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# Sensing of ultra-ppm level NO<sub>2</sub> gas via synergistic effects of Cr doping and e-beam irradiation in WO<sub>3</sub> nanostructures†

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Herein, we report the effects of Cr doping and electron beam irradiation (EBI) on WO<sub>3</sub> thin films towards low-level detection of NO<sub>2</sub>. Structural, morphological, and defect correlations are presented to provide a comprehensive understanding of the film's properties. XRD and Raman analyses confirmed the formation of well-crystallized WO<sub>3</sub> films due to doping and irradiation. Photoluminescence analysis validated the presence of oxygen vacancy (V<sub>o</sub>) defects in the films which was further corroborated by XPS studies. The EBI Cr–WO<sub>3</sub> film exhibited a higher concentration of V<sub>o</sub> defects and smaller grain sizes, contributing to enhanced sensing performance. A maximum sensor response of 3.43 was achieved at an NO<sub>2</sub> concentration of 5 ppm. Theoretical limit of detection (LOD) calculations indicated ultra-low NO<sub>2</sub> detection capability at 0.4 ppm (400 ppb), highlighting the potential of the prepared film for applications in environmental monitoring, healthcare diagnostics, and industrial settings.

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## 1. Introduction

Recently, ambient air pollution has emerged as a severe global issue, adversely affecting human health in multiple ways, such as causing heart disease, stroke, and lung diseases like asthma and cancer, ultimately leading to death.<sup>1–3</sup> According to a 2019 survey conducted by the World Health Organization (WHO), a striking 99% of the global population resided in areas where air quality did not meet WHO guidelines.<sup>4</sup> The primary pollutants contributing to this situation include particulate matter (PM), nitrogen dioxide (NO<sub>2</sub>), sulfur dioxide (SO<sub>2</sub>), and carbon monoxide (CO). Among them, NO<sub>2</sub> is a significant air pollutant primarily released from vehicles, industrial operations, fossil fuel combustion and nuclear power plants.<sup>5</sup> The Occupational Safety and Health Administration (OSHA) recommends that the permissible exposure limit for NO<sub>2</sub> must be maintained below 5 ppm.<sup>6</sup> Exceeding this limit can result in severe cardiovascular and respiratory diseases, while also contributing to environmental issues such as acid rain and photochemical smog, posing serious risks to both human health and the ecosystem. Moreover, recent research highlights a correlation between

ambient NO<sub>2</sub> levels and increased disease severity in COVID-19 patients.<sup>7</sup>

To overcome these problems, the field of materials science has developed various metal oxide semiconductor (MOS)-based gas sensors using SnO<sub>2</sub>,<sup>8</sup> WO<sub>3</sub>,<sup>9</sup> ZnO,<sup>10</sup> TiO<sub>2</sub>,<sup>11</sup> CuO.<sup>12</sup> MOS sensors are favoured for their excellent sensitivity, chemical stability, and ease of fabrication. They typically operate on the chemiresistive principle, wherein the adsorption of gas molecules onto the surface alters the electrical resistance of the sensing layer through surface reactions and modulation of charge carriers.<sup>13,14</sup> A more detailed explanation of the sensing mechanism is provided in Section 3.6.2. In recent years, tungsten oxide (WO<sub>3</sub>)-based gas sensors have garnered considerable attention within the research community owing to their superior catalytic properties, structural stability, and the presence of oxygen vacancy defects.<sup>15–17</sup> To further enhance the sensing performance of WO<sub>3</sub>, various modification strategies such as doping, composite formation, heterojunction engineering, and irradiation have been employed.<sup>18–22</sup> For example, Hingangavkar *et al.*<sup>3</sup> synthesized WO<sub>3</sub>-rGO composite to detect NO<sub>2</sub> at a working temperature of 150 °C. The sensor exhibited a response of 2.52 towards 100 ppm NO<sub>2</sub>. Rossinyol *et al.*<sup>23</sup> prepared Cr-doped WO<sub>3</sub> thin films and tested them towards a low concentration of NO<sub>2</sub> at 230 °C. Sensors marked the slow recovery of 16 min towards 0.1 ppm NO<sub>2</sub>. Similarly, Mathankumar *et al.*<sup>24</sup> fabricated Y-doped WO<sub>3</sub> sensors for NO<sub>2</sub> detection. Increased oxygen vacancies were observed upon doping, leading to a substantial improvement in sensor response. On the other hand, Lee *et al.*<sup>25</sup> investigated the

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NO<sub>2</sub> sensing behaviour of WO<sub>3</sub> films decorated with carbon nanofibers and subsequently subjected to UV irradiation. They reported a detection limit of 1 ppm at room temperature, though the sensor demonstrated a prolonged recovery time of 7 minutes. Zhang *et al.*<sup>22</sup> utilized laser irradiation to enhance the NO<sub>2</sub> sensing performance of WO<sub>3</sub> films at 200 °C, achieving a 4.3-fold increase in sensor response. These findings suggest that doping, composite formation, and irradiation significantly alter the properties of WO<sub>3</sub>, resulting in enhanced sensor performance. Nevertheless, considering the hazardous effects of NO<sub>2</sub>, there remains a pressing need for WO<sub>3</sub>-based sensors with improved sensitivity and rapid response/recovery characteristics at optimal operating temperatures.

In this study, we present a novel approach to enhance the NO<sub>2</sub> sensing characteristics of WO<sub>3</sub> films through the combined effects of chromium doping and electron beam irradiation. The films were evaluated for their response to low concentrations (1–5 ppm) of NO<sub>2</sub> at an optimal operating temperature of 200 °C. Furthermore, we established structure–property correlations related to the sensing features.

