# **Materials** Advances



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# PAPER

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# Enhancing thermometric precision: modulating the temperature of maximum sensitivity via erbium dopant addition in Ba<sub>2</sub>GdV<sub>3</sub>O<sub>11</sub>:Tm<sup>3+</sup>/Yb<sup>3+</sup> nano phosphors<sup>†</sup>

Ikhlas Kachou,<sup>a</sup> Kamel Saidi, 🕩 <sup>a</sup> Christian Hernández-Álvarez,<sup>b</sup> Mohamed Dammak ( \* and Inocencio R. Martín \*

Developing luminescence sensors often prioritizes maximizing relative sensitivity to achieve optimal performance. However, a critical parameter often overlooked is the temperature at which maximum sensitivity occurs. In this study, we delve into this crucial aspect by exploring the impact of erbium doping in Tm<sup>3+</sup>/Yb<sup>3+</sup> co-doped Ba<sub>2</sub>GdV<sub>3</sub>O<sub>11</sub> nano phosphors. The crystal structure, microscopic morphology, and luminescence mechanism of BGVO:Yb<sup>3+</sup>/Tm<sup>3+</sup> and Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> up conversion nanoparticles, as well as the temperature sensing characteristics are investigated. Under 975 nm laser excitation, the BGVO:Yb<sup>3+</sup>/Tm<sup>3+</sup> and BGVO:Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> nano phosphors exhibited strong blue and green upconversion luminescence, respectively. The luminescence intensity ratio (LIR) approach was used to analyze the temperature-dependent luminescence spectra in the 300-600 K temperature range. The thermometry strategies were based on thermally coupled energy levels (TCLs) and nonthermally coupled energy levels (NTCLs) of  $Er^{3+}$  and  $Tm^{3+}$  for temperature sensing performance. In the  $Tm^{3+}/Yb^{3+}$  codoped samples, the relative sensitivity typically peaks around 350 K, attributed to TCLs (1.7% K<sup>-1</sup>, 700 nm/800 nm) with generally lower relative sensitivity compared to non-TCLs (5.39% K<sup>-1</sup>, 700 nm/475 nm). However, non-TCL sensitivities in the 300–600 K range lack a clear maximum. In contrast, Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> samples exhibit distinct maxima in non-TCL sensitivities within this temperature range (1.91% K<sup>-1</sup>, 700 nm/550 nm), offering precise temperature determination for specific applications. Our findings underscore the potential of erbium doping to modulate temperature sensitivity peaks, crucial for optimizing performance in tailored luminescence nanosensors and offering fresh concepts for investigating alternative superior optical temperature sensing nano materials.

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# Introduction

Temperature is a fundamental parameter in thermodynamics, essential for scientific inquiry, climatology, industrial production, and routine tasks. However, traditional temperature sensors often face limitations due to susceptibility to electromagnetic interference and potential errors from component heating.<sup>1-5</sup> Consequently, the development of non-contact temperature detection methods, such as fluorescence thermometry, has gained momentum.<sup>6</sup> A key spectral characteristic utilized in fluorescence thermometry is the

luminescence intensity ratio (LIR), which undergoes changes with temperature variations.<sup>7-10</sup> LIR thermometry, also known as selfreferencing radiometric thermometry, reduces reliance on external factors like excitation intensity fluctuations or electromagnetic interference.<sup>11</sup> Luminescence nano thermometry predominantly employs optical sensors based on Ln3+ ions due to their advantageous properties.<sup>12</sup> These sensors commonly utilize radiometric techniques, particularly assessing band intensity ratios.<sup>7,13,14</sup> While ions like Er<sup>3+</sup>, Tm<sup>3+</sup>, and Nd<sup>3+</sup> are frequently utilized due to their thermally coupled levels (TCLs), their relative sensitivity  $(S_r)$  to temperature is limited to specific energy differentials between TCLs, resulting in relatively low Sr values. To address this limitation, recent research has focused on utilizing band intensities from nonthermally coupled levels (non-TCLs) of different Ln<sup>3+</sup> ions, offering theoretically limitless temperature sensitivity.<sup>7,13,14</sup>

Luminescent nanothermometers use a wide range of materials, including inorganic fluorides, phosphates, vanadate doped with optically active d-block metal ions, and lanthanide

<sup>&</sup>lt;sup>a</sup> Laboratoire de Physique Appliquée, Groupe des Matériaux Luminescents, Faculté des Sciences de Sfax, Département de Physique, Université de Sfax, BP 1171 Sfax, Tunisia. E-mail: madidammak@yahoo.fr, mohamed.dammak@fss.usf.tn

<sup>&</sup>lt;sup>b</sup> Universidad de La Laguna, Departamento de Física, MALTA-Consolider Team and IMN. Apdo. Correos 456, E-38206, San Cristóbal de La Laguna, Santa Cruz de Tenerife, Spain

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ions, along with carbon dots and quantum dots.<sup>8,15,16</sup> These matrices offer strong photoluminescence signals and can be employed to create highly efficient phosphors due to their low phonon energies and thermal stability. Lanthanide-based inorganic compounds, including nanoparticles (NPs), possess a multitude of advantageous characteristics, such as broad-range emission induced by UV, visible, or NIR irradiation, long luminescence lifetimes, narrow emission bands, large spectral shifts, stability in the (+3) oxidation state, resistance to photobleaching, and high-temperature treatment.<sup>17-19</sup> These advantageous characteristics are primarily attributed to several factors: the forbidden nature of 4f-4f transitions in Ln<sup>3+</sup> ions, crystal-field effects, and the shielding of 4f electrons by 5s and 5p electrons.<sup>20</sup> Moreover, lanthanide ions (Ln<sup>3+</sup>) exhibit up conversion (UC) luminescence, which involves the emission of higher-energy photons due to multiphoton absorption of lower-energy radiation, typically in the NIR range.<sup>21,22</sup> Lanthanide ions serve as both luminescence activators (e.g.,  $Eu^{3+}$ ,  $Tb^{3+}$ , Nd<sup>3+</sup>, Er<sup>3+</sup>, and Tm<sup>3+</sup>) and sensitizers (mostly Ce<sup>3+</sup> and Yb<sup>3+</sup>) due to their advantageous spectroscopic characteristics, including long luminescence lifetimes, wide range of emission states, large Stokes shift, and the capacity for upconversion.<sup>23</sup> Upconverting nanomaterials find applications in biological sensing, particularly for subcutaneous sensing, owing to their small size and the use of NIR laser excitation to avoid damaging UV radiation.<sup>24</sup>

This study investigates the influence of erbium  $(Er^{3+})$  doping in  $Tm^{3+}/Yb^{3+}$  codoped  $Ba_2GdV_3O_{11}$  (BGVO) nano phosphors, focusing on its impact on luminescence and thermometric properties. While maximizing relative sensitivity is a primary goal in luminescence-based temperature sensing, the temperature at which maximum sensitivity occurs and the temperature uncertainty are equally crucial but often overlooked. By incorporating erbium ions as a third dopant alongside  $Tm^{3+}$  and  $Yb^{3+}$  ions, the study reveals insights into how this addition can significantly affect luminescence and thermometric properties.

