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# A comprehensive insight into deep-level defect engineering in antimony chalcogenide solar cells†

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Antimony chalcogenides ( $\text{Sb}_2\text{X}_3$ ,  $\text{X} = \text{S}$  and  $\text{Se}$ ) are intriguing materials for the fabrication of next-generation, flexible/wearable, lightweight, and tandem photovoltaic (PV) devices. Recently, the power conversion efficiency (PCE) of 10.75% and 11.66% has been demonstrated in (single junction)  $\text{Sb}_2\text{X}_3$  and  $\text{Sb}_2\text{X}_3/\text{Si}$  (tandem) solar cells, respectively. However, the inevitable presence ( $>10^{16} \text{ cm}^{-3}$ ) of deep-level defects (especially  $\text{Sb}_\text{S}$  and  $\text{Sb}_\text{Se}$  antisites) induces Fermi-level ( $E_\text{F}$ ) pinning, accelerates Shockley–Read–Hall (SRH) recombination, and shortens the carrier lifetime. Unambiguously, these defects result in sluggish charge transport and high open-circuit voltage ( $V_\text{OC}$ ) deficits in the corresponding  $\text{Sb}_2\text{X}_3$  solar cells. Therefore, a comprehensive understanding of the deep-level defects and their passivation strategies can be instrumental in reducing the  $V_\text{OC}$  deficits and boosting the PCE values. In this regard, the present review highlights the expanding toolbox of defect-engineering strategies for  $\text{Sb}_2\text{X}_3$  films, laying a solid foundation for improving the PCE of  $\text{Sb}_2\text{X}_3$  solar cells.

## 1. Introduction

The unique quasi-one-dimensional (Q-1D) crystal structure, economical fabrication, and power conversion efficiency (PCE) evolution have stimulated significant scientific and industrial interest towards antimony chalcogenide ( $\text{Sb}_2\text{X}_3$ ,  $\text{X} = \text{S}, \text{Se}$ )

photovoltaics (PV). Furthermore,  $\text{Sb}_2\text{X}_3$  has the advantages of earth-abundance, low eco-toxicity, high absorption coefficient ( $>10^5 \text{ cm}^{-1}$ ), low Urbach energy ( $\sim 30 \text{ meV}$ ), low exciton binding energy, readily tunable bandgap ( $E_\text{g} \sim 1.1\text{--}1.7 \text{ eV}$ , depending on the  $\text{Se/S}$  ratio), superior physicochemical stability, balanced (ambipolar) charge transport, and ultra-flexibility, satisfying most of the requirements of an ideal PV material.<sup>1–7</sup> Due to the rapid progress in film deposition strategies and composition and device engineering, PCE values of 8.00%,<sup>8</sup> 10.57%,<sup>9</sup> and 10.75%<sup>10</sup> have been reported for  $\text{Sb}_2\text{S}_3$ ,  $\text{Sb}_2\text{Se}_3$ ,

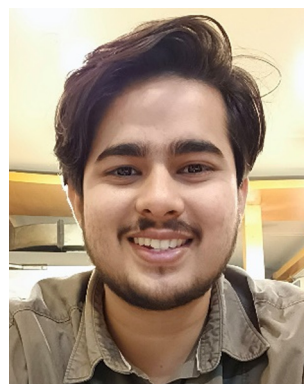
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and  $\text{Sb}_2(\text{S,Se})_3$  solar cells, respectively. Recently, a PCE of 11.66% was demonstrated in an  $\text{Sb}_2(\text{S,Se})_3/\text{Si}$  tandem solar cell.<sup>11</sup>  $\text{Sb}_2\text{S}_3$  ( $E_g \sim 1.7$  eV) and  $\text{Sb}_2\text{Se}_3$  ( $E_g \sim 1.1$  eV) are promising materials for the fabrication of the top-cell and bottom-cell in tandem solar cells, respectively.<sup>12–15</sup> In addition to generic power production (terawatt levels),  $\text{Sb}_2\text{X}_3$  solar cells are well suited for integration in futuristic, lightweight, flexible, and wearable electronic devices.<sup>15–18</sup>

A comprehensive understanding of the defect-formation mechanism and control of defects *via* passivation strategies has been a cornerstone in the successful technological deployment of established semiconductor technologies. However, this profound understanding is limited in the case of  $\text{Sb}_2\text{X}_3$  solar cells, which severely suffer from large open circuit voltage ( $V_{\text{OC}}$ ) deficits ( $>0.6$  eV) owing to their high density of (electrically active) deep defects. These defects are localized in real space, dictate the carrier transport, and restrict the PCE of the corresponding devices from reaching the theoretical limit ( $\sim 30\%$ ).<sup>18,19</sup>

To date, various reviews have covered different critical aspects of  $\text{Sb}_2\text{X}_3$  PV, with particular emphasis on presenting either a generic overview (of anisotropy, crystal-structure, band-structure, photophysical properties, deposition techniques, and recent advances),<sup>18,20–26</sup> doping,<sup>27</sup> device and interfacial engineering,<sup>19,28,29</sup> or commercialization prospects,<sup>30</sup> while the defect engineering aspect has largely remained unexplored. Recently, Wijesinghe *et al.*<sup>31</sup> published a review emphasizing defect engineering in  $\text{Sb}_2\text{Se}_3$  solar cells. However, to the best of our knowledge, the present review is unique, highlighting the origin of deep-level defects in  $\text{Sb}_2\text{X}_3$  (covering  $\text{Sb}_2\text{S}_3$ ,  $\text{Sb}_2\text{Se}_3$ , and  $\text{Sb}_2(\text{S,Se})_3$ ), their passivation strategies, and strategies for boosting their PCE beyond the state-of-the-art. This work is organized into five main sections. Firstly, we briefly elucidate the charge transport properties of  $\text{Sb}_2\text{X}_3$  materials. Next, we summarize the defect engineering investigations, discussing the origin of deep-level defects and their influence on the performance of  $\text{Sb}_2\text{X}_3$  solar cells. Subsequently, we present the engineering strategies for defect passivation and PCE



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improvement in  $\text{Sb}_2\text{X}_3$  solar cells. Finally, we present a summary and technological outlook, outlining potential challenges and future research directions. We anticipate that this work will inspire the development of novel strategies for defect regulation in  $\text{Sb}_2\text{X}_3$  films and solar cells, laying a solid foundation for improving their PCE.

## 2. Charge carrier dynamics in $\text{Sb}_2\text{X}_3$

Although the anisotropic structural and optoelectronic properties of  $\text{Sb}_2\text{X}_3$  have been investigated comprehensively, there is no consensus on its charge-carrier dynamics, which critically affects the performance of  $\text{Sb}_2\text{X}_3$  solar cells. Specifically, the nature of the carrier transport remains ambiguous, regarding if it is band-like or thermally activated hopping. Moon's group performed several time-resolved studies to understand the carrier dynamics involved in  $\text{Sb}_2\text{Se}_3$  nanostructures.<sup>32</sup> Time-resolved terahertz (tr-THz) spectroscopy revealed that the photoexcited carriers in  $\text{Sb}_2\text{Se}_3$  undergo effective mobility loss within 30 ps owing to the carrier localization, while time-resolved photoluminescence (TRPL) suggested the recombination of photogenerated carriers in a lengthier (ns) time scale. Furthermore, the carrier localization was attributed to the soft/deformable and anharmonic crystal structure, strong electron–phonon coupling and carrier interaction with crystal distortions in  $\text{Sb}_2\text{Se}_3$ .<sup>33</sup> Zhu's group conducted an in-depth investigation on the carrier dynamics in  $\text{Sb}_2\text{S}_3$  single crystals and thin films *via* transient absorption spectroscopy (TAS).<sup>34</sup> They concluded that the red-shifted PL (Stokes shift of 0.6 eV), picosecond carrier trapping, polarized trap emission, and hopping transport were strongly correlated with the intrinsic self-trapping of photogenerated carriers by lattice deformations in  $\text{Sb}_2\text{S}_3$ . The self-trapping model rationalizes the high  $V_{\text{OC}}$ -deficits in  $\text{Sb}_2\text{S}_3$  solar cells, pinning the upper limit of the  $V_{\text{OC}}$  at 0.8 V and PCE at 16%. Grad *et al.*<sup>35</sup> performed time-resolved two-photon photoemission experiments to study the photoexcited carrier dynamics in single-crystal  $\text{Sb}_2\text{S}_3$ . They proposed a model of self-trapping of free charge carriers by optical phonons and the formation of intrinsic traps. Consistent with the above-mentioned hypothesis, Tao *et al.* reported a Stokes-shifted ( $\approx 0.5$  eV) broadband emission and (barrierless) ultrafast ( $\approx 20$  ps) carrier self-trapping in  $\text{Sb}_2\text{Se}_3$  and  $\text{Sb}_2(\text{S},\text{Se})_3$ , regardless of their crystallinity and stoichiometry.<sup>36</sup>

To counter the proposed intrinsic self-trapping model, various studies revealed that carriers are localized by extrinsic (bulk and interface) defects. Yang's group investigated the anisotropic photoconductivity in  $\text{Sb}_2\text{Se}_3$  single crystals *via* tr-THz spectroscopy.<sup>37</sup> The temporal evolution of photoconductivity revealed that electron trapping resulted in a variation in photoconductivity anisotropy, which was accompanied by a decline in the photoconductivity magnitude, while electron–hole recombination only reduced the magnitude but did not affect the anisotropy. Photoconductivity decay and the accompanying anisotropy reversal revealed that at low pump intensities (relevant to the solar insolation), electron trapping outpaced charge recombination and dictated the photoconductivity dynamics. Leng's group investigated the charge carrier

transfer and trap state localization process in  $\text{Sb}_2\text{Se}_3$  films and heterostructures ( $\text{CdS}/\text{Sb}_2\text{Se}_3$  and  $\text{SnO}_2/\text{Sb}_2\text{Se}_3$ ) *via* visible-pump terahertz-probe spectroscopy (VPTPS).<sup>38</sup> In the case of the  $\text{CdS}/\text{Sb}_2\text{Se}_3$  heterostructure, the band edge electron transfer time ( $\sim 50$  ps) was almost double that for the trapping process ( $\sim 30$  ps), revealing that the intrinsic trap-assisted carrier recombination dominates the loss mechanisms.

Walsh's group investigated the tendency of polaron-trapping and its effect on the charge-carrier transport in  $\text{Sb}_2\text{X}_3$  using density functional theory (DFT) and Boltzmann transport calculations.<sup>6</sup> The modeling of electron and hole polarons in  $\text{Sb}_2\text{X}_3$  indicated the intrinsic formation of large polarons, in contrast to prior suggestions of small polarons (*i.e.*, self-trapped carriers). As illustrated in Fig. 1(a), the isotopically averaged mobilities (for both electrons and holes) exceeded  $10 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  at room temperature and declined with an increase in temperature, confirming the band-like transport in  $\text{Sb}_2\text{S}_3$  and  $\text{Sb}_2\text{Se}_3$ . It was also revealed that the intrinsic mobility is limited by scattering from polar optical phonons at low and moderate defect densities. In the case of high charge defect concentrations ( $> 10^{18} \text{ cm}^{-3}$ ), impurity scattering dominated. Yang's group employed Tr-THz and TAS to investigate both the free and trapped carrier dynamics in  $\text{Sb}_2\text{Se}_3$  films.<sup>39</sup> The results revealed that the trapped carriers remain mobile and reach charge-collecting interfaces prior to recombination. Interestingly, in addition to free carriers, trapped carriers were also found to contribute to the photocurrent in  $\text{Sb}_2\text{Se}_3$  solar cells, as schematically depicted in Fig. 1(d).

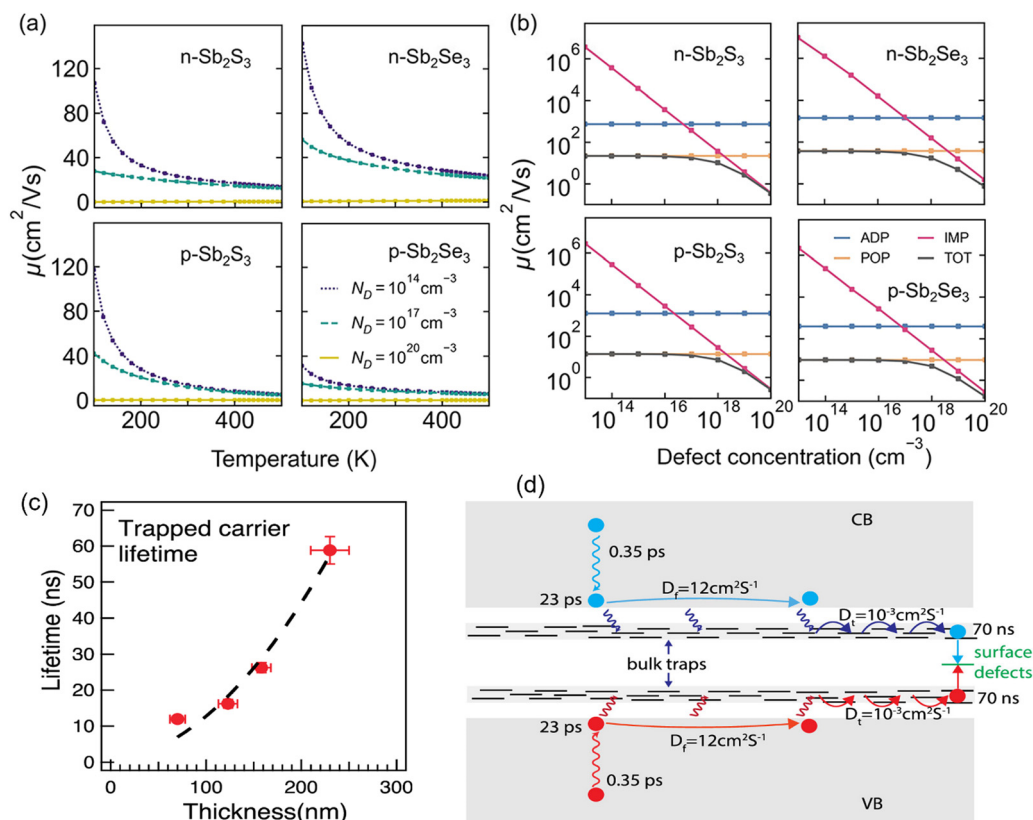
Although  $\text{Sb}_2\text{S}_3$  and  $\text{Sb}_2\text{Se}_3$  are isomorphous and exhibit similar material properties, they exhibit minor differences in their electronic properties. Generally,  $\text{Sb}_2\text{S}_3$  exhibits intrinsic to weakly n-type conductivity, in contrast to weakly p-type conductivity in  $\text{Sb}_2\text{Se}_3$ . Strong electron–phonon coupling-induced (intrinsic) self-trapping is prominent in  $\text{Sb}_2\text{S}_3$ , limiting the  $V_{\text{OC}}$  and FF values in  $\text{Sb}_2\text{S}_3$  solar cells. However, it is conceivable that there is ambiguity regarding the self-trapping in  $\text{Sb}_2\text{Se}_3$  and  $\text{Sb}_2(\text{S},\text{Se})_3$  and the extent to which it impedes carrier transport in these materials.<sup>35,40</sup> The suppressed or non-existent self-trapping in  $\text{Sb}_2\text{Se}_3$  and  $\text{Sb}_2(\text{S},\text{Se})_3$  facilitates improved charge transport, which is reflected in their better device performance than their  $\text{Sb}_2\text{S}_3$  counterparts (Table S1, ESI†). Notably, the upper limit of the  $V_{\text{OC}}$  ( $\sim 0.8$  eV) projected by the self-trapping model for  $\text{Sb}_2\text{S}_3$  solar cells was experimentally reported by Maity *et al.*<sup>41</sup> and Peng *et al.*<sup>42</sup> In this regard, any further improvement in  $V_{\text{OC}}$  will provide evidence for the validity and efficacy of this model. Engineering strategies to suppress electron–phonon coupling (such as stiffening the elastic properties by strain) can counter the ultrafast self-trapping in  $\text{Sb}_2\text{X}_3$  films.

## 3. Defects in $\text{Sb}_2\text{X}_3$ materials and their implication in the performance of solar cells

Computational and experimental studies on  $\text{Sb}_2\text{X}_3$  thin films revealed the highly complicated and sensitive (to film deposition







**Fig. 1** Calculated average mobility ( $\mu$ ) values for electrons and holes in  $\text{Sb}_2\text{S}_3$  and  $\text{Sb}_2\text{Se}_3$  as a function of (a) temperature and (b) defect concentration ( $N_D$ ). Calculated total (TOT) and component mobilities as a function of bulk defect concentration at 300 K. ADP, acoustic deformation potential; POP, polar optical phonon; and IMP, ionized impurity. Adapted under the guidelines of the Creative Commons CC BY license from ref. 6 Copyright 2022, the American Chemical Society. (c) Dependence of trapped carrier lifetime ( $\tau_t$ ) on film thickness. (d) Schematic illustration of the free and trapped charge carrier transport. The free carrier cooling (0.35 ps) and subsequent trapping (23 ps) time constants are shown. The trapped carriers hop among the bulk trap sites with a narrow energy distribution. Surface recombination leads to the depopulation of the trapped carriers. Adapted with permission from ref. 39 Copyright 2019, the American Chemical Society.

conditions) nature of defects. The following defects in  $\text{Sb}_2\text{X}_3$  films dictate the performance of  $\text{Sb}_2\text{X}_3$  solar cells: (1) intrinsic zero-dimensional (0D) defects, including vacancies, interstitials and antisites; (2) 0D defects caused by impurities (doped or permeable ions); and (3) 2D defects (grain boundaries, interfaces, and surface defects). In addition, defects in the transport layer and poor contact between the transport layer and electrodes also affect charge extraction, resulting in a higher series resistance ( $R_s$ ). Moreover, severe defects such as microcracks and scratches on the surface cause a rough surface topography, lowering the FF in devices. However, in this review, we focus on the point (0D) defects in the absorber layers. Point defects can be classified as deep or shallow, depending on whether the thermal activation energies for the electrons ( $E_{th}$ ) are higher or lower than  $k_B T$  ( $k_B$  is the Boltzmann constant and  $T$  is the absolute temperature). Point defects with deep-energy levels are prone to tightly trapping free electrons and holes, resulting in defect-mediated Shockley–Read–Hall (SRH) recombination, and thus curtaining the average lifetime ( $\tau_{eff}$ ) of the charge carriers and  $V_{OC}$  of PV devices. In contrast, shallow defects are benign and instrumental tools for regulating the carrier density and electrical conductivity. Solar energy conversion requires the

photoinduced generation of long-lived charge carriers, which necessitate their separation and collection at the respective electrodes to generate a discernible  $V_{OC}$  and  $J_{SC}$ . Thus, mitigating all the possible charge carrier loss channels is indispensable for attaining excellent PCE values.<sup>38</sup>

Suppressing the rate of trap-assisted SRH recombination enables devices to achieve higher  $V_{OC}$  values before the  $J_{SC}$  is canceled completely by recombination. The  $V_{OC}$  can be formulated in terms of non-radiative voltage loss ( $\Delta V_{OC}^{NR}$ ) from the maximum achievable voltage ( $\Delta V_{OC}^{NR}$ ),<sup>43</sup> as follows:

$$V_{OC} = V_{OC}^{\max} - \Delta V_{OC}^{NR} \quad (1)$$

$$= V_{OC}^{\max} - \frac{k_B T}{e} \log(\text{EQE}_{EL}) \quad (2)$$

where  $\text{EQE}_{EL}$  is the external electroluminescence quantum efficiency,  $T$  is the absolute temperature, and  $e$  is the elementary charge. Naturally, reducing the defect density and the non-radiative recombination in solar cells is the key to reducing  $\text{EQE}_{EL}$ .