## 2. Experimental section

### 2.1 Materials and methods

Ammonium metatungstate hydrate (AMT) [(NH<sub>4</sub>)<sub>6</sub>H<sub>2</sub>W<sub>12</sub>O<sub>40</sub>·xH<sub>2</sub>O, 99.99%] from Sigma-Aldrich, chromium(III) nitrate non-

ahydrate [Cr(NO<sub>3</sub>)<sub>3</sub>·9H<sub>2</sub>O, 99%] from Sigma-Aldrich, polyethylene glycol 400 (PEG 400) [H(OCH<sub>2</sub>CH<sub>2</sub>)<sub>n</sub>OH] from Molychem, nitric acid [HNO<sub>3</sub>, 69%] from Merck Life Science Pvt. Ltd were used without further processing.

To synthesize WO<sub>3</sub> films, AMT was dissolved in doubly distilled water to achieve 0.01 M concentration. To improve the adhesion properties, PEG 400 and a few drops of conc. HNO<sub>3</sub> were added maintaining the solution's pH at ~3. The solution was stirred for about 30 minutes in a magnetic stirrer to get a homogeneous mixture. For Cr-doped WO<sub>3</sub> (Cr-WO<sub>3</sub>), the same procedure was followed, incorporating a chromium precursor into the AMT solution. The films were then deposited onto a soda lime glass substrate by employing the spray pyrolysis technique. Spray pyrolysis is a low-cost, scalable, and versatile technique that enables uniform deposition of oxide thin films with controlled composition and doping. Its simplicity and adaptability make it highly suitable for fabricating gas sensing materials. The deposition temperature was set at 400 °C and air pressure was adjusted to 0.4 bar to get the uniform films with an approximate thickness of 600 nm.

Electron beam irradiation (EBI) was conducted using a LINAC accelerator, with an 8 MeV electron beam irradiated onto the Cr-WO<sub>3</sub> film. For clarity and ease of reference, the WO<sub>3</sub>, Cr-WO<sub>3</sub>, and electron beam irradiated Cr-WO<sub>3</sub> films are denoted as sample A, sample B, and sample C, respectively, throughout the manuscript. The detailed composition and treatment conditions are summarized in Table 1. Fig. 1 illustrates the film synthesis procedure employed in the present study.

### 2.2 Materials characterization details

X-ray diffraction (XRD) analysis [Rigaku SmartLab] was performed over a 2θ range of 20°–80° at 40 kV and 30 mA using Cu-Kα radiation (λ = 1.54 Å) to determine structural and phase information. Raman spectroscopy [Horiba Lab Ram HR] was

Table 1 Summary of sample names with composition and treatment details

Sample name	Composition	Post-treatment
Sample A	Pure WO <sub>3</sub>	None
Sample B	WO <sub>3</sub> + 3 wt% Cr	None
Sample C	WO <sub>3</sub> + 3 wt% Cr	Electron beam irradiation (8 MeV, 10 kGy)

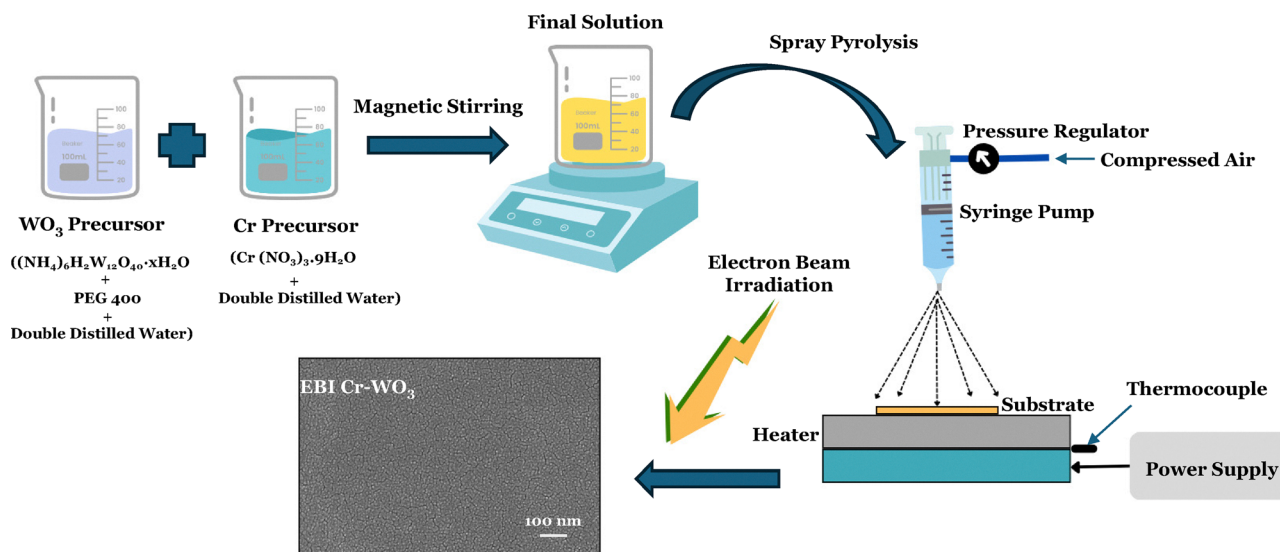


Fig. 1 Synthesis process of the sensing layer in this study.



implemented utilizing a 532 nm laser source to gain additional insights into the structure, defects, and the presence of various modes of vibrations in the samples. UV-Visible spectrophotometry [Shimadzu 1800] was employed to determine the bandgap of the samples (Section S2, ESI<sup>†</sup>). Photoluminescence (PL) measurements [JASCO FP-8500] were performed using a xenon lamp excitation source at a wavelength of 256 nm to analyze the defect states present in the samples. Morphological features were examined *via* field emission scanning electron microscopy (FESEM) [ZEISS ULTRA55]. X-ray photoelectron spectroscopy (XPS) [Kratos AXIS ULTRA] with Al-K $\alpha$  source ( $\lambda = 1486.6$  eV) was used to identify the oxidation state of the elements and perform composition studies on the samples.