We present here the inorganic nano phosphors BGVO:Yb<sup>3+</sup>/ Tm<sup>3+</sup> and tridoped Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup>, which were created for temperature sensing and exhibit strong, perceptible UC luminescence ( $\lambda_{ex} = 975$  nm). The produced nanomaterial exhibits both phononassisted emission from ytterbium ions (above  $\approx 900$  nm) and UC luminescence from thulium ions. The material's luminescence characteristics were examined in the 400–960 nm range, which includes the first biological window. The relative sensitivities and temperature resolutions for the produced nanomaterial were determined using the band intensity ratios of the emission bands of Yb<sup>3+</sup>, Er<sup>3+</sup>, and Tm<sup>3+</sup> ions situated in that wide spectrum range.

# Synthesis procedure

#### Synthesis of BaGdV<sub>3</sub>O<sub>11</sub> doped Yb<sup>3+</sup>, Tm<sup>3+</sup> and Er<sup>3+</sup>

The Ba<sub>2</sub>GdV<sub>3</sub>O<sub>11</sub> structures doped with 2%Er<sup>3+</sup>, 1%Tm<sup>3+</sup> and 15%Yb<sup>3+</sup> were synthesized using the Pechini sol–gel technique. Nano phosphors are prepared by acquiring multiple chemicals, including barium nitrate [Ba(NO<sub>3</sub>)<sub>2</sub> (99.0%)], gadolinium nitrate hexahydrate [Gd(NO<sub>3</sub>)<sub>3</sub>·6H<sub>2</sub>O (99.9%)], erbium nitrate

pentahydrate [(Er(NO<sub>3</sub>)<sub>3</sub>·5H<sub>2</sub>O) (99.9%)], thulium nitrate pentahydrate [Tm(NO<sub>3</sub>)<sub>3</sub>·5H<sub>2</sub>O (99.9%)], ytterbium nitrate pentahydrate [Yb(NO<sub>3</sub>)<sub>3</sub>·5H<sub>2</sub>O (99.9%)], ammonium metavanadate [NH<sub>4</sub>VO<sub>3</sub> (99.96%)] and citric acid [C<sub>6</sub>H<sub>8</sub>O<sub>7</sub> (99.0%)]. To initiate the process, all the precursors are placed into an Erlenmeyer flask and stirred using a magnetic stirrer at 70 °C. Following that, citric acid is introduced, causing the solution to transition from green to blue in color. The mixture is gently swirled and then heated to 80 degrees Celsius until it undergoes hydrolysis, transforming into a sol, and ultimately forming a gel. Subsequently, the gel is subjected to calcination through annealing at 700 °C. The final product is further ground for additional analysis.

#### Characterization

The XRD patterns were measured in the  $2\theta$  range of  $(20-50^{\circ})$  on an X-ray diffractometer (Bruker D8, Germany) with Cu K $\alpha$  radiation (1.5405 Å), working at 40 kV and 30 mA for phase detection. A UV vis-NIR spectrometer (PerkinElmer Lambda 950) was used to measure UV-vis-NIR absorption. The samples' morphologies were examined using a scanning electron microscope (JEOL, JSM 6510LV). The sample is positioned within a homemade nano heater chamber and here heating is regulated with an accuracy of  $\pm 0.5$  K. Utilizing 975 nm diode laser stimulation at a consistent pump power of 30 mW, emission spectra were captured using a monochromator spectrometer (Horiba Jobin Yvon, iHR320).

### Results and discussion

#### Structural properties and particle morphology

Phase composition and crystallinity of luminous materials play a crucial role in their performance. In general, luminophores with high crystallinity exhibit both stronger and weaker luminescent characteristics. The crystallinity of the powder composed of all-nano phosphors was assessed using X-ray diffraction (XRD). Fig. S1 (ESI<sup>†</sup>) displays the powder X-ray diffraction (XRD) pattern of the BYVO:2%Er3+/1%Tm3+/ 15%Yb<sup>3+</sup> and BGVO:1%Tm<sup>3+</sup>/15%Yb<sup>3+</sup> nano phosphor. The XRD patterns are closely matched the JCPDS card no. 81-1778 for all the synthesized nano phosphors.<sup>25,26</sup> Rietveld refinement displayed in Fig. 1a and b was done based on the XRD data of all samples. The Rietveld findings indicate that these compounds crystallized in a monoclinic system with the space group  $P2_1/c$ . The unit cell parameters are calculated to be a =12.376 Å, b = 7.732 Å, c = 11.179 Å and  $\beta = 103.486^{\circ}$ . The absence of impurity peaks in the prepared nano phosphors is confirmed by the XRD pattern. The XRD patterns remained unchanged even when the Er<sup>3+</sup>, Yb<sup>3+</sup>, and Tm<sup>3+</sup> ions were substituted for the Y<sup>3+</sup> site, indicating the successful incorporation of these ions into the lattice structure. This outcome can be attributed to the similarity in ionic radii among  $Y^{3+}$  (CN = 6; r = 0.90 Å),  $\text{Er}^{3+}$  (CN = 6; r = 0.95 Å), Yb<sup>3+</sup> (CN = 6; r = 0.95 Å), and Tm<sup>3+</sup>  $(CN = 6; r = 0.88 \text{ Å}) \text{ ions.}^{27}$ 

A transmission electron microscopy (TEM) technique was employed to examine the surface morphology and distribution of grain sizes. In Fig. 2, we have provided TEM images of both

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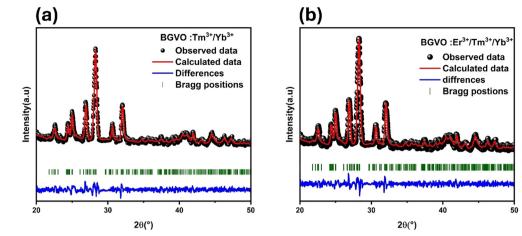


Fig. 1 (a) Rietveld refinement of XRD data of BGVO:1%Tm<sup>3+</sup>/15%Yb<sup>3+</sup> and (b) Rietveld refinement of XRD data of BGVO:2%Er<sup>3+</sup>/1%Tm<sup>3+</sup>/15%Yb<sup>3+</sup>.

doped and undoped samples. It is evident that the irregular shapes observed in the images are due to the aggregation of grains with sizes ranging from 300 to 500 nm.