$V_{OC}$  and FF are closely related to the defect-assisted recombination losses caused by bulk traps and interface defects (improper energy-level alignment, lattice-mismatching, and





dangling bonds). The FF of a solar cell is defined as the product of its current and voltage values at the maximum power output relative to the product of  $J_{SC}$  and  $V_{OC}$ . Therefore, the FF is strongly affected not only by the series and shunt resistance but also by carrier recombination (at the maximum power point conditions). The latter effect is well illustrated by the semiempirical expression, showing the dependency of FF on the diode ideality factor ( $\eta$ ) and  $V_{OC}$ ,<sup>44</sup> as follows:

$$FF = \frac{v_m}{v_m + 1} \frac{v_{oc} - \ln(v_m + 1)}{v_{oc}(1 - e^{-v_{oc}})} \quad (3)$$

$$v_{oc} = \frac{V_{OC}}{\eta k_B T} \quad (4)$$

$$v_m = v_{oc} - \ln(v_{oc} + 1 - \ln v_{oc}) \quad (5)$$

The  $V_{OC}$  of a solar cell is directly related to the splitting of the (electron and hole) quasi-Fermi levels in its semiconducting absorber (and thus excess-charge carrier densities,  $\Delta n$  and  $\Delta p$ , excited under specific illumination conditions, usually at one sun). In turn,  $\Delta n$  ( $\approx \Delta p$ ) directly depends on the effective free-carrier lifetime ( $\tau_{eff}$ ) through  $\tau_{eff} = \Delta n / u_{eff}$ , where  $u_{eff}$  is the effective recombination rate.<sup>45,46</sup> Notably,  $\tau_{eff}$  is often used to assess the electronic quality of semiconductors and is usually extracted experimentally by TRPL, TAS or transient photovoltage (TPV) studies. Given that the recombination rates are additive in nature,  $\tau_{eff}$  depends on the bulk ( $\tau_{bulk}$ ) and surface ( $\tau_{surface}$ ) recombination lifetimes, as follows:<sup>47</sup>

$$\frac{1}{\tau_{eff}} = \frac{1}{\tau_{bulk}} + \frac{1}{\tau_{surface}} \quad (6)$$

The bulk carrier lifetime ( $\tau_{bulk}$ ) depends on the defect-induced (SRH) recombination lifetime ( $\tau_{SRH}$ ), band-to-band radiative recombination lifetime ( $\tau_{rad}$ ) and Auger recombination lifetime ( $\tau_{Auger}$ ), as follows:<sup>48</sup>

$$\frac{1}{\tau_{bulk}} = \frac{1}{\tau_{rad}} + \frac{1}{\tau_{SRH}} + \frac{1}{\tau_{Auger}} \quad (7)$$

In  $Sb_2X_3$  solar cells, the rate of interface recombination is faster ( $\tau_{surface} \sim 0.1 - 1$  ns) than that at the bulk ( $\tau_{bulk} \sim 5 - 60$  ns). However, biexponential fitting of the decay curve (TRPL, TAS, and TPV) to get individual values of longer and shorter lifetimes and correlate the values with the  $\tau_{bulk}$  and  $\tau_{surface}$  is challenging.<sup>18,31</sup>

To intuitively assess the influence of individual types of deep-level defects on the device performance, the  $(\sigma N_T)^{-1}$  values are evaluated and compared. According to the trap-assisted-SRH recombination model, the  $\tau_{SRH}$  value associated with a specific defect depends on the carrier thermal velocity ( $v_{th}$ ), capture cross section ( $\sigma$ ), and trap density ( $N_T$ ), as follows:<sup>49,50</sup>

$$\tau_{SRH} = \frac{1}{v_{th} \sigma N_T} \quad (8)$$

Deep-level transient spectroscopy (DLTS), thermally stimulated current (TSC), and admittance spectroscopy are the most frequently used techniques for characterizing defects in

semiconductors.<sup>51</sup> In general, TSC is limited by its low sensitivity, shallow detection depth, and complicated operation.<sup>52,53</sup> Alternatively, admittance spectroscopy can analyze faster emission processes, and thus is effective in probing shallow defects and dopant energy levels. However, although this technique facilitates the detection of single trap species (majority-carrier trap only), its sensitivity decreases for deeper trap levels.<sup>54,55</sup> Compared with the two above-mentioned methods, DLTS is a powerful tool to probe (bulk) deep-level defects (located  $>0.3$  eV from the band-edges) in semiconductors with higher sensitivity, wider observation range of trap depth, and more versatile and convenient operation and analysis. It employs the transient capacitance of the p-n junction (during emission of carriers from the traps) at different temperatures as the probe to monitor the changes in the charge state of deep-level traps. This technique has been widely explored for the characterization of defects in mature PV technologies (e.g., Si,<sup>56</sup> GaAs,<sup>57,58</sup> CIGS,<sup>59</sup> and CdTe<sup>60,61</sup>), and emerging perovskite solar cells.<sup>62-64</sup>

However, conventional DLTS has the major limitation of unreliable minority-carrier trap detection. In this regard, a novel optical-DLTS (O-DLTS) technique has emerged, which can effectively probe the majority and minority-carrier traps simultaneously. Deconvolution of DLTS spectra reveals distinctive (positive) peaks and (negative) valleys, which correspond to the majority and minority carrier traps, respectively. As a standard procedure, the defect activation energy ( $E_a$ ) and capture cross-section ( $\sigma$ ) are calculated using the Arrhenius plot line-fitting (APL) method. The Shockley-Read-Hall emission of carriers from the trap states is a thermal activation process. The emission rate ( $v$ ) of carriers from the defect is essentially the Arrhenius equation, as follows:

$$v = v_0 \exp\left(-\frac{E_a}{k_B T}\right) \quad (9)$$

where  $v_0$  is the attempt-to-escape frequency (enveloping information of the carrier capture cross section),  $k_B$  is the Boltzmann constant, and  $T$  is the absolute temperature. In the case where  $v_0$  and  $E_a$  are not dependent on the temperature, the Arrhenius plot, i.e.,  $\log(v)$  vs.  $T^{-1}$  converges to a straight line. The values of  $v_0$  and  $E_a$  can be calculated from the Y-intercept, and slope of the plot respectively. In the case of a non-Arrhenius process,  $v_0$  and  $E_a$  are dependent on the temperature, and consequently the Arrhenius plot is curved. An in-depth understanding of the DLTS instrumentation, operation, and data interpretation can be obtained in the literature.<sup>65-67</sup> Notably, DLTS measurements reveal a lower number of defect types than that predicted by theoretical calculations. Fundamentally, this inconsistency arises because all possible defect types are considered in theoretical calculations, but most defects are unstable (with high formation energies), and thus cannot be probed experimentally. Moreover, DLTS is effective in detecting only the deep-level defects in the depletion region and insensitive to shallow-level defects or interface defects.

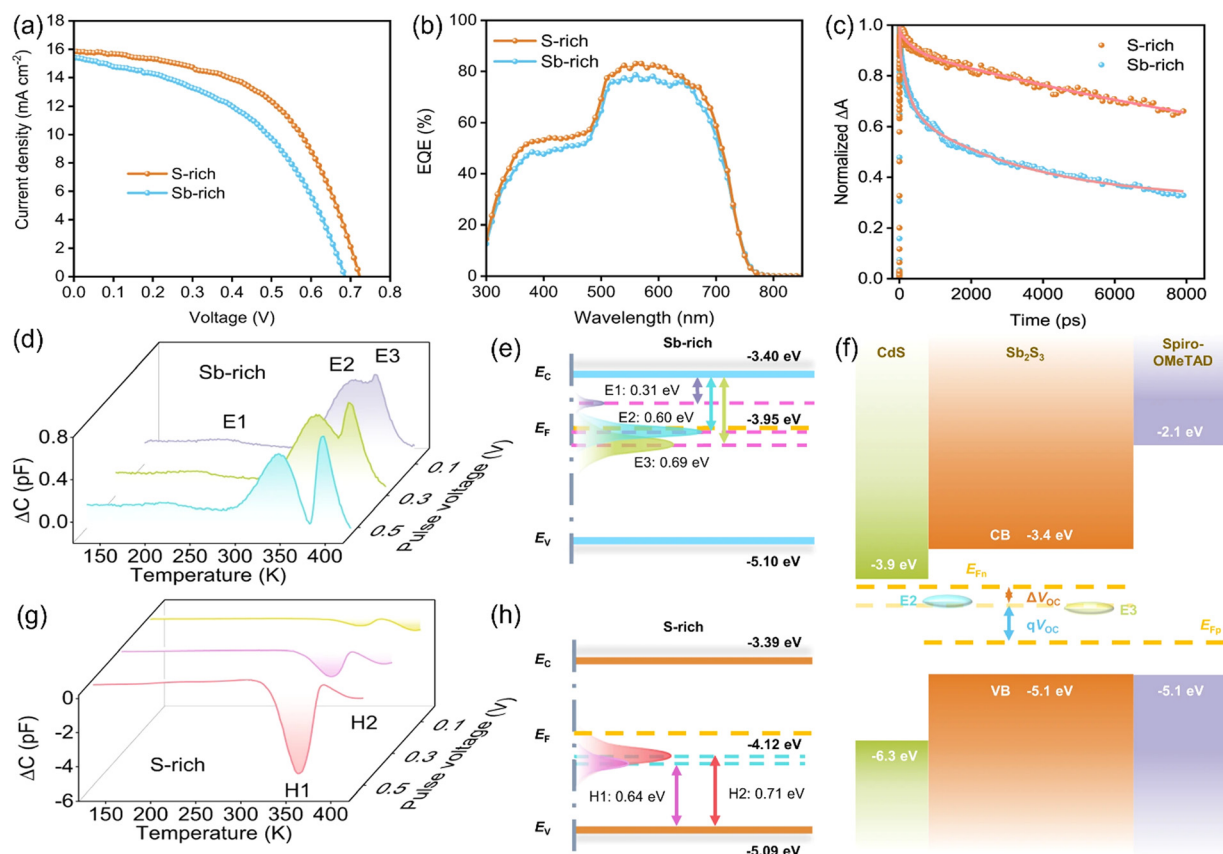
The distinctive Q-1D crystal structure and low structural symmetry in  $Sb_2X_3$  induce unconventional and complicated defect physics. When defects are formed on non-equivalent



sites (two non-equivalent Sb sites, *i.e.*, Sb1 and Sb2, and three non-equivalent X sites, *i.e.*, X1, X2, and X3), the formation energies and charge-state transition levels are different, resulting in multiple inherent defects.<sup>68,69</sup> Unambiguously, the defect types are closely related to the growth conditions, in which  $V_{\text{S}}/V_{\text{Se}}$  (S/Se-vacancy) and  $\text{Sb}_{\text{S}}/\text{Sb}_{\text{Se}}$  (Sb-antisites) defects are formed under Sb-rich conditions, whereas  $V_{\text{Sb}}$  (Sb-vacancy) and  $\text{Se}_{\text{Sb}}/\text{Sb}_{\text{S}}$  (Se/S-antisites) are dominant under Se- or S-rich conditions.<sup>25,26</sup> Interestingly, these defects are located at around the mid-gap with densities ( $N_{\text{T}}$ ) in the range of  $10^{15}$ – $10^{17}$   $\text{cm}^{-3}$ , and  $\tau_{\text{eff}}$  is measured to be in the range of 1.3–67 ns.<sup>70,71</sup> Huang *et al.*<sup>68</sup> performed an in-depth study on the defects in  $\text{Sb}_2\text{S}_3$  via first principles calculations. Owing to presence of two non-equivalent Sb sites (Sb1 and Sb2), and three Se sites (Se1, Se2 and Se3), there are two Sb vacancies ( $V_{\text{Sb1}}$  and  $V_{\text{Sb2}}$ ), three Se vacancies ( $V_{\text{Se1}}$ ,  $V_{\text{Se2}}$  and  $V_{\text{Se3}}$ ), two  $\text{Se}_{\text{Sb}}$  antisites ( $\text{Se}_{\text{Sb1}}$  and  $\text{Se}_{\text{Sb2}}$ ), and three  $\text{Sb}_{\text{Se}}$  antisites ( $\text{Sb}_{\text{Se1}}$ ,  $\text{Sb}_{\text{Se2}}$  and  $\text{Sb}_{\text{Se3}}$ ). The cation-replace-anion ( $\text{Sb}_{\text{Se}}$ ) and anion-replace-cation ( $\text{Se}_{\text{Sb}}$ ) antisite defects can have high concentrations and be dominant in  $\text{Sb}_2\text{S}_3$ . Interestingly, two Se replacing one Sb ( $2\text{Se}_{\text{Sb}}$ ) antisite defect can also have a high concentration under Se-rich conditions owing to the large space between different  $[\text{Sb}_4\text{Se}_6]_n$  ribbons and produce a shallow

acceptor level. Computational studies revealed that cation–anion antisite defects have low formation energies and are located near the middle of the bandgap, acting as recombination centers.<sup>6,29,72,73</sup>

Lian *et al.*<sup>74</sup> conducted an O-DLTS study to probe the deep-level defects in Sb-rich and S-rich  $\text{Sb}_2\text{S}_3$  films and solar cells. The results showed that the Sb-rich film displayed three donor defects (E1, E2 and E3, acting as electron traps), with the energy level of 0.31, 0.60, and 0.69 eV below the CBM, respectively (Fig. 2e). In contrast, the S-rich film exhibited only two acceptor defects (H1 and H2, acting as hole traps) with the energy levels of 0.64 and 0.71 eV above the VBM, respectively (Fig. 2h). The E1, E2, E3, H1, and H2 carrier traps were attributed to the Sb-interstitial ( $\text{Sb}_{\text{i}}$ ), S-vacancy ( $V_{\text{S}}$ ),  $\text{Sb}_{\text{S}}$  antisite, Sb-vacancy ( $V_{\text{Sb}}$ ), and  $\text{Sb}_{\text{Sb}}$  antisite defects, respectively. The E2 ( $V_{\text{S}}$ ) and E3 ( $\text{Sb}_{\text{S}}$ ) traps in the Sb-rich  $\text{Sb}_2\text{S}_3$  films had a larger capture cross-section, higher trap density, and located closer to the Fermi energy level ( $E_{\text{F}}$ ) compared to H1( $V_{\text{Sb}}$ ) and H2 ( $\text{Sb}_{\text{Sb}}$ ) in the S-rich  $\text{Sb}_2\text{S}_3$ . There is a high possibility that the electron quasi- $E_{\text{F}}$  is pinned near the E2 and E3 traps in Sb-rich  $\text{Sb}_2\text{S}_3$ . In contrast, S-rich  $\text{Sb}_2\text{S}_3$  displayed a reduced capture cross section and defect density, resulting in suppressed recombination and



**Fig. 2** (a)  $J$ - $V$  characteristics and (b) EQE spectra of the fabricated  $\text{Sb}_2\text{S}_3$  solar cells. (c) Transient kinetic decay (scatter) and fittings according to biexponential decay function (solid lines) monitored at 545 nm of Sb-rich and S-rich films. (d) DLTS signals of Sb-rich and (g) S-rich  $\text{Sb}_2\text{S}_3$  films at different pulse voltages ranging from 0.1 to 0.5 V, synergized with an identical pulse-width optical pulse. (e) Conduction band ( $E_{\text{C}}$ ), valence band ( $E_{\text{V}}$ ), Fermi level ( $E_{\text{F}}$ ), and trap energy level ( $E_{\text{T}}$ ) of Sb-rich and (h) S-rich  $\text{Sb}_2\text{S}_3$  films. (f) Schematic diagram of  $V_{\text{OC}}$  derived from splitting of the electron and hole quasi-Fermi levels. Adapted under the guidelines of the Creative Commons CC BY license from ref. 74 Copyright 2019, Nature Springer.

**Table 1** Defect state, energy level ( $E_T$ ), cross-section ( $\sigma$ ), defect density ( $N_T$ ) and carrier density ( $N_S$ ) of the detected defects in  $Sb_2X_3$  thin-films and solar cells. Capture cross-sections of  $\sim 10^{-16}$ – $10^{-17}$  cm<sup>2</sup> indicate that defects may act as effective trapping or recombination centers

Material	Deposition technique	Defect characterization technique	Trap	Type	$E_T$ (eV)	$\sigma$ (cm <sup>2</sup> )	$N_T$ (cm <sup>-3</sup> )	Ref.
$Sb_2S_3$	Chemical bath deposition (CBD) Hydrothermal	O-DLTS	H1	S <sub>Sb2</sub>	$E_V + 0.500$	$1.24 \times 10^{-17}$	$1.85 \times 10^{13}$	8
			H3	S <sub>Sb1</sub>	$E_V + 0.776$	$6.59 \times 10^{-16}$	$1.57 \times 10^{14}$	
			H1	V <sub>Sb</sub>	$E_C - 0.572$	$2.98 \times 10^{-17}$	$4.22 \times 10^{13}$	82
			H2	S <sub>Sb2</sub>	$E_C - 0.671$	$3.81 \times 10^{-16}$	$2.69 \times 10^{14}$	
			H1	V <sub>Sb</sub>	$E_V + 0.576$	$2.52 \times 10^{-17}$	$6.67 \times 10^{13}$	81
		DLTS	H2	S <sub>Sb</sub>	$E_V + 0.701$	$4.37 \times 10^{-16}$	$1.08 \times 10^{14}$	
			H1	V <sub>Sb</sub>	$E_V + 0.536$	$2.39 \times 10^{-16}$	$3.71 \times 10^{15}$	5
			H2	S <sub>Sb</sub>	$E_V + 0.709$	$4.49 \times 10^{-14}$	$3.67 \times 10^{15}$	
			H1	V <sub>Sb</sub>	$E_V + 0.537$	$1.68 \times 10^{-16}$	$3.52 \times 10^{15}$	
			H2	S <sub>Sb</sub>	$E_V + 0.697$	$2.62 \times 10^{-14}$	$4.50 \times 10^{14}$	
Thiourea-treated $Sb_2S_3$			E1	S <sub>i</sub>	$E_C - 0.325$	$1.12 \times 10^{-17}$	$1.16 \times 10^{14}$	
$Sb_2S_3$			H1	S <sub>Sb</sub>	$E_V + 0.445$	$6.14 \times 10^{-15}$	$2.16 \times 10^{13}$	77
			H2		$E_V + 0.698$	$2.31 \times 10^{-16}$	$4.25 \times 10^{14}$	
(NH <sub>4</sub> ) <sub>2</sub> S-treated $Sb_2S_3$			H2		$E_V + 0.708$	$6.85 \times 10^{-16}$	$6.29 \times 10^{13}$	
$Sb_2(S,Se)_3$		O-DLTS	E1	V <sub>S</sub>	$E_C - 0.325$	$6.90 \times 10^{-17}$	$6.57 \times 10^{14}$	83
			H1	S <sub>Sb</sub>	$E_V + 0.546$	$8.33 \times 10^{-17}$	$1.20 \times 10^{15}$	
			H2		$E_V + 0.673$	$1.14 \times 10^{-16}$	$2.26 \times 10^{15}$	
			E1	V <sub>S</sub>	$E_C - 0.255$	$6.39 \times 10^{-18}$	$1.02 \times 10^{16}$	84
			H1	V <sub>Sb</sub>	$E_V + 0.321$	$2.15 \times 10^{-16}$	$1.15 \times 10^{14}$	
			H2	S <sub>Sb2</sub>	$E_V + 0.467$	$5.08 \times 10^{-16}$	$1.55 \times 10^{14}$	
			H3	S <sub>Sb1</sub>	$E_V + 0.579$	$7.15 \times 10^{-17}$	$3.14 \times 10^{16}$	
			H4	S <sub>Sb3</sub>	$E_V + 0.781$	$8.42 \times 10^{-15}$	$2.10 \times 10^{16}$	
			H1	S <sub>Sb</sub>	$E_V + 0.664$	$4.05 \times 10^{-16}$	$1.26 \times 10^{14}$	76
			H2		$E_V + 0.713$	$4.47 \times 10^{-15}$	$1.19 \times 10^{14}$	
$Sb_2(S,Se)_3$		DLTS O-DLTS	H3		$E_V + 0.741$	$2.39 \times 10^{-14}$	$7.30 \times 10^{13}$	
			H1	S <sub>Sb2</sub>	$E_V + 0.507$	$1.27 \times 10^{-15}$	$1.46 \times 10^{15}$	80
			H2	S <sub>Sb1</sub>	$E_V + 0.689$	$4.95 \times 10^{-17}$	$3.42 \times 10^{15}$	
			H3	S <sub>Sb3</sub>	$E_V + 0.762$	$4.45 \times 10^{-14}$	$7.62 \times 10^{15}$	
			H1	S <sub>Sb2</sub>	$E_V + 0.502$	$1.05 \times 10^{-17}$	$2.24 \times 10^{14}$	
			H3	S <sub>Sb3</sub>	$E_V + 0.766$	$1.28 \times 10^{-15}$	$1.21 \times 10^{15}$	
			E1	S <sub>Sb2</sub>	$E_C - 0.57$	$2.48 \times 10^{-12}$	$1.45 \times 10^{15}$	85
			E2	S <sub>Sb1</sub>	$E_C - 0.71$	$4.39 \times 10^{-12}$	$4.78 \times 10^{15}$	
			H1	S <sub>Sb</sub>	$E_V + 0.761$	$3.97 \times 10^{-15}$	$6.45 \times 10^{15}$	
			E1	S <sub>Sb2</sub>	$E_C - 0.584$	$1.99 \times 10^{-13}$	$6.49 \times 10^{14}$	
$Sb_2(S,Se)_3$		DLTS	H1	S <sub>Sb</sub>	$E_V + 0.784$	$1.6 \times 10^{-14}$	$2.97 \times 10^{15}$	
			H1	S <sub>Sb1</sub>	$E_V + 0.500$	$1.99 \times 10^{-17}$	$6.28 \times 10^{12}$	10
			H2	S <sub>Sb2</sub>	$E_V + 0.671$	$5.91 \times 10^{-17}$	$1.88 \times 10^{13}$	
			E1	V <sub>S2</sub>	$E_C - 0.747$	$1.60 \times 10^{-16}$	$2.47 \times 10^{13}$	49
			H1	S <sub>Sb</sub>	$E_V + 0.290$	$6.61 \times 10^{-17}$	$6.55 \times 10^{12}$	
			H2	S <sub>Sb1</sub>	$E_V + 0.638$	$3.18 \times 10^{-15}$	$4.15 \times 10^{12}$	
			H3	S <sub>Sb2</sub>	$E_V + 0.483$	$1.48 \times 10^{-16}$	$4.64 \times 10^{14}$	
			H4	S <sub>Sb2</sub>	$E_V + 0.682$	$3.67 \times 10^{-14}$	$8.63 \times 10^{14}$	
			H5	S <sub>Sb3</sub>	$E_V + 0.790$	$4.87 \times 10^{-14}$	$9.75 \times 10^{14}$	
			H1	S <sub>Sb</sub>	$E_V + 0.560$	$5.98 \times 10^{-18}$	$4.97 \times 10^{15}$	86
$Sb_2Se_3$	CBD	Admittance Spectroscopy	H2		$E_V + 0.770$	$2.68 \times 10^{-16}$	$6.40 \times 10^{15}$	
			E1	S <sub>Sb</sub>	$E_C - 0.563$	$1.40 \times 10^{-16}$	$1.97 \times 10^{15}$	
			H1	V <sub>Se</sub>	$E_V + 0.203$	$2.43 \times 10^{-19}$	$5.37 \times 10^{14}$	87
			H2	S <sub>Se</sub>	$E_V + 0.364$	$3.85 \times 10^{-16}$	$1.27 \times 10^{15}$	
			H1	S <sub>Se1</sub>	$E_V + 0.609$	$3.23 \times 10^{-17}$	$1.01 \times 10^{13}$	9
		Admittance spectroscopy	H2	S <sub>Se2</sub>	$E_V + 0.691$	$3.30 \times 10^{-15}$	$4.27 \times 10^{12}$	
			H1	Bulk trap	$E_V + 0.286$	$3.2 \times 10^{-17}$	$2.46 \times 10^{14}$	88
			H2		$E_V + 0.188$	$9.9 \times 10^{-21}$	$1.36 \times 10^{15}$	
			H3	Interface trap	$E_V + 0.570$	$4.1 \times 10^{-16}$	$1.15 \times 10^{16}$	
			H1	S <sub>Se</sub>	$E_V + 0.111$	—	$1.3 \times 10^{15}$	89
$Sb_2Se_3$	Vapor transport deposition (VTD)	O-DLTS	H2	—	$E_V + 0.578$	—	—	
			H1	V <sub>Sb</sub>	$E_V + 0.48$	$1.5 \times 10^{-17}$	$1.2 \times 10^{15}$	90
			H2	S <sub>Se</sub>	$E_V + 0.71$	$4.9 \times 10^{-13}$	$1.1 \times 10^{14}$	
			E1	S <sub>Se</sub>	$E_C - 0.61$	$4.0 \times 10^{-13}$	$2.6 \times 10^{14}$	
			E1	V <sub>Se</sub>	$E_C - 0.206$	$5.34 \times 10^{-21}$	$2.74 \times 10^{14}$	91
		Temperature-dependent (dark) conductivity and admittance spectroscopy	E2	S <sub>Se2</sub>	$E_C - 0.569$	$1.24 \times 10^{-16}$	$8.02 \times 10^{14}$	
			H1	S <sub>Se2</sub>	$E_V + 0.531$	$3.79 \times 10^{-18}$	$3.48 \times 10^{14}$	
			E1	V <sub>Se2</sub>	$E_C - 0.359$	$1.03 \times 10^{-18}$	$2.70 \times 10^{14}$	15
			E2	S <sub>Se2</sub>	$E_C - 0.609$	$1.54 \times 10^{-16}$	$1.73 \times 10^{15}$	
			H1	S <sub>Se2</sub>	$E_V + 0.691$	$1.37 \times 10^{-16}$	$8.76 \times 10^{14}$	
$Sb_2Se_3$	Spin coating	DLTS	H1	V <sub>Sb</sub>	$E_V + 0.479$	$2.07 \times 10^{-16}$	$6.10 \times 10^{14}$	92
			H2	S <sub>Se</sub>	$E_V + 0.713$	$1.53 \times 10^{-14}$	$1.00 \times 10^{15}$	





