12.0/12.2



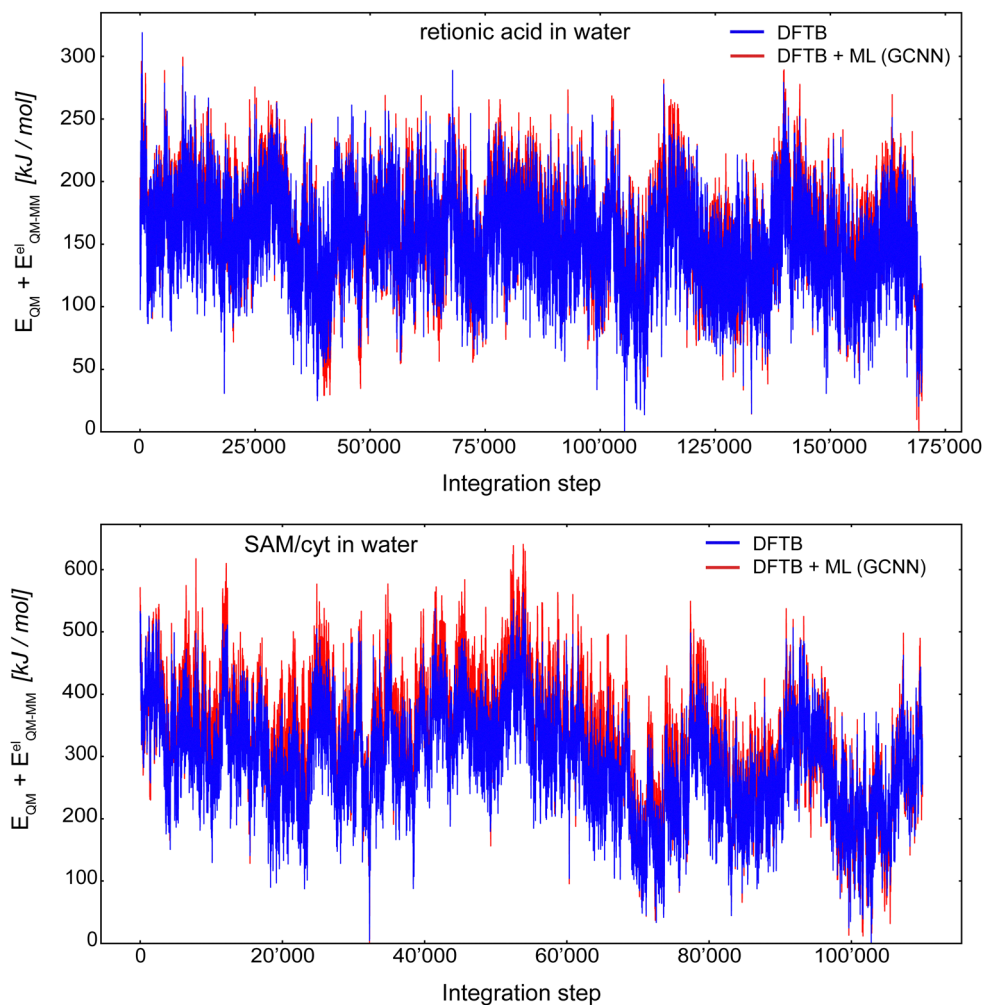


Fig. 10 MD simulation of retinoic acid in water (top) and SAM/cyt in water (bottom) using an integration time step of 0.5 fs. The energy ($E_{\text{QM}} + E_{\text{QM-MM}}^{\text{el}}$) trajectory is extracted from 170 000 (retinoic acid) and 110 000 (SAM/cyt) consecutive steps performed by the DFTB+GCNN model. The first 200 steps were discarded as equilibration. The pure DFTB energy is shown in blue and the GCNN Δ -correction is shown in red. Note that the GCNN model was only trained on the initial 7000 QM/MM MD simulation steps (not shown here) and that no adaptive on-the-fly re-training of the GCNN model was performed.

outliers get averaged in the MAE assessment, they might be less tolerable in an actual simulation, where the results of the next step depend directly on the results of the previous step. For this reason, it is crucial to test the performance of the developed (QM)ML/MM models in a prospective MD simulation. We performed therefore (QM)ML/MM MD simulations using the Δ -learning GCNN model for the two test systems with larger conformational flexibility as in ref. 34: (i) retinoic acid and (ii) (close to) the transition state of SAM with cytosine in water. Note that the models were only trained on the initial 7000 steps of the QM/MM MD simulations and no adaptive re-training during the (QM)ML/MM production runs was performed. Shen and Yang³⁵ have shown that adaptive neural networks can use on-the-fly corrections of the model to further improve the model performance. However, this comes with an increase in the training cost of the model and changing energies/forces (*i.e.*, the estimated properties of a configuration may not be the same at the beginning *versus* the end of the simulation).

