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Aromaticity and photophysics of tetrasila- and tetragerma-annelated tetrathienylenes as new representatives of the hetero[8]circulene family†

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The electronic structure, absorption and emission spectra, aromaticity and photophysical behavior of the recently synthesized tetrasilatetrathia[8]circulene and tetragermatetrathia[8]circulene compounds have been studied computationally. Both compounds demonstrate a specific bifacial aromaticity, which is unusual for hetero[8]circulenes; the inner eight-membered core sustains an expected strong paratropic magnetically-induced ring current, while the outer perimeter contains saturated Si(Et)₂ and Ge(Et)₂ moieties which break the conjugation between the thiophene rings. The overall magnetically-induced ring current for both studied circulenes is close to zero because of the strong local diatropic currents in each thiophene ring that compensate the paratropic counterpart. The electronic absorption and emission spectra of tetrasilatetrathia[8]circulene and tetragermatetrathia[8]circulene demonstrate a clear visible vibronic progression. The 0–0 band is the most active one in the absorption spectra, while in the fluorescence spectra the 0–1 band composed of several normal vibrations is more intense compared with the 0–0 band in excellent agreement with experiment. Accounting for spin–orbit coupling effects, an analysis of the photophysical constants for the two compounds demonstrates: (1) a clear manifestation of the internal heavy atom effect on the inter-system crossing efficiency; (2) one to two order domination of non-radiative rates over the fluorescence rate; and (3) that the S₁–S₀ internal conversion is extremely slow and can not compete with the fluorescence, while the S₁–T_n inter-system crossing is a main deactivation channel of the S₁ excited state. These results provide new insight into the electronic structure and photophysics of tetrasilatetrathia[8]circulene and tetragermatetrathia[8]circulene as novel standalone representatives of hetero[8]circulenes – tetraannelated derivatives of tetrathienylene.

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1. Introduction

Among the numerous representatives of hetero[8]circulenes the heteroannelated derivatives of tetraphenylene constitute the main and most studied ones.^{1–3} Some of them, like azaoxa[8]circulenes and tetraoxa[8]circulenes, demonstrate an intense blue fluorescence^{4,5} that is useful for fabrication of blue fluorescent organic light-emitting diodes (OLEDs) and extremely stable white

exciplex-based OLEDs.⁶ At the same time, tetraoxa[8]circulenes and other heteroannelated tetraphenylenes (like tetraaza[8]circulenes and tetrathia[8]circulenes) demonstrate an ability to form continuous infinite one- and two-dimensional ribbons and sheets.^{7–11} These graphene-like materials sustain semi-conductive properties and high dynamic stability, which makes them promising organic materials for use as semiconductors for modern organic electronics applications.^{7,11} The electronic structure, spectra and aromaticity of heteroannelated derivatives of tetraphenylene have been well studied both experimentally and theoretically^{1–3,12} and even some novel hypothetical circulenes have been predicted as upcoming targets of synthesis surveys.⁵

Heteroannelated derivatives of tetrathienylene (cycloocta[1,2-c:3,4-c':5,6-c'':7,8-c''']tetrathiophene) constitute another much smaller class of hetero[8]circulenes. Among them only octathia[8]circulene (8S)^{13,14} and tetrathiatetraselena[8]circulene (4S4Se)¹⁵ have been prepared and characterized. Because of the high molecular symmetry and poor solubility of the 8S and 4S4Se circulenes their electronic absorption and emission spectra

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have only been sparsely studied (mainly by quantum-chemical methods),^{16,17} while the vibrational IR and Raman spectra have been well characterized and investigated in great detail, accounting for crystal packing effects and anharmonicity corrections.^{16,18,19} Both 8S and 4S4Se species have been used as promising organic semiconductors for fabrication of organic field-effect transistors with a maximum hole mobility of $9 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$.¹⁵ Due to the planar structure and very high molecular symmetry, the 8S and 4S4Se species also demonstrate supramolecular assembly into 2D and 3D arrays, which is extremely important for modulation of the charge-transport properties in such systems.^{20,21}

Very recently, Miyake *et al.* synthesized novel representatives of heteroannulated tetrathienylenes called tetrasilatetraphia[8]circulene²² (4S4SiEt₂) and tetragermatetraphia[8]circulene²³ (4S4GeEt₂) by a palladium-catalyzed silylation/germylation process with subsequent rhodium-catalyzed intramolecular dehydrogenative cyclization. A similar technique has been also applied for the synthesis of tetrabenzotetraphia[8]circulenes²⁴ (4S), which at the same time can be considered as tetrathiaannulated derivatives of tetraphenylene and tetrabenzoannulated derivatives of tetrathienylene (Fig. 1). Important to note is that the 4S compound was prepared for the first time by Wong *et al.* from octabromotetraphenylene by S-bridging,²⁵ and one year later by Miyake *et al.* who independently prepared the 4S compound by a reverse strategy through (–CH=CH–) bridging of tetrathienylene.²⁴

Circulene 4S has been comprehensively studied both theoretically and experimentally during the last ten years including the electronic and vibrational spectra, crystal structure, aromaticity *etc.*,^{26–29} while the novel 4S4SiEt₂ and 4S4GeEt₂ compounds^{22,23} still represent a challenge for computational chemistry. Especially important to note is that the photophysics of 4S4SiEt₂ and 4S4GeEt₂ has been studied experimentally by Miyake *et al.* in great detail^{22,23} including estimation of the rate constants for the internal conversion (IC), intersystem crossing (ISC), fluorescence and phosphorescence processes. That has motivated us to focus the present paper on high-level computations of the photophysical properties for the 4S4SiEt₂ and 4S4GeEt₂ compounds in order to explain and support the available experimental data.^{22,23}

Moreover, the aromaticity of heteroannulated derivatives of tetrathienylene is still insufficiently studied and thus we have applied computations, in particular the gauge-including magnetically induced currents (GIMIC) methodology,^{30,31} to explain the distribution of diatropic and paratropic ring currents in the 4S4SiEt₂ and 4S4GeEt₂ molecules as novel representatives of fully annulated tetrathienylenes.

