Nanoscale Horizons

The home for rapid reports of exceptional significance in nanoscience and nanotechnology rsc.li/nanoscale-horizons

Pm 1D EU Nd La Ho Yb d Tm Sc Gd

ROYAL SOCIETY | Celebrating OF CHEMISTRY | IVPT 2019 ISSN 2055-6756





Nanoscale Horizons

REVIEW



Cite this: Nanoscale Horiz., 2019, 4, 10

Received 2nd July 2018, Accepted 13th August 2018

DOI: 10.1039/c8nh00154e

rsc.li/nanoscale-horizons

NIR light-activated upconversion semiconductor photocatalysts

Qingyong Tian,^{†ab} Weijing Yao,^{†a} Wei Wu^b*^{ac} and Changzhong Jiang^a

Harvesting of near infrared (NIR) light in the abundant and environmentally friendly solar spectrum is particularly significant to enhance the utilization rate of the cleanest energy on earth. Appreciating the unique nonlinear optical properties of upconversion materials for converting low-energy incident light into high-energy radiation, they become the most promising candidates for fabricating NIR light-active photocatalytic systems by integrating with semiconductors. The present review summarizes recent NIR light-active photocatalytic systems based on a sequence of NaYF₄-based, fluoride-based, oxide-based and Ln³⁺ ion-doped semiconductor-based photocatalytic mechanisms and enhancement effects for efficient photo-redox performance of different upconversion semiconductor photocatalysts. We envision that this review can inspire multidisciplinary research interest in rational design and fabrication of efficient full-spectrum active (UV-visible-NIR) photocatalytic systems and their wider applications in solar energy conversion.

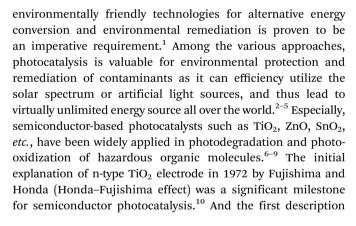
1. Introduction

Currently, steadily worsening environmental pollution is an overwhelming global crisis that threatens the health of human beings and sustainable development of natural ecosystems. To solve these urgent issues, the exploitation of green and

Wuhan University, Wuhan 430072, P. R. China. E-mail: weiwu@whu.edu.cn; Fax: +86-27-68778433; Tel: +86-27-68778529

^b College of Materials Science & Engineering, Zhengzhou University, Zhengzhou 450052, P. R. China

† Contributions are equal.





Qingyong Tian

Qingyong Tian received his BS degree at Hubei University in 2013. He is currently pursuing a PhD degree under the supervision of Prof. Changzhong Jiang and Prof. Wei Wu in the School of Physics and Technology at Wuhan University. His research focuses on the design and synthesis of multifunctional nanostructured materials for applications in photocatalysts and advanced energy devices.



Weijing Yao

Weijing Yao received her BS in 2015 from Inner Mongolia Agricultural University. She is currently pursuing a PhD degree under the supervision of Prof. Wei Wu in the School of Printing and Packaging at Wuhan University. Her research involves the design and synthesis of upconversion materials for anti-counterfeiting applications.



View Article Online

^a School of Printing and Packaging and School of Physics and Technology,

^c Suzhou Research Institute of Wuhan University, Suzhou 215000, P. R. China

Review

of TiO₂ suspension for the photodegradation of polychlorinated biphenyl and cyanide by Cary and Bard in 1980 opened a new gate for TiO₂-based semiconductors for the elimination of environmental pollution.^{11,12} TiO₂ possesses the advantages of cost-effectiveness and pollution-free and sustainable development and has been widely investigated in recent decades. However, the wide bandgap of TiO_2 (>3.0 eV) has been a major limitation for efficient utilization of solar photons.^{13–15} The light absorption threshold is confined in the ultraviolet (UV) region (300-400 nm), which only accounts for 5% of the full spectrum. Therefore, the losses in solar conversion efficiency have resulted in under-utilization of visible (43%, 400-700 nm) and NIR light (52%, 700-2500 nm). Narrow bandgap semiconductors, such as CdS, Fe2O3, g-C3N4, etc., could broaden the absorption threshold to the visible region and demonstrate efficient photocatalytic performance.^{16,17} NIR-light-active semiconductors with narrower bandgap could further absorb NIR light, while the lower-lying conduction band (CB) and higherlying valence band (VB) present fairly inert photodecomposition or photo-oxidization abilities to organic dyes. Carbon quantum dots (CQDs) with both up- and down-converted photoluminescence (PL) and photoinduced electron transfer ability could give rise to intense absorption of NIR light.^{18,19} The surface plasmon resonance (SPR) peaks of noble-metal plasmonic nanoparticles can be red-shifted to the NIR region by rational adjustment of sizes, morphologies and alloving compositions.^{20,21} Their tunable optical characteristics make CODs and noble-metal plasmonic nanoparticles alternative co-catalysts to harvest NIR photons and enhance the NIRinduced photocatalytic performance. However, the inherent photocorrosion susceptibility and high processing cost make them far from practical propositions. Therefore, much effort should be made to introduce NIR light into semiconductorbased photocatalyst systems.^{22,23}

One increasingly appealing and practical approach is coupling photocatalytic semiconductors with nonlinear optical materials that could utilize the abundant NIR light in photocatalysis, to improve the utilization of solar energy, and finally enhance the photoreactions.^{24,25} Of especial note are the anti-Stokes luminescence mechanism-based upconversion materials, with the unique property to sequentially absorb two or multiple low-energy photons and radiate out specific types of photons with shorter wavelength, which are considered as potentially efficient candidates for exploitation of NIR excitation. Lanthanide ion (Ln³⁺)-doped upconversion materials with abundant f-orbital configurations have narrow fluorescence emission bands via intra-4f (4f'-4f') or 4f-5d (4f'-4f''-15d) orbital transitions.^{26,27} The involved symmetries of inner electronic quantum states would yield metastable energy levels with long lifetimes that favor the incidence of successive populations in the energized states of Ln³⁺ ions. Under NIR laser irradiation, the pump photons in the ground state of sensitizer ions will be sequentially excited to the excited states of activator ions. Then, the population will relax non-radiatively or radiatively to the ground state of activator ions to achieve the upconversion emissions. The perfect shielding effect of outer 5s and 5p shells of Ln³⁺ prevents the influence of surrounding environment on electronic transitions, resulting in perfect physical and chemical stability against photobleaching and photochemical degradation. Doping with different Ln³⁺ ions in an inorganic substrate will induce designable upconversion emissions, from the UV, passing through the entire visible and extending to the NIR region under irradiation.²⁸⁻³⁰ Fig. 1 presents the occupation of Ln³⁺ ion-doped upconversion materials with large absorption cross section for converting the NIR light in the solar spectrum into visible light. The main adsorption centers of the Ln³⁺ ions, *i.e.*, Nd³⁺, Yb³⁺, Ho³⁺ and Er³⁺, are located at 808 (Nd³⁺: ${}^{4}I_{9/2} - {}^{4}F_{5/2}$), 980 (Yb³⁺: ${}^{2}F_{7/2} - {}^{2}F_{5/2}$), 1160 (Ho³⁺: ${}^{5}I_{8} - {}^{5}I_{6}$)



Wei Wu

Wei Wu received his PhD in 2011 from the Department of Physics, Wuhan University, China. He then joined the group of Prof. Daiwen Pang at Wuhan University (2011) and Prof. V. A. L. Roy at City University of Hong Kong (2014) as a postdoctoral fellow. Now he is a full professor and Director of Laboratory of Printable Func-Nanomaterials tional and Printed Electronics, School of Printing and Packaging, Wuhan University. He has published over

80 papers, which have received over 3000 citations. He is also the topical editor of Frontiers in Materials. His research interests include the synthesis and application of functional nanomaterials, printed electronics and intelligent packaging.



Changzhong Jiang

and co-author, more than 100 publications in various reputed international journals, such as Physical Review Letters, Nano Letters, ACS Nano, and Advanced Materials. His research interests include the synthesis and application of low-dimension nanomaterials, magnetic materials and ion beam modification of materials.

of

Chang Zhong Jiang received his BS

in 1983 from Huazhong University

of Science and Technology and MS

in 1990 from Wuhan University.

He obtained his PhD in 1999 from

Université Claude Bernard Lyon 1,

France. He has been a full professor in the Department of

Physics, Wuhan University since

2001, and he is also the Director

Application, Wuhan University.

He has published, as an author

Ion Beam

Center for

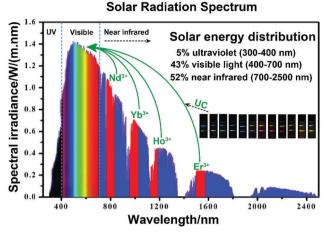


Fig. 1 The schematic conversion processing of NIR light into visible light $via \ln^{3+}$ ions, referring to standard AM 1.5 solar radiation spectrum, and the tunable emissions of upconversion materials.

and 1530 (Er³⁺: ⁴I_{15/2}-⁴I_{13/2}) nm, respectively.³¹ These ions will act as sensitizers that absorb the centered NIR photons and emit UV or visible photons through multiple photon absorption or energy transfer processes. For instance, tunable emissions varying from blue to red colors can be realized by fine-tuning the concentration of sensitizer Yb³⁺ ions and changing the doped activator Tm³⁺ and Er³⁺ ions.³² The spectral power from the normalized air mass (AM) 1.5 spectrum yield is nearly 25 W m⁻² ranging from 1480 to 1580 nm when Er^{3+} acts as the activator.³³ It will be extremely attractive to combine upconversion materials as frequency conversion elements with UV- or visible-active semiconductor photocatalysts to further harvest solar energy in a broader spectral range, and thus promoting the utilization of solar energy efficiently.

2. The Ln³⁺ ion-doped upconversion mechanism

The upconversion emissions are based on the transition of 4f electronic configurations between rare earth (RE) ions (Sc, Y and La are technically d-shield elements).³⁴ Each kind of RE ion has its definite energy level positions, and the upconversion processes between different RE ions are different from each other, while all of them represent similar property of chemical elements attributed to the dominant trivalent oxidation state. The characteristic energy level positions of different Ln³⁺ ions are illustrated in Fig. 2, which was well defined by experimental measurement by Dieke (Dieke diagram) and Carnall et al.31,35 The energy level positions give detailed illustration of absorption and excitation spectra of Ln³⁺ ion-doped upconversion materials with distinguishable spectroscopic fingerprints. The upconversion process has been recognized to be grouped into three stages: excited state absorption (ESA), energy transfer (ET), and phonon avalanche (PA).

