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Homogeneous in-plane WSe₂ P–N junctions for advanced optoelectronic devices†

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Conventional doping schemes of silicon (Si) microelectronics are incompatible with atomically thick two-dimensional (2D) transition metal dichalcogenides (TMDCs), which makes it challenging to construct high-quality 2D homogeneous p–n junctions. Herein, we adopt a simple yet effective plasma-treated doping method to seamlessly construct a lateral 2D WSe₂ p–n homojunction. WSe₂ with ambipolar transport properties was exposed to O₂ plasma to form WO_x on the surface in a self-limiting process that induces hole doping in the underlying WSe₂ via electron transfer. Different electrical behaviors were observed between the as-exfoliated (ambipolar) region and the O₂ plasma-treated (p-doped) region under electrostatic modulation of the back-gate bias (V_{BG}), which produces a p–n in-plane homojunction. More importantly, a small contact resistance of 710 $\Omega \mu\text{m}$ with a p-doped region transistor mobility of $\sim 157 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ was achieved due to the transformation of Schottky contact into Ohmic contact after plasma treatment. This effectively avoids Fermi-level pinning and significantly improves the performance of photodetectors. The resultant WSe₂ p–n junction device thus exhibits a high photoresponsivity of $\sim 7.1 \times 10^4 \text{ mA W}^{-1}$ and a superior external quantum efficiency of $\sim 228\%$. Also, the physical mechanism of charge transfer in the WSe₂ p–n homojunction was analyzed. Our proposed strategy offers a powerful route to realize low contact resistance and high photoresponsivity in 2D TMDC-based optoelectronic devices, paving the way for next-generation atomic-thickness optoelectronics.

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1. Introduction

Two-dimensional (2D) materials, such as transition metal dichalcogenides (TMDCs), hold promise as fundamental building blocks for next-generation flexible electronic and optoelectronic devices because of their tunable bandgap, high mobility, good chemical stability, and ability to form high-quality interfaces in van der Waals (vdW) heterostructures.^{1–5} Among all TMDCs, WSe₂ is an intriguing 2D material that offers bipolar transport characteristics with thickness-dependent band gaps of 1.7 eV (direct, monolayer) and approxi-

mately 1.2 eV (indirect, multilayer), as well as theoretical electron and hole mobilities of about 250 and 270 $\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, respectively.^{6–8} Owing to its unique physical properties, WSe₂ has opened the door to versatile electronic and optoelectronic devices with excellent optoelectronic performance. In recent years, new artificial structures based on van der Waals (vdW) vertical heterostructures or lateral p–n heterojunctions have attracted great interest in the application of WSe₂ optoelectronic devices.^{9–14} For instance, Jo and co-workers demonstrated that the photoresponsivity of WSe₂/h-BN-based p–n heterojunction photodetectors was remarkably improved by a triphenylphosphine n-doping method.¹⁵ Guo and co-workers reported that a WSe₂–ZnO p–n heterojunction photodetector fabricated by combining p-type WSe₂ and n-type ZnO exhibited an ultra-high photoresponsivity of $4.83 \times 10^3 \text{ A W}^{-1}$ under 405 nm light illumination.¹⁶ Liu *et al.* reported that a photodetector based on a WSe₂–Bi₂Te₃ p–n heterojunction can give rise to a fast response time of $\sim 210 \mu\text{s}$ and a high photoresponsivity of $\sim 20.5 \text{ A W}^{-1}$ under 633 nm illumination.¹⁷ These results have demonstrated that the successful construction of heterojunctions can significantly improve the photo-response characteristics of TMDC-based devices, which, however, still suffer from unavoidable residues and complex flake alignment processes, *i.e.*, difficult to localize the target

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material onto another flake accurately.¹⁸ In addition, the presence of a vdW gap between junctions may inhibit carrier charge transfer, which is generally undesirable for the device performance. As an alternative strategy, homogeneous junctions offer intrinsic advantages over heterogeneous junctions by providing clean and self-aligned interfaces.

Since the Fermi level (E_F) of WSe₂ resides in the middle of the band gap, its ambipolar nature allows for electron or hole dominant transport *via* selective doping, which can be used as an ideal candidate for the implementation of p–n photodiodes, photovoltaic cells, and complementary logic systems.^{19–23} However, the *in situ* synthesis of high-quality homogeneous WSe₂ p–n junction-based photodetectors still faces many hurdles, such as the complex fabrication process, slow photoresponsivity, low photodetectivity, *etc.* On the other hand, the formation of a Schottky barrier (SB) at the WSe₂–metal interface inevitably results in a high contact resistance (R_c).^{24,25} Furthermore, the physical metal deposition process may cause potential lattice disorders or even local destruction of the material, leading to Fermi-level pinning or poor band alignment, ultimately limiting the efficiency of photodiodes.^{26–28} Consequently, carrier transport can only be carried out by tunneling through thin SBs, resulting in a low current density. To overcome the inherent performance limits of 2D TMDCs and to improve the contact problems on WSe₂ functional devices, various doping schemes, work function engineering techniques, electrical tuning techniques, *etc.* have been explored.^{29–35} However, these methods are relatively complicated and require high temperature or vacuum process conditions, which greatly hinder their practical application in constructing high-performance WSe₂-based photodetectors. Most recently, transition metal oxides (TMOs) used in organic electronics and 2D material-based electronic devices have been demonstrated as effective p-type contacts and dopants.^{36,37} This doping scheme has been widely developed and applied in

complementary metal–oxide–semiconductor (CMOS) technology. Despite many encouraging outcomes so far, the fabrication of a 2D TMDC p–n homojunction with low R_c is still limited by technical challenges.

In this work, we demonstrate an efficient and reliable approach by performing O₂ plasma treatment to easily permit the amplification of hole transport, and thus achieve *in situ* p-type semiconductor characteristics in WSe₂. The under-stoichiometric oxidation of WSe₂ into WO_x highly induces hole doping in the neighboring (or underlying) WSe₂ *via* electron transfer from the underlying WSe₂ to the top oxidized WO_x layer. Moreover, our results revealed that the transformation of the SB to Ohmic contact arising from plasma treatment greatly increases the hole mobility from ~22 to ~157 cm² V⁻¹ s⁻¹ in a p-type doped WSe₂ transistor. As a result, the as-fabricated in-plane p–n homojunction exhibits a superior external quantum efficiency of ~228%, an excellent photoresponsivity of ~7.1 × 10⁴ mA W⁻¹, and a photodetectivity of ~3 × 10³ Jones under 532 nm light illumination. The resultant device performance of our method proves it to be an ideal strategy for constructing high-performance p–n junction optoelectronic devices by using self-oxidizing 2D TMDCs, which may be applied to other TMDCs for future multifunctional electronic and optoelectronic applications.

