

Cite this: *Chem. Sci.*, 2019, 10, 5353

All publication charges for this article have been paid for by the Royal Society of Chemistry

Origin of the overall water splitting activity of Ta₃N₅ revealed by ultrafast transient absorption spectroscopy†

Dharmapura H. K. Murthy,^a Hiroyuki Matsuzaki,^a Zheng Wang,^b Yohichi Suzuki,^a Takashi Hisatomi,^b Kazuhiko Seki,^a Yasunobu Inoue,^{cd} Kazunari Domen^{b*bc} and Akihiro Furube^{bae}

Tantalum nitride (Ta₃N₅) is one of the few visible light absorbing photocatalysts capable of overall water splitting (OWS), by which the evolution of both H₂ and O₂ is possible. Despite favourable energetics, realizing the OWS or efficient H₂ evolution in Ta₃N₅ prepared by the nitridation of tantalum oxide (Ta₂O₅) or Ta foil remains a challenge even after 15 years of intensive research. Recently our group demonstrated OWS in Ta₃N₅ when prepared by the short time nitridation of potassium tantalate (KTaO₃). To obtain a mechanistic insight on the role of Ta precursor and nitridation time in realizing OWS, ultrafast dynamics of electrons (3435 nm probe) and holes (545 nm probe) is investigated using transient absorption spectroscopy. Electrons decay majorly by trapping in Ta₃N₅ prepared by the nitridation of Ta₂O₅, which do not show OWS. However, OWS activity in Ta₃N₅ prepared by 0.25 hour nitridation of KTaO₃ is particularly favoured by the virtually absent electron and hole trapping. On further increasing the nitridation time of KTaO₃ from 0.25 to 10 hour, trapping of both electron and hole is enhanced which concurrently results in a reduction of the OWS activity. Insights from correlating the synthesis conditions—structural defects—carrier dynamics—photocatalytic activity is of importance in designing novel photocatalysts to enhance solar fuel production.

Received 15th January 2019

Accepted 25th April 2019

DOI: 10.1039/c9sc00217k

rsc.li/chemical-science

1. Introduction

Considering ever increasing energy demand, there is an urgent need to promote and enhance the production of renewable energy sources. This approach will pay dividends in protecting our environment from a serious crisis. In this direction, using solar energy to generate H₂ from water using a photocatalytic process constitutes one of the methods.^{1–9} Over the past decades, this approach has been intensively investigated to enhance the solar to hydrogen (STH) conversion efficiency. For the commercial use of the solar H₂, a stable photocatalyst

consistently demonstrating a higher STH efficiency is desired. To address this challenge, novel visible light absorbing photocatalysts are being synthesized and tested for photocatalytic water splitting activity.

Tantalum nitride (Ta₃N₅) is an n-type, visible light absorbing nitride-based photocatalyst. The conduction band (CB) and valence band (VB) of Ta₃N₅ are characterized by the Ta 5d and N 2p orbitals, respectively.^{10,11} Ta₃N₅ is commonly prepared by the nitridation reaction of tantalum (Ta) precursors such as Ta foil or tantalum oxide (Ta₂O₅). The VB and CB positions in Ta₃N₅ are thermodynamically suitable to drive water oxidation and reduction to generate O₂ and H₂, respectively. Hence, Ta₃N₅ is capable of overall water splitting (OWS) upon one-step photoexcitation. However, Ta₃N₅ is usually known as a photoanode material with efficient O₂ evolution,^{10,12–17} but not for either H₂ evolution or OWS ability.

Recently, OWS activity with an apparent quantum efficiency (AQE) of 70% is demonstrated in an ultraviolet light absorbing SrTiO₃ photocatalyst.¹⁸ As a further advancement, Ta₃N₅ forms one of the few promising visible light absorbing photocatalysts capable of OWS. Nevertheless, achieving OWS in Ta₃N₅ prepared by nitriding Ta foil or Ta₂O₅ still remains a challenge. Recently our group employed potassium tantalate (KTaO₃) as a Ta precursor for the nitridation reaction to prepare Ta₃N₅. In this approach, Ta₃N₅ prepared by 0.25 hour nitridation of

^aNational Institute of Advanced Industrial Science and Technology (AIST), Tsukuba Central 2, 1-1-1 Umezono, Tsukuba, Ibaraki 305-8568, Japan. E-mail: hiroyuki-matsuzaki@aist.go.jp

^bCentre for Energy & Environmental Science, Shinshu University, 4-17-1 Wakasato, Nagano-shi, Nagano 380-8553, Japan

^cDepartment of Chemical System Engineering, School of Engineering, The University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-8656, Japan. E-mail: domen@chemsys.t.u-tokyo.ac.jp

^dJapan Technological Research Association of Artificial Photosynthetic Chemical Process (ARPCHEM), 2-11-9 Iwamotocho, Chiyoda-ku, Tokyo 101-0032, Japan

^eDepartment of Optical Science, Tokushima University, 2-1 Minamijosanjima-cho, Tokushima 770-8506, Japan. E-mail: furube.akihiro@tokushima-u.ac.jp

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c9sc00217k



KTaO₃ demonstrated the highest OWS activity. This is the first report on realizing visible light OWS activity without the need of sacrificial reagents for Ta₃N₅.¹⁹ However, the efficiency of OWS activity was found to decrease with increase in the nitridation time from 0.25 to 10 hour. These observations emphasize the need to elucidate the role of the Ta precursor and the nitridation time in determining the efficiency of OWS.

As stated earlier, the H₂ evolution in Ta₃N₅ (prepared by the nitridation Ta foil or Ta₂O₅) is apparently less known and typically smaller when compared to O₂ evolution. In a few reports that demonstrate H₂ evolution in Ta₃N₅, the size of Ta₃N₅ particles was relatively small (between 5 to 8 nm (ref. 20) and in one case it is 2 nm (ref. 21)). In addition to the smaller size, Ta₃N₅ particles were typically present in a composite form with either SiO₂ (ref. 20) or graphene oxide.²¹ In these reports, the effect of size on the band bending/optical properties and the expected electronic interaction between the Ta₃N₅ with the counterpart in the composite is not considered. In another report showing H₂ evolution, the surface of the Ta₃N₅ was modified with magnesium layer.²² Thus, without the modification of either structural/surface properties of Ta₃N₅ or by inducing electronic interaction with another material, H₂ evolution was not realized. Based on these observations, it is agreeable to state that efficient generation of H₂ in Ta₃N₅ is fundamentally less pronounced compared to O₂.^{19,21,23} However, the origin of this observation is unclear.