2. Computational details

The structure of the 4S4SiEt₂ and 4S4GeEt₂ molecules in the ground singlet electronic state (*S*₀) and first excited state (*S*₁) has been optimized at the DFT³² and TD DFT³³ level of theory, respectively, in the gas phase using the B3LYP^{34,35} hybrid functional and the 6-31G(d,p)³⁶ basis set within the Gaussian 16 program package.³⁷ The vibronic absorption and emission spectra of the 4S4SiEt₂ and 4S4GeEt₂ molecules have been calculated within the Franck–Condon approximation through the time-domain formalism of Lax and Heller within the double harmonic parallel-mode approximation.^{38,39}

Magnetically induced current densities (current strengths) have been calculated using the GIMIC method.^{30,31} The NMR shielding calculations required for GIMIC computations have been performed at the B3LYP/def2-TZVP^{34,35,40} level of theory using the Turbomole program package.⁴¹ Ring-current strengths (*I*, nA T^{–1}) and current strength pathways for the selected chemical bonds have been obtained by integrating the current density that flows through planes intersecting the chemical bonds. Additionally, the nucleus-independent chemical shift (NICS)⁴² descriptors have been calculated for the 4S4SiEt₂ and 4S4GeEt₂ molecules at the center (NICS(0)) and 1 Å above the center (NICS(1)) of each ring using the gauge-independent atomic orbital (GIAO)⁴³ method at the B3LYP/6-311++G(d,p)^{34,35,44,45} level of theory. The Gaussian 16 program package³⁷ was used for the NICS calculations.

The photophysical rate constants have been computed for the *S*₁ state geometries based on extended multi-configuration quasidegenerate second order perturbation theory (XMC-QDPT2)⁴⁶

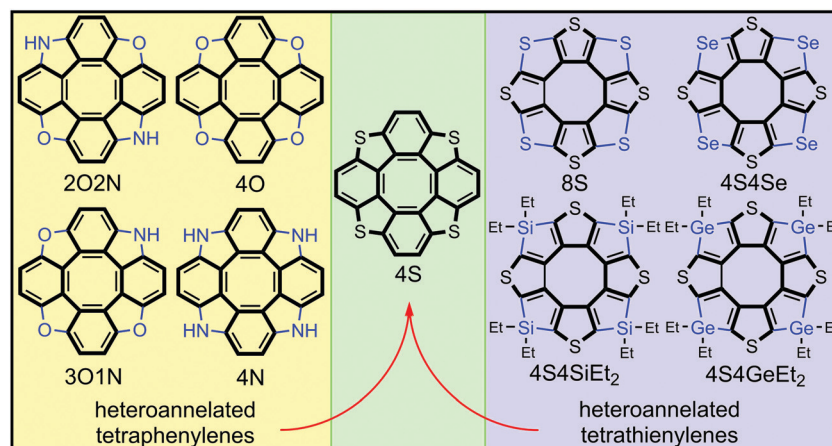


Fig. 1 The structure of some representatives of heteroannulated derivatives of tetraphenylenes and tetrathienylenes.



calculations of the singlet and triplet excitation energies. The number of active electrons (e), the number of active orbitals (o), and the number of states (s) of the state-average (SA) complete active space self-consistent field (CASSCF) calculations have been assumed as 10e, 10o, and 5s. The XMC-QDPT2 calculations have been carried out using the Firefly program package.⁴⁷

The spin-orbit coupling matrix elements $\langle S_1 | H_{SO} | T_i \rangle$ between the S_1 state and the energetically low-lying i -th triplet states T_i have been calculated at the *ab initio* CASSCF level of theory using the Gamess-US software.⁴⁸ The IC rate constants have been computed by the original algorithm published in ref. 49 using the nonadiabatic coupling matrix elements calculated within the Turbomole package⁴¹ at the TD-DFT/B3LYP/def2-TZVP^{34,35,40} level of theory in the framework of perturbation theory. The spin-orbit coupling matrix elements required for the ISC rate calculations were taken from the previous CASSCF calculations. We assume that the vibrational frequencies and the Huang–Rhys factors are the same for the IC and ISC channels, something that was tested in our previous work.⁴⁹ More details on the photo-physical calculations can be found in the ESI.†

3. Results and discussion

3.1. Aromaticity (NICS and GIMIC criteria)

The aromaticity of hetero[8]circulenes is still a challenge for computational chemistry. The most acceptable concept for today is that hetero[8]circulenes are nonaromatic species in general because of the competitive diatropic (I_{dia}) and paratropic (I_{para}) currents in the outer (so-called rim) and inner (so-called hub) subsystems.^{28,30,50} Thus, the resulting magnetically-induced current (MIC) is close to zero ($|I_{\text{dia}}| \approx |I_{\text{para}}|$) implying a non-aromatic character of most hetero[8]circulenes – derivatives of tetraphenylene and tetrathienylene. However, the studied 4S4SiEt₂ and 4S4GeEt₂ molecules contain sp³-hybrid Si and Ge atoms in the rim-subsystem that break conjugation in the outer perimeter. That is why the common double aromaticity concept can not be applied to the 4S4SiEt₂ and 4S4GeEt₂ molecules. Indeed, the current estimated along the Si(Ge)–C bonds is strictly equal to zero (Fig. 2).