2. Results and discussion

Fig. 1(a) shows the schematic illustration of a pristine and plasma-treated WSe₂. The WSe₂ flake was treated with physical O₂ plasma bombardment where the Se atoms were replaced by O atoms, which then forms a WO_x layer on the surface of the WSe₂ flake. The same sample was repeatedly subjected to plasma treatment, and the transfer characteristics of the as-fabricated WSe₂ FET at cumulative plasma durations of 5 to 120 s were measured (Fig. S1†). A change in the majority of carriers, from electrons to holes (n-dominant ambipolar to strong p-type), is observed with an on/off ratio reaching 10⁸ after 120 s of O₂ plasma treatment as shown in Fig. S1.† As a result, the topmost layers of the WSe₂ flake were completely oxidized into WO_x, consistent with previous reports where WO_x is used as an effective doping layer.²³ In addition, the thickness of the WSe₂ flakes increased by ~1 nm after O₂ plasma treatment (Fig. 1(b)), which is an indication of the formation of a thin layer of WO_x atop the WSe₂ flake. The oxidation depth of the WSe₂ flake was subsequently assessed using the secondary-ion mass spectrometry (SIMS) technique. As shown in Fig. S2,† oxidation is found to be a self-limiting process. The depth of the oxide layer (WO_x, including WO⁻, WO₂⁻, and WO₃⁻) is determined to be approximately 1.2 nm, which is restricted only in the very top layer, therefore offering an easy and efficient control of the oxide thickness. To further validate the formation of the WO_x layer, the surface composition of the plasma-treated WSe₂ was further analyzed using X-ray photoelectron spectroscopy (XPS), as presented in Fig. 1(c). The XPS core spectrum of the pristine WSe₂ shows



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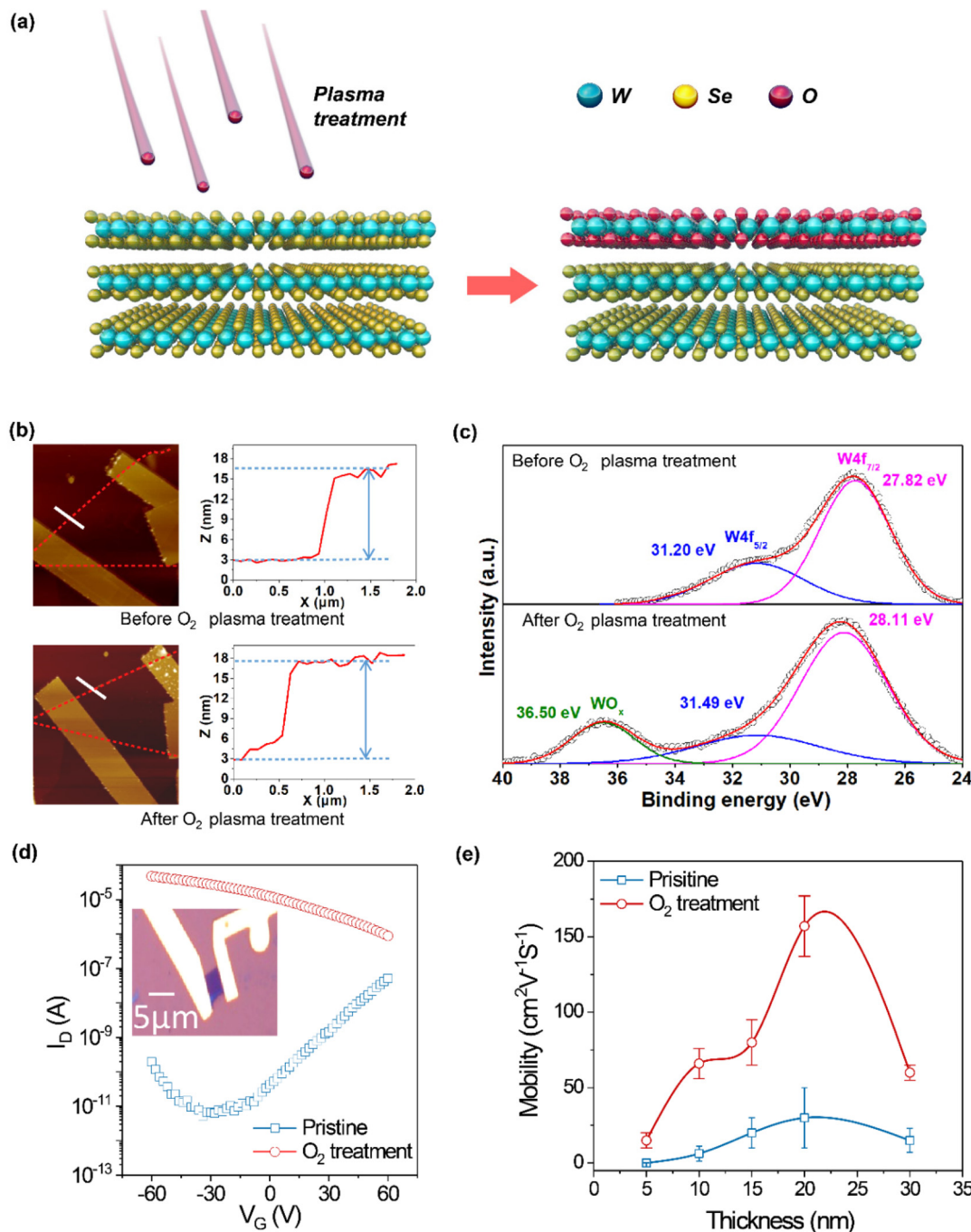