### 2.3 Gas sensing measurements

Films with dimensions of 1 cm  $\times$  1 cm were used for the sensing tests. Silver paste electrodes (Alfa Aesar) were applied to the film surface, maintaining a 5 mm distance between them. The film was placed inside a chamber on a graphite block equipped with a lamp that served as a heater. A constant gas flow of 500 sccm was maintained within the chamber using programmable mass flow controllers (MFCs, Alicat Scientific). The sensor's operating temperature was optimized between room temperature and 250 °C. At lower temperatures, the response was negligible due to poor recovery behaviour. A significant and stable response, along with faster recovery, was observed at 200 °C, which was therefore selected for all subsequent measurements.

The sensor film was alternately exposed to dry air (79% N<sub>2</sub> + 21% O<sub>2</sub>) and the target gas (NO<sub>2</sub>) to record measurements. Dry air was purged until stabilization was achieved, after which NO<sub>2</sub> was introduced until a stable response was noted. Current-voltage (*I-V*) characteristics were recorded using a Keithley 2450 source meter to evaluate the sensor's response to the target gas. The diagrammatic representation of the gas sensor set-up used in the present study is shown in Fig. 2.

## 3. Results and discussion

### 3.1 Structural insights from XRD

X-ray diffraction (XRD) analysis was performed to investigate the structural modifications in the WO<sub>3</sub> film caused by chromium doping and electron beam irradiation. Fig. 3(a) presents the XRD patterns of samples A, B and C respectively. XRD pattern exhibits a prominent peak (200) at  $\sim 24.63^\circ$  and several minor peaks suggesting the polycrystalline nature of the films. No secondary phase related to WO<sub>3</sub>/Cr was detected, indicating that the deposited samples were free from impurities. All the diffraction peaks were matched with ICDD 43-1035, according to which the samples exhibit a monoclinic ( $\gamma$ ) structure.<sup>26–28</sup>

Various structural parameters derived from XRD patterns using the eqn (S1)–(S3) (ESI<sup>†</sup>) and are summarised in Table 2. The systematic increase in crystallite size (*D*) values from 12 nm to 16 nm is evident, accompanied by a concurrent reduction in dislocation density ( $\delta$ ) and strain ( $\varepsilon$ ) values. Due to the slight difference in ionic radii of the host and dopant ions ( $r_{\text{Cr}^{3+}} = 0.62$  Å;  $r_{\text{W}^{6+}} = 0.60$  Å), only minor reductions in strain and dislocation density were observed upon Cr incorporation into the WO<sub>3</sub> matrix.<sup>29,30</sup> However, the introduction of EBI into the Cr–WO<sub>3</sub> matrix induced further reduction in the strain and dislocation defects along with the left shift of the major peak (200) (Fig. 3(b)). The higher energy electron beam can cause point defects or vacancies, resulting in lattice expansion as substantiated by increased lattice plane distance (*d*) value (Table 2) and associated shift in the XRD peak.<sup>31</sup> The reduced lattice strain due to EBI can be attributed to a localized stress relaxation mechanism, likely driven by defect redistribution and structural reorganization.<sup>32,33</sup> During e-beam treatment, the transfer of momentum and energy induces the formation of vacancies and self-interstitials, which are subsequently partially removed *via* a diffusion-controlled mechanism. This defect reduction process leads to strain relaxation and improved crystallinity, as evidenced by the decrease in full-width at half-maximum (FWHM) values and the corresponding increase in crystallite size. Specifically, the FWHM values of the (200)



Fig. 2 Schematic diagram illustrating the configuration of the gas sensing system.



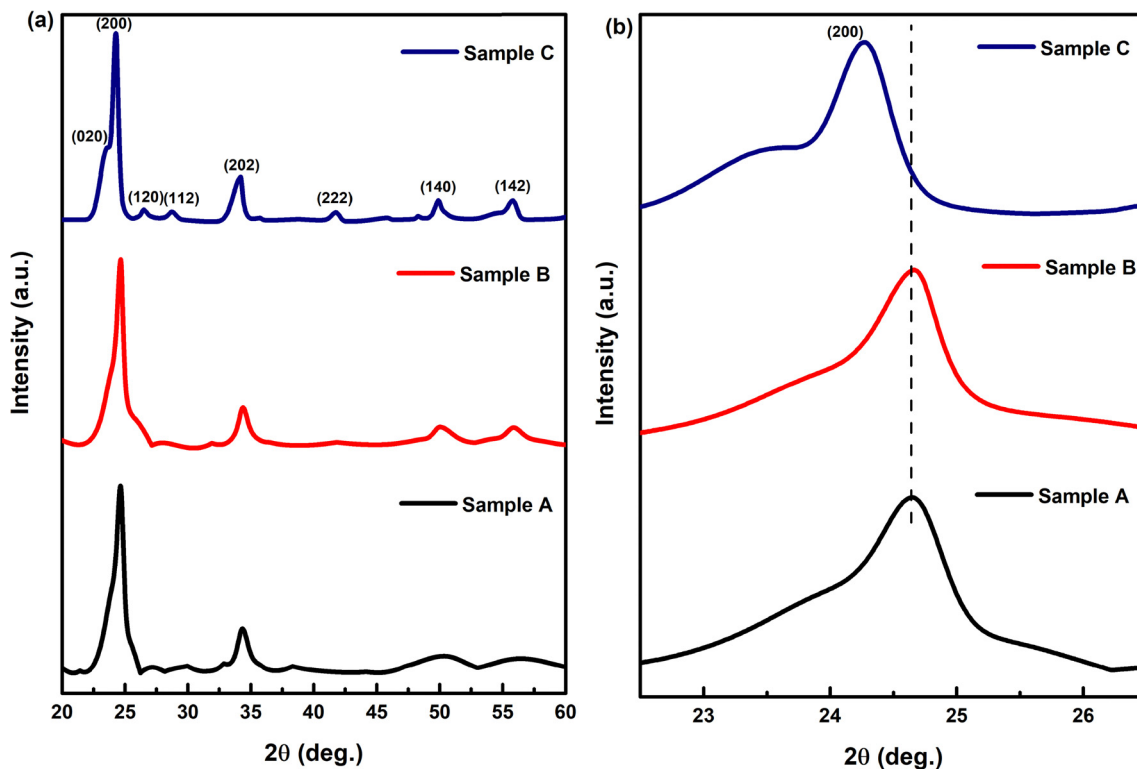


Fig. 3 (a) XRD patterns of samples A, B and C. (b) Enlarged view of peak (200).