At the same time each thiophene ring possesses its own strong local diatropic current ($I_{\text{dia}} = 8.6$ and 8.3 nA T^{-1} for the 4S4SiEt₂ and 4S4GeEt₂ molecules, respectively) in contrast to the totally-conjugated circulenes like tetraoxa[8]circulene, octathia[8]circulene, *etc.*⁵⁰ It is here rather surprising that the local diatropic MIC strength (I_{dia}) is almost the same by absolute value as the paratropic MIC inside the planar eight-membered core ($I_{\text{para}} = -9.2$ and -8.7 nA T^{-1} for the 4S4SiEt₂ and 4S4GeEt₂ molecules, respectively). This fact supports the same overall non-aromatic nature of the 4S4SiEt₂ and 4S4GeEt₂ molecules ($I_{\text{tot}} = -0.6$ and -0.4 nA T^{-1} , respectively) as for most of the hetero[8]circulenes.⁵⁰ However, the topology of the MIC is principally different. Generally, the totally-conjugated hetero[8]circulenes do not sustain significant MIC along the radial C–C bonds, *i.e.* the diatropic and paratropic currents circulate in the rim and hub perimeters in opposite directions almost independently of each other (the rim and hub subsystems are

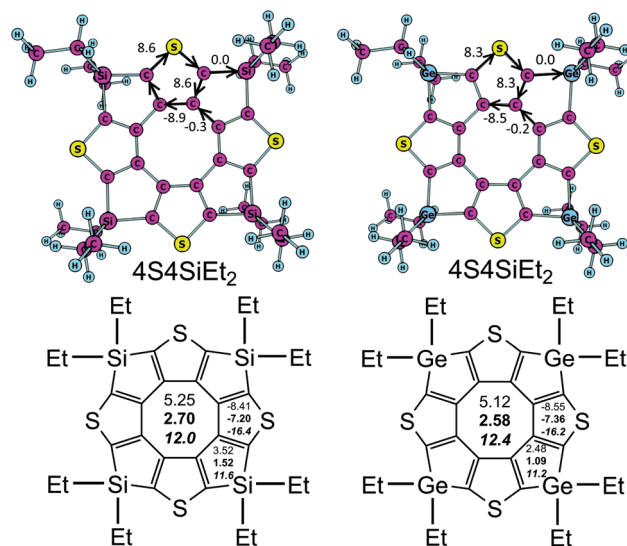


Fig. 2 MIC strengths (top, numbers in nA T^{-1}) and NICS(0), NICS(1) (in bold), and NICS(1)_{zz} indices (bottom part, numbers in ppm) for the 4S4SiEt₂ and 4S4GeEt₂ molecules.

separated by the C–C radial bonds). In the case of 4S4SiEt₂ and 4S4GeEt₂ circulenes the rim-system is broken by the sp³-Si(Ge) atoms, which induces MIC circulation along the radial C–C bonds (Fig. 3).

The GIMIC results for the 4S4SiEt₂ and 4S4GeEt₂ molecules are in good qualitative agreement with the NICS calculations (Fig. 2). The inner cyclooctatetraene core sustains significantly positive NICS(0), NICS(1) and NICS(1)_{zz} values for both studied molecules meaning an antiaromatic nature of this ring. At the same time, the NICS(0), NICS(1) and NICS(1)_{zz} values for the thiophene rings are significantly negative (*i.e.* these rings are locally aromatic), but for the Si/Ge-containing rings the NICS(0) and NICS(1) descriptors are weakly positive (*i.e.* these rings are just weakly antiaromatic). Despite the fact that NICS(1)_{zz} is significantly positive for the Si/Ge-containing rings, the two other xx and yy components of the magnetic shielding tensor are significantly negative (about -4 ppm) and compensate the zz one. Actually, Si/Ge-containing rings do not sustain

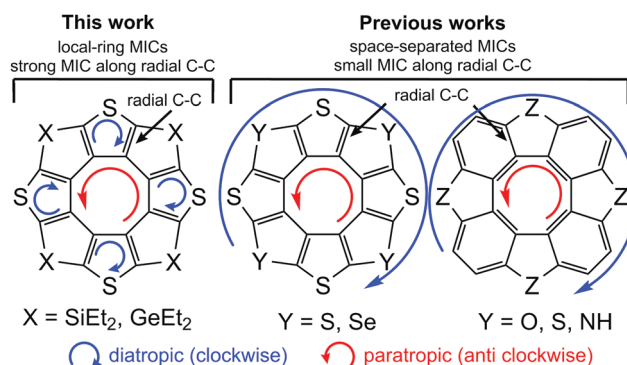


Fig. 3 Principal difference between the ring-current topology in the totally conjugated heterocirculenes and the studied 4S4SiEt₂ and 4S4GeEt₂ compounds.