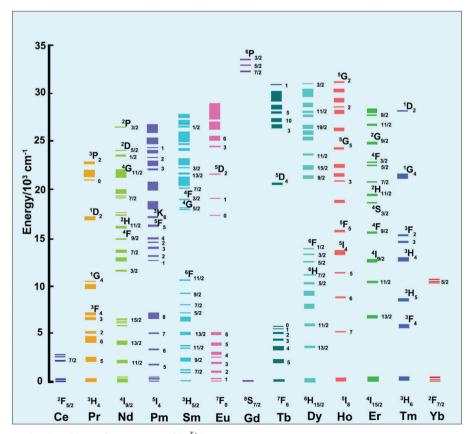


Fig. 2 Characteristic energy level positions of a sequence of Ln³⁺ ions

2.1 Excited state absorption (ESA)

ESA is the basic process of upconversion luminescence first proposed by Bloembergen et al. in 1959.³⁶ The mechanism is that one type of Ln^{3+} ion is excited from the ground state (GS) to an excited state with high energy level through sequential twoor multiphoton absorption. Typically, as illustrated in Fig. 3a, the Ln³⁺ ion absorbed one photon hv₁ and is promoted from GS to intermediate metastable (electronic) state (E_1) . Because of the long relaxation time of metastable energy level, the excited ions will stay or decay between the intra-band for a long time. And then, if another energy-matched photon matches the energy interval between E_1 and upper excited state E_2 , the ions at state E₁ will absorb the photon energy and populate the state E2, achieving successive absorption of two photons. Three or more multiphoton absorption processes will be deduced by analogy. Finally, the excited ions relax non-radiatively or radiatively to the GS, inducing the upconversion emissions. Additionally, supposing the incident photons are sequentially trapped by individual Ln³⁺ ions, then the upconversion efficiency of the ESA process is independent of the Ln³⁺ doping concentration. The usual doping concentration of Ln³⁺ ions is less than 1%.³⁷ Under continuous excitation, the intensity of upconversion luminescence (I_{uc}) is proportional to excitation power (I_{ex}) :

$$I_{\rm em} \propto (I_{\rm ex})^n$$
 (1)

E3

E2

E2

where *n* is the number of pump photons.

2.2 Energy transfer upconversion (ETU)

E2

processes, respectively

The ETU process generally occurs between two types of ions (donor and acceptor) in close proximity via dipole-dipole interaction. High concentration of dopants is required in ETU process-dominated upconversion materials, as the dopant content determines the distance between the proximal ions. ETU can be classified into three categories based on diverse

F1 G G М (f) EMU (e) PA Fig. 3 Principal diagram for the upconversion processes of Ln³⁺-doped crystals. (a) ESA, (b) SET, (c) CR, (d) CU, (e) PA and (f) EMU. The red, green, and purple lines stand for photon excitation, energy transfer, and emission

energy transfer types: successive energy transfer (SET), cross relaxation (CR) and cooperative upconversion (CU).

2.2.1 Successive energy transfer (SET). The SET process is illustrated in Fig. 3b. The donor ions in the excited state should match the energy requirement with proximal acceptor ions in the ground state. Consequently, the donor ions will transfer energy to the acceptor ions and excite them into the excited state (E1), albeit themselves will return to the GS through nonradiative relaxation. The excited acceptor ions may be promoted to higher excited state (E_2) via successive energy transfer.

2.2.2 Cross relaxation (CR). The CR process between one or two types of neighboring ions is displayed in Fig. 3c. Both Ion₁ and Ion_2 (Ion_1) are populating the intermediate energy level $(E_1 \text{ and } E_2)$. Then, Ion_2 (Ion_1) would be motivated to higher intermediate energy level (E3) in accordance with the nonradiatively energy transfer processes of proximal Ion₁, accompanied by the relaxation of Ion₁ to the ground state.

2.2.3 Cooperative upconversion (CU). Three ions are involved in the CU process (Fig. 3d), where the two excited ions (Ion₁ and Ion_2) transfer the energy to the third ion (Ion_3) simultaneously and emit Ion₃ from the GS to higher energy level, while Ion₁ and Ion2 relax non-radiatively to the GS. The CU process involves the interaction between three ions, and the acceptor ions require a matched metastable energy level with donor ions; therefore, the CU process always occurs with low probability.31

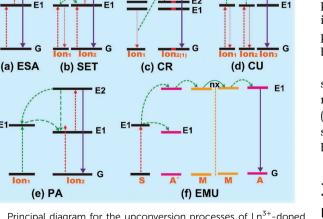
2.3 Phonon avalanche (PA)

The PA process was proposed firstly from an investigation of the upconversion spectrum of Pr³⁺-doped lanthanum chloride (LaCl₃) host by Chivian.³⁸ The PA process is an unconventional looping process of ESA and CR, which contains four procedures of energy transfer (Fig. 3e). (I) The Ion_2 initially emerges to the intermediate level (E_1) by GS energy transition. (II) This is followed by energy transfer by ESA or ETU process to populate the higher excited energy level (E_2) where the upconversion luminescence occurred. And then (III) a CR process occurred among activated Ion₂ and proximal ground-state Ion₁, causing two populations in the metastable state E_1 . (IV) Finally, the looping process readily populates and exponentially increases the population of excited state E2 through ESA process, invoking intense upconversion radiation like an avalanche process. The PA process is highly dependent on the pump intensity, which should be higher than the threshold value of metastable energy level.³⁹

In addition, with the rapid development of core-shell nanostructured upconversion materials, a new energy migrationmediated upconversion (EMU) mechanism was proposed recently (Fig. 3f).^{40,41} This involves successive energy transfer from the sensitizers (S) to accumulators (A') and migrators (M), and trapped by activators (A), finally inducing the upconversion process.

3. The Ln³⁺ ion-doped upconversion materials

The upconversion efficiency of Ln³⁺ ion-doped materials is reliant on the crystal structures, local site symmetry, crystal



field strength, and phonon energy of the host matrixes. It will be a key factor to choose appropriate host matrixes to achieve favorable optical performance of upconversion materials with high upconversion luminescence efficiency. Currently, the majority of upconversion materials are doped with Ln³⁺ ions into appropriate dielectric host matrix. The doped Ln³⁺ ions are randomly distributed in the cation sites, which are offered by the ordered array of host matrixes.42 The host matrixes generally do not constitute the exciting energy levels but will provide suitable crystal field for the doped Ln³⁺ ions that enables corresponding emissions. Host matrixes with low phonon energy are preferred for effective upconversion luminescence, because multi-phonon supported non-radiative relaxations between the nearest spaced energy levels can be minimized, thus greatly prolonging the lifetime of intermediate levels. The non-radiative relaxation probability (W_p) induced by multi-phonon relaxation is expressed as:43,44

$$W_{\rm p} = C \left[\frac{\exp(hw/KT)}{\exp(hw/KT)^{-1}} \right]^p {\rm e}^{-\alpha \Delta E}$$
(2)

where P is the order of phonons $(P = \Delta E/hw)$, ΔE is the gap between two energy levels, w is the phonon frequency, h is the Boltzmann constant, and α and C are constants related to matrix materials. P will increase with decreasing phonon energy, resulting in the decrease of non-radiative relaxation possibility, thus enhancing the upconversion luminescence efficiency. Furthermore, the increased lifetime of metastable energy levels will prohibit the quenching of radiative transitions in the upconversion process. The phonon energy depends on the crystal lattice of host matrixes, and charge and diameter of anions implanted in the crystal. When the phonon energy matches the frequency of excitation or emission, the absorption of the crystal lattice can reduce the emission efficiency, and thus the host matrix is desirable to occupy with lower phonon energy. The most investigated host matrixes include low phonon energy (fluorides $\sim 355 \text{ cm}^{-1}$, chlorides $\sim 260 \text{ cm}^{-1}$, oxides $\sim 600 \text{ cm}^{-1}$, iodides $\sim 144 \text{ cm}^{-1}$, bromides $\sim 172 \text{ cm}^{-1}$, etc.) and high phonon energy (silicates $\sim 1000 \text{ cm}^{-1}$, borates ~1100 cm⁻¹, phosphates, etc.) materials, $^{35,45-47}$ as listed in Table 1. Considering the effect of the photon energy of host matrixes on the upconversion luminescence, the efficiency increases in the order of oxides < fluorides < chlorides < bromides < iodides. Comparatively, heavy halides are hygroscopic, and the structural stability of the matrixes decreased in inversed proportional of oxides > fluorides > chlorides. Therefore, a series of investigations were carried out, expecting to find new-fashioned host matrixes with advantages of high upconversion efficiency (like fluorides and chlorides) and high stability (like oxides), to achieve practical applications. Typically, the hexagonal NaYF₄ crystal with low phonon energy and excellent chemical stability is accepted to be the most efficient host material hitherto.

In the following section, the effectiveness of commonly investigated upconversion materials with different host matrixes that are applied in the field of photocatalysis, including a sequence of $NaYF_4$ -based, fluoride-based, oxide-based and

Table 1 The phonon energy of various host matrixes

| Host matrix | Phonon energy (cm ⁻¹) | Host matrix | Phonon energy (cm ⁻¹) |
|-------------|--------------------------------------|-------------|--------------------------------------|
| Fluorides | ~355 | Silicates | ~ 1000 |
| Chlorides | ~ 260 | Borates | ~ 1100 |
| Iodides | ~ 144 | Phosphates | \sim 1300 |
| Bromides | ~ 172 | Bismuthates | ~ 1192 |
| Oxides | ~ 600 | | |

 ${\rm Ln}^{3+}$ ion-doped semiconductor-based photocatalysts, is discussed in detail.

4. The upconversion semiconductorbased photocatalysts

4.1 The mechanism of upconversion semiconductor-based photocatalysts

The proposed photocatalytic mechanism of semiconductorbased upconversion photocatalysts is illustrated in Fig. 4. The upconversion materials serve as the frequency conversion element and harvest two or multiple NIR light photons and then convert them into UV or visible light photons *via* an efficient fluorescent resonance energy transfer (FRET) process. Subsequently, the converted photons can be re-absorbed by the surrounding semiconductors and inject the electrons from VB to CB of semiconductors, thereby increasing the yield of photoinduced carriers. The direct energy migration between the combination of UC materials and semiconductors under NIR light irradiation could extend the utilization of the solar spectrum and promote the photocatalytic efficiency.

4.2 The NaYF₄-based upconversion photocatalysts

RE ion-doped sodium fluoride matrix contains two different types of crystal phases, including cubic and hexagonal phases (Fig. 5a and b). In the cubic phase, the host F^- ions form a high-symmetry cubic crystal structure, with Na⁺ and RE³⁺ ions randomly occupying the cation sites. Comparatively, the Na⁺ and RE³⁺ ions are selectively occupying the two types of relatively low-symmetry cation sites that are contained in regular crystal sublattices of F^- ions in hexagonal NaYF₄.⁴²

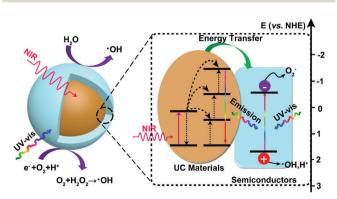


Fig. 4 The proposed photocatalytic enhancement mechanism for semiconductor-based upconversion photocatalysts.

Generally, host crystal lattices with low site symmetry of doping lanthanoids are preferable for upconversion luminescence over their higher symmetry counterparts, because low-symmetry lattices ordinarily have a crystal field with lots of uneven components surrounding the doping ions. The uneven components strengthen the electronic coupling among 4f electronic orbits and their higher electronic configuration and subsequently promote 4f-4f transition possibilities of the doping ions.⁴⁸ The density functional theory (DFT)-calculated charge densities for cubic and hexagonal NaREF4 crystals are presented in Fig. 5c and d. In the hexagonal phase, the Ln3+ dopants increase the asymmetry distribution and improve the charge transfer, resulting in an enhancement of polarizability and formation energy. Therefore, hexagonal phase NaREF₄ nanocrystals exhibited an order of magnitude augmentation in upconversion efficiency when compared with cubic phase.⁴⁹ Especially, the hexagonal phase NaYF₄, having the advantages of wide energy bandgap (\sim 8 eV)⁵⁰ and a very low phonon energy (\sim 360 cm⁻¹ \approx 45 meV),^{51,52} has been considered as one of the most efficient upconversion host matrixes for decades. And most photocatalysts have been combined with hexagonal phase NaYF₄ to achieve efficient NIR-active photocatalytic systems.

