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Anisotropy-dependent chirality transfer from cellulose nanocrystals to β -FeOOH nanowhiskers

Polysaccharides are closely associated with biogenic iron oxide in biomineralization. Using cellulose nanocrystals as both templates and chiral ligands, we demonstrate that this interaction not only guides mineralization but also induces chirality transfer at the molecular level. Oriented attachment of β -FeOOH nanoparticles on the surface of cellulose nanocrystals leads to highly aligned chiral nanowhiskers with an anisotropy-dependent g-factor, driven by enhanced dipole-dipole interactions. Incorporation of β -FeOOH also modifies the magnetic response of cellulose helical photonic films, highlighting strong coupling between the mineral phase and the chiral matrix.

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Anisotropy-dependent chirality transfer from cellulose nanocrystals to β -FeOOH nanowhiskers†

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Chiral iron oxides and hydroxides have garnered considerable interest owing to the unique combination of chirality and magnetism. However, improving their g -factor, which is critical for optimizing the chiral magneto-optical response, remains elusive. We demonstrated that the g -factor of β -FeOOH could be boosted by enhancing the anisotropy of nanostructures during a biomimetic mineralization process. Cellulose nanocrystals were used as both mineralization templates and chiral ligands, driving oriented attachment of β -FeOOH nanoparticles and inducing the formation of highly aligned chiral nanowhiskers. Circular dichroism spectra and time-dependent density-functional theory proved that chirality transfer was induced from cellulose nanocrystals to β -FeOOH through ligand–metal charge transfer. Interestingly, chirality transfer was significantly enhanced during the elongation of nanowhiskers. A nearly 34-fold increase in the g -factor was observed when the aspect ratio of nanowhiskers increased from 2.6 to 4.4, reaching a g -factor of 5.7×10^{-3} , superior to existing dispersions of chiral iron oxides and hydroxides. Semi-empirical quantum calculations revealed that such a remarkable improvement in the g -factor could be attributed to enhanced dipolar interactions. Cellulose nanocrystals exert vicinal actions on highly anisotropic β -FeOOH with a large dipole moment, increasing structural distortions in the coordination geometry. This mechanism aligns with the static coupling principle of one-electron theory, highlighting the strong interaction potential of supramolecular templates. Furthermore, paramagnetic β -FeOOH nanowhiskers alter the magnetic anisotropy of cellulose nanocrystals, leading to a reversed response of helical photonic films to magnetic fields, promising for real-time optical modulation.

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Introduction

Chiral inorganic nanomaterials have captured the attention of researchers due to their interesting optical properties and diverse applications.^{1–3} Imposing symmetry breaking of organic matter on inorganic nanomaterials is the most effective means for inducing chirality transfer.^{4–6} The fundamental mechanism is related to interfacial electronic coupling, which has allowed the emergence of chiroptical properties that are otherwise hard to achieve in traditional chiral molecules.^{4,6} Among various chiroptical inorganic nanoparticles (NPs), iron oxides stand out as an intriguing platform owing to their unique spin and magnetic characteristics.^{7–9} Despite the progress made in the synthesis of chiral iron oxides using amino acids as chiral agents,^{2,8,10} challenges persist in (1) understanding the relationship between structural parameters and the g -factor, (2) extending to supramolecular templates for the potential increase of asymmetry, and (3) exploring novel anisotropy-dependent chiromagnetic properties. In particular, while

different strategies have been proposed to improve the g -factor, such as the engineering of transition dipoles,⁶ exciton coupling,¹¹ and long-range order,³ none have been proven effective for iron oxides and hydroxides.

Polysaccharides are intricately linked with biogenic iron oxyhydroxides,¹² particularly in the form of aligned goethite¹³ and magnetite nanowhiskers (NWs).¹⁴ Primarily associated with hard biomineralized tissues, the potential for chirality transfer from polysaccharides to iron oxides remains elusive. This connection between polysaccharides and iron oxides offers a great opportunity to combine chirality with magnetism through a biomimetic concept for generating chiromagnetic effects.^{7,10} Combining two earth-abundant materials in a hybrid chiral system also mirrors the elegant strategies employed by living organisms to maximize the use of limited resources, highlighting the sustainable advantages of such applications. In contrast to small chiral molecules such as amino acids, structural polysaccharides may not only provide chiral interfaces for facilitating chirality transfer but also act as supramolecular templates, offering significant potential for strengthened interactions.^{15,16} This aspect, often overlooked in theories focusing on excited state coupling,¹⁷ could be better understood through one-electron theory based on the static coupling mechanism.¹⁸