The diffuse reflectance spectra are displayed in Fig. 3 of BGVO co-doped Tm<sup>3+</sup>/Yb<sup>3+</sup> and tri-doped Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> in the range (200–1100 nm). The spectrum presents an intense broad band peaking at 280 nm, which is attributed to the O<sup>2–</sup>–V<sup>5+</sup> entity. However, this spectrum shows four distinct absorption peaks at 525, 648, 682 and 796 nm. These peaks correspond to the transitions of the Er<sup>3+</sup> and Tm<sup>3+</sup> ions, respectively. The  ${}^{2}F_{7/2} \rightarrow {}^{2}F_{5/2}$  (Yb<sup>3+</sup>) absorption transition produces a broad absorption band with a peak at 975 nm.<sup>7</sup>

# Luminescence measurements

#### Characteristics of emissions at ambient temperature

After being stimulated by a 975 nm laser, the synthesized nanomaterial exhibits a pronounced, easily discernible blue up-conversion (anti-Stokes) luminescence to the naked eye. Fig. 4a represents the emission spectra for the sample co-doped with  $Yb^{3+}/Tm^{3+}$  when it is excited with 975 nm wavelength. The spectrum is the result of several sharp bands typical of  $Tm^{3+}$  ions in the spectral region of 450–850 nm and one

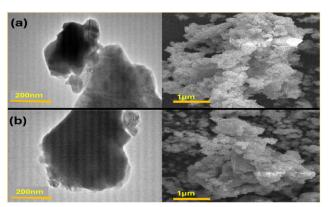


Fig. 2 SEM images for BGVO doped (a)  $1\% Tm^{3+}/15\% Yb^{3+}$  and (b)  $2\% Er^{3+}/1\% Tm^{3+}/15\% Yb^{3+}.$ 

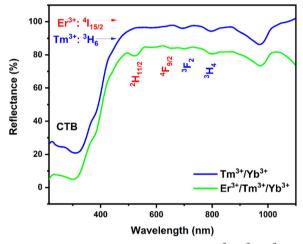


Fig. 3 The reflectance spectra of BGVO doped  $Tm^{3+}$ ,  $Er^{3+}$ ,  $Yb^{3+}$ 

band more than 900 nm typical of Yb<sup>3+</sup> ions that is partially filtered by the short-pass 950 nm filter used.<sup>2</sup> All detected bands are associated with the 4f-4f intra-configurational radiative transition of Tm<sup>3+</sup>:<sup>1</sup>G<sub>4</sub>  $\rightarrow$  <sup>3</sup>H<sub>6</sub> ( $\approx$ 470 nm), <sup>1</sup>G<sub>4</sub>  $\rightarrow$  <sup>3</sup>F<sub>4</sub> ( $\approx$ 640 nm), <sup>3</sup>F<sub>2,3</sub>  $\rightarrow$  ${}^{3}\text{H}_{6} (\approx 700 \text{ nm}), {}^{3}\text{H}_{4} \rightarrow {}^{3}\text{H}_{6} (\approx 800 \text{ nm}), \text{ and } \text{Yb}^{3+} : {}^{2}\text{F}_{5/2} \rightarrow {}^{2}\text{F}_{7/2}$ (>900 nm).<sup>28-30</sup> This system employs Yb<sup>3+</sup> as a sensitizer, pumping the excited states of Tm<sup>3+</sup> (UC emitter) primarily via energy transfer UC processes when absorbed NIR light is transferred to neighboring Tm<sup>3+</sup>. The BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup> nanomaterial demonstrates heightened near-infrared (NIR) emission at low excitation densities, offering potential utility in applications such as deep tissue imaging, optical signal amplification, and various other fields.<sup>31,32</sup> In another way, the UC emission spectra shown in Fig. 4b of the BGVO tri-doped Tm<sup>3+</sup>, Er<sup>3+</sup>, and Yb<sup>3+</sup> nano phosphors with a 975 nm laser diode were detected in the 450-960 nm wavelength range. Five UC emission peaks can be detected at 470, 487, 523, 538, 655, and 800 nm, associated to the  ${}^{1}G_{4} \rightarrow {}^{3}H_{6}$  $(Tm^{3+}), {}^{2}H_{11/2} \rightarrow {}^{4}I_{15/2} (Er^{3+}), {}^{4}S_{3/2} \rightarrow {}^{4}I_{15/2} (Er^{3+}),$  the combined  ${}^{1}G_{4} \rightarrow {}^{3}F_{4} (Tm^{3+})/{}^{4}F_{9/2} \rightarrow I_{15/2} (Er^{3+}) and {}^{3}H_{4} \rightarrow {}^{3}H_{6} (Tm^{3+})$ transitions, respectively, and of Yb<sup>3+</sup>: ${}^{2}F_{5/2} \rightarrow {}^{2}F_{7/2} (>900 \text{ nm}).$ 

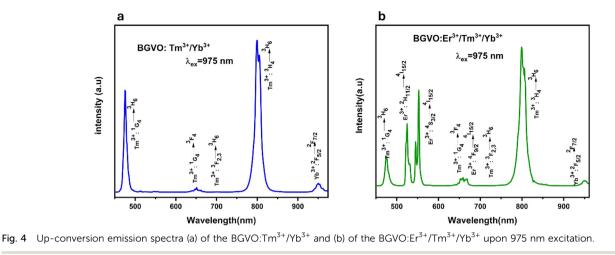


Fig. S2 (ESI<sup>†</sup>) shows these results and a rise time and a decay time for the  $Tm^{3+}/Yb^{3+}$  and  $Er^{3+}/Tm^{3+}/Yb^{3+}$  doped BGVO nano phosphors have been measured using the 975 nm excitation. These curves are characteristics of upconversion processes, and the temporal evolution depends on the lifetimes of the involved intermediate levels.

To evaluate the transition mechanism of different emission peaks, we also examine the correlation between the luminous intensity of the UP (upconversion) and the power of the pump. The intensity of each emission peak in the upconversion process is determined by a power function of the power of infrared excitation. The formula is as follows:

$$I = K \cdot P^n \tag{1}$$

*P* denotes the pump energy, and *I* denotes the UC fluorescence intensity, while *n* signifies the number of excitation photons required to generate a single UC photon. Fig. 5a illustrates the correlation between the natural logarithm of UC fluorescence intensity and the natural logarithm of BGVO: $Tm^{3+}/Yb^{3+}$  pump power. The slope of the curve is 2.3 (475 nm), and 1.2 (800 nm), indicating that the blue UC is a three-photon process, while NIR emission of BGVO: $Tm^{3+}/Yb^{3+}$  is a two-photon process. Fig. 4b demonstrates the presence of blue and red emissions upon 975 nm excitation for the BGVO: $Er^{3+}/Tm^{3+}/Yb^{3+}$  tri doped nano phosphor. The relationship between the natural logarithm of UC intensity and the natural logarithm of pump power is represented in Fig. 5b. The slope of the curve is 2.10 (475 nm), 1.26 (545 nm) and 1.48 (650 nm), and 1.12 (800 nm), indicating that all emission peaks are a two-photon process, except for the 475 nm one that can be produced by a three-photon process. These values (see Fig. S2, ESI†) could indicate that in the codoped sample there is transfer from the Er to Tm ions.