Table 2 Structural parameters evaluated from XRD

Sample name	$d$ spacing (Å)	Crystallite size $D$ (nm)	Dislocation density $\delta$ ( $\times 10^{15}$ lines per $m^2$ )	Strain $\epsilon$ ( $\times 10^{-2}$ )
Sample A	3.61	12	7.16	1.37
Sample B	3.61	13	5.88	1.24
Sample C	3.66	16	3.96	1.04

peak were found to decrease progressively from 0.69 for sample A to 0.62 for sample B and 0.51 for sample C.

### 3.2 Structural insights from Raman

Structure confirmation and defects induced were further verified *via* Raman spectra as illustrated in Fig. 4(a). According to group theory, monoclinic  $WO_3$  with  $P2_1/n$  space group and  $C_{2h}^5$  point group has phonon zone centers given by,<sup>34</sup>

$$\Gamma_{\text{phonons}} = 24(A_g + A_u + B_g + B_u) \quad (1)$$

In the above equation, 48 zone centers were assigned to Raman active modes ( $A_g$  and  $B_g$ ) and remaining related to acoustic and infrared (IR) active modes. The Raman modes associated with  $A_g$  symmetry are the most commonly reported  $WO_3$  modes in the literature.<sup>35–37</sup> This predominance arises from the fact that  $B_g$  modes exhibit significantly lower intensity, making them difficult to detect in Raman spectra. In the present study, characteristic monoclinic  $WO_3$  peaks were observed at approximately 112, 134, 273, 327, 716, 805, and 961  $cm^{-1}$ , corresponding to various vibrational modes within the  $WO_3$  lattice. The peaks observed at approximately 716 and

805  $cm^{-1}$  correspond to the asymmetric and symmetric stretching vibrations ( $\nu_{as}$  and  $\nu_s$ ) of the O–W–O bonds, which represent the most prominent modes of monoclinic  $WO_3$ .<sup>35</sup> Peaks at 273 and 327  $cm^{-1}$  are linked to the bending vibrations ( $\delta$ ) of the O–W–O bonds, while those below 200  $cm^{-1}$  are associated with lattice vibrational modes.<sup>38,39</sup> Additionally, a peak at 961  $cm^{-1}$  was consistently detected in all the films, indicating a symmetric stretching vibration of the terminal W=O bonds. Table 3 depicts the Raman peak assignments and corresponding peak positions in samples A, B and C.

**3.2.1 Impact of doping and irradiation.** Upon Cr doping, no additional Raman peaks were observed, confirming the effective substitution of  $Cr^{3+}$  ions into the  $W^{6+}$  lattice. Moreover, doping had minimal impact on the peak profile. However, after electron beam irradiation of the Cr- $WO_3$  film, distinct sharp peaks emerged. The presence of sharp and well-defined Raman peaks typically indicates a high degree of crystallinity and structural order in the material. For  $WO_3$ , characteristic peaks at 716 and 805  $cm^{-1}$  are associated with a well-ordered monoclinic structure.<sup>40</sup> Fig. 4(b) presents the Voigt deconvolution of these peaks, revealing that their full-width at half-maximum (FWHM) values decrease with Cr incorporation and further narrow upon irradiation. The lowest recorded FWHM values of 48  $cm^{-1}$  and 27  $cm^{-1}$  for the irradiated sample suggest an enhancement in crystallinity. This improvement may be attributed to irradiation-induced stress relaxation, as observed in the XRD analysis, which contributes to the increased structural order of the films.

In addition, a slight shift of 1–6  $cm^{-1}$  in the peak positions was observed among the doped and irradiated samples. This





Fig. 4 (a) Raman spectra of samples A, B and C. (b) Enlarged view of deconvoluted O–W–O stretching vibrations.

Table 3 Raman peak assignments corresponding to different peak positions

Peak assignments	Peak position (cm <sup>-1</sup> )		
	Sample A	Sample B	Sample C
Lattice modes	111, 134	112, 134	111, 134
$\delta(\text{O-W-O})$	270, 321	273, 322	273, 327
$\nu_{\text{as}}(\text{O-W-O}), \nu_{\text{s}}(\text{O-W-O})$	722, 803	717, 804	716, 805
$\nu_{\text{s}}(\text{W=O})$	962	959	961

shift can be attributed to multiple factors, including internal stress arising from the interaction between  $\text{WO}_3$  atoms and the laser, atomic coupling effects, phonon confinement, and point defect-induced diffusion.<sup>41,42</sup> These mechanisms influence vibrational spectra, leading to subtle but measurable peak shifts. In the present case, the synergistic effect of doping and irradiation is suggested to induce point defects, such as oxygen vacancies or stress-induced distortions, which plausibly contribute to the observed peak shifts. This behaviour is consistent with previously observed trends in oxygen-deficient  $\text{WO}_3$ .<sup>43</sup>

### 3.3 Photoluminescence (PL) properties

To understand the defect contribution in the film structure, PL was conducted. Fig. 5(a)–(c) shows the deconvoluted PL spectra of the films, excited at a wavelength of 256 nm. Violet and blue emissions (2.64–3.13 eV) originate from the band–band transitions and vacant oxygen-related defects. In the present study, the bandgap determined from UV-visible spectroscopy is

approximately 3.10 eV (Fig. S1(b), ESI<sup>†</sup>), which is larger than the prominent peak observed at 2.69 eV in the PL spectrum. Therefore, a major peak at 2.69 eV is assigned to electron and hole recombination from the defect energy levels such as oxygen vacancies ( $V_{\text{o}}$ ), tungsten interstitials ( $W_{\text{inter}}$ ), self-trapped excitons, and other defects in the bandgap. Upadhyay *et al.*<sup>44</sup> and Bhargava *et al.*<sup>45</sup> reported similar findings from their study. They have ascribed the blue emission to the presence of oxygen vacancies and defects located in the conduction band minimum within the fundamental bandgap. Emission at 2.64 eV is probably due to defects related to  $\text{Cr}^{3+}$  and  $\text{WO}_3$  transitions.<sup>46</sup> Lower energy emissions at 2.21 eV (green) and 2.13 eV (orange) are mediated by oxygen vacancy defects ( $V_{\text{o}}^+/V_{\text{o}}^{++}$ ) and deformation centers created in the crystal lattice, respectively.<sup>47</sup> Fig. 6(a) and (b) shows a graphical representation of the peak widths and peak positions of all the samples corresponding to different emission centers.