The photodegradation performance of NaYF₄-based semiconductor photocatalysts is presented in Table 2. As mentioned above, TiO₂ attracts extensive interest as a photocatalyst because the numerous advantages like abundance, chemical stability, environmental friendliness, and high photocatalytic activities. Incorporating upconversion materials with TiO₂ will extend the absorption edge up to the NIR region and finally achieve higher photocatalytic efficiency.^{53–57} For example, core–shell NaYF₄:Yb,Tm@TiO₂ nanoparticles show excellent activities under both UV and NIR irradiation for photodegradation of RhB dyes.⁵⁵ And the photocatalytic activities of NaYF₄:Yb,Tm@TiO₂ nanoparticles vary with the sizes of the

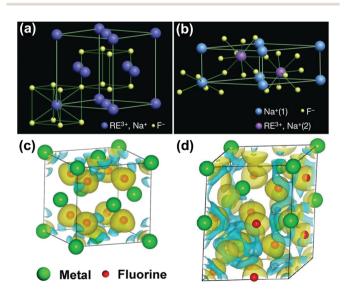


Fig. 5 (a) Cubic and (b) hexagonal type crystal phases of Ln^{3+} -doped NaYF₄ and DFT-calculated charge density for (c) cubic and (d) hexagonal NaYF₄. (Reproduced with permission.⁴² Copyright 2010, Springer Nature.)

NaYF₄ core (Fig. 6a-d), the optimum core size being about 50 nm, which is beneficial to the photocatalytic performance of the TiO₂ shell. In addition, the size-dependent upconversion mechanism of NaYF4:Yb3+,Er3+ nanoparticles has been experimentally and theoretically demonstrated, confirming the competition process between the upconversion population and non-radiative relaxation.⁵⁸ Another group further synthesized shell thickness-controllable NaYF4@TiO2 core-shell composites by varying the ratio of Yb³⁺,Tm³⁺ co-doped NaYF₄ microrods and Ti precursors (Fig. 6e-h), proving that optimum shell thickness is beneficial to the photocatalytic performance.⁵⁷ As displayed in Fig. 6i, when a thin TiO₂ shell is coated, only a small part of the UV radiation would be captured by the TiO₂ layer and most of it will be scattered and wasted. On the contrary, a thick TiO₂ shell will greatly inhibit the upconverted emissions penetrating to the external surface and largely increase the recombination possibilities of photogenerated carriers, which is also detrimental to photocatalytic performance. Furthermore, other upconversion material-semiconductor composite photocatalytic systems are rationally designed. For example, NaYF4:Yb/Tm-BiOCl,⁵⁹ NaYF4:Yb/Er-MoS2,⁶⁰ NaYF4:Yb/Tm-g-C₃N₄,^{61,62} NaYF₄:Yb/Tm-CdS,^{63,64} NaYF₄:Yb/Er-CdSe,⁶⁵ NaYF₄:Yb/ Tm-ZnO,^{66,67} and NaYF₄:Yb/Tm-MOFs⁶⁸ have been designed and constructed as NIR-responsive photocatalysts. Ternary heterogeneous materials are also developed and applied as highly efficient UV-visible-NIR full-spectra active photocatalysts.^{69,70}

Unfortunately, upconversion nanomaterials always suffer the drawbacks of low quantum yields and intense luminescence quenching, mainly resulting from the comparatively insensitive extinction coefficients of lanthanide dopants.²⁷ Innovative methods have been investigated for enhancing the upconversion intensity in lanthanide-doped nanocrystals, including surface plasmon coupling, surface passivation, energy transfer modulation, etc.²⁷ Especially, the surface plasmon coupling of noble metals with upconversion material-semiconductor system presents fascinating NIR-driven photocatalytic performance.71-74 Surface plasmon coupling of noble metals with upconversion materials improves the upconversion luminescence from three aspects as follows: (1) enhancing the absorption of the sensitizer through electric-field coupling, (2) improving the radiative decay rate of the activator, and (3) increasing energy transfer from the sensitizer to the activator, as displayed in Fig. 7a. The SPR-induced local electric field could enhance the photon flux and affect the radiative decay rates of the activators, thereby leading to enhanced upconversion luminescence and high quantum yield.75,76 In our previous work, finite-difference time-domain analysis (FDTD) was used to stimulate the electromagnetic field intensity of upconversion and plasmonic enhanced photocatalysts under each of UV, visible and NIR light irradiation (Fig. 7b-e). The SPR effect of Ag nanoparticles leads to higher photocarrier yield and improves the efficient transfer of photocarriers and energy, resulting in quintupled enhancement of photocatalytic efficiency.72 The plasmon and upconversion enhanced broadband photocatalysts NYF@TiO2-Au are synthesized with significantly superior photocatalytic performance to the commercial P25 TiO2.71 Fig. 7f illustrates the NIR-active photocatalytic mechanism.

Table 2 Photodegradation performance of NaYF₄-based semiconductor photocatalysts

| Photocatalyst (mg) | Light source | Pollutant ^{<i>a</i>} (concentration, mg L ⁻¹ /volume, mL) | Degradation ratio/time (h) | Ref |
|---|---|---|-------------------------------|----------|
| NaYF ₄ :Yb/Tm@TiO ₂ (0.5) | 980 nm laser (10 W cm ⁻²) | MB (15/0.5) | 65%/14 | 53 |
| $NaYF_4:Yb/Tm@TiO_2(20)$ | 980 nm laser (0.5 A) | $RhB(10^{-5} M/10)$ | 22%/1 (83%/6) | 54 |
| 4 2 2 () | ARC lamp (69920) | Phenol (10^{-5} M/50) | 17%/1 (53%/6) | |
| $NaYF_4:Yb,Tm@TiO_2(7)$ | 50 W Xe lamp | RhB (5 \times 10 ⁻⁵ M/50) | 100%/0.5 | 55 |
| NaYF ₄ :Yb/Tm@TiO ₂ | 1 W 980 nm laser | MB (20/0.5) | 68%/24 | 56 |
| 4 2 | | MO (15/0.5) | 69%/24 | |
| | | RhB (10/0.5) | 76%/76 | |
| NaYF ₄ :Yb/Tm@TiO ₂ (10) | 1.5 W 980 nm laser | MB (5/1) | 90%/12 | 57 |
| NaYF ₄ :Yb/Tm-BiOCl (20) | 300 W Xe lamp, $\lambda > 780$ nm | RhB (10/50) | 100%/2.7 | 59 |
| | ······································ | MB (10/50) | 58%/2.7 | |
| MoS_2 -NaYF ₄ :Yb/Er (20) | 980 nm laser | RhB(25/10) | 61%/12 | 60 |
| NaYF ₄ :Yb/Tm/g- C_3N_4 (1) | 1 W 980 nm laser | MB (15/0.5) | 83%/6 | 61 |
| (4114) 111/S 031(4 (1) | 300 W Xe lamp | 1112 (10,010) | 89%/1.5 | 01 |
| $NaYF_4:Yb/Tm/g-C_3N_4$ (10) | 1 W 980 nm laser | RhB (10/100) | 22%/7 | 62 |
| 14114.10/111/g 0314 (10) | 50 W Xe lamp | 101D (10,100) | 80%/1.75 | 02 |
| NaYF ₄ :Yb/Tm–CdS (20) | 2 W 976 nm laser | RhB (10/10) | 13%/3 | 63 |
| (40) | 300 W Tu lamp | (50/20) | 72%/2.5 | 05 |
| NaYF ₄ :Yb/Tm@C@CdS (10) | 50 W Xe lamp | (30/20) RhB (10 ⁻⁵ M/50) | 92%/1.2 | 64 |
| Marr4.10/111@0@000 (10) | 50 W Xe lamp | $\frac{M}{M} (10^{-5} \text{ M/50})$ | 95%/1.2 | 04 |
| NaYF ₄ :Yb/Er-CdSe (0.5) | 2 W 1560 nm laser | MB (15/0.5) | 65%/30 | 65 |
| 10114.10/11-003e(0.3) | 2 W 980 nm laser | MB (15/0.5) | 75%/30 | 05 |
| NaYF4:Yb/Tm-ZnO (0.5) | 2 W 980 nm laser | RhB (20/0.5) | 65%/30 | 66 |
| | 980 nm laser (50 mW cm $^{-2}$) | | | |
| $NaYF_4:Yb/Tm@SiO_2-ZnO(50)$ | 50 W Xe lamp | RhB (10/50) | 17%/7 89%/1.7 | 67 |
| NaYF ₄ :Yb/Tm/MOFs (20) | Xe lamp (100 mW cm ^{-2}) | RhB (10/50) | 85%/1.5 | 68 |
| | 2 W 980 nm laser | | | 68 69 |
| NaYF ₄ :Yb/Tm@TiO ₂ /RGO (5) | | MB(10/5) MO(15/5) | 81.1%/12 (90%/1) | 69 |
| | 300 W Xe lamp | MO (15/5) Phenol (20/5) | 45.9%/12 (90%/1) | |
| N-VE $V_{\rm c}/{\rm D}_{\rm c}$ | 2 W 000 mm lanen | | 29.8%/12 (61%/1) | 70 |
| NaYF ₄ :Yb/Er@TiO ₂ -Ag ₆ Si ₂ O ₇ (5) | 2 W 980 nm laser | MB (10/10) | 45%/5 | 70 |
| | 300 W Xe lamp | | 100%/2 min | -1 |
| $NaYF_4$ TiO ₂ -Au (10) | 8 W UV lamps | MO (20/100) | 95%/1 | 71 |
| | 300 W Xe lamp ($\lambda > 420$ nm) | | 83%/8 | |
| | 2 W 980 nm laser | | 80%/20 | |
| | 150 W solar simulator | $\mathbf{P} = \mathbf{P} \left(1 0 / 1 0 \right)$ | 80%/2 | |
| $NaYF_4:Yb/Tm@SnO_2/Ag(5)$ | 2 W 980 nm laser | RhB (10/10) | 100%/5 | 72 |
| | 300 W Xe lamp | $\mathbf{p} \in \mathcal{C}$ $(1 \circ 1^{-5} 1 1^{-5} 1)$ | 100%/0.17 | |
| NaYF ₄ :Yb/Tm@TiO ₂ /Ag | 50 W Xe lamp | R6G (10^{-5} M/50) | 96%/2 | 73 |
| NaGdF ₄ :Yb/Er-mTiO ₂ | AM 1.5 solar light | RhB | 99%/0.8 | 77 |
| NaGdF ₄ :Er/Yb@BiFeO ₃ (20) | 980 nm laser (80 mW mm ^{-2}) | MO (10/20) | 55%/3.5 | 78 |
| | 500 W Xe lamp ($\lambda > 420$ nm) | | 80%/15 | |
| NaLuF4:Gd,Yb,Tm@SiO2@TiO2:Mo | 500 W Xe lamp | RhB (20/50) | 100%/3.5 | 79 |
| | 980 nm laser | | 90%/3.5 | |
| NaGdF ₄ :Yb:Er/Ag/TiO ₂ (8) | 1.5 W 980 nm laser | RhB (5/5) | 40%/18 | 80 |
| (40) | 500 W UV-vis light | (40/40) | 97%/0.5 | |
| NaYF ₄ :Yb/Tm@NaYF ₄ /TiO ₂ (10) | 8 W 980 nm laser | RhB (1.43/5) | 41%/6 | 85 |
| | 250 W Xe lamp | | 73%/0.42 | |
| NaYF4:Tm/Yb@Yb/Nd@N-TiO2 | Xe lamp (100 mW cm ^{-2}) | RhB (10^{-5} M/50) | 96%/1.5 | 86 |

^a MB represents methylene blue; RhB represents Rhodamine B; MO represents methyl orange; R6G represents Rhodamine 6G.