A key approach to examine the role of Ta precursor and the nitridation time on the efficiency of OWS is investigating the decay dynamics of photogenerated charge carriers. Whilst there has been extensive research/progress in the development of Ta₃N₅ from the materials perspective, the dynamics of photo-generated charge carriers and its impact on the efficiency of the photocatalytic reaction have not been explored. The fundamental process in OWS in Ta₃N₅ is the generation of electron-hole pairs upon absorption of light. Then, these carriers must migrate towards the surface of the Ta₃N₅ to get involved in the reaction with water. However, due to structural/electronic defects in Ta₃N₅, efficient charge transport towards the surface is affected by the undesirable trapping and/or recombination. To unveil these photoinduced processes in detail, femtosecond time-resolved diffuse reflectance (fs-TDR) spectroscopy is employed, which is a pump-probe method. Here, excitation of Ta₃N₅ with an ultrafast laser pulse leads to the formation of holes and electrons in the VB and CB, respectively. The dynamics of these transient species can be followed by monitoring their transient absorption spectrum using an optical probe pulse from visible to infrared. This technique has been widely used to examine charge carrier dynamics in various types of photocatalysts and to correlate the obtained photo-physical parameters with that of photocatalytic efficiency.^{24–29} Note that TDR is essentially similar to the transient absorption spectroscopy (TAS) except that the TA signal in the former is monitored in diffuse reflection mode because of the opaque nature of the Ta₃N₅ powder samples. The ultrafast dynamics of electrons and holes and their decay pathway was found to be strongly influenced by the Ta precursor used to prepare Ta₃N₅ and the nitridation time. Results from this systematic

investigation help in understanding the structure–property relationship and provide rational insights on Ta₃N₅ material design to improve the efficiency of OWS. As mentioned earlier, this work provides mechanistic insight on explaining OWS through charge carrier dynamics, while a comprehensive study of the morphological characterization and photocatalytic performance is reported elsewhere.¹⁹

2. Results and discussion

Fig. 1 shows the ground-state optical absorption spectra for the Ta₃N₅ prepared by nitridation of Ta₂O₅ or KTaO₃ in both visible and NIR region. The absorption peak at 545 nm (Fig. 1A) is attributed to the fundamental transition (N 2p to Ta 5d) in Ta₃N₅. Though the magnitude of absorption at 545 nm increases with nitridation time, the energetic position was not influenced by the choice of Ta precursor or the nitridation time. As shown in Fig. 1B, an enhancement in the magnitude of NIR absorption is noticed with an increase in the nitridation time from 0.25 to 10 hour indicating the presence of defect states within the bandgap of Ta₃N₅. For Ta₃N₅ prepared by 0.25 hour nitridation of KTaO₃, noticing a virtually absent NIR absorption suggests a significant reduction in the concentration of defect states. This conclusion is in good agreement with our previous study using scanning transmission electron microscopy (STEM), which revealed the single-crystal nature of this sample.¹⁹ Previous reports have different opinions in relating the NIR absorption with the electronic nature of these defect states.^{17,30–33} Hence, we consider that the NIR absorption in Ta₃N₅ is due to the formation of defects related to charged nitrogen vacancies and/or reduced Ta vacancy.

In the following parts, Section 2.1 will correlate the charge carrier dynamics in Ta₃N₅ synthesized by the nitridation of Ta₂O₅ and KTaO₃ (0.25 hour). In Section 2.2, the effect of KTaO₃ nitridation time on the charge carrier dynamics in Ta₃N₅ will be presented. In the last Section 2.3, plausible approaches to enhance the efficiency of OWS in Ta₃N₅ will be outlined.

2.1 Origin of overall water splitting in Ta₃N₅ prepared by the 0.25 hour nitridation of KTaO₃

The efficiency of OWS was found to be highest in Ta₃N₅ prepared by 0.25 hour nitridation of KTaO₃. On the other hand, 0.25 hour nitridation of Ta₂O₅ neither forms Ta₃N₅ phase nor shows the OWS. To understand the origin of OWS activity, charge carrier dynamics is compared between the Ta₃N₅ prepared by 0.25 hour nitridation of KTaO₃ and long-time nitridation of Ta₂O₅. We selectively monitor the dynamics of free electrons and VB holes at 3435 and 545 nm wavelengths, respectively. The responsible optical transitions probed at these wavelengths are depicted in Scheme 1. It has been consistently demonstrated that probing at 3435 nm (IR region) typically monitors the intra-band transition close to the CB levels. Thus, 3435 nm probe offers unique information on the dynamics of free/shallowly trapped electrons.^{24,29,34,35} With regard to the hole dynamics, 545 nm probe is employed. As depicted in Fig. 1A, the absorption peak observed at 545 nm is



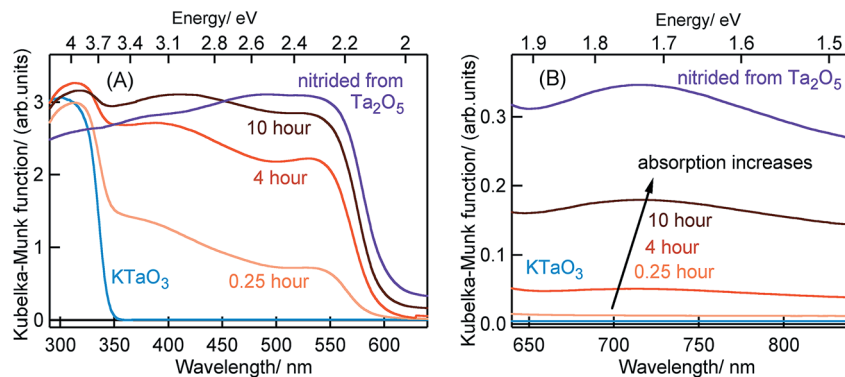
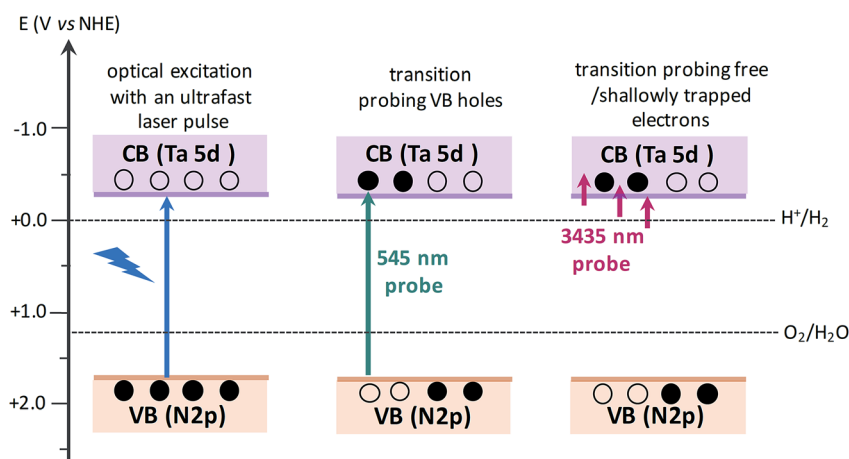


Fig. 1 Ground-state optical absorption spectra measured in the diffuse reflection mode in the (A) visible and (B) NIR region for KTaO_3 , Ta_3N_5 nitrided from KTaO_3 for different times as indicated and Ta_3N_5 nitrided from Ta_2O_5 . The arrow in the (B) indicates an enhancement in the NIR absorption with an increase in the nitridation time of KTaO_3 .

due to the bandgap transition from N 2p (VB) to Ta 5d (CB). Monitoring the dynamics at 545 nm yields direct information on the occupancy of holes in the VB, where the recovery of the ground-state bleaching signal offers information on the hole dynamics.