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Supramolecular templates with multiscale architecture can induce helical organization to manipulate chirality,^{3,19} with cellulose nanocrystals (CNCs) serving as one such example. As a renewable resource, CNCs, assembled from β -1,4-linked anhydro-D-glucose chains with asymmetric carbon atoms, were proven to be twisted.²⁰ A chiral nematic phase can be produced through evaporation-induced self-assembly (EISA) of CNCs at room temperature.²¹ This hierarchical chiral template has been widely used for the synthesis of hybrid chiral photonic films,^{22–25} chiral inorganic films,^{26–29} and helical assembly of inorganic NPs.^{30–32} It is worth noting that in these previous efforts, chirality transfer from CNCs to inorganic NPs occurred at the micro- and nano-scales. The asymmetric packing of inorganic NPs is the main reason for the observation of chiroptical properties.

In this work, oriented attachment of β -FeOOH NPs onto CNCs was demonstrated to form highly aligned chiral NWs. Chirality transfer from CNCs to β -FeOOH occurs at the molecular level, and the induced optical activity of β -FeOOH was confirmed by circular dichroism (CD) spectra, in agreement with time-dependent density-functional theory (TD-DFT). An enhancement of the g -factor of β -FeOOH by up to 34-fold during its oriented growth was observed, which was elaborated by one-electron theory. Dipole moment (DM) calculations showed that CNCs introduce field perturbations through dipolar interactions with the crystal lattice of β -FeOOH, resulting in structural distortions of β -FeOOH, beneficial to chiroptical properties. Orderly arranged paramagnetic β -FeOOH NWs modify CNCs' magnetic anisotropy, leading to a unique magneto-optical response of helical photonic films.

Results and discussion

CNCs were directly employed to induce oriented mineralization of iron oxides under mild conditions. In a typical synthesis, certain amounts of CNC solutions were added dropwise to a diluted Fe^{3+} solution. After a 36-hour room-temperature reaction, high-contrast NWs were observed, exhibiting a strong orientation along the axis of CNCs (Fig. 1a). Low-magnification TEM images confirmed the alignment, displaying two substances (CNCs and β -FeOOH) with distinctly different contrasts simultaneously (Fig. S1†). Concurrently, small NPs were also present (Fig. 1b and S2a–d†). Their sizes closely match the diameters of NWs (Fig. S2e and f†), aligning well with the elementary nanofibrils in CNCs.²¹ A high-resolution transmission electron microscopy (HRTEM) image revealed interplanar spacing distances of approximately 0.330 nm and 0.255 nm, corresponding to the lattice planes of (30–1) and (21–1) for β -FeOOH, respectively (Fig. 1c, inset). Fast Fourier transform (FFT) analysis confirmed the growth direction of NWs aligned along the b -axis (Fig. 1c, inset).^{33,34} Further examination through high-angle annular dark field scanning transmission electron microscopy (HAADF-STEM) and energy dispersive X-ray spectroscopy (EDXS) mapping revealed the homogeneous dispersion of β -FeOOH on CNCs (Fig. 1d). In a solution lacking CNCs, oval NPs with diameters of approximately 15–30 nm were obtained (Fig. S3†), implying the critical role of CNCs in the confined growth of β -FeOOH. The template effect of CNCs is likely favoured by structural matching³⁵ because the Fe–Fe distance of 3.03 Å along the b -axis of β -FeOOH perfectly matches the O–O distance of 2.99 Å in cellulose (Fig. 1e). This geometric alignment is accompanied by the strong binding of O in the d -



Fig. 1 (a and b) TEM images of β -FeOOH adhering to the surface of cellulose nanocrystals (CNCs). The white and yellow arrows indicate the presence of β -FeOOH NWs and NPs, respectively. (c) An HRTEM image with the upper and lower insets showing the lattice fringes and FFT analysis of the central square part, respectively. The yellow line is perpendicular to (020) planes, parallel to the b -axis and the growth direction of NWs. (d) An HAADF-STEM image and EDXS mappings. (e) Structural matching between cellulose and β -FeOOH.



glucose units with Fe^{3+} (Fig. S4†). The binding of Fe^{3+} with sulfate half-ester groups, primarily located at C_6 , is also possible, although they constitute only a minor percentage.³⁶ The alignment of $\beta\text{-FeOOH}$ with CNCs suggests a level of structural control akin to that exerted by chitin over iron oxides during mineralization,^{37,38} although the precise pathways may differ.