**3.3.1 Impact of doping and irradiation.** Doping or irradiation did not introduce any distinct new defect levels in the crystal lattice. However, minor variations in defect density were observed among the doped and irradiated samples. The contribution of defects can also be assessed through peak width or FWHM of the emission peaks, as defect-related emissions are often associated with broader peaks or increased FWHM values. The blue emission at 2.69 eV was observed to narrow upon doping and irradiation, suggesting a possible reduction in  $V_{\text{o}}$ ,  $W_{\text{inter}}$ , or self-trapped excitons within the  $\text{WO}_3$  lattice. Additionally, emissions in the yellow-orange spectral range were suppressed for samples B and





Fig. 5 (a)–(c) PL spectra of samples A, B, and C.



Fig. 6 Plots of (a) peak width vs. emission centers and (b) peak position vs. emission centers of samples A, B, and C.

C, which may indicate reduced lattice deformations. In contrast, a slight broadening of the green emission was noted upon irradiation. FWHM calculations for the peak at 2.21 eV suggest an increase in  $(V_o^+/V_o^{++})$  defect concentration in the irradiated sample,

with values measured at 0.16 for samples A and B and 0.18 for sample C. These results are consistent with the presence of a higher concentration of  $(V_o^+/V_o^{++})$  defect centers in sample C compared to the others.



### 3.4 Morphological investigation

Fig. 7(a)–(c) present the FESEM micrographs of samples A, B and C, along with the corresponding grain distribution histogram. All films exhibited high homogeneity with no visible cracks on the surface. The undoped film possesses smaller grains, while upon doping, evenly distributed and well-defined grains became evident. Chromium induces the formation of larger nanocrystalline structures with irregular grain shapes, as illustrated in Fig. 7(b). This increase in grain size can be attributed to the role of  $\text{Cr}^{3+}$  in enhancing grain boundary mobility and reducing defect density during film growth, thereby promoting grain coalescence. A similar morphology is reported by Wang *et al.*<sup>48</sup> for  $\text{WO}_3\text{-TiO}_2$  composite films. EBI led to grain fragmentation, inducing fine structures, which is visible in the morphological features shown in Fig. 7(c). The average grain sizes of the films were estimated using ImageJ software, with values determined to be  $64 \pm 8$  nm,  $72 \pm 11$  nm, and  $51 \pm 8$  nm for samples A, B, and C, respectively. The irradiated sample (sample C) exhibited a smaller grain size than the others. Smaller grains provide a higher surface-to-volume ratio, facilitating faster electron exchange between grain boundaries.<sup>9</sup> This enhanced exchange increases the material's reactivity and gas sensitivity, as more adsorption sites are available for gas interaction.

### 3.5 Compositional analysis

The X-ray photoelectron spectroscopy (XPS) technique was utilized to analyze the elemental composition and oxidation states of the samples. Fig. 8(a)–(c) display the core level spectra of W 4f, while Fig. 8(d)–(f) show the O 1s core level spectra. Additionally, Fig. 8(g) and (h) illustrate the Cr 2p core level

spectra for samples B and C. These spectra were calibrated using the reference C 1s peak at 284.8 eV.

**3.5.1 Tungsten core level spectra.** In Fig. 8(a)–(c), the W 4f core level spectra are displayed as a doublet, with two distinct components:  $4f_{7/2}$  and  $4f_{5/2}$  with the tungsten oxidation state of +6. These components were located at binding energies of 36.16 eV and 38.30 eV, respectively for sample A. For sample B, the peak positions were noted at 35.88 eV and 38.03 eV, while sample C exhibited peaks at 35.94 eV and 38.08 eV.<sup>49,50</sup> Upon doping and e-beam treatment, a slight shift in the peaks was evident, indicating the formation of sub-stoichiometric  $\text{WO}_3$ .<sup>31</sup> Furthermore, the satellite peak at 42.35 eV is associated with the W  $5p_{3/2}$  component.

**3.5.2 Oxygen core level spectra.** The O 1s spectra were analyzed, revealing distinct peaks corresponding to different oxygen species (Fig. 8(d)–(f)). The primary peak ( $\text{O}_I$ ), observed for all three samples at around 530.65 eV, is attributed to  $\text{O}^{2-}$  ions bonded to  $\text{W}^{6+}$  ions. The secondary peak ( $\text{O}_{II}$ ) was noticed at approximately 531.43 eV, which corresponds to  $\text{O}^-$  ions or oxygen vacancies present within the lattice structure. Similarly, a third peak ( $\text{O}_{III}$ ) at 532.52 eV, likely originates from chemisorbed species such as  $\text{O}_2^-$ .<sup>51</sup> Table 4 summarizes the peak positions and percentage area ratios of  $\text{O}_I$ ,  $\text{O}_{II}$ , and  $\text{O}_{III}$ . A peak shift is observed in samples B and C, which is consistent with the W 4f core level spectra.