Fig. 6 shows the schematic diagram of the energy levels of Yb<sup>3+</sup>, Er<sup>3+</sup> and Tm<sup>3+</sup>, as well as energy transfer and non-radiative transition processes. It can be seen from the figure that more effective energy conversion successively occurs from Yb<sup>3+</sup> to Er<sup>3+</sup> and Tm<sup>3+</sup> under 975 nm excitation. For the UP mechanism of Yb<sup>3+</sup> and Er<sup>3+</sup> ions, the <sup>2</sup>F<sub>7/2</sub> ground state of Yb<sup>3+</sup> ion absorbs a 975 nm photon energy and transitions to the <sup>2</sup>F<sub>5/2</sub> excited state. The <sup>2</sup>F<sub>5/2</sub> excited state of the Yb<sup>3+</sup> ion transfers energy to the <sup>4</sup>I<sub>11/2</sub> energy level of Er<sup>3+</sup>, and the electrons at the <sup>4</sup>I<sub>11/2</sub> level of Er<sup>3+</sup> rapidly relax to the <sup>4</sup>I<sub>13/2</sub> level by non-radiative transition. Er<sup>3+</sup> at the <sup>4</sup>I<sub>13/2</sub> level can transfer energy with one Yb<sup>3+</sup> in the excited state and be excited to the <sup>4</sup>I<sub>9/2</sub> level. Er<sup>3+</sup> at the <sup>4</sup>I<sub>9/2</sub> level rapidly transitions to the ground state to produce upconverted red light. Simultaneously, Yb<sup>3+</sup> reaches the excited state under 975 nm excitation. The energy transfer between the excited Yb<sup>3+</sup> and Er<sup>3+</sup>

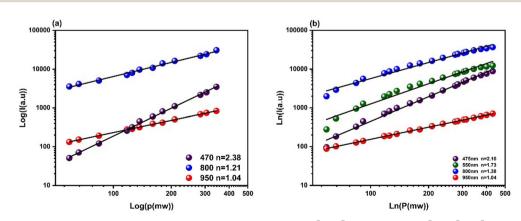


Fig. 5 Dependence of the intensity with the pump power at 975 nm of (a) BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup> and (b) BGVO:Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup>.

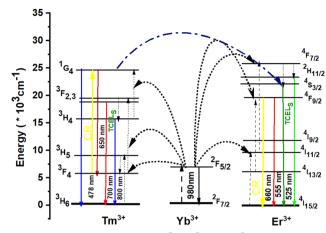


Fig. 6 Energy levels diagram of  $Tm^{3+}$ ,  $Yb^{3+}$ , and  $Er^{3+}$  ions in the BGVO host.

makes the  $\text{Er}^{3+}$  of the  ${}^{4}\text{I}_{11/2}$  level transition to the  ${}^{2}\text{H}_{11/2}$  and  ${}^{4}\text{S}_{3/2}$ levels.  $\text{Er}^{3+}$  at the  ${}^{2}\text{H}_{11/2}$  and  ${}^{4}\text{S}_{3/2}$  levels rapidly transitions to the ground state to produce upconverted green light. For the UC mechanism of  $Yb^{3\scriptscriptstyle +}$  and  $Tm^{3\scriptscriptstyle +}$  ions, the  $^2F_{7/2}$  ground state of Yb<sup>3+</sup> ion absorbs a 975 nm photon energy and transitions to the  $^2\!F_{5/2}$  excited state, the  $^2\!F_{5/2}$  excited state of the  $Yb^{3+}$  ion transfers energy to the <sup>3</sup>H<sub>5</sub> ground state of the Tm<sup>3+</sup> ion. The electrons at the <sup>3</sup>H<sub>5</sub> level of Tm<sup>3+</sup> rapidly relax to the <sup>3</sup>F<sub>4</sub> level by non-radiative transition. Tm<sup>3+</sup> at the <sup>3</sup>F<sub>4</sub> level can transfer energy with one Yb<sup>3+</sup> in the excited state and be excited to <sup>3</sup>F<sub>2,3</sub> levels.  $Tm^{3+}$  at the  ${}^{3}F_{2,3}$  levels rapidly relaxes to the  ${}^{3}H_{4}$  level or transitions to the ground state to produce upconverted red light, respectively. Tm<sup>3+</sup> at the <sup>3</sup>H<sub>4</sub> level can transfer energy with one Yb<sup>3+</sup> in the excited state and be excited to the <sup>1</sup>G<sub>4</sub> level. Tm<sup>3+</sup> at the <sup>1</sup>G<sub>4</sub> level transitions to the ground state to produce UC blue fluorescence. The resonant cross relaxation process (CR) <sup>3</sup>F<sub>4</sub>  $(Tm^{3+}) + {}^{4}F_{9/2} (Er^{3+}) \rightarrow {}^{1}G_{4} (Tm^{3+}) + {}^{4}I_{15/2} (Er^{3+})$  between  $Tm^{3+}$ and Er<sup>3+</sup> ions leads to the population of the <sup>1</sup>G<sub>4</sub> state of Tm<sup>3+</sup> ions and depopulation of the <sup>4</sup>F<sub>9/2</sub> state of Er<sup>3+</sup> ions, and then to the  ${}^{2}H_{11/2}/{}^{4}S_{3/2}$  state through the multiphonon-assisted relaxations.

#### **Temperature sensing**

Temperature is an important parameter in scientific research, healthcare, and manufacturing, and reliable temperature measurement is necessary for practical applications. First, it is important to note that BGVO was chosen as a host for the optically active lanthanide ions ( $Yb^{3+}$ ,  $Tm^{3+}$  and  $Er^{3+}$ ) primarily due to its resistance to high-temperature treatment, and the ability to dope with various lanthanide ions. This resulted in a product that was well-crystallized, nano-sized, and displayed bright and intense luminescence. Changes in temperature affect the intensity of the emission bands of  $Er^{3+}$ ,  $Tm^{3+}$  and  $Yb^{3+}$  ions that are observed.

**Codoped BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup> luminescent sensors.** Fig. 7 displays the UC spectra as the temperature increases. The intensity of the Tm<sup>3+</sup> bands gradually decrease due to the thermal heating process, except for the thermalized band around 700 nm, whose intensity increases. The Yb<sup>3+</sup> band also shows a similar effect of decreasing intensity. In contrast to Tm<sup>3+</sup> UC emissions, the rate of decrease is much lower. The alteration of color with varying temperature aligns with the chromaticity coordinates demonstrated in the CIE diagram (Fig. 7). These findings suggest that the BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup> materials can also find applications as visual optical temperature sensors, leveraging their thermochromic properties. Furthermore, the temperature-sensitive and adjustable luminescence of BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup> materials can be harnessed for safety signage purposes in high-temperature settings.