Time-dependent TEM observations acquired insight into the mechanism of oriented mineralization. To visualize the growth of $\beta\text{-FeOOH}$ clearly, high-magnification TEM images were taken. Under such imaging conditions, CNCs could hardly be seen due to their extremely low contrast. However, because the size of $\beta\text{-FeOOH}$ NWs is much smaller than that of CNCs, it is practically feasible to differentiate $\beta\text{-FeOOH}$ NWs from CNCs. After 6 hours of reaction, the resultant material was labelled as CNC-Fe-6, with similar abbreviations used for other composites obtained at different reaction times. The surface of CNCs exhibited scattered small NPs in CNC-Fe-6 (Fig. 2a and S5a†), which aligned along the long axis and subsequently transitioned into NWs at a later stage (CNC-Fe-12) (Fig. 2b and S5b†). Additionally, a few NPs free from CNCs were also observed (Fig. S5c and S5c†). With prolonged reaction time, the coverage of NWs increased, but NPs persisted (Fig. 2c–e and S5d†), continuing to align linearly (Fig. S6†). The aspect ratios (ARs) of $\beta\text{-FeOOH}$ exhibited continuous growth before experiencing a slight drop at 48 h (Fig. S7†). X-ray diffraction (XRD) patterns confirmed that both NPs and NWs could be assigned to $\beta\text{-FeOOH}$ (JCPDS 42-1315) (Fig. 2f). These results indicated that oriented attachment of $\beta\text{-FeOOH}$ to CNCs constitutes a fundamental mechanism in the formation of aligned NWs. Fe^{3+} interacts with functional groups on CNCs, promoting the nucleation of small NPs, which is favoured by the electrostatic attraction between Fe^{3+} and CNCs, as confirmed by the sharp decline of electrokinetic potential after the addition of Fe^{3+} (Table S1†). Small NPs may also form directly in the solution but tend to assemble with those already anchored to CNCs,

creating linear chains that eventually grow into NWs (Fig. 2a–e, inset, and S8†).

Quantum mechanical calculations based on semi-empirical PM6 theory revealed that the DM of $\beta\text{-FeOOH}$ increases as ARs become larger (Fig. S9†). The dipole–dipole interaction potential (E_{dd})³⁹ between CNCs and $\beta\text{-FeOOH}$ is strengthened according to the variations of DMs (μ') and angles (θ') formed between the vectors of DMs and the line connecting the centre (Fig. 2g and S10†), encouraging oriented attachment of $\beta\text{-FeOOH}$ on CNCs. However, dipolar interactions between CNCs and $\beta\text{-FeOOH}$ could limit particle mobility, which explains the presence of small NPs even at the end of reactions, differing from oriented attachment in the solution.^{40–42}

The UV-vis spectra of CNC dispersions exhibited a distinct UV absorption profile⁴³ (Fig. 3a). The appearance of new absorption peaks during the reaction implied the formation of $\beta\text{-FeOOH}$ (Fig. 3a). Specifically, the peaks centred at 216, 268, and 350 nm are attributed to $\text{O } 2p \rightarrow \text{Fe } 3d$ ligand–metal charge transfer (LMCT) transitions.⁴⁴ The peak at 450 nm results from the double excitation of the Fe–Fe pair transition, while the peak at 538 nm is ascribed to a single ion transition.^{45,46} Additionally, the broad peak beyond 600 nm may be related to other forbidden d–d transitions. The redshift in the absorption edge, which occurred as NWs grew, can be attributed to changes in the energy levels of $\beta\text{-FeOOH}$ with varying ARs.⁴⁷

In the CD spectra, CNC dispersions exhibited a conspicuous positive signal (Fig. 3b). However, this signal significantly diminished after a 6-hour reaction (CNC-Fe-6) (Fig. 3b). With the reaction's extension, negative peaks began to emerge, gradually increasing in intensity and shifting bathochromically from 12 to 48 hours (Fig. 3b). Compared with time-dependent TEM observations (Fig. 2a–e), the emergence of NWs appeared to be critical for an apparent negative CD signal. In contrast, the dispersion of $\beta\text{-FeOOH}$ NPs measuring 15–30 nm exhibited a bare CD signal, albeit with similar absorption spectra (Fig. S11†).



