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Na_{2.4}Al_{0.4}Mn_{2.6}O₇ anionic redox cathode material for sodium-ion batteries – a combined experimental and theoretical approach to elucidate its charge storage mechanism†

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Here we report the synthesis *via* ceramic methods of the high-performance Mn-rich Na_{2.4}Al_{0.4}Mn_{2.6}O₇ oxygen-redox cathode material for Na-ion batteries, which we use as a testbed material to study the effects of Al substitution and subsequent Na excess in the high-capacity, anionic redox-based cathode material, Na₂Mn₃O₇. The material shows a stable electrochemical performance, with a specific capacity of 215 mA h g⁻¹ in the 1.5–4.7 V voltage range at C/20 and a capacity retention of 90% after 40 cycles. Using a combination of electrochemical and structural analysis together with hybrid density functional theory calculations we explain the behaviour of this material with changes in Mn/anionic redox reactions and associated O₂ release reactions occurring during electrochemical cycling (Na⁺ ion insertion/extraction), and compare these findings to Na₂Mn₃O₇. We expect that these results will advance understanding of the effect of dopants in Mn-rich cathode materials with oxygen redox activity to pave their way towards high-performance sodium-ion batteries.

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Introduction

Sodium-ion batteries (SIBs) are a promising candidate for next-generation batteries, as sodium is an abundant element with an easy-access source, which translates into a low-cost and sustainable alternative energy storage technology. Moreover, the chemical similarities of Li and Na coupled with the already well-established knowledge in lithium-ion battery (LIB) technology allow for rapid implementation of SIBs in applications such as the grid or low-cost transportation.^{1,2} Among the cathode materials researched in the field of SIBs, layered Na–Mn–O ternary compounds have been particularly relevant due to their high capacity, abundance, non-toxicity and low cost.

The Na_xMnO₂ (*x* > 0.5) composition has been the most studied, showing a maximum theoretical capacity of 244 mA h g⁻¹ based on the Mn^{3+/4+} redox-active couple. Nonetheless, these materials generally present low capacity retention, *e.g.*, β-NaMnO₂ shows 53% capacity retention at C/20 after 100 cycles, with an initial capacity of 175 mA h g⁻¹ (when cycled between 4.2 V and 2.0 V);³ and α-NaMnO₂ shows 70% capacity retention at C/30 after 20 cycles, with an initial capacity of 185 mA h g⁻¹ (when cycled between 3.8 V and 2.0 V).⁴ The poor cycling stability of Mn-based compounds is generally attributed to structural instabilities associated with the Jahn–Teller (JT) effect of the Mn³⁺ ions formed during cycling, particularly under a low working voltage range (1.5–2.1 V).⁵

Na₂Mn₃O₇ (or Na_{0.57}Mn_{0.86}[]_{0.14}O₂) (*P* $\bar{1}$ space group) has been extensively studied in the past years due to the large anionic redox contribution (75 mA h g⁻¹ above 4.0 V vs. Na⁺/Na) to the overall charge stored in the material (theoretical capacity = 155 mA h g⁻¹), which results in high energy densities.^{6,7} This large anionic redox capacity arises from non-bonding 2p orbitals of the oxygen atoms neighbouring the unoccupied Mn site in the transition metal (TM) layers.^{8–11} Despite its large anionic redox capacity, Na₂Mn₃O₇ does not fully use its redox-active oxygen ions. Earlier studies on Na₂Mn₃O₇ revealed that three non-bonding oxygen atoms per formula unit are redox-active, where theoretically, a maximum of only two of them can be oxidized/reduced during cycling.^{6,8} This shows that, once

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loading of 2.4–3.0 mg cm⁻². Sodium metal disks (15 mm in diameter) (Alfa Aesar Merck) were used as reference electrodes, and 1 M NaPF₆ (99%, Alfa Aesar) in the anhydrous organic solution of ethylene carbonate/diethylene carbonate (EC : DEC 1 : 1 w/w%, Gotion) as electrolyte. The electrolyte was dried under molecular sieves for a week before use.

Galvanostatic measurements were carried out on a battery tester (Neware) in the voltage range of 1.5–4.7 V vs. Na⁺/Na at selected current rates (C/20 to 5/C, 1C = 160 and 155 mA g⁻¹ for Na_{2.4}Al_{0.4}Mn_{2.6}O₇ and Na₂Mn₃O₇, respectively).

Cyclic voltammetry (CV) experiments were conducted using an IviumStat potentiostat (Alvatek), at a scan rate of 0.1 mV s⁻¹ in the voltage range of 1.5–4.7 V vs. Na⁺/Na.

Galvanostatic intermittent titration technique (GITT) measurements during charge/discharge were carried out by applying a pulse of ±1.7 μA for 1 h followed by an open circuit relaxation of 2 h. The procedure was carried out in the voltage range of 1.5–4.7 V vs. Na⁺/Na. The diffusion coefficient for Na_{2.4}Al_{0.4}Mn_{2.6}O₇, *D*, was calculated using eqn (S1),† where we assumed a constant molar mass (320.81 g mol⁻¹) and cell volume (169.81 cm³ mol⁻¹) during the cycling process. The surface area was calculated using two different approaches: (1) using the average particle size value obtained from SEM images (0.77 μm diameter particles), whereby we assumed that particles were spherical and had a smooth surface and an intimate contact with the electrolyte; and (2) considering a flat surface contact area (electrode area).

Electrochemical impedance spectroscopy (EIS) studies were carried out on Na_{2.4}Al_{0.4}Mn_{2.6}O₇ and Na₂Mn₃O₇ with two different electrochemical cell set-ups. EIS data were collected in the 100 kHz to 10 mHz frequency range with an amplitude of 10 mV and after a 1 h voltage pulse using a Biologic VMP-300 potentiostat and an Ivium Octostat potentiostat. EIS data on the first cycle were collected using a three-electrode cell (EL-cell), where the active electrode materials were cycled using metallic Na disks as counter and reference electrodes. The reference is a pre-assembled and commercial Na ring (EL-cell) enclosed in the cell around the glass fibre separator. On the other hand, EIS data on the 100th cycle were collected using two-electrode coin cells.

All the electrochemical testing was conducted at room temperature.