2.1.1 The effect of Ta precursor on the free electron dynamics probed at 3435 nm. Fig. 2A shows the transients corresponding to free electrons (3435 nm probe) in Ta_3N_5 prepared by the nitridation of KTaO_3 for 0.25 hour. In the Ta_3N_5 prepared by the nitridation of KTaO_3 for 0.25 hour, there still exists a large amount of unreacted KTaO_3 phase. Hence, it is essential to disregard the possibility of free electron generation from the unreacted KTaO_3 . The onset in the optical absorption of KTaO_3 is ≈ 345 nm (≈ 3.6 eV). The visible light excitation employed in this study does not allow bandgap transition in KTaO_3 to generate free electrons (Fig. S1†). Hence, the free electron dynamics presented in Fig. 2A originates from the Ta_3N_5 phase formed after nitridation, not from the unreacted KTaO_3 .

Considering the TA signal is monitored in the diffuse reflection mode in the current study, a direct comparison of the signal magnitude is not straightforward due to the likely differences in the scattering/surface properties of different types of powder samples under study. Hence, we primarily rely on the decay dynamics rather than the signal magnitude while comparing between two different samples. The TA signal magnitude at 3435 nm probe for Ta_3N_5 prepared by 0.25 hour nitridation of KTaO_3 can be found in ESI (Fig. S2).† To understand the decay pathways of free electrons in Ta_3N_5 (0.25 hour nitridation of KTaO_3), the effect of pump fluence on the half lifetime, $\tau_{1/2}$, (*i.e.*, the time at which initial concentration of photogenerated electrons reach to half of its initial value) is compared. The $\tau_{1/2}$ value decreases from 200 to 5 ps upon increasing the fluence from 0.014 to 0.12 μJ per pulse. A qualitatively similar observation is noticed also for hole dynamics probed at 545 nm (Fig. S3†). Thus, the dynamics of both electron and hole is very sensitive to pump fluence from early ps time. In addition, matching of both free electron and hole



Scheme 1 Schematics illustrating the nature of electronic transitions responsible for the transient absorption signal when probed at 545 and 3435 nm, corresponding to VB holes and free electrons, respectively. When the VB holes are formed due to light absorption, a ground-state bleaching signal is expected at 545 nm probe. The redox potentials of H_2 and O_2 evolution reaction^{15,36} with respect to CB and VB positions of Ta_3N_5 shows the capability of Ta_3N_5 to participate in overall water splitting.



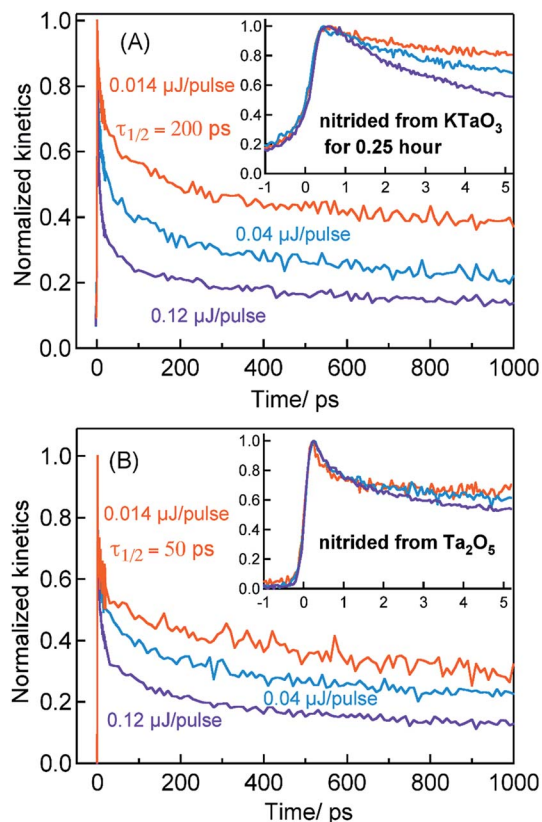


Fig. 2 The fs-TDR time profiles corresponding to free electron dynamics probed at 3435 nm for Ta_3N_5 nitrided from (A) KTaO_3 for 0.25 hour and (B) from Ta_2O_5 . 440 nm pump is used to conduct these measurements.

dynamics is noticed (Fig. S4[†]). Fitting the transients in Fig. 2A to a second-order type recombination model showed good agreement (Fig. S5[†]). Taking these observations collectively, it is inferred that photogenerated carriers in Ta_3N_5 (0.25 hour nitridation of KTaO_3) decay primarily *via* second-order type recombination of free (untrapped) electrons with holes. If the charge carriers decay *via* trapping to defect states (generally related to first-order type recombination), fluence-independent dynamics is expected, which is however not observed.

Fig. 2B displays the normalized electron kinetics in Ta_3N_5 prepared by the nitridation of Ta_2O_5 . The TA signal magnitude at 3435 nm probe for Ta_3N_5 prepared by nitridation of Ta_2O_5 can be found in ESI (Fig. S2).[†] Despite increasing the pump fluence by a factor of approximately nine, the electron dynamics in early 5 ps (Fig. 2B inset) is virtually unaffected suggesting the occurrence of electron trapping. A similar observation is noticed also at 580 nm pump (Fig. S6[†]). The $\tau_{1/2}$ value in Fig. 2B decreases from 50 to 7 ps upon increasing the fluence from 0.014 to 0.12 μJ per pulse. A matching between the free electron and hole dynamics was not observed (Fig. S7[†]) and the transients could not be fitted to the second-order type electron-hole recombination model (Fig. S5[†]). Combining these observations, it is concluded that the major part of electrons in Ta_3N_5 prepared by the nitridation of Ta_2O_5 decays *via* trapping rather than second-order type direct electron-hole recombination.

2.1.2 Electron transfer process from Ta_3N_5 prepared by 0.25 hour nitridation of KTaO_3 to the Rh cocatalyst. To reveal the mechanism of OWS, the effect of Rh cocatalyst loading and the water interface on charge carrier dynamics in Ta_3N_5 prepared by 0.25 hour nitridation of KTaO_3 is carried out. Typically, metal cocatalyst particles are loaded on the surface of the photocatalyst to enhance the overall efficiency of photocatalytic reaction. In the case of Ta_3N_5 , Rh metal cocatalyst particles are photodeposited on the surface by the photo-deposition process. The Rh cocatalyst acts as an electron acceptor to facilitate the reduction reaction with water in generating H_2 . A faster electron (3435 nm probe) decay starting from ≈ 5 ps was noticed for Ta_3N_5 prepared by 0.25 hour nitridation of KTaO_3 after Rh cocatalyst loading (Fig. S8[†]) suggests electron transfer process to the Rh cocatalyst. To further corroborate this process, the effect of water interface on the electron transfer to Rh cocatalyst is investigated. Considering water is prone to absorb IR probe light, 545 nm probe which corresponds to VB holes is employed. Fig. 3 compares the VB hole dynamics (545 nm probe) in the presence of water interface for Ta_3N_5 prepared by 0.25 hour nitridation of KTaO_3 with and without Rh cocatalyst loading. Clearly, a longer hole lifetime is noticed from ≈ 5 ps after loading Rh cocatalyst. As concluded from Fig. 2A, electrons decay majorly *via* recombination with VB holes. However, when electrons are transferred to Rh cocatalyst present on the surface, the number of electrons that recombine with VB holes is decreased. As a result, the efficiency of electron-hole recombination is reduced eventually resulting in a prolonged lifetime for VB holes due to slow recovery of the ground-state bleaching. A longer hole lifetimes is particularly beneficial to promote water oxidation reaction to generate O_2 and to realise OWS. Observing electron transfer process corroborates the notion that electron trapping by the structural/electronic defects is an inefficient process, which is in agreement with the virtually absent ground-state NIR absorption (Fig. 1B) and the single-crystal nature of the Ta_3N_5 .