Fig. 1 (a) Schematic illustration of O₂ plasma treatment of bulk WSe₂. (b) The AFM images of a WSe₂ flake before and after O₂ plasma treatment, and the right part shows the AFM height profile revealing the thickness of the ~13 nm WSe₂ flake. (c) The XPS spectral comparison shows the existence of one extra peak at a binding energy of 36.50 eV corresponding to WO_x and a slight shift in the binding energies of W peaks after O₂ plasma treatment. (d) Transfer characteristics of a WSe₂ FET with V_d = 1 V before and after O₂ plasma treatment. The inset shows the optical image of the as-fabricated WSe₂ FET. (e) The mobility obtained from WSe₂ FETs of different thicknesses.

two main peaks at 27.82 and 31.20 eV corresponding to W 4f_{7/2} and W 4f_{5/2}, respectively.^{38,39} In comparison, the W 4f_{7/2} (28.11 eV) and W 4f_{5/2} (31.49 eV) peaks of the plasma-treated WSe₂ showed a slight shift to higher binding energies. Moreover, a clear additional peak can be observed at the binding energy of 36.50 eV corresponding to WO_x (x ≤ 3),⁴⁰ which verifies the presence of WO_x at the top of WSe₂ flakes. These results

demonstrate that oxygen radicals are efficiently doped into WSe₂ crystals during plasma exposure. Furthermore, the transport characteristics of the pristine and plasma-treated WSe₂-based FETs were investigated under an applied gate voltage from -60 to 60 V. The optical microphotograph of the plasma treated WSe₂ based-FET is presented in the inset of Fig. 1(d). The pristine WSe₂-based FET exhibits n-type dominated ambi-



polar characteristics, as shown in Fig. 1(d). However, after O₂ plasma treatment, the transfer curve shows a heavily p-doped characteristic (Fig. 1(d)). These results confirm the modulation of the as-fabricated WSe₂ channels with O₂ plasma, which act as p-type dopants when exposed to controlled plasma conditions.⁴¹ The field-effect mobility (μ_{FE}) can be calculated from the following formula $\mu_{FE} = g_m L_c h / C_{ox} W_{ch} V_{SD}$,⁴² where L_{ch} and W_{ch} are the channel length and width, respectively, $C_{ox} = 1.15 \times 10^{-8}$ F cm⁻² for 300 nm SiO₂, $g_m = \partial I_D / \partial V_{BG}$ is the transconductance, and V_{SD} is the source–drain voltage. The maximum μ_{FE} of the as-fabricated WSe₂ FETs is found to be about 157 cm² V⁻¹ s⁻¹, which is much higher than previously reported results (see Table 1). Fig. 1(e) shows the μ_{FE} of the pristine and O₂-treated WSe₂ devices fabricated with different thicknesses from 5 nm to 30 nm, where the FET mobility is significantly increased by up to a factor of 8 after O₂ plasma-induced p-type doping. We believe that the high mobility obtained may be attributed to the lowering of the contact resistance (R_c) by plasma doping.

It is well-known that a high contact resistance R_c would significantly affect the electrical and optical performances of TMDC-based FETs.³⁰ Hence, we fabricated devices with the same channel width and different channel lengths ranging from 2 to 5 μ m to evaluate the contact R_c of the p-doped WSe₂ channels by the transmission line method (TLM). Fig. 2(a) shows the optical images of the WSe₂ devices with varied channel lengths for examining contact properties. The enlarged schematic in Fig. 2(a) illustrates the carrier pathways (yellow arrows) in the channel of the plasma-treated device, where the electron transfer from the underlying WSe₂ to the surface WO_x contributes to the formation of electron transport paths (red arrows) from the channel to the doping layer, which is consistent with previous reports.^{23,43,44} In particular, a thin WO_x layer formed after the oxygen plasma treatment will attract electrons from the neighboring (or underlying) WSe₂, thereby inducing the heavily p-doped WSe₂, *i.e.*, the formation of a conducting channel of holes. The transfer characteristics of these devices are shown in Fig. S3.† Fig. 2(b) and (c) show the output characteristics obtained from the pristine and O₂ plasma-treated WSe₂-based FETs, respectively. The non-linear

increase of hole current with increasing drain bias (V_D) indicates a large SB between the metal and pristine WSe₂ layers (see Fig. 2(b)). In contrast, the excellent linear characteristics of the output curve of the O₂-plasma treated WSe₂ layers in Fig. 2(c) suggest the conversion of Schottky contact into Ohmic-like behavior. This may be attributed to the shift of the Fermi level towards the valence band maximum (VBM) of WSe₂: p-type dominant conversion. We believe that this phenomenon contributes to the transition from Schottky emission to direct tunneling for easier hole injection. For Schottky emission,⁴⁵ $I = I_{SE} \exp(A\sqrt{V})$, where I_{SE} and A (>0) are two parameters related to Schottky emission behavior. For F–N tunneling or FNT,⁴⁶ $I = I_{FN} V^2 \exp(B/V)$, where I_{FN} and B (<0) are two parameters related to F–N tunneling behavior. To better elucidate the transition from Schottky emission to direct tunneling, the I - V curves are plotted in the forms of $\ln I$ -vs.- $V^{1/2}$ and $\ln(I/V^2)$ -vs.- V^{-1} , respectively. Fig. S4† shows that before plasma treatment, the channel is dominated by Schottky emission, while after plasma treatment, the F–N tunneling becomes dominant. To further quantify the variation in R_c after plasma treatment, Fig. 2(d) shows the linear fitting of the total resistance (R_{total}) as a function of channel length (L). R_{total} could be described as $R_{total} = LR_{ch}/W + 2R_c$,⁴⁷ where R_c is the contact resistance, W is the channel width, L is the channel length, and R_{ch} is the 2D channel resistance. The value of R_c can be directly extracted from the y-intercept of the linear fit to R_{total} . Fig. 2(e) shows R_c ranging from 0.71 to 6.73 k Ω μ m under different gate voltages ranging from -60 to 60 V. Considering the small thickness of the channel, uniform doping leads to a uniform current distribution in the vertical direction. A high gate bias leads to an increase in carrier density and the narrowing of the potential barrier along the metal and the material interface. The elevated hole concentration and narrowed potential barrier aid in the injection of carriers into the WSe₂ channel and thereby achieve a transition from Schottky emission to direct tunneling. As a result, a lower contact resistance is obtained.⁴⁷ In particular, the lowest R_c of $\sim 710 \Omega \mu$ m measured at $V_{BG} = -60$ V can be successfully achieved in the plasma-treated WSe₂ FETs, suggesting a reduction in R_c after the plasma treatment. It should be noted that such a small R_c highlights the importance of the self-terminating oxidation process, resulting in the thinnest oxide and therefore a transparent tunnel barrier for increased hole injection between the metal and WSe₂ as a key factor for the excellent device performance.