Table 1 (continued)

Material	Deposition technique	Defect characterization technique	Trap	Type	$E_T$ (eV)	$\sigma$ (cm <sup>2</sup> )	$N_T$ (cm <sup>-3</sup> )	Ref.
(2.21%) Te-doped Sb <sub>2</sub> S <sub>3</sub> (5.23%) S-doped Sb <sub>2</sub> S <sub>3</sub>			E1	Sb <sub>Se</sub>	$E_c - 0.628$	$7.05 \times 10^{-17}$	$6.49 \times 10^{14}$	
			H1	V <sub>sb</sub>	$E_v + 478$	$2.71 \times 10^{-15}$	$5.34 \times 10^{13}$	
			E1	Sb <sub>Se</sub>	$E_c - 0.637$	$9.48 \times 10^{-16}$	$3.32 \times 10^{14}$	
			H1	V <sub>sb</sub>	$E_v + 0.500$	$6.80 \times 10^{-14}$	$6.93 \times 10^{14}$	
			H2	Se <sub>Sb</sub>	$E_v + 0.721$	$8.19 \times 10^{-14}$	$1.43 \times 10^{15}$	
			E1	Sb <sub>Se</sub>	$E_c - 0.633$	$2.54 \times 10^{-16}$	$6.60 \times 10^{14}$	

prolonged carrier lifetimes. The trap state distribution is schematically shown in Fig. 2d–h and the parameters are listed in Table 1. The TAS decay profile (Fig. 2c) revealed that the S-rich Sb<sub>2</sub>S<sub>3</sub> films exhibit much longer carrier lifetimes (18.7 ns) than the Sb-rich Sb<sub>2</sub>S<sub>3</sub> films (3.8 ns). The reduced defect type (and density) contributed to the prolonged carrier lifetimes, improved  $V_{OC}$  and PCE. The S-deficit in the Sb-rich Sb<sub>2</sub>S<sub>3</sub> resulted in the inevitable formation of  $V_S$ . Subsequently, the excess Sb preferentially filled the  $V_S$ , rather than entering the interstitial site, given that the formation energy for Sb<sub>S</sub> was lower than that for Sb<sub>i</sub>. Therefore, the E2 ( $V_S$ ) and E3 (Sb<sub>S</sub>) defects were predominant in the Sb-rich Sb<sub>2</sub>S<sub>3</sub>. However, in the S-rich Sb<sub>2</sub>S<sub>3</sub>, initially, the S atoms entered the crystal structure to passivate  $V_S$ . Under the S-rich conditions, the formation of Sb<sub>S</sub> and Sb<sub>i</sub> defects was suppressed due to their increased formation energy, while the formation of a large amount of H1 ( $V_{Sb}$ ) was induced due to its low formation energy. Furthermore, some S atoms may even occupy the  $V_{Sb}$  and form the H2 (S<sub>Sb</sub>) antisite to maintain structural stability. It was also revealed that the Sb<sub>i</sub> defect did not have a critical influence on the carrier lifetimes, indicating the high tolerance of the Q-1D crystal structure. The space in (Sb<sub>4</sub>S<sub>6</sub>)<sub>n</sub> ribbons can easily accommodate impurities/defects to a certain extent. Cai *et al.*<sup>75</sup> performed a first principles study on Sb<sub>2</sub>S<sub>3</sub> and concluded that the formation energies of the dominant acceptor defects ( $V_{Sb2}$ , Sb<sub>S2</sub>, and S<sub>Sb2</sub>) and donor defect ( $V_{S2}$ ) in intrinsic Sb<sub>2</sub>S<sub>3</sub> are similar. This results in  $E_F$  level pinning near the middle of the band gap. These results are consistent with the experimentally observed high resistivity values ( $\sim 10^8 \Omega \text{ cm}$ ) in undoped Sb<sub>2</sub>S<sub>3</sub> thin films. It is widely accepted that the high resistivity in Sb<sub>2</sub>S<sub>3</sub> originates from the compensation between its intrinsic donor ( $V_S$ ) and acceptors ( $V_{Sb}$ , Sb<sub>S</sub>, and S<sub>Sb</sub>), which have comparably high densities and low formation energies. This compensation also limits the improvement in conductivity through direct extrinsic doping.

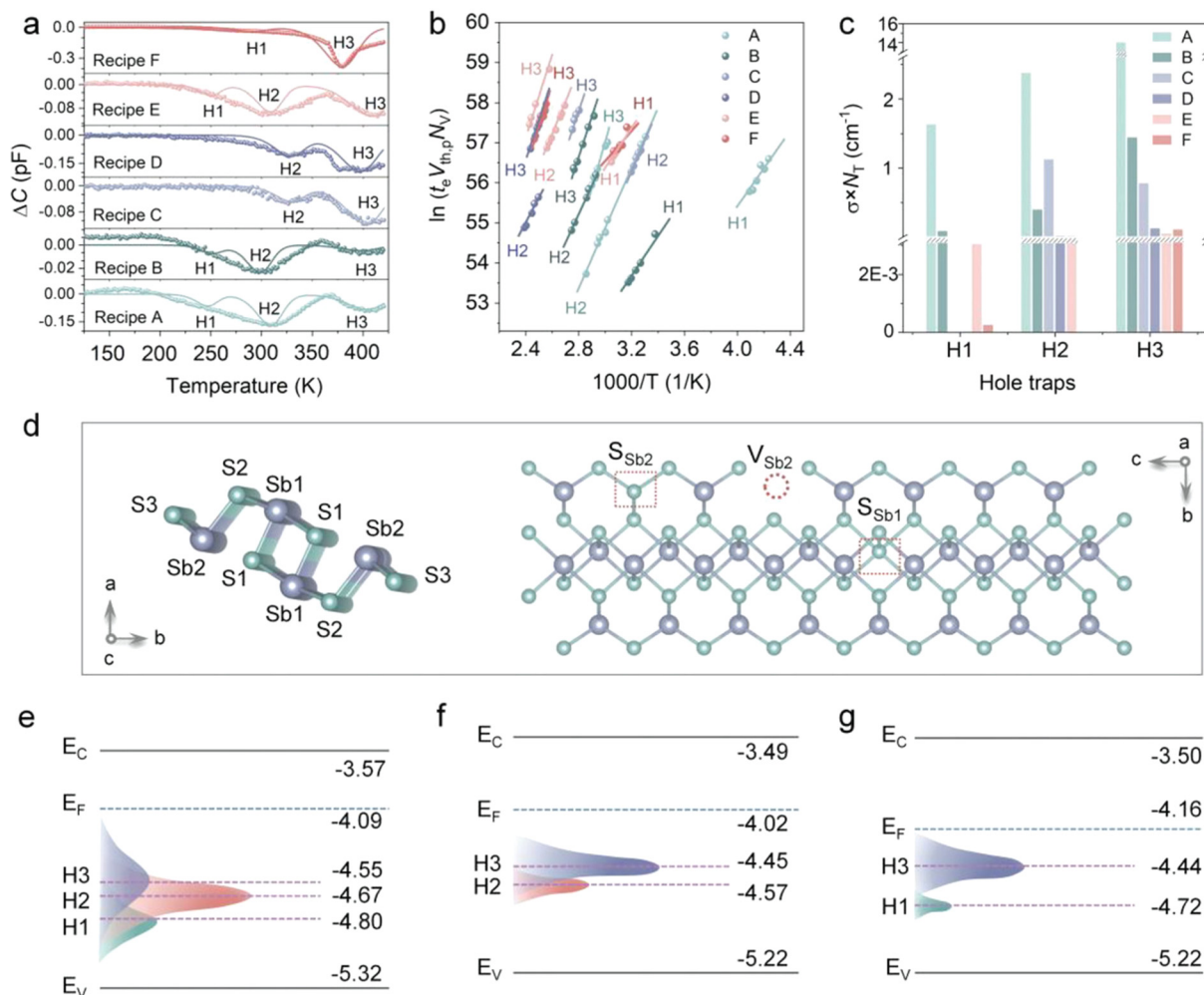
Wang *et al.*<sup>8</sup> fabricated Sb<sub>2</sub>S<sub>3</sub> solar cells *via* chemical bath deposition (CBD) employing six different routes. It was demonstrated that using multiple S-sources (sodium thiosulfate and thioacetamide) accelerated the release of S<sup>2-</sup> ions, leading to the formation of S-rich Sb<sub>2</sub>S<sub>3</sub> films with better morphology (average grain size  $\sim 2.8 \mu\text{m}$ ), crystallographic orientation, prolonged carrier lifetimes and less defects. Energy dispersive X-ray (EDX) analysis revealed atomic ratios of S/Sb of  $>1.5$  for all the Sb<sub>2</sub>S<sub>3</sub> films, indicating the (desirable) S-richness in the films. The defects in the n-Sb<sub>2</sub>S<sub>3</sub> films were characterized *via*

O-DLTS. The O-DLTS signal revealed only negative peaks, suggesting the presence of minority traps (holes), *i.e.*, H1, H2 and H3, with  $E_a$  of  $\approx 0.51$ , 0.66, and 0.77 eV, respectively (Fig. 3). H1 and H3 were attributed to the antisite defects (S<sub>Sb2</sub> and S<sub>Sb1</sub>, respectively), while H2 to  $V_{Sb2}$ . All three traps were found to be located below the intrinsic  $E_F$ . Compared to the H2 and H3 traps, the H1 trap was located further away from  $E_F$ , which was generally submerged in electrons and remained passive. The H1 trap exhibited a lower trap density and smaller capture cross-section, and thus a less detrimental effect on carrier trapping. On the contrary, the deep-level H3 (S<sub>Sb1</sub>) trap was located closer to  $E_F$  and exhibited a larger capture cross-section than the other traps, and thus was the most active in charge trapping. The devices fabricated using the modified recipe were free from the H2 deep-level defect ( $V_{Sb2}$ ) and demonstrated a record PCE of 8.0% in Sb<sub>2</sub>S<sub>3</sub> solar cells. Furthermore, bi-exponential fitting of the TAS curve revealed a prolonged carrier lifetime of 9298 ps in the FTO/CdS/Sb<sub>2</sub>S<sub>3</sub> stack, indicating efficient transport and reduced recombination. Furthermore, an excellent diode ideality factor of 1.31 and low reverse saturation current were obtained, suggesting the substantial suppression of non-radiative (SRH) recombination in the device.

Huang *et al.*<sup>76</sup> demonstrated *in situ* sulfurization by introducing tartaric acid (TTA) additive in the hydrothermal deposition process of Sb<sub>2</sub>S<sub>3</sub>. The consensus is that reducing the density of Sb<sub>S</sub> defect is one key aspect to enhance the PCE of Sb<sub>2</sub>S<sub>3</sub> solar cells. S atoms easily volatilized from the Sb<sub>2</sub>S<sub>3</sub> thin films during annealing to form  $V_S$ . Given that the formation energy of Sb<sub>i</sub> is much higher than that of Sb<sub>S</sub>, excess Sb preferentially filled the  $V_S$  voids, thus forming Sb<sub>S</sub> antisite defects. *In situ* sulfurization induced S-richness in the annealed Sb<sub>2</sub>S<sub>3</sub> films. Consequently, this increased the formation energy of Sb<sub>S</sub> defects. To accommodate the stability of the structure in S-rich Sb<sub>2</sub>S<sub>3</sub> films, some of the S atoms occupied  $V_{Sb}$  to form S<sub>Sb</sub> defects. *In situ* sulfurization successfully compensated the S-loss occurring during annealing and suppressed the formation of Sb<sub>S</sub>. Three hole traps (H1, H2 and H3) were obtained for the control Sb<sub>2</sub>S<sub>3</sub> sample, which are attributed to Sb<sub>S</sub> defects. In the case of the sulfurized sample, only H1 and H3 traps were obtained. The passivation of the Sb<sub>S</sub> defects and improvement in crystallinity led to an improved PCE (6.31%) in the *in situ*-sulfurized device, outperforming the control device (PCE of 5.46%).

Huang *et al.*<sup>77</sup> demonstrated a novel (NH<sub>4</sub>)<sub>2</sub>S-induced hydrothermal sulfurization process for the fabrication of Sb<sub>2</sub>S<sub>3</sub> solar cells. (NH<sub>4</sub>)<sub>2</sub>S undergoes hydrolysis to produce H<sub>2</sub>S, a strong





**Fig. 3** (a) O-DLTS signals and high-resolution evaluation simulations for Sb<sub>2</sub>S<sub>3</sub> solar cells based on recipes A–F. (b) Corresponding Arrhenius plots obtained from the O-DLTS signals. (c) Histogram of the calculated  $\sigma \times N_T$  values of different hole traps in Sb<sub>2</sub>S<sub>3</sub> solar cells based on recipes A–F. (d) Schematic diagram of an Sb<sub>4</sub>S<sub>6</sub> ribbon with five nonequivalent atomic sites and S<sub>Sb1</sub>, S<sub>Sb2</sub>, and V<sub>Sb2</sub> defects in the Sb<sub>2</sub>S<sub>3</sub> lattice. (e)–(g) Energy levels and defect levels of the devices based on recipes B, D, and F, respectively. Adapted with permission from ref. 8 Copyright 2022, Wiley-VCH.

sulfurization agent. The XRD patterns revealed improvements in crystallinity, while the SEM micrographs suggested no improvement in morphology on adopting the sulfurization strategy. DLTS analysis revealed the presence of two minority carrier (hole) traps, *i.e.*, H1 and H2, in the control films, with the activation energy of 445 and 698 meV, respectively, which were both attributed to Sb<sub>s</sub> antisite defects. The sulfurization process annihilated H1, given that only H2 was detected in the sulfurized films. The PCE of the solar cell (FTO/CdS/Sb<sub>2</sub>S<sub>3</sub>/spiro-OMeTAD/Au) improved from 6.01% to 6.92% on adopting the sulfurization strategy. The reverse (leakage) current density ( $J_0$ ) value decreased from  $4.85 \times 10^{-2}$  mA cm<sup>-2</sup> to  $2.07 \times 10^{-5}$  mA cm<sup>-2</sup> in the respective devices, indicating a significant improvement in diode quality and rectification characteristics. Choi *et al.*<sup>78</sup> investigated the trap-state distribution in CBD-deposited and post-deposition thioacetamide (TA, CH<sub>3</sub>CSNH<sub>2</sub>)-treated Sb<sub>2</sub>S<sub>3</sub> films using capacitance transient-based DLTS, where Sb<sub>2</sub>S<sub>3</sub>/Au-based Schottky diodes (500 mm diameter) were fabricated for the measurement. TA decomposed into volatile

acetonitrile (CH<sub>3</sub>CN) and hydrogen sulfide (H<sub>2</sub>S) gas on annealing at temperatures exceeding 150 °C in an inert atmosphere. The released H<sub>2</sub>S readily reacted with the surface of the Sb<sub>2</sub>S<sub>3</sub> film, passivating the V<sub>S</sub> defects and reducing the oxide phase (oxidized sulfur, sulfates, and Sb<sub>2</sub>O<sub>3</sub>). Interestingly, the untreated sample demonstrated a single hole trap with an activation energy of  $E_v + 0.52$  eV, capture cross-section of  $1.34 \times 10^{-17}$  cm<sup>2</sup> and density of  $(2-5) \times 10^{14}$  cm<sup>-3</sup>, while no trap was found for the TA-treated sample. TA treatment was found to be instrumental in boosting the  $V_{OC}$  and PCE values from 570.5 meV and 5.5% (in the control device) to 645.7 meV and 7.5% (in champion device), respectively. This spectacular improvement was attributed to the annihilation of V<sub>S</sub> defects (near the Sb<sub>2</sub>S<sub>3</sub>/HTL) by H<sub>2</sub>S-assisted sulfurization. H<sub>2</sub>S-assisted sulfurization has also been found to be instrumental in passivating V<sub>S/Se</sub> defects, reducing the (detrimental) Sb<sub>2</sub>O<sub>3</sub> phase (at the surface) and boosting the carrier lifetimes and mobility in Sb<sub>2</sub>Se<sub>3</sub> films.<sup>79</sup>

Similar to TA, thiourea (TU, SC(NH<sub>2</sub>)<sub>2</sub>) treatment has also been used as an effective sulfurization strategy (by releasing



H<sub>2</sub>S during decomposition at temperatures of >150 °C) for passivating V<sub>S/Se</sub> defects in Sb<sub>2</sub>X<sub>3</sub> films. Qi *et al.*<sup>5</sup> performed an investigation on the trap physics in hydrothermally deposited Sb<sub>2</sub>S<sub>3</sub> films using DLTS. The as-deposited films exhibited two minority (hole) traps, *i.e.*, H1 and H2 (attributed to V<sub>Sb</sub> and Sb<sub>S</sub>, respectively). These defects possessed high ionization energies of 0.536 and 0.709 eV, respectively, and served as severe recombination centers. H2 was highly effective in charge trapping and V<sub>OC</sub> loss (owing to its high density and large capture cross-section), while H1 was less effective (low capture cross-section of  $\sim 10^{-16}$  cm<sup>2</sup>). TU treatment on Sb<sub>2</sub>S<sub>3</sub> films led to the suppression of the H2 (Sb<sub>S</sub>) defect density by one order, suggesting the alleviation of S-deficiency at the surface.