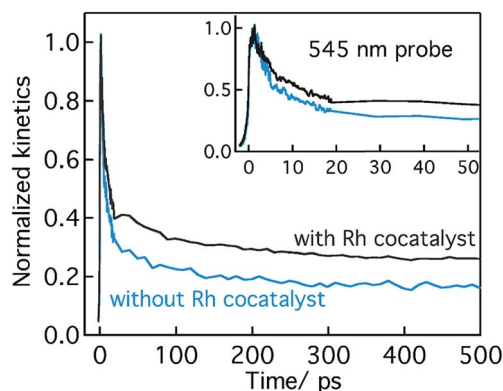


Fig. 3 The fs-TDR time profiles corresponding to VB holes probed at 545 nm for 0.25 hour nitridation of KTaO_3 with and without Rh cocatalyst loading. Both these transients are recorded in the presence of water under the same pump fluence (440 nm pump; 0.23 μJ per pulse). Note that these experiments are conducted in water media and signs of these transients are inverted for easy comparison.



formed upon nitriding KTaO_3 for 0.25 hour as reported previously by using STEM.¹⁹

2.1.3 On the inadequate H_2 evolution in Ta_3N_5 prepared by the nitridation of Ta_2O_5 . For Ta_3N_5 prepared by the nitridation of Ta_2O_5 , due to electron trapping within the bulk, efficient migration of electrons towards the surface cannot be anticipated. As a result, electron transfer to the Rh was indeed not observed (Fig. S8†). Instead, hole transfer to Rh cocatalyst was observed as evidenced by the longer electron lifetime. We do not completely understand why electron transfer to the Rh cocatalyst was not observed here. From XPS study, the Rh cocatalyst was confirmed to be metallic but not as an oxide (Fig. S9†). One of the plausible mechanisms is due to the occurrence of faster/efficient electron trapping is expected to reduce the electron mobility and the number of free electrons available for transfer to the Rh cocatalyst. This is the likely reason behind noticing inadequate H_2 evolution (compared to O_2) and consequently the absence of OWS activity in Ta_3N_5 prepared by the nitridation of Ta_2O_5 . This observation is in contrast to the Ta_3N_5 prepared by 0.25 hour nitridation of KTaO_3 which shows OWS. This comparison further corroborates the effect of Ta precursor on the electron dynamics and on the electron transfer to the Rh cocatalyst.

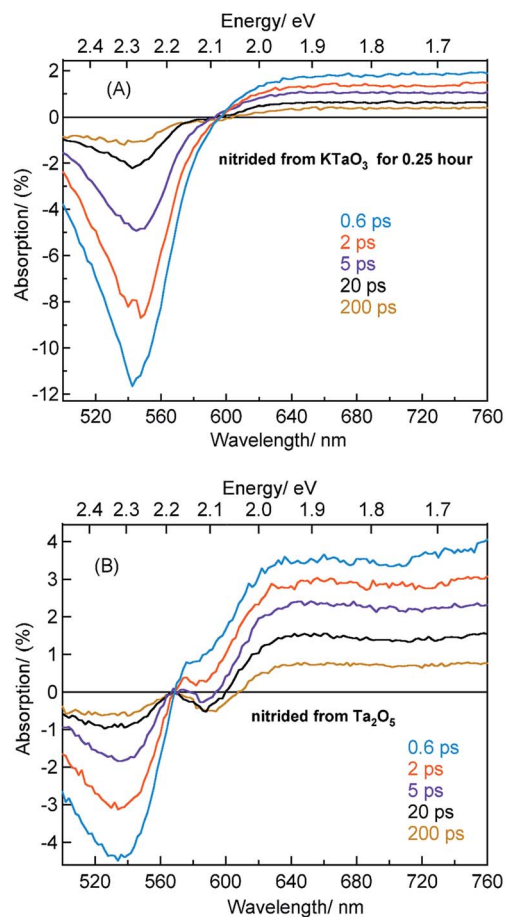


Fig. 4 TA spectra in the visible region at different delay times for Ta_3N_5 prepared by the nitridation of (A) KTaO_3 for 0.25 hour and (B) Ta_2O_5 . 440 nm pump with a fluence of $0.48 \mu\text{J}$ per pulse is used while conducting these measurements.

2.1.4 The effect of Ta precursor on the hole trapping process. Fig. 4 displays the TA spectra of Ta_3N_5 prepared by the nitridation of KTaO_3 (for 0.25 hour) and Ta_2O_5 in the visible region at different delay times starting from 0.6 to 200 ps. The TA spectra provide useful information on the presence of defect states within the bandgap. Before discussing these results, it is important to associate the spectral features with the possible transitions and/or defect states. Typically, the TA spectrum of Ta_3N_5 shows three prominent features: (i) a bleaching (negative) signal around 545 nm, (ii) a second bleaching signal at 590 nm, and (iii) a broad positive TA signal between 650 to 750 nm. Let us discuss each one of this contribution.

The bleaching signal at 545 nm is due to the depopulation of the ground state due to band gap transition (Fig. 1A and Scheme 1) and is correlated to the VB free hole dynamics. The second bleaching (negative) signal at 590 nm probe cannot be explained unless considering the possibility of an optical transition from a filled defect state. To this end, the electronic transition from a positively charged nitrogen vacancies (V_N^+) to the CB is tentatively proposed. A previous prediction using first principles density functional theory calculations is in good agreement with this assignment.^{30,37} However, this transition is not discernible in Fig. 1A, which is likely due to low absorption coefficient typical to transitions involving defect states. By correlating the energetic positions of bleaching at 590 nm with 545 nm, the V_N^+ states are situated at 0.17 eV above the VB maxima. When holes from the VB are trapped to V_N^+ states, V_N^+ is converted to V_N^{2+} or V_N^{3+} , which reduces the density of V_N^+ states. As a result, possibility of the V_N^+ to the CB optical transition at 590 nm reduces and eventually shows bleaching signal (Fig. S10†). We observe the 590 nm bleaching already at 0.6 ps delay time which indicates the ultrafast hole trapping process from the VB to V_N^+ . Thus, the presence of hole trapping process can be tentatively related to 590 nm bleaching. In an earlier report from Ziani *et al.* also observed the bleaching signal at 590 nm for Ta_3N_5 (nitrided from Ta foil) thin films.³⁸ In a recent microsecond TAS study, trapped holes are found to be probed at 590 nm in Ta_3N_5 (nitrided from Ta_2O_5).³⁴ Combining these observations, the presence of bleaching signal at 590 nm can be correlated to the formation of hole trapping nitrogen vacancies above the VB. The broad positive TA signal between 650 to 750 nm is tentatively assigned to shallowly trapped electrons (Fig. S11†) and will be discussed in Section 2.2.