The p–n junction is crucial for the separation of photon-induced electron–hole (e–h) pairs towards achieving photo-detectors with high responsivity or solar cells with high quantum efficiency.^{48,49} Through oxidation doping, the partial doping of a WSe₂ FET results in a lateral p–n homojunction. To determine their optoelectronic performance, WSe₂ lateral junctions were fabricated. More detailed information regarding the device preparation process can be found in the Experimental section. Fig. 3(a) shows the schematic diagram and the optical image (top-view) of the fabricated WSe₂ devices. These WSe₂ devices were then exposed to O₂ plasma

Table 1 Comparison of the field-effect mobility with other WSe₂-based electronic devices

Materials	Field-effect mobility (cm ² V ⁻¹ s ⁻¹)	Ref.
WSe₂ pn-junction	157	This work
WSe ₂ /MoS ₂ heterojunction	8.42	6
WSe ₂ (1-x)Te _{2x}	46	7
WSe ₂ /ReS ₂ pn-junction	16.5	10
WSe ₂ pn-junction	36.28	11
Monolayer WSe ₂	132	13
WSe ₂ /h-BN heterojunction	25.6	15
WO _x /WSe ₂ heterojunction	40	23
WSe ₂ flake	2.2	24
WSe ₂ /Graphene heterojunction	84	25
Monolayer WSe ₂	82	28
WSe ₂ flake	150	35





Fig. 2 Contact properties of the WSe₂ FETs obtained by TLM measurement. (a) Optical micrograph of the TLM devices. The enlarged section on the right side shows the hole transport path occurring in the whole channel, which can lead to low contact resistance and thus higher mobility. (b) Schottky contact can be clarified from the output characteristics of the pristine WSe₂ FETs. (c) By applying O₂ plasma treatment, Ohmic contact was achieved in WSe₂ FETs. (d) R_{total} as a function of channel length to extract R_c using a TLM linear fitting curve. (e) R_c versus bottom gate voltage in which the lowest R_c obtained was $\sim 710 \Omega \cdot \mu\text{m}$ at a V_{BG} of -60 V .

under different operating conditions (see the Experimental section). Note that a PMMA layer is employed as a mask to protect areas with electrodes 1 to 2, leaving areas with electrodes marked with 3 to 4 exposed for O₂ plasma treatment. The O₂ molecule was then chemisorbed onto the top surface of the exposed WSe₂, in which the top layer of the Se atom was substituted by an oxygen atom. Finally, the WO_x layer was formed on top of the WSe₂ layer. As a result of p-doping due to O₂ plasma (formation of a WO_x layer on the top of WSe₂), a p-type

(n-type) dominant conduction is observed in the O₂ plasma-treated (plasma-protected) region. The typical transfer curves of the pristine WSe₂ channel (through electrodes 1 to 2) were measured to maintain n-type behaviour (Fig. 3(b)). The electrical performance of the as-prepared plasma-treated channel was then measured through electrodes 3 to 4 (Fig. 3(c) and Fig. S5†), yielding p-type characteristics. With the capability of achieving both unipolar p-type (plasma-treated) and n-type (pristine) conduction behaviours, a lateral p–n homojunction





Fig. 3 (a) 3D schematic structure and optical photograph (right side) of an O₂ plasma-treated WSe₂ FET device showing channels 1–2, 2–3, and 3–4 as n-WSe₂, pn-WSe₂, and p-WSe₂, respectively. (b) and (c) The transfer curves of n- and p-WSe₂ channels measured at V_d = 1 V after O₂ plasma treatment, which are represented in both linear and logarithmic scales. (d) The output curve of an O₂ plasma-treated WSe₂ p–n junction device obtained at V_g = –40 V is represented on a linear scale. (e) The output curve is represented on a logarithmic scale to calculate the p–n junction ideality factor as 1.4.

can then be fabricated using a single WSe₂ flake. Fig. 3(d) and (e) show the corresponding I_d – V_d output curves of the partially doped WSe₂ device (channel: electrodes 2–3) with V_g = –40 V in the linear (left) and log (right) scales, in which little to no current flow is observed under reverse source–drain bias (V_d < 0 V). However, a much higher current flow can be observed when a forward bias is applied (V_d > 0 V), indicating the forward rectifying diode behavior. Furthermore, an ideality factor (*n*)^{50,51} of 1.4 was calculated from the dark current (Fig. 3(e)), indicating the formation of a near-ideal p–n junction.^{52,53} It is known that *n* = 1 indicates an ideal p–n junction

diode, *n* ranging between 1 and 2 implies the existence of a tunneling current, and *n* = 2 indicates the existence of defects or interface traps that drives the recombination process.⁵⁴ Thus, our good ideality factor of ~1.4 can be ascribed to the low charge trap density at the homojunction interface.⁵⁵ The above results prove the successful preparation of a p–n junction between electrodes 2 and 3. However, it should be noted that the p–n WSe₂ device is formed with a Schottky junction in the n-type part of WSe₂ (contact 2), p–n junctions in the channel (electrodes 2–3), and Ohmic-like contact in the p-type part of WSe₂ (contact 3).



