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Effect of charge selective contacts on the quasi Fermi level splitting of CuGa_3Se_5 thin film photocathodes for hydrogen evolution and methylviologen reduction†

Ye Cheng,^a Chengcan Xiao,^a Behzad Mahmoudi,^b Roland Scheer,^c A. Wouter Maijenburg ^b and Frank E. Osterloh ^{*a}

The copper chalcopyrite CuGaSe_2 and the defect-related phase CuGa_3Se_5 are promising photocathode materials for solar hydrogen generation. Existing devices exhibit photocurrents nearing 68% (CuGa_3Se_5)–86% (CuGaSe_2) of their theoretical limit but they are plagued by photovoltage losses that reduce their energy conversion efficiency. To evaluate the reasons, we determine the light intensity dependent quasi-Fermi level splitting (QFLS) in p- CuGa_3Se_5 /liquid junctions for the first time, using the vibrating Kelvin probe surface photovoltage technique. The QFLS, or internal photovoltage, corresponds to the maximum electrochemical work a photoelectrode can perform under a given illumination condition. In the presence of water or methylviologen (2+/+) as electron acceptors, the QFLS of Mo/p- CuGa_3Se_5 /liquid junctions reaches 0.22 to 0.29 V, respectively, under 49 mW cm^{-2} 400 nm illumination, while the QFLS of a FTO/p- CuGa_3Se_5 photoelectrode is only 0.15–0.16 V. The lower voltage of the latter is attributed to a Schottky junction at the back contact, which limits majority charge carrier (hole) extraction from the semiconductor. Photovoltage losses also result from Fermi level pinning of the minority carriers at surface states 0.5 eV above the CGSe valence band. This problem can be overcome by chemical bath deposition of a CdS overlayer, which functions as a selective contact for electron extraction from CuGa_3Se_5 and which raises the QFLS to 0.44 V at 49 mW cm^{-2} . No significant QFLS enhancement occurs upon adsorption of Cd^{2+} ions to the CuGa_3Se_5 electrode surface, suggesting that Cd^{2+} adsorption alone does not remove the Fermi level pinning effect. Overall, these results provide a better understanding of the effect of surface treatments and charge selective contacts on the photovoltage of CuGa_3Se_5 photoelectrodes and indicate pathways to improve its solar fuel conversion efficiency.

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Broader context

The photoelectrochemical water splitting reaction is a promising pathway to hydrogen fuel from solar energy. The copper chalcopyrites function as photocathodes for solar hydrogen generation but are plagued by photovoltage losses that reduce their energy conversion efficiency. To evaluate the reasons, we determine the light intensity dependent quasi-Fermi level splitting (QFLS) in p- CuGa_3Se_5 /liquid junctions for the first time, using the Vibrating Kelvin Probe Surface Photovoltage Spectroscopy (VK-SPS) technique. The QFLS, or internal photovoltage, corresponds to the maximum electrochemical work a photoelectrode can perform under a given illumination condition. It cannot be observed with standard photoelectrochemical measurements but is accessible with VK-SPS. We find that on the photoelectrode front side, the internal photovoltage is limited by Fermi level pinning of the minority carriers at surface states 0.5 eV above the p- CuGa_3Se_5 valence band. Higher photovoltages are obtained after bath deposition of a CdS overlayer which functions as a selective contact for electron extraction from CuGa_3Se_5 . On the back side photovoltage losses result from a Schottky junction between FTO and p- CuGa_3Se_5 , which traps photoholes at the interface. The photovoltage can be raised by replacing the FTO with a molybdenum substrate, which has a lower barrier for hole transfer. These observations provide a better understanding of the effect of surface treatments and charge selective contacts on the solar fuel conversion efficiency of CuGa_3Se_5 photoelectrodes.

^a Department of Chemistry, University of California, Davis, CA 95616, USA. E-mail: fosterloh@ucdavis.edu

^b Center for Innovation Competence (ZIK) “SiLi-nano”, Martin Luther University Halle-Wittenberg, Karl-Freiherr-von-Fritsch-Str. 3, D-06120, Halle (Saale), Germany

^c Institute of Physics, Martin Luther University Halle-Wittenberg, Von-Danckelmann-Platz 3, D-06120, Halle (Saale), Germany

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Introduction

Due to their tunable band gaps (1.0–1.85 eV), high absorption coefficients and usability in the polycrystalline state, copper chalcopyrites, $\text{Cu}(\text{In}_x\text{Ga}_{1-x})\text{Se}_2$ (CIGS) ($x = 0, \dots, 1$), are established light absorbers in thin film photovoltaic (PV) applications.^{1–3} Recently, these materials have also been applied as photocathodes for photoelectrochemical proton reduction (Fig. 1).^{4–15} For this use, the copper-poor phase CuGa_3Se_5 (CGSe) has emerged as a champion.^{13,14} Its structure is obtained from the zinc blende (ZnS) crystal structure by replacing Zn atoms with equal amounts of Cu and Ga (Fig. 1a) and by then removing one formula unit of Cu_2Se . This increases the band gap to 1.86 eV, which provides a larger driving force for proton reduction,¹¹ allowing photocurrents of 68% of the theoretical limit (-12.1 mA cm^{-2} measured at -0.40 V vs. RHE).¹⁴ However, the photovoltage of CuGa_3Se_5 photocathodes (0.2 V, which corresponds to only 11% of the band gap) continues to be very low.^{13,14} The Sivula group reported that photovoltage losses in the related $\text{CuIn}_{0.3}\text{Ga}_{0.7}\text{S}_2$ photocathode are caused by surface Ga and In vacancies.¹⁶ Similar surface defects may limit the operation of the CuGa_3Se_5 (CGSe) photocathode. Indeed, it has been shown that surface treatment with CdS boosts proton reduction with copper and silver chalcopyrites.^{5,10,11,15,17–19}

To better understand the performance of CuGa_3Se_5 , we have analyzed the photochemical charge separation properties of several CGSe/liquid junctions with vibrating Kelvin probe liquid surface photovoltage spectroscopy (VK-SPV). In this technique, the contact potential difference (CPD) of the sample is measured with a commercially available vibrating Kelvin probe (*Besocke Delta Phi*). Illumination disturbs the distribution of the charge carriers in the sample and changes the CPD, producing a surface photovoltage (SPV), as shown in eqn (1) and Fig. 1c. This SPV signal equals the Quasi Fermi Level Splitting (QFLS) energy, according to eqn (2).^{20–24}

$$\text{SPV} = \text{CPD}(\text{light}) - \text{CPD}(\text{dark}) \quad (1)$$

$$\text{SPV} = \frac{E_{\text{F,n}} - E_{\text{F,p}}}{e} = \frac{\text{QFLS}}{e} \quad (2)$$

The QFLS, or internal photovoltage, is an important property of a photoelectrode. It corresponds to the maximum electric

potential the photoelectrode can generate under a given illumination condition.^{25–30} The QFLS cannot be obtained from standard photoelectrochemical measurements due to the lack of a direct connection to the solid–liquid interface.^{31–38} VK-SPV overcomes this problem, making QFLS data available in a contactless way.²¹

Once the QFLS is known, eqn (2) can also be used to calculate the minority carrier (electron) potential $E_{\text{F,n}}$ from the majority carrier (hole) potential $E_{\text{F,p}}$. The latter is accessible from open circuit potential measurements *via* eqn (3).^{39–41} In this equation, E_{C} is the potential of the counter electrode.