Fig. 4A shows the TA spectra of Ta_3N_5 prepared by the nitridation of KTaO_3 for 0.25 hour. The free VB hole dynamics shows time-dependent change as observed at 545 nm bleaching. The virtually absent second bleaching at 590 nm indicates that the density of V_N^+ defect states is too low to promote hole trapping. This conclusion is in good agreement with noticing a pronounced pump fluence dependent free VB hole dynamics probed at 545 nm (Fig. S3†). The occurrence of hole trapping is found to be an incompetent process in Ta_3N_5 prepared by 0.25 hour nitridation of KTaO_3 .

Fig. 4B shows the TA spectra of Ta_3N_5 prepared by the nitridation of Ta_2O_5 , which is different compared to Fig. 4A in terms of the second bleaching signal at 590 nm in addition to



the one at 545 nm. Effect of pump fluence on the dynamics of VB free holes (545 nm probe) is found to be less prominent particularly in early 5 ps time scale (Fig. S12[†]). This observation suggests the presence of hole trapping V_N^- defect states above the VB. Noticing the pronounced second bleaching signal at 590 nm with time indeed support the occurrence of hole trapping. In short, the presence of second bleaching at 590 nm indicates an enhancement in the formation of hole trapping V_N^- defect states.

Conclusion from Section 2.1 is briefly summarized as follows. Depending on the Ta precursor employed for the nitridation reaction to prepare Ta_3N_5 , a distinctive electron and hole dynamics is rationalized. In Ta_3N_5 prepared by the nitridation of Ta_2O_5 , trapping of both electrons and holes is observed. In particular, due to efficient electron trapping, electron transfer to the Rh cocatalyst is not observed which is essential to realize H_2 evolution and consequently OWS activity.

In the case of Ta_3N_5 prepared by the 0.25 hour nitridation of $KTaO_3$, charge carriers are free and the decay is primarily *via* second-order recombination between CB electrons and VB holes. The decay of both electrons and holes by trapping is found to be ineffective. This conclusion is in good agreement with noticing virtually absent ground-state NIR absorption (Fig. 1B) and with the single-crystal nature of the Ta_3N_5 without surface defects.¹⁹ Hence, efficient migration of both electrons and holes toward the surface is ensured eventually prompting the OWS activity.

2.2 The effect of nitridation time on the charge carrier dynamics in Ta_3N_5 prepared by the nitridation of $KTaO_3$

The efficiency of OWS is found to be high for 0.25 hour nitrided Ta_3N_5 . However, the OWS activity was reduced by a factor of approximately five with an increase in the nitridation time from 0.25 to 10 hour.¹⁹ Hence, investigating the effect of nitridation time on the carrier dynamics is essential to understand the origin of this phenomena.

2.2.1 Comparing the free electron dynamics in Ta_3N_5 prepared by 0.25 and 10 hour nitridation of $KTaO_3$.

The TA signal magnitude at 3435 nm probe for Ta_3N_5 prepared by 10 hour nitridation of $KTaO_3$ can be found in ESI (Fig. S2).[†] As depicted in Fig. 5A, despite changing the pump fluence by a factor of approximately nine, the free electron dynamics (3435 nm probe) in early 5 ps for Ta_3N_5 prepared by 10 hour nitridation of $KTaO_3$ is barely affected. Comparing between Fig. 2A and 5A, a pronounced electron trapping in Ta_3N_5 nitrided from $KTaO_3$ for 10 hour compared to 0.25 hour can be perceived. This notion is further corroborated by noticing: (i) a negligible influence on the free electron dynamics upon loading Rh cocatalyst (Fig. S13[†]), and (ii) a slow recovery of the VB free holes due to a reduction in the number of free electrons (because of trapping) available for recombination (Fig. S14[†]). These observations indicate the occurrence of electron trapping in Ta_3N_5 nitrided from $KTaO_3$ for 10 hour. This conclusion is further supported by noticing an enhancement in the ground-state NIR absorption (Fig. 1B) due to defect formation within the bandgap.

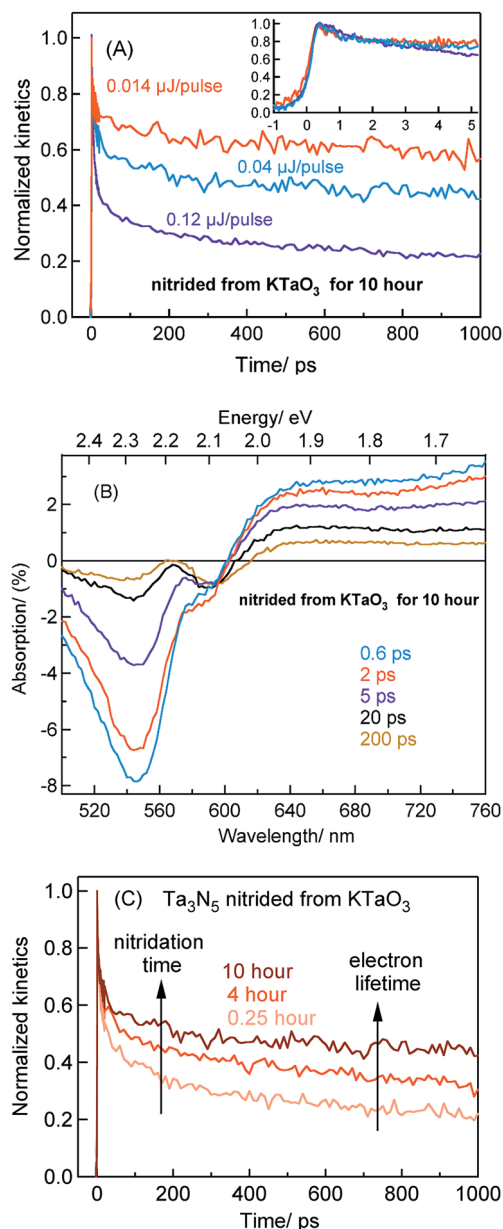


Fig. 5 (A) The fs-TDR time profiles corresponding to free electron dynamics (440 nm pump and 3435 nm probe) in Ta_3N_5 prepared by the 10 hour nitridation of $KTaO_3$. (B) TA spectra (440 nm pump; 0.48 μJ per pulse) of Ta_3N_5 prepared by the 10 hour nitridation of $KTaO_3$. (C) Effect of nitridation time on the free electron dynamics (440 nm pump; 0.04 μJ per pulse) in Ta_3N_5 prepared by the nitridation of $KTaO_3$.