Fig. 10 shows the MD energy ($E_{\text{QM}} + E_{\text{QM-MM}}^{\text{el}}$) trajectories for retinoic acid in water (top panel) and SAM/cyt in water (bottom panel) using an integration step of 0.5 fs. The (QM)ML/MM MD simulation of retinoic acid in water was carried out for 170 000 consecutive steps, while the simulation of SAM/cyt in water was performed for 110 000 consecutive steps. For both systems, the first 200 steps were discarded as equilibration. Thus, it was possible to propagate these systems stably using the Δ -learning GCNN model based on a fraction of the otherwise required QM/MM MD simulation steps. For both systems, the energy fluctuations during the (QM)ML/MM MD simulation were comparable to the ones with the Δ -learning HDNNP model: For retinoic acid, $\sigma_{\text{DFTB+GCNN}} = 37.1 \text{ kJ mol}^{-1}$ and $\sigma_{\text{DFTB+HDNNP}} = 55.2 \text{ kJ mol}^{-1}$, and for SAM/cyt, $\sigma_{\text{DFTB+GCNN}} = 94.9 \text{ kJ mol}^{-1}$ and $\sigma_{\text{DFTB+HDNNP}} = 73.3 \text{ kJ mol}^{-1}$. Of course, the possibility that in a longer simulation an ill-represented structure may be encountered cannot be fully excluded. Interestingly, even though the accuracy of the GCNN model on the validation/test sets of the SAM/cyt in water was considerably lower than that with



the HDNNP model, we still observe a stable and robust simulation over 110 000 (QM)ML/MM MD steps based solely on the initial 7000 QM/MM MD simulation steps.

5 Conclusion

In this work, we investigated the use of GCNN models in (QM)ML/MM MD simulations for condensed-phase systems with DFT accuracy for the QM subsystem, and compared the results with the previously developed HDNNPs for the same systems. In the first step, we evaluated different GCNN architectures capable of incorporating long-range effects, with and without a Δ -learning scheme. While no large improvements could be found with the more complex GCNN architectures, the Δ -learning GCNN using DFTB as the baseline yielded the most accurate description of the energies and forces. Based on these observations, we focused on the basic GCNN with Δ -learning for the remainder of the study. Next, we assessed the influence of different parameters on the performance of the GCNN models. Here, we found that the inclusion and the correct weighting of the QM and MM gradients in the loss function were crucial to improve the model performance. Interestingly, we observed that increasing the loss weights of the QM gradients does not lead to an improved accuracy when predicting them due to overfitting. Instead, the relative loss weights of $w_E = 1$, $w_{F_{\text{QM}}} = 0.1$ and $w_{F_{\text{MM}}} = 10$ provided the model with the best ability to generalize new data points.

In order to reduce the computational requirements of the GCNN models, we also investigated different neighborhood reduction schemes in the creation of the GCNN edges. While these schemes decrease indeed the computational costs without significantly affecting the accuracy of the QM energies, the performance on the QM and MM forces is clearly worse. Thus, we do not recommend using these schemes in (QM)ML/MM MD simulations, as the predicted forces are directly used to propagate the system in time.

An interesting observation was made regarding the order of the data points in the training set. In order to mimic the future application, the training set was chosen as the first 70% of frames from a QM/MM trajectory. This leads to a time-based ordering of the training data points. When using a random ordering of the training set (without changing the actual training/validation/test split!), a similar model performance was reached, but the random ordering converges faster and thus leads to a reduction in the required training time.