$$V_{\text{OC}} = \frac{E_{\text{F,p}}}{e} - E_{\text{C}} \quad (3)$$

Therefore, the combination of SPV and OCP measurements under illumination with a variable intensity makes it possible to obtain absolute $E_{\text{F,n}}$ and $E_{\text{F,p}}$ values for a given solid–liquid junction. In the following, this approach is used for the first time to characterize the QFLS in CGSe photocathodes, in the presence of methylviologen and/or water as electron acceptors, using Mo and FTO as back contacts, and before and after surface modification with CdS. We find that the internal photovoltage in CGSe is not limited by the water reduction overpotential, but rather by the presence of surface states located 0.5 eV above the valence band edge of the material. The FTO back contact further decreases the photovoltage of the CGSe by formation of a Schottky junction, which promotes electron–hole recombination. On the other hand, a CdS front contact promotes both $E_{\text{F,n}}$ and $E_{\text{F,p}}$. This is attributed to improved electron–hole separation at the CGSe/CdS p–n junction and the oxidizing CdS valence band maximum that rejects photoholes from CGSe. These results are a step towards reducing photovoltage losses in CGSe photoelectrodes.

Results and discussion

CGSe thin film photocathodes were grown by physical vapor deposition as reported before.¹⁴ The samples, noted as Mo–CGSe and FTO–CGSe, respectively, consist of 2 μm thick CuGa_3Se_5 on molybdenum- or FTO-coated soda-lime glass. Photographs, SEM, XRD and optical absorption spectra of the

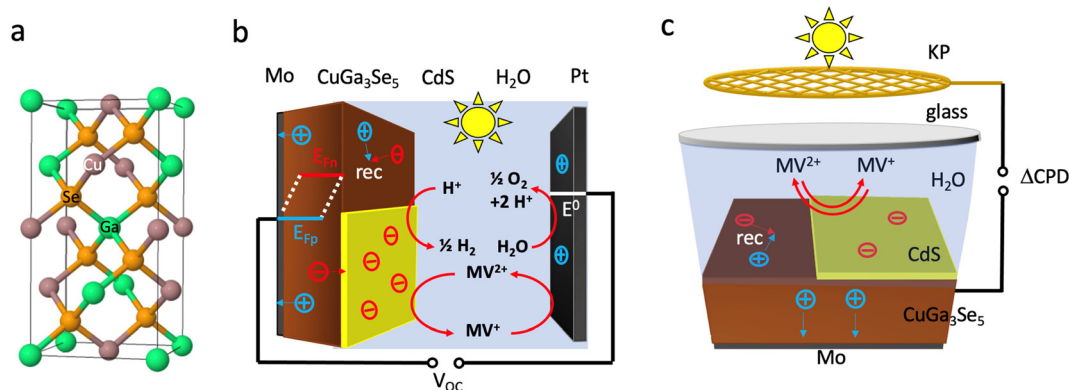


Fig. 1 (a) Crystal structure of CuGaSe_2 , the parent phase of CuGa_3Se_5 , (b) PEC cell, (c) liquid SPV configuration. MV = methylviologen.



films are shown in Fig. S1 (ESI[†]).¹⁴ Linear sweep voltammetry scans for the two electrodes immersed in a 0.05 M sodium phosphate buffer (with 0.1 M Na₂SO₄) under 400 nm illumination with an intensity of 81 mW cm⁻² are compared in Fig. 2a. Cathodic photocurrents are attributed to the hydrogen evolution reaction (HER), as confirmed by gas bubble formation. The photocurrent reaches -3 mA cm⁻² at 0 V_{RHE} for the Mo-CGSe electrode, which is approximately one-third of the photocurrent achieved in a 0.5 M H₂SO₄ electrolyte under AM 1.5G solar irradiance reported previously by the Maijenburg group.¹⁴ This difference is due to the lower proton concentration in the sodium phosphate electrolyte (pH 7) used here, which slows down the hydrogen evolution reaction (HER). The slow HER also causes the capacitive current (sawtooth shape), which represents trapping of photogenerated charge carriers at the CGSe/electrolyte interface.⁴² Based on the standard water reduction potential $E_{(H^+/H_2)}^0$ and the photocurrent onset potential E_{on} , the external photovoltage $V_{ph}(ext, LSV)$ ⁴³ of the Mo-CGSe/NaPi system is $E_{on} - E_{(H^+/H_2)}^0 = 0.32 V_{RHE} - 0.0 V_{RHE} = 0.32 V$, similar to the previously reported value.¹⁴ In comparison, the FTO-CGSe/NaPi electrode has a lower photocurrent and a later photo-onset, corresponding to a lower external photovoltage $V_{ph}(ext, LSV) = E_{on} - E_{(H^+/H_2)}^0 = 0.23 V_{RHE} - 0.0 V_{RHE} = 0.23 V$.

In order to determine the degree to which the photoelectrode performance is limited by the water reduction overpotential, the electrolyte was replaced with a solution of 0.01 M methylviologen chloride in 0.05 M sodium phosphate buffer. MV²⁺ is a fast and reversible electron acceptor with $E^0 = -0.03 V_{RHE}$ at pH 7.⁴⁴ As can be seen from Fig. 2b, the photoelectrochemical response of the Mo-CGSe and FTO-CGSe films is nearly unchanged. Photocurrents are -1.0 mA cm⁻² and -2.4 mA cm⁻² at -0.03 V_{RHE} for the FTO-CGSe and Mo-CGSe electrodes, respectively, and the photovoltage values are $V_{ph}(ext, LSV) = E_{on} - E_{(MV^{2+}/MV^+)}^0 = 0.25 V$ and 0.38 V, respectively. Again, the FTO-CGSe electrode underperforms in comparison to the Mo-CGSe electrode.

The reason for the different performances of both electrodes can be understood with the energy diagrams in Fig. 3a and b. The FTO/CGSe interface contact produces a Schottky junction due to the workfunction difference of the materials. The resulting hole transfer barrier from CGSe to FTO causes trapping of holes at the interface and reduces the quasi-Fermi level splitting and photovoltage of the FTO/CGSe photoelectrode. The Mo/CGSe contact in Fig. 3b is better suited for majority carrier extraction, because the higher work function of Mo reduces the barrier for hole transfer from CGSe to Mo.