significant π -electronic MICs and that is why NICS(1)_{zz} indices are not really a good local aromaticity criterion for these rings. The overall balance between the NICS indices looks the same as obtained by the GIMIC method: local paratropic currents inside the cyclooctatetraene core are compensated by the local diatropic currents inside the thiophene rings, while Si/Ge-containing rings just weakly contribute with an antiaromatic component. Actually, the NICS data emphasize the local-ring topology of the MICs in the 4S4SiEt₂ and 4S4GeEt₂ molecules obtained by GIMIC and this model is principally different from the (C–C-radial)-separated topology of concentric MICs in most hetero[8]circulenes.^{28,30,50} An additional observation is that the outer ethyl substituents are not important in the GIMIC and NICS calculations; the deviations in I_{tot} for model 4S4SiH₂ and 4S4GeH₂ are less than 0.1 nA T^{−1} and the differences in the NICS values are of the order of 10^{−2} ppm. This fact is expected because of the negligible effect of σ -type ethyl groups on the π -system of tetrathienylene.

3.2. Vibronic absorption and fluorescence spectra

4S4SiEt₂ and 4S4GeEt₂ circulenes are characterized by similar absorption and fluorescence spectra with clearly resolved vibronic structure.^{22,23} In order to explain these spectra and assign the vibronic bands we have initially optimized both the 4S4SiEt₂ and 4S4GeEt₂ circulenes using the TDDFT/B3LYP/6-31G(d,p) method in the first excited singlet state (S_1), which is responsible for the fluorescence emission. The results are summarized in Table 1. By orbital symmetry type, the S_1 state of 4S4SiEt₂ and 4S4GeEt₂ can be attributed to the $\pi\pi^*$ nature corresponding to the HOMO–LUMO configuration (Fig. S1, ESI†).

The S_1 state is structurally closely similar to the ground state – the main differences relate to an elongation of the radial C–C and C–S bonds, while the remaining skeletal bonds (C–C bonds in the inner COT core and Si/Ge–C bonds) become slightly shorter upon excitation into the S_1 state (Table S1, ESI†). Such structural behavior is principally different from that for the totally-conjugated tetraoxa[8]circulene and mixed azaoxa[8]circulenes, which change bond alternation in the inner COT core upon excitation into the S_1 state.⁵¹

Based on the calculated Hessians for the ground (S_0) state and gradients of the excited (S_1) states of real Et-substituted and model H-substituted compounds the absorption ($S_0 \rightarrow S_1$) and fluorescence ($S_1 \rightarrow S_0$) vibronic spectra were calculated within the Franck–Condon approximation using the Lax–Heller model.^{38,39} This model shows a good result for the cases of

slightly-shifted potential energy surfaces (PES) of the S_0 and S_1 states. Indeed, the energy difference between the vertical and 0–0 electronic transitions is only 0.04 eV and 0.05 eV for the 4S4SiEt₂ and 4S4GeEt₂ molecules, respectively, something that allows us to use the simple Lax–Heller model.^{38,39}

At the first stage, we have defined the vibronic activity of the outer ethyl groups. We have found that these groups do not significantly contribute to the main promotive skeletal vibrational modes both in absorption and fluorescence vibronic spectra and only some new low-frequency promotive modes appear in the spectra of the model 4S4SiH₂ and 4S4GeH₂ molecules. All of these vibrational modes are characterized by very small Huang–Rhys factors and actually are overlapped by the 0–0 transition. The calculated vs. experimental vibronic absorption and fluorescence spectra are presented in Fig. 4 and 5, respectively, while the shapes of the promotive modes are shown in the ESI† (Fig. S2 and S3 for the real 4S4SiEt₂ and 4S4GeEt₂ molecules and for the model H-substituted 4S4SiH₂ and 4S4GeH₂ species). As can be seen from Fig. 4, the 0–0 electronic transition takes place at 380 and 370 nm for the 4S4SiEt₂ and 4S4GeEt₂ molecules and it produces the first most intense band in the absorption spectra of both compounds (exp.: 384 and 365 nm, respectively). Just three vibrational modes promote the vibronic progression in the absorption spectra of the 4S4SiEt₂ (4S4GeEt₂) molecule in the region 320–360 nm (Table 2). The modes at 794 (763) and 1329 (1320) cm^{−1} correspond to the “breathing” symmetrical stretching vibrations of the inner COT core, while the mode 1329 (1320) cm^{−1} is additionally accompanied by an in-phase stretching of C–C radial bonds. These specific “breathing” vibrational modes are also active in the vibronic absorption spectra of other hetero[8]circulenes⁵¹ indicating that the rigid planar COT skeleton is mainly responsible for the experimentally observed vibronic structure in electronic spectra of hetero[8]circulenes as a unique class of heteroannulated cyclooctatetraenes. The third active mode at 1139 (1117) cm^{−1} corresponds to the in-plane out-of-phase bending vibrations of the radial C–C bonds but it is much less intense in the vibronic absorption spectrum compared with the COT “breathing” vibrations (Table 2).