Liang *et al.*<sup>15</sup> fabricated flexible Sb<sub>2</sub>Se<sub>3</sub> solar cells (Mo/Sb<sub>2</sub>Se<sub>3</sub>/CdS/ZnO/AZO/Ag) on a polyimide (PI) substrate and investigated the influence of an ultra-thin PbSe layer (at the Mo/Sb<sub>2</sub>Se<sub>3</sub> interface) on the device performance. The control device

(without a PbSe interlayer), *i.e.*, C-Sb<sub>2</sub>Se<sub>3</sub>, was found to underperform (PCE of 3.04%) compared to the device with a PbSe interlayer, *i.e.*, Pb-Sb<sub>2</sub>Se<sub>3</sub> (PCE of 8.43%). The improved performance of the latter was attributed to the synergistic influence of the PbSe layer in suppressing the unfavorable *[hk0]*-orientated grains, alleviating the Schottky barrier height (at the Mo/Sb<sub>2</sub>Se<sub>3</sub> interface) and mitigating the bulk-defects in the absorber layer. The TPV study revealed longer carrier lifetimes ( $\sim 0.78$  ms) in the Pb-Sb<sub>2</sub>Se<sub>3</sub> devices in comparison to the C-Sb<sub>2</sub>Se<sub>3</sub> devices ( $\sim 0.02$  ms). The defect dynamics in both devices were investigated using DLTS (Fig. 4c–f). The Arrhenius plots (Fig. 4d) were found to be straight lines, suggesting that the activation energy  $E_a$  (slope of the plots) and capture cross-sections (calculated from the Y-intercept) were strictly independent of temperature. Two electron traps (E1 and E2) and one hole trap (H2) were identified from the DLTS signal in both devices (Fig. 4c). Trap E1 was attributed to the V<sub>Se2</sub> vacancy

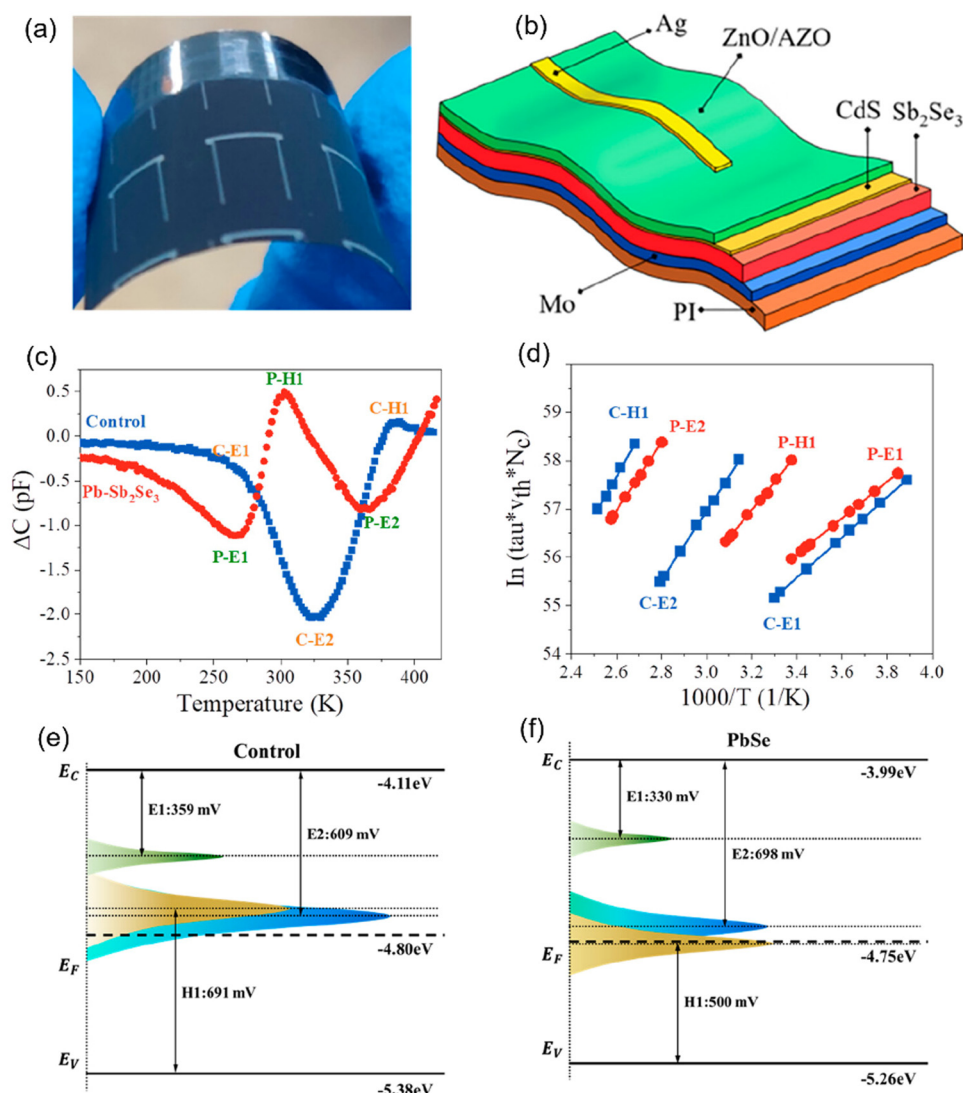


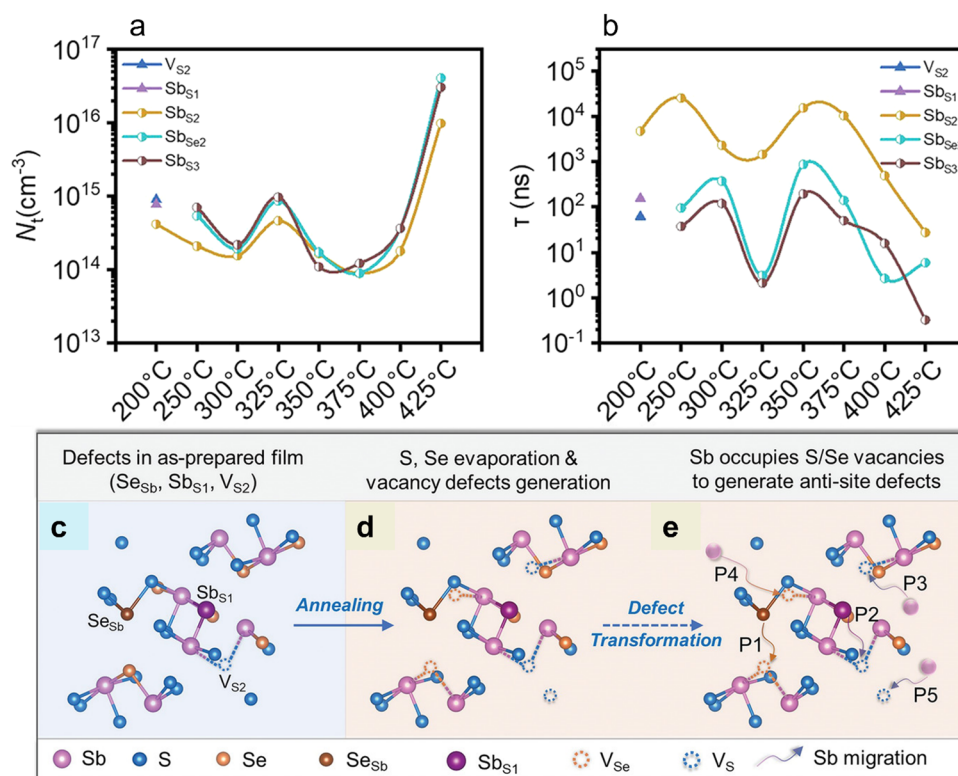
Fig. 4 (a) Photograph and (b) schematic device structure of flexible Sb<sub>2</sub>Se<sub>3</sub> solar cells fabricated on PI in a substrate configuration of PI/Mo/Sb<sub>2</sub>Se<sub>3</sub>(Pb-Sb<sub>2</sub>Se<sub>3</sub>)/CdS/ZnO/AZO/Ag. (c) DLTS signals of Pb-Sb<sub>2</sub>Se<sub>3</sub> and the control devices. (d) Arrhenius plots obtained from DLTS signals. (e) and (f) Energy states and defect levels of the control device and Pb-Sb<sub>2</sub>Se<sub>3</sub> solar cells. Adapted with permission from ref. 15 Copyright 2023, the American Chemical Society.



defect.  $V_{Se}$  exhibited a low formation energy under both Se-rich and Se-poor conditions and acted as shallow-level electron traps. Also,  $V_{Se}$  exhibited a lower density and capture cross-section (listed in Table 1), and thus ruled out as a stringent recombination center. E2 and H1 were found to be located close and deeper in the forbidden gap, resulting in ambiguous interpretations. These traps could be either attributed to the donor-acceptor pair of the same (amphoteric) defect ( $Sb_{Se}$ ) or [ $Sb_{Se} + Se_{Sb}$ ] antisite complexes. The E2 and H1 defects in the Pb-Sb<sub>2</sub>Se<sub>3</sub> solar cells were attributed to  $V_{Se3}$  and  $Pb_{Sb}$ , respectively. H1 was found to be located much shallower and with (one order) lower capture cross-section in Pb-Sb<sub>2</sub>Se<sub>3</sub> than in C-Sb<sub>2</sub>Se<sub>3</sub>, suggesting its impact in alleviating the carrier lifetimes, and the  $V_{OC}$  was limited in the Pb-Sb<sub>2</sub>Se<sub>3</sub> device.

Che *et al.*<sup>49</sup> investigated the influence of the annealing process on the defect formation mechanism in hydrothermally deposited n-type Sb<sub>2</sub>(S,Se)<sub>3</sub> films. High-temperature annealing (200–425 °C) led to a two-step defect transformation process, *i.e.*, the formation of  $V_{Se}/V_S$  vacancy defects (with high formation energies), followed by the migration of Sb ions to fill the vacancy defects (forming antisite defects with low formation energies). O-DLTS analysis revealed the presence of one electron trap (E1) and five-hole traps (H1, H2, H3, H4, and H5) in the samples. The former was assigned to S-vacancy ( $V_{S2}$ ), while the latter to antisite defects ( $Se_{Sb}$ ,  $Sb_{S1}$ ,  $Sb_{S2}$ ,  $Sb_{Se2}$ , and  $Sb_{S3}$ , respectively). H2–H5 acted as amphoteric defects. E1 ( $V_{S2}$ ), H1

( $Se_{Sb}$ ) and H2 ( $Sb_{S1}$ ) defects with high formation energies appeared in the as-deposited (unannealed) film (Fig. 5c). After post-annealing, these traps gradually decreased and transformed into H3–H5 antisite defects (*i.e.*,  $Sb_{S2}$ ,  $Sb_{Se2}$ , and  $Sb_{S3}$ ), respectively, with lower formation energies. It was concluded that post-annealing provided a driving force for the formation of relatively low-energy defects. S-volatilization left sulfur vacancies ( $E1$ ,  $V_S$ ), which were subsequently occupied by the Sb atoms, as depicted by the P3 process (Fig. 5e), thus creating Sb-substituted S (H3,  $Sb_{S2}$ ) anti-site defect. Furthermore, owing to the instability (high formation energy) of the  $Se_{Sb}$  (H1) defect, the Se atoms (in  $Se_{Sb}$ ) migrate to occupy neighboring  $V_{S(e)}$  under the thermal driving force, leading to the mitigation of H1, as shown by the P1 process. Se-volatilization forms  $V_{Se}$ , which gets occupied by the neighboring Sb atoms, forming Sb-antisite defect (H4,  $Sb_{Se2}$ ), as shown by the P4 process (Fig. 5e). Concurrently, the further loss of S (during annealing) creates more  $V_S$ , which prompts the Sb atoms to occupy these unstable  $V_{S3}$  (P5 process in Fig. 5e), leading to the creation of a new low-formation energy defect, *i.e.*,  $Sb_{S3}$  (H5). Interestingly, the elevated annealing temperature propelled the Sb atoms in the high-formation energy defect H2 ( $Sb_{S1}$ ) to migrate to the surrounding  $V_{S(e)}$  and form defects  $Sb_{S(e)}$  (depicted by the P2 process).  $V_{S2}$ ,  $Sb_{S1}$ ,  $Sb_{Se2}$ , and  $Sb_{S3}$  are all located near the  $E_F$ -level, which significantly influenced the carrier lifetimes and  $V_{OC}$ . The transient absorption spectroscopy (TAS) results revealed that the Sb<sub>2</sub>(S,Se)<sub>3</sub> film



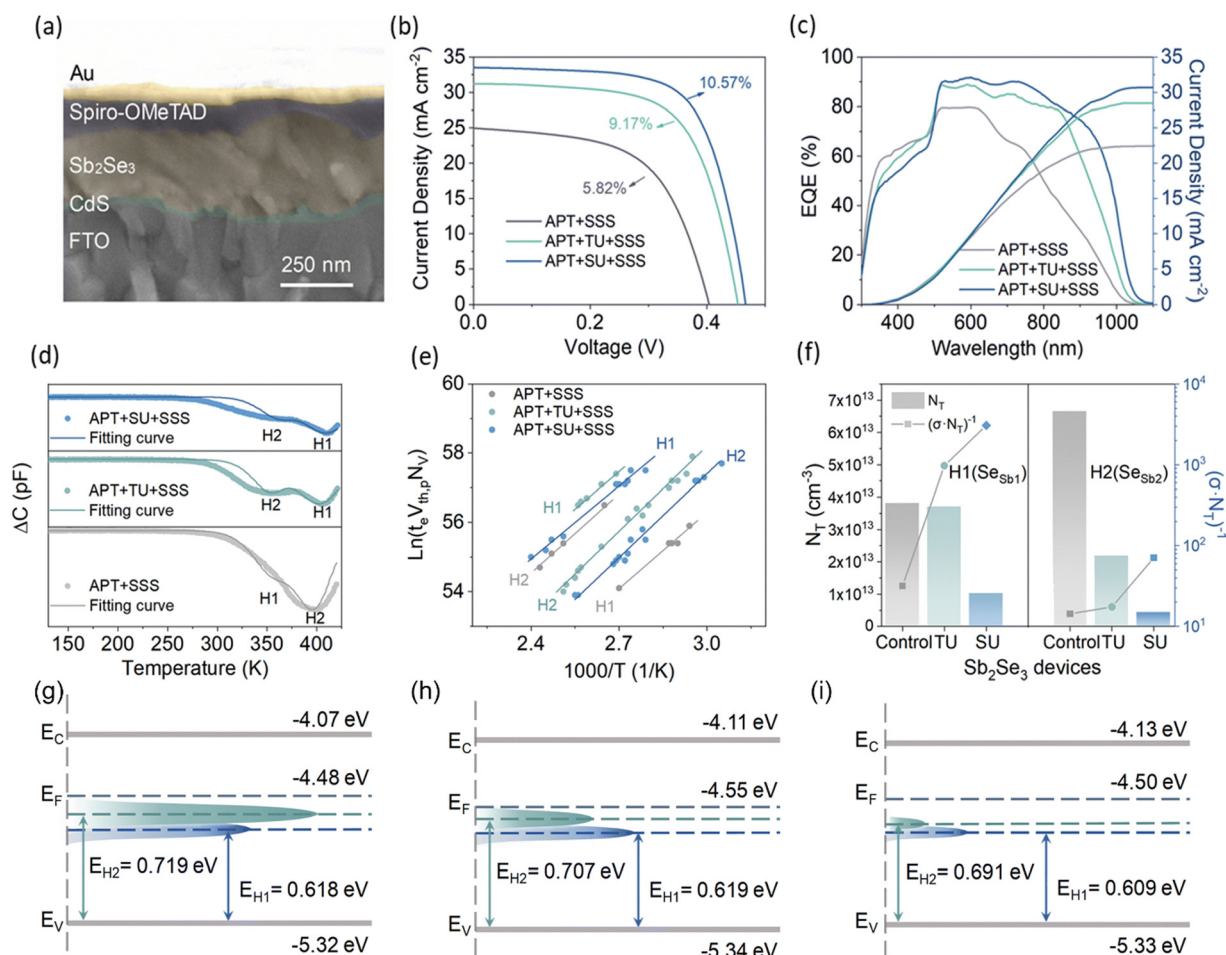
**Fig. 5** (a) Dependence of defect density ( $N_T$ ) on annealing temperature and (b) lifetime associated with specific defects obtained from the O-DLTS signals of the devices. Schematic of (c) three types of point defects in the as-deposited films (without post-annealing). (d) Loss of S and Se atoms to generate vacancy defects during annealing and (e) formation mechanism of antisite defects in the Sb<sub>2</sub>(S,Se)<sub>3</sub> films. Adapted with permission from ref. 49 Copyright 2022, Wiley-VCH.

annealed at 350 °C for 10 min exhibited the longest carrier lifetime (7.0 ns) compared to the films annealed at other temperatures and for higher durations. It was concluded that the  $\text{Sb}_2(\text{S,Se})_3$  films annealed at 350 °C exhibited the lowest defect concentrations and longest carrier lifetimes, thus leading to a high photovoltaic performance (PCE of 9.7%).

Tang *et al.*<sup>80</sup> investigated the defect properties in hydrothermally deposited  $\text{Sb}_2\text{S}_3$  and  $\text{Sb}_2(\text{S,Se})_3$  films using O-DLTS. Three-hole traps (H1, H2 and H3) were detected in all the samples. Also, similar  $E_a$  values were obtained for the traps in all the samples, suggesting similar origins for each type of defect. Specifically, H1, H2 and H3 were located at about 0.50 eV, 0.68 eV and 0.76 eV above the valence band and assigned to antisite defects ( $\text{Sb}_{\text{Sb}2}$ ,  $\text{Sb}_{\text{Sb}1}$ , and  $\text{Sb}_{\text{Sb}3}$ ), respectively. Notably, the capture cross section exhibited a decreasing trend (up to two orders) with an increase in the Se-content in the films. Also, the H2 defect disappeared in the films with a higher Se-content, indicating that the defect properties are strongly correlated with the elemental composition and stoichiometry.

An increase in Se-content not only induced preferred  $[hk1]$ -oriented growth, but also favored efficient carrier transport by alleviating the defect density and capture cross-section, leading to a PCE of 10.10% in the device employing the  $\text{Sb}_2(\text{S,Se})_3$  thin-film (with 29% Se-content), outperforming the  $\text{Sb}_2\text{S}_3$ -based device (PCE of 6.02%).

Zhao *et al.*<sup>9</sup> conducted O-DLTS measurements in Se-rich  $\text{Sb}_2\text{Se}_3$  films, which were synthesized using the CBD technique. Antimony potassium tartrate (APT) and sodium selenosulfate (SSS) were used as the Sb and Se sources, respectively, while thiourea (TU) and selenourea (SU) were used as additives. The films were synthesized *via* three routes, *i.e.*, APT + SSS (no additive), APT + SSS + TU (additive) and APT + SSS + SU (additive). The additives were found to be instrumental in regulating the reaction kinetics and defect properties of the  $\text{Sb}_2\text{Se}_3$  films (Fig. 6d). The detailed defect parameters extracted from the Arrhenius curves (Fig. 6e) and DLTS signals are summarized in Table 1. Two types of hole traps (H1 and H2) were detected in all the samples, which are attributed to



**Fig. 6** (a) Cross-sectional SEM image of  $\text{Sb}_2\text{Se}_3$  planar heterojunction solar cell. (b)  $J$ - $V$  curves and (c) EQE spectra of the  $\text{Sb}_2\text{Se}_3$  solar cells fabricated via three different routes. Plots (d)–(i) show the deep-level defect profile in the  $\text{Sb}_2\text{Se}_3$  films. (d) O-DLTS signals from  $\text{Sb}_2\text{Se}_3$  devices. (e) Arrhenius plots obtained from O-DLTS signals. (f) Values of  $N_T$  (left) and  $(\sigma N_T)^{-1}$  (right) for all the  $\text{Sb}_2\text{Se}_3$  samples at the H1 ( $\text{Se}_{\text{Sb}1}$ ) and H2 ( $\text{Se}_{\text{Sb}2}$ ) defects. Deep-level defect profile of (g) control- $\text{Sb}_2\text{Se}_3$  film, (h) TU- $\text{Sb}_2\text{Se}_3$  film and (i) SU- $\text{Sb}_2\text{Se}_3$  film. Adapted with permission from ref. 9 Copyright 2022, The Royal Society of Chemistry Publishing.

Se-antisites ( $\text{Se}_{\text{Sb1}}$  and  $\text{Se}_{\text{Sb2}}$ , respectively). Under Se-rich conditions, the  $\text{Se}_{\text{Sb1}}$  and  $\text{Se}_{\text{Sb2}}$  antisite defects have relatively low formation energies ( $< 1$  eV), making the replacement of Sb with Se feasible. The Se-antisite defects form deep-level defects,

inducing  $E_{\text{F}}$ -level pinning and serving as proactive SRH-recombination centers. Furthermore, the Se-antisite defects are specifically amphoteric defects with similar carrier capture cross-sections for both electrons and holes, thus severely

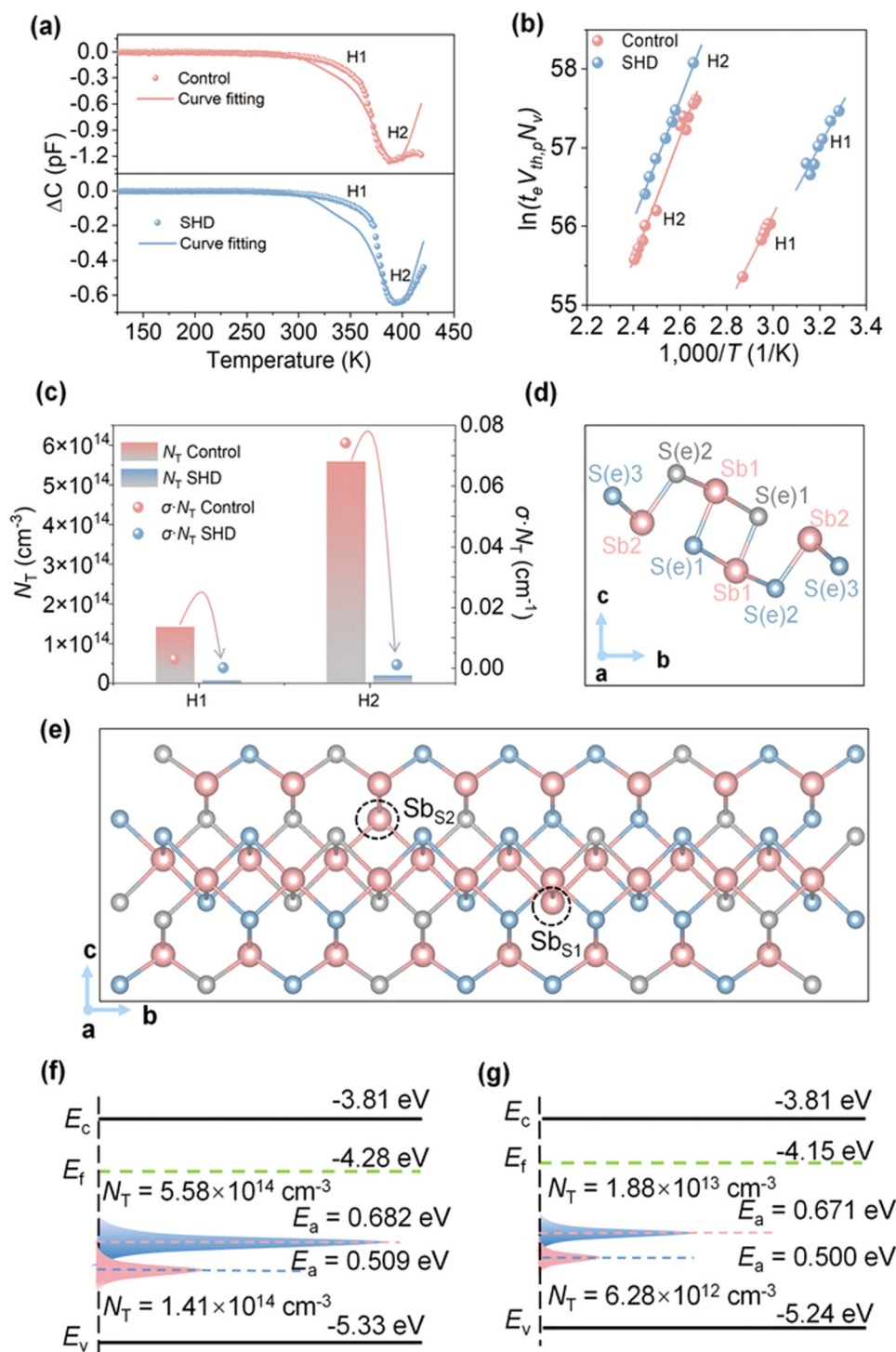


Fig. 7 Deep-level defect analysis. (a) O-DLTS signals and high-resolution evaluation simulations for  $\text{Sb}_2(\text{S,Se})_3$  solar cells without and with ethanol additive and (b) corresponding Arrhenius plots. (c) Values of  $N_T$  and  $(\sigma \cdot N_T)$  at the H1 ( $\text{Sb}_{\text{S2}}$ ) and H2 ( $\text{Sb}_{\text{S1}}$ ) defects. (d) Schematic diagram of an  $\text{Sb}_4\text{S}_6$  ribbon with five non-equivalent atomic sites. (e)  $\text{Sb}_{\text{S2}}$  and  $\text{Sb}_{\text{S1}}$  defects in  $\text{Sb}_2(\text{S,Se})_3$  lattice. Energy states and defect levels of the devices (f) without and (g) with ethanol. Adapted with permission from ref. 10 Copyright 2023, Wiley-VCH.





impeding the carrier transport and limiting the final PV performance. The H2 trap level is located deeper in the band gap than H1, implying that H2 is the dominant defect in these films. The control  $\text{Sb}_2\text{Se}_3$  device had a higher trap density ( $6.65 \times 10^{13} \text{ cm}^{-3}$ ) of H2, which is consistent with its poor film quality. After regulating the CBD reaction in the deposition processes, the values of  $N_T$  significantly decreased. Particularly, the H2 defect density of the SU- $\text{Sb}_2\text{Se}_3$  device was reduced by one order of magnitude, alleviating the SRH recombination. The mitigation of the non-radiative recombination pathways is consistent with the  $J$ - $V$  measurements in the dark and at different light intensities. It was concluded that the sensitive manipulation of the reaction kinetics *via* the additive-assisted CBD technique offers an effective route for the fabrication of high-quality  $\text{Sb}_2\text{Se}_3$  films and high-PCE (10.57%) solar cells.