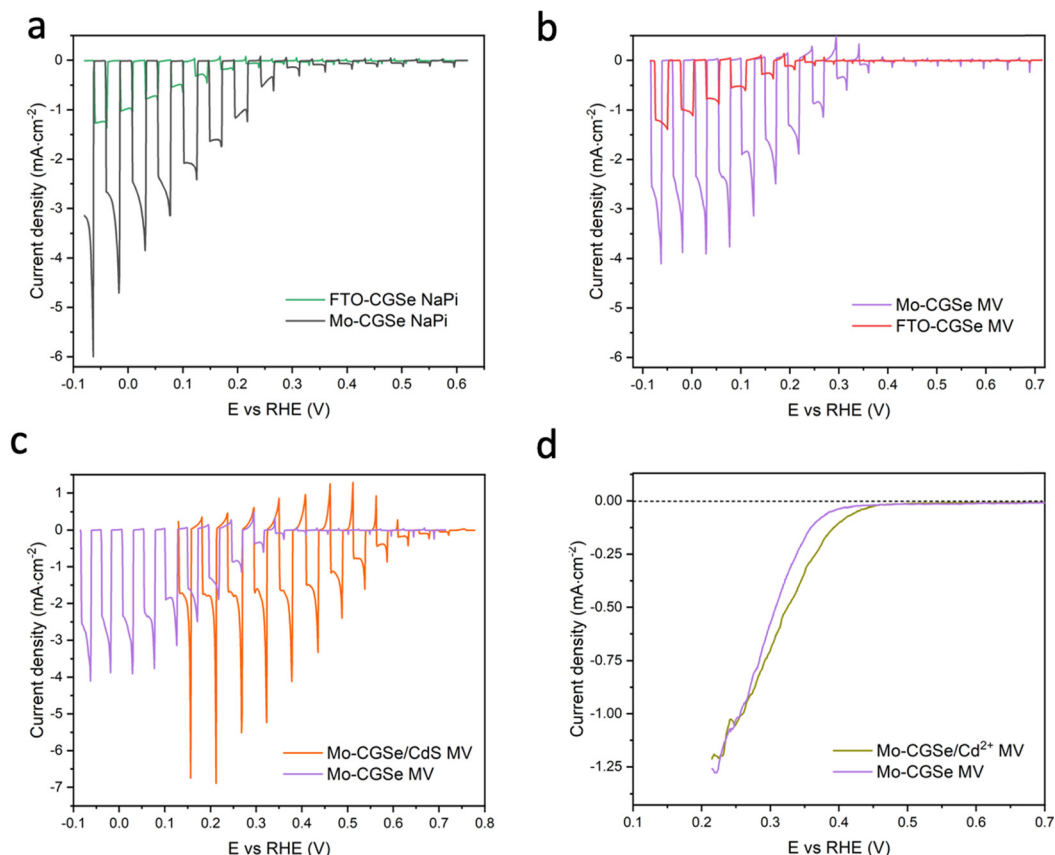


Fig. 2 Linear sweep voltammetry (LSV) scans under chopped illumination (400 nm, 81 mW cm⁻²) for (a) Mo-CGSe and FTO-CGSe in NaPi, (b) Mo-CGSe and FTO-CGSe in MV, (c) Mo-CGSe and Mo-CGSe/CdS in MV, (d) Mo-CGSe and Mo-CGSe/Cd²⁺ in MV. Continuously illuminated scans are shown here to highlight the difference in photocurrent onset potentials. Chopped scans are shown in Fig. S2 (ESI[†]). Electrolytes were de-aerated by continuous N₂ bubbling. NaPi is 0.1 M Na₂SO₄/0.05 M sodium phosphate buffer (pH = 7) and MV is 0.01 M MVCl₂/0.05 M sodium phosphate buffer (pH = 7).



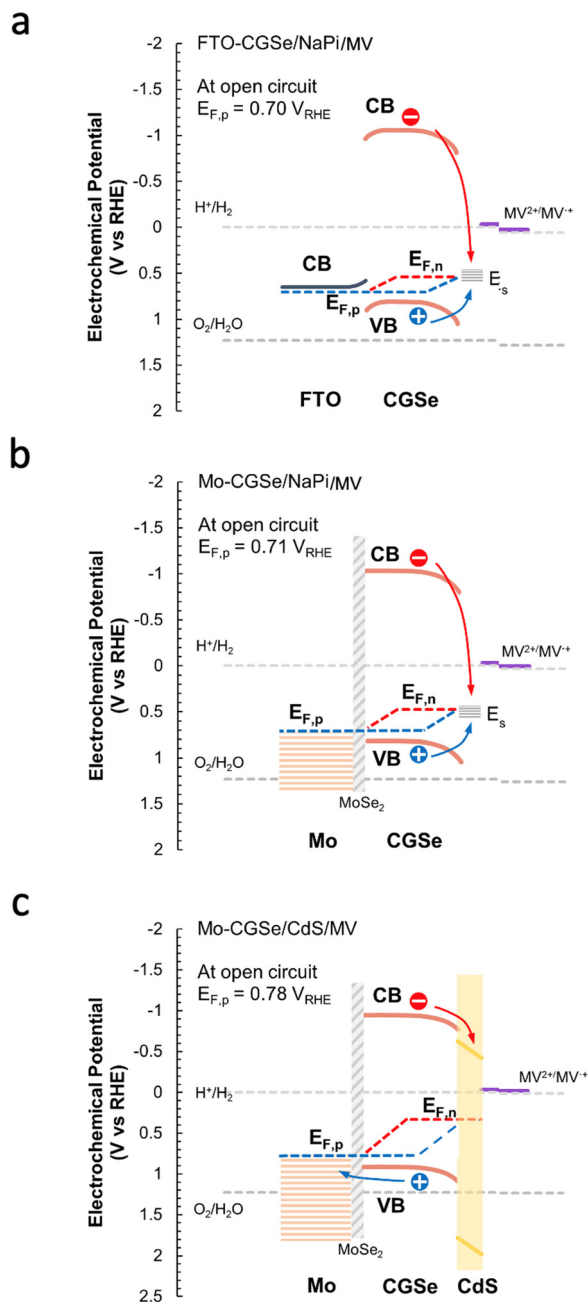


Fig. 3 Energy diagrams for (a) FTO–CGSe/NaPi/MV, (b) Mo–CGSe/NaPi/MV, (c) Mo–CGSe/CdS/MV junctions showing band bending and QFLS under illumination at open circuit. The work function of solvent-cleaned FTO is 4.48 to 4.73 eV,⁴⁵ while the work function of the Mo-substrate (110) is 4.9–5.0 eV.⁴⁶ Using the E_F values for CGSe (0.53–0.60 V_{RHE}) from Fig. S4 (ESI[†]), the Schottky barrier height of FTO/CGSe is 0.2–0.5 eV and that of Mo/CGSe is \sim 0 eV (near ohmic contact). CGSe band edges are based on Mahmoudi *et al.*¹⁴ Standard reduction potentials were taken from ref. 44 and adjusted to RHE, as described in the Experimental section.

Additionally, molybdenum is well-known to form low-defect contacts with chalcopyrites due to the formation of a passivating MoSe_2 layer at the interface.^{47–49}

To probe the QFLS of these junctions, surface photovoltage measurements were employed. Full SPV spectra for each photoelectrode in contact with an aqueous 0.01 M methylviologen

electrolyte and sodium phosphate buffer (NaPi) are shown in Fig. 4. All samples generate positive photovoltage signals, in agreement with the charge transfer scheme in Fig. 4a. Generally, the photovoltage begins at 1.5 eV, approximately 0.35 eV below the optical band gap (1.85 eV), which is attributed to the excitation of sub-bandgap defects, likely at the CGSe interfaces. A slightly higher signal at 1.3–1.5 eV for FTO–CGSe (shown by a grey square) indicates a higher defect concentration for this sample, possibly due to Ga_{Cu} and Ga_{Se} , and copper vacancies V_{Cu} .^{50,51} The SPV signal above 1.8 eV is due to band gap excitation of CGSe while the double peaks at 2.0 eV and 3.0 eV are a result of the emission characteristics of the Xe light source (Fig. S3, ESI[†]).