In the vibronic fluorescence spectra the 0–0 transition is less intense than the 0–1 band promoted by the low-frequency mode at 666 (626) cm^{−1} for the 4S4SiEt₂ (4S4GeEt₂) molecule. Actually, both these 0–0 and 0–1 transitions form the first experimentally-observed band at 397 and 377 nm in the fluorescence spectra of the 4S4SiEt₂ and 4S4GeEt₂ compounds, respectively. The next most intense band in the experimental spectrum (412 and 392 nm, respectively) corresponds to the manifold of active modes in the region of 1100–1300 cm^{−1} combined with a double excitation of the low-frequency mode 666 (628) cm^{−1} (Table 2). All the low-frequency vibrational modes involve the bending vibrations of the CH bonds of the ethyl substituents (Fig. S2 and S4), but for the model 4S4SiH₂ and 4S4GeH₂ compounds these modes transform into the Si(Ge)–H bending vibrations of different symmetry. The active modes in the region 1100–1300 cm^{−1} mainly correspond to the asymmetrical stretching and bending vibrations of the radial C–C bonds (Fig. S3, ESI†).

Table 1 Energies (eV) and oscillator strengths (in parenthesis) for the vertical electronic transitions of the 4S4SiEt₂ and 4S4GeEt₂ circulenes calculated using the TDDFT/B3LYP/6-31G(d,p) method at the S_0 - and S_1 -state starting geometries

State	4S4SiEt ₂	4S4GeEt ₂
S_0 geometry		
S_1	3.302 (0.003); exp.: 3.23	3.485 (0.003); exp.: 3.40
S_2	3.34 (0.0001)	3.52 (0.0001)
S_1 geometry		
S_1	2.89 (0.001); exp.: 3.23	3.22; exp.: 3.40



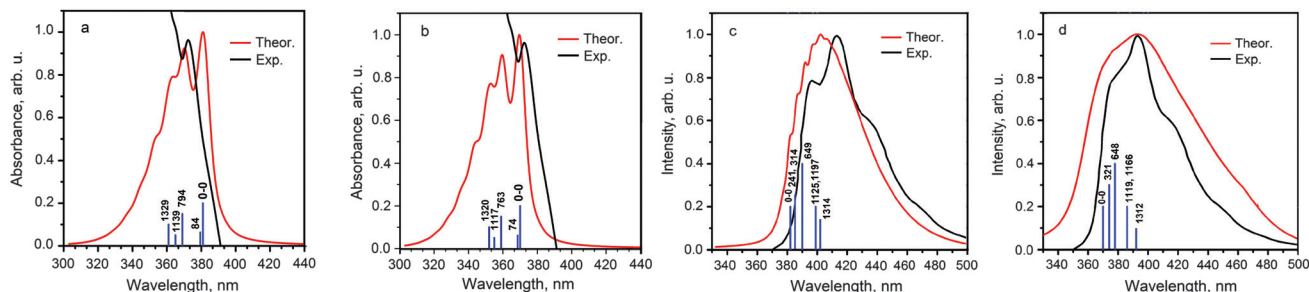


Fig. 4 Theoretically calculated vs. experimental vibronic absorption (a and b) and fluorescence (c and d) spectra of the 4S4SiEt₂ (a and c) and 4S4GeEt₂ (b and d) molecules. Vertical sticks correspond to the energy of promotive vibrational modes (cm⁻¹) relative to the 0–0 transition. Experimental curves are taken from ref. 22 and 23.