Wu *et al.*<sup>81</sup> investigated the structure–property–performance relationship in  $\text{Sb}_2\text{S}_3$  thin-film solar cells *via* the in-depth characterization of the crystal orientation, carrier transport and device performance. The study revealed that by merely controlling the post-annealing process, the grain growth, orientation and defect physics can be regulated. Normal grain growth occurred when the  $\text{Sb}_2\text{S}_3$  film was subjected to a low-temperature treatment process, yielding  $[hk1]$ -oriented films. Alternatively, direct annealing at high temperature promoted abnormal grain growth and favored  $[hk0]$ -oriented  $\text{Sb}_2\text{S}_3$  film growth. The device using the  $[hk1]$ -oriented  $\text{Sb}_2\text{S}_3$  thin-film demonstrated a PCE of 6.82%, outperforming the devices based on an  $[hk0]$ -oriented absorber (6.27%). The O-DLTS results suggested that the  $[hk1]$ -oriented  $\text{Sb}_2\text{S}_3$  film (with normal grain growth) exhibited only one type of hole trap ( $S_{\text{sb}}$ ), while the  $[hk0]$ -oriented film (with abnormal grain growth) demonstrated two types of hole traps ( $S_{\text{sb}}$  and  $V_{\text{sb}}$ ). The dark  $J$ - $V$  measurements revealed a higher leakage current in the latter device. Furthermore, the TAS analysis suggested a shorter average lifetime ( $t_{\text{avg}}$  of  $\sim 3924$  ps) in the device with  $[hk1]$ -oriented  $\text{Sb}_2\text{S}_3$  films than that with the  $[hk0]$ -oriented films ( $t_{\text{avg}}$  of  $\sim 4818$  ps). The shorter lifetime and better performance were attributed to the rapid electron extraction by the ETL (CdS) layer and lower density of deep traps in the devices with the  $[hk1]$ -oriented  $\text{Sb}_2\text{S}_3$  films than their  $[hk0]$  counterparts.

Chen *et al.*<sup>10</sup> demonstrated a solvent-assisted hydrothermal deposition (SHD) technique for the deposition of high-quality  $\text{Sb}_2(\text{S,Se})_3$  films. Ethanol was used as an additive and found to be instrumental in regulating the reaction kinetics during the deposition procedure and stoichiometry in  $\text{Sb}_2(\text{S,Se})_3$  films. The devices fabricated using the ethanol-SHD technique demonstrated a substantial reduction in trap density and dark current density and higher recombination resistance ( $R_{\text{rec}}$ ) than the control devices (Fig. 7c). The FTO/ $\text{Sb}_2(\text{S,Se})_3$ /Au structure was adopted to measure the conductivity of the  $\text{Sb}_2(\text{S,Se})_3$  films, revealing an improvement in conductivity from  $1.65 \times 10^{-4} \text{ S cm}^{-1}$  to  $4.07 \times 10^{-4} \text{ S cm}^{-1}$  on adopting the novel route. The deep level defects in the films were probed using O-DLTS. Two-hole traps (H1 and H2) were obtained in the control films, which were attributed to  $\text{Sb}_{\text{S}_2}$  and  $\text{Sb}_{\text{S}_1}$  defects, respectively (Fig. 7a). The films synthesized using the SHD protocol exhibited one order of

suppression in trap density of H1 and H2 defects, greatly benefiting the carrier transport. Consequently, the fabricated device (FTO/CdS/ $\text{Sb}_2(\text{S,Se})_3$ /spiro-OMeTAD/Au) delivered a PCE of 10.75%, a record for  $\text{Sb}_2\text{X}_3$  solar cells.

Li *et al.*<sup>85</sup> investigated the role of potassium iodide (KI) surface treatment on  $\text{Sb}_2(\text{S,Se})_3$  thin films and solar cells. KI treatment was found to successfully improve the crystallinity and morphology of the  $\text{Sb}_2(\text{S,Se})_3$  films, inhibit the deep-level defects and improve the band alignment for efficient charge transport. It was revealed that I diffused into the  $\text{Sb}_2(\text{S,Se})_3$  crystal and induced the benign  $[211]$ -crystal orientation. The as-deposited films were Sb-rich, with  $\text{Sb}_{\text{S}}$  and  $\text{Sb}_{\text{Se}}$  antisites as the two dominant defects. The diffused I interacted strongly with the Sb atoms, forming Sb–I bonding. Therefore, I passivated the  $V_{\text{S/Se}}$  defects, and thus inhibited the formation of the detrimental  $\text{Sb}_{\text{S}}$  and  $\text{Sb}_{\text{Se}}$  antisite defects. Fig. 8(a) illustrates the DLTS signal from the control and KI-treated device. Two majority (electron) traps and one minority (hole) trap were identified for the control device, which were denoted as E1 ( $\text{Sb}_{\text{Se}2}$ ) and E2 ( $\text{Sb}_{\text{Se}1}$ ) and H1 ( $\text{Sb}_{\text{S}}$ ), respectively. KI treatment suppressed the density of E1 and H1, while annihilating the E2 defect (defect parameters summarized in Table 1). Furthermore, the KI-treated device exhibited a lower dark (reverse) current density and higher PCE (9.22%) than its untreated counterpart (PCE of 8.19%).

Wen *et al.*<sup>90</sup> demonstrated a vapor transport deposition (VTD) technique to deposit  $\text{Sb}_2\text{Se}_3$  thin films with improved film crystallinity, lengthier carrier lifetimes, and reduced bulk and interfacial defects. The champion device (ITO/CdS/ $\text{Sb}_2\text{Se}_3$ /Au) demonstrated a PCE of 7.6%, outperforming the device fabricated *via* RTE (PCE of 5.6%). As envisaged in Fig. 9(c), the power value ( $\alpha$ ) for both the devices was  $\sim 0.9$ , *i.e.*, closer to unity (first-order), suggesting that trap-assisted SRH recombination is the dominant loss mechanism. The diode ideality factor ( $\eta$ ) decreased from 1.51 (for RTE-fabricated device) to 1.23 (for VTD-fabricated device), suggesting the mitigation of SRH recombination on switching the film deposition route (Fig. 9d). The CdS/ $\text{Sb}_2\text{Se}_3$  interfacial defect density was quantitatively characterized *via* the capacitance–voltage ( $C$ - $V$ ) profiling and deep-level capacitance profiling (DLCP) techniques. The interfacial defect density of the RTE-fabricated and VTD-fabricated devices was calculated to be  $2.1 \times 10^{11} \text{ cm}^{-2}$  and  $2.8 \times 10^{10} \text{ cm}^{-2}$ , respectively. Interestingly, during VTD deposition, Cd diffusion into the  $\text{Sb}_2\text{Se}_3$  film transformed p-type  $\text{Sb}_2\text{Se}_3$  into n-type, forming a buried homojunction at the CdS/ $\text{Sb}_2\text{Se}_3$  interface. This led to suppressed interfacial defects, benefiting the device performance. The DLTS study revealed one electron trap (E1) and two hole traps (H1 and H2) in the  $\text{Sb}_2\text{Se}_3$  films. The E1, H1 and H2 defects were attributed to the  $\text{Sb}_{\text{Se}}$ ,  $V_{\text{Sb}}$  and  $\text{Se}_{\text{Sb}}$  defects, respectively. The VTD deposition process suppressed the density of the  $[\text{Sb}_{\text{Se}} + \text{Se}_{\text{Sb}}]$  antisite defect pair by more than an order of magnitude. As envisaged in Fig. 9(h), the H1 state was under  $E_F$  and submerged in electrons, and thus the H1 defects always remained inert. In contrast, the H2 and E1 traps were above the  $E_F$  and active in trapping holes and electrons, respectively. Moreover, owing



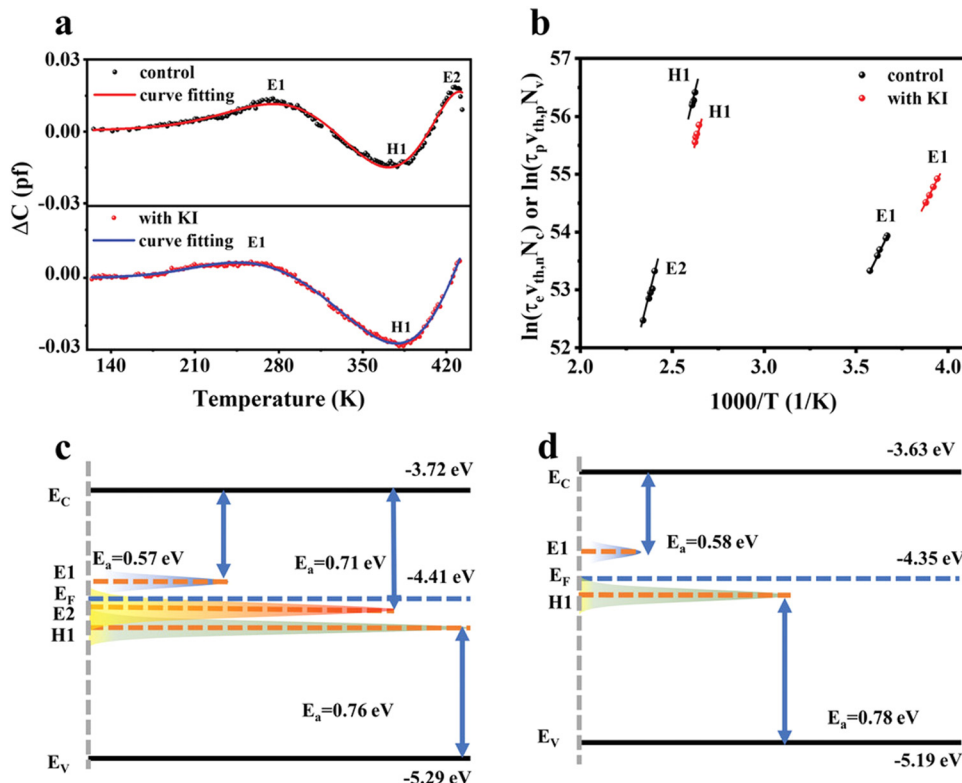


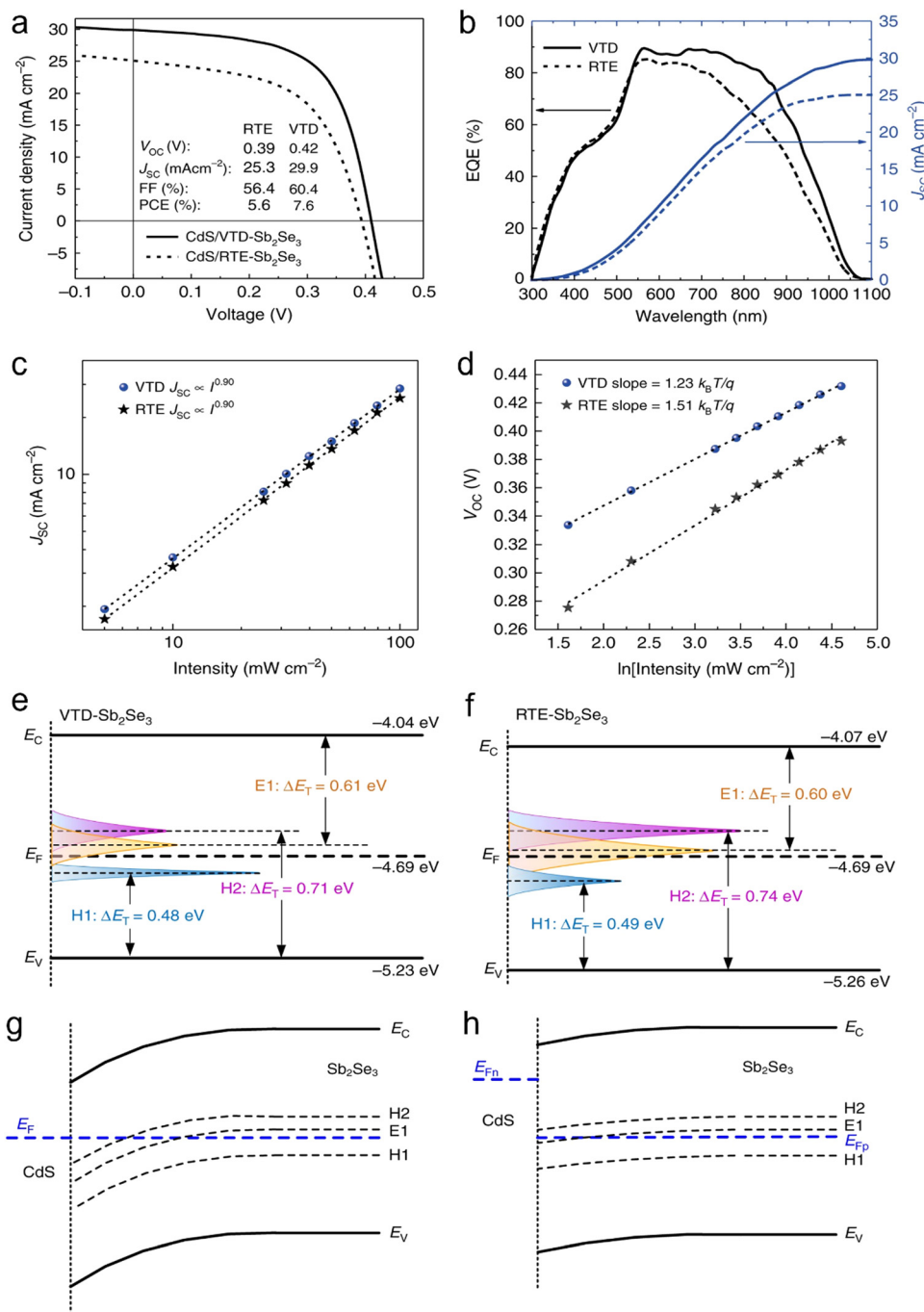
Fig. 8 (a) DLTS signals for the device without and with KI post-treatment. (b) Corresponding Arrhenius plots obtained from the DLTS signals. Position distribution of CBM ( $E_C$ ), VBM ( $E_V$ ), Fermi level ( $E_F$ ) and defect energy level for  $Sb_2(S,Se)_3$  films (c) without and (d) with KI post-treatment. Adapted with permission from ref. 85 Copyright 2022, Wiley-VCH.

to the higher defect density of H2 and E1 than the carrier concentration ( $\sim 10^{13} \text{ cm}^{-3}$ ) in the  $Sb_2Se_3$  layer, the (hole) quasi  $E_F$ -level ( $E_{FP}$ ) would be more likely to be pinned near the E1 and H2 levels.

Ma *et al.*<sup>92</sup> investigated the influence of (0–4%) Te and S doping on  $Sb_2Se_3$  films and investigated their deep-defect profiles using DLTS. It was revealed that the incorporation of Te was decisive in regulating the atomic ratio of Se/Sb, given that a sustained decline in Se content was observed with an increase in Te doping. Different crystal structures of  $Sb_2Se_3$  and  $Sb_2Te_3$  inhibited the formation of the alloy- $Sb_2(Se,Te)_3$  film, and therefore instead of going into the lattice, Te settled at the grain boundaries and surface. Se-rich  $Sb_2Se_3$  favored the formation of  $Se_{Sb}$  and  $V_{Sb}$  defects, while the Sb-rich films benefited the formation of  $Sb_{Se}$  and  $V_{Se}$  defects. The undoped  $Sb_2Se_3$  demonstrated two-hole traps (H1 and H2) and one-electron trap (E1), acting as acceptor and donor defects, respectively. These defects were attributed to the antimony vacancy ( $V_{Sb}$ ), selenium antisite ( $Se_{Sb}$ ) and antimony antisite ( $Sb_{Se}$ ) defects, respectively. Te-doping suppressed the hole-trap densities by eliminating the  $Se_{Sb}$  defects (H2) and mitigating the  $V_{Sb}$  defects (H1). However, for a high Te doping ( $\sim 4\%$ ),  $Sb_{Se}$  (E1) exhibited an order of increment in magnitude, and a new electron trap (E2,  $V_{Se}$ ) also started to appear in the DLTS signal. Interestingly, S-doping was found to have detrimental effects on the  $Sb_2Se_3$  film and solar cell performance, given that S-incorporation could not emulate Te in suppressing the defect-formation process.

Duan *et al.*<sup>91</sup> demonstrated a novel injection vapor deposition (IVD) technique for the deposition of high crystallinity, [001]-oriented  $Sb_2Se_3$  films, with a minimal deep-level defect density. The  $Sb_2Se_3$ -based devices fabricated *via* IVD demonstrated a lower trap density, suppressed recombination losses and better performance (PCE of 10.12%) than the devices fabricated *via* close-spaced sublimation (CSS) and co-evaporation (CoE) method (PCE of 9.31% and 3.96%, respectively). Capacitance–voltage ( $C$ – $V$ ) profiling, deep-level capacitance profiling (DLCP) measurements, and minority carrier injection DLTS measurements were carried out to probe the interface and deep traps in the  $Sb_2Se_3$  solar cells. Subsequently, bulk trap density ( $N_T$ ) values of  $2.3 \times 10^{16}$ ,  $1.5 \times 10^{16}$ , and  $5 \times 10^{16} \text{ cm}^{-3}$  were obtained for the devices with IVD- $Sb_2Se_3$ , CSS- $Sb_2Se_3$  and CoE- $Sb_2Se_3$  films, respectively. The corresponding devices exhibited an interface defect density of  $1.32 \times 10^{12}$ ,  $9.30 \times 10^{12}$  and  $1.20 \times 10^{12} \text{ cm}^{-2}$ , respectively. Furthermore, the DLTS analysis revealed one (type) electron trap (E1), two electron traps (E1 and E2), and two electron traps and one hole trap (E1, E2, and H1) in the CSS- $Sb_2Se_3$ , IVD- $Sb_2Se_3$ , and CoE- $Sb_2Se_3$  films (Fig. 10a–c), respectively. The IVD-E2, CSS-E1, and CoE-E2 defects were attributed to  $Se_{Sb2}$ , while CoE-H1 to the  $Se_{Sb1}$  antisite traps. The capture cross section for the IVD-E1 trap was only  $\sim 10^{-21} \text{ cm}^2$ , which was more than two orders of magnitude smaller than that for the other traps. It also exhibited a lower trap density of  $2.74 \times 10^{14} \text{ cm}^{-3}$ , thus having a less detrimental effect on the charge transport and device





**Fig. 9** (a)  $J$ - $V$  characteristics and (b) EQE plots of the VTD- and RTE-fabricated  $\text{Sb}_2\text{Se}_3$  solar cells. Light intensity-dependent (c)  $J_{SC}$  and (d)  $V_{OC}$  of the devices. Energy states and defect level of (e) VTD-fabricated and (f) RTE-fabricated devices. Energy band diagrams at CdS/ $\text{Sb}_2\text{Se}_3$  interface (g) in the dark and (h) under illumination. Reproduced under the guidelines of the Creative Commons CC-BY license from ref. 90 Copyright 2018, Springer Nature.

parameters. Interestingly, it was observed that during the CSS deposition of the  $\text{Sb}_2\text{Se}_3$  absorber layers, the actual temperature of the substrate reached as high as 370–380 °C. The contribution of source radiative heating to the surface thermal energy became significant and [001]-oriented  $\text{Sb}_2\text{Se}_3$  nanorods could be obtained in tens of seconds owing to the strong anisotropy in chemical-bonding. In contrast to CSS, the surface thermal energy in IVD is mainly attributed to the heating of the substrate. According to the terrace-ledge-kink (TLK) growth

model, an increase in the growth temperature and surface thermal energy increases the critical radius for crystal growth and mitigates the density of stable nuclei, resulting in the formation of compact, [001]-oriented  $\text{Sb}_2\text{Se}_3$  films with reduced donor-acceptor defect complexes.