Finally, we compared the Δ -learning HDNNP and Δ -learning GCNN models with each other for five different test systems in water. While both models perform at a similar accuracy, the GCNN model reaches slightly lower MAE values for most of the five test systems. The Δ -learning GCNN model can also be used to perform stable (QM)ML/MM MD simulations as the corresponding HDNNP model. However, the two model types differ drastically in their architecture and come with advantages and disadvantages, which should be considered when choosing an adequate model for a (QM)ML/MM simulation. The symmetry

functions in HDNNPs include a cosine term for all the inter-atomic angles within an atomic environment, which results in a computational scaling of $O(N^3)$, where N is the combined number of atoms in the QM and MM environments. The GCNN model, on the other hand, only depends on the edge-update operations to describe the atomic environments resulting in a computational scaling of $O(N^2)$. Note that, by using a finite cutoff for the atomic environments and edge-update operations, the scaling can be reduced to $O(N(\log N)^2)$ and $O(N \log N)$, respectively. Additionally, the number of angle terms in the HDNNP symmetry functions scales exponentially with the number of element types (which can be partially resolved by introducing weighted symmetry functions), whereas the GCNN atom types are encoded directly in an embedding vector and do not directly change the scaling of the model. Another important difference is that HDNNPs are based upon individual neural-network potentials (NNPs) for each atomic environment, whose evaluation can be easily distributed over multiple GPUs and the memory requirement for each individual NNP does not increase drastically with the increasing QM or MM system size. However, in contrast, GCNNs are based upon iterative and overlapping message passing operations and thus, a distributed evaluation is not as straightforward. Additionally, for each message passing operation within the GCNN graph, all edges have to be transformed using a dense layer. Therefore, the memory requirement for the GCNN evaluation scales as $O(N_e n_f)$, where N_e is the number of edges. This means that the memory increases with the increasing QM or MM system size. Note that this can be critical when using a batched training procedure on a GPU.

In summary, the HDNNP based Δ -learning (QM)ML/MM setup appears to be best suited for condensed-phase systems with a limited number of different element types in the QM zone and the surrounding MM zone. Examples are the reactions of organic molecules in a solvent – similar to the five test systems investigated in this paper – or reactions at the interface of a mono-atomic surface. In these cases, the limited number of symmetry functions and the parallelizability over multiple GPUs/CPUs favor the HDNNP setup. On the other hand, the GCNN based Δ -learning (QM)ML/MM setup is most useful for condensed-phase systems with a larger number of element types, as the model scaling of the GCNN is much less affected by the number of element types than the HDNNP. Examples are the reactions within the active site of a metalloenzyme, proteins with covalently bound or interacting ligands, reactions involving organometallic catalysts, or reactions catalyzed at the interface of a poly-atomic surface.

Data and software availability

All structures used in the training, validation and test sets were taken from ref. 34 (freely available at <https://www.research-collection.ethz.ch/handle/20.500.11850/512374>). The GCNNs were implemented using Tensorflow/keras.⁵⁰ The models were trained in Python and then exported to the C++ GROMOS code as described in ref. 34 for the HDNNPs.



Conflicts of interest

There are no conflicts to declare.