2.2.2 The effect of nitridation time on the hole trapping.

To investigate the effect of increasing the nitridation time on the hole dynamics, TA spectra of Ta_3N_5 prepared by 10 hour nitridation of $KTaO_3$ is recorded (Fig. 5B). The presence of a second bleaching signal at 590 nm (which is absent in Fig. 4A) indicates the formation of hole trapping V_N^- defect states. A similar observation is also noticed for Ta_3N_5 prepared by 4 hours nitridation of $KTaO_3$ (Fig. S15[†]). To rule out any potential contribution to the bleaching signal at 590 from the unreacted $KTaO_3$, TA spectra of $KTaO_3$ is recorded. As expected, no such



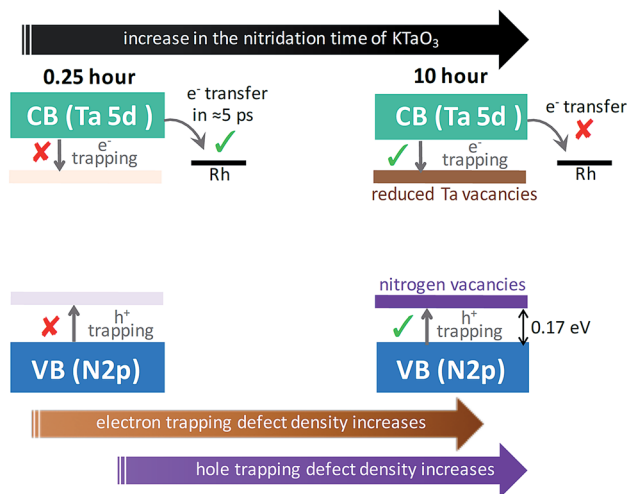
distinct bleaching signal at 590 nm reminiscent to that of Ta_3N_5 is detected for KTaO_3 (Fig. S16†).

To further rationalize the effect of nitridation time on the hole trapping process, the electron lifetime probed at 3435 nm is compared as a function of KTaO_3 nitridation time. As shown in Fig. 5C, the electron lifetime gets longer with an increase in the nitridation time. This observation can be explained as follows. Due to hole trapping, fewer holes are available for electrons to decay by recombination and hence a longer electron lifetime can be expected. Thus, from Fig. 5B and C, an enhancement in the efficiency of hole trapping process with an increase in the nitridation time is inferred.

Note that, we have not considered the effect of dark carrier (which depends on the relative position of Fermi level with respect to the CB and the defect density) on the decay dynamics of electrons and holes in Ta_3N_5 prepared by the nitridation of KTaO_3 . This is due to experimental difficulty in determining the precise concentration of dark carrier or Fermi level position in these powder samples which is in a composite form with KTaO_3 . In this regard, a further systematic investigation of the effect of nitridation time on the relative shift in the Fermi level (with respect to the CB) and determining the corresponding dark carrier density would be helpful.

To put the results from Fig. 5 in perspective, it is instructive to understand the plausible mechanism by which defects are formed and their relation with the nitridation time. In this regard, a tentative model is proposed. The temperature employed during the nitridation reaction of KTaO_3 is 1173 K. The nitrogen source used for nitriding the Ta precursor is NH_3 gas, which decomposes into N_2 and H_2 under the experimental condition. Ideally, the Ta 5d orbital of the CB is present in Ta^{5+} states, which ensures to maintain charge neutrality and avoid electron trapping. However, during the nitridation reaction, the time exposed to H_2 increases with nitridation time. This H_2 gas acts as a reducing agent to convert the Ta^{5+} states to form Ta^{4+} and/or Ta^{3+} , which are proposed to act as electrons trapping centres. This notion is in agreement with a previous report demonstrating the formation of reduced Ta vacancies using synchrotron-excited X-ray photoelectron spectroscopy (XPS).³⁹ Similarly, the formation of nitrogen (V_{N}) vacancies which are associated to hole trapping centres becomes inevitable, particularly for nitridation time exceeding 0.25 hour. Maintaining the ideal Ta_3N_5 stoichiometric ratio *via* synchronous substitution of nitrogen to oxygen (present in Ta precursor) is the key to control the defect formation.

Based on the conclusions drawn from Fig. 5, the formation of both electron and hole trapping centres within the bandgap of Ta_3N_5 is pronounced with an increase in the nitridation time. The efficiency of charge transport from the bulk to the surface is negatively affected by carrier trapping. As a consequence, inefficient electron transfer to the Rh cocatalyst is observed and the number of free carriers available at the surface of Ta_3N_5 for photocatalytic reaction will be reduced. Due to these two combined effects, the efficiency of OWS is found to be reduced with increase in the nitridation time from 0.25 to 10 hour (Scheme 2). From our previous STEM study, the location of charge trapping defects in long time (>0.25 hour) nitrided Ta_3N_5



Scheme 2 Proposed model to explain the effect of nitridation time on the formation of defects within the bandgap and its impact on the efficiency of OWS reaction. The Fermi level positions of Ta_3N_5 and Rh metal are not considered while drawing the energetic position of Rh metal cocatalyst with respect to the CB.

is found to be predominantly on the surface of Ta_3N_5 .¹⁹ Hence, removing these surface defect states either by chemical treatment or by surface modification may further enhance the efficiency of OWS activity for long-time nitrided Ta_3N_5 . A future work is necessary to unravel the electronic nature of the bulk and surface defect states in detail.

2.3 How to further enhance the OWS activity of Ta_3N_5 (0.25 hour nitrided from KTaO_3) under the visible light?

The AQE value at 500 nm is ten times smaller compared to 420 nm.¹⁹ Hence, one of the challenges is to enhance the AQE under visible light excitation. To discern whether the low AQE in visible light is due to inefficient charge generation, free electron dynamics (3435 nm probe) is compared between 580 and 440 nm excitation (Fig. S17†). At 580 nm excitation, a similarity in terms of both carrier lifetime and pump fluence dependence is noticed to that of 440 nm pump. In addition, electron transfer to the Rh cocatalyst is also noticed at 580 nm pump (Fig. S18†). Besides that, Ta_3N_5 is widely considered to have an indirect bandgap. Note that 580 nm pump is very close to the band edge of the Ta_3N_5 , however, neither free electron lifetime nor electron photogeneration is affected. Hence, indirect bandgap nature of the Ta_3N_5 does not sufficiently explain the lower activity under the visible light excitation. Alternatively, the absorption coefficient value of Ta_3N_5 at 530 nm is lower by a factor of five compared to 400 nm.^{31,40,41} Upon 0.25 hour nitridation of KTaO_3 , the actual amount of Ta_3N_5 phase formed on KTaO_3 is $\approx 2\%$.¹⁹ Thus, a low AQE in the visible region is attributed to a combination of the lower absorption coefficient and the presence of only 2% of Ta_3N_5 which can potentially absorb visible light, but not due to inefficient charge photogeneration. Thus, increasing the effective density of Ta_3N_5 on KTaO_3 is expected to further improve the AQE under visible light excitation. To this end, multiple nitridation of KTaO_3 (while keeping the



nitridation time as 0.25 hour) to expose the unreacted KTaO_3 surface in the direction of NH_3 flow forms one of the approaches.

3. Conclusions

A detailed ultrafast TAS study allowed us to elucidate a correlation between the synthesis procedure – structural defects – carrier dynamics – efficiency of OWS. Selectively probing the dynamics of electrons and holes allowed us to develop a phenomenological model that describes the role of Ta precursor employed to prepare Ta_3N_5 and the effect of nitridation time in the context of OWS.