Fig. 2 (a–e) TEM images of mineralized CNCs obtained at different reaction times (6–48 h) and (f) the corresponding XRD patterns. The inset shows the progression of oriented attachment. (g) The calculated dipole–dipole interaction potential (E_{dd}) between CNCs and $\beta\text{-FeOOH}$ with different aspect ratios (ARs). The inset displays the geometry used for the calculation of E_{dd} when AR = 3.





Fig. 3 (a) UV-vis and (b) CD spectra of dispersions of CNCs and CNC-Fe-*x* (*x* = 6–48). (c) Atomistic models used for TD-DFT calculations. (d) Simulated CD and absorption spectra of Fe₃O₆-3Glu.

The above results indicated the chirality transfer from CNCs to β-FeOOH. To elucidate the mechanism of chirality transfer, we conducted calculations using the TD-DFT method.⁴⁸ The geometry of a D-glucopyranose unit extracted from the structure of CNCs has a considerable effect on the simulated CD spectra, primarily due to the different orientations of –OH groups (Fig. S12†). However, the calculation of two D-glucopyranose units linked by β-(1–4) glycosidic bonds extracted from different positions showed similar CD spectra, indicating that two glucose monomers may represent cellulose quite well (Fig. S13†). Thus, to simplify the search for excited states, we substituted the complex cellulose molecules with two glucose monomers linked by β-(1–4) glycosidic bonds, denoted as Glu (Fig. 3c). The calculated absorption peak of Glu is in the UV-vis range and the apparent Cotton effect can be observed because of the σ to σ* and n to σ* transitions according to the hole–electron distribution and charge differential density plot (Fig. S14†).

We then established Fe₃O₆, an elementary Fe–O structure in the akaganéite unit cell, bonded with three Glu units through C₃–O, denoted as Fe₃O₆-3Glu (Fig. 3c). The formation of Fe–O bonds between CNCs and β-FeOOH was confirmed by both FT-IR and XPS spectra. The broad peaks between 3500 and 3350 cm^{−1} are attributed to the stretching vibrations of –OH,⁴⁹ which exhibited an obvious redshift in CNC-Fe-36 compared with CNCs and β-FeOOH (Fig. S15†), implying hydrogen bonding interactions between CNCs and β-FeOOH. The FT-IR bands at 703 cm^{−1} for β-FeOOH corresponding to Fe–O bonds⁵⁰ shifted to 696 cm^{−1} in CNC-Fe-36 (Fig. S15†), suggesting the presence of Fe–O–C bonds in the composites.⁵¹ XPS spectra of CNC-Fe-36 showed the characteristic peaks from C, O and Fe elements (Fig. S16a†). High-resolution Fe 2p XPS spectra (Fig. S16b†) showed Fe 2p_{1/2} and 2p_{3/2} peaks and two satellite peaks at 725.5, 711.9, 732.6 and 720.2 eV, respectively, corresponding to Fe³⁺ in β-FeOOH.^{52–54} In the XPS spectra of C 1s (Fig. S16c†), an obvious shift of the peak related to C–O bonds toward low binding energies in CNC-Fe-36 compared to CNCs could be attributed to

the lower electronegativity of Fe atoms. The O 1s spectra of CNC-Fe-36 (Fig. S16d†) included a new Fe–O–C peak at 531.7 eV,⁵⁵ confirming the existence of Fe–O–C bonds. The peak related to Fe–O–H presented an obvious shift toward high binding energies, a result also observed for β-FeOOH/graphene oxide composites.⁵⁵ A typical lattice oxygen (Fe–O) peak^{52,56} at 530.3 eV was observed for CNC-Fe-36 compared to 531.3 eV for β-FeOOH.

The simulated absorption spectra of Fe₃O₆-3Glu differed from those of Glu (Fig. 3d and S13c†), extending into the longer wavelength range, closely correlated with the CD spectra (Fig. 3d). Fe₃O₆-3Glu generated chiroptical signals at 253, 349, and 439 nm due to LMCT, and the peaks at 600 nm and 800 nm are the result of d–d transitions according to hole–electron distributions^{57,58} (Fig. S17 and S18†). The HOMOs of Fe₃O₆-3Glu are situated not only at O atoms of Glu but also at lattice O atoms. LUMOs are mainly located at the d and p orbitals of the Fe₃O₆ centre, with a small portion distributed over O atoms in the chiral ligands (Fig. S19†). No CD signals could be observed for Fe₃O₆ due to its symmetrical structure (Fig. S20†). Despite the use of the simplified model, these results confirmed that glucose monomers could impart chirality to β-FeOOH through Fe–O–C bonds (Fig. 3d). Although there are noticeable similarities in the shapes of the experimental and calculated CD spectra, deviations are still apparent, likely stemming from the use of small molecular structures.⁵

According to the Rosenfeld equation, $R_{\text{oa}} = \text{Im}(|\mu| \cdot |m| \cos \theta)$, optical activity (R_{oa}) depends on the transition electric dipole moment (μ), the transition magnetic dipole moment (m), and the angle between them (θ). When θ exceeds 90°, R_{oa} becomes negative, while it remains positive when θ is less than 90°. TD-DFT calculations indicated that interfacing CNCs with β-FeOOH led to the formation of new excited states (Fig. S21, Tables S2 and S3†), which are responsible for the shift from the positive peak observed in CNCs to the negative peak seen in β-FeOOH NWs upon chirality transfer (Fig. 3b).

We did not observe any evidence of helical packing of β-FeOOH NWs with long-range order on the surface of CNCs. This can be explained by the fact that the direct growth of β-FeOOH NWs on the CNC surface occurs through random nucleation sites, which may disrupt the helical order. This contrasts with the self-assembly of NPs on a supramolecular template, where a more organized structure is typically achieved.^{3,30} Additionally, the presence of small NPs alongside the NWs further complicates the establishment of long-range order. FFT analysis of the top, middle, and bottom regions of a β-FeOOH NW revealed that the *b*-axis direction remained consistent at different positions (Fig. S22†), suggesting that no twisting occurred.

We employed the *g*-factor to assess the optical activity of CNC/β-FeOOH dispersions obtained at different reaction times to eliminate the influence of concentration. Our primary focus was on the negative peak occurring at approximately 350 nm, as it represents the most intense transition associated with chirality transfer. The *g*-factor exhibited an increase from 12 to 36 hours, followed by a negligible change at the 48-hour mark (Fig. 4a). This change in the *g*-factor closely paralleled the variations in the ARs of β-FeOOH (Fig. 4a). Notably, as ARs increased from 2.6 to 4.4, we observed a nearly 34-fold increase





Fig. 4 (a) The change of the *g*-factor and ARs during the reaction. (b) Schematic illustration of the division of two groups in one-electron theory. (c) Space-filling models of iron oxide clusters with different ARs connected with β -1,4-linked anhydro-D-glucose from three monomers through Fe–O–C₃. (d) The effect of ARs on Δd and σ^2 .

in the *g*-factor (from 1.7×10^{-4} to 5.7×10^{-3}) (Fig. 4a). The absolute value of 5.7×10^{-3} is one order of magnitude higher than that of CdSe nanorods (5×10^{-4}),⁵⁹ representing a fairly high *g*-factor achieved for the dispersion of chiral iron oxides, which is typically in the range from 1×10^{-4} to 2×10^{-3} . For example, chiral superparamagnetic Fe₃O₄ supraparticles with a *g*-factor of $\sim 2 \times 10^{-3}$ were fabricated by Li *et al.*² Polyacetylene macromolecular helicity endowed Fe₃O₄ nanoparticles with a dissymmetric factor of $\sim 1.3 \times 10^{-3}$.⁶⁰ Magnetic nanoparticles decorated with D-cysteine molecules were demonstrated to have a *g*-factor of 1.69×10^{-3} by Visheratina *et al.*⁶¹ In addition, regioselective magnetization endowed one-dimensional Zn_x-Cd_{1-x}S-Ag₂S/Au@Fe₃O₄ hybrid nanorods with a *g*-factor of 0.15×10^{-3} .⁶²

The enlargement of the *g*-factor of CdSe nanorods was explained by a non-degenerate coupled-oscillator model⁵⁹ for surface-related CD and linear polarization^{63,64} for excitonic CD. Since anisotropic structures often possess large DMs for strong interactions with chiral molecules, one-electron theory⁶⁵ based on the static coupling mechanism (ESI, Section S2†) provides a comprehensive explanation for the correlation between the *g*-factor and ARs. In this theory, the perturbation of electronic states of achiral chromophores by vicinal actions results in chiroptical activity. The rotatory power and *g*-factor, therefore, depend on vicinal actions, which are proportional to the amount of atomic displacement distorting the structural framework.

Applying this theory to our case, we regarded β -FeOOH as the chromophoric group and CNCs as the vicinal groups (Fig. 4b). To estimate the effect of vicinal actions, the distortion of coordination geometry was characterized by changes in the bond length (Δd) and bond angle variance (σ^2).⁶⁶ An increase in Δd and σ^2 as calculated by the PM6 method was observed when the ARs of β -FeOOH increased (Fig. 4c and d) or when increasing the number of glucose monomers while maintaining the same binding site with Fe₃O₆ (Fig. S23†). Because the increase of ARs

of iron oxide clusters and the number of glucose monomers accompanied the increase of DMs (Fig. S9 and S24†), strengthened dipolar interactions (Fig. 2g) play a significant role in enhancing vicinal actions. This effect provides a strong short-range field perturbation for inducing structural distortions, ultimately contributing to a larger *g*-factor.^{67,68} It is important to mention that oriented attachment may also result in an augmented density of crystal defects^{69,70} (Fig. S25†), further distorting the coordination geometry.

The chiral β -FeOOH could co-assemble with CNCs to produce iridescent photonic films. A consistent directional arrangement was observed on the film surface of CNCs (Fig. S26a†), which diminished for CNC-Fe-6 (Fig. S26b†), probably due to the disturbance of interfibrillar interactions by Fe³⁺. The arrangement gradually improved from CNC-Fe-12 to CNC-Fe-36, but the trend reversed for CNC-Fe-48 (Fig. S26c–f†). A longer reaction likely transformed Fe³⁺ into thin β -FeOOH NWs, which could conform to the assembly of CNCs. However, when β -FeOOH NWs became too large, the conformality decreased, affecting the assembly again (Fig. S26f†). The observed optimal loading of inorganic NPs is common in CNC-based chiral photonic films.^{15,16} Polarized optical microscopy (POM) images (Fig. S26,† insets) showed obvious birefringence, consistent with the anisotropic structure. The colour of the images gradually changed from green to yellow and then to orange, indicating a change in the packing of CNCs.⁷¹ The cross-sectional SEM images revealed a lamellar architecture (Fig. S27†), a characteristic feature of the cholesteric arrangement.⁷²

All these films exhibited positive CD peaks due to the preferential reflection of left-handed circularly polarized light (CPL) by a left-handed helical structure⁷³ (Fig. S28†). Slight changes in CD spectra during sample rotation and “front-back” measurements confirmed that linear dichroism (LD) and linear birefringence (LB) as potential artifacts could be ruled out (Fig. S29†).^{74,75} As more β -FeOOH was incorporated, CD peaks redshifted about 150 nm from 6 to 48 hours (Fig. S28†) due to the higher average refractive indices⁷⁶ and the larger helical pitch.²³ The redshift was similarly observed when increasing the film thickness (Fig. S30†), which, however, can be attributed to the varied assembly kinetics.^{77,78} In line with the optimal helical organization (Fig. S26†), the highest CD intensity was observed for CNC-Fe-36 (Fig. S28†). In diffusion reflectance circular dichroism (DRCD) spectra, in addition to the peaks associated with the photonic bandgap, the characteristic peaks from chiral β -FeOOH NWs were also evident due to the selective absorption of diffused CPL (Fig. S31†), which is consistent with the CD spectra from the solution (Fig. 3b).

To control the optical properties of CNC-based materials, an external magnetic field represents an effective way. CNCs with anisotropic diamagnetism can be regulated directly,^{79–81} while the addition of magnetic NPs allows more facile structural control.^{23,73,82} These efforts focused on the interference with the dynamic self-assembly of CNCs by the external magnetic field. Little attention was paid to the modulation of the optical properties of CNC-based photonic films after EISA, which holds great promise for the real-time control of light-matter interactions.^{83–86} We studied the optical modulation of photonic



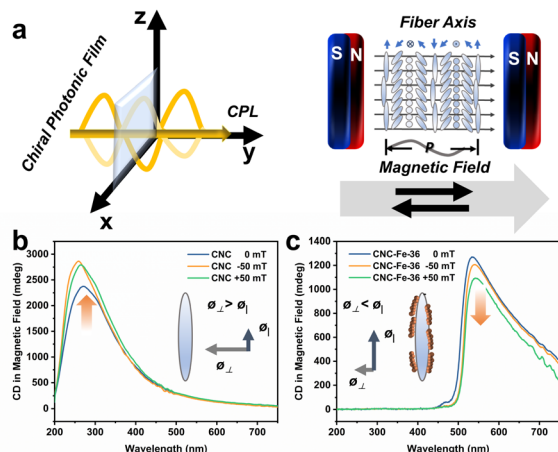


Fig. 5 (a) Schematic of the experimental set-up of CD spectral measurement under a magnetic field parallel or antiparallel to the direction of incident light and the helical axis of chiral photonic films. The effect of the magnetic field on the CD spectra of (b) CNC and (c) CNC-Fe-36 films. Insets show the magnetic anisotropy of CNCs and β -FeOOH modified CNCs, respectively.

films through a magnetic field. Different from previous work on the magnetic effect on chiral NPs^{6,87} or molecules,^{88,89} here, a helical medium dominated the response. The real-time modulation of CD spectra with a weak magnetic field (50 mT) aligned parallel or anti-parallel to the incident light and the helical axis was performed (Fig. 5a). We observed that intensity increases for CNC films (Fig. 5b) and decreases for CNC/ β -FeOOH films (Fig. 5c and S32[†]). Cellulose exhibits negative diamagnetic anisotropy,⁸¹ meaning $\theta_{\parallel} - \theta_{\perp} < 0$, where $\theta_{\parallel}$ and $\theta_{\perp}$ represent the diamagnetic susceptibilities in the directions parallel and perpendicular to the fiber axis, respectively (Fig. 5b, inset). The magnetic field enhances selective reflection by aligning the magnetic dipole at different domains with the helical axis, consistent with the emergence of well-defined absorption bands for CNC films (Fig. S33a[†]). Conversely, the aligned paramagnetic β -FeOOH NWs (Fig. S34[†]) may alter magnetic anisotropy⁹⁰ (Fig. 5c, inset). The magnetic field then tends to align the magnetic dipole perpendicular to the incident light, reducing the absorption of CNC-Fe-*x* films due to misalignment with the helical axis (Fig. S33b–f[†]).

Conclusions

In summary, oriented attachment of β -FeOOH NPs on CNCs led to the formation of highly aligned chiral NWs. Their *g*-factor increased substantially along the elongation process, and the observed correlation between the *g*-factor and ARs could be attributed to enhanced dipolar interactions, responsible for the induced structural distortions. The static coupling mechanism based on one-electron theory can be readily extended to comprehend chirality transfer in various hybrid systems. The observed anisotropy-dependent chiro-magnetic properties hold great promise for technological advances in field-controlled optical devices.

Experimental section

Materials

Iron trichloride hexahydrate and potassium hydroxide were purchased from Sinopharm Chemical Reagent Co., Ltd. 6 wt% CNC solution was purchased from CelluForce. Deionized (DI) water (Millipore, Milli-Q grade) was used in all experiments. All chemical solvents and reagents were used without further purification.

Synthesis of CNC-Fe-*x* dispersions and films

We prepared CNC-Fe-*x* dispersions (where *x* = 6, 12, 24, 36, and 48, representing the reaction time) by adding 1 mL of a 2 wt% CNC solution and 250 μ L of a 10 mg per mL KOH solution to 50 mL of a 20 mM FeCl₃ solution. The mixture was continuously stirred for various durations (6, 12, 24, 36, and 48 hours) at room temperature. The solution was centrifuged three times at 8000 rpm for 10 minutes. After centrifugation, the resulting material was redispersed in 2.5 mL of DI water and placed in Petri dishes with a diameter of 30 mm. Different films were formed through evaporation-induced self-assembly (EISA) at room temperature. The synthesis of β -FeOOH NWs followed a similar process to CNC-Fe-*x*, without the addition of CNCs.

Theoretical calculation

Binding energies. Cellulose was represented by six glucose monomers linked by β -(1–4) glycosidic bonds. All computational procedures were carried out within Gaussian 16, utilizing the m062x/6-31g(d) functional/basis set. To account for the solvation effects of water, the structure was optimized with a continuum solvation model incorporating the density implicit solvent (SMD). Binding energy, ΔG , was calculated as follows:

$$\Delta G = G_{ab} - G_a - G_b$$

$$G = E + C$$

where G_{ab} is the Gibbs free energy of the system after binding and G_a and G_b are the Gibbs free energy of the two components (*a* and *b*) before binding. *E* is the electron energy of the substance, and *C* is the correction to Gibbs free energy.

TD-DFT. To facilitate the exploration of excited states, two glucose monomers linked by β -(1–4) glycosidic bonds (Glu), extracted from the structure of CNCs (cif: 810597), were used as surrogates. An elementary Fe–O structure from the akaganéite unit cell, denoted as Fe₃O₆, was linked with 3 Glu through C₃–O–Fe, forming Fe₃O₆–3Glu. The geometries of these clusters were optimized using density functional theory (DFT) at the b3lyp/6-31g(d) level, accounting for the solvent effects of water within Gaussian 16. For the calculation of UV-vis and CD spectra, the CAM-B3LYP method and a 6-31G(d) basis set were employed to compute 60 excited states through time-dependent DFT (TD-DFT). Simultaneously, HOMOs, LUMOs, and other MOs were derived from the calculations. A comprehensive examination of electron excitation characteristics was conducted, involving hole-electron analysis, facilitated by Multiwfn.



Dipole moments (DMs). For modelling cellulose, a three antiparallel chain model, each containing three β -1,4-linked D-glucose units, was constructed. Iron oxide clusters with 16, 24, and 32 Fe atoms were created, corresponding to ARs of 1, 2, and 3, respectively. The geometries were first optimized by using the Merck molecular force field (MMFF) to acquire a relaxed configuration. Models optimized using the MMFF were then further optimized using PM6 according to semi-empirical theory for the calculations of DMs. The DMs and the electrostatic potential of Glu_{1-3} , where the subscript represents the number of glucose monomers, were calculated within Gaussian 16 after optimization at the level of 6-31+g(d,p).

Dipole-dipole interactions. The interaction energies between two permanent dipoles, denoted as μ (CNCs) and μ' (β -FeOOH with ARs of 1, 2, or 3), can be calculated using the following equation:

$$E_{\text{dd}} = \frac{\mu \times \mu'}{r^3} (2\cos\theta\cos\theta' + \sin\theta\sin\theta'\cos\zeta)$$

In this equation, E_{dd} represents the pair of permanent dipole-dipole interactions. The variable “ r ” denotes the distance between the centers of two dipoles, while θ and θ' represent the angles formed by μ and μ' with the central line. The variable “ ζ ” indicates the angle between the planes formed by the axes of dipoles μ and μ' with the line of centers. In this specific scenario, $\zeta = 0$ since all the lines are in the same plane. The positions of central axes and the values of r , $\cos(\theta)$, $\cos(\theta')$, $\sin(\theta)$, and $\sin(\theta')$ can be determined by assuming that the dipoles are situated in a 3D coordinate system.

Calculations of structural distortions. Iron oxide clusters with different ARs were connected with β -1,4-linked anhydro-D-glucose from three monomers through Fe–O–C₃ bonds. There are 1, 2 and 3 binding sites for iron oxide clusters with ARs of 1, 2 and 3, respectively. These models were optimized at the PM6 level within Gaussian 16. The bond length change (Δd) and bond angle variance (σ^2) were calculated according to

$$\Delta d = \sum_{i=1}^6 (d_i - d_0)^2 / (6d^2)$$

and

$$\sigma^2 = \sum_{i=1}^{12} (\theta_i - 90)^2 / 11,$$

where d_i is the specific Fe–O bond length, d_0 is the average Fe–O bond length and θ_i is the O–Fe–O bond angle in Fe–O octahedrons directly connected with glucose monomers. 6 is the number of bonds and 11 is the number of angles minus 1 in each octahedron. Average Δd and σ^2 were calculated for models containing iron oxide clusters with ARs of 2 and 3. For Fe_3O_6 connected with Glu_{1-3} , structure optimization was performed within Gaussian 16 at the level of 6-31+g(d,p). Δd and σ^2 were calculated according to

$$\Delta d = \sum_{i=1}^{11} (d_i - d_0)^2 / (11d^2)$$

and

$$\sigma^2 = \sum_{i=1}^{15} (\theta_i - 109.467)^2 / 14,$$

where d_i is the specific Fe–O bond length, d_0 is the average Fe–O bond length and θ_i is the O–Fe–O bond angle in Fe–O tetrahedrons. 11 is the number of bonds and 14 is the number of angles minus 1 in Fe_3O_6 .

Characterization. X-ray diffraction (XRD) patterns were recorded using a Rigaku D/max 2550 instrument with Cu K α radiation ($\lambda = 1.5418 \text{ \AA}$) in the range of 10–70° (2θ) at a scan rate of 2° min^{−1}. Scanning electron microscopy (SEM) was conducted using a JSM-7800F instrument at an accelerating voltage of 3 kV and a working distance of 9.8–10 mm. Platinum was sputtered onto the sample surface to enhance its conductivity. Transmission electron microscopy (TEM) images were obtained with an FEI Tecnai F20 instrument operating at 200 kV. Typically, 10 μL of 0.05 wt% sample solution was added dropwise onto a 300-mesh carbon support film before drying at room temperature. The zeta potential of 0.1