Acknowledgements

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References

- 1 S. Riniker, *J. Chem. Inf. Model.*, 2018, **58**, 565–578.
- 2 P. S. Nerenberg and T. Head-Gordon, *Curr. Opin. Struct. Biol.*, 2018, **49**, 129–138.
- 3 A. Warshel and M. Levitt, *J. Mol. Biol.*, 1976, **103**, 227–249.
- 4 A. J. Mulholland, P. D. Lyne and M. Karplus, *J. Am. Chem. Soc.*, 2000, **122**, 534–535.
- 5 H. M. Senn and W. Thiel, *Angew. Chem., Int. Ed.*, 2009, **48**, 1198–1229.
- 6 G. Groenhof, *Methods Mol. Biol.*, 2013, **924**, 43–66.
- 7 H. Lin and D. G. Truhlar, *Theor. Chem. Acc.*, 2007, **117**, 185–199.
- 8 C. Panosetti, A. Engelmann, L. Nemeč, K. Reuter and J. T. Margraf, *J. Chem. Theory Comput.*, 2020, **16**, 2181–2191.
- 9 E. Brunk and U. Röthlisberger, *Chem. Rev.*, 2015, **115**, 6217–6263.
- 10 L. W. Chung, W. M. Sameera, R. Ramozzi, A. J. Page, M. Hatanaka, G. P. Petrova, T. V. Harris, X. Li, Z. Ke, F. Liu, H. B. Li, L. Ding and K. Morokuma, *Chem. Rev.*, 2015, **115**, 5678–5796.
- 11 A. S. Christensen, T. Kubař, Q. Cui and M. Elstner, *Chem. Rev.*, 2016, **116**, 5301–5337.
- 12 J. J. Stewart, *J. Mol. Model.*, 2013, **19**, 1–32.
- 13 M. Rupp, A. Tkatchenko, K.-R. Müller and O. A. Von Lilienfeld, *Phys. Rev. Lett.*, 2012, **108**, 1–5.
- 14 J. Behler, *Phys. Chem. Chem. Phys.*, 2011, **13**, 17930–17955.
- 15 K. Hansen, G. Montavon, F. Biegler, S. Fazli, M. Rupp, M. Scheffler, O. A. von Lilienfeld, A. Tkatchenko and K. R. Müller, *J. Chem. Theory Comput.*, 2013, **9**, 3404–3419.
- 16 S. Lorenz, A. Groß and M. Scheffler, *Chem. Phys. Lett.*, 2004, **395**, 210–215.
- 17 J. Behler, *J. Phys.: Condens. Matter*, 2014, **26**, 183001.
- 18 J. S. Smith, B. Nebgen, N. Lubbers, O. Isayev and A. E. Roitberg, *J. Chem. Phys.*, 2018, **148**, 241733.
- 19 J. Behler, *J. Chem. Phys.*, 2016, **145**, 170901.
- 20 J. S. Smith, O. Isayev and A. E. Roitberg, *Chem. Sci.*, 2017, **8**, 3192–3203.
- 21 A. Grisafi, J. Nigam and M. Ceriotti, *Chem. Sci.*, 2021, **12**, 2078–2090.
- 22 K. T. Schütt, P. J. Kindermans, H. E. Saucedo, S. Chmiela, A. Tkatchenko and K. R. Müller, *Adv. Neural Inf. Process. Syst.*, 2017, 991–1001.
- 23 O. T. Unke and M. Meuwly, *J. Chem. Theory Comput.*, 2019, **15**, 3678–3693.
- 24 D. P. Kovács, C. Van Der Oord, J. Kucera, A. E. A. Allen, D. J. Cole, C. Ortner and G. Csányi, *J. Chem. Theory Comput.*, 2021, **17**, 7696–7711.
- 25 V. Zaverkin and J. Kästner, *J. Chem. Theory Comput.*, 2020, **16**, 5410–5421.
- 26 T. W. Ko, J. A. Finkler, S. Goedecker and J. Behler, *Nat. Commun.*, 2021, **12**, 398.
- 27 Z. Li, J. R. Kermode and A. De Vita, *Phys. Rev. Lett.*, 2015, **114**, 1–5.
- 28 B. Anderson, T. S. Hy and R. Kondor, *Adv. Neural Inf. Process. Syst.*, 2019, **32**, 14537–14546.
- 29 T. Xie and J. C. Grossman, *Phys. Rev. Lett.*, 2018, **120**, 145301.
- 30 A. Grisafi and M. Ceriotti, *J. Chem. Phys.*, 2019, **151**, 204105.
- 31 O. T. Unke, S. Chmiela, M. Gastegger, K. T. Schütt, H. E. Saucedo and K.-R. Müller, *Nat. Commun.*, 2021, **12**, 7273.
- 32 C. Chen, W. Ye, Y. Zuo, C. Zheng and S. P. Ong, *Chem. Mater.*, 2019, **31**, 3564–3572.
- 33 R. Ramakrishnan, P. O. Dral, M. Rupp and O. A. Von Lilienfeld, *J. Chem. Theory Comput.*, 2015, **11**, 2087–2096.
- 34 L. Bösel, M. Thürlemann and S. Riniker, *J. Chem. Theory Comput.*, 2021, **17**, 2641–2658.
- 35 L. Shen and W. Yang, *J. Chem. Theory Comput.*, 2018, **14**, 1442–1455.
- 36 J. Zeng, T. J. Giese, S. Ekesan and D. M. York, *J. Chem. Theory Comput.*, 2021, **17**, 6993–7009.
- 37 K. T. Schütt, O. T. Unke and M. Gastegger, *Proceedings of the 38th International Conference on Machine Learning*, 2021, vol. 139, pp. 9377–9388.
- 38 J. Klicpera, J. Groß and S. Günnemann, 2020, arXiv:2003.03123.
- 39 V. G. Satorras, E. Hoogeboom and M. Welling, 2021, arXiv:2102.09844.
- 40 A. Vaswani, N. Shazeer, N. Parmar, J. Uszkoreit, L. Jones, A. N. Gomez, L. Kaiser and I. Polosukhin, *Adv. Neural Inf. Process. Syst.*, 2017, 5999–6009.
- 41 J. Behler, *J. Chem. Phys.*, 2011, **134**, 074106.
- 42 M. Elstner, *Theor. Chem. Acc.*, 2006, **116**, 316–325.
- 43 B. Hourahine, B. Aradi, V. Blum, F. Bonafé, A. Buccheri, C. Camacho, C. Cevallos, M. Y. Deshayé, T. Dumitric, A. Dominguez, S. Ehlert, M. Elstner, T. Van Der Heide, J. Hermann, S. Irle, J. J. Kranz, C. Köhler, T. Kowalczyk, T. Kubař, I. S. Lee, V. Lutsker, R. J. Maurer, S. K. Min, I. Mitchell, C. Negre, T. A. Niehaus, A. M. Niklasson, A. J. Page, A. Pecchia, G. Penazzi, M. P. Persson, J. Åezáč, C. G. Sánchez, M. Sternberg, M. Stöhr, F. Stuckenberg, A. Tkatchenko, V. W. Yu and T. Frauenheim, *J. Chem. Phys.*, 2020, **152**, 124101.
- 44 R. E. Buló, B. Ensing, J. Sikkema and L. Visscher, *J. Chem. Theory Comput.*, 2009, **5**, 2212–2221.
- 45 J. Gilmer, S. S. Schoenholz, P. F. Riley, O. Vinyals and G. E. Dahl, *34th International Conference on Machine Learning, ICML 2017*, 2017, vol. 3, pp. 2053–2070.