From the point of view of temperature sensing, we have focused on two diverse band intensity ratios, corresponding to the thermally coupled levels (TCLs) of Tm<sup>3+</sup> and the non-TCLs of Yb<sup>3+</sup> and Tm<sup>3+</sup>. Herein, the thermalization processes between the <sup>3</sup>F<sub>2,3</sub> and <sup>3</sup>H<sub>4</sub> excited levels of Tm<sup>3+</sup> are well-established in the literature.<sup>33</sup> These TCLs are separated by relatively large energy (usually  $\Delta E \approx 1500\text{-}2000 \text{ cm}^{-1}$ ), resulting in effective thermalization in a high-temperature range and good relative sensitivity. As the temperature increases the lower energy *I*<sub>1</sub> band (800 nm) decreases, and the intensity of the higher energy *I*<sub>2</sub> band (700 nm) increases, according to the Boltzmann-type

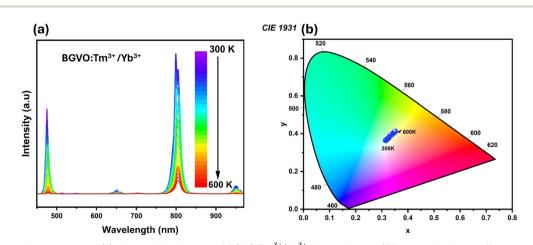


Fig. 7 (a) Upconversion spectra, and (b) chromaticity diagram of BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup> obtained upon 975 nm excitation at different temperatures.

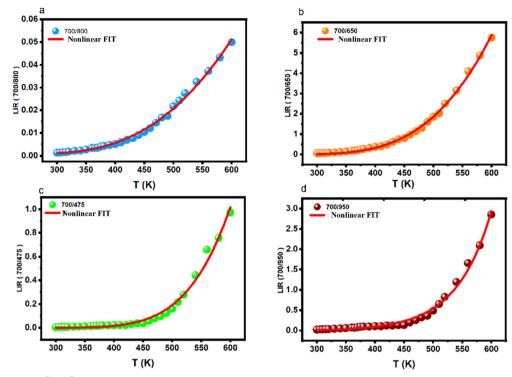


Fig. 8 LIR variations of Tm<sup>3+</sup>/Yb<sup>3+</sup>; (a) 700/800 nm, (b) 700/650 nm, (c) 700/475 nm and (d) 700/950 nm as a function of temperature.

distribution.34

$$\text{LIR}_{\text{TCL}} = \frac{I_2}{I_1} = A \times \exp\left(\frac{-\Delta E}{K_{\text{B}}T}\right)$$
(2)

where *A* is a parameter that depends on the host lattice, *T* is the absolute temperature,  $K_{\rm B}$  is the Boltzmann constant, and  $\Delta E$  is the energy gap between these two thermalized states. The temperature variation of the LIR (700/800 nm) value within the range of 300–603 K is illustrated in Fig. 8a. By employing eqn (2), we established a correlation between the calculated 700/800 nm band intensity ratio and temperature, yielding  $\Delta E = 1527$  cm<sup>-1</sup> and  $R_2 = 0.996$ , as shown in Fig. 8a.

Meanwhile, in the case of uncoupled levels (NTCL), it is challenging to populate them through thermal excitation because of their significant energy gaps. Consequently, traditional LIR is not suitable for NTCL. Experimental LIR data can be effectively modeled using a flowing function.<sup>35,36</sup>

$$\text{LIR}_{\text{NTCL}} = \frac{I_2}{I_1} = B \times \exp\left(\frac{-C}{K_{\text{B}}T}\right) + A \tag{3}$$

the LIR<sub>NTCL</sub> is linked to the non-thermally coupled levels (non-TCLs) of Tm<sup>3+</sup>, specifically referring to the intensity ratio of the 700/650 nm and 700/475 nm bands as displayed in Fig. 8b and c with ( $R^2 = 0.997$  and  $R^2 = 0.987$ ), respectively.

In contrast, the LIR parameter, linked to  $Yb^{3+}$  and  $Tm^{3+}$  NTCL, which corresponds to the 700/950 nm band intensity ratio, exhibited a correlation with temperature through the

utilization of the following multi-exponential function.

$$LIR = \frac{I_1}{I_2} = A_1 \times \exp\left(\frac{-T}{C_1}\right) + A_2 \times \exp\left(\frac{-T}{C_1}\right) + A_3 \times \exp\left(\frac{-T}{C_1}\right) + B$$
(4)

where  $A_1, A_2, A_3, C_1, C_2, C_3$  and *B* are fitting constants. The empirical function was used in this context due to the lack of a suitable physical model that could explain the observed temperaturedependent variations in the LIR parameter. It is worth mentioning that a similar empirical exponential function was employed by Marcin Runowski<sup>2</sup> to establish a correlation between the LIR of NTCL containing Tm<sup>3+</sup> and Yb<sup>3+</sup> ions at various temperatures. With the LIR derived from dual center Yb<sup>3+</sup> and Tm<sup>3+</sup> associated with the 700/950 nm band intensity ratio shown in Fig. 8d, we have achieved excellent fitting results with  $R^2 = 0.9906$ .

Sensitivity was a crucial and notable attribute when it came to quantitatively evaluating the appropriateness of materials for use as optical sensors in real-world applications. Moreover, the absolute sensitivity  $(S_a)$  and relative sensitivity  $(S_r)$  can be computed using the following formula. These sensitivities are defined by these two equations:

$$S_{\rm a} = \frac{\rm dLIR}{\rm dT} \tag{5}$$

$$S_{\rm r} = \frac{1}{\rm LIR} \frac{\delta \rm LIR}{\delta T} \tag{6}$$

In all cases, regardless of whether it is TCL, NTCL, singlecenter, or dual-center, the absolute sensitivity increases with

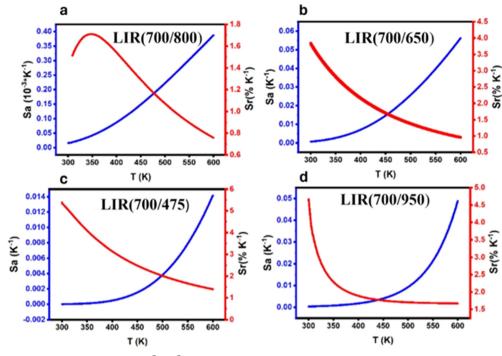


Fig. 9 The calculated  $S_a$  and  $S_r$  values based on Tm<sup>3+</sup>/Yb<sup>3+</sup> (a) LIR 700/800 nm, (b) LIR 700/650 nm, (c) LIR 700/475 nm and (d) LIR 700/950 nm as a function of temperature.

the rise in temperature, with a maximum of  $0.38 \times 10^{-3}$ ,  $0.056 \times 10^{-3}$ ,  $0.014 \times 10^{-3}$ , and  $0.048 \times 10^{-3}$ % K<sup>-1</sup> for 700/ 800 nm, 700/650 nm, 700/475 nm, and 700/950 nm, respectively. It's observed that in the case of TCL, the relative sensitivity decreases with an increase in temperature, reaching a maximum equal to 1.71% K<sup>-1</sup> at 348 K as it can be seen on Fig. 9a. In the cases of NTCL, Fig. 9b–d show that the relative sensitivity ratio decreases with increasing temperature. It's obvious that the  $S_r$  of the NTCL is higher than that of the TCL, and when comparing the NTCL in the single center (Tm<sup>3+</sup>/ Tm<sup>3+</sup>) and dual center (Yb<sup>3+</sup>/Tm<sup>3+</sup>), we can conclude that  $S_r$  in the single center (Tm<sup>3+</sup>/Tm<sup>3+</sup>) has excellent relative sensitivity, which is equal to 5.39% K<sup>-1</sup> at 300 K (Fig. 9c).