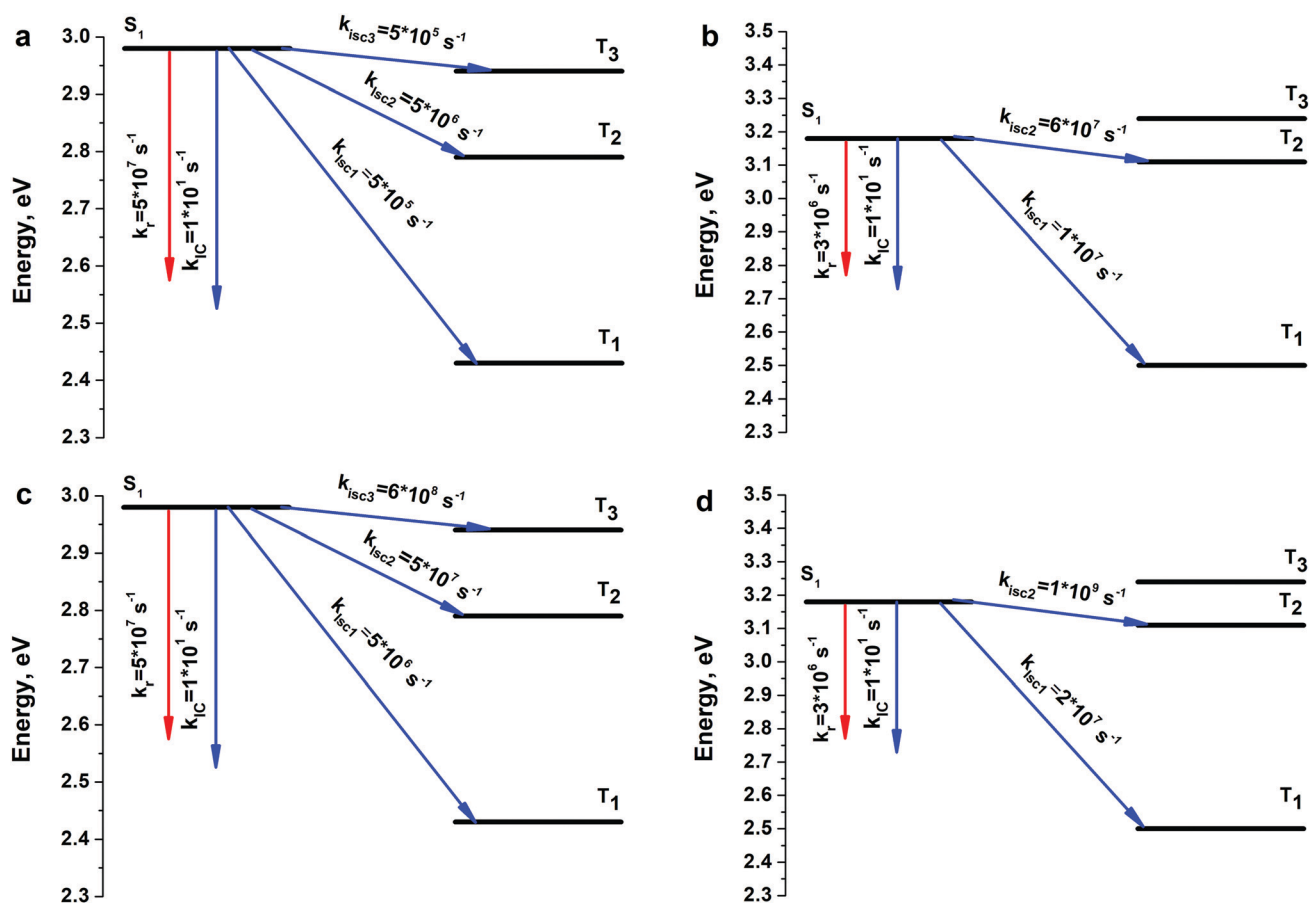


Fig. 5 Calculated rate constants for the main photophysical processes of S₁ state deactivation for the 4S4SiH₂ (a and c) and 4S4GeH₂ (b and d) molecules without (a and b) and with (c and d) accounting for spin–vibronic interactions.

Table 2 Promotive modes active in the absorption and fluorescence spectra of the 4S4SiEt₂ and 4S4GeEt₂ molecules and their H-substituted models 4S4SiH₂ and 4S4GeH₂ (the Huang–Rhys factors are presented in parentheses)

	Absorption (S ₀ –S ₁)	Fluorescence (S ₁ –S ₀)
4S4SiEt ₂	84(0.1); 794(0.5); 1139(0.06); 1329(0.31)	241(0.3); 314(0.4); 369(0.4); 397(0.2); 649(0.6); 1125(0.1); 1197(0.20); 1314(0.2)
4S4GeEt ₂	74(0.1); 763(0.65); 1117(0.08); 1320(0.31)	223(0.36); 321(0.49); 343(0.43); 362(0.2); 648(0.75); 1119(0.14); 1166(0.30); 1312(0.2)
4S4SiH ₂	784(0.75); 1152(0.07); 1326(0.29)	250(0.36); 323(0.49); 373(0.43); 397(0.2); 666(0.75); 1137(0.14); 1193(0.30); 1312(0.2)
4S4GeH ₂	754(0.65); 1138(0.1); 1317(0.29)	180(0.43); 297(0.46); 321(0.25); 341(0.2); 628(0.33); 1127(0.18); 1185(0.30); 1331(0.2)



One can stress that the ethyl groups play a more important role in vibronic fluorescence spectra compared with the absorption spectra because of the high vibronic activity of low-frequency bending vibrations of the CH groups in the ethyl moieties. But even the H-substituted models 4S4SiH₂ and 4S4GeH₂ adequately describe the experimentally observed fluorescence spectra because the Ge(Si)–H bending vibrations demonstrate a similar vibronic activity (Table 2).