- 46 G. Corso, L. Cavalleri, D. Beaini, P. Liò and P. Velickovic, *Adv. Neural Inf. Process. Syst.*, 2020, **33**, 13260–13271.
- 47 N. Schmid, C. D. Christ, M. Christen, A. P. Eichenberger and W. F. van Gunsteren, *Comput. Phys. Commun.*, 2012, **183**, 890–903.
- 48 K. Meier, N. Schmid and W. F. V. Gunsteren, *J. Comput. Chem.*, 2012, **33**, 2108–2117.
- 49 F. Neese, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2018, **8**, 4–9.
- 50 M. Abadi, A. Agarwal, P. Barham, E. Brevdo, Z. Chen, C. Citro, G. S. Corrado, A. Davis, J. Dean, M. Devin, S. Ghemawat, I. Goodfellow, A. Harp, G. Irving, M. Isard, Y. Jia, R. Jozefowicz, L. Kaiser, M. Kudlur, J. Levenberg, D. Mané, R. Monga, S. Moore, D. Murray, C. Olah, M. Schuster, J. Shlens, B. Steiner, I. Sutskever, K. Talwar, P. Tucker, V. Vanhoucke, V. Vasudevan, F. Viégas, O. Vinyals, P. Warden, M. Wattenberg, M. Wicke, Y. Yu and X. Zheng, *TensorFlow: Large-Scale Machine Learning on Heterogeneous Systems*, 2015.
- 51 I. G. Tironi, R. Sperb, P. E. Smith and W. F. van Gunsteren, *J. Chem. Phys.*, 1995, **102**, 5451–5459.
- 52 P. Ramachandran, B. Zoph and Q. V. Le, 2017, arXiv: 1710.05941.
- 53 J.-b. Cordonnier and A. Loukas, 2020, arXiv:2006.16362.
- 54 D. P. Kingma and J. Ba, *Adam: A Method for Stochastic Optimization*, 2017.
- 55 E. Fix and J. L. Hodges, *Int. Stat. Rev.*, 1989, **57**, 238.
- 56 P. Niggli, *XXIV. Die topologische Strukturanalyse. I.*
- 57 V. A. Blatov, *Crystal. Rev.*, 2004, **10**, 249–318.