In addition to H_2 evolution being fundamentally low compared to O_2 , the OWS has not been realized for Ta_3N_5 prepared by nitriding Ta_2O_5 . The answer to this longstanding question was found to originate from the absence of electron transfer to the Rh cocatalyst by the virtue of efficient electron trapping.

When Ta_3N_5 is prepared by employing KTaO_3 as a Ta precursor for nitridation reaction, OWS was observed for the first time. For 0.25 hour nitrided Ta_3N_5 , which shows the highest OWS activity, charge carriers decayed by second-order type electron-hole recombination but not by trapping. Due to single-crystal and defect-free nature of 0.25 hour nitrided Ta_3N_5 , efficient migration of both electrons and holes towards the surface was achieved. Thus, to realize OWS, avoiding carrier trapping process was found to be indispensable.

When the nitridation time of KTaO_3 was increased from 0.25 to 10 hour, the efficiency of OWS reduced by a factor of approximately five. This was attributed to the formation of defects which promoted both electron and hole trapping eventually affecting the efficient migration of charge carriers to the surface. These transient optical studies yielded key insights into the factors determining the efficiency of solar driven OWS by Ta_3N_5 .

4. Experimental section

4.1 Synthesis of Ta_3N_5 by the nitridation of KTaO_3

KTaO_3 particles were fabricated by a conventional solid-state reaction method. Ta_2O_5 (99.9%; Kojundo Chemical Laboratory Co., Ltd.) and K_2CO_3 (99.5%; Kanto Chemical Co., Inc.) were mixed at a Ta : K molar ratio of 1 : 1.05. Excess K was added to compensate for losses by volatilization at high temperatures. The mixture was thoroughly ground in an agate mortar for 90 min in the presence of a small amount of ethanol as a dispersant. After drying, the resulting mixture was transferred into an alumina crucible and calcined at 1173 K for 1 hour and then at 1423 K for 10 hour in static air. The KTaO_3 obtained in this manner was washed with ultrapure water at 343 K for 2 hour and centrifuged twice to remove any residual K_2CO_3 . The powder was then completely dried by heating at 343 K overnight. Subsequently, the as-prepared KTaO_3 was subjected to a nitridation process to obtain Ta_3N_5 grown on KTaO_3 . KTaO_3 (0.5 g) was transferred into an alumina tube and nitrided at

1173 K for 0.25, 4 and 10 hour under a flow of NH_3 gas at 100 mL min^{-1} .

4.2 Synthesis of Ta_3N_5 by the nitridation of Ta_2O_5

The commercially available Ta_2O_5 (99.9%; Kojundo Chemical Laboratory Co., Ltd.) was subjected to the nitridation at 1173 K for 20 hours under a flow of NH_3 gas at 100 mL min^{-1} .

4.3 Photodeposition of a Rh cocatalyst on Ta_3N_5

A Rh cocatalyst for H_2 evolution was loaded on Ta_3N_5 photocatalysts by a previously-reported photodeposition process.¹⁹ A Rh core was photodeposited using $\text{RhCl}_3 \cdot 3\text{H}_2\text{O}$ (Kanto Chemical Co., Inc.) as the metal precursor. This was accomplished by dispersing the photocatalyst powder in 150 mL of an aqueous methanol solution (10 vol%) containing the metal precursor. The pH of this solution was not adjusted but the temperature was maintained at 288 K by circulating cooling water. The suspension was evacuated to completely remove dissolved air and then exposed to visible light ($\lambda \geq 420$ nm) with continuous stirring. The photodeposition of Rh was conducted approximately for 3 hours.

4.4 Femtosecond time-resolved diffuse reflectance (TDR) spectroscopy

In fs-transient diffuse reflectance (fs-TDR) measurements, a femtosecond Ti:sapphire laser with a regenerative amplifier (Spectra-Physics, Solstice, wavelength of 800 nm, pulse width of 100 fs, pulse energy of 3.5 mJ per pulse, and repetition rate of 1 kHz) was used as a light source. The output from the laser was split into four paths for the excitation of two optical parametric amplifiers (OPAs: Spectra-Physics, TOPAS Prime), the white-light-continuum generation by focusing the fundamental light (800 nm) into a sapphire plate, and the second and third harmonic generations of the fundamental light (800 nm) by using BBO ($\beta\text{-BaB}_2\text{O}_4$) crystals. For generating the 440 or 580 nm pump pulse, light from one of the OPA was used. For the probe pulse, a white light continuum covering from 500 nm to 1600 nm and a 3435 nm probe light generated from the other OPA with a difference-frequency generation crystal was used. The time resolution of the system was about 140 fs. The powder samples are taken in 1 mm quartz cuvettes. The diameter of the pump beam on the sample was about 0.5 mm as observed with a charge-coupled-device (CCD) camera. Amplified Si photodetector is used to measure the TA spectra in the visible region. Liquid-nitrogen-cooled mercury-cadmium-telluride (HgCdTe) photodetector is used for IR probe (3435 nm) experiments. The diffusely reflected light from the sample was passed through a grating monochromator (Princeton Instruments, Acton SP2150) for data acquisition. The transient absorption intensity of the TDR measurements is presented as percentage absorption, where absorption (%) = $100(1 - R/R_0)$, using R and R_0 as the intensities of the diffusely reflected light with and without excitation, respectively. A more detailed description of the fs-TDR setup is available elsewhere.^{42,43}



4.5 X-ray photoelectron spectroscopy (XPS)

The XPS spectra were obtained using X-ray photoelectron ULVAC-PHI, INC PHI Quantera 2 spectrometer.

Conflicts of interest

There are no conflicts of interest to declare.

Acknowledgements

This work is supported by the “Research Project for Future Development: Artificial Photosynthetic Chemical Process (ARPCHEM)” (METI/NEDO, Japan: 2012-2022).