Nuclear magnetic resonance (NMR)

²³Na magic-angle spinning (MAS) NMR spectra of pristine Na_{2.4}Mn_{2.6}Al_{0.4}O₇ and Na₂Mn₃O₇ materials were obtained at 9.4 T on a Bruker Avance III 400 MHz spectrometer working at a Larmor frequency of 105.8 MHz. Powdered samples were packed into 2.5 mm rotors, in an argon-filled MBraun glovebox ([H₂O] and [O₂] ≤ 0.1 ppm) and kept static or spun up to 30 kHz. ²³Na MAS NMR spectra were acquired using a single pulse or Hahn echo pulse sequence, with a recycle delay of 0.5 s. Spectra referencing was done using a secondary solid reference of NaCl (δ_{iso} = 7.2 ppm). Due to the nature of the samples and the experimental parameters used, the relative intensities may not match contributions in the materials and therefore spectra are not quantitative. Additionally, the distorted baseline of the

single pulse experiments can affect the observed position of peaks compared to the flatter baseline of the Hahn echo experiments. *Ex situ* ²³Na MAS NMR spectra were obtained at 9.4 T, using a Hahn echo pulse sequence and 30 kHz MAS. 1 kHz line broadening was used in spectra processing. *Ex situ* samples were run in powder form in coin cells (active material : carbon black weight ratio of 80 : 20) using around 8 mg of active material per cell. Na metal was used as a counter and reference electrode and 1 M NaPF₆ in EC : DEC (1 : 1 vol%) as the electrolyte. After cycling the cells to a defined voltage, the cells were opened, and powder was recovered. Powders were washed three times with DEC inside an argon-filled glovebox (MBraun), followed by the removal of excess solvent and drying under vacuum at room temperature in the glovebox antechamber before the NMR measurement.

Ex situ synchrotron X-ray absorption near-edge spectroscopy (XANES)

X-ray absorption near edge structure (XANES) data were collected at the B18 beamline at Diamond Light Source (Harwell, UK). For the measurement, cast electrodes were cycled at different states of charge, and then, extracted from the coin cells and rinsed three times with DEC in an argon-filled glovebox (MBraun) ([H₂O] and [O₂] ≤ 0.1 ppm). After this, electrodes were dried at 80 °C under vacuum in the glovebox antechamber and sealed in individual aluminium laminated pouches under vacuum to minimise air exposure during data acquisition. Measurements were performed at the Mn K-edge energies above and below the absorption edges (*ca.* 6540 eV) and electrode data were collected in fluorescence mode, using a 36-element Ge detector. Mn metal foil was used as an internal standard and reference data for MnO, Mn₃O₂ and MnO₂ powders were collected on transmission mode using gas-filled ionisation chambers to measure both incident and transmitted intensities. Three repetitions of each sample were measured, totalling a 15 min measurement. Athena software in the Demeter package was used to perform energy calibration, background subtraction and normalization.^{28,29} Linear combination fitting (LCF) was performed using the Athena software within an energy range of -20 eV below to +30 eV above the edge. Each sample spectra were fitted for every combination of the standard spectra. The combination of standards that resulted in the best fit (smallest *R*-factor and reduced χ²) was chosen as the most likely representation of the sample.

Ab initio calculations

All the calculations performed in this work employed DFT as implemented in the Vienna *Ab initio* Simulation Package code.^{30,31} Interactions between core and valence electrons were described using the projector augmented wave (PAW) method.³² Na (2p⁶3s¹), Mn (3p⁶3d⁵4s²), Al (3s²3p¹), and O (2s²2p⁴) electron configurations were treated as the valence electrons. Convergence to plane wave energy was checked, with a cut-off of 500 eV found to be sufficient to converge the total energy to within 0.01 eV atom⁻¹. Brillouin zones for all compounds were sampled such that the *k*-points were converged with an accuracy of the total energy in 0.001 eV per atom, which corresponds to *k*-



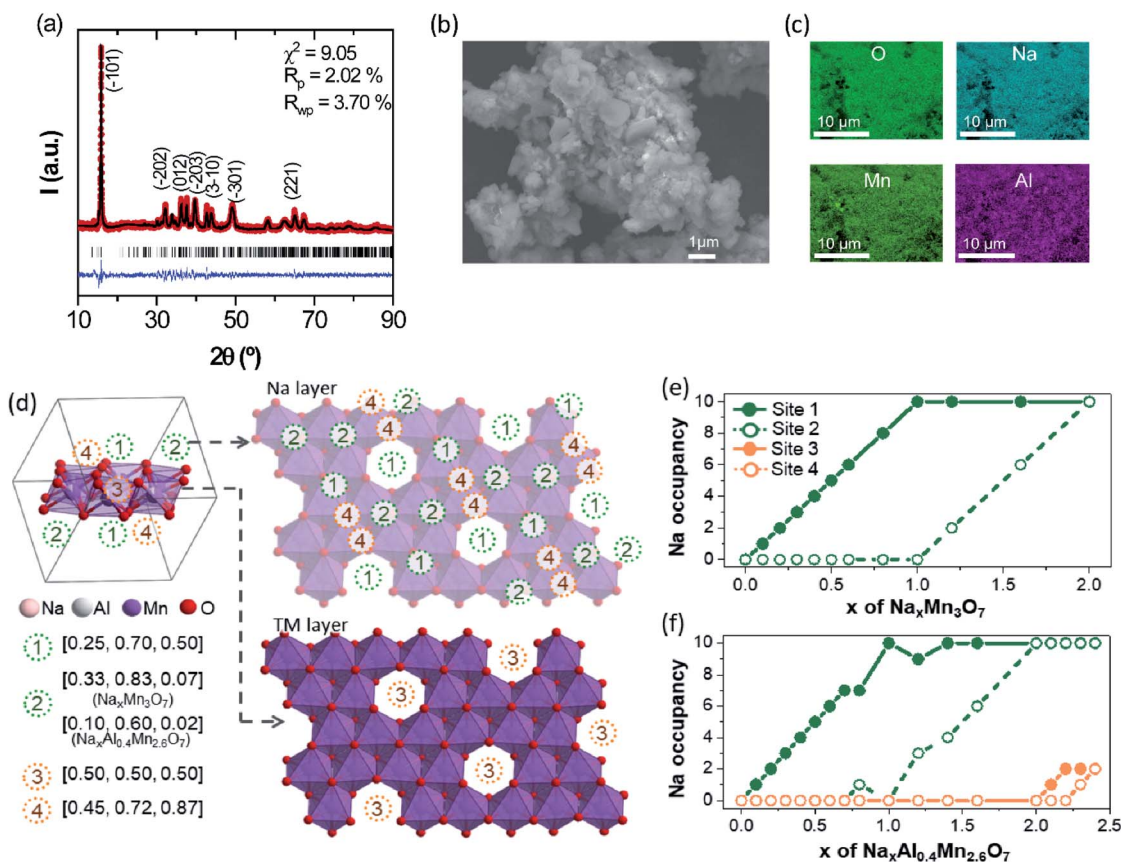


Fig. 1 (a) XRD data of layered $\text{Na}_{2.4}\text{Mn}_{2.6}\text{Al}_{0.4}\text{O}_7$ and corresponding Le Bail fit using $P1$ space group. Experimental data are presented in red, calculated data are shown in black and the difference between these two is shown in blue. Tick marks indicate the position of the calculated Bragg reflections. Miller indices are shown for the most intense reflections. (b) FE-SEM image of $\text{Na}_{2.4}\text{Mn}_{2.6}\text{Al}_{0.4}\text{O}_7$. (c) EDS mapping images showing O, Na, Mn and Al elements. (d) Schematics of predicted Na sites in $\text{Na}_x\text{Mn}_3\text{O}_7$ and $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. Top views of the Na and TM layers of supercells used in this study are shown. Number of occupied Na sites of the optimized (e) $\text{Na}_x\text{Mn}_3\text{O}_7$ and (f) $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ supercell structures, showing Na^+ ion insertion sequences during sodiation.