From a sensing perspective, oxygen vacancies serve as trap centres for incoming gas molecules and are a potential factor in enhancing gas sensing characteristics.<sup>50</sup> Similarly, chemisorbed species influence the sensing features as reported by Khader *et al.*<sup>52</sup> The present analysis indicates an increase in oxygen defects ( $\text{O}_{II} + \text{O}_{III}$ ) resulting from the synergistic effects

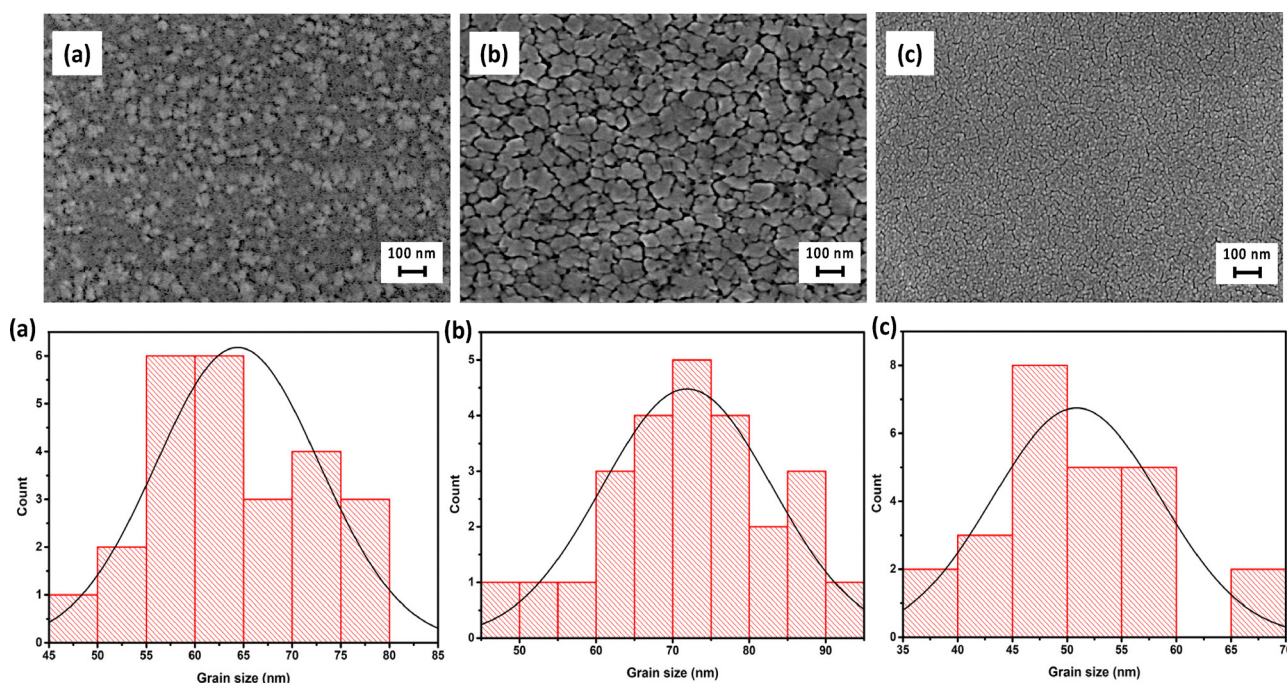


Fig. 7 (a)–(c) FESEM images and grain distribution histograms of samples A, B, and C, respectively.





Fig. 8 Core-level spectra of (a)–(c) W 4f for samples A, B, and C, respectively; (d)–(f) O 1s for samples A, B, and C, respectively; and (g) and (h) Cr 2p for samples B and C, respectively.

of doping and irradiation. This enhancement suggests a higher concentration of oxygen-related defects in the films, particularly in sample C, which implies more favourable conditions for sensing.

Table 4 O 1s core-level spectra analysis of samples A, B, and C

Sample name	Peak	Position (eV)	Area ratio (%)
Sample A	O <sub>I</sub>	530.65	53
	O <sub>II</sub>	531.43	7
	O <sub>III</sub>	532.52	40
Sample B	O <sub>I</sub>	530.23	50
	O <sub>II</sub>	530.77	28
	O <sub>III</sub>	531.73	22
Sample C	O <sub>I</sub>	530.67	41
	O <sub>II</sub>	531.11	30
	O <sub>III</sub>	532.31	29

**3.5.3 Chromium core level spectra.** The confirmation of chromium is evident from the presented spectra in Fig. 8(g) and (h), which show the Cr 2p<sub>3/2</sub> level at a binding energy of 576.50 eV. It represents the +3 oxidation state of chromium without any other impurity phase.<sup>53</sup> The minor peak adjacent to the main peak in the irradiated sample is often referred to as a satellite peak and is attributable to the Cr metallic state.<sup>54</sup>

### 3.6 Gas sensing characteristics

**3.6.1 NO<sub>2</sub> sensing measurements.** Gas sensing measurements for all films were conducted in response to NO<sub>2</sub> at an optimized operating temperature of 200 °C. Fig. 9(a)–(c) display the transient response curves of the films at NO<sub>2</sub> concentrations of 1, 3, and 5 ppm. The characteristics of the curves are influenced by the type of semiconducting material and the gas to which it is exposed. Given that WO<sub>3</sub> displays n-type conductivity, its resistance rises upon exposure to an oxidizing gas like NO<sub>2</sub>, and then decreases when it is returned to an air atmosphere.<sup>55</sup> All the films showed a noticeable response for different NO<sub>2</sub> concentrations. The key parameters that determine the sensor's efficiency are given in Table 5. The sensor response is a dimensionless quantity, defined as the relative ratio of resistances, and is evaluated using the equation,<sup>2,3</sup>

$$S = \left( \frac{R_g - R_a}{R_a} \right) \quad (2)$$

Here,  $R_a$  represents the resistance variation caused by the film's exposure to air, while  $R_g$  indicates the resistance change due to its exposure to NO<sub>2</sub>.