3.3. Photophysical properties

In ref. 22 and 23 the main photophysical constants like the internal-conversion rate (k_{IC}), intersystem crossing rate (k_{ISC}), fluorescence rate (k_f) and fluorescence quantum yield (ϕ_f) were experimentally estimated for both 4S4SiEt₂ and 4S4GeEt₂ circulenes. Actually, it was the first time a detailed experimental estimation was given for the photophysical properties of the hetero[8]circulenes. In previous sections we have stated that the 4S4SiH₂ and 4S4GeH₂ models adequately describe the electronic and spectroscopic parameters of the real 4S4SiEt₂ and 4S4GeEt₂ systems. Thus, in this section we are focusing on the photophysical properties of the model 4S4SiH₂ and 4S4GeH₂ compounds in comparison with experimentally measured constants.^{22,23} One can see from Fig. 5 that three (T_1 , T_2 and T_3) and two (T_1 and T_2) triplet excited states lie below the S_1 level for the 4S4SiH₂ and 4S4GeH₂ molecules, respectively (the same pattern was also computed for the Et-substituted systems). All these triplet states participate in non-radiative quenching of the S_1 state through multichannel ISC (k_{ISC1} , k_{ISC2} and k_{ISC3}). Calculations of the corresponding spin–orbit coupling (SOC) matrix elements show that the S_1 – T_1 radiationless transition is the main deactivation channel of the S_1 state energy without accounting for spin–vibronic interactions (Table 3). One can also observe a clear internal heavy atom effect⁵² moving from the Si- to the Ge-containing circulene ($\langle\phi(S_1)|H_{SO}|\phi(T_n)\rangle$ (4S4GeH₂) > $\langle\phi(S_1)|H_{SO}|\phi(T_n)\rangle$ (4S4SiH₂)). However, the total ISC rate ($\sum k_{ISCi}$) estimated as a sum of partial $k_{ISC1,2,3}$ rates was found to be considerably underestimated compared with the experimental data (6.0×10^6 s^{−1} vs. exp.: 4.6×10^8 for 4S4SiEt₂ and 7.0×10^7 s^{−1} vs. exp.: 1.3×10^8 s^{−1} for 4S4GeEt₂, Table 3). In order to clarify the $\sum k_{ISCi}$ rate, spin–vibronic interactions were taken into account through the distortion of geometries along the out-of-plane E_g normal modes (Fig. 6). As follows from the symmetry selection rules for the SOC operator,^{52–54} only the out-of-plane E_g vibrations affect the corresponding SOC matrix elements. As a result, all the $\langle\phi(S_1)|H_{SO}|\phi(T_n)\rangle$ matrix elements have been increased (especially $\langle\phi(S_1)|H_{SO}|\phi(T_3)\rangle$), which produces an enhancement of $\sum k_{ISCi}$ by one–two orders of magnitude, providing more reliable agreement with experiment (Table 3). Finally, taking account of spin–vibronic interactions provides excellent agreement with respect to the fluorescence quantum yield (0.07 vs. exp.: 0.012 for 4S4SiEt₂, 0.003 vs. exp.: 0.0076 for 4S4GeEt₂), which means a prevailing role of ISC quenching over the IC process.

At the same time our calculations are considerably different (by seven orders of magnitude) from experimental data with

Table 3 Energies of the S_1 and T_n excited electronic states, SOC matrix elements between them and the corresponding rate constants for the IC (k_{IC}), ISC (k_{ISC}) and fluorescence (k_f) processes, including the resulting fluorescence quantum yields (ϕ_f)

Parameter	Without spin–vibronic interactions		With spin–vibronic interactions	
	4S4SiH ₂	4S4GeH ₂	4S4SiH ₂	4S4GeH ₂
S_1 , eV (f)	3.04(0.14)	3.18(0.007)	—	—
S_1 (exp.), eV	3.23	3.40	—	—
T_1 , eV	2.43	2.5	—	—
T_2 , eV	2.79	3.11	—	—
T_3 , eV	2.94	3.24	—	—
$\langle\phi(S_1) H_{SO} \phi(T_1)\rangle$, cm ^{−1}	0.30	2.27	0.89	3.35
$\langle\phi(S_1) H_{SO} \phi(T_2)\rangle$, cm ^{−1}	0.08	0.58	0.21	0.13
$\langle\phi(S_1) H_{SO} \phi(T_3)\rangle$, cm ^{−1}	0.02	0.12	0.43	0.8
k_{ISC1} , s ^{−1}	5×10^5	1×10^7 s ^{−1}	5×10^6	2×10^7
k_{ISC2} , s ^{−1}	5×10^6	6×10^7 s ^{−1}	5×10^7	1×10^9
k_{ISC3} , s ^{−1}	5×10^5	—	6×10^8	—
$\sum k_{ISCi}$, s ^{−1}	6×10^6	7×10^7	7×10^8	1×10^9
$\sum k_{ISCi}$ (exp.), s ^{−1}	4.6×10^8	1.3×10^8	—	—
k_{IC} , s ^{−1}	1×10^1	1×10^1	1×10^1	1×10^1
k_{IC} (exp.), s ^{−1}	6.3×10^8	2.7×10^8	—	—
k_f , s ^{−1}	5×10^7	3×10^6	5×10^7	3×10^6
k_f (exp.), s ^{−1}	1.3×10^7	3×10^6	—	—
ϕ_f	0.90	0.02	0.07	0.003
ϕ_f (exp.)	0.012	0.0076	—	—

respect to the internal conversion rates ($k_{IC} \sim 10^1$ s^{−1}).^{22,23} Experimental estimations provide a 10^8 s^{−1} order of magnitude for the internal conversion process, which means that this channel is quite active in the S_1 state deactivation. In our opinion, ref. 22 and 23 give an incorrect estimation of the k_{ISC} rate by measuring the decay of the T–T absorption during energy transfer from the 4S4SiH₂ donor to the β -carotene acceptor.^{55,56} This method only means that the T_1 state lies lower than the S_1 state for the studied molecules and that only the T_1 state of circulene is able to transfer energy to the T_1 state of β -carotene. However, the ISC processes into the T_2 and T_3 states also quench the S_1 state energy by a commensurate rate with the S_1 – T_1 process. Most likely, both k_{ISC2} and k_{ISC3} rate constants of the order 10^7 – 10^8 s^{−1} were erroneously included in the IC channel by subtracting ϕ_f and ϕ_{ISC} (S_1 – T_1) from unity. Indeed, the calculated total rate constant for the non-radiative quenching ($k_{IC} + \sum k_{ISCi}$) is comparable with the experimental data for both compounds, while the ratio between the IC and ISC rates is strongly different compared with our calculations (theor.: $\sum k_{ISCi} \gg k_{IC}$, but exp.: $\sum k_{ISCi} \approx k_{IC}$). One more argument in favor of our interpretation that the experimental k_{ISC} rate is underestimated in ref. 22 and 23 is that the opposite internal heavy atom effect was obtained^{22,23} when comparing k_{ISC} values for 4S4SiEt₂ (4.6×10^8 s^{−1}) and 4S4GeEt₂ (1.3×10^8 s^{−1}). This contradicts the fundamental mass-dependent property of SOC effects (SOC increases with the rise of the atomic number of an element)⁵² and also the experimentally estimated radiative phosphorescence lifetimes ($\tau_{phos}(4S4SiEt_2) = 235$ ms > $\tau_{phos}(4S4GeEt_2) = 98$ ms), which correctly reproduce the internal heavy atom effect. We should stress here that the total phosphorescence lifetimes ($\tau_{phos,tot.}(4S4SiEt_2) = 33$ s, $\tau_{phos,tot.}(4S4GeEt_2) = 1.2$ s), estimated taking into account