Under NIR light irradiation, Er^{3+} and Tm^{3+} co-doped NaYF₄ excites UV and visible emissions, which overlap with the intrinsic light absorption of TiO₂ and SPR band of Au nanoparticles. The excited UV emissions will assist the TiO₂ to generate reductive electrons and highly oxidative holes. The excited visible emissions will interact with the SPR effect of Au nanoparticles to facilitate the transfer of hot electrons to nearby TiO₂ layer and minimize the electron–hole recombination. Finally, the broadband, effective nanohybrid photocatalyst improves the efficient electron and energy transfer among constituent components that leads to excellent photocatalytic activity. However, the direct contact of noble metal nanoparticles with upconversion materials is prejudicial to the fluorescence efficiency, because direct electron transfer or

FRET occurs.²⁷ For boosting the fluorescence emissions, a minimum distance between them should be maintained of at least 5 nm, which is propitious for decreasing the non-radiative processes and increasing the radiative decay rate.⁷⁶

Other Ln^{3+} ion-doped sodium fluoride-based phosphors such as NaGdF₄ and NaLuF₄ also have been investigated as promising frequency conversion candidates for NIR photocatalysts.⁷⁷⁻⁸⁰ For example, β -NaLuF₄ and tetragonal-phase LiLuF₄ are reported as novel host matrixes that may show higher yields of upconversion efficiency than β -NaYF₄,⁸¹⁻⁸³ ascribed to the distinctive electronic states at the top valence band of Lu³⁺ ions and the smaller lattice-cell volume counterparts.

Besides, to construct core-shell structural upconversion nanocrystals through epitaxial growth of small lattice-mismatched

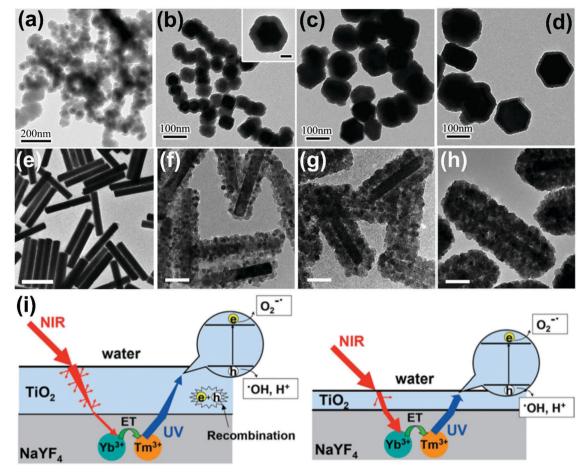


Fig. 6 TEM images of NaYF₄:Yb,Tm nanoplates of different sizes: (a) 30 nm, (b) 50 nm, (c) 100 nm, (d) 130 nm covered with amorphous TiO₂ shell. (Reproduced with permission.⁵⁵ Copyright 2014, Royal Society of Chemistry.) TEM images of (e) NaYF₄:Yb³⁺,Tm³⁺ microrods and (f–h) NaY-F₄:Yb³⁺,Tm³⁺@TiO₂ composites with different thickness of TiO₂ shell (scale bars: (e) 1 μ m, (f–h) 100 nm). (Reproduced with permission.⁵⁷ Copyright 2013, Royal Society of Chemistry.) (i) The effect of various TiO₂ shells on photocatalytic mechanisms (left: thick shell; right: moderate shell). (Reproduced with permission.⁵⁷ Copyright 2013, Royal Society of Chemistry.)

layers is excepted as a fascinating approach for improving the upconversion features, by reducing the non-radiative decay losses. The outer layers would confine the doping ions in the interior core crystals and thus inhibit the surface energy loss, resulting in improved upconversion yields. For example, excellent upconversion efficiency improvements of about 7- and 79-fold were realized after epitaxial coating of NaYF4 shell on the NaYF4: Yb,Er(Tm) core, respectively.⁸⁴ Attributed to valid surface passivation, Huang et al. achieved record high upconversion quantum yield of 7.6% for LiLuF4:Yb,Tm@LiLuF4 core/shell upconversion nanoparticles.⁸¹ Until now, there are only two studies applied to the core/shell structured NaYF4:Yb/Tm@NaYF4 and NaYF4:Tm/ Yb@Yb/Nd luminescent cores as energy transducers that combine with TiO₂ semiconductor for environmental purification.^{85,86} Therefore, there are still bright prospects for development of core/shell upconversion nanoparticles for photocatalytic application.

4.3 The fluoride-based upconversion photocatalysts

Because of the closely arranged energy gaps of Ln³⁺ ions, host matrixes with low photonic energy are needed to overcome the non-radiative loss between the excited levels for realizing

sensitive fluorescent transitions. Yttrium fluoride (YF₃) takes advantage of low vibration energy ($\sim 355 \text{ cm}^{-1}$) that could reduce the quenching of electron transitions within the 4f shell, and has been selected as an appropriate candidate for matrix materials.²⁵ A model of the orthorhombic YF₃ structure is displayed in Fig. 8a, where the Ln³⁺ ions will mainly occupy the interstitial sites of YF3 crystals at low doping concentration and substitute the sites of Y³⁺ ions at high doping concentration.^{87,88} Various trivalent lanthanide ion-doped YF₃ crystals with microand nanostructural morphologies (Fig. 8b-h) have been synthesized with hydrothermal process, reverse micelle method, thermal decomposition method, micro-emulsion method and microwave method by alternatively parameterizing the reaction conditions.⁸⁹⁻⁹³ Additionally, the sizes, morphologies and phases could be manipulated for colorful emitting luminescence properties of Ln³⁺ ion-doped YF₃ nanocrystals which could act as frequency conversion elements for wide application in NIR-driven photocatalysis.

Also, Ln³⁺ ion-doped fluoride nanocrystals are coupled with conventional semiconductors to fabricate NIR-active photocatalysts.⁹⁴ For instance, peanut-like YF₃:Yb³⁺,Tm³⁺ nanocrystals

Review

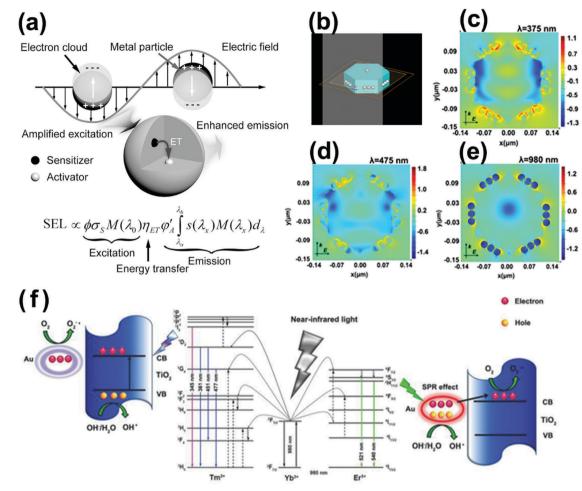


Fig. 7 (a) Plausible mechanism for plasmonic enhancement of upconversion emissions. (Reproduced with permission.²⁷ Copyright 2014, John Wiley and Sons.) (b–e) FDTD-calculated local electric field enhancement for NaYF₄@SnO₂@Ag nanoparticles at different incident wavelengths of 375, 475 and 980 nm. (Reproduced with permission.⁷² Copyright 2017, American Chemical Society.) (f) Schematic photocatalytic mechanism of NYF@TiO₂–Au under NIR excitation. (Reproduced with permission.⁷¹ Copyright 2015, John Wiley and Sons.)

with high specific surface were coated with uniform TiO₂ shell, expanding the light utilization threshold to the NIR region.95 Bi^{3+} ions were further doped in YF₃ by substituting Y³⁺ ions to tailor the energy transfer process between Yb³⁺ with Er³⁺/Tm³⁺. The Bi³⁺ ions possessing larger ionic radius will increase the cell volume of matrixes and adjust the crystal field symmetry, therefore prolonging the lifetimes of the metastable energy states by changing their wave functions, and finally enhancing the upconversion emissions.^{96,97} Additionally, the vacancies and point defects in Ln^{3+} ion-doped calcium fluoride (CaF₂) matrix will induce charge compensation effects, thus increasing the energy transfer rates and the upconversion probability.98,99 The non-toxic, stable, abundant and nonhygroscopic features of CaF2 have received increasing attention in fabrication of NIR-driven photocatalysts.^{100–105} For example, by integrating Er³⁺/Tm³⁺/Yb³⁺ tri-doped CaF₂ nanocrystal with magnetic ZnFe₂O₄ and ZnO, a multi-component and multi-functional NIR photocatalyst, ETY-FCZ, was fabricated.¹⁰⁴ The ETY-FCZ presents continuous degradation of salicylic acid (SA) (Fig. 8i) and MO (Fig. 8j) with 71.9% and 43.73% removal rates when exposed to visible-NIR light irradiation for 180 min, respectively. Fig. 8k presents the proposed photocatalytic mechanism of the ETY-FCZ system, where the upconversion emitted UV-visible photons boosted the generation rates of electron-hole pairs of the ZnO-ZnFe₂O₄ heterostructures. Examples of fluoride-based upconversion photocatalysts based on the recent literature are presented in Table 3.

4.4 The RE sesquioxide-based upconversion photocatalysts

Although the phonon energy of RE sesquioxide materials (Y_2O_3 , Gd_2O_3 , Lu_2O_3 , La_2O_3 , etc.) is high, their excellent chemical/ thermal stability makes them promising luminescence matrix candidates. Y_2O_3 is regarded as the best sesquioxide matrix for lanthanoids with smallest phonon energy ($\sim 380 \text{ cm}^{-1}$) among other metallic oxides and similar ionic radius (90 pm) to lanthanide elements. The ionic radii of Ln^{3+} are listed in Table 4, where the ionic radii (assuming six-coordination) reduce from 106.1 pm for La^{3+} ions to 84.8 pm for Lu^{3+} ions accompanied with an increase of atomic number. The similar ionic radii of Ln^{3+} ions to host cations will hinder the

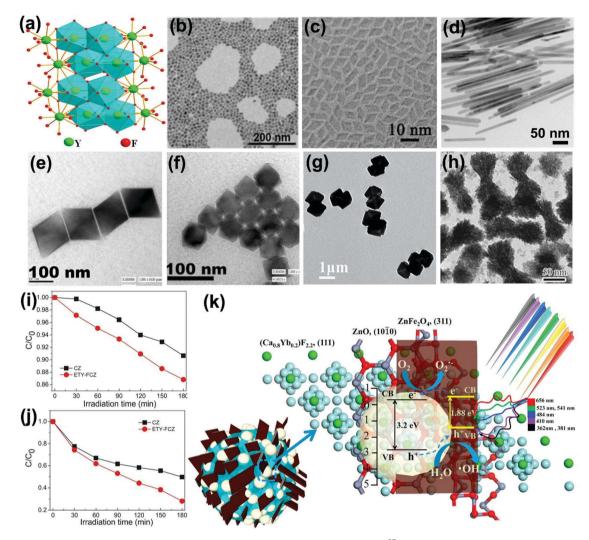


Fig. 8 (a) Schematic representation of orthorhombic structure. (Reproduced with permission.⁸⁷ Copyright 2017, American Chemical Society.) Typical YF₃ morphologies of (b) nanospheres (reproduced with permission;⁸⁸ copyright 2014, Royal Society of Chemistry), (c) rhombic shapes (reproduced with permission;⁸⁹ copyright 2014, Royal Society of Chemistry), (d) nanorods (reproduced with permission;⁹¹ copyright 2017, American Chemical Society), (e) quadrilateral-shaped particles (reproduced with permission;⁹⁰ copyright 2005, American Chemical Society), (g) dumbbell-shaped particles (reproduced with permission;⁹² copyright 2017, Royal Society of Chemistry) and (h) peanut-like nanocrystals (reproduced with permission;⁹⁵ copyright 2013, Elsevier BV). The photocatalytic performances of as-fabricated NIR-driven photocatalyst ETY-FCZ for degradation of (i) SA and (j) MO, and (k) the proposed photocatalytic mechanism (reproduced with permission;¹⁰⁴ copyright 2015, Elsevier BV).