The sensor response was boosted with the addition of chromium and further enhanced by irradiation, consistent across all NO<sub>2</sub> concentrations tested. As shown in Fig. 9(d), the calibration curves reveal that sample B (Cr-doped WO<sub>3</sub>) outperforms the undoped film (sample A), and the irradiated Cr-doped film (sample C) shows the highest response. The rate of response and recovery exhibited oscillatory behaviour, with sample C achieving significantly faster kinetics. This improvement arises from a true synergy between doping and irradiation. Cr<sup>3+</sup> incorporation creates additional oxygen vacancies and modifies the electronic structure to increase adsorption sites and carrier density, while subsequent electron beam irradiation refines crystallinity and optimizes defect distribution. Collectively, these





Fig. 9 (a)–(c) Transient response curves of samples A, B, and C at an operating temperature of 200 °C. (d) Corresponding calibration curves.

Table 5 Sensor parameters evaluated from the response curves

Sample name	NO <sub>2</sub> conc. (ppm)	Response time (s)	Recovery time (s)	Sensor response
Sample A	1	189	32	0.64
	3	85	120	1.02
	5	50	90	1.06
Sample B	1	450	188	1.12
	3	210	128	2.17
	5	75	95	2.78
Sample C	1	184	28	1.07
	3	58	23	2.65
	5	15	30	3.43

coupled effects yield a superior sensor response of 3.43 and rapid response/recovery times of 15 s/30 s at 5 ppm NO<sub>2</sub>. This represents an overall 3.2-fold enhancement relative to the undoped film. The response and recovery characteristic curve for sample C is shown in Fig. 10(a). Theoretical limit of detection (LOD) was estimated for sample C using eqn (S4) (ESI<sup>†</sup>). It suggests that the possible lowest detection limit is 0.4 ppm (400 ppb)

demonstrating the capability of sensing ultra-ppm concentration of NO<sub>2</sub> gas. The linear regression plot used to calculate LOD is shown in Fig. 10(b).

The enhancement in the sensor response, as well as response and recovery features, can be attributed to the following three factors:

(i) High crystallinity of the samples as evidenced by XRD and Raman analyses. Peak width analysis indicated enhanced crystallinity due to the combined effects of doping and irradiation, with sample C exhibiting higher crystalline nature. Improved crystallinity enhances the sensing performance of oxide thin films by promoting better structural order and reducing defects, thereby facilitating efficient charge transport and gas–solid interactions. This observation is consistent with previous reports emphasizing the positive influence of crystallinity on the sensitivity of metal oxide sensors.<sup>56,57</sup> (ii) The presence of point defects such as oxygen vacancies was confirmed through PL and XPS studies. PL spectroscopy revealed green emission bands associated with these vacancies, indicating a higher concentration of charged oxygen vacancies in sample C. Furthermore, XPS analysis supported this observation by confirming the presence of oxygen





Fig. 10 (a) Response and recovery kinetics of sample C at 5 ppm  $\text{NO}_2$  and 200 °C. (b) Linear regression graph illustrating the sensor response as a function of  $\text{NO}_2$  concentration for sample C at 200 °C.

vacancies and chemisorbed oxygen species in the samples. Notably, the concentration of oxygen vacancies rose significantly from 7% in sample A to 28% in sample B and reached 30% in sample C. These oxygen vacancies and related species enhanced surface reactivity and provided additional active sites for gas adsorption. (iii) The smaller grain sizes of the samples ranging from 51 nm to 72 nm as seen in morphology studies. FESEM analysis showed that sample C had the smallest grains, with an average size of 51 nm. According to the  $D$ - $2L$  model (where  $D$  represents grain size and  $L$  denotes the thickness of the electron charge layer formed on the sensor surface), smaller grains typically favour the higher gas sensitivity.<sup>58</sup> Together, these results suggest that the observed structural and defect-related features played a key role in the enhanced sensor performance of sample C.

Selectivity, repeatability, and stability are critical factors that determine sensor performance in practical applications. In this study, the sensing responses of sample C were evaluated for nitrogen dioxide ( $\text{NO}_2$ ) and its common interfering gases including ammonia ( $\text{NH}_3$ ), carbon monoxide (CO), and methane ( $\text{CH}_4$ ),

which are typically associated with the combustion processes. The recorded sensor responses were 3.43 for  $\text{NO}_2$ , 1.9 for  $\text{NH}_3$ , 0.21 for CO, at 5 ppm concentration and 0.22 for 50 ppm  $\text{CH}_4$  as shown in Fig. 11(a). The results indicate that  $\text{NH}_3$  is the most significant interferent, causing a signal attenuation of  $\sim 45\%$ , while CO and  $\text{CH}_4$  are the least significant interferents, resulting in a signal attenuation of  $\sim 94\%$  at respective gas concentrations. In conclusion, sample C exhibited strong resistance to interference from  $\text{NH}_3$ , CO, and  $\text{CH}_4$ . Furthermore, repeatability test was conducted after one week, involving three cycles of  $\text{NO}_2$  purging at a concentration of 5 ppm. Fig. 11(b) illustrates the reproducible characteristics of sample C, which exhibited consistent responses across consecutive measurements. However, it is noteworthy that the sensor response decreased by approximately 5% over the week. Thus, the present study confirms the good reproducibility and initial stability of sample C.

**3.6.2  $\text{NO}_2$  sensing mechanism.** The sensing mechanism of n-type metal oxides, such as EBI Cr- $\text{WO}_3$  film, upon exposure to air and  $\text{NO}_2$  can be described as follows. When air is introduced



Fig. 11 (a) Selectivity histogram and (b) repeatability test performed after a week for sample C.



onto the film surface at an optimal working temperature of 200 °C, oxygen ions are drawn from the conduction band of the sensor film. Typically,  $O^-$  ions dominate within the temperature range of 150 °C to 300 °C.<sup>59</sup> This interaction leads to the formation of an electron charge layer ( $w$ ) on the film surface and a potential barrier ( $\Delta\phi$ ) between the film grains, thereby reducing the film's resistance. The gas sensing mechanism and related energy band diagram are illustrated in Fig. 12, and the equations describing the oxygen adsorption phenomena are presented in eqn (3)–(5).