(where $M = e.g., \text{Mg}$ and Al) typically show peaks that can be assigned to details such as layering arrangements, defects, and Mn oxidation states.^{3,6,44} ^{23}Na MAS NMR spectra of $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ and $\text{Na}_2\text{Mn}_3\text{O}_7$ are shown in Fig. S7†. The spectrum for $\text{Na}_2\text{Mn}_3\text{O}_7$ (Fig. S7a†) shows an isotropic peak at ~ 1500 ppm. This is consistent with a large Fermi contact interaction that is expected for Na within a predominately Mn^{4+} setting, as expected for this material.^{3,44} Two sharp peaks close to 0 ppm surrounded by a spinning sideband manifold are also observed. Resonances close to 0 ppm suggest the presence of Na environments within a diamagnetic impurity phase.³ Although no secondary phases were observed in the XRD data, one explanation could be a low-level amorphous phase. The ^{23}Na MAS NMR spectrum for $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ (Fig. S7b†) shows a wide spectral envelope, consisting of broadened intensity centred around 1000 ppm, and spinning sideband manifolds centred around 500 and 0 ppm, the latter again showing the presence of a diamagnetic impurity. Comparison of spectra recorded at variable MAS rate (Fig. S8†) shows that the broad sideband manifold has isotropic peaks at 730 ppm, 500 ppm and 330 ppm. Comparison of both static NMR spectra (Fig. S9†) confirms the lower shifts for $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ and the presence

of two main groups at ~ 1000 and ~ 500 ppm. The lower shifts observed for $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ indicate a weaker Fermi contact interaction and the isotropic peaks at 500 and 330 ppm are consistent with the shift range for Na^+ ions within a local Mn^{3+} setting. In summary, the comparison of the ^{23}Na spectra for both materials shows the extra local environments for Na^+ ions and lower shifts corresponding to Mn^{3+} upon Al^{3+} ion substitution, where the former and latter results corroborate our DFT data (Fig. 1e, f and 4) and XANES data (Fig. S10†), respectively. However, the spectral broadening prevents more detailed analysis.

To gain further insight into the crystal structure of $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, DFT calculations were used to determine the atomic positions of additional Na^+ ions within the structure, assuming a disordered distribution of Al^{3+} and Mn^{4+} ions within the 2i sites (0.07, 0.07, 0.22) in the vacancy-ordered TM layer.^{8,42} The insertion sequences of $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ upon discharge were then predicted by modelling partially desodiated $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ (see Computational details for the construction of model structures, Fig. S1 and S2 in ESI†). Fig. 1d displays the symmetrically inequivalent Na sites predicted from $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, represented in the unit cell and supercell



Table 1 List of Na sites in $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, showing the neighbouring structures and coordination of each Na site. The formation energies of Na vacancies are also shown for reference. All structures and vacancy formation energies were calculated using HSE06 functional

Na site	$\text{Na}_2\text{Mn}_3\text{O}_7$		$\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$						
	1	2	1'	1''	1'''	2'	2''	3	4
Neighbouring structures	9 MnO_6 octahedra	6 MnO_6 octahedra	9 MnO_6 octahedra	7 MnO_6 and 2 AlO_6 octahedra	8 MnO_6 and 2 AlO_6 octahedra	4 MnO_6 octahedra	3 MnO_6 and 2 AlO_6 octahedra	4 MnO_6 and 2 AlO_6 octahedra	2 MnO_6 and 2 AlO_6 octahedra
Coordination	Tetrahedral	Octahedral	Prismatic	Prismatic	Prismatic	Prismatic	Prismatic	Prismatic	Octahedral
Vacancy formation energy (eV)	5.30	4.80	5.17	4.74	5.08	4.31	4.81	3.79	4.10

with ten formula units. In the supercell, four different Na sites were observed: three different 2i sites in the Na layer (Na1 (0.25, 0.70 and 0.50), Na2 (0.10, 0.6, 0.02) and Na4 (0.45, 0.72 and 0.87)), which may accommodate up to ten Na^+ ions per site; and one 1 h site in the TM layer, Na3 (0.5, 0.5, 0.5), which may uptake five Na^+ ions. These sites are identical to those observed in $\text{Na}_2\text{Mn}_3\text{O}_7$,^{6,39} except for the Na2 site which is found in a slightly different position *i.e.*, Na2 (0.33, 0.83, 0.07) due to excess Na in sites Na3 and Na4, which may repel some pre-existing Na^+ ions in the Na2 site.

Further computational studies show that Na1 and Na2 sites of $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ are occupied sequentially until $x = 2$, with a slight difference in Na active sites near $x = 1$ compared to $\text{Na}_x\text{Mn}_3\text{O}_7$ (Fig. 1e and f). Further Na^+ ion insertion occurs into the vacancy sites in the TM layer (Na3 site) until Na content reaches $x = 2.2$. Na^+ ion insertion into vacancies in the TM layer may increase the TM–Na repulsion within layers and Na–Na

repulsion between adjacent layers, increasing the cell volume, as observed in the XRD data (Fig. 1a). Last, Na^+ ion insertion occurs into site Na4 up to $x = 2.4$. The presence of Na in sites Na3 and Na4 causes slight displacements of Na sites such that the Td (site Na1) and Oh (site Na2) sites become trigonal prismatic (Table 1). The structure of $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, with Na^+ ions (*ca.* 0.2 Na) located in the TM layer resembles that of other cathode materials with demonstrated oxygen redox activity such as Li-rich layered oxides^{45–47} and other Na layered compounds.^{48,49}

Charge carrier pathways and structural changes

To provide a better insight into the changes in diffusion pathways after Al^{3+} ion substitution, we represented the potential Na^+ ion diffusion pathways of both $\text{Na}_x\text{Mn}_3\text{O}_7$ and $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. This was done by classifying the active Na sites in terms of their shape and displaying them on the TM layer



Fig. 2 (a) Changes in number and type of Na sites of $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ under differing Na content. Results are plotted only for $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ as $\text{Na}_x\text{Mn}_3\text{O}_7$ shows the same trends. Planar view of (b) $\text{Na}_x\text{Mn}_3\text{O}_7$ and (c) $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ with respect to Na content, showing the active Na sites and associated Na^+ ion diffusion pathways. (d) Schematics of suggested Na^+ ion diffusion pathways of $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ when excess Na^+ ions are intercalated into $\text{Na}_2\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$.