Society of Chemistry) and (h) peanutas-fabricated NIR-driven photocatalys permission;¹⁰⁴ copyright 2015, Elsevie generation of crystal defects and tating the upconversion efficie

generation of crystal defects and lattice stress, and thus facilitating the upconversion efficiency.¹⁰⁶ More importantly, the distance of interacting ions (*D*) dominates the electric dipole– dipole energy transfer efficiency (*P*) between sensitizer and activator ions, which follows the relationship:¹⁰⁷

$$P \sim \frac{1}{D^6} \tag{3}$$

where the nearest interaction distances of RE ions confined in the host lattice of Y_2O_3 , Gd_2O_3 , Lu_2O_3 and La_2O_3 are presented in Table 5. Comparatively, it can be predicted that a more efficient energy transfer could be achieved with the same lanthanide dopants in the Lu_2O_3 host than in the Gd_2O_3 and La_2O_3 hosts, and their upconversion efficiency may obey the sequence of Lu₂O₃ > Gd₂O₃ > La₂O₃. As mentioned above, the low vibration energy of Y_2O_3 matrix improves the probability of radiative transitions among electronic energy levels of lanthanide dopants, inducing the highest interacting energy transfer rate in sesquioxides.

Schematic illustrations of RE sesquioxide unit cells modeled according to the structural parameters in Table 5 are shown in Fig. 9.^{108–110} The nanocrystalline systems of RE₂O₃ (RE = Y, Gd, Lu) prefer the formation of cubic structure with space group of *Ia*3 (No. 206) (Fig. 9a and b), while La₂O₃ is predisposed to unconventional hexagonal phase with space group of *P*3*m*1 (No. 164) (Fig. 9c and d). Theoretically, the polymorphism of RE sesquioxides with lighter atoms (La–Sm) generally means that they adopt the hexagonal structure and the heavier atoms

 Table 3
 Photodegradation performances of fluoride-based semiconductor photocatalysts

| Photocatalyst (mg) | Light source | Pollutants (concentration, mg L ⁻¹ /volume, mL) | Degradation ratio/time (min) | Ref. |
|---|--|--|---------------------------------|------|
| YF ₃ :Yb,Tm/TiO ₂ (0.5) | 1.5 W 980 nm laser | MB (15/0.5) | 70%/1800 | 94 |
| YF_3 :Yb,Tm(a)TiO_2 (20) | 50 W Xe lamp | MO (50/20) | 85%/70 | 95 |
| $NaY(WO_4)_2/TiO_2/YF_3:Yb,Tm$ (20) | 1000 W mercury lamp | MO (10/20) | 100%/60 | 96 |
| | | SA (10/20) | 90%/60 | |
| CaF_2 :Er,Yb@TiO ₂ (50) | 500 W Xe lamp (720 nm $< \lambda < 1$ 100 nm) | MB (10/50) | 60%/720 | 100 |
| 2, 2, 2, 2, , | | MO (10/50) | 32%/720 | |
| BiVO ₄ /CaF ₂ :Er,Tm,Yb (0.5) | 980 nm laser (10 W cm $^{-2}$) | MB (15/0.5) | 75%/420 | 101 |
| $CaWO_4$ (a) $(TiO_2/CaF_2:Er/Tm/Yb)$ (20) | 980 nm laser (2 A) | MO (10/20) | 20.29%/180 | 102 |
| | 1000 W mercury lamp ($\lambda \ge 780$ nm) | | 46.57%/180 | |
| $CaWO_4$ (TiO_2/CaF_2 : Er/Tm/Yb) (20) | 980 nm laser $(2 A)$ | MO (10/20) | 20.29%/180 | 103 |
| | 1000 W mercury lamp | | 99.02%/60 | |
| CaF ₂ :Er/Tm/Yb@ZnFe ₂ O ₄ /ZnO (20) | 1000 W mercury lamp ($\lambda \ge 400 \text{ nm}$) | MO (10/20) | 71.9%/180 | 104 |
| 2 0 21 () | | SA (10/20) | 43.73%/180 | |
| CaF_2 :Yb@BiVO ₄ (20) | 980 nm laser (2 A) | MO (10/10) | 10%/360 | 105 |
| 2 0 | 1000 W mercury lamp | | 41%/180 | |

Table 4 The ionic radii of Ln³⁺

| Element | La | Ce | Pr | Nd | Pm | Sm | Eu | Gd |
|---|----|-------------|----|----|----|----|----|----|
| Atomic number Ionic radius ^{<i>a</i>} (10 ⁻¹² m, pm) | | 58 103.4 | | | | | | |
| Element | Tb | Dy | Но | Er | Tm | Yb | Lu | _ |
| Atomic number Ionic radius (10 ⁻¹² m, pm) | | | | | | | | |
| ^{<i>a</i>} Assuming six-coordination. | | | | | | | | |

(Ho–Lu) prefer the cubic, whereas the intermediate atoms (Sm–Dy) can be stabilized in both cubic and monoclinic structures (Fig. 9e).¹⁰⁸ The lighter Ln³⁺ ions with larger ionic radii prefer the formation of hexagonal crystals in virtue of higher electron cloud distortion induced by increased dipole polarizability.⁴² The results will offer us reasonable guidance for accommodating appropriate Ln³⁺ dopants in morphological construction engineering.

RE sesquioxide nanocrystals incorporated with semiconductors were employed for synthesis of photocatalysts for environmental cleaning and energy conversion.^{111–116} For example, the Yb³⁺,Tm³⁺ co-doped Y₂O₃ upconversion luminescence agent was composited with TiO₂ and showed higher infrared-to-UV photocatalytic activity than non-doped TiO₂ under irradiation of solar light.¹¹⁴ Another chain-like $Er/Tm/Yb/Y_2O_3$ composite demonstrated higher upconversion luminescent intensity than flake-like $Er/Tm/Yb/Y_2O_3$ upconversion nanoparticles, leading to better photocatalytic performance for degradation of MB dyes when coupling with 3C-SiC nanocrystals.¹¹³ Furthermore, the Gd³⁺ ions with magnetic properties are used to fabricate optically and magnetically active bifunctional materials for advanced multifunctional photocatalysts.^{112,117}

4.5 The Ln³⁺ ion-doped upconversion photocatalysts

In recent years, the doping of ions into a semiconductor matrix, including non-metal ions (P^{3-} , S^{2-} , Cl^- , *etc.*) and metal ions (Zn^{2+} , Cu^{2+} , Fe^{3+} , *etc.*), has become popular and a prime motivation to extend the optical absorption threshold, or trap the photogenerated electrons by induced oxygen vacancy and/or surface defect, finally improving the photoactivities.^{4,9,118,119} Similarly, the current strategies of doping with Ln^{3+} ions into semiconductors such as TiO₂, ZnO, BiVO₄, *etc.*, have proved to be an interesting and not yet fully explored highway for optimizing the photocatalytic performance.^{120–128} As displayed in Fig. 10, 4f orbital electronic states of Ln^{3+} ions bring the sub-bandgap levels that settled down below the conduction band edge of semiconductors. The electrons will populate the inserted electronic states of Ln^{3+} ions instead of intrinsic conduction

| Table 5 Structural parameters of RE2O3 (RE: Y, La, Gd, Lu) nanocrystals | | | | | | | |
|---|--|--|---|---|--|--|--|
| Hosts | Y_2O_3 | Gd_2O_3 | Lu_2O_3 | La_2O_3 | | | |
| Phonon energy (cm ⁻¹) | ~380 | ~600 | ~618 | ~ 400 | | | |
| Bandgap (eV) | 5.6 | 5.4 | 6 | 5.5 | | | |
| Ionic radius ^{<i>a</i>} (10^{-12} m, pm) | 90 | 93.8 | 84.8 | 106.1 | | | |
| Crystal system | Cubic | Cubic | Cubic | Hexagonal | | | |
| Space group | Ia3 (206) | Ia3 (206) | Ia3 (206) | $P\bar{3}m\bar{1}$ (164) | | | |
| Lattice parameters | a = b = c = 10.6073 Å | a = b = c = 10.8012 Å | $a = \dot{b} = \dot{c} = 10.3946 \text{ Å}$ | a = b = 3.9198 Å c = 6.1139 Å | | | |
| | $\alpha = \beta = \gamma = 90^{\circ}$ | $\alpha = \beta = \gamma = 90^{\circ}$ | $\alpha = \beta = \gamma = 90^{\circ}$ | $\alpha = \beta = 90^\circ, \gamma = 120^\circ$ | | | |
| Nearest RE distance (Å) | 3.52 (Y1-Y2) | 3.59 (Gd1–Gd2) | 3.44 (Lu1–Lu2) | 3.83 (La–La) | | | |
| | 3.53 (Y2-Y2) | 3.61 (Gd2–Gd2) | 3.46 (Lu2–Lu2) | | | | |

^{*a*} Assuming six-coordination.

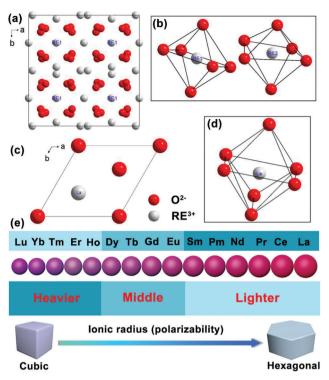


Fig. 9 Schematic illustration of the unit cell projected on a-b plane and the RE coordination polyhedra: (a and b) RE₂O₃ (RE = Y, Gd, Lu), (c and d) La₂O₃ (reproduced with permission;¹⁰⁸ copyright 2018, Elsevier BV). (e) lonic radius (or polarizability) induced phase transition from cubic to hexagonal of Ln³⁺ ion-doped upconversion crystals.