To obtain light intensity dependent QFLS data, SPV measurements were repeated under transient illumination (Fig. 5). A 400 nm LED light source was used to ensure band gap excitation of CuGa_3Se_5 . As before, positive SPV signals arise from the transfer of majority carriers (holes) into the electric contact at the CGSe backside and of minority carriers (electrons) into surface states. Signals are mostly reversible and increase with light intensity, as expected for the QFLS of a junction.²⁷ Baseline drifts indicate trapping of charge carriers at the interfaces. Overall, the Mo–CGSe electrodes produce a larger SPV signal than the FTO–CGSe electrodes, regardless if water (NaPi) or methylviologen is used as the electron acceptor. This confirms the improved charge carrier separation for Mo–CGSe electrodes as already seen in the former PEC measurements (Fig. 2) and as postulated by the energy schemes in Fig. 3.

Together with the light intensity dependent open circuit potentials (Fig. S4, ESI[†]), the SPV data was used to construct the QFL plots in Fig. 6. In both MV and NaPi electrolytes, $E_{\text{F,p}}$ becomes more oxidizing with increasing light intensity, as expected for a photocathode. On the other hand, $E_{\text{F,n}}$ stays constant and remains close to the energy of the surface states (likely gallium vacancies¹⁶), approximately 0.5 eV above the valence band edge. This phenomenon is known as Fermi level pinning and has been documented for many semiconductor electrodes,^{16,43,52–55} incl. chalcopyrites.⁵⁴

Additionally, it can be seen that the FTO–CGSe electrode produces $E_{\text{F,n}}$ potentials \sim 0.1 V (0.063 V for NaPi and 0.108 V for MV) less reducing than the Mo–CGSe electrode. This is a result of the Schottky junction in Fig. 3a, which promotes electron–hole recombination at the FTO–CGSe interface. Despite the better photovoltage of the Mo–CGSe electrode, $E_{\text{F,n}}$ remains too weak to reduce protons nor MV^{2+} even under the highest illumination intensity. This explains the need for an applied potential that is more negative than 0.3 V_{RHE} to produce a cathodic photocurrent (Fig. 2a and b). As mentioned above, surface modification of Cu-based chalcogenide photocathodes with a thin layer of n-type CdS reduces the photovoltage loss and gives larger cathodic photocurrents.^{5,15,19,56,57}

To test the effect of this treatment on the present system, a thin layer of CdS was deposited on the surface of the Mo–CGSe electrode by the reported chemical bath deposition (CBD) method.¹⁵ Photographs, SEM and EDX mapping data of the resulting electrode are shown in Fig. S5 (ESI[†]). The CdS layer is



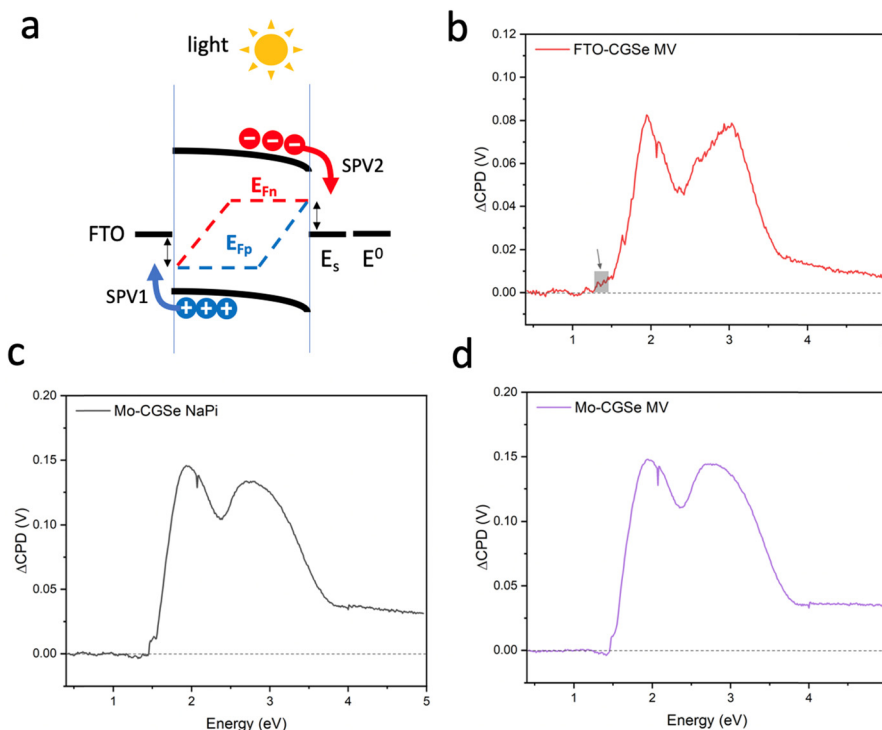


Fig. 4 SPV spectra of the CGSe-electrodes. (a) Charge transfer leading to the SPV signal for a p-type semiconductor–solid–liquid junction. Spectra for the (b) FTO–CGSe/MV, (c) Mo–CGSe/NaPi, and (d) Mo–CGSe/MV; 0.01 M MVCl₂ in 0.05 M sodium phosphate buffer (pH = 7). SPV discontinuities at 1.45 and 2.07 eV are due to monochromator filter-changes. SPV onsets are found at 1.5 eV, except for (b) which has a slightly lower onset at 1.3 eV, indicating a sub-band gap defect.

uniform and retains the sub-micrometer morphology of the underlying Mo–CGSe substrate.

Indeed, the PEC performance of the resulting Mo–CGSe/CdS for MV²⁺ reduction is enhanced significantly (Fig. 2c). The photo-onset for MV²⁺ reduction shifts from 0.35 V_{RHE} to 0.65 V_{RHE}, increasing the photovoltage to 0.68 V (from 0.38 V for the bare Mo–CGSe photoelectrode). Also, the photocurrent more than doubles at 0.35 V_{RHE}, but the saturated photocurrent does not exceed -2.5 mA cm^{-2} . This is attributed to the low MVCl₂ concentration (0.01 M) in the electrolyte and the slow mass transport resulting from it. Additionally, large ($> 4 \text{ mA cm}^{-2}$) transient cathodic and smaller anodic ($> 1 \text{ mA cm}^{-2}$) current spikes indicate a strong capacitive current component that is attributed to electron trapping/de-trapping in the CdS layer of the p–n–CGSe/CdS junction.^{58–60}

The sequence of SPV and OCP measurements was repeated for the Mo–CGSe/CdS–MV junction (Fig. 5 and Fig. S4, ESI†) to produce the QFL data in Fig. 6e. It can be seen that the CdS layer raises the $E_{F,n}$ potentials above the energy E_s of the surface states over the entire illumination range and even in the dark (E_F). This proves that the CdS layer removes or modifies the Fermi level pinning states at the CGSe electrode surface. Additionally, $E_{F,p}$ levels are found to be 0.02–0.10 V more oxidizing than before, especially at the lower illumination intensity, and that on average the QFLS is increased by 40%. This illustrates the better charge separation at the CGSe–CdS interface, resulting from the introduction of the p–n–junction,

as shown in Fig. 3c. While the junction promotes electron extraction *via* an electric field, the deep CdS valence band, relative to CGSe, acts as a hole mirror. Overall, the QFL data agree well with the observed increase in the photoelectrochemical MV reduction behavior of the CdS-modified Mo–CGSe electrode.