When  $NO_2$ , an oxidizing gas, is introduced, it extracts electrons from the conduction band, leading to an increase in the width of the electron charge layer and the height of the potential barrier. This process results in a rise in the sensor film's resistance. The presence of a high concentration of oxygen defects ( $V_o \sim 30\%$ ), as estimated from XPS analysis, facilitates the accumulation of additional  $NO_2$  adsorbates and enhances chemical interactions with available  $O^-$  ions, further expanding the electron charge layer and potential barrier. Consequently, the resistance increases rapidly, improving the sensor's response.<sup>60</sup> The equations governing  $NO_2$  adsorption are provided in eqn (6)–(8).

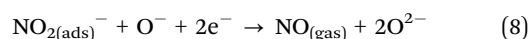
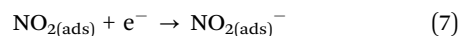
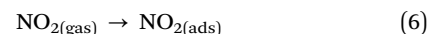


Table S1 (ESI†) provides the comparison data of reported  $NO_2$  sensors in the literature. While many  $WO_3$ -based sensors require high gas concentrations (e.g., 200 ppm for mesoporous  $WO_3$ ) to achieve high responses, the present study demonstrates that Cr doping and electron beam irradiation significantly enhance  $NO_2$  sensing at a much lower concentration (5 ppm). The EB-irradiated Cr- $WO_3$  sensor achieves a response of 3.43, outperforming materials like Ag-doped  $WO_3$  (1.5),  $WO_3$  nanoflowers (2.25), and Pd-rGO (1.05). Although mesoporous  $WO_3$  shows the highest response (99.78), it requires elevated temperatures (250–300 °C) and high  $NO_2$  levels (200 ppm), making the present work more suitable for low-concentration detection. However, the operating temperature (200 °C) still needs further reduction, which could potentially be achieved using catalytic additives (e.g., noble metals),  $WO_3$  composites, or alternative irradiation techniques.

## 4. Conclusion

In this study,  $WO_3$  films were successfully synthesized *via* spray pyrolysis, and Cr doping along with EBI was employed to enhance their  $NO_2$  sensing performance. XRD and Raman analyses confirmed the formation of well-crystallized  $WO_3$ ,

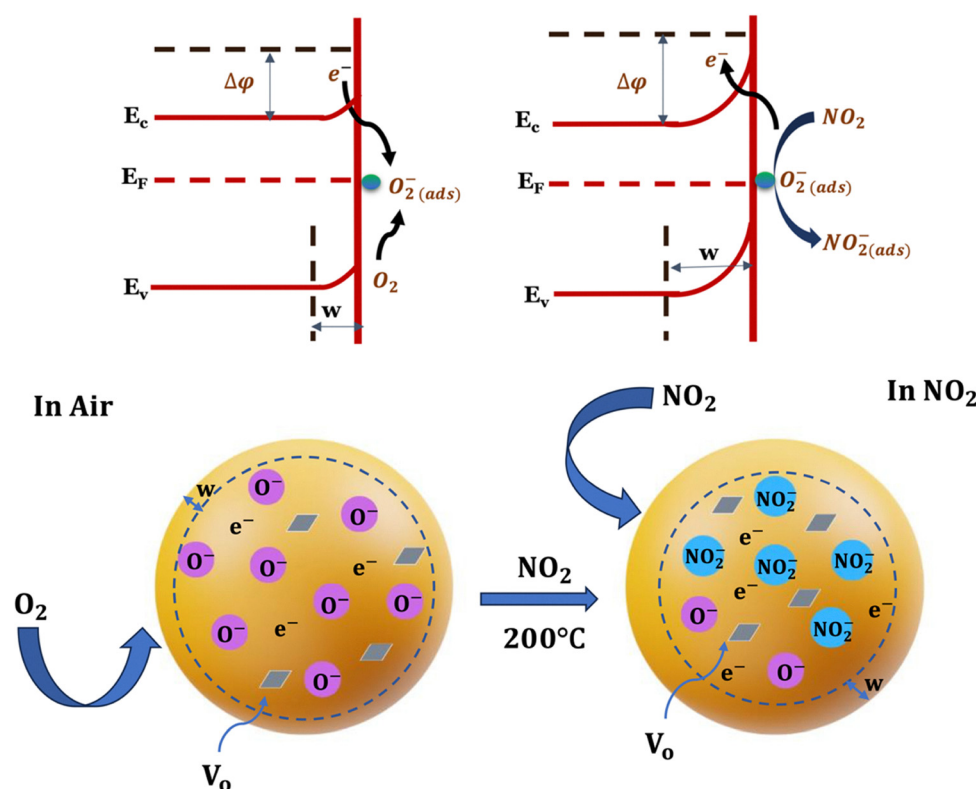


Fig. 12 Mechanism of gas sensing and band bending in EBI Cr- $WO_3$  film.



films. Preliminary investigations of oxygen defects using Raman and PL spectroscopy were further confirmed by XPS, which revealed a higher concentration of  $V_o$  defects in sample C. The sensor response of sample C exhibited a 3.2-fold increase compared to sample A at 5 ppm  $\text{NO}_2$ , attributed to the synergistic effects of Cr doping and EBI. Additionally, the sensor demonstrated rapid response and recovery times of 15 s and 30 s, respectively, at an optimal operating temperature of 200 °C. The enhanced sensing performance can be attributed to the presence of oxygen vacancy defects, reduced grain size, and improved crystallinity. While this work establishes a strong foundation, further studies exploring a wider range of Cr doping concentrations and EBI dosages could yield deeper insights into defect-driven sensing mechanisms, aiding in sensor optimization. Moreover, evaluating the sensor's response to a broader spectrum of target gases would help confirm its selectivity and expand its applicability across diverse gas sensing scenarios. Overall, this study presents a promising strategy for enhancing the sensing performance of  $\text{WO}_3$ -based gas sensors through the combined effects of doping and irradiation, advancing their suitability for environmental monitoring and industrial applications.

## Conflicts of interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Data availability

The data supporting this article have been included as a part of the ESI.†

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