Temperature resolution ( $\delta T$ ) is a crucial parameter for evaluating the effectiveness of a temperature sensor. It denotes the smallest temperature variation that the sensor is capable of detecting. It is calculated by dividing the relative uncertainty in determining the ratiometric FIR value,  $\delta$ FIR/FIR, by the relative thermal sensitivity  $S_r$ .

$$\delta T = \frac{1}{S_{\rm r}} \cdot \frac{\delta {\rm LIR}}{{\rm LIR}} \tag{7}$$

The value of  $\delta T$  is influenced by the material's effectiveness, which is determined by  $S_{\rm r}$ , as well as the conditions of optical thermal detection and calibration, such as the signal-to-noise ratio. To optimize the resolution limit, one possibility is to employ enhanced acquisition methods and improve measurement performance. One way to achieve this is by increasing the integration time and taking the average of multiple measurements to minimize experimental noise.

Based on eqn (6), we have displayed all the computed temperature resolutions as a function of temperature in the range 300–600 K (see Fig. 10), the values obtained are below 0.5 K. This shows that both LIR methods demonstrate great accuracy over the whole temperature range. This statement offers an approximate sense of the thermal resolution that may be attained in temperature measurements with this luminous substance. It is worth noticing that the LIR (700/475 nm) (Fig. 10c), offers the best thermal resolution values (0.1 K). Based on the gathered data, it can be stated that the developed thermometer is appropriate for temperature sensing in both biological and industrial systems. The results reveal that the BGVO-based nano phosphor displays improved temperature resolution compared to the other ceramics, within the required operating temperature range.<sup>37</sup>

**Tridoped BGVO Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> luminescent sensors.** By including the Er<sup>3+</sup> as an additional luminescent center to augment the overall emission bands, the FIR of multiple emission bands can be calculated using various FIR fitting techniques, the approach towards multi-rate optical thermometry can be actualized, thus facilitating the objective of self-referenced temperature measurement.

To investigate how up-conversion luminescence changes with temperature, we collected emission spectra from the BGVO: $\mathrm{Er}^{3^+}/\mathrm{Tm}^{3^+}/\mathrm{Yb}^{3^+}$  nano phosphors. The up-conversion emission spectra at various temperatures are presented in Fig. 11a. It is evident that the temperature has a noticeable impact on the relative intensities of the measured upconversion (UC) emission bands of  $\mathrm{Er}^{3^+}$  and  $\mathrm{Tm}^{3^+}$  ions. The gradual decrease in luminescence intensity with increasing

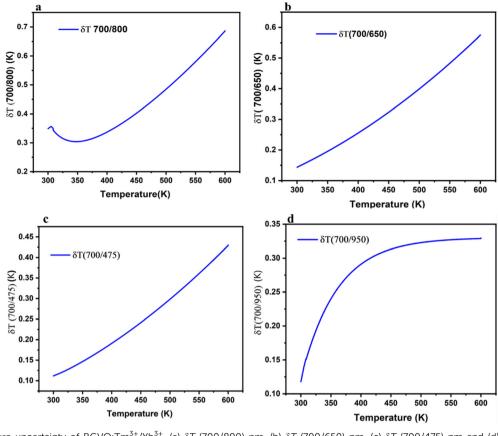


Fig. 10 Temperature uncertainty of BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup>, (a)  $\delta T$  (700/800) nm, (b)  $\delta T$  (700/650) nm, (c)  $\delta T$  (700/475) nm and (d)  $\delta T$  (700/950) under 975 nm.

temperature can be attributed to nonradiative relaxation, energy transfer, and crossover processes, all playing a role in this effect. It's important to highlight that the emission band around ~700 nm (corresponding to the  ${}^{3}F_{2,3} \rightarrow {}^{3}H_{6}$  transition of Tm<sup>3+</sup>) is initially absent (close to the noise level) at room temperature but becomes visible and intensifies as the temperature increases. The 950 nm band of Yb<sup>3+</sup> exhibits a similar trend of decreasing intensity as well. The determined FIR values of the emission bands originating from TCLs of  $\text{Er}^{3+}$  ions ( ${}^{2}\text{H}_{11/2}/{}^{4}\text{S}_{3/2}$  *i.e.* 525/550), and Tm<sup>3+</sup> ions ( ${}^{3}\text{F}_{2,3}/{}^{3}\text{H}_{4}$  *i.e.* 700/800) are well-fitted with a value of  $R^{2}$  = 0.999 and 0.995 based on eqn (1), as depicted in Fig. 12a and b, respectively. Furthermore, the 700/475, 700/650, 700/550 and 700/950 ratios are analyzed using nonlinear fits (see Fig. 12c–f).

In all analyzed cases, regardless of the energy level type (TCL, NTCL, single-center, or dual-center), the absolute

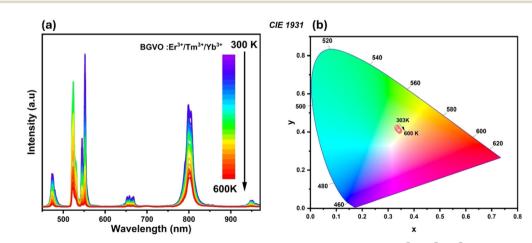


Fig. 11 (a) Temperature-dependence of the emission spectra; (b) chromaticity diagram of the BGVO:Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> sample.

sensitivity increased with rising temperature. The maximum  $S_{\rm a}$  observed was:  $11 \times 10^{-3} \ {\rm K}^{-1}$  (525/550 nm), 0.22  $\times 10^{-3} \ {\rm K}^{-1}$  (700/800 nm),  $10.0 \times 10^{-3} \ {\rm K}^{-1}$  (700/475 nm), 0.8  $\times 10^{-3} \ {\rm K}^{-1}$  (700/650 nm), 0.6  $\times 10^{-3} \ {\rm K}^{-1}$  (700/550 nm), and 17.0  $\times 10^{-3} \ {\rm K}^{-1}$  (700/950 nm) shown in Fig. 13a–f. Notably, the  $S_{\rm a}$  values were significantly higher for  ${\rm Er}^{3+}/{\rm Tm}^{3+}/{\rm Yb}^{3+}$  tri-doped BGVO nano phosphors compared to  ${\rm Tm}^{3+}/{\rm Yb}^{3+}$  co-doped samples. This observation highlights the role of  ${\rm Er}^{3+}$  doping in enhancing this crucial parameter.