Next, we explored the optoelectronic properties of p–n WSe₂ homojunction photodetectors. The transfer curves of the as-fabricated photodetector were measured in the dark and 532 nm illumination. By comparison, a photocurrent response characteristic under light irradiation is observed in Fig. 4(a), which may be mainly attributed to the strong light absorption in the WSe₂ p–n junction. This then generates a larger density of photoinduced e–h pairs which leads to a reduction of the Schottky barrier, thereby allowing the photogenerated carriers to transmit more effectively along the metal/WSe₂ interface under an externally applied bias. To further investigate the performance of WSe₂ p–n junction photodetectors and to study their dynamic optical response, time (*T*)- and gate voltage (*V_G*)-dependent photocurrent measurements under a fixed source–drain bias of 1 V were performed as shown in Fig. 4(b). Here, the time (*T*)- and source–drain voltage (*V_D*)-dependent photocurrent of the WSe₂ p–n junction photodetectors was also measured at the fixed gate bias of 40 V (Fig. 4(c)). As *V_G* swept from +40 to –40 V, the photocurrent enhancement was observed under various *V_D* and *V_G* conditions. In particular, the maximum photocurrent is obtained under a gate bias voltage of 40 V. Meanwhile, prominent variations in the gate-dependent photocurrent can be obtained at both *V_G* < 0 V and *V_G* > 0 V sides. These results suggest that the photocurrent response can be effectively modulated by the gate electric field. In addition, the photocurrent is observed to slightly decrease with time under continuous light irradiation, indicating possible charge trapping effects. Higher gate voltage (40 V) leads to an increase in carrier concentration while the potential barrier at the contact interface becomes narrower. Thus, the charge trap states are relatively more easily filled with carriers, leading to a slight reduction in photocurrent. In contrast, at a gate voltage of 0 V, when the Fermi level is in equilibrium, the carrier density is relatively low. As such, carriers are difficult to be captured by trap states, so minimal photocurrent changes are observed in time-dependent photocurrent measurements (Fig. S6†). The photocurrent characteristics of the junction can be further understood through the energy band diagrams, as shown in Fig. 4(d)–(f). Fig. 4(d) shows the band diagram of the WSe₂ p–n junction in the equilibrium state. The Fermi level of the p-type area is near to the valence band, whereas the Fermi level of the n-type part is near to the conduction band. Under reverse bias, the energy barrier at the junction becomes higher as the width of the depletion region increases. Consequently, electrons cannot transmit effectively, and the p–n junction is in the OFF state. Under forward bias, the energy barrier becomes lower, thus allowing electrons to easily pass through the junction, and the p–n junction is in the ON state. This typical p–n junction behavior is consistent with the output curve as shown in Fig. 3(d). Furthermore, the mechanism of photocurrent generated in the dark (Fig. 4(e)) and light (Fig. 4(f)) states was also proposed. When the negative gate voltage is applied in the dark state (Fig. 4(e)), the Fermi level drops below the valence band of the p-type branch, which enables holes to flow easily from the junction into the metal. As a result,

there is no current in the p–n junction. When a positive gate voltage is applied, the Fermi level moves into the conduction band in the n-type branch, which allows electrons to easily transfer from the junction into the metal. Therefore, the designed lateral p–n junction shows no current when electrons are dominant carriers. Under light irradiation, the generation of e–h pairs takes place near the depletion region, and the photoexcited carriers are swept from the junction in opposite directions by an externally applied bias, resulting in a net increase in photocurrent, as shown in Fig. 4(f). As a result, internal hole and electron transport combined with the resulting overdose e–h pairs from light absorption increases the current, which is much higher than the dark state (as shown in Fig. 4(e)) even at low gate voltages. In both positive and negative gate biases, the photocurrent is high as the substantial photogenerated e–h pairs can be readily separated and driven toward the cathode and anode contacts in the *in situ* WSe₂ lateral p–n junction.

Next, we further investigated the photoresponse of the designed WSe₂ p–n junction by measuring the *I_D*–*V_D* characteristics at the gate bias of –40 V and 40 V, respectively. Under different powers of light illumination (6, 16, 31, and 42 mW), the WSe₂ p–n junction exhibited ideal diode characteristics as shown in Fig. 4(g) and (h). The photoresponse figures of merit such as photoresponsivity (*R*), external quantum efficiency (EQE), and detectivity (*D*^{*}) are sequentially evaluated. The spectral responses of WSe₂ photodetectors were expressed through *R*, which is defined as the ratio of the photocurrent and the incident laser power, *i.e.*, $R = I_{\text{ph}}/P_{\text{Laser}}$, where *I_{ph}* is the photocurrent = $|I_{\text{light}} - I_{\text{dark}}|$, and *P_{Laser}* is the incident laser power per unit area. Here, $I_{\text{ph}} = |I_{\text{light}} - I_{\text{dark}}| = 6 \times 10^{-6}$ A, $P = L \cdot W / (\pi \cdot r^2) \cdot P_m = 3.5 \times 2.5 / (3.14 \times 30^2) \times 2.7 \times 10^{-5} = 8.4 \times 10^{-8}$ W, in which *L* and *W* are the channel length and channel width, *r* is the radius, and *P_m* is the power measured by a power meter. Fig. 4(i) shows the photoresponsivities of the WSe₂ photodetector as a function of incident laser power. Another crucial figure of merit for the photodetector is the EQE, which is given by $\text{EQE} (\%) = 100 \cdot hcR / e\lambda$, where *h* is Planck's constant (6.63×10^{-34} J s), *c* is the speed of light (3×10^8 m s^{–1}), *e* is the unit charge, and *λ* is the wavelength of incident light (*λ* = 532 nm). The EQE of a photodetector represents the ratio of the total number of charge carriers generated to the number of incident photons. Through the above equations, the *R* and EQE of our p–n junction are calculated to be $\sim 7.1 \times 10^4$ mA W^{–1} and $\sim 228\%$, respectively, outperforming other reported results as summarized in Table 2.^{10,12,14–17,19,56–59} The high values of *R* and EQE also demonstrate that the plasma-treated WSe₂ p–n junction exhibits a good photoresponsivity. In addition, *D*^{*} is another figure of merit for photodetector devices to evaluate the detector sensitivity, which is generally calculated using the following equation of $D^* (\text{Jones}) = (R \cdot A^{1/2}) / (I_{\text{dark}} \cdot 2e)^{1/2}$. The *D*^{*} value obtained for our designed WSe₂ p–n junction photodetector was about 3×10^3 Jones. These results prove the effectiveness of O₂ plasma treatment for achieving high-performance WSe₂-based optoelectronic devices.