(Fig. 2). In $\text{Na}_x\text{Mn}_3\text{O}_7$, active Na sites at low sodiation levels ($0 < x \leq 1.0$) are characterized by prismatic Na1 sites above unoccupied Mn sites (Fig. 2a and b). Based on the distribution of active Na sites, two diffusion pathways can be suggested: in-plane 1D diffusion along with the prismatic Na1 sites, and interplanar diffusion through unoccupied Mn sites. As Na^+ ions continue to intercalate into $\text{Na}_x\text{Mn}_3\text{O}_7$ ($1.0 < x \leq 2.0$), repulsion between Na^+ ions increases and pushes Na1 sites toward unoccupied Mn sites, changing their coordination to tetrahedral. The so-formed tetrahedral Na1 sites are likely to be fixed under $x \geq 1.0$ (Fig. 1e) and block the diffusion pathways, which slows down the diffusion kinetics. Na^+ ion diffusion in highly sodiated $\text{Na}_x\text{Mn}_3\text{O}_7$, therefore, tends to follow 1D diffusion pathways through octahedral Na2 sites, which is expected to be more sluggish than prismatic–prismatic Na^+ ion diffusion at low sodiation content ($0 < x \leq 1.0$).⁵⁰

In the case of $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, Na^+ ion diffusion at high voltages ($0 < x \leq 0.7$) follows two diffusion pathways, similar to $\text{Na}_x\text{Mn}_3\text{O}_7$ under $0 < x \leq 1.0$ (Fig. 2c). However, with an additional supply of Na^+ ions ($0.7 < x \leq 1.2$), both prismatic Na1 and Na2 sites become active, unlike in $\text{Na}_x\text{Mn}_3\text{O}_7$ (Fig. 1f). This is speculated to be due to the higher electronegativity of Al compared to Mn, which destabilizes O 2p orbitals of Al–O bonds and, in turn, makes Na^+ ions stable on site 2 above AlO_6 octahedra. These

active Na2 sites bridge 1D diffusion pathways between Na1 sites and allow a broader range of diffusion along the in-planar direction, resulting in enhanced rate performance of $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. Under highly sodiated conditions ($1.2 < x \leq 2.0$), similar to $\text{Na}_x\text{Mn}_3\text{O}_7$, Na1 sites are fixed as tetrahedral sites near unoccupied Mn sites, causing Na^+ ion diffusion to occur *via* octahedral Na2 sites. Once all Na1 and Na2 sites are occupied, excess Na^+ ions intercalate into the unoccupied Mn site (*i.e.*, Na3) and nearby Na2 site (*i.e.*, Na4) (Fig. 2d). This process is speculated to occur with the cooperative motion of pre-existing Na1 or Na2 as most free spaces are already occupied by Na1 and Na2.

The above discussion suggests that Na^+ ion diffusion pathways are widened due to the high electronegativity of Al^{3+} ions. Moreover, subsequent structural analysis on the Al-substituted $\text{Na}_2\text{Mn}_3\text{O}_7$ cathode revealed that Al shortens the lattice along the in-planar direction of the TM layers, while increasing the interlayer spacing of these TM layers (see Discussion on the structural changes after Al substitution section and Fig. S11 and S12 in ESI†). Since most Na^+ ion diffusion pathways of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ cathodes are aligned along the in-planar direction, these structural changes will shorten the diffusion pathways while widening the diffusion channels, further improving the Na^+ ion diffusion kinetics. This may explain the enhanced rate performance observed for $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$

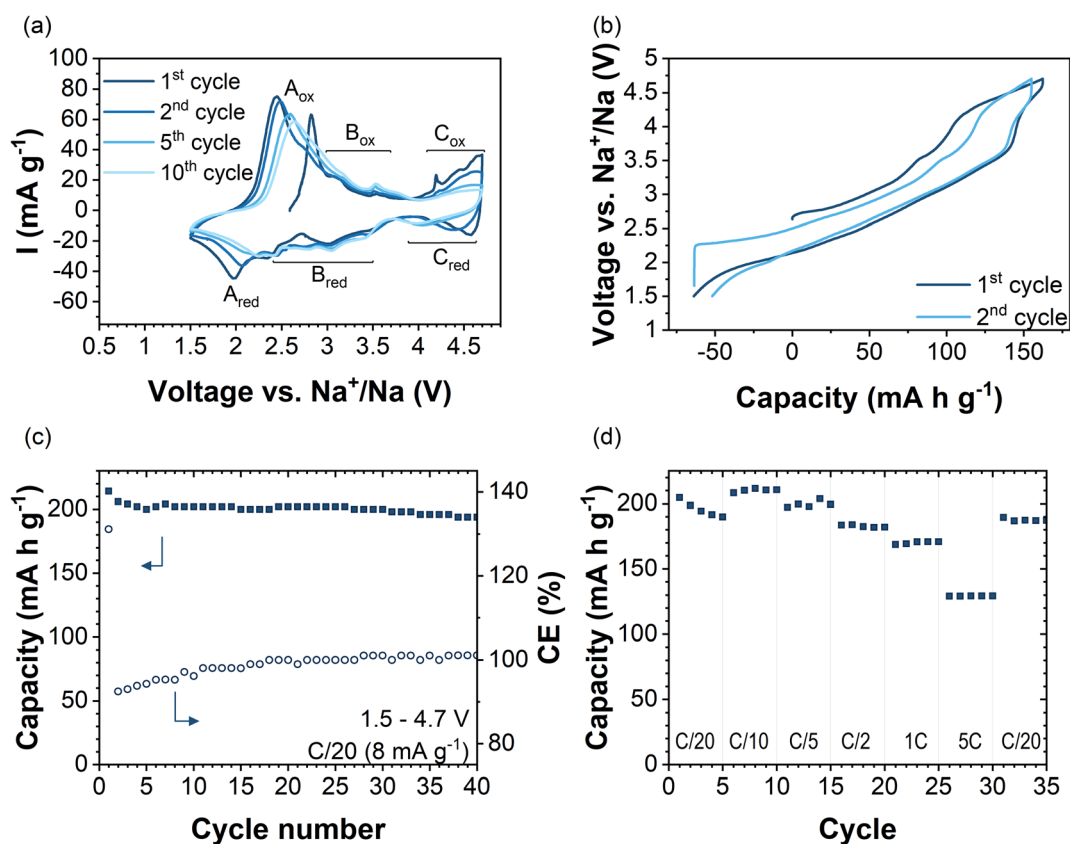


Fig. 3 (a) CV curves of $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ at a scan rate of 0.1 mV s^{-1} in the voltage range 1.5–4.7 V vs. Na^+/Na . (b) Galvanostatic charge–discharge voltage profiles in the voltage range 1.5–4.7 V vs. Na^+/Na at C/20 (8 mA g^{-1}) for cycles 1 and 2. (c) Discharge capacity (solid squares) and coulombic efficiency (open circles) as a function of cycle number in the voltage range 1.5–4.7 V vs. Na^+/Na at a current density of C/20 (8 mA g^{-1}). (d) Rate capability tests in the voltage range of 1.5–4.7 V vs. Na^+/Na using current densities ranging from C/20 to 5C.