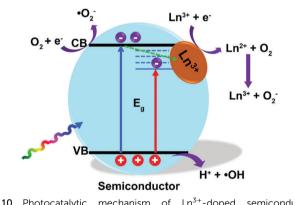


Fig. 10 Photocatalytic mechanism of Ln^{3+} -doped semiconductor photocatalysts.

band edge of semiconductors, thus narrowing the bandgap energy of photocatalysts and extending the optical absorption threshold.^{129,130} For example, the bandgap energy values of B-BiVO₄ with different doping concentrations of Eu³⁺ ions varied from 0 to 1.0% were calculated by using the Kubelka–Munk function. The bandgaps were tailored from 2.36 eV to 2.26 eV by raising the dopant dosage of Eu³⁺ ions, inducing approximately nine times higher photocatalytic activity (0.8 Eu–B-BiVO₄) compared with pure BiVO₄.¹³¹

Additionally, the lanthanides are well known for being capable of forming complexes with a variety of Lewis bases,

like amines, aldehydes, alcohols, and thiols, taking advantage of their functional groups and the f-orbitals of the lanthanoids.^{34,132} Meanwhile, the substitution of high valence state ions (such as Ti⁴⁺) with trivalent Ln³⁺ ions will break the charge balance, which may be compensated by negative functional groups of dyes.¹³³ This is desirable to increase the adsorption capacity of organic dyes with photocatalysts due to the negative functional group, therefore lying good preparatory foundations for subsequent photocatalytic reactions.¹³⁴ For example, three-fold higher adsorption constants k were achieved in Pr^{3+}/Eu^{3+} ion-doped TiO₂ photocatalysts that were analyzed according to Langmuir's theory.¹³⁵ The Ln³⁺ sites can act as electron acceptors to trap the photogenerated electrons and minimize undesired electronhole recombination, where the trapped electrons are capable of combining with oxygen to form superoxide radicals for photooxidation of adsorbed organics. Especially, for Gd³⁺ ions with special half-filled electronic configuration of 7f-electrons, it is easier to transfer the trapped electrons to the interfacial adsorbed oxygen.130,136

It is well known that doping of Ln³⁺ ions will improve the photocatalytic performance of semiconductors by regulating the band energy and promoting the separation of carriers. Meanwhile, the loading concentration and type of Ln³⁺ dopant ions are also important parameters that affect the photocatalytic performance. High dopant dosage will increase the surface barrier of the semiconductors and decrease the space charge region. A large electric field is beneficial for promoting the separation of photoinduced carriers. But if the doping concentration is too high, the space charge region will become narrow that substantially reduces the penetration depth of light into semiconductors, and hence greatly increases the recombination possibilities of photoinduced carriers.^{130,137} Moreover, owing to the generally large ionic radii of Ln³⁺ ions (Table 4), for example, larger than Ti^{4+} (~68 pm), small and isolated clusters of lanthanide oxide will form and cover the surface of semiconductors at high doping concentration, which will aggregate as carrier capture centers, thus deactivating the photocatalytic activity.¹³⁸ Therefore, there must be an optimum concentration level of dopants to balance the thickness of the space charge layer substantially equivalent to the light penetration depth. For instance, La³⁺ ion-doped TiO₂ at a concentration of 0.5% was proved to possess the most efficient photodegradation rate for MO counterparts.¹³⁹ And Dy³⁺-doped ZnO nanoparticles with a doping concentration of 3% were found to endow the highest decolorization efficiency of C. I. Acid Red 17 under visible light irradiation.¹⁴⁰

In the activator ion (*i.e.*, Tm^{3+} , Er^{3+}) and sensitizer ion (*i.e.*, Yb^{3+} , Nd^{3+}) co-doped semiconductor system, the sensitizer will accept the NIR photons and transfer the energy to the nearest activator ions for emitting UV-visible photons by the upconversion process. As mentioned above, the upconverted photons favor the donation to the valence band of semiconductor followed by oxidative hole (h⁺) generation by extraction of a reductive electron, thus boosting the yield of photocarriers and improving the photocatalytic efficiency.^{141,142}

5. Outlooks and perspectives

In summary, the synthesis of upconversion materials with effective quantum yields and fabrication of upconversion-based semiconductor photocatalysts are considerable challenges. Despite the significant developments in this direction, there are several tough challenges associated with upconversion materials that limit their applications in basic sciences, including lower upconversion yields and confined excitation wavelengths. However, the reported highest upconversion efficiency is $3.0 \pm 0.3\%$ for Yb³⁺,Er³⁺ co-doped bulk NaYF₄ materials.¹⁴³ Nevertheless, the micro- or nanomaterials possess lower upconversion efficiency than bulk materials because of the synergistic effect of surface state defects induced by the small particle sizes and adhering water reducing the optical efficiency. Therefore, much effort is required to design and develop high-efficiency upconversion materials. To achieve this purpose, the first method is to select host materials with low phonon energy, which prefer to decrease multi-phonon relaxation between closely spaced energy levels, thus generating more effective upconversion. The host materials provide a surrounding crystalline field for Ln³⁺ ions but do not determine the emissions. The type of lanthanide dopant and doping concentration are vital for tailoring the upconverted luminescence. In addition, the SPR effect of noble metals is proved theoretically and experimentally to boost the upconversion luminescence by enhancing the absorption of the sensitizer through electric-field coupling, improving the radiative decay rate of the activator, and increasing energy transfer from the sensitizer to the activator. The intrinsic features of SPR are greatly beneficial to the photocatalytic performance of semiconductor photocatalysts by boosting the yield of photogenerated electron/hole pairs and promoting their effective separation. In addition, the SPR of noble metal nanoparticles could effectively promote the efficiency of upconversion materials. Interestingly, researchers have recently discovered localized surface plasmon resonances (LSPRs) in heavily doped semiconductor nanoparticles, quantum dots and amorphous two-dimensional nanomaterials.144-146 The occurrence of LSPR originates from collective oscillations of excess free carriers, the concentration of which $(10^{19}-10^{21} \text{ cm}^{-3})$ can be regulated by adjusting stoichiometric ratios, vacancy or doping concentrations, phase structures, etc.¹⁴⁷ Therefore, the LSPR-induced electric field in semiconductors provides a powerful strategy to enhance the upconversion efficiency. More importantly, the unique interactions based on the process of plasmonic energy transfer will further improve the NIR light-active photocatalytic performance.148,149 The marriage of upconversion with semiconductor LSPR could be a novel frontier in fundamental investigations and practical solar energy conversion applications.

Also, epitaxial growth of a shell with small lattice mismatch around a core can reduce non-radiative decay losses of the surface fluorescence, which would offer a useful way to improve the fluorescence intensity of upconversion crystals. In a coreshell structure, the dopant RE ions are incorporated into the interior core layer of particles. The shell layer can effectively suppress energy loss on the crystal surface, resulting in higher luminescence efficiency. These attractive features of core–shell structures make them suitable candidates for photocatalytic application.

The preceding discussion is concentrated entirely on Yb³⁺ ion (absorption cross section $\approx 10^{-20}$ cm²)-sensitized upconversion materials with the photo-absorption band centered at 980 nm. By paving attention to the drawbacks of confined excitation wavelength, innovative strategies that incorporate replaceable sensitizers have been proposed for tuning the excitation wavelengths of upconversion crystals. Nd³⁺ is ideal as an alternative sensitizer since it has a one order of magnitude larger absorption cross section ($\approx 10^{-19} \text{ cm}^2$) at $\approx 800 \text{ nm}$ than Yb³⁺ at \approx 980 nm.²⁷ Benefiting from the unique features of deep tissue penetration depth and low photothermal effect of ≈ 800 nm light (first NIR window), the Nd³⁺-sensitized upconversion nanoparticles have attracted considerable interdisciplinary attention in biosensing, bioimaging, drug delivery, therapy, and three-dimensional displays.¹⁵⁰ Unfortunately, there is no literature focusing on the Nd³⁺-sensitized upconversion materials as energy conversion elements that are composited with semiconductor photocatalysts, for further utilizing solar light centered at 800 nm. The emergence of Nd³⁺-sensitized upconversion nanoparticles would open other possible avenues of future works on upconversion semiconductor-based photocatalytic systems. We anticipate new breakthroughs in the construction of novel and highly efficient NIR-light-activated photocatalysts.

At present, almost all the considered NIR light-active photocatalytic systems are applied in organic dye degradation. It will be of significant importance to exploit potential effective photocatalysts for solving the challenges faced by the community of upconversion-based materials for their more efficient and wider applications to energy production and storage such as water splitting, CO_2 conversion, *etc.*^{151,152}

Author contributions

All authors contributed during the preparation of the manuscript. All authors have given approval to the final version of the manuscript.

Conflicts of interest

The authors declare that they have no competing interests.

Acknowledgements

This work was supported by the NSFC (51471121), Hubei Provincial Natural Science Foundation (2014CFB261), Basic Research Plan Program of Shenzhen City (JCYJ20160517104459444), Natural Science Foundation of Jiangsu Province (BK20160383) and Wuhan University.

References

- 1 H. Tong, S. Ouyang, Y. Bi, N. Umezawa, M. Oshikiri and J. Ye, *Adv. Mater.*, 2012, **24**, 229–251.
- 2 W. Wu, J. Changzhong and V. A. Roy, *Nanoscale*, 2015, 7, 38–58.
- 3 Z. Wu, S. Yang and W. Wu, Nanoscale, 2016, 8, 1237-1259.
- 4 J. Liu, Z. Wu, Q. Tian, W. Wu and X. Xiao, *CrystEngComm*, 2016, **18**, 6303–6326.
- 5 C. Chen, W. Ma and J. Zhao, *Chem. Soc. Rev.*, 2010, **39**, 4206-4219.
- 6 L. Sun, W. Wu, S. Yang, J. Zhou, M. Hong, X. Xiao, F. Ren and C. Jiang, ACS Appl. Mater. Interfaces, 2014, 6, 1113–1124.
- 7 Q. Tian, W. Wu, L. Sun, S. Yang, M. Lei, J. Zhou, Y. Liu, X. Xiao, F. Ren, C. Jiang and V. A. L. Roy, *ACS Appl. Mater. Interfaces*, 2014, 6, 13088–13097.
- 8 W. Wu, S. Zhang, X. Xiao, J. Zhou, F. Ren, L. Sun and C. Jiang, ACS Appl. Mater. Interfaces, 2012, 4, 3602–3609.
- 9 J. Liu, S. Yang, W. Wu, Q. Tian, S. Cui, Z. Dai, F. Ren, X. Xiao and C. Jiang, ACS Sustainable Chem. Eng., 2015, 3, 2975–2984.
- 10 A. Fujishima and K. Honda, Nature, 1972, 238, 37-38.
- 11 J. H. Carey, J. Lawrence and H. M. Tosine, *Bull. Environ. Contam. Toxicol.*, 1976, **16**, 697–701.
- 12 B. Kraeutler and A. J. Bard, *J. Am. Chem. Soc.*, 1978, **100**, 5985–5992.
- 13 G. L. Amy, L. Linsebigler and J. T. Yates, Jr., *Chem. Rev.*, 1995, **95**, 735–758.
- 14 S. G. Kumar and L. G. Devi, *J. Phys. Chem. A*, 2011, **115**, 13211–13241.
- 15 H. Park, Y. Park, W. Kim and W. Choi, J. Photochem. Photobiol., C, 2013, 15, 1-20.
- 16 Q. Tian, W. Wu, J. Liu, Z. Wu, W. Yao, J. Ding and C. Jiang, Dalton Trans., 2017, 46, 2770–2777.
- 17 W. Wu, L. Sun, S. Yang, J. Zhou, M. Hong, X. Xiao, F. Ren and C. Jiang, *ACS Appl. Mater. Interfaces*, 2014, **6**, 1113–1124.
- 18 H. Li, Z. Kang, Y. Liu and S.-T. Lee, J. Mater. Chem., 2012, 22, 24230–24253.
- 19 H. Li, R. Liu, S. Lian, Y. Liu, H. Huang and Z. Kang, Nanoscale, 2013, 5, 3289-3297.
- 20 J. Wu, Z. Zhang, B. Liu, Y. Fang, L. Wang and B. Dong, *Sol. RRL*, 2018, **2**, 1800039.
- 21 M. Hu, J. Chen, Z.-Y. Li, L. Au, G. V. Hartland, X. Li, M. Marquez and Y. Xia, *Chem. Soc. Rev.*, 2006, 35, 1084–1094.
- 22 A. Kubacka, M. Fernandez-Garcia and G. Colon, *Chem. Rev.*, 2012, **112**, 1555–1614.
- 23 Z. Zhang and J. T. Yates, Jr., Chem. Rev., 2012, 112, 5520-5551.
- 24 W. Yang, X. Li, D. Chi, H. Zhang and X. Liu, *Nanotechnology*, 2014, **25**, 482001.
- 25 X. Liu, H. Chu, J. Li, L. Niu, C. Li, H. Li, L. Pan and C. Q. Sun, *Catal. Sci. Technol.*, 2015, 5, 4727–4740.
- 26 X. Li, F. Zhang and D. Zhao, *Chem. Soc. Rev.*, 2015, 44, 1346–1378.
- 27 S. Han, R. Deng, X. Xie and X. Liu, Angew. Chem., Int. Ed., 2014, 53, 11702–11715.