Lastly, we use the SPV/OCP approach to differentiate the effect of the p–n–junction from the effect of surface passivation. It is well-known that photovoltage losses at CGSe photocathodes can be reduced by the application of cadmium salts.^{61–63} For example, surface pretreatment with cadmium ions before the deposition of buffer layers was reported to improve the performance of Cu(In,Ga)Se₂, CuIn_xSe_{2–x} or CuGa₃Se₅-based solar cells.^{61,63} This was mainly attributed to the incorporation of Cd²⁺ at the surface and the removal of the surface oxide. To investigate this possibility, the Mo–CGSe electrode was soaked in a 0.1 M aqueous cadmium acetate solution at 65 °C for 30 min, followed by washing with water. Samples treated this way contain small amounts of Cd²⁺ on their surface, as confirmed by X-ray photoelectron spectroscopy (XPS) (Fig. S6, ESI†). This compares to observations in the literature.⁴⁴

The LSV of the resulting electrode in MV electrolyte is shown in Fig. 2d. Compared to the non-treated sample, the LSV curve shows a 0.05 mV anodic shift of the photocurrent onset and a small boost in photocurrent, which disappears at more reducing potentials. This shows that the surface



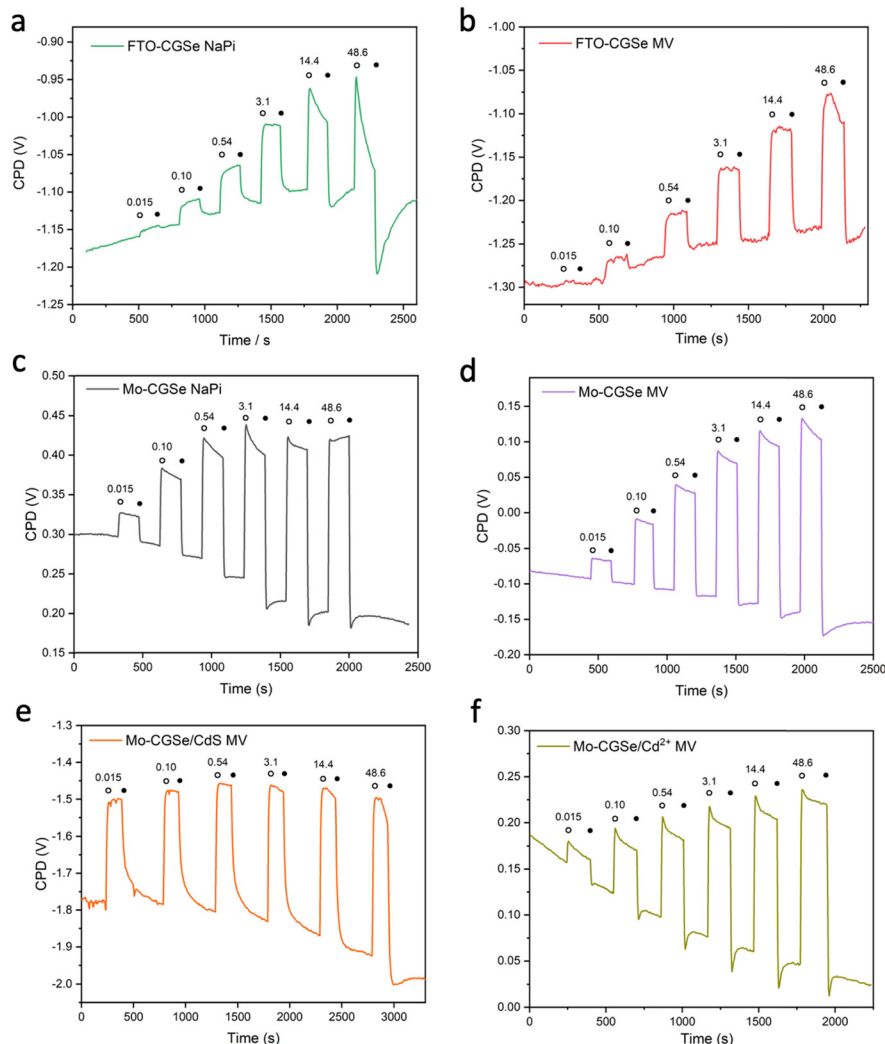


Fig. 5 Surface photovoltage (SPV) data as a function of light intensity (mW cm^{-2}) from a 400 nm LED. (a) FTO-CGSe/NaPi, (b) FTO-CGSe/MV, (c) Mo-CGSe/NaPi, (d) Mo-CGSe/MV, (e) Mo-CGSe/CdS/MV, and (f) Mo-CGSe/Cd²⁺/MV.

treatment with Cd²⁺ improves the photocathode performance to a small extent, but not nearly as much as a CdS overlayer. Using the data in Fig. 5f and Fig. S4f (ESI[†]), the QFL plot for the Cd-ion-modified Mo-CGSe electrode was constructed in Fig. 6d. No significant changes occur with the quasi Fermi levels and, importantly, no increase in the QFLS (the internal photovoltage) is seen. This suggests that Cd²⁺ adsorption alone does not suppress charge carrier recombination at the CGSe surface and does not remove the Fermi level pinning effect. The observed small improvement in the PEC behavior in Fig. 2d must therefore be due to other factors, *e.g.* a change of the surface charge or of the band edge potentials. However, to improve the internal photovoltage of the electrode, a full CdS layer and a p-/n-junction are needed.

Conclusions

For the first time, surface photovoltage (SPV) measurements were used to measure the quasi-Fermi level splitting (QFLS) in

CuGa₃Se₅ thin film photocathodes in contact with different electrolytes, back contacts, and before and after surface treatments. For the non-treated electrodes, Fermi level pinning near the surface states, 0.5 eV above the CGSe valence band, is confirmed as the cause for the voltage loss during photoelectrochemical water and methyl viologen reduction. The surface states both limit the band bending at the CGSe/liquid interface and they serve as electron-hole recombination centers. Application of a CdS surface layer improves performance, as established by prior works, and increases the internal photovoltage to 0.44 V (MV) under 49 mW cm^{-2} illumination at open circuit. This increase is due to the electron selective properties of the CGSe/CdS interface (the CdS valence band maximum is too oxidizing to accept holes from CGSe) and the formation of a p-/n-junction. Molybdenum is confirmed to be a better back contact for CGSe than FTO, based on the increased internal photovoltage (from 0.16 V to 0.22 V in NaPi electrolyte under 49 mW cm^{-2} illumination). This confirms the hypothesis that the photovoltage of FTO/CGSe is limited by the poor hole selectivity of the Schottky junction at the interface. Overall, these