discharge process.⁶⁰ Nyquist plots for $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_2$ and $\text{Na}_2\text{Mn}_3\text{O}_7$ showed larger charge-transfer resistance values for the former, becoming even more evident after 100 cycles (Fig. S17 and S18†). EIS data for other Al-substituted compounds, such as $\text{P2-Na}_{0.5}\text{Mn}_{0.5-x}\text{Al}_x\text{Co}_{0.5}\text{O}_2$ ¹⁵ and Al-doped $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$ ⁶¹ have also shown an increase in charge transfer resistance with Al^{3+} ion substitution, which has been attributed to the formation of a thicker cathode-electrolyte interface (CEI) layer at high states of charge and/or a more unstable surface.^{15,61,62}

Charge compensation mechanism

Ex situ XRD data were collected on cycled electrodes at different states of charge during the first cycle, to examine the structural arrangements occurring in $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_2$ (Fig. S19†). XRD data on the electrode charged to 4.7 V (red data) shows broadened peaks and changes in the relative intensities, suggesting an increase in disorder and atomic occupations with respect to the pristine electrode. Furthermore, an increase in the cell volume parameters attributed to an increase of the interlayer distance was observed. This behaviour is typical in layered oxides, and it is ascribed to the increased repulsion of the oxygen ions in adjacent layers in the absence of shielding effect of Na^+ ions in the Na-layer.^{63–65} No extra phases were observed, indicating a solid-solution behaviour. On the other hand, XRD data on the electrode that was charged and discharged (4.7–1.5 V voltage range) shows two extra peaks at $\approx 38.5^\circ$ and 42.8° 2θ (indicated with a diamond symbol in Fig. S19†). These are in good agreement with the orthorhombic $\text{Na}_{0.44}\text{MnO}_2$ phase (*Pbma* space group), previously observed during discharge in $\text{Na}_2\text{Mn}_3\text{O}_7$ (Fig. S20†).⁶⁶ Table S2† summarises the crystallographic data obtained by indexing the XRD data shown in Fig. S19† using the Le Bail method.

Ex situ ²³Na solid-state NMR spectra for $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ electrodes (pristine, charged and discharged) are shown in Fig. S21.† The spectrum for the pristine cathode is similar to that of the $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ powder with distinct features at ~ 500 ppm and ~ 730 ppm, and with an added intensity around 0 ppm due to residual electrolyte. The resonance at 330 ppm is hidden by a spinning sideband. In the charged electrode material, the 500 and 730 ppm features are not observed, and the spectrum instead shows intensity at a higher shift of 910 ppm. Although this shift is lower than the 1500 ppm observed for the undoped material, the increase in shift compared to the pristine doped material is consistent with the removal of Na^+ ions from the structure, resulting in a more Mn^{4+} ion rich local environment.^{3,44} The spectrum for the discharged material, whilst broadened, is similar to the pristine electrode material. For comparison, *ex situ* ²³Na SSNMR spectra were also obtained for $\text{Na}_2\text{Mn}_3\text{O}_7$ electrode materials (Fig. S22†). As for the $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ materials, there is an intense peak close to 0 ppm in the spectra, due to residual electrolyte and the diamagnetic component discussed previously. The pristine electrode material shows a main peak at ~ 1500 ppm, consistent with the $\text{Na}_2\text{Mn}_3\text{O}_7$ powder. Spectra obtained at different MAS rates show that the shift of this peak

may be temperature-dependent. The spectrum of the charged electrode material shows small added features at 930 ppm and 650 ppm (with one peak a spinning sideband of the other). The spectrum for the discharged electrode is different with the peak at ~ 1500 ppm not observed and the spinning sideband pattern of the broader component of the spectrum shifted to lower ppm and over a wider shift range. This is consistent with the discharged material having an increased sodium content, and therefore different local Na environments when compared to the pristine material.

Although new Na environments can be detected with NMR, it has not been possible to decide their nature. Therefore, to gain further understanding of the charge/discharge behaviour in $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ and its $\text{Na}_2\text{Mn}_3\text{O}_7$ parent material, we examined *via* computational methods the insertion sequences of Na atoms and relate the results to the voltage profiles of the cathode materials upon cycling.

For this purpose, the Na sites of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_2\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ occupied during charging/discharging were predicted by optimizing the geometries of partially sodiated $\text{Na}_x\text{Mn}_3\text{O}_7$ and $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ supercells. Upon cycling of the sodium layered electrodes, the local bonding between Na atoms and TM changes largely, whereas TM–O bonds mostly maintain their octahedral coordination. For such systems, the electrochemical potentials are primarily determined by the bond dissociation energy between active Na atoms and TM layers. From this perspective, the difference in Na^+ ion insertion sequences between $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ and their differing Na bonding environment during sodiation can result in differing voltage profiles. To verify this relationship, we first summarized all local bonding environments of Na sites in $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ structures (Table 1) and assessed how strongly Na^+ ions are bounded in each coordination environment by comparing their Na vacancy formation energies (Fig. 4a). Comparison in Na vacancy formation energies reveals that, of two Na sites in $\text{Na}_2\text{Mn}_3\text{O}_7$, Na^+ ions in site Na1 are more strongly bonded to TM layers, requiring higher potentials to remove them compared to those in site Na2. This matches the two voltage plateaus observed, where Na^+ ions from site Na2 are first removed at low voltage ($\text{Na}_2\text{Mn}_3\text{O}_7 \rightarrow \text{NaMn}_3\text{O}_7$), followed by the removal of Na^+ ions from site Na1 at high voltage ($\text{NaMn}_3\text{O}_7 \rightarrow \text{Mn}_3\text{O}_7$).

For $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, the presence of Al^{3+} ions and Na^+ ion excess changes the structure such that Na^+ ions in Na1 and Na2 sites have three and two different variations of local bonding environment, respectively (represented by prime symbols in Table 1). Consequently, Na vacancy formation energies may vary even for the same Na site: 4.74–5.17 eV for Na1 site and 4.31–4.81 eV for Na2 site. This shows that Na^+ ion extraction from these sites during charging occurs throughout various voltage values, resulting in sloping voltage profiles of $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ in the range of $0.4 < x < 2.0$ (Fig. 3b). The relatively high Na vacancy formation energies from Na1 and Na2 indicate that the removal of these Na sites will require high charge voltages (4.2–5 V), which corresponds to the high voltage plateau in Fig. 3b and redox peaks in region C in the CV data (Fig. 3a). On the other hand, the Na vacancy formation energies of sites Na3 and