- 28 S. Gai, C. Li, P. Yang and J. Lin, *Chem. Rev.*, 2014, **114**, 2343–2389.
- 29 X. Yu, M. Li, M. Xie, L. Chen, Y. Li and Q. Wang, *Nano Res.*, 2010, **3**, 51–60.
- 30 X. Yu, L. Chen, M. Li, M. Xie, L. Zhou, Y. Li and Q. Wang, *Adv. Mater.*, 2008, 20, 4118–4123.
- 31 S. V. Eliseeva and J. C. Bunzli, *Chem. Soc. Rev.*, 2010, **39**, 189–227.
- 32 F. Wang and X. Liu, J. Am. Chem. Soc., 2008, 130, 5642-5643.
- 33 A. Shalav, B. S. Richards and M. A. Green, Sol. Energy Mater. Sol. Cells, 2007, 91, 829–842.
- 34 S. Bingham and W. A. Daoud, *J. Mater. Chem.*, 2011, **21**, 2041–2050.
- 35 D. Chen, Y. Wang and M. Hong, *Nano Energy*, 2012, 1, 73–90.
- 36 N. Bloembergen, Phys. Rev. Lett., 1959, 2, 84-85.
- 37 X. Li, F. Zhang and D. Zhao, Nano Today, 2013, 8, 643-676.
- 38 J. S. Chivian, W. E. Case and D. D. Eden, *Appl. Phys. Lett.*, 1979, 35, 124–125.
- 39 G. Chen, H. Qiu, P. N. Prasad and X. Chen, *Chem. Rev.*, 2014, **114**, 5161–5214.
- 40 F. Wang, R. Deng, J. Wang, Q. Wang, Y. Han, H. Zhu, X. Chen and X. Liu, *Nat. Mater.*, 2011, **10**, 968–973.
- 41 Q. Su, S. Han, X. Xie, H. Zhu, H. Chen, C.-K. Chen, R.-S. Liu, X. Chen, F. Wang and X. Liu, *J. Am. Chem. Soc.*, 2012, **134**, 20849–20857.
- F. Wang, Y. Han, C. S. Lim, Y. Lu, J. Wang, J. Xu, H. Chen, C. Zhang, M. Hong and X. Liu, *Nature*, 2010, 463, 1061–1065.
- 43 K. K. Pukhov, F. Pelle and J. Heber, *Mol. Phys.*, 2003, 101, 1001–1006.
- 44 C. B. Layne, W. H. Lowdermilk and M. J. Weber, *Phys. Rev. B: Solid State*, 1977, 16, 10–20.
- 45 X. Liu, M. Li, X. Wang, F. Huang, Y. Ma, J. Zhang, L. Hu and D. Chen, *J. Lumin.*, 2014, **150**, 40–45.
- 46 O. Ravi, S. J. Dhoble, B. Ramesh, G. Devarajulu, C. M. Reddy, K. Linganna, G. R. Reddy and B. D. P. Raju, *J. Lumin.*, 2015, **164**, 154–159.
- 47 Y. Jia, T.-Y. Sun, J.-H. Wang, H. Huang, P. Li, X.-F. Yu and
 P. K. Chu, *CrystEngComm*, 2014, 16, 6141–6148.
- 48 W. Zheng, P. Huang, D. Tu, E. Ma, H. Zhu and X. Chen, *Chem. Soc. Rev.*, 2015, 44, 1379–1415.
- 49 G. S. Yi and G. M. Chow, Adv. Funct. Mater., 2006, 16, 2324–2329.
- 50 K. Chong, T. Hirai, T. Kawai, S. Hashimoto and N. Ohno, J. Lumin., 2007, 122–123, 149–151.
- 51 C. Li, Z. Quan, J. Yang, P. Yang and J. Lin, *Inorg. Chem.*, 2007, **46**, 6329–6337.
- 52 M. Haase and H. Schäfer, *Angew. Chem., Int. Ed.*, 2011, **50**, 5808–5829.
- 53 Y. Tang, W. Di, X. Zhai, R. Yang and W. Qin, *ACS Catal.*, 2013, **3**, 405–412.
- 54 W. Wang, M. Ding, C. Lu, Y. Ni and Z. Xu, *Appl. Catal., B*, 2014, **144**, 379–385.
- 55 W. Su, M. Zheng, L. Li, K. Wang, R. Qiao, Y. Zhong, Y. Hu and Z. Li, *J. Mater. Chem. A*, 2014, **2**, 13486–13491.

- 56 D.-X. Xu, Z.-W. Lian, M.-L. Fu, B. Yuan, J.-W. Shi and H.-J. Cui, *Appl. Catal.*, *B*, 2013, **142–143**, 377–386.
- 57 Y. Zhang and Z. Hong, Nanoscale, 2013, 5, 8930-8933.
- 58 X. Chen, Y. Zhu, D. Zhou, W. Xu, J. Zhu, G. Pan, Z. Yin, H. Wang, S. Cui and H. Song, *J. Mater. Chem. C*, 2017, 5, 2451–2458.
- 59 L. Bai, W. Jiang, C. Gao, S. Zhong, L. Zhao, Z. Li and S. Bai, *Nanoscale*, 2016, **8**, 19014–19024.
- 60 M. Chatti, V. N. Adusumalli, S. Ganguli and V. Mahalingam, *Dalton Trans.*, 2016, **45**, 12384–12392.
- 61 M. Z. Huang, B. Yuan, L. Dai and M. L. Fu, J. Colloid Interface Sci., 2015, 460, 264–272.
- 62 E. Cheng, W. Yin, S. Bai, R. Qiao, Y. Zhong and Z. Li, *Mater. Lett.*, 2015, **146**, 87–90.
- 63 C. Li, F. Wang, J. Zhu and J. C. Yu, *Appl. Catal.*, B, 2010, 100, 433–439.
- 64 M. Tou, Y. Mei, S. Bai, Z. Luo, Y. Zhang and Z. Li, *Nanoscale*, 2016, **8**, 553–562.
- 65 X. Guo, C. Chen, D. Zhang, C. P. Tripp, S. Yin and W. Qin, *RSC Adv.*, 2016, 6, 8127–8133.
- 66 X. Guo, W. Song, C. Chen, W. Di and W. Qin, *Phys. Chem. Chem. Phys.*, 2013, 15, 14681–14688.
- 67 M. Tou, Z. Luo, S. Bai, F. Liu, Q. Chai, S. Li and Z. Li, *Mater. Sci. Eng.*, C, 2017, 70, 1141–1148.
- 68 M. Li, Z. Zheng, Y. Zheng, C. Cui, C. Li and Z. Li, ACS Appl. Mater. Interfaces, 2017, 9, 2899–2905.
- 69 W. Wang, Y. Li, Z. Kang, F. Wang and J. C. Yu, *Appl. Catal.*, B, 2016, **182**, 184–192.
- 70 Q. Tian, W. Yao, Z. Wu, J. Liu, L. Liu, W. Wu and C. Jiang, J. Mater. Chem. A, 2017, 5, 23566–23576.
- 71 Z. Xu, M. Quintanilla, F. Vetrone, A. O. Govorov, M. Chaker and D. Ma, *Adv. Funct. Mater.*, 2015, 25, 2950–2960.
- 72 Q. Tian, W. Yao, W. Wu, J. Liu, Z. Wu, L. Liu,
 Z. Dai and C. Jiang, ACS Sustainable Chem. Eng., 2017, 5, 10889–10899.
- 73 Y. Ma, H. Liu, Z. Han, L. Yang and J. Liu, *J. Mater. Chem. A*, 2015, **3**, 14642–14650.
- 74 J. Wang, H. Huang, D. Zhang, M. Chen, Y. Zhang, X. Yu, L. Zhou and Q. Wang, *Nano Res.*, 2015, 8, 2548–2561.
- 75 W. Deng, F. Xie, H. T. Baltar and E. M. Goldys, *Phys. Chem. Chem. Phys.*, 2013, **15**, 15695–15708.
- 76 K. Aslan, I. Gryczynski, J. Malicka, E. Matveeva, J. R. Lakowicz and C. D. Geddes, *Curr. Opin. Biotechnol.*, 2005, 16, 55–62.
- 77 P. Padhye, S. Sadhu, M. Malik and P. Poddar, *RSC Adv.*, 2016, 6, 53504–53518.
- 78 J. Zhang, Y. Huang, L. Jin, F. Rosei, F. Vetrone and J. P. Claverie, ACS Appl. Mater. Interfaces, 2017, 9, 8142–8150.
- 79 D. Yin, L. Zhang, X. Cao, J. Tang, W. Huang, Y. Han, Y. Liu,
 T. Zhang and M. Wu, *RSC Adv.*, 2015, 5, 87251–87258.
- 80 N. Prakash, D. Thangaraju, R. Karthikeyan, M. Arivanandhan, Y. Shimura and Y. Hayakawa, *RSC Adv.*, 2016, 6, 80655–80665.
- 81 P. Huang, W. Zheng, S. Zhou, D. Tu, Z. Chen, H. Zhu, R. Li,
 E. Ma, M. Huang and X. Chen, *Angew. Chem., Int. Ed.*, 2014, 53, 1252–1257.
- 82 S. Zeng, J. Xiao, Q. Yang and J. Hao, *J. Mater. Chem.*, 2012, 22, 9870–9874.