Liu *et al.*<sup>86</sup> employed the controlled co-sublimation of mixed  $\text{Sb}_2\text{Se}_3$  and  $\text{Sb}_2\text{S}_3$  powders to deposit  $\text{Sb}_2(\text{S,Se})_3$  films, with a V-shape (graded) bandgap. In contrast to the most widely used hydrothermal deposition method for  $\text{Sb}_2\text{X}_3$  films, CSS



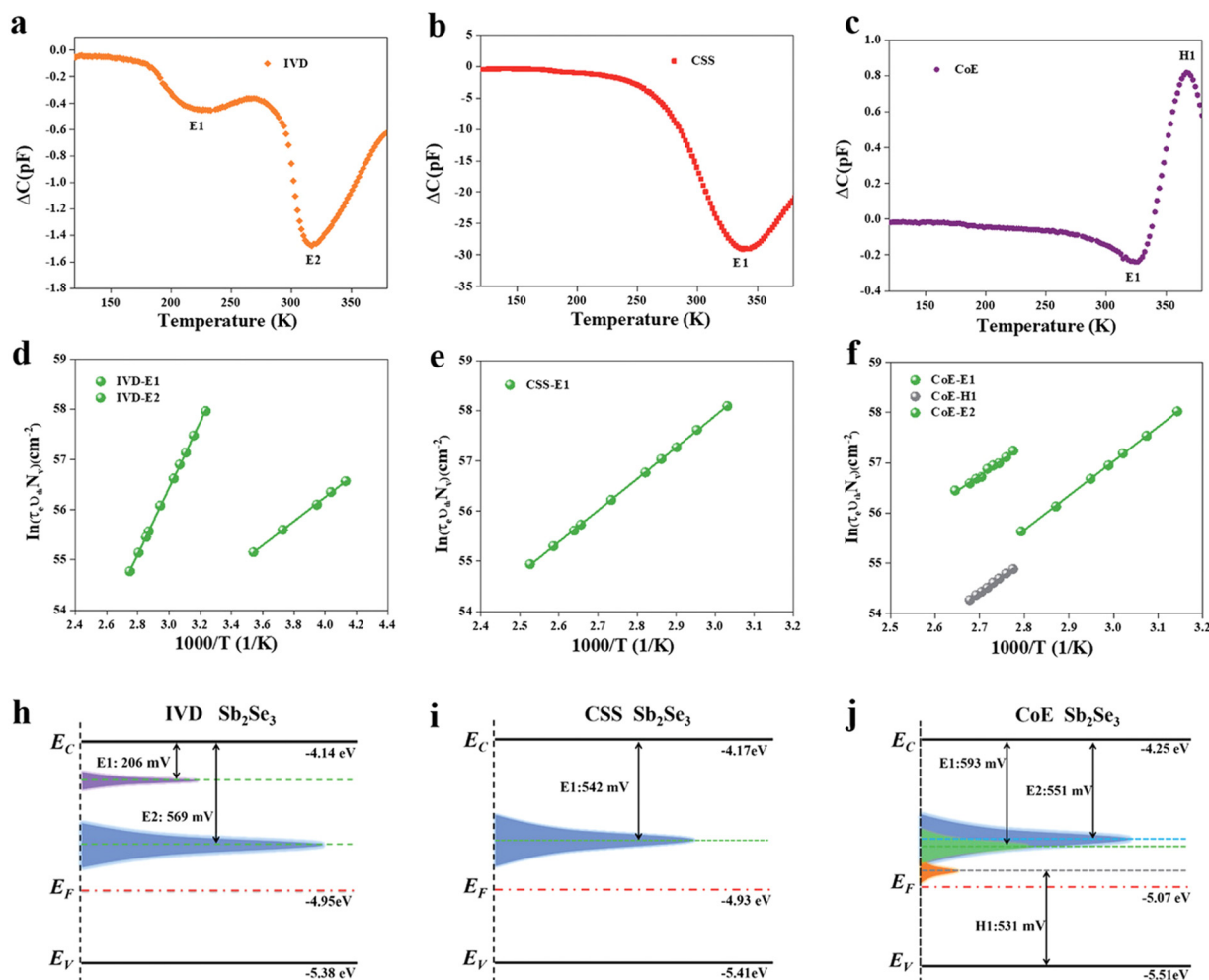


Fig. 10 DLTS analysis of solar cells based on IVD-, CSS-, and CoE-Sb<sub>2</sub>Se<sub>3</sub> absorbers. DLTS signals of (a) IVD-Sb<sub>2</sub>Se<sub>3</sub>, (b) CSS-Sb<sub>2</sub>Se<sub>3</sub>, and (c) CoE-Sb<sub>2</sub>Se<sub>3</sub> solar cells. Arrhenius plots obtained from DLTS signals for (d) IVD-Sb<sub>2</sub>Se<sub>3</sub>, (e) CSS-Sb<sub>2</sub>Se<sub>3</sub>, and (f) CoE-Sb<sub>2</sub>Se<sub>3</sub> solar cells. Deep-level defect profile of (g) IVD-Sb<sub>2</sub>Se<sub>3</sub>, (h) CSS-Sb<sub>2</sub>Se<sub>3</sub>, and (i) CoE-Sb<sub>2</sub>Se<sub>3</sub> solar cells. Adapted with permission from ref. 91 Copyright 2022, Wiley-VCH.

deposition offers the key advantages of fast deposition rates ( $\approx \mu\text{m min}^{-1}$ ), large-area film uniformity, high throughput, and scalability for commercial manufacturing. In this work, instead of the toxic CdS, compact TiO<sub>2</sub> was used as the ETL and record PCE value of 9.02% was obtained for the CdS-free Sb<sub>2</sub>X<sub>3</sub> solar cells. An increment in the S/Se ratio was found to be successful in improving the grain size and favorable  $[hk1]$ -orientations, simultaneously suppressing the detrimental bulk defects in the Sb<sub>2</sub>(S,Se)<sub>3</sub> films. The incorporation of S in Sb<sub>2</sub>(S,Se)<sub>3</sub> enlarged the band gap and  $V_{\text{bi}}$ , both contributing to the enhancement in  $V_{\text{OC}}$ . The TPV decay measurement revealed carrier recombination lifetimes of 5.47 ms, 6.24 ms, and 3.27 ms in the devices employing Sb<sub>2</sub>Se<sub>3</sub>, Sb<sub>2</sub>(S,Se)<sub>3</sub>-47%, and Sb<sub>2</sub>(S,Se)<sub>3</sub>-76% films, respectively. The electron-only devices were fabricated with the architecture of FTO/TiO<sub>2</sub>/Sb<sub>2</sub>(S,Se)<sub>3</sub>/PCBM/Ag, and lower  $V_{\text{TFL}}$  (and  $N_{\text{T}}$ ) was obtained in the Sb<sub>2</sub>(S,Se)<sub>3</sub>-47% based device than its counterparts. The capacitance-frequency analysis revealed  $N_{\text{T}}$  values of  $8.49 \times 10^{15}$  and  $1.08 \times 10^{16} \text{ cm}^{-3}$  in the Sb<sub>2</sub>(S,Se)<sub>3</sub>-47% and Sb<sub>2</sub>(S,Se)<sub>3</sub>-76% films, respectively. The DLTS signal for the Sb<sub>2</sub>(S,Se)<sub>3</sub>-47% device revealed two negative peaks, which

were attributed to two minority (hole) traps, i.e., H1 and H2 (both assigned to Sb<sub>s</sub> antisite defects) (Fig. 11d). Alternatively, the Sb<sub>2</sub>(S,Se)<sub>3</sub>-76% device exhibited one electron trap (E1, S<sub>Sb</sub>) and one hole trap (H2, Sb<sub>S</sub>) (Fig. 11g). Particularly, the E1 and H2 traps were located close to the  $E_{\text{F}}$ -level, and thus highly effective in  $E_{\text{F}}$ -level pinning and concomitant  $V_{\text{OC}}$  deficit. It was speculated that an increase in S-content induced S<sub>Sb</sub> antisite defects. Furthermore, a lower S-content passivated the V<sub>Se</sub> defects and inhibited the Sb<sub>S</sub> defect density in the Sb<sub>2</sub>(S,Se)<sub>3</sub> absorber films.

Han *et al.*<sup>93</sup> investigated the role of post-annealing SbCl<sub>3</sub> treatment on Sb<sub>2</sub>S<sub>3</sub> thin films and the corresponding solar cells (FTO/TiO<sub>2</sub>/Sb<sub>2</sub>S<sub>3</sub>/spiro-OMeTAD/Au). The hole trap density was investigated by analyzing the dark  $J$ - $V$  characteristics of the hole-only device (FTO/PEDOT:PSS/Sb<sub>2</sub>S<sub>3</sub>/spiro-OMeTAD/Au). Fig. 12(e) indicates the linear response of all the hole-only devices at low voltages treated with different concentrations of SbCl<sub>3</sub>. On exceeding the kink point (referred to as trap field limit voltage,  $V_{\text{TFL}}$ ), the current demonstrated a non-linear increase, suggesting that the trap states were filled by the



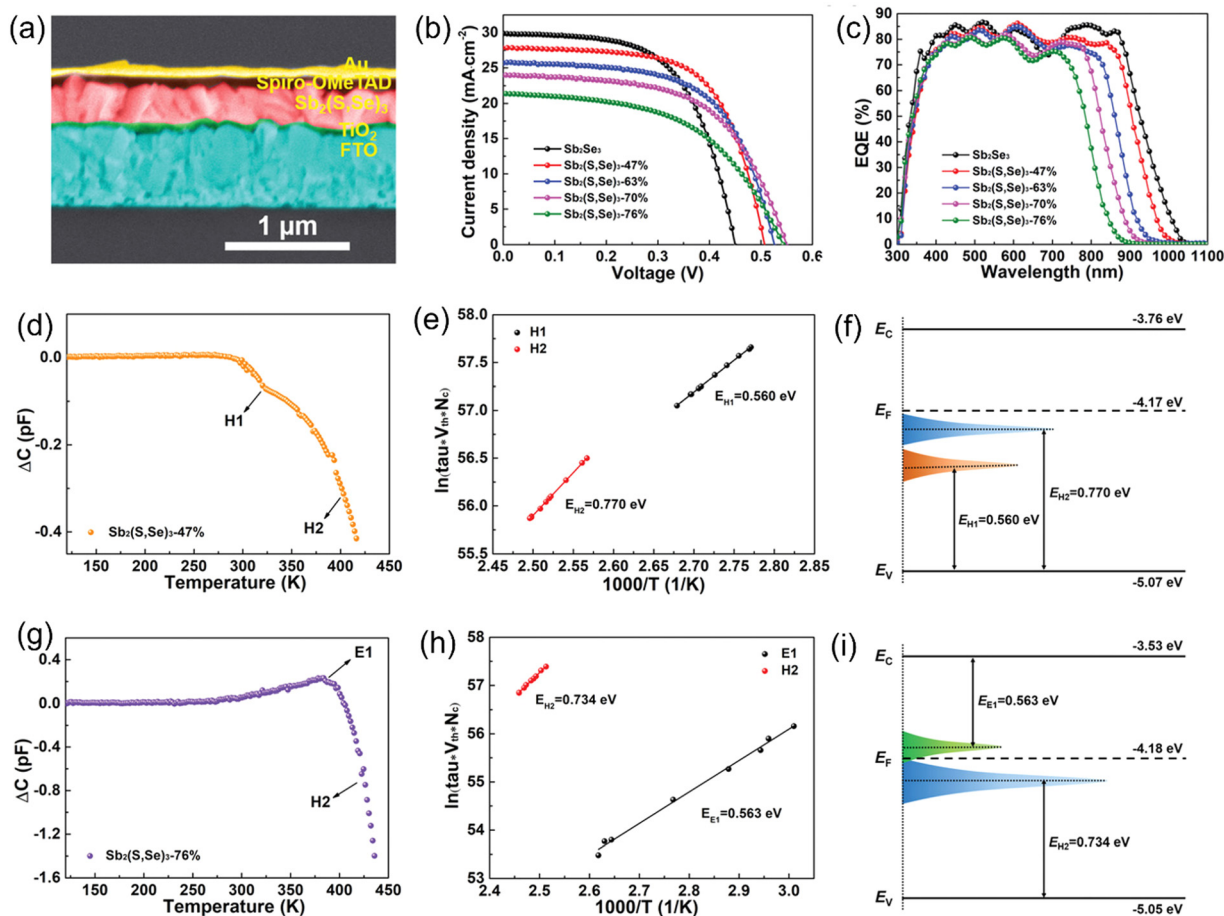


Fig. 11 (a) Cross-sectional SEM image of the  $\text{Sb}_2(\text{S,Se})_3$  solar cells. (b) Current density–voltage ( $J$ – $V$ ) curves and (c) external quantum efficiency (EQE) spectra of the  $\text{Sb}_2(\text{S,Se})_3$  solar cells with different S-contents. (d) and (g) DLTS signals and (e) and (h) Arrhenius plots obtained from the DLTS signals. (f) and (i) Conduction ( $E_C$ ), valence ( $E_V$ ) band edges, Fermi ( $E_F$ ), and defect energy levels of the  $\text{Sb}_2(\text{S,Se})_3$  films with 47% and 76% S-content, respectively. Adapted with permission from ref. 86 Copyright 2022, Wiley-VCH.

injected carriers. The trap density ( $N_T$ ) depends linearly on  $V_{\text{TFL}}$  ( $N_T = \frac{2\epsilon\epsilon_0 V_{\text{TFL}}}{eL^2}$ , where  $e$  is the elementary charge of the electron,  $L$  is the  $\text{Sb}_2\text{S}_3$  film thickness,  $\epsilon$  ( $= 6.67$ ) is the relative permittivity of  $\text{Sb}_2\text{S}_3$ , and  $\epsilon_0$  is the vacuum permittivity). The  $V_{\text{TFL}}$  of the control device was 0.75 V, corresponding to a trap-state density of  $2.81 \times 10^{16} \text{ cm}^{-3}$ . An increase in the  $\text{SbCl}_3$  concentration from 15 to 30  $\text{mg mL}^{-1}$  led to a decrease in  $V_{\text{TFL}}$  from 0.57 to 0.40 V, and  $N_T$  from  $2.13 \times 10^{16}$  to  $1.50 \times 10^{16} \text{ cm}^{-3}$ . A further increase in  $\text{SbCl}_3$  concentration to 60  $\text{mg mL}^{-1}$  led to an increase in  $V_{\text{TFL}}$  to 0.65 V and  $N_T$  to  $2.43 \times 10^{16} \text{ cm}^{-3}$ . The TPC and TPV decay were analyzed under short-circuit conditions and open-circuit voltage to reveal the influence of  $\text{SbCl}_3$  treatment on the charge transfer and recombination lifetimes, respectively. Compared to the control device (1.03  $\mu\text{s}$ ), the devices treated with 30  $\text{mg mL}^{-1}$   $\text{SbCl}_3$  demonstrated the most efficient charge transfer (0.31  $\mu\text{s}$ ). The same device also exhibited a lengthier carrier recombination lifetime (89.34  $\mu\text{s}$ ) than the control devices (28.35  $\mu\text{s}$ ). The ideality factor ( $\eta$ ) for the control device (1.71) dropped to the lowest value of 1.59 in the case of 30  $\text{mg mL}^{-1}$   $\text{SbCl}_3$  treatment. All these results unambiguously suggest inhibited SRH recombination in the device with

30  $\text{mg mL}^{-1}$   $\text{SbCl}_3$  treatment, delivering a PCE of 7.10% (outperforming the control device with a PCE of 4.37%). Li *et al.*<sup>94</sup> fabricated electron-only devices (FTO/buffer layer/ $\text{CdS}/\text{Sb}_2\text{S}_3/\text{PCBM}/\text{Ag}$ ) to estimate the (electron) trap density in the  $\text{Sb}_2\text{S}_3$  layer.  $V_{\text{TFL}}$  values of 0.69, 0.39, and 0.29 V were calculated for the devices without, with  $\text{SnO}_2$ , and with Mg-doped  $\text{SnO}_2$  as the buffer layers, respectively. Estimated trap densities of  $5.66 \times 10^{15}$ ,  $3.2 \times 10^{15}$  and  $2.38 \times 10^{15} \text{ cm}^{-3}$  were obtained in the corresponding devices, respectively. The reduction in trap states was attributed to the lowering of the lattice mismatch at  $\text{CdS}/\text{Sb}_2\text{S}_3$  on introducing the buffer layers and suppression of the grain boundary defects.

Hu *et al.*<sup>88</sup> systematically identified and characterized the defects in  $\text{Sb}_2\text{Se}_3$  solar cells using temperature-dependent admittance spectroscopy measurements. DC reverse bias-dependent admittance measurements (in the frequency range 20–10<sup>5</sup> Hz) revealed the presence of three types of defects with energy levels varying in the range of 0.3–0.4, 0.2–0.6 and 0.5–0.6 eV above the VBM, which were denoted as D1, D2, and D3, respectively. D1 and D2 were identified as bulk defects, while D3 as an interface or near-interface defect. D1 exhibited a large cross-section, but low trap density, leading to moderate ( $\sim \mu\text{s}$ )



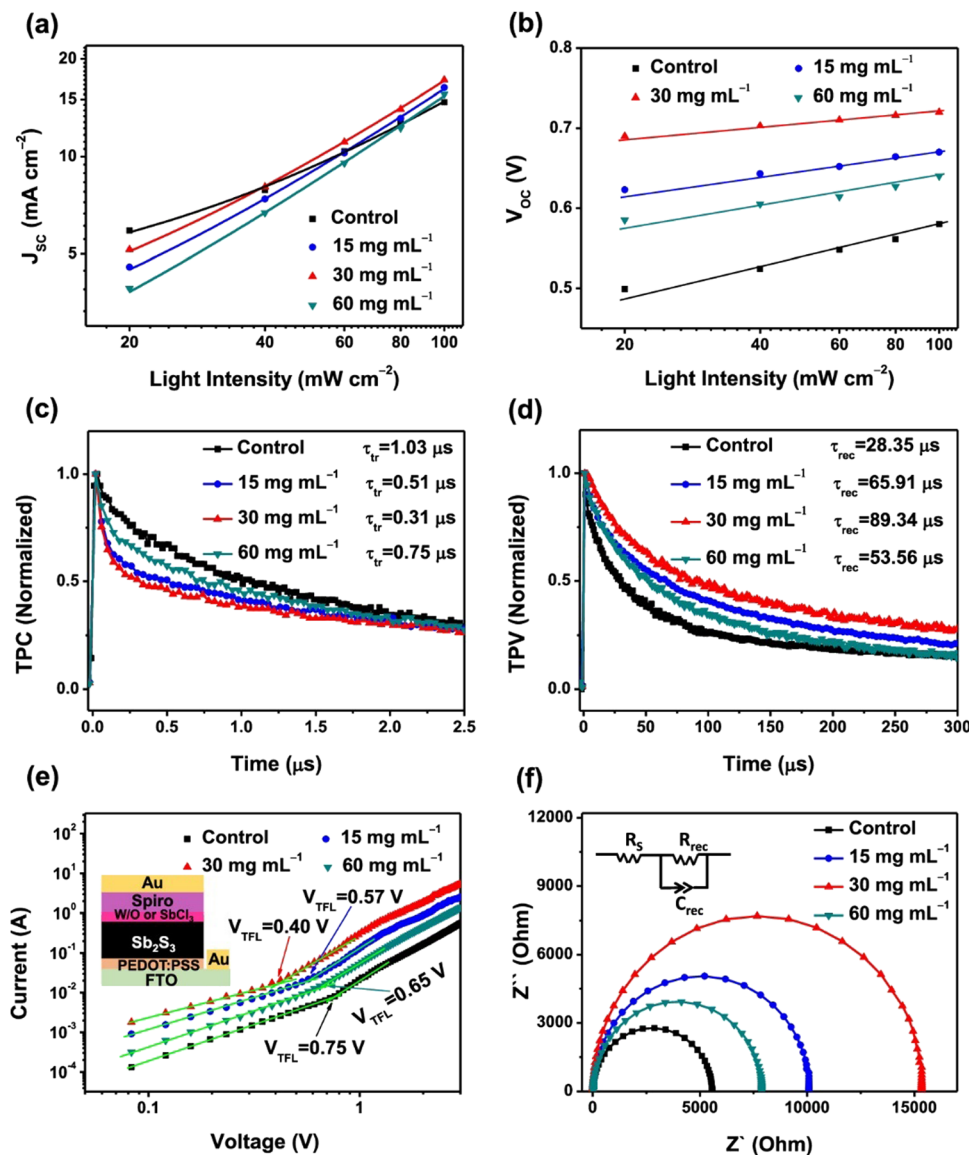


Fig. 12 (a)  $J_{\text{sc}}$  and (b)  $V_{\text{oc}}$  as a function of light intensity in the control  $\text{Sb}_2\text{S}_3$  and  $\text{SbCl}_3$ -treated  $\text{Sb}_2\text{S}_3$  devices with different  $\text{SbCl}_3$  concentrations. (c) TPC and (d) TPV decay characteristics of the control  $\text{Sb}_2\text{S}_3$  and  $\text{SbCl}_3$ -treated  $\text{Sb}_2\text{S}_3$  devices. (e) Dark  $J$ - $V$  measurements of the hole-only devices, revealing  $V_{\text{TFL}}$  kink point behavior. (f) Nyquist plots for the devices measured in the dark at an applied voltage close to the  $V_{\text{oc}}$  (inset: equivalent circuit diagram employed to fit the Nyquist plots). Adapted with permission from ref. 93 Copyright 2020, the American Chemical Society.

carrier lifetimes for holes. D2 was assigned as a shallow acceptor level, with an insignificant capture cross-section ( $10^{-20} \text{ cm}^2$ ) and lengthier ( $\sim \text{ms}$ ) carrier-lifetimes for holes. D3 exhibited a large cross-section for holes, in conjugation with high defect density, leading to short lifetimes for holes ( $\sim \text{ns}$ ), thus serving as a prominent recombination center. Pan *et al.*<sup>87</sup> performed quantitative investigation of the majority carrier defect density in all vacuum-deposited  $\text{Sb}_2(\text{S,Se})_3$  solar cells *via* the admittance spectroscopy technique. The  $E_a$  values of the devices with CSS- and CBD-CdS were 203 and 364 meV, respectively. Two-hole traps (H1 and H2) were identified, which were attributed to  $V_{\text{Se}}$  and  $\text{Sb}_{\text{Se}}$  defects, respectively. The  $\text{Sb}_2(\text{S,Se})_3$  device with CBD-CdS (PCE of 5.05%) exhibited a higher defect density, capture cross-section and shorter capture lifetime than its CSS-CdS counterpart (PCE of 7.12%). According to this study,

it was concluded that the PCE of state-of-the-art  $\text{Sb}_2\text{X}_3$  solar cells can be improved by adopting higher crystallinity CSS-CdS films as the ETL instead of the widely used CBD-CdS films.