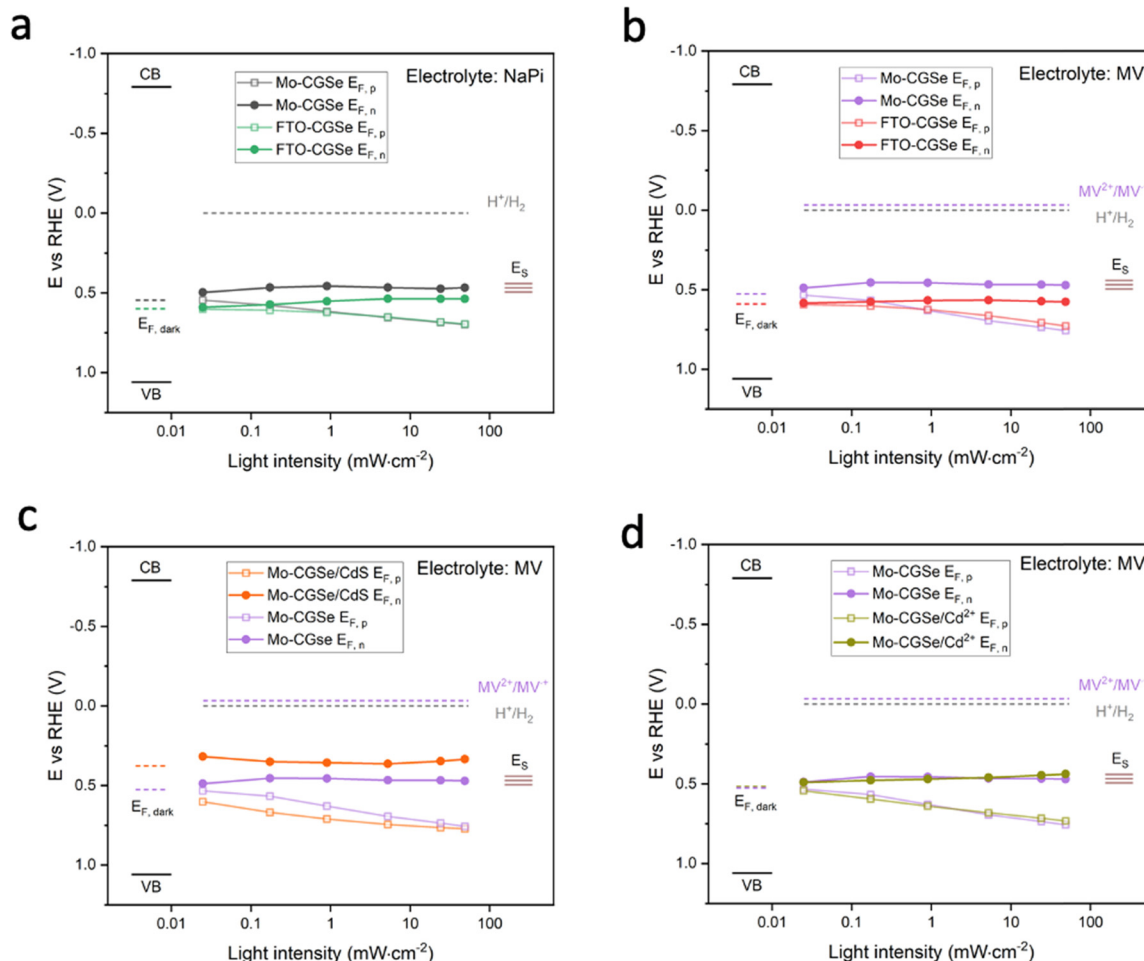


Fig. 6 Quasi-Fermi level (QFL) plots versus the logarithm of light intensity (400 nm LED). (a) Mo-CGSe/NaPi versus FTO-CGSe/NaPi. (b) Mo-CGSe/MV versus FTO-CGSe/MV. (c) Mo-CGSe/CdS/MV versus Mo-CGSe/MV. (d) Mo-CGSe/Cd²⁺/MV versus Mo-CGSe/MV.

findings provide a better understanding of the photoelectrochemistry of CGSe photoelectrodes. Additionally, this work confirms SPV as a powerful method to measure internal photovoltage in semiconductor/liquid junctions and to identify efficient charge-selective contacts.

Experimental section

Chemicals

Sodium phosphate monobasic monohydrate ($\geq 99.0\%$, Sigma-Aldrich), sodium phosphate dibasic heptahydrate (98.0%, Sigma-Aldrich), sodium sulfate ($\geq 99.0\%$, Sigma-Aldrich) and methyl viologen dichloride hydrate (98%, Acros Organics) were used for the preparation of the electrolytes. Cadmium acetate dihydrate (analytical reagent, Mallinckrodt), thiourea (99%, Alfa Aesar) and ammonium hydroxide (29.7%, certified ACS plus, Fisher Chemicals) were used for the cadmium sulfide chemical bath deposition. Potassium ferricyanide (99.2%, Sigma) and potassium hexacyanoferrate(n) trihydrate (98.0–102.0%, Sigma-Aldrich) were used as received. Water was purified to a resistivity of 18 M Ω cm by a Nanopure system.

CuGa₃Se₅ (CGSe) films were prepared by single-stage co-evaporation of high-purity Cu, Ga and Se onto 2×2 cm² pieces of Mo-coated soda-lime glass containing Na (Guardian Industries Corp.) and FTO-coated soda-lime glass (Sigma-Aldrich). High-vacuum physical vapor deposition was performed simultaneously on FTO and Mo while rotating substrates at 10 rpm to have a homogeneous deposition. The nominal temperature of the substrates was kept at 625 °C. The atomic ratio of (Cu/Ga) was close to 0.33 ($0.33 < [\text{Cu}/\text{Ga}] < 0.37$) for all the samples which were analyzed later by energy dispersive X-ray spectroscopy (EDX) using an EDAX Genesis installed on a JEOL SEM-4701F with a cold field emitter using an accelerating voltage of 15 kV.

Cadmium acetate treatment

Inspired by previous works,^{62,63} the film was soaked in an aqueous 0.1 M cadmium acetate solution at 65 °C for 30 min. The film was then rinsed with water to remove excess reagent.

Cds layer deposition

CdS layers were formed on the surface of CGSe films by the chemical bath deposition (CBD) method.¹⁵ Prior to CdS deposition, the surface of the CGSe films was pretreated with an



aqueous solution containing 2 M NH_4OH and 7.5 mM $\text{Cd}(\text{CH}_3\text{COO})_2$ at 80 °C for 10 min. CBD of CdS was performed by immersing the Cd^{2+} -pretreated films in a bath solution containing 0.375 M $\text{SC}(\text{NH}_2)_2$, 7.5 mM $\text{Cd}(\text{CH}_3\text{COO})_2$ and 2 M NH_4OH at 65 °C for 5 min. After CdS deposition, the samples were annealed in air at 300 °C for 60 min.

Scanning electron microscopy (SEM) and energy dispersive X-ray spectroscopy (EDX) were recorded using a Scios Dual-Beam FIB/SEM.

X-ray photoelectron spectroscopy (XPS) was conducted using a Supra XPS spectrometer with an Al $K\alpha$ source that generates X-rays at 1487 eV. All analysis was done in the ultra-high vacuum (UHV) analytical chamber with a pressure of 10^{-7} mbar. The survey scan was conducted at a constant pass energy of 40 eV with a scan step of 0.5 eV, and the high-resolution core-level spectra were recorded at a constant pass energy of 40 eV with a scan step of 0.1 eV. The spectra were electrostatically corrected based on the position of C 1s (284.8 eV).