- 83 L. Wang, M. Lan, Z. Liu, G. Qin, C. Wu, X. Wang, W. Qin,
 W. Huang and L. Huang, *J. Mater. Chem. C*, 2013, 1, 2485–2490.
- 84 G.-S. Yi and G.-M. Chow, Chem. Mater., 2007, 19, 341-343.
- 85 F. Zhang, C.-L. Zhang, H.-Y. Peng, H.-P. Cong and H.-S. Qian, Part. Part. Syst. Charact., 2016, 33, 248–253.
- 86 K. Wang, W. Qincheng, F. Wang, S. Bai, S. Li and Z. Li, *Inorg. Chem. Front.*, 2016, 3, 1190–1197.
- 87 H. Guan, Y. Sheng, Y. Song, C. Xu, X. Zhou, K. Zheng, Z. Shi and H. Zou, *J. Phys. Chem. C*, 2017, **121**, 23080–23095.
- 88 H. Li, Q. Huang, Y. Wang, K. Chen, J. Xie, Y. Pan, H. Su, X. Xie, L. Huang and W. Huang, *J. Mater. Chem. C*, 2017, 5, 6450–6456.
- 89 Y. Zhu, W. Xu, S. Cui, M. Liu, C. Lu, H. Song and D.-H. Kim, J. Mater. Chem. C, 2016, 4, 331–339.
- 90 J.-L. Lemyre and A. M. Ritcey, Chem. Mater., 2005, 17, 3040–3043.
- 91 G. Murali, R. K. Mishra, J. M. Lee, Y. C. Chae, J. Kim, Y. D. Suh, D.-k. Lim and S. H. Lee, *Cryst. Growth Des.*, 2017, 17, 3055–3061.
- 92 G. Murali, S. Kaur, Y. C. Chae, M. Ramesh, J. Kim, Y. D. Suh, D.-K. Lim and S. H. Lee, *RSC Adv.*, 2017, 7, 24255–24262.
- 93 H. Qiu, G. Chen, R. Fan, L. Yang, C. Liu, S. Hao, M. J. Sailor, H. Agren, C. Yang and P. N. Prasad, *Nanoscale*, 2014, 6, 753–757.
- 94 W. Qin, D. Zhang, D. Zhao, L. Wang and K. Zheng, *Chem. Commun.*, 2010, 46, 2304–2306.
- 95 Q.-L. Ye, X. Yang, C. Li and Z. Li, *Mater. Lett.*, 2013, **106**, 238–241.
- 96 S. Huang, Z. Lou, Z. Qi, N. Zhu and H. Yuan, *Appl. Catal.*, *B*, 2015, 168–169, 313–321.
- 97 N. Niu, F. He, S. Gai, C. Li, X. Zhang, S. Huang and P. Yang, J. Mater. Chem., 2012, 22, 21613–21623.
- 98 N.-N. Dong, M. Pedroni, F. Piccinelli, G. Conti, A. Sbarbati, J. E. Ramírez-Hernández, L. M. Maestro, M. C. Iglesias-dela Cruz, F. Sanz-Rodriguez, A. Juarranz, F. Chen, F. Vetrone, J. A. Capobianco, J. G. Solé, M. Bettinelli, D. Jaque and A. Speghini, ACS Nano, 2011, 5, 8665–8671.
- 99 I. Masahiko, G. Christelle, G. Yannick, L. Kheirreddine, F. Tsuguo and B. Georges, *J. Phys.: Condens. Matter*, 2004, 16, 1501–1521.
- 100 S. Huang, L. Gu, C. Miao, Z. Lou, N. Zhu, H. Yuan and A. Shan, J. Mater. Chem. A, 2013, 1, 7874–7879.
- 101 S. Huang, N. Zhu, Z. Lou, L. Gu, C. Miao, H. Yuan and A. Shan, *Nanoscale*, 2014, **6**, 1362–1368.
- 102 S. Huang, Z. Lou, A. Shan, N. Zhu, K. Feng and H. Yuan, J. Mater. Chem. A, 2014, 2, 16165–16174.
- 103 S. Huang, Z. Lou, N. Zhu and H. Yuan, *Catal. Commun.*, 2015, **61**, 6–10.
- 104 S. Huang, H. Wang, N. Zhu, Z. Lou, L. Li, A. Shan and H. Yuan, *Appl. Catal.*, B, 2016, 181, 456–464.
- 105 X. Liu, W. Di and W. Qin, *Appl. Catal.*, *B*, 2017, **205**, 158–164.
- 106 M. Haase and H. Schafer, *Angew. Chem., Int. Ed.*, 2011, **50**, 5808–5829.

- 107 D. L. Dexter, J. Chem. Phys., 1953, 21, 836-850.
- 108 J. Tang, Y. Zhang, J. Gou, Z. Ma, G. Li, Y. Man and N. Cheng, *J. Alloys Compd.*, 2018, 740, 229–236.
- 109 S. Som and S. K. Sharma, J. Phys. D: Appl. Phys., 2012, 45, 415102.
- 110 E. W. Barrera, M. C. Pujol, F. Diaz, S. B. Choi, F. Rotermund, K. H. Park, M. S. Jeong and C. Cascales, *Nanotechnology*, 2011, 22, 075205.
- 111 Z. E. Liu, J. Wang, Y. Li, X. Hu, J. Yin, Y. Peng, Z. Li, Y. Li, B. Li and Q. Yuan, ACS Appl. Mater. Interfaces, 2015, 7, 19416–19423.
- 112 J. Leng, J. Chen, D. Wang, J.-X. Wang, Y. Pu and J.-F. Chen, *Ind. Eng. Chem. Res.*, 2017, **56**, 7977–7983.
- 113 L. Chen, F. He, C. Shi, C. He, E. Liu, L. Ma and N. Zhao, *J. Alloys Compd.*, 2018, **740**, 719–726.
- 114 T. Li, S. Liu, H. Zhang, E. Wang, L. Song and P. Wang, J. Mater. Sci., 2010, 46, 2882–2886.
- 115 H. Zhou, S. Xu, D. Zhang, S. Chen and J. Deng, *Nanoscale*, 2017, **9**, 3196–3205.
- D. Yin, F. Zhao, L. Zhang, X. Zhang, Y. Liu, T. Zhang, C. Wu,
 D. Chen and Z. Chen, *RSC Adv.*, 2016, 6, 103795–103802.
- 117 Y. Liu, D. Wang, J. Shi, Q. Peng and Y. Li, *Angew. Chem.*, *Int. Ed.*, 2013, **52**, 4366–4369.
- 118 Z. Wu, S. Yang and W. Wu, CrystEngComm, 2016, 18, 2222–2238.
- 119 S. Yang, B. Zhou, Z. Ding, H. Zheng, L. Huang, J. Pan, W. Wu and H. Zhang, *J. Power Sources*, 2015, 286, 124–129.
- 120 S. Obregon and G. Colon, Chem. Commun., 2012, 48, 7865–7867.
- 121 L. Shan and Y. Liu, J. Mol. Catal. A: Chem., 2016, 416, 1-9.
- 122 S. Ganguli, C. Hazra, M. Chatti, T. Samanta and V. Mahalingam, *Langmuir*, 2016, **32**, 247–253.
- 123 K. Fu, J. Huang, N. Yao, X. Xu and M. Wei, *Ind. Eng. Chem. Res.*, 2015, **54**, 659–665.
- 124 C. Regmi, Y. K. Kshetri, S. K. Ray, R. P. Pandey and S. W. Lee, *Appl. Surf. Sci.*, 2017, **392**, 61–70.
- 125 S. Obregon, S. W. Lee and G. Colon, *Dalton Trans.*, 2014, 43, 311–316.
- 126 P. M. Martins, V. Gomez, A. C. Lopes, C. J. Tavares,
 G. Botelho, S. Irusta and S. Lanceros-Mendez, *J. Phys. Chem. C*, 2014, **118**, 27944–27953.
- 127 S. Huang, Y. Feng, L. Han, W. Fan, X. Zhao, Z. Lou, Z. Qi, B. Yu and N. Zhu, *RSC Adv.*, 2014, 4, 61679–61686.
- 128 M. Samadi, M. Zirak, A. Naseri, E. Khorashadizade and A. Z. Moshfegh, *Thin Solid Films*, 2016, **605**, 2–19.
- 129 C. Wang, Y. Ao, P. Wang, J. Hou and J. Qian, *Appl. Surf. Sci.*, 2010, **257**, 227–231.

- 130 A.-W. Xu, Y. Gao and H.-Q. Liu, J. Catal., 2002, 207, 151–157.
- 131 M. Wang, Y. Che, C. Niu, M. Dang and D. Dong, *J. Hazard. Mater.*, 2013, **262**, 447–455.
- 132 K. T. Ranjit, I. Willner, S. H. Bossmann and A. M. Braun, *Environ. Sci. Technol.*, 2001, **35**, 1544–1549.
- 133 A. S. Weber, A. M. Grady and R. T. Koodali, *Catal. Sci. Technol.*, 2012, **2**, 683.
- 134 Q. Tian, W. Wu, S. Yang, J. Liu, W. Yao, F. Ren and C. Jiang, *Nanoscale Res. Lett.*, 2017, **12**, 221.
- 135 K. T. Ranjit, H. Cohen, I. Willner, S. Bossmann and A. M. Braun, *J. Mater. Sci.*, 1999, **34**, 5273–5280.
- 136 W. Choi, A. Termin and M. R. Hoffmann, *J. Phys. Chem.*, 1994, **98**, 13669–13679.
- 137 K. M. Parida and N. Sahu, J. Mol. Catal. A: Chem., 2008, 287, 151–158.
- 138 J. Choi, H. Park and M. R. Hoffmann, *J. Phys. Chem. C*, 2010, **114**, 783–792.
- 139 D. Xu, L. Feng and A. Lei, *J. Colloid Interface Sci.*, 2009, **329**, 395–403.
- 140 A. Khataee, R. Darvishi Cheshmeh Soltani, Y. Hanifehpour, M. Safarpour, H. Gholipour Ranjbar and S. W. Joo, *Ind. Eng. Chem. Res.*, 2014, 53, 1924–1932.
- 141 K. Feng, S. Huang, Z. Lou, N. Zhu and H. Yuan, *Dalton Trans.*, 2015, **44**, 13681–13687.
- 142 J. Méndez-Ramos, P. Acosta-Mora, J. C. Ruiz-Morales, M. Sierra, A. Redondas, E. Ruggiero, L. Salassa, M. E. Borges and P. Esparza, *Opt. Mater.*, 2015, 41, 98–103.
- 143 J. C. Boyer and F. C. van Veggel, *Nanoscale*, 2010, 2, 1417–1419.
- 144 X. Liu and M. T. Swihart, Chem. Soc. Rev., 2014, 43, 3908–3920.
- 145 W. Liu, Q. Xu, W. Cui, C. Zhu and Y. Qi, Angew. Chem., Int. Ed., 2017, 56, 1600–1604.
- 146 K. Manthiram and A. P. Alivisatos, J. Am. Chem. Soc., 2012, 134, 3995–3998.
- 147 A. Agrawal, S. H. Cho, O. Zandi, S. Ghosh, R. W. Johns and D. J. Milliron, *Chem. Rev.*, 2018, **118**, 3121–3207.
- 148 D. Zhou, D. Liu, W. Xu, Z. Yin, X. Chen, P. Zhou, S. Cui,
 Z. Chen and H. Song, ACS Nano, 2016, 10, 5169–5179.
- 149 Z. Zhang, Y. Liu, Y. Fang, B. Cao, J. Huang, K. Liu and B. Dong, *Adv. Sci.*, 2018, 1800748.
- 150 X. Xie, Z. Li, Y. Zhang, S. Guo, A. I. Pendharkar, M. Lu, L. Huang, W. Huang and G. Han, *Small*, 2017, **13**, 1602843.
- 151 Y. Chen, C. Lu, L. Tang, S. Wei, Y. Song and J. Wang, *Sol. Energy Mater. Sol. Cells*, 2016, **149**, 128–136.
- 152 J. Hong, Z. S. Hong, X. Cheng, C. W. Bo, W. J. Chen, L. X. Feng and Q. J. Rong, *Adv. Energy Mater.*, 2015, **5**, 1401041.