The variation of the  $S_r$  parameter with temperature exhibits similar behavior for both  $Tm^{3+}/Yb^{3+}$  and  $Er^{3+}/Tm^{3+}/Yb^{3+}$  codoped samples in the case of TCLs (see Fig. 13a and b). However, for non-thermally coupled energy levels (NTCL), a significant difference is observed. Unlike  $\text{Tm}^{3+}/\text{Yb}^{3+}$  co-doped samples,  $\text{Er}^{3+}/\text{Tm}^{3+}/\text{Yb}^{3+}$  co-doped samples exhibit a well-defined maximum  $S_r$  value ( $S_{\text{rmax}}$ ) at specific temperatures. This behavior holds promise for the design of high-sensitivity sensors within a well-defined temperature range.

The LIR values of non-TCLs for single centers (700/650), (700/475), and dual centers  $(Tm^{3+}/Er^{3+})$  (700/550) and  $(Tm^{3+}/Yb^{3+})$  (700/950) at temperatures ranging from 300 K to 600 K are displayed in Fig. 13c–f. Therefore, eqn (3) can be successfully utilized to fit the data for (700/650), (700/475), and (700/550), while eqn (4) is suitable for fitting (700/950). The resulting

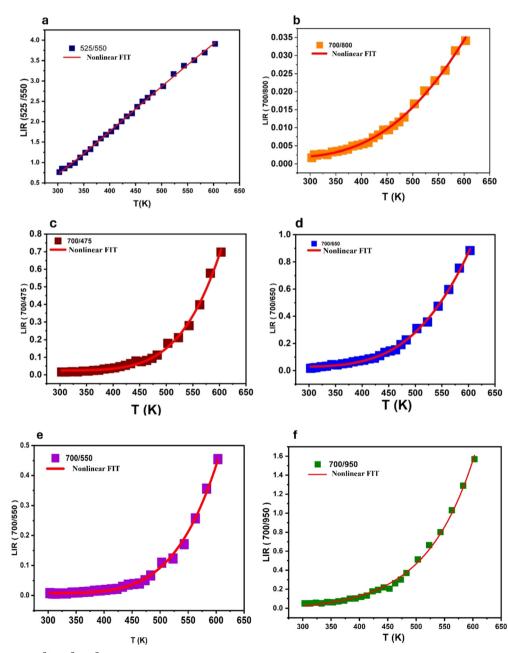


Fig. 12 LIR variations of Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup>; (a) 525/550 nm, (b) 700/800 nm, (c) 700/475 nm, (d) 700/650 nm, (e) 700/550 nm and (f) 700/950 nm as a function of temperature.

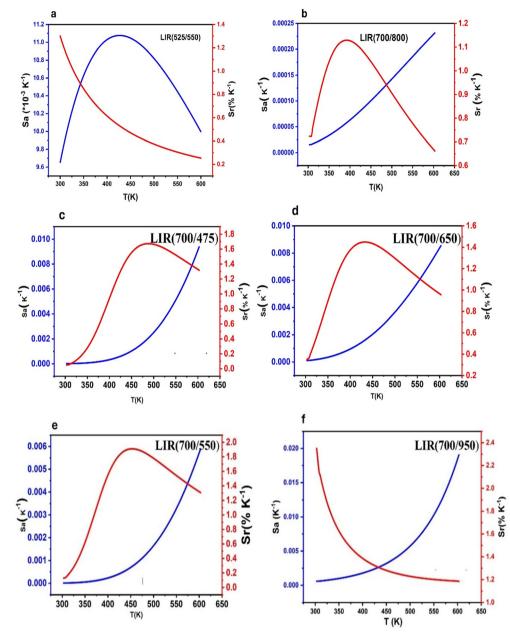


Fig. 13 The calculated  $S_a$  and  $S_r$  values based on  $Er^{3+}/Tm^{3+}/Yb^{3+}$ : (a) 700/800 nm, (b) 700/650, (c) 700/475, (d) 700/550 nm, (e) 525/550 nm and (f) 700/950 nm as a function of temperature.

variations of  $S_a$  and  $S_r$  with temperature are depicted in Fig. 13c–e.

Remarkably, the utilization of NTCL resulted in a significantly higher  $S_{\rm rmax}$  value of 2.35% K<sup>-1</sup> at 300 K. This achievement was made possible by employing the combination of upconversion (UC) emission bands at 700/950 nm, originating from the transitions of Tm<sup>3+</sup>/Yb<sup>3+</sup> in the tridoped BGVO:Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> nano phosphor.

The maximum  $S_r$  values ( $S_{rmax}$ ) for the  $Er^{3+}/Tm^{3+}/Yb^{3+}$  codoped samples were determined to be: 1.3, 1.12, 1.67, 1.45, 1.91 and 2.35 K<sup>-1</sup> for emission wavelengths of (525/550 nm), (700/800 nm), (700/475), (700/650), (700/550), and (700/950), respectively (shown in Fig. 13a–f). These values are slightly lower than those observed in the  $\text{Tm}^{3+}/\text{Yb}^{3+}$  co-doped samples, but with the advantage of a well-defined  $S_{\text{rmax}}$  at a specific temperature.

Fig. 14a–f illustrates the temperature uncertainties ( $\delta T$ ) for all the samples. Notably, the addition of Er<sup>3+</sup> doping to BGVO:Tm<sup>3+</sup>/Yb<sup>3+</sup> significantly improves the temperature resolution. In particular, for the 700/550 and 700/445 ratios,  $\delta T$  reaches an impressive value of 0.02 K. This observation further emphasizes the crucial role of Er<sup>3+</sup> doping in enhancing the performance of these optical temperature sensors.

Table 1 shows that compared to other lanthanide-based optical thermometers, the  $Ba_2GdV_3O_{11}$ : $Er^{3+}/Tm^{3+}/Yb^{3+}$  and  $Ba_2GdV_3O_{11}$ : $Tm^{3+}/Yb^{3+}$  sensors have higher  $S_r$  values, in a wide

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temperature sensing range (300–600 K). This makes them perfect for a wide range of luminescence thermometry uses.

Our findings underscore the potential of erbium doping to modulate temperature sensitivity peaks, and improve the temperature resolution, crucial for optimizing performance in tailored luminescence sensors. This research highlights the profound influence of erbium  $(Er^{3^+})$  doping on the temperature sensitivity characteristics of thermometric systems, paving the way for improved precision in temperature measurements. This adjustment not only sheds light on the underlying mechanisms of tunability, but also unlocks novel opportunities for customizing thermometric devices to meet specific application needs.