Fig. 4 (a) Typical photoresponse of the WSe₂ p-n homojunction FET devices under 532 nm illumination at different gate biases demonstrating clear photovoltaic behavior. (b) and (c) Time-dependent photocurrent measurements of pn-WSe₂ achieved under 532 nm laser irradiation at the applied $V_d = 1$ V with various V_g values and fixed $V_g = 40$ V with various V_d values. (d–f) The corresponding energy band diagrams show the carrier transport paths in the WSe₂ p-n homojunction device with and without light irradiation at the applied gate bias. (g) and (h) The photocurrents of the as-fabricated WSe₂ p-n homojunction under different powers of light illumination (5.35, 16, 31 and 42 mW) at gate biases of -40 V and 40 V, respectively. (i) The photoresponsivity (R) of the WSe₂ p-n junction obtained under gate biases of -40 V and 40 V as a function of light power. The responsivity was linear with the incident optical power.



Table 2 Comparison of R with other WSe_2 -based devices

Materials	Light irradiation conditions (nm)	R (mA W^{-1})	Response time (ms)	Ref.
WSe_2 p-n junction	532	7.1×10^4	42	This work
$\text{WSe}_2/\text{ReS}_2$ p-n junction	532	3×10^3	0.004	10
WSe_2 p-n junction	520	80	0.106	12
$\text{WSe}_2/\text{MoSe}_2$ heterojunction	1550	127		14
$\text{WSe}_2/\text{h-BN}$ junction	520	1.27×10^9	38.2–97	15
WSe_2/ZnO junction	405	4.83×10^6	0.01	16
$\text{WSe}_2/\text{Bi}_2\text{Te}_3$ p-n junction	633	2.05×10^4	0.21	17
p-Doped WSe_2	735	600	0.008	19
Multi-layer WSe_2	532	100	0.01	56
$\text{WSe}_2/\text{BP}/\text{MoS}_2$ heterojunction	532	6.32×10^3		57
Multi-layer $\text{WSe}_2/\text{pentacene}$	638	1.93×10^4		58
$\text{WSe}_2/\text{MoS}_2$ heterojunction	532	3.4		59

3. Conclusion

In summary, lateral p–n homojunctions were successfully fabricated from bulk WSe_2 via oxygen plasma doping and demonstrated in gate-modulated photodetector applications. After the plasma treatment, the lowest R_c of $\sim 710 \Omega \mu\text{m}$ of the WSe_2 p–n junction can be achieved, which is attributed to the highly p-doped WSe_2 that converts the Schottky barrier to Ohmic contact between the metal and WSe_2 interface. Thanks to the low R_c and excellent mobility of $\sim 157 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, our plasma-treated WSe_2 p–n junction exhibited an exceptional photoresponsivity of $\sim 7.1 \times 10^4 \text{ mA W}^{-1}$, a superior external quantum efficiency of $\sim 228\%$, and a detectivity of $\sim 3 \times 10^3$ Jones under 532 nm light illumination wavelength. Our work provides an industry-compatible plasma doping technique to fabricate homogeneous in-plane p–n junctions for advanced 2-D optoelectronic device applications. Going forward, homogeneous complementary 2-D inverters can be fabricated by integrating the plasma-treated and as-exfoliated regions of a single WSe_2 layer to create basic building blocks for logic circuits. Our results contribute toward the realization of the local plasma treatment of WSe_2 flakes to tune the electrical properties of WSe_2 and the feasibility of fabricating WSe_2 -based optoelectronic logic applications.

4. Experimental section

Fabrication of pristine WSe_2 -based FETs and p-type WSe_2 -based FETs

The synthesis process of the pristine WSe_2 -based device is given as follows: at first, the WSe_2 film was exfoliated from the crystalline bulk material (HQ Graphene, the Netherlands) onto a 300 nm-thick SiO_2/Si substrate. Next, electrodes were patterned using electron beam lithography (EBL), and then deposited with Cr/Au (5 nm/70 nm) using an electron beam/thermal evaporator. To obtain p-type WSe_2 -based FETs, the as-fabricated pristine WSe_2 -based FETs were treated with O_2 plasma (20 W, 120 s). Our used XPS model is Thermo Fisher (Microlab350) with magnesium light.

Fabrication of the WSe_2 p–n junction

The fabrication process of the WSe_2 p–n junction is as follows: first, the pristine WSe_2 device was spin-coated with PMMA (Microchem, A4 495) at a speed of 3000 rpm (60 s). The device was then annealed at 150 °C for 2 min on a hot plate. Afterward, another layer PMMA (Microchem, A6 950) was spin-coated on the top of PMMA A4 with the same process. Furthermore, to obtain the p–n junction structure, we used EBL along with a developer (mixed DI water and IPA solution at a ratio of 1 : 3) to open a channel near the electrode, followed by rinsing with IPA and drying under the N_2 gas blowing. Finally, O_2 plasma treatment is employed to treat the as-fabricated device to obtain the WSe_2 p–n junction structure. The electrical transfer performances of the devices were tested in a vacuum probe station (Keithley 4200) at 0.1 Pa. The photocurrent response performance was measured using a light source of 532 nm.

Photocurrent measurement

Solid-state laser diodes were used as illumination sources. The incident power of the lasers (P_m) was measured with an optical power meter (Newport model 1918-C) and a thermopile sensor (919P-003-10). We measured the photoresponse characteristics with a monochromator light source. The device was placed on the sample holder, and a continuous wave laser beam generated from the solid-state laser diode and controller (Thorlabs LPSC-520-FC and CLD1010LP for 520 nm (visible) and Thorlabs LPSC-852-FC and CLD1010LP for 852 nm (near IR)) was focused onto the device through a 40× objective lens (Olympus LUCPlanFLN, 40×, NA = 0.6).

Conflicts of interest

The authors declare no conflict of interest.