Li *et al.*<sup>95</sup> fabricated an  $\text{Sb}_2\text{Se}_3/\text{CdS}$  core-shell structured nano-array solar cell, with a record FF of 70.3% and PCE of 9.2% (Fig. 13d). Briefly, [001]-oriented  $\text{Sb}_2\text{Se}_3$  nano-arrays were deposited using CSS and a four-step model was proposed for the growth of  $\text{Sb}_2\text{Se}_3$  nanorods on an Mo/MoSe<sub>2</sub> substrate including surface absorption, film growth, splitting, and nanorod array growth stage. DLCP measurements suggested that the  $\text{Sb}_2\text{Se}_3$  films deposited using CSS had a lower bulk-defect density ( $N_d$  of  $\sim 2 \times 10^{15} \text{ cm}^{-3}$ ) than the thermally evaporated films ( $N_d$  of  $\sim 2 \times 10^{17} \text{ cm}^{-3}$ ). Capacitance-voltage ( $C$ - $V$ ) profiling suggested that the  $\text{Sb}_2\text{Se}_3/\text{CdS}$  core-shell interface was much more defective (interface defect density of

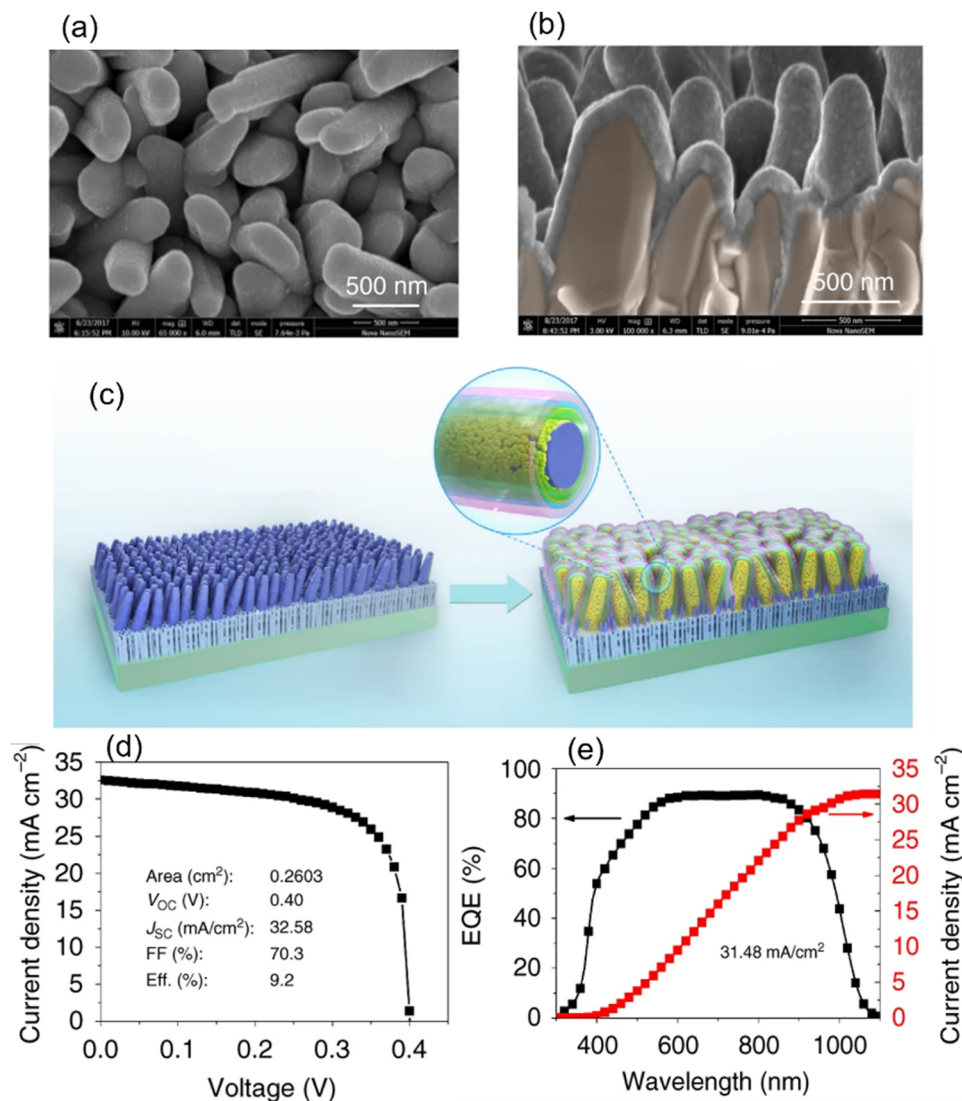


Fig. 13 (a) Top-view and (b) cross-sectional SEM images of CdS buffer deposited on Sb<sub>2</sub>Se<sub>3</sub> nanorod arrays. (c) Schematic of the Sb<sub>2</sub>Se<sub>3</sub> nanorod arrays on Mo-coated glass and fabricated Sb<sub>2</sub>Se<sub>3</sub>/CdS (core/shell) nanorod array solar cells. (d) *J*–*V* curve and (e) EQE spectrum and integrated *J*<sub>SC</sub> curve of the champion device. Adapted under the guidelines of the Creative Commons BB CY license from ref. 95 Copyright 2019, Springer Nature.

$\sim 3 \times 10^{12} \text{ cm}^{-2}$ ) than the planar CdS/Sb<sub>2</sub>Se<sub>3</sub>, ZnO/Sb<sub>2</sub>Se<sub>3</sub> or TiO<sub>2</sub>/Sb<sub>2</sub>Se<sub>3</sub> interfaces in the superstrate configurations. An ultrathin ALD-grown TiO<sub>2</sub> layer was used at the Sb<sub>2</sub>Se<sub>3</sub>/CdS interface to passivate the surface defects in the Sb<sub>2</sub>Se<sub>3</sub> layer and suppress the detrimental current-leakage paths. Furthermore, as illustrated in Fig. 13(e), the champion device exhibited an excellent EQE (>85%, between 550 and 900 nm), suggesting efficient long-range carrier transport along the [001]-direction.

## 4. Strategies toward high PCE Sb<sub>2</sub>X<sub>3</sub> solar cells

### 4.1. Deep-level defect passivation and *E<sub>F</sub>*-level depinning in photoactive Sb<sub>2</sub>X<sub>3</sub> films

Mitigating deep-defects (particularly that with a large capture cross section) is a crucial strategy to elevate the minority carrier

lifetime and boost the *V*<sub>OC</sub> (and PCE) of fabricated Sb<sub>2</sub>X<sub>3</sub> solar cells. A short lifetime ( $\sim 0.01$ –1 ns) and long lifetime ( $\sim 5$ –60 ns) are commonly observed in Sb<sub>2</sub>X<sub>3</sub> solar cells.<sup>29</sup> The former is attributed to interface recombination, while the latter to bulk recombination. Yang *et al.*<sup>34</sup> studied the (longer) carrier lifetimes in an Sb<sub>2</sub>S<sub>3</sub> thin film and single crystal using TAS and concluded that the longer lifetime originates from the self-trapped exciton (STE) owing to the soft lattice in Q-1D-structured Sb<sub>2</sub>X<sub>3</sub>. The device PCE (and *V*<sub>OC</sub>) was constrained by the shortest lifetimes. The simulation results revealed that 2 orders of magnitude improvement in carrier lifetime can gain an  $\sim 150$  mV increase in *V*<sub>OC</sub> in Sb<sub>2</sub>X<sub>3</sub> solar cells, as illustrated in Fig. 14.<sup>29</sup>

As shown in Fig. 15(a), a high density of multiple types of deep traps exists in Sb<sub>2</sub>X<sub>3</sub> thin films. These traps serve as scavengers for photogenerated carriers and accelerate trap-assisted SRH recombination, and thus detrimental to the





Fig. 14 Simulated  $V_{OC}$  of  $Sb_2S_3$  and  $Sb_2Se_3$  solar cells as a function of carrier lifetime by SCAPS. The shadow area exhibits the generally accepted lifetimes and  $V_{OC}$  of the  $Sb_2S_3$  and  $Sb_2Se_3$  solar cells. Adapted with permission from ref. 29 Copyright 2020, the American Chemical Society.

performance of  $Sb_2X_3$  solar cells. Solution-processed  $Sb_2X_3$  films usually contain a minuscule amount of oxide impurities (mainly  $Sb_2O_3$ ). These oxide impurities form deep traps, and thus limit the  $V_{OC}$  values. In this case, given that  $Sb_2O_3$  is recognized as a center of electron-hole pair recombination,  $Sb_2O_3$ -free,  $Sb_2X_3$  films are advantageous for achieving high-performance  $Sb_2X_3$  solar cells. The  $Sb_2X_3$  films synthesized using vacuum-assisted techniques are relatively oxide free and exhibit a higher degree of phase purity.<sup>96</sup> However, recent studies suggest that controlled  $O_2$ -incorporation can help in passivating bulk and interface defects, improving the device performance.<sup>97</sup> In addition, *in situ* sulfurization,<sup>98,99</sup> post-deposition thioacetamide (TA) treatment,<sup>78</sup> ammonium sulfide  $((NH_4)_2S)$  treatment,<sup>77</sup> thiourea (TU) treatment,<sup>5</sup> selenization,<sup>100–102</sup> alkali-metal doping,<sup>103–105</sup> and alcohol-vapor annealing<sup>106</sup> are key strategies to passivate deep-defects and boost the PCE of  $Sb_2X_3$  solar cells.

Another issue with  $Sb_2X_3$  solar cells is Fermi level ( $E_F$ ) pinning by the traps located near the mid-bandgap. Lian *et al.*<sup>74</sup> investigated the deep-level defects in Sb-rich and S-rich  $Sb_2S_3$  films *via* O-DLTS. It was revealed that S-vacancy ( $V_S$ ) and Sb-antisite ( $Sb_S$ ) are the two main electron traps in S-deficient films. These traps have a high density and capture cross-section and contribute to  $E_F$ -level pinning. The electron quasi- $E_F$  level is pinned by these traps, resulting in the inefficient extraction of trapped photo-excited carriers. S-rich films have two main hole traps, which have been identified as Sb-vacancy ( $V_{Sb}$ ) and S-antisite ( $S_{Sb}$ ). These traps are not electrically active compared to that in S-deficient films. Based on a first principles study, Cai *et al.*<sup>75</sup> and Jiang *et al.*<sup>107</sup> concluded that the formation energies of the dominant acceptor defects ( $V_{Sb2}$ ,  $Sb_{S2}$ , and  $S_{Sb2}$ ) and donor defect ( $V_{S2}$ ) in intrinsic  $Sb_2S_3$  are similar. This results in  $E_F$ -level pinning at the near-middle of the band gap. These results are in agreement with the experimentally observed high resistivity in undoped  $Sb_2S_3$  thin films.<sup>108</sup> Wen *et al.*<sup>90</sup> attributed the hole quasi  $E_F$ -level pinning in  $Sb_2Se_3$  to the  $[Sb_{Se} + Se_{Sb}]$  anti-site complex. In general, these

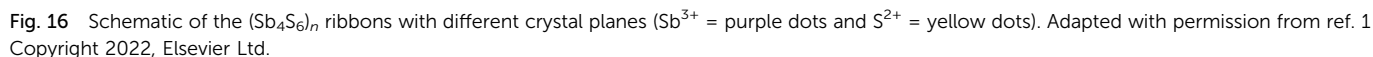
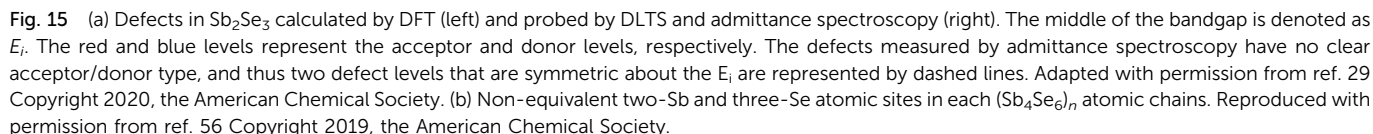
defects exhibit a higher density than the carrier density ( $\sim 10^{13} \text{ cm}^{-3}$ ) in  $Sb_2Se_3$  films. These defects are located slightly above the intrinsic  $E_F$ -level, and thus highly active trap centers; in contrast, the traps below the intrinsic  $E_F$ -level are occupied, and thus inactive. The VTD process has been found to be more effective in depositing  $Sb_2X_3$  films with an improved density of  $[Sb_{Se} + Se_{Sb}]$  anti-site pairs, thus depinning the hole quasi  $E_F$ -level and inducing greater splitting of the quasi-  $E_F$ -levels during illumination.  $Sb_2X_3$  films should be made anion (S,Se) rich to passivate the bulk defects and promote the charge transport efficiency.

## 4.2. Crystallographic orientation engineering

Owing to their unique quasi-1D crystal structure and the associated anisotropy in carrier mobility, the orientation of crystallites/ribbons (schematically illustrated in Fig. 16) plays a decisive role in the device performance. In this case,  $[hk1]$ -oriented  $Sb_2X_3$  films offer excellent carrier transport, suppressed bulk recombination, and trap-free grain boundaries, facilitating high-PCE devices. In this regard, the growth of  $Sb_2X_3$  should be regulated from the 2D (layer-like) mode to 3D (island-like) mode. Liang *et al.*<sup>109</sup> demonstrated the growth of highly  $[hk1]$ -oriented  $Sb_2Se_3$  nanorod arrays perpendicular to a selenized tungsten (W) back-contact in substrate-structured solar cells. The resulting ultrathin film of tungsten selenide ( $WSe_2$ ) also acted as the HTL, promoting efficient hole extraction, and consequently leading to a PCE of 8.46%. Wang *et al.*<sup>110</sup> reported the preparation of highly  $[221]$ -oriented  $Sb_2Se_3$  films on a randomly oriented ZnO buffer layer. The ZnO layer was deposited *via* a green and scalable spray pyrolysis technique, and the device demonstrated a low density of interface defects and highly improved efficiency and stability. Zhou *et al.*<sup>111</sup> employed a simple and fast ( $\sim 1 \mu\text{m min}^{-1}$ ) RTE process to orient the crystal growth perpendicular to the substrate and produced an  $Sb_2Se_3$  thin-film solar with a certified PCE of 5.6%. Li *et al.*<sup>112</sup> introduced a seed screening strategy to deposit  $[211]$ - and  $[221]$ -oriented  $Sb_2Se_3$  films on a  $TiO_2$  film. It was revealed that lying seeds (with  $[hk0]$ -orientation) bonded with the substrate through weak vdW forces, while the standing seeds ( $[hk1]$ -orientation) bonded with the substrate by covalent bonds. The  $[hk0]$ -oriented seeds easily re-evaporated from the  $TiO_2$  surface at high substrate temperature, leaving tightly bonded  $[hk1]$ -oriented seeds and assisting the growth of  $[221]$ - and  $[211]$ -oriented  $Sb_2Se_3$  grains. Liu *et al.*<sup>113</sup> revealed that devices based on thermally deposited hexagonal-CdS outperformed that with solution-processed cubic-CdS. The former was found to be more conducive to the growth of  $[221]$ -oriented, compact, and less-defective  $Sb_2X_3$  films.

Park *et al.*<sup>114</sup> manipulated the microstructure (morphology and crystallite orientation) of the  $Sb_2Se_3$  absorber layer by regulating the substrate temperature during co-evaporation. Consequently,  $[hk1]$ -oriented  $Sb_2Se_3$  nanopillar arrays were obtained at 315 °C. According to this study, it was concluded that a narrow window exists for the substrate temperature, where vertical  $[hk1]$ -oriented  $Sb_2Se_3$  rod arrays can be grown. The  $[hk1]$ -oriented crystallites assist in maximizing the carrier





carrier transport from the inter-ribbon (hopping) to intra-ribbon (band-like) mode, improving the carrier collection efficiency. Notably, the  $[hk1]$ -oriented  $\text{Sb}_2\text{X}_3$  films not only ensured better charge transport, but also bonded strongly with the substrate below (*via* covalent bonds), in contrast to the  $[hk0]$ -oriented films (where ribbons lie flat on the substrate surface and weak van der Waals bonding links the substrate and  $\text{Sb}_2\text{X}_3$ ).

Zhao *et al.*<sup>9</sup> demonstrated an outstanding improvement in average grain size in CBD Sb<sub>2</sub>Se<sub>3</sub> films (from 170 to 435 nm) by using additives to manipulate the reaction kinetics and deposition process. Wang *et al.*<sup>117</sup> developed a fast chemical approach (FCA) for spin coating Sb<sub>2</sub>S<sub>3</sub> films with an average grain size exceeding 12 μm. Han *et al.*<sup>106</sup> revealed that alcohol vapor (ethanol, methanol, and isopropanol) post-annealing

induces a significant enhancement in the crystallinity, morphology and carrier lifetimes of  $\text{Sb}_2\text{S}_3$  films (as illustrated in Fig. 17). This annealing strategy improved the charge collection efficiency and PCE of  $\text{Sb}_2\text{S}_3$  solar cells. Alcohol vapor induced solvent stitching *via* partial dissolution of the GBs, creating a melt phase, which acted as a “stitching adhesive” for enhanced grain growth.