The repeatability (*R*) is another crucial metric for assessing the accuracy of the temperature sensing techniques employed. The thermometric parameters, specifically the LIR values, were measured multiple times as the sample was cycled from a low to a high temperature, as illustrated in Fig. S3(a–d) and S4(a–f) (ESI†). The repeatability (*R*) was calculated using the following

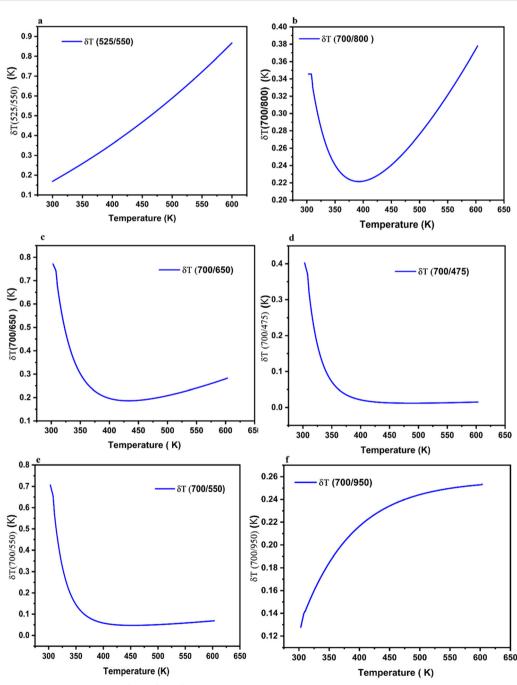


Fig. 14 Temperature uncertainty of BGVO: $Er^{3+}/Tm^{3+}/Yb^{3+}$ : (a)  $\delta T$  (525/550) nm, (b)  $\delta T$  (700/800) nm, (c)  $\delta T$  (700/475) nm, (d)  $\delta T$  (700/650), (e)  $\delta T$  (700/ 550) nm, and (f)  $\delta T$  (700/950) under 975 nm excitation.

Table 1 Sensitivities and temperature range for the temperature-sensing luminescence material

Materials	T-range (K)	$S_{ m r}~(\%~{ m K}^{-1})$ at 300 K	$S_{ m r}~(\%~{ m K}^{-1})$ at $T_{ m max}$	Ref.
NaBiF <sub>4</sub> :Tm <sup>3+</sup> /Yb <sup>3+</sup>	303-443	0.7	1.1	38
$Y_2Ti_2O_7:Tm^{3+}/Yb^{3+}$	300-400	0.4	0.87	39
$GdVO_4$ :Tm <sup>3+</sup> /Yb <sup>3+</sup>	297-673	0.4	1.54	40
$NaY_{2}F_{7}:Tm^{3+}/Yb^{3+}$	300-567	0.5	1.2	41
$Na_3GdV_2O_8Tm^{3+}/Yb^{3+}$	298-565	1.34	2.01	42
$Ba_2GdV_3O_{11}:Er^{3+}/Yb^{3+}$	298-573	1.1	1.1	43
$Gd_{2}O_{3}:Yb^{3+}/Er^{3+}$	300-900	0.83	0.20	44
$NaSrY(MoO_4)_3: Er^{3+}/Yb^{3+}$	298-335	1.19	1.19	45
YPVO <sub>4</sub> :2%Er <sup>3+</sup> /15%Yb <sup>3+</sup>	300-440	0.937	0.937	46
$Ba_2YV_3O_{11}:Er^{3+}/Yb^{3+}/Bi^{3+}$	300-440	1.43	1.43	37
$Er^{3+}/Yb^{3+}Na_3Gd(VO_4)_2$	300-600	0.83	0.83	47
$Ba_2GdV_3O_{11}:Tm^{3+}/Yb^{3+}$	303-603	5.397	5.397 (700/475)	This wor
$Ba_2GdV_3O_{11}:Er^{3+}/Tm^{3+}/Yb^{3+}$	303-603	0.2	2.35 (700/950)	This worl

formula:

$$R_{\rm p}(100\%) = \left(1 - \frac{\max|{\rm Mi}(T)_{\rm c} - M(T)_{\rm c}|}{{\rm FIR}_{\rm c}}\right) \times 100$$
 (8)

where  $\operatorname{Mi}(T)_c$  is the measured parameter (LIR or band centroid) in the *i*th cycle and  $M(T)_c$  is the mean value of the  $M(T)_c$  over 10 cycles. The LIR values determined change reversibly with temperature, and the LIR values for all LIRs were above 97% over the 300–600 K temperature range, confirming the good repeatability and reliability of the thermometric methods applied.

## Conclusion

In this study, the influence of erbium  $(Er^{3+})$  doping on the temperature sensitivity profile of upconversion luminescent nanoparticles based on Ba<sup>2</sup>GdV<sup>3</sup>O<sup>11</sup> co-doped with Tm<sup>3+</sup> and Yb<sup>3+</sup> was investigated. Both BGVO:Yb<sup>3+</sup>/Tm<sup>3+</sup> and BGVO:Er<sup>3+</sup>/ Tm<sup>3+</sup>/Yb<sup>3+</sup> nanoparticles exhibited strong blue and green upconversion luminescence, respectively, under 975 nm laser excitation. The fluorescence intensity ratio (FIR) method was employed to analyze the temperature-dependent luminescence spectra across a broad temperature range (300-600 K). Our analysis was focused on both thermally coupled energy levels (TCLs) and non-thermally coupled energy levels (NTCLs) of Er<sup>3+</sup> and Tm<sup>3+</sup> for their impact on temperature sensing performance. The key finding of this work is the ability to modulate the temperature of maximum sensitivity through Er doping. In  $Tm^{3+}/Yb^{3+}$  co-doped samples, the peak sensitivity (1.7% K<sup>-1</sup>) occurred around 350 K primarily due to TCLs. Interestingly, non-TCL sensitivities lacked a clear maximum within the 300-600 K range. In contrast, Er<sup>3+</sup>/Tm<sup>3+</sup>/Yb<sup>3+</sup> tri-doped samples displayed a distinct maximum in non-TCL sensitivity within this temperature range, offering superior temperature determination for specific applications. Notably, Er<sup>3+</sup> doping also improved the temperature resolution from 0.1 K to 0.02 K. Remarkably, the utilization of non-TCLs resulted in a significantly higher  $S_{\rm rmax}$  value of 2.35% K<sup>-1</sup> at 300 K. This achievement was made possible by employing the combination of upconversion (UC) emission bands at 700/950 nm, originating from the transitions of Tm<sup>3+</sup>/Yb<sup>3+</sup> in the tridoped BGVO Er<sup>3+</sup>/ Tm<sup>3+</sup>/Yb<sup>3+</sup> nano phosphor. Thus, erbium doping has emerged

as a powerful tool for fine-tuning temperature sensitivity in luminescent nanosensors. This precise control over sensitivity translates to enhanced temperature resolution, a critical factor for optimizing sensor performance. By revealing the underlying mechanism, our research unlocks the potential for designing customized thermometric devices with superior accuracy.

# Data availability

The data supporting this article have been included as part of the ESI.†

## Conflicts of interest

There are no conflicts to declare.

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