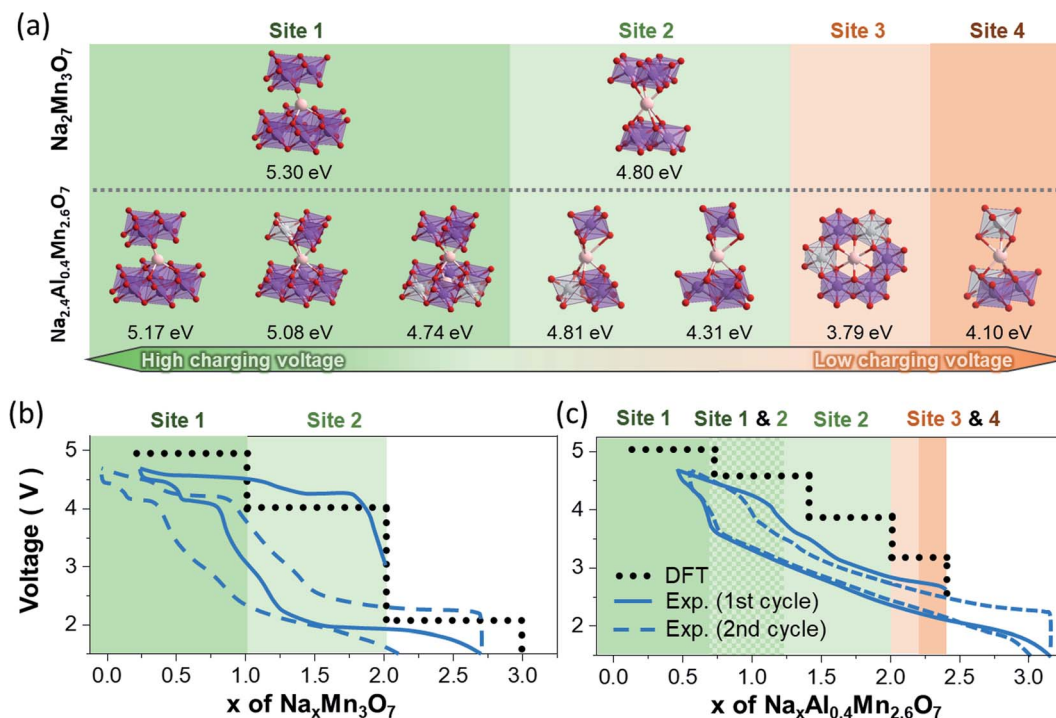


Fig. 4 (a) Local bonding environment of Na sites and their predicted Na vacancy formation energies in $\text{Na}_x\text{Mn}_3\text{O}_7$ and $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. Theoretical and experimental voltage profiles of (b) $\text{Na}_x\text{Mn}_3\text{O}_7$ and (c) $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. Electrochemically active Na sites for each Na content range are highlighted in differing colours.

Na4 are lower than those of sites Na1 and Na2, which is due to electrostatic repulsions with adjacent Na atoms in Na1 and Na2 sites. Such weakly bonded Na atoms are likely to be removed first during charging, which is represented as the low voltage anodic peaks in region B of the cyclic voltammetry data (Fig. 3a). Additional theoretical voltage calculations reveal that the model structures resemble well the first charge curves of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ (Fig. 4b and c): (1) predicted voltage profiles agree with the overall shape of the charging curves, including the two-stepwise profiles of $\text{Na}_2\text{Mn}_3\text{O}_7$ and the sloping curves in the $0.4 \leq x \leq 2.0$ region and the distinct low voltage plateau in the $2.0 \leq x \leq 2.4$ region of $\text{Na}_x\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. (2) When experimental voltage values are averaged over the x range of each voltage plateau, the averaged voltage values are close to the predicted voltages by less than 0.5 V. However, the above predictions are no longer consistent with the experimental data as charging and discharging processes proceed, due to structural changes which are not considered in these calculations.

Furthermore, we simulated the first desodiation process of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ and analysed its effect on electronic and atomic structures. Before desodiation, the valence band maxima (VBM) of both materials are predominantly formed by O 2p states and the conduction band minima are occupied by p-d hybridized states of O and Mn atoms (Fig. S23†).

Density of states calculations show that, after removing Na atoms, electron holes are likely to be centred on the 2p states of oxygen. Additional analysis of the local Madelung potentials of

O revealed that O atoms from Mn–O–Mn or Mn–O–Al near the unoccupied Mn sites in the TM layers (also referred to as non-bonding oxygen atoms⁸) have lower Madelung potentials (23.93–25.88 V) than those of other O atoms (27.39–29.86 V) (Table S3†). Considering that the energy to generate a positive hole is lower for lower Madelung energy,^{38,67} electron holes tend to be localized on 2p states of O near the unoccupied Mn sites. Such hole localizations can also be predicted from the partial density of states of O atoms (Fig. S24†), where high-energy electrons are concentrated on the non-bonding O atoms. As a result, the desodiation of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ causes the localized holes at 2p states of non-bonding O atoms, which can be observed after removing a Na atom from $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ (Fig. 5a–d). Such hole localization behaviours occur for all Na vacancies in $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, regardless of the type of Na sites (Fig. S25 and S26†). The above results indicate that the desodiation of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ proceeds by redox reactions from 2p states of non-bonding O atoms,⁶⁸ which is in good agreement with previous reports on $\text{Na}_2\text{Mn}_3\text{O}_7$.^{8,59}

Upon further desodiation, electron holes continuously accumulate in non-bonding O atoms and, in turn, induce some structural distortion in the lattice of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. To understand this coupling between excess holes and structural changes, we first analysed the structural changes near the non-bonding O atoms. Fig. 5e shows the changes in bond length between non-bonding O atoms in $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ before and after desodiation. In the case of $\text{Na}_2\text{Mn}_3\text{O}_7$, all O–O bonds near the unoccupied Mn sites have





Fig. 5 Projected density of states calculated for (a) $\text{Na}_{1.9}\text{Mn}_3\text{O}_7$ and (b) $\text{Na}_{2.3}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. Insets show the magnified density of states near hole levels. Partial charge densities (green isosurfaces) of (c) $\text{Na}_{1.9}\text{Mn}_3\text{O}_7$ and (d) $\text{Na}_{2.3}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$ calculated for the hole bands located ≈ 1.5 eV above the VBM. (e) Distributions of O–O bond lengths near the unoccupied Mn sites before and after the desodiation of $\text{Na}_2\text{Mn}_3\text{O}_7$ and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$. Distribution of hole polarons after desodiation are shown as insets, with labels for the corresponding types of O listed in Table 2. In the insets, the lengths of contracted O–O bonds are indicated by numbers.

similar bond lengths of ~ 2.8 Å. However, after the weakly bounded Na atoms in site Na2 (see Fig. 4a) are extracted upon charge, the resultant NaMn_3O_7 shows contracted O–O bonds with bond lengths ranging from 2.46 to 2.80 Å. Subsequent

calculations on the localized holes of NaMn_3O_7 reveal that, of four different O sites in $\text{Na}_2\text{Mn}_3\text{O}_7$ (Table 2), the three-coordinated O2 (O coordinated by two Mn and one Na) is the most preferential hole site because of its low Madelung

Table 2 List of O sites in NaMn_3O_7 , $\text{Na}_2\text{Mn}_3\text{O}_7$, $\text{Na}_{1.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, and $\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$, showing neighbouring atoms and oxidation states of O sites. Relative formation energies of O vacancies are shown for reference