References

- M. G. Walter, E. L. Warren, J. R. McKone, S. W. Boettcher, Q. X. Mi, E. A. Santori and N. S. Lewis, *Chem. Rev.*, 2010, **110**, 6446–6473.
- X. B. Chen, S. H. Shen, L. J. Guo and S. S. Mao, *Chem. Rev.*, 2010, **110**, 6503–6570.
- F. E. Osterloh, *Chem. Soc. Rev.*, 2013, **42**, 2294–2320.
- Q. Wang, T. Hisatomi, Q. X. Jia, H. Tokudome, M. Zhong, C. Z. Wang, Z. H. Pan, T. Takata, M. Nakabayashi, N. Shibata, Y. B. Li, I. D. Sharp, A. Kudo, T. Yamada and K. Domen, *Nat. Mater.*, 2016, **15**, 611–615.
- T. Yamada and K. Domen, *ChemEngineering*, 2018, **2**, 36, DOI: 10.3390/chemengineering2030036.
- T. Hisatomi, J. Kubota and K. Domen, *Chem. Soc. Rev.*, 2014, **43**, 7520–7535.
- A. Kudo and Y. Miseki, *Chem. Soc. Rev.*, 2009, **38**, 253–278.
- Q. Wang, T. Hisatomi, M. Katayama, T. Takata, T. Minegishi, A. Kudo, T. Yamada and K. Domen, *Faraday Discuss.*, 2017, **197**, 491–504.
- D. M. Fabian, S. Hu, N. Singh, F. A. Houle, T. Hisatomi, K. Domen, F. E. Osterloh and S. Ardo, *Energy Environ. Sci.*, 2015, **8**, 2825–2850.
- E. Nurlaela, A. Ziani and K. Takanabe, *Mater. Renew. Sustain. Energy*, 2016, **5**, 18, DOI: 10.1007/s40243-40016-40083-z.
- C. Zhen, R. Z. Chen, L. Z. Wang, G. Liu and H. M. Cheng, *J. Mater. Chem. A*, 2016, **4**, 2783–2800.
- G. J. Liu, S. Ye, P. L. Yan, F. Q. Xiong, P. Fu, Z. L. Wang, Z. Chen, J. Y. Shi and C. Li, *Energy Environ. Sci.*, 2016, **9**, 1327–1334.
- X. J. Feng, T. J. LaTempa, J. I. Basham, G. K. Mor, O. K. Varghese and C. A. Grimes, *Nano Lett.*, 2010, **10**, 948–952.
- M. Harb, P. Sautet, E. Nurlaela, P. Raybaud, L. Cavallo, K. Domen, J. M. Basseta and K. Takanabe, *Phys. Chem. Chem. Phys.*, 2014, **16**, 20548–20560.
- W. J. Chun, A. Ishikawa, H. Fujisawa, T. Takata, J. N. Kondo, M. Hara, M. Kawai, Y. Matsumoto and K. Domen, *J. Phys. Chem. B*, 2003, **107**, 1798–1803.
- M. Zhong, T. Hisatomi, Y. Sasaki, S. Suzuki, K. Teshima, M. Nakabayashi, N. Shibata, H. Nishiyama, M. Katayama, T. Yamada and K. Domen, *Angew. Chem., Int. Ed.*, 2017, **56**, 4739–4743.
- Y. W. Wang, D. Z. Zhu and X. X. Xu, *ACS Appl. Mater. Interfaces*, 2016, **8**, 35407–35418.
- T. H. Chiang, H. Lyu, T. Hisatomi, Y. Goto, T. Takata, M. Katayama, T. Minegishi and K. Domen, *ACS Catal.*, 2018, **8**, 2782–2788.
- Z. Wang, Y. Inoue, T. Hisatomi, R. Ishikawa, Q. Wang, T. Takata, S. Chen, N. Shibata, Y. Ikuhara and K. Domen, *Nat. Catal.*, 2018, **1**, 756–763.
- X. Liu, L. Zhao, K. Domen and K. Takanabe, *Mater. Res. Bull.*, 2014, **49**, 58–65.
- M. Xiao, B. Luo, M. Q. Lyu, S. C. Wang and L. Z. Wang, *Adv. Energy Mater.*, 2018, **8**, 8.
- S. S. Chen, Y. Qi, Q. Ding, Z. Li, J. Y. Cui, F. X. Zhang and C. Li, *J. Catal.*, 2016, **339**, 77–83.
- K. Maeda, N. Nishimura and K. Domen, *Appl. Catal., A*, 2009, **370**, 88–92.
- Y. Suzuki, D. H. K. Murthy, H. Matsuzaki, A. Furube, Q. Wang, T. Hisatomi, K. Domen and K. Seki, *J. Phys. Chem. C*, 2017, **121**, 19044–19052.
- S. R. Pendlebury, X. L. Wang, F. Le Formal, M. Cornuz, A. Kafizas, S. D. Tilley, M. Gratzel and J. R. Durrant, *J. Am. Chem. Soc.*, 2014, **136**, 9854–9857.
- R. B. Singh, H. Matsuzaki, Y. Suzuki, K. Seki, T. Minegishi, T. Hisatomi, K. Domen and A. Furube, *J. Am. Chem. Soc.*, 2014, **136**, 17324–17331.
- A. Furube, T. Asahi, H. Masuhara, H. Yamashita and M. Anpo, *J. Phys. Chem. B*, 1999, **103**, 3120–3127.
- R. Godin, Y. Wang, M. A. Zwijnenburg, J. W. Tang and J. R. Durrant, *J. Am. Chem. Soc.*, 2017, **139**, 5216–5224.
- D. H. K. Murthy, H. Matsuzaki, Q. Wang, Y. Suzuki, K. Seki, T. Hisatomi, T. Yamada, A. Kudo, K. Domen and A. Furube, *Sustainable Energy Fuels*, 2019, **3**, 208–218.
- T. Jing, Y. Dai, X. C. Ma, W. Wei and B. B. Huang, *RSC Adv.*, 2015, **5**, 59390–59397.
- B. A. Pinaud, P. C. K. Vesborg and T. F. Jaramillo, *J. Phys. Chem. C*, 2012, **116**, 15918–15924.
- B. A. Pinaud, A. Vailionis and T. F. Jaramillo, *Chem. Mater.*, 2014, **26**, 1576–1582.
- J. J. Wang, A. B. Ma, Z. S. Li, J. H. Jiang, J. Y. Feng and Z. G. Zou, *Phys. Chem. Chem. Phys.*, 2015, **17**, 23265–23272.
- J. J. M. Vequizo, M. Hojamberdiev, K. Teshima and A. Yamakata, *J. Photochem. Photobiol., A*, 2018, **358**, 315–319.
- A. Yamakata, J. J. M. Vequizo and M. Kawaguchi, *J. Phys. Chem. C*, 2015, **119**, 1880–1885.
- E. Nurlaela, Y. Sasaki, M. Nakabayashi, N. Shibata, T. Yamada and K. Domen, *J. Mater. Chem. A*, 2018, **6**, 15265–15273.
- J. J. Wang, A. B. Ma, Z. S. Li, J. H. Jiang, J. Y. Feng and Z. G. Zou, *Phys. Chem. Chem. Phys.*, 2015, **17**, 8166–8171.
- A. Ziani, E. Nurlaela, D. S. Dhawale, D. A. Silva, E. Alarousu, O. F. Mohammed and K. Takanabe, *Phys. Chem. Chem. Phys.*, 2015, **17**, 2670–2677.
- S. Khan, M. J. L. Santos, C. F. Malfatti, J. Dupont and S. R. Teixeira, *Chem.–Eur. J.*, 2016, **22**, 18501–18511.



- 40 J. M. Morbec, I. Narkeviciute, T. F. Jaramillo and G. Galli, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2014, **90**, 10.
- 41 H. Hajibabaei, O. Zandi and T. W. Hamann, *Chem. Sci.*, 2016, **7**, 6760–6767.
- 42 T. Asahi, A. Furube, H. Fukumura, M. Ichikawa and H. Masuhara, *Rev. Sci. Instrum.*, 1998, **69**, 361–371.
- 43 A. Furube, Z. S. Wang, K. Sunahara, K. Hara, R. Katoh and M. Tachiya, *J. Am. Chem. Soc.*, 2010, **132**, 6614–6615.