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References

- H. Xu, X. Han, X. Dai, W. Liu, J. Wu, J. Zhu, D. Kim, G. Zou, K. A. Sablon, A. Sergeev, Z. Guo and H. Liu, *Adv. Mater.*, 2018, **30**, 1706561.
- P. Gant, P. Huang, D. Pérez de Lara, D. Guo, R. Frisenda and A. Castellanos-Gomez, *Mater. Today*, 2019, **27**, 8–13.
- N. Flory, P. Ma, Y. Salamin, A. Emboras, T. Taniguchi, K. Watanabe, J. Leuthold and L. Novotny, *Nat. Nanotechnol.*, 2020, **15**, 118–124.
- H. K. Ng, D. Xiang, A. Suwardi, G. Hu, K. Yang, Y. Zhao, T. Liu, Z. Cao, H. Liu, S. Li, J. Cao, Q. Zhu, Z. Dong, C. K. I. Tan, D. Chi, C.-W. Qiu, K. Hippalgaonkar, G. Eda, M. Yang and J. Wu, *Nat. Electron.*, 2022, **5**, 489–496.
- T. Liu, D. Xiang, H. K. Ng, Z. Han, K. Hippalgaonkar, A. Suwardi, J. Martin, S. Garaj and J. Wu, *Adv. Sci.*, 2022, **9**(20), 2200816.
- G. H. Shin, C. Park, K. J. Lee, H. J. Jin and S. Y. Choi, *Nano Lett.*, 2020, **20**(8), 5741–5748.
- P. Yu, J. Lin, L. Sun, Q. L. Le, X. Yu, G. Gao, C. H. Hsu, D. Wu, T. R. Chang, Q. Zeng, F. Liu, Q. J. Wang, H. T. Jeng, H. Lin, A. Trampert, Z. Shen, K. Suenaga and Z. Liu, *Adv. Mater.*, 2017, **29**(4), 1603991.
- S. Ahmed, X. Ding, P. P. Murmu, N. Bao, R. Liu, J. Kennedy, L. Wang, J. Ding, T. Wu, A. Vinu and J. Yi, *Small*, 2020, **16**, 1903173.
- S. Jia, Z. Jin, J. Zhang, J. Yuan, W. Chen, W. Feng, P. Hu, P. M. Ajayan and J. Lou, *Small*, 2020, **16**, 2002263.
- A. Varghese, D. Saha, K. Thakar, V. Jindal, S. Ghosh, N. V. Medhekar, S. Ghosh and S. Lodha, *Nano Lett.*, 2020, **20**(3), 1707–1717.
- J. Sun, Y. Wang, S. Guo, B. Wan, L. Dong, Y. Gu, C. Song, C. Pan, Q. Zhang, L. Gu, F. Pan and J. Zhang, *Adv. Mater.*, 2020, **32**, 1906499.
- Y. Tang, Z. Wang, P. Wang, F. Wu, Y. Wang, Y. Chen, H. Wang, M. Peng, C. Shan, Z. Zhu, S. Qin and W. Hu, *Small*, 2019, **15**, 1805545.
- A. S. Bandyopadhyay, N. Adhikari and A. B. Kaul, *Chem. Mater.*, 2019, **31**(23), 9861–9874.
- H. Xue, Y. Wang, Y. Dai, W. Kim, H. Jussila, M. Qi, J. Susoma, Z. Ren, Q. Dai, J. Zhao, K. Halonen, H. Lipsanen, X. Wang, X. Gan and Z. Sun, *Adv. Funct. Mater.*, 2018, **28**, 1804388.
- S. H. Jo, D. H. Kang, J. Shim, J. Jeon, M. H. Jeon, G. Yoo, J. Kim, J. Lee, G. Y. Yeom, S. Lee, H. Y. Yu, C. Choi and J. H. Park, *Adv. Mater.*, 2016, **28**(24), 4824–4831.
- N. Guo, L. Xiao, F. Gong, M. Luo, F. Wang, Y. Jia, H. Chang, J. Liu, Q. Li, Y. Wu, Y. Wang, C. Shan, Y. Xu, P. Zhou and W. Hu, *Adv. Sci.*, 2020, **7**, 1901637.
- H. Liu, X. Zhu, X. Sun, C. Zhu, W. Huang, X. Zhang, B. Zheng, Z. Zou, Z. Luo, X. Wang, D. Li and A. Pan, *ACS Nano*, 2019, **13**(11), 13573–13580.
- R. Cheng, D. Li, H. Zhou, C. Wang, A. Yin, S. Jiang, Y. Liu, Y. Chen, Y. Huang and X. Duan, *Nano Lett.*, 2014, **14**(10), 5590–5597.
- T. Wang, K. Andrews, A. Bowman, T. Hong, M. Koehler, J. Yan, D. Mandrus, Z. Zhou and Y. Q. Xu, *Nano Lett.*, 2018, **18**(5), 2766–2771.
- D. Li, M. Chen, Z. Sun, P. Yu, Z. Liu, P. M. Ajayan and Z. Zhang, *Nat. Nanotechnol.*, 2017, **12**, 901–906.
- Y. Ren, X. Yang, L. Zhou, J. Y. Mao, S. T. Han and Y. Zhou, *Adv. Funct. Mater.*, 2019, **29**, 1902105.
- M. V. Gustafsson, M. Yankowitz, C. Forsythe, D. Rhodes, K. Watanabe, T. Taniguchi, J. Hone, X. Zhu and C. R. Dean, *Nat. Mater.*, 2018, **17**, 411–415.
- X. Liu, D. Qu, Y. Yuan, J. Sun and W. J. Yoo, *ACS Appl. Mater. Interfaces*, 2020, **12**(23), 26586–26592.
- S. J. Yang, K. T. Park, J. Im, S. Hong, Y. Lee, B. W. Min, K. Kim and S. Im, *Nat. Commun.*, 2020, **11**, 1574.
- H. L. Tang, M. H. Chiu, C. C. Tseng, S. H. Yang, K. J. Hou, S. Y. Wei, J. K. Huang, Y. F. Lin, C. H. Lien and L. J. Li, *ACS Nano*, 2017, **11**(12), 12817–12823.
- S. Fan, S. J. Yun, W. J. Yu and Y. H. Lee, *Adv. Sci.*, 2020, **7**, 1902751.
- S. W. LaGasse, P. Dhakras, K. Watanabe, T. Taniguchi and J. U. Lee, *Adv. Mater.*, 2019, **31**, 1901392.
- G. Dastgeer, M. F. Khan, G. Nazir, A. M. Afzal, S. Aftab, B. A. Naqvi, J. Cha, K. A. Min, Y. Jamil, J. Jung, S. Hong and J. Eom, *ACS Appl. Mater. Interfaces*, 2018, **10**(15), 13150–13157.
- H. G. Ji, P. Solis-Fernandez, D. Yoshimura, M. Maruyama, T. Endo, Y. Miyata, S. Okada and H. Ago, *Adv. Mater.*, 2019, **31**, 1903613.
- M. Si, C. Jiang, W. Chung, Y. Du, M. A. Alam and P. D. Ye, *Nano Lett.*, 2018, **18**(6), 3682–3687.
- G. Dastgeer, M. F. Khan, J. Cha, A. M. Afzal, K. H. Min, B. M. Ko, H. Liu, S. Hong and J. Eom, *ACS Appl. Mater. Interfaces*, 2019, **11**(11), 10959–10966.
- R. Zhou, V. Ostwal and J. Appenzeller, *Nano Lett.*, 2017, **17**(8), 4787–4792.
- H. Liu, D. Fu, X. Li, J. Han, X. Chen, X. Wu, B. Sun, W. Tang, C. Ke, Y. Wu, Z. Wu and J. Kang, *ACS Nano*, 2021, **15**(5), 8244–8251.
- A. M. Dibos, Y. Zhou, L. A. Jauregui, G. Scuri, D. S. Wild, A. A. High, T. Taniguchi, K. Watanabe, M. D. Lukin, P. Kim and H. Park, *Nano Lett.*, 2019, **19**(6), 3543–3547.
- E. V. Calman, L. H. Fowler-Gerace, D. J. Choksy, L. V. Butov, D. E. Nikonov, I. A. Young, S. Hu, A. Mishchenko and A. K. Geim, *Nano Lett.*, 2020, **20**(3), 1869–1875.
- Y. Huang, K. Yu, H. Li, K. Xu, Z. Liang, D. Walker, P. Ferreira, P. Fischer and D. E. Fan, *Adv. Mater.*, 2020, **32**, 2003439.