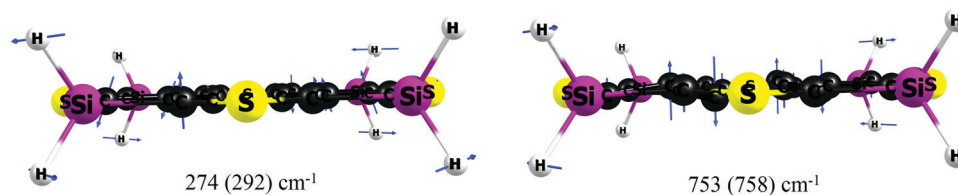


Fig. 6 Scaled displacement vectors for the two out-of-plane E_g modes responsible for the enhancement of the SOC matrix elements upon geometry distortion of the $4S4SiH_2$ ($4S4GeH_2$) molecule. Both modes are quite similar and their effects on the SOC matrix elements are almost identical (deviations are less than 0.02 cm^{-1}).

phosphorescence quantum yields (*i.e.* accounting for non-radiative quenching), are in excellent agreement with our calculations by the second-order approximate coupled-cluster (CC_2) method⁵⁷ (31 s and 2 s, respectively), indicating a high-accuracy of the theoretical calculations of the SOC effects. A more detailed computational study of the phosphorescence of $4S4SiEt_2$ and $4S4GeEt_2$ can be found in ref. 58.

Finally, we want to stress that despite the quantitative mismatch between the calculated and experimental k_{ISC} and k_{IC} rates, the general conclusion obtained from theory and experiment is the same: ISC is the main deactivation channel of the S_1 state energy compared with the fluorescence and internal conversion processes. The main reason lies in the quite large SOC matrix elements between the S_1 and T_n ($n = 1-3$) states despite the similar $\pi\pi^*$ nature of both the singlet and triplet counterparts. For common organic molecules the SOC matrix elements between S_1 and T_n states are of the order of 0.1 cm^{-1} .^{52,59} It means that $4S4SiEt_2$ and $4S4GeEt_2$ molecules possess quite specific photophysical behavior due to their unique circulenoid structure and due to the presence of heavy Si and Ge atoms in the structure of these molecules.

4. Conclusions

In the present work we have comprehensively studied the electronic structure, spectra, aromaticity and photophysical properties of two recently synthesized derivatives of tetrathienylene named tetrasilatetrathia[8]circulene ($4S4SiEt_2$) and tetragermatetrathia[8]circulene ($4S4GeEt_2$). Both studied molecules possess double aromaticity characterized by a local magnetically-induced current in contrast to the space-separated global current model for most hetero[8]circulenes. Generally, both $4S4SiEt_2$ and $4S4GeEt_2$ circulenes can be considered as globally non-aromatic species in terms of the magnetically-induced current criterion; the local diatropic currents inside the four thiophene rings are completely canceled by the paratropic currents inside the inner eight-membered core producing a close to zero net current strength typical for non-aromatic species. Both the $4S4SiEt_2$ and $4S4GeEt_2$ species demonstrate clearly visible vibronic structure in the absorption and emission spectra due to the activity of the stretching and bending C–C vibrations of the tetrathienylene skeleton. At the same time, the low-frequency CH bending vibrations play a crucial role for the fluorescence vibronic spectra. Accounting for spin–vibronic interactions through geometry distortion along the several out-of-plane E_g modes provides

reliable agreement between the calculated and experimental photophysical constants. At the same time, neglecting the spin–vibronic interactions gives an underestimation of the spin–orbit coupling matrix element between the S_1 and T_1 states, which subsequently leads to a strong overestimation of the fluorescence rate and quantum yield. Moreover, the fact that several triplet excited states lie lower in energy than the S_1 state plays a crucial role in the balance between the internal conversion and inter-system crossing rates. Accounting for this assumption we have concluded that the inter-system crossing channel constitutes a main deactivation process responsible for the deactivation of the S_1 state energy. Actually, this paper contains for the first time a detailed photophysical investigation of representatives of the hetero[8]circulene family and provides a close comparison with experimental data. We hope that the current study will help to generalize and interconnect the aromaticity and spectral theory of hetero[8]circulenes as a unique class of totally-annulated cyclooctatetraenes.

Conflicts of interest

There are no conflicts to declare.

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