O site	$\text{Na}_2\text{Mn}_3\text{O}_7$				NaMn_3O_7							
	1	2	3	4	1	2	2'	3	5			
Adjacent atoms	2 Mn and 2 Na	2 Mn and Na	3 Mn and Na	3 Mn and 2 Na	2 Mn and 2 Na	2 Mn and Na	2 Mn and Na	3 Mn and Na	3 Mn			
Oxidation state	2–	2–	2–	2–	2–	2–	<i>n</i> -(oxidized)	2–	2–			
Vacancy formation energy (eV)	0.143	0.0	1.372	1.297	0.54	0.39	0.0	1.07	1.10			
O site	$\text{Na}_{2.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$											
	1	2	3	4	5	6	7	8				
Adjacent atoms	2 Mn and 2 Na	2 Mn and 3 Na	Mn, Al and 3 Na	2 Mn, Al, and Na	2 Mn, Al, and 2 Na	3 Mn and Na	3 Mn and 2 Na	3 Mn and 3 Na				
Oxidation state	2–	2–	2–	2–	2–	2–	2–	2–				
Vacancy formation energy (eV)	0.0	0.152	0.504	0.811	1.244	1.133	1.228	1.275				
O site	$\text{Na}_{1.4}\text{Al}_{0.4}\text{Mn}_{2.6}\text{O}_7$											
	1	2	2'	3	4	5	5'	6	7	8	9	10
Adjacent atoms	2 Mn and 2 Na	2 Mn and Na	2 Mn and Na	Mn, Al, and Na	2 Mn, Al, and Na	Mn, Al, and 2 Na	Mn, Al, and 2 Na	2 Mn and Al	3 Mn and Na	3 Mn and Na	2 Mn, Al, and 2 Na	3 Mn and 2 Na
Oxidation state	2–	2–	<i>n</i> -(oxidized)	<i>n</i> -(oxidized)	2–	2–	<i>n</i> -(oxidized)	2–	2–	2–	2–	2–
Vacancy formation energy (eV)	0.174	0.636	0.482	0	0.535	0.31	0.061	0.52	1.845	1.141	0.972	1.694



potential (Tables S3 and S4[†]) and its proximity to Na vacancies. Such selective partial oxidization of O^{2-} to O^{n-} causes the contraction of bond distance to occur locally near O2 atoms (Fig. 5e and S27[†]). Additional calculations showed that the partially oxidized O^{n-} acted as main sites for the formation of O–O dimer species: (1) the vacancy formation energies of O sites showed that the partially oxidized O^{n-} required the lowest energy to be removed compared to other O sites (Table 2); and (2) nudged elastic band calculations showed that the formation of the O_2 dimer from the non-bonding O^{n-} species was a thermodynamically driven process (with reaction enthalpy of -0.29 eV), with a low activation barrier of 0.78 eV (Fig. S28[†]). Together with Mn ion migration to the Na layer, the kinetic barrier for O–O dimerization can be further lowered.⁶⁹ The above findings, *i.e.*, shortened bond distances of O2–O2 bonds, their low vacancy formation energies, and low activation barriers for O_2 dimer formation, suggest the formation of “peroxo-like” O–O dimers (and associated O_2 gas release) after the first charge of $Na_2Mn_3O_7$. The formation of peroxo-like dimers has been demonstrated with resonant inelastic X-ray scattering (RIXS) in other layered oxides such as $P2-Na_xLi_yMn_{1-y}O_2$ ($y = 0.2, 0.25$) with non-bonding O atoms.⁷⁰

The partial oxidation of non-bonding O atoms and associated O–O bond shortening can also be observed in $Na_{2.4}Al_{0.4}Mn_{2.6}O_7$ (Fig. 5e). Overall, O–O bonds near the unoccupied Mn sites are slightly distorted compared to $Na_2Mn_3O_7$ such that the O–O bond lengths vary from 2.78 to 2.84 Å. After desodiation, the O–O bonds of $Na_{1.4}Al_{0.4}Mn_{2.6}O_7$ are shortened to the lowest length of 2.42 Å. Such contraction in O–O bonds arise mainly from their partially oxidized O sites, *i.e.*, O2', O3, and O5' (Table 2 and Fig. 5e), where each oxygen atom has differing coordination environment; two Mn and Na atoms for O2'; Mn, Al and Na atoms for O3; and Mn, Al, and two Na atoms for O5'. Madelung potentials of Al-coordinated O (17.37 and 20.85 V for O3 and O5, respectively) are lower than that of two-Mn-coordinated O2' (20.86 V), indicating that, upon further charging, excess holes are likely to be localized in O atoms adjacent to Al. Additional calculations on the O vacancy formation energies show that among all O atoms nearby the unoccupied Mn sites (O1, O2, O2', O3, O5' and O5), partially oxidized O atoms, especially those coordinated by Al (O3 and O5'), are more readily removed upon charging. This suggests that owing to the preferential hole formation excess in Al-coordinated O atoms during charging, $Na_{2.4}Al_{0.4}Mn_{2.6}O_7$ can also suffer from the formation of O–O dimers and associated O loss upon charging.

Calculations on the Na and O vacancies of $Na_2Mn_3O_7$ and $Na_{2.4}Al_{0.4}Mn_{2.6}O_7$ enabled us to reveal the early stage of the anionic redox reactions during charge. Na^+ ion extraction during charging partially oxidizes non-bonding oxygen atoms adjacent to the unoccupied Mn sites, causing the oxygen atoms to have localized excess holes (Fig. 5a–d). The distribution of excess holes can be modified by external dopants, such as Al, where holes are preferentially located on non-bonding oxygen atoms coordinated by Al. Next, the so-formed electron holes in oxygen atoms shorten the O–O bond distance up to 2.4 Å (Fig. 5e). Compared to other oxygen atoms, those comprising

the shortened O–O bonds can be readily removed from their original sites (Table 2), which suggest the formation of oxygen dimers with shorter bond lengths, *e.g.*, peroxo-like O–O dimers, and O_2 gas. Song *et al.* hypothesised on the possibility of O_2 gas evolution from mostly surface but also bulk $Na_2Mn_3O_7$ particles together with a decrease of crystallinity and increase of particle stress, although without providing any further experimental evidence.⁶

The above findings of contracted O–O bonds and their weak binding to the TM layer were observed in both $Na_2Mn_3O_7$ and $Na_{2.4}Al_{0.4}Mn_{2.6}O_7$, suggesting that these are susceptible to O loss upon the first charge. Such O loss can also be supported from experiments: both cathodes display shortened high voltage plateaus during the first discharge process and subsequent charge–discharge cycles (Fig. 4b and c), which indirectly indicates a decrease in the amount of redox-active O. Furthermore, additional calculations on the electronic structures of $Na_2Mn_3O_7$ and $Na_{2.4}Al_{0.4}Mn_{2.6}O_7$ show that the removal of non-bonding O atoms leads to the formation of localized excess electrons on two adjacent Mn^{4+} ions (Fig. 6). The so-formed excess electrons reduce both Mn^{4+} ions to Mn^{3+} , where excess electrons positioned in Mn with lower Madelung potentials exhibit higher energy (Table S4 and Fig. S29–S31[†]). This suggests that O loss at high voltages causes the reduction of Mn^{4+} to Mn^{3+} to maintain charge neutrality, which can explain the mixed oxidation states of Mn at high voltages observed in XANES spectra. Fig. S10a and b[†] illustrates XANES spectra of Mn K-edge at various states of charge, together with the MnO