Photoelectrochemical (PEC) linear sweep voltammetry (LSV) measurements

Photoelectrochemical measurements were conducted using a Gamry Reference 600 Potentiostat connected to a typical three-electrode system, with a Pt counter electrode and a calomel reference electrode (3.5 M KCl). All the PEC measurements were performed with stirring and continuous N_2 purging. LSV scans were performed with a 10 mV s^{-1} scan rate in the cathodic direction. The potentials in each measurement were converted to the normal hydrogen electrode (NHE) by calibration in 10 mM $\text{K}_3\text{Fe}(\text{CN})_6/\text{K}_4\text{Fe}(\text{CN})_6$. Potentials were adjusted to RHE with the following formula: $V_{\text{RHE}} = V_{\text{NHE}} + 0.0592 \times \text{pH}$. The same 400 nm LED described for the SPV measurements below was used as the light source.

The electrolytes used were as follows: NaPi: 0.1 M Na_2SO_4 with 0.05 M sodium phosphate buffer (pH = 7); MV: 0.01 MVCl_2 with 0.05 M sodium phosphate buffer (pH = 7). The pH of the electrolytes was measured with a pH meter (Fisher Scientific accumet AE150).

Open circuit potential (OCP) measurements

A two-electrode setup was used for the OCP measurements by using a calomel electrode (3.5 M KCl) as the counter/reference electrode and the CGSe electrode as the working electrode. All OCP measurements were performed with stirring and continuous N_2 purging. The potentials in each measurement were converted into the values *versus* RHE as described above. The same 400 nm LED light source was used.

Surface photovoltage (SPV) measurements

Contact potential differences (CPD) were measured using a circular (2.5 mm diameter), semitransparent vibrating gold mesh disk (Kelvin Probe S, Besocke Delta Phi) that was mounted inside of a home-built vacuum chamber and controlled by a Kelvin Control 7 Oscillator/amplifier (Besocke Delta Phi). Samples were placed approximately 1 mm underneath the Kelvin probe. A 150 W Xe lamp passing through an Oriel

Cornerstone 130 monochromator was used as the light source for the full SPV spectra. For transient SPV data, a 400 nm LED (LZC-00UA00, LedEngin) was used as the light source, and the irradiation power was controlled by a power supply (Naweisz NP6005). Light intensity was measured with a photometer (International Light IL1400BL) equipped with a SEL005 detector. SPV data was recorded every 5 seconds by measuring the contact potential difference (CPD) value at each step. For liquid SPV measurements, 15 μL of the respective liquid electrolyte was dropped on the surface of the sample and then covered with a microscope cover glass (Fisher Scientific, thickness 0.17–0.25 mm). During the measurements, the chamber was continuously purged with a water-saturated N_2 flow (flow rate of 0.02–0.10 SLPM, standard liter per minute) to suppress the evaporation of the electrolyte.

Author contributions

The experiments were designed by F. E. O. and conducted by Y. C. and C. X. CuGa_3Se_5 films were prepared by B. M. and A. W. M. with assistance by R. S. The manuscript was written by F. E. O. and Y. C. with additional contributions by all other authors. All authors have given approval to the final version of the manuscript.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 S. Niki, M. Contreras, I. Repins, M. Powalla, K. Kushiya, S. Ishizuka and K. Matsubara, *Prog. Photovoltaics Res. Appl.*, 2010, **18**, 453–466.
- 2 M. Nakamura, K. Yamaguchi, Y. Kimoto, Y. Yasaki, T. Kato and H. Sugimoto, *IEEE J. Photovolt.*, 2019, **9**, 1863–1867.
- 3 M. Bär, W. Bohne, J. Röhrich, E. Strub, S. Lindner, M. C. Lux-Steiner, C. H. Fischer, T. P. Niesen and F. Karg, *J. Appl. Phys.*, 2004, **96**, 3857–3860.
- 4 B. Marsen, B. Cole and E. L. Miller, *Sol. Energy Mater. Sol. Cells*, 2008, **92**, 1054–1058.
- 5 M. Moriya, T. Minegishi, H. Kumagai, M. Katayama, J. Kubota and K. Domen, *J. Am. Chem. Soc.*, 2013, **135**, 3733–3735.
- 6 N. Gaillard, D. Prasher, J. Kaneshiro, S. Mallory and M. Chong, Materials Research Society Symposium Proceedings, 2013, vol. 1558, mrss13-1558-z1502-1507.



- 7 T. J. Jacobsson, V. Fjällström, M. Sahlberg, M. Edoff and T. Edvinsson, *Energy Environ. Sci.*, 2013, **6**, 3676–3683.
- 8 Y. Chen, X. Feng, M. Liu, J. Su and S. Shen, *Nanophotonics*, 2016, **5**, 524–547.
- 9 N. Gaillard, D. Prasher, M. Chong, A. Deangelis, K. Horsley, H. A. Ishii, J. P. Bradley, J. Varley and T. Ogitsu, *ACS Appl. Energy Mater.*, 2019, **2**, 5515–5524.
- 10 H. Kobayashi, N. Sato, M. Orita, Y. Kuang, H. Kaneko, T. Minegishi, T. Yamada and K. Domen, *Energy Environ. Sci.*, 2018, **11**, 3003–3009.
- 11 H. Kumagai, T. Minegishi, N. Sato, T. Yamada, J. Kubota and K. Domen, *J. Mater. Chem. A*, 2015, **3**, 8300–8307.
- 12 J. Kim, T. Minegishi, J. Kobota and K. Domen, *Jpn. J. Appl. Phys.*, 2011, **51**, 015802.
- 13 C. P. Muzzillo, W. E. Klein, Z. Li, A. D. DeAngelis, K. Horsley, K. Zhu and N. Gaillard, *ACS Appl. Mater. Interfaces*, 2018, **10**, 19573–19579.
- 14 B. Mahmoudi, F. Caddeo, T. Lindenberg, T. Schneider, T. Hölscher, R. Scheer and A. W. Maijenburg, *Electrochim. Acta*, 2021, **367**, 137183.
- 15 L. Zhang, T. Minegishi, M. Nakabayashi, Y. Suzuki, K. Seki, N. Shibata, J. Kubota and K. Domen, *Chem. Sci.*, 2015, **6**, 894–901.
- 16 Y. Liu, M. Bouri, L. Yao, M. Xia, M. Mensi, M. Grätzel, K. Sivula, U. Aschauer and N. Guijarro, *Angew. Chem., Int. Ed.*, 2021, **60**, 23651–23655.
- 17 I. S. Khan, C. P. Muzzillo, C. L. Perkins, A. G. Norman, J. L. Young, N. Gaillard and A. Zakutayev, *J. Phys. Energy*, 2021, **3**, 024001.
- 18 S. Ishizuka, R. Okamoto and S. Ikeda, *Adv. Mater. Interfaces*, 2022, **9**, 2201266.
- 19 D. Yokoyama, T. Minegishi, K. Maeda, M. Katayama, J. Kubota, A. Yamada, M. Konagai and K. Domen, *Electrochem. Commun.*, 2010, **12**, 851–853.
- 20 S. L. Jang, *Solid-State Electron.*, 1991, **34**, 373–377.
- 21 T. Dittrich and S. Fengler, *Surface Photovoltage Analysis of Photoactive Materials*, World Scientific Publishing Europe Ltd., London, 2020.
- 22 R. Chen, F. Fan, T. Dittrich and C. Li, *Chem. Soc. Rev.*, 2018, **47**, 8238–8262.
- 23 D. K. Schroder, *Meas. Sci. Technol.*, 2001, **12**, R16–R31.
- 24 L. Kronik and Y. Shapira, *Surf. Sci. Rep.*, 1999, **37**, 1–206.
- 25 W. Shockley and H. J. Queisser, *J. Appl. Phys.*, 1961, **32**, 510–519.
- 26 H. Gerischer, *J. Electroanal. Chem.*, 1977, **82**, 133–143.
- 27 P. Würfel, *Physics of Solar Cells*, Wiley-VCH, Weinheim, 2005.
- 28 R. Krol, *Photoelectrochem. Hydrogen Production*, 2012, vol. 102, pp. 13–67.
- 29 R. Memming, *Semiconductor Electrochemistry*, 2015, pp. 169–266, DOI: [10.1002/9783527688685.ch7](https://doi.org/10.1002/9783527688685.ch7).
- 30 B. Miao, A. Iqbal and K. H. Bevan, *J. Phys. Chem. C*, 2019, **123**, 28593–28603.
- 31 M. X. Tan, C. N. Kenyon and N. S. Lewis, *J. Phys. Chem.*, 1994, **98**, 4959–4962.
- 32 O. Kruger, C. N. Kenyon, M. X. Tan and N. S. Lewis, *J. Phys. Chem. B*, 1997, **101**, 2840–2849.
- 33 K. Lobato, L. M. Peter and U. Würfel, *J. Phys. Chem. B*, 2006, **110**, 16201–16204.
- 34 S. Zhang, X. Yang, K. Zhang, H. Chen, M. Yanagida and L. Han, *Phys. Chem. Chem. Phys.*, 2011, **13**, 19310–19313.
- 35 F. D. Lin and S. W. Boettcher, *Nat. Mater.*, 2014, **13**, 81–86.
- 36 M. R. Nellist, F. A. L. Laskowski, F. Lin, T. J. Mills and S. W. Boettcher, *Acc. Chem. Res.*, 2016, **49**, 733–740.
- 37 M. R. Nellist, F. A. L. Laskowski, J. Qiu, H. Hajibabaei, K. Sivula, T. W. Hamann and S. W. Boettcher, *Nat. Energy*, 2018, **3**, 46–52.
- 38 F. A. L. Laskowski, S. Z. Oener, M. R. Nellist, A. M. Gordon, D. C. Bain, J. L. Fehrs and S. W. Boettcher, *Nat. Mater.*, 2020, **19**, 69–76.
- 39 H. Gerischer, *Electrochim. Acta*, 1990, **35**, 1677–1699.
- 40 N. S. Lewis, *Inorg. Chem.*, 2005, **44**, 6900–6911.
- 41 R. Han, M. A. Melo Jr, Z. Zhao, Z. Wu and F. E. Osterloh, *J. Phys. Chem. C*, 2020, **124**, 9724–9733.
- 42 J. Wang and F. E. Osterloh, *J. Mater. Chem. A*, 2014, **2**, 9405–9411.
- 43 A. J. Bard, A. B. Bocarsly, F. R. F. Fan, E. G. Walton and M. S. Wrighton, *J. Am. Chem. Soc.*, 1980, **102**, 3671–3677.
- 44 J. R. White and A. J. Bard, *J. Electroanal. Chem. Interfacial Electrochem.*, 1986, **197**, 233–244.
- 45 M. G. Helander, M. T. Greiner, Z. B. Wang, W. M. Tang and Z. H. Lu, *J. Vac. Sci. Technol., A*, 2011, **29**, 011019.
- 46 D. R. Lide, *CRC Handbook of Chemistry and Physics*, CRC Press, Taylor and Francis, Boca Raton, FL, 88 (Internet Version 2008), 2008.
- 47 N. Kohara, S. Nishiwaki, Y. Hashimoto, T. Negami and T. Wada, *Sol. Energy Mater. Sol. Cells*, 2001, **67**, 209–215.
- 48 D. Abou-Ras, G. Kostorz, D. Bremaud, M. Kalin, F. V. Kurdesau, A. N. Tiwari and M. Dobeli, *Thin Solid Films*, 2005, **480**, 433–438.
- 49 J. H. Yoon, J. H. Kim, W. M. Kim, J. K. Park, Y. J. Baik, T. Y. Seong and J. H. Jeong, *Prog. Photovoltaics*, 2014, **22**, 90–96.
- 50 D. Habib and G. E. H. Moussa, *World J. Condens. Matter Phys.*, 2016, **6**, 27–34.
- 51 B. Mahmoudi, F. Caddeo, T. Lindenberg, T. Schneider, T. Hölscher, R. Scheer and A. W. Maijenburg, *Electrochim. Acta*, 2020, 137183, DOI: [10.1016/j.electacta.2020.137183](https://doi.org/10.1016/j.electacta.2020.137183).
- 52 C. L. Hinkle, M. Milojevic, B. Brennan, A. M. Sonnet, F. S. Aguirre-Tostado, G. J. Hughes, E. M. Vogel and R. M. Wallace, *Appl. Phys. Lett.*, 2009, **94**, 162101.
- 53 B. Klahr, S. Gimenez, F. Fabregat-Santiago, T. Hamann and J. Bisquert, *J. Am. Chem. Soc.*, 2012, **134**, 4294–4302.
- 54 Y. Liu, F. Le Formal, F. Boudoire, L. Yao, K. Sivula and N. Guijarro, *J. Mater. Chem. A*, 2019, **7**, 1669–1677.
- 55 M. F. Lichterman, S. Hu, M. H. Richter, E. J. Crumlin, S. Axnanda, M. Favaro, W. Drisdell, Z. Hussain, T. Mayer, B. S. Brunschwig, N. S. Lewis, Z. Liu and H.-J. Lewerenz, *Energy Environ. Sci.*, 2015, **8**, 2409–2416.
- 56 T. Hayashi, R. Niishiro, H. Ishihara, M. Yamaguchi, Q. X. Jia, Y. B. Kuang, T. Higashi, A. Iwase, T. Minegishi, T. Yamada, K. Domen and A. Kudo, *Sustainable Energy Fuels*, 2018, **2**, 2016–2024.



- 57 H. Kaneko, T. Minegishi and K. Domen, *Chem. – Eur. J.*, 2018, **24**, 5697–5706.
- 58 R. Krol, in *Photoelectrochemical Hydrogen Production*, ed. R. van de Krol and M. Grätzel, Springer, US, 2012, ch. 3, vol. 102, pp. 69–117.
- 59 L. M. Peter, J. Li and R. Peat, *J. Electroanal. Chem. Interfacial Electrochem.*, 1984, **165**, 29–40.
- 60 P. Salvador and C. Gutiérrez, *J. Electroanal. Chem. Interfacial Electrochem.*, 1984, **160**, 117–130.
- 61 M. A. Contreras, M. J. Romero, B. To, F. Hasoon, R. Noufi, S. Ward and K. Ramanathan, *Thin Solid Films*, 2002, **403–404**, 204–211.
- 62 B. Lei, W. W. Hou, S. H. Li, W. B. Yang, C. H. Chung and Y. Yang, *Sol. Energy Mater. Sol. Cells*, 2011, **95**, 2384–2389.
- 63 I. S. Khan, C. P. Muzzillo, C. L. Perkins, A. G. Norman, J. L. Young, N. Gaillard and A. Zakutayev, *J. Phys. Energy*, 2021, **3**, 024001.