- 37 Z. Zheng, J. Chen, Y. Wang, X. Wang, X. Chen, P. Liu, J. Xu, W. Xie, H. Chen, S. Deng and N. Xu, *Adv. Mater.*, 2018, **30**, 1705318.
- 38 R. Zhang, D. Drysdale, V. Koutsos and R. Cheung, *Adv. Funct. Mater.*, 2017, **27**, 1702455.
- 39 X. Yu, N. Guijarro, M. Johnson and K. Sivula, *Nano Lett.*, 2018, **18**(1), 215–222.
- 40 H. Medina, J. G. Li, T. Y. Su, Y. W. Lan, S. H. Lee, C. W. Chen, Y. Z. Chen, A. Manikandan, S. H. Tsai, A. Navabi, X. Zhu, Y. C. Shih, W. S. Lin, J. H. Yang, S. R. Thomas, B. W. Wu, C. H. Shen, J. M. Shieh, H. N. Lin, A. Javey, K. L. Wang and Y. L. Chueh, *Chem. Mater.*, 2017, **29**, 1587.
- 41 B. Liu, Y. Ma, A. Zhang, L. Chen, A. N. Abbas, Y. Liu, C. Shen, H. Wan and C. Zhou, *ACS Nano*, 2016, **10**(5), 5153–5160.
- 42 C. Liu, G. Li, R. Di Pietro, J. Huang, Y.-Y. Noh, X. Liu and T. Minari, *Phys. Rev. Appl.*, 2017, **8**, 034020.
- 43 Z. Li, S. Yang, R. Dhall, E. Kosmowska, H. Shi, I. Chatzakis and S. B. Cronin, *ACS Nano*, 2016, **10**(7), 6836–6842.
- 44 T. D. Ngo, M. S. Choi, M. Lee, F. Ali, Y. Hassan, N. Ali, S. Liu, C. Lee, J. Hone and W. J. Yoo, *Adv. Sci.*, 2022, **9**, 2202465.
- 45 Y. Tian, C. Guo, S. Guo, T. Yu and Q. Liu, *Nano Res.*, 2014, **7**, 953–962.
- 46 Y. Tian, L. Jiang, X. Zhang, G. Zhang and Q. Zhu, *AIP Adv.*, 2018, **8**, 035105.
- 47 F. Xia, V. Perebeinos, Y. M. Lin, Y. Wu and P. Avouris, *Nat. Nanotechnol.*, 2011, **6**, 179–184.
- 48 S. Ross, P. Klement, A. M. Jones, N. J. Ghimire, J. Yan, D. G. Mandrus, T. Taniguchi, K. Watanabe, K. Kitamura, W. Yao, D. H. Cobden and X. Xu, *Nat. Nanotechnol.*, 2014, **9**, 268–272.
- 49 J. Guo, D. Xie, B. Yang and J. Jiang, *Solid-State Electron.*, 2018, **144**, 1.
- 50 M. Y. Li, C. H. Chen, Y. Shi and L. J. Li, *Mater. Today*, 2016, **19**, 322–335.
- 51 Z. Yang, L. Liao, F. Gong, F. Wang, Z. Wang, X. Liu, X. Xiao, W. Hu, J. He and X. Duan, *Nano Energy*, 2018, **49**, 103.
- 52 T. Shen, J. C. Ren, X. Liu, S. Li and W. Liu, *J. Am. Chem. Soc.*, 2019, **141**, 3110.
- 53 A. Li, Q. Chen, P. Wang, Y. Gan, T. Qi, P. Wang, F. Tang, J. Z. Wu, R. Chen, L. Zhang and Y. Gong, *Adv. Mater.*, 2019, **31**, 1805656.
- 54 N. N. Mude, R. N. Bukke and J. Jang, *Adv. Mater. Technol.*, 2022, 2101434.
- 55 Z. Cheng, *et al.*, *J. Mater. Chem. C*, 2022, **10**, 1511–1516.
- 56 N. R. Pradhan, J. Ludwig, Z. Lu, D. Rhodes, M. M. Bishop, K. Thirunavukkuarasu, S. A. McGill, D. Smirnov and L. Balicas, *ACS Appl. Mater. Interfaces*, 2015, **7**(22), 12080–12088.
- 57 H. Li, L. Ye and J. Xu, *ACS Photonics*, 2017, **4**, 823.
- 58 H. Im and S. Kim, *Adv. Electron. Mater.*, 2021, **7**, 2100003.
- 59 P. Lin, L. Zhu, D. Li, L. Xu, C. Pan and Z. Wang, *Adv. Funct. Mater.*, 2018, **28**, 1802849.