#### 4.4. Selection of ETL and HTL materials and (ETL/ $\text{Sb}_2\text{X}_3$ , $\text{Sb}_2\text{X}_3$ /HTL) interface engineering

All the high-PCE- $\text{Sb}_2\text{X}_3$  solar cells exclusively use the toxic CdS as the ETL and expensive and long-term unstable Spiro-OMeTAD as the HTL (Table S1, ESI†). In this regard, there have been serious attempts to either substitute CdS and Spiro-OMeTAD completely or reduce their usage. The devices based on an oxide ETL ( $\text{TiO}_2$ ,  $\text{SnO}_2$ , and  $\text{ZnO}$ ) underperform (showing an inferior PCE) than their CdS-based counterparts. Doping  $\text{O}^{118}$  and  $\text{Al}^{92}$  in the CdS layers effectively modifies the band alignment at CdS/ $\text{Sb}_2\text{Se}_3$  from a “cliff-like” to “spike-like” alignment, which is well suited for efficient electron extraction. Double buffer layers, *e.g.*,  $\text{TiO}_2/\text{CdS}$ ,<sup>119,120</sup>  $\text{SnO}_2/\text{CdS}$ ,<sup>120,121</sup>

$\text{Zn(O,S)}/\text{CdS}$ ,<sup>103,122,123</sup>  $\text{ZnSnO}_3/\text{CdS}$ ,<sup>5</sup>  $\text{Mg:SnO}_2/\text{CdS}$ ,<sup>94</sup> and  $\text{ZnO}/\text{TiO}_2$ ,<sup>124</sup> successfully suppresses the shunt paths and improve the cascaded band-alignment for efficient carrier transport. Han *et al.*<sup>125</sup> demonstrated the synergistic role of zinc halide ( $\text{ZnCl}_2$ ,  $\text{ZnBr}_2$ , and  $\text{ZnI}_2$ ) treatment of  $\text{TiO}_2$  in modifying the  $\text{TiO}_2$  layer and  $\text{TiO}_2/\text{Sb}_2\text{S}_3$  interface. In particular,  $\text{ZnCl}_2$  treatment curtailed the roughness and trap density in a  $\text{TiO}_2$  film, simultaneously increasing the electron concentration and mobility in the layer. The improvement in the  $\text{TiO}_2/\text{Sb}_2\text{S}_3$  interface quality and device performance on  $\text{ZnCl}_2$  treatment were mainly attributed to the passivation of undercoordinated  $\text{Ti}^{4+}$  and  $\text{O}^{2-}$  (on  $\text{TiO}_2$  surface) by coordinating with  $\text{Cl}^-$  and  $\text{Zn}^{2+}$ , respectively, and improved crystallinity of  $\text{Sb}_2\text{S}_3$  film (with enlarged grains and diminished trap states).  $\text{Zn(O,S)}$  has emerged as a promising ETL for  $\text{Sb}_2\text{X}_3$  solar cells, given that the S-atoms (at the surface of  $\text{Zn(O,S)}$ ) serve as chemical bridge to facilitate the quasi-epitaxial growth of  $\text{Sb}_2\text{X}_3$  films.<sup>50,126</sup> The quasi-vertical growth of [211]- and [121]- oriented  $\text{Sb}_2\text{S}_3$  films has been found to be conducive on CdS substrates, whereas the S-deficit on the CdS surface creates a beneficial bonding environment between the exposed undercoordinated  $\text{Cd}^{2+}$

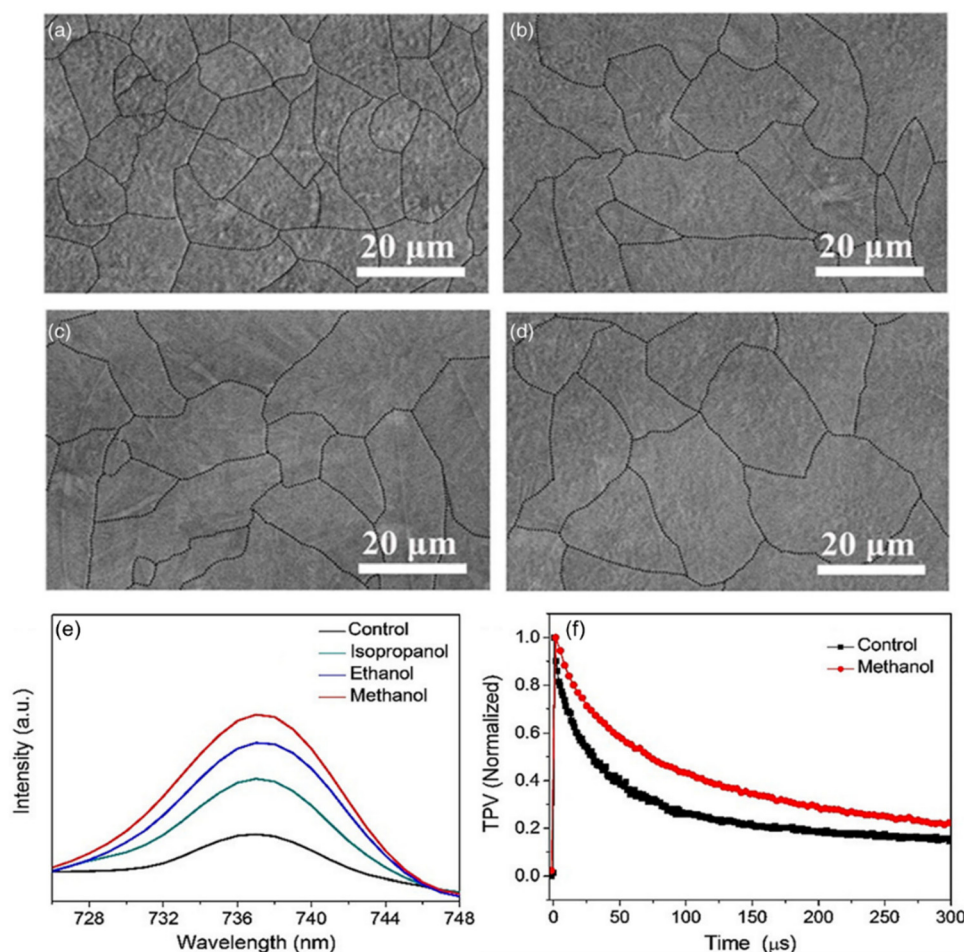


Fig. 17 SEM micrographs of (a) control  $\text{Sb}_2\text{S}_3$  film, (b) iso-annealed  $\text{Sb}_2\text{S}_3$  film, (c) eth-annealed  $\text{Sb}_2\text{S}_3$  film, and (d) meth-annealed  $\text{Sb}_2\text{S}_3$  film. (e) PL spectra of the control, iso-annealed, eth-annealed, and meth-annealed  $\text{Sb}_2\text{S}_3$  films and (f) TPV decay characteristics of the control and meth-annealed  $\text{Sb}_2\text{S}_3$  solar cells. Adapted with permission from ref. 106 Copyright 2019, Wiley-VCH.



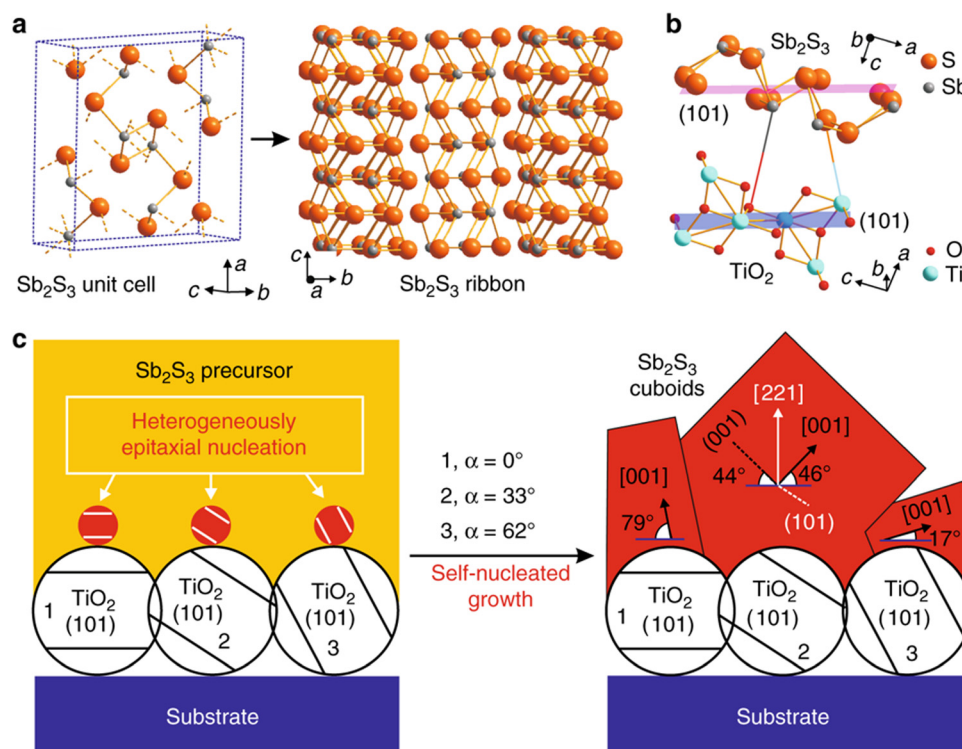


and S atoms in the  $\text{Sb}_2\text{S}_3$  molecules.<sup>127</sup> Also, the [101]-facet grains of  $\text{TiO}_2$  has been found to be effective in inducing the quasi-epitaxial growth of  $[hk1]$ -oriented  $\text{Sb}_2\text{S}_3$  films.<sup>128</sup> Solution processing has been found to be successful in the growth of  $\text{Sb}_2\text{S}_3$  monolayers (with preferentially oriented, single-crystalline cuboids) on  $\text{TiO}_2$  nanoparticle films (schematically presented in Fig. 18).<sup>129</sup>

Li *et al.*<sup>130</sup> skillfully utilized Cu-diffusion (from CuSCN) in the  $\text{Sb}_2\text{Se}_3$  absorber layer. This diffusion not only inverted the GBs (in  $\text{Sb}_2\text{Se}_3$ ) but also increased p-type doping in CuSCN (owing to the increment in  $V_{\text{Cu}}$ , which acts as shallow acceptors). CuSCN has been found to successfully improve the built-in potential and carrier collection efficiency and alleviate the back surface recombination losses. CuSCN and  $\text{CuSbSe}_2$  have emerged as the top choice as the HTL in  $\text{Sb}_2\text{X}_3$  solar cells, owing to the negligible lattice mismatch and lower trap density at the CuSCN/ $\text{Sb}_2\text{X}_3$  and  $\text{CuSbSe}_2$ / $\text{Sb}_2\text{X}_3$  interfaces. MnS has shown promising results as the HTL in  $\text{Sb}_2\text{X}_3$  solar cells, with a better performance compared to its spiro-OMeTAD-based counterparts in terms of PCE and stability. Furthermore, owing to the better band-alignment (with  $\text{Sb}_2\text{X}_3$ ), physicochemical robustness, and low price (4\$ per g) of MnS, it seems to be a better alternative to Spiro-OMeTAD powder (500\$ per g) in  $\text{Sb}_2\text{X}_3$  solar cells.<sup>131</sup> Liu *et al.*<sup>132</sup> employed an ultrathin poly(methyl methacrylate) (PMMA) film as a hole-selective

tunneling layer to modify the  $\text{Sb}_2\text{S}_3$ /spiro-OMeTAD interface in  $\text{Sb}_2\text{S}_3$ -based nano-array solar cells. As illustrated in Fig. 19, the PMMA layer was found to be effective in blocking the interfacial recombination channels and boosting the overall device performance.

$\text{SbCl}_3$  treatment has been used in the fabrication of  $\text{Sb}_2\text{S}_3$ <sup>93</sup> and  $\text{Sb}_2\text{Se}_3$ <sup>133</sup> solar cells.  $\text{SbCl}_3$  treatment of  $\text{Sb}_2\text{S}_3$  films was found to be highly effective in mitigating the trap states (at the interface and grain boundaries) and suppressing the non-radiative recombination. This strategy was successful in simultaneously improving the PCE and stability of  $\text{Sb}_2\text{S}_3$  solar cells. In the case of  $\text{Sb}_2\text{Se}_3$  solar cells,  $\text{SbCl}_3$  treatment (on the CdS layer) improved the morphology of the CdS layer, transformed the band-offset at CdS/ $\text{Sb}_2\text{Se}_3$  from cliff-like to spike-like, and induced  $[hk1]$ -oriented growth in  $\text{Sb}_2\text{Se}_3$  films. It was found that 1 nm thick (and sputtered)  $\text{Si}_3\text{N}_4$ <sup>134</sup> and 2 nm thick (and evaporated)  $\text{CeO}_2$ <sup>135</sup> ultrathin interface layers (sandwiched between CdS and  $\text{Sb}_2\text{Se}_3$ ) showed promising results in inducing columnar  $[hk1]$ -oriented growth, ameliorating the crystallinity and surface evenness of  $\text{Sb}_2\text{Se}_3$  thin films. Controlled oxygen doping in CdS (O:CdS) has been found to be successful in suppressing parasitic absorption (by CdS layer), improving the CdS/ $\text{Sb}_2\text{Se}_3$  band-alignment, mitigating the interface recombination and improving electron transport.<sup>136</sup> Potassium doping in CdS (K:CdS) has been reported to simultaneously improve



**Fig. 18** Schematic illustrations of (a) orthorhombic crystal structure of  $\text{Sb}_2\text{S}_3$ , consisting of  $(\text{Sb}_4\text{S}_6)_n$  ribbons stacked parallel to the [001] direction. (b) Atomic configuration at the  $\text{Sb}_2\text{S}_3/\text{TiO}_2$  interface. (101) planes of  $\text{Sb}_2\text{S}_3$  grow epitaxially over anatase  $\text{TiO}_2$  template. (c) Nucleation/growth model for the growth of  $\text{Sb}_2\text{S}_3$  single-crystalline cuboids on  $\text{TiO}_2$  nanoparticle film. The most competitive  $\text{TiO}_2$  nanoparticle orientation for the  $\text{Sb}_2\text{S}_3$  single-crystals to nucleate/grow features its (101) plane tilting at an angle ( $\alpha$ ) on the substrate plane close to  $\alpha = 33^\circ$  (particle 2), but the nucleation/growth at the other sites around (e.g., particle 1 of  $\alpha = 0^\circ$  and particle 3 of  $\alpha = 62^\circ$ ) is relatively much slower or suppressed due to their unfavorable  $\text{TiO}_2$  (101) plane orientations. Adapted under the guidelines of the Creative Commons CC BY license from ref. 129 Copyright 2019, Springer Nature.



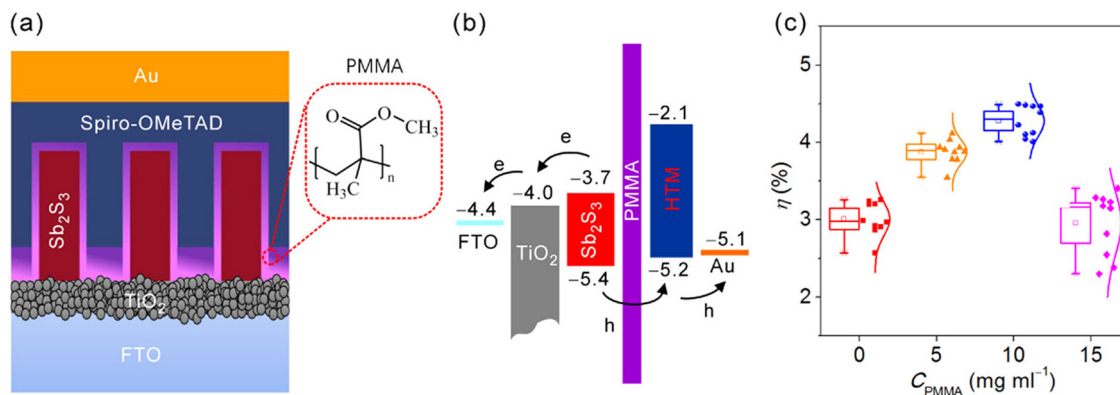


Fig. 19 Schematic of (a) configuration of the  $\text{Sb}_2\text{S}_3/\text{TiO}_2$  nano-array solar cell with PMMA interfacial layer and (b) its energy level diagram. (c) Statistical distribution of PCE ( $\eta$ ) without ( $C_{\text{PMMA}} = 0 \text{ mg mL}^{-1}$ ) and with ( $C_{\text{PMMA}} = 5, 10$  and  $15 \text{ mg mL}^{-1}$ ) PMMA interfacial layer, illustrating the dependence of the device performance on PMMA concentration. Adapted with permission from ref. 132 Copyright 2023, the American Chemical Society.

the crystallinity of  $\text{CdS}$  and  $\text{Sb}_2\text{S}_3$  films, decreasing the surface roughness of  $\text{Sb}_2\text{S}_3$  films and improving all the device parameters of  $\text{Sb}_2\text{S}_3$  solar cells. An ultrathin  $\text{PbS}$  interface layer successfully formed an  $\text{Sb-S-Pb}$  bonding bridge and cascaded band alignment, accelerating the hole transport and inhibiting carrier recombination at the  $\text{Sb}_2(\text{S,Se})_3/\text{carbon}$  interface.<sup>137</sup>

#### 4.5. Doping strategies

The low built-in potential caused by a low intrinsic doping density ( $\sim 10^{12} \text{ cm}^{-3}$ ) and poor carrier collection efficiency due to sluggish carrier mobility lead to a mediocre performance in  $\text{Sb}_2\text{X}_3$  solar cells. However, efficient doping in  $\text{Sb}_2\text{X}_3$  is very challenging. Owing to the distinctive Q-1D structure, the dopant atoms commonly settle in the inter-ribbon space or at the GBs and become electrically inactive.  $\text{Ti}$ <sup>138</sup> and  $\text{Zn}$ <sup>139</sup> doping has been shown to be successful in improving the crystallinity and absorbance of  $\text{Sb}_2\text{S}_3$  thin films, leading to a higher  $J_{\text{sc}}$  in  $\text{Sb}_2\text{S}_3$  solar cells. Particularly, Zn doping was found to successfully boost the carrier (electron) concentration in  $\text{Sb}_2\text{S}_3$  films and solar cell performance. Zn was found to coordinate with S at the grain boundaries. The low hole mobility in the pristine  $\text{Sb}_2\text{S}_3$  film was elevated from  $1.0 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  to  $205 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  on Sn-doping.<sup>140</sup> Ni-doped  $\text{Sb}_2\text{S}_3$  films demonstrated better photon harvesting than pristine  $\text{Sb}_2\text{S}_3$  films. Ag,<sup>141</sup> Co,<sup>142</sup> Mn,<sup>143</sup> Sm,<sup>144</sup> Fe,<sup>142</sup> Cu,<sup>145</sup> Sn,<sup>146</sup> Bi,<sup>147</sup> Se,<sup>148</sup> and alkali metal (Li, Na, K, Rb, and Cs)<sup>149</sup> doping has also been found to successfully enhance the carrier density and conductivity in  $\text{Sb}_2\text{S}_3$  films.<sup>150–152</sup> Ni settled at the grain boundaries and formed  $\text{Ni}_x\text{S}_y$  at lower temperatures, which provided nucleation centers for the  $[hk1]$ -oriented growth of  $\text{Sb}_2\text{S}_3$  ribbons.<sup>153</sup> An improvement of up to six orders in p-type conductivity (from  $10^{-8}$  to  $10^{-2} \text{ S cm}^{-1}$ ) was demonstrated in  $\text{Sb}_2\text{S}_3$  films on C doping.<sup>154</sup> The increased electrical conductivity was attributed to the formation of  $\text{C}_\text{S}$  substitutional defects in the  $\text{Sb}_2\text{S}_3$  crystal structure. C doping has also been used to implement  $\text{p}^+$  doping in  $\text{Sb}_2(\text{S,Se})_3$  films.<sup>155</sup>

Pb doping has been shown to successfully enhance the carrier (hole) concentration and conductivity of  $\text{Sb}_2\text{S}_3$  films by four and three orders, respectively.<sup>156</sup> The  $\text{Sb}^{3+}$  in the  $\text{Sb}_2\text{S}_3$  crystal was partially replaced by  $\text{Pb}^{2+}$ , forming acceptor energy

levels for effective hole-doping. I doping induced n-doping in  $\text{Sb}_2\text{S}_3$  via substitution of Se by the I atoms, leading to a one order increase in conductivity.<sup>157</sup> An improvement of three orders in p-type conductivity and four orders in hole density in  $\text{Sb}_2\text{S}_3$  films was obtained on Sn doping.<sup>157</sup> The conductivity of  $\text{Sb}_2\text{S}_3$  films slightly improved upon doping with transition metals and lanthanides ( $\text{Lu}^{3+}$ ,  $\text{Ho}^{3+}$ ,  $\text{Nd}^{3+}$ ,  $\text{Sm}^{3+}$ , and  $\text{Gd}^{3+}$ ) and co-doping with  $\text{Lu}^{3+}/\text{Yb}^{3+}$  and  $\text{Lu}^{3+}/\text{Er}^{3+}$ .<sup>158,159</sup> Fe doping in  $\text{Sb}_2\text{S}_3$  was found to improve the carrier concentrations and mobility of the majority carriers. Costa *et al.*<sup>160</sup> reported that  $\text{Fe:Sb}_2\text{S}_3$  films exhibit p-type conductivity, while Li *et al.*<sup>161</sup> reported that these films have n-type conductivity. Te doping has been implemented to introduce n-type doping in  $\text{Sb}_2\text{S}_3$  films.  $\text{Te:Sb}_2\text{S}_3$  films were deposited via thermal evaporation,<sup>162</sup> sputtering,<sup>163</sup> and spin coating.<sup>92</sup> However, three independent studies suggested the Te-dopant settled at the lattice in the inter ribbon spacing and in the grain boundaries, respectively. S,<sup>92</sup> Na<sup>164</sup> and Mg<sup>161</sup> doping has also been implemented in  $\text{Sb}_2\text{S}_3$  thin films. S substitutes Se in the lattice, while Na and Mg prefer to settle in the grain boundaries or inter-ribbon space (where they are electrically inert).

## 5. Concluding remarks

Considering the overall assessment, the expanding toolbox of defect-engineering strategies is anticipated to play a pivotal role in the foreseeable development of  $\text{Sb}_2\text{X}_3$  PV. The dominant deep-level defects in  $\text{Sb}_2\text{X}_3$  solar cells are anion-vacancies ( $\text{V}_\text{S}$  and  $\text{V}_{\text{Se}}$ ) and cation anti-sites ( $\text{Sb}_\text{S}$  and  $\text{Sb}_{\text{Se}}$ ). It seems conceivable that the regulation of S/Se volatilization (especially during the annealing step) is the key to controlling the deep defects. The high formation energy of  $\text{V}_{\text{S/Se}}$  induces (thermally driven) the migration of neighboring Sb atoms to form  $\text{Sb}_{\text{S/Se}}$  defects. Therefore, controlled annealing of as-deposited (amorphous) films under an S/Se-rich environment or in an  $\text{H}_2\text{S}/\text{CS}_2$  environment can be a straightforward strategy to passivate  $\text{V}_{\text{S/Se}}$  before they transform into detrimental  $\text{Sb}_{\text{S/Se}}$  defects. Exploring novel low-temperature synthesis protocols for the synthesis of  $\text{Sb}_2\text{X}_3$  films can also be instrumental in curtaining the anti-site



defects. Nanocrystal seed- and seed-layer assisted growth has also been found to successfully deposit  $\text{Sb}_2\text{X}_3$  films with desirable characteristics. Vapor transport deposition techniques have been found to be more effective in depositing  $\text{Sb}_2\text{X}_3$  films with diminished anti-site defects compared to the conventional thermal evaporation and solution processing methods. A slight S/Se-rich stoichiometry and  $[hk1]$ -oriented growth are crucial for obtaining  $\text{Sb}_2\text{X}_3$  films with a low-defect density, better photoconductivity and adhesion with the substrates than S/Se-deficient and  $[hk0]$ -oriented films. In addition, doping strategies, and halide salt treatment of ETL/ $\text{Sb}_2\text{X}_3$  and  $\text{Sb}_2\text{X}_3$ /HTL interfaces have also been found to be instrumental in regulating the defect dynamics and PV performance. This work is anticipated to spur greater control of the engineering of deep defects in  $\text{Sb}_2\text{X}_3$  thin films, eliminating their adverse effects on the device performance and fabrication of efficient  $\text{Sb}_2\text{X}_3$  solar cells.

## Author contributions

S. B. – conceptualization, drafting original manuscript. S. S. – visualization, methodology. A. K. C. – writing, editing. N. S. P. – reviewing, editing, formatting. A. G. P. – methodology, editing. K. R. – funding acquisition, supervision.

## Conflicts of interest

There are no conflicts to declare.

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