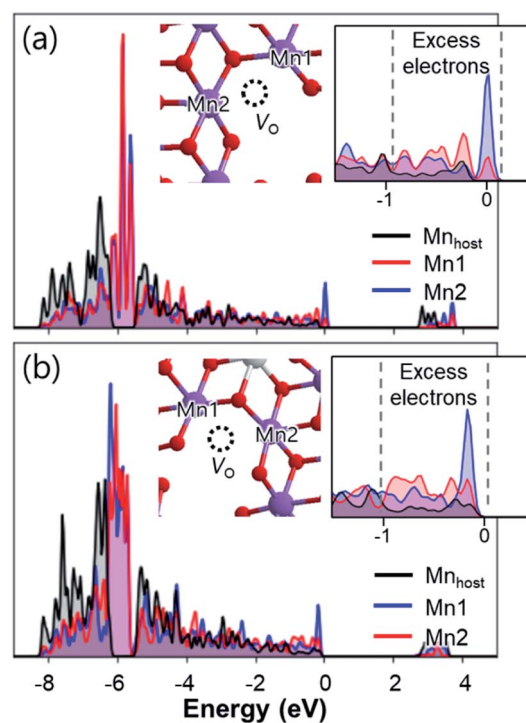


Fig. 6 Projected density of states of Mn atoms before (Mn_{host}) and after (Mn1 and Mn2) the formation of O vacancies in (a) $Na_2Mn_3O_7$ and (b) $Na_{2.4}Al_{0.4}Mn_{2.6}O_7$. Insets show O vacancy positions, reduced Mn atoms, and magnified density of states near excess electrons.



(Mn²⁺), Mn₂O₃ (Mn³⁺) and MnO₂ (Mn⁴⁺) references. Linear combination fit analyses (LCF) were used to semi-quantify the oxidation state of Mn at the different points of charge and discharge and the results are shown in Fig. S10c.† A report of the goodness of fit parameters (*R*-factor and reduced χ^2) along with the per cent that each model contributes to the fit is summarised in Table S5.† A similar reduction of the TM at highly charged states due to the oxygen redox activity has also been observed for Ni^{4+/3+} in Li_{1.2}Ni_{0.4}Ru_{0.4}O₂.⁷¹ Considering that the Mn³⁺O₆ octahedra are Jahn–Teller active,⁷² such reduction of Mn⁴⁺ ions might have detrimental effects on the cycling stability of the studied cathodes.

Although the above analyses have focused on O loss and subsequent effect on Mn reduction, electron holes formed in O atoms upon desodiation can cause other structural changes. In addition to O₂ gas release, these excess holes can shorten the O–O bond lengths to form various O₂ dimer species, such as peroxy-like O₂ⁿ⁻ or peroxide O₂²⁻-species.¹⁰ Previous studies on Li-ion layered cathodes have reported that the shortened O–O bond can trigger TM migration to octahedral sites in the Li layer and subsequent formation of spinel-like phases.⁷³ Other studies have also suggested the possibility of cation dopant migration to the tetrahedral site of the Na layer in layered oxides, which prevents further TM migration and stabilizes cyclability of cathodes.⁷⁴

Conclusions

In summary, we have studied the effects of Al³⁺ ion substitution and Na⁺ ion excess in the Mn-rich, Na₂Mn₃O₇ cathode material with triclinic *P*1 space group. We synthesised Na_{2.4}Al_{0.4}Mn_{2.6}O₇ (Na_{0.68}Mn_{0.74}Al_{0.11}[]_{0.15}O₂ in conventional Na_xMO₂ notation) as a reference compound by a ceramic method and its crystal structure was characterised with X-ray diffraction. Given the highly disordered nature of the cathode material, we used DFT calculations to provide an accurate crystallographic model of the structure. We then assessed the electrochemical properties of the Na_{2.4}Al_{0.4}Mn_{2.6}O₇ cathode in Na half-cells using different electrochemical methods, including galvanostatic cycling, cyclic voltammetry, electrochemical impedance spectroscopy, and galvanostatic intermittent titration. Cyclic voltammetric data showed the presence of both Mn^{3+/Mn⁴⁺} and O^{2-/Oⁿ⁻}-redox-active species in the material during Na cycling, which were further confirmed using XANES spectroscopy and computational studies. On the other hand, galvanostatic measurements showed an initial specific capacity of 215 mA h g⁻¹ in the 1.5–4.7 voltage range at C/20 and a capacity retention of 90% after 40 cycles. Experimental load curves were compared to computationally generated ones to gain insight into the Na⁺ ion insertion/extraction sequences generated upon cycling. Subsequent examinations on the modelled structures revealed that owing to the high electronegativity of Al³⁺ compared to the Mn⁴⁺ counterpart, Na_{2.4}Al_{0.4}Mn_{2.6}O₇ allows Na⁺ ions to diffuse throughout various directions, parallel to the TM layers. The shorter Al–O bond lengths and lower oxidation state of Al³⁺ ions distort the Na⁺ ion migration pathways such that these are shorter and have wider diffusion channels, which

further improve Na⁺ ion diffusion kinetics. These results suggest selection criteria of dopants for high-rate performance layered oxides cathodes, where these should have: (1) higher electronegativity, (2) smaller ionic radius, and (3) lower oxidation states than TM elements. Therefore, we expect that this concept can be extended for developing future fast charging layered oxide cathodes.

To understand the underlying chemistry under oxygen loss and capacity fading, we performed partial charge density calculations on electron holes of desodiated Na_{2.4}Al_{0.4}Mn_{2.6}O₇, and confirmed oxygen redox reactions to stem from 2p states of non-bonding O atoms, analogously to Na₂Mn₃O₇. In conjunction with O vacancy formation energies, energy barrier calculations suggested that the O atoms with electron holes during desodiation form “peroxy-like” O–O dimer species by O–O bond shortening. Furthermore, we showed that the preferential positions for O–O dimer formation can be adjusted by Al³⁺ ion substitution in the TM layer. Such behaviour arises from the change in electrostatic energy (Madelung potential) of O atoms neighbouring Al atoms. This suggests that one can adjust the O–O dimer distribution of desodiated cathodes by careful dopant selection, *i.e.*, by using dopants with low oxidation states and large ionic radii that lower the Madelung potential of adjacent oxygen atoms, one can force O–O dimer species to form nearby dopants, and *vice versa*. Considering that O–O dimers are a major cause of O₂ gas release¹⁰ and TM migration into the alkali metal layer,⁷³ these studies may guide future studies on developing cathode materials that are resistive to oxygen loss and structural degradation. In conjunction with future calculations on structural changes during anionic redox (*e.g.*, O, cation, and dopant migration), these findings may enable researchers to uncover the effect of dopants on anionic redox reactions.

Conflicts of interest

There are no conflicts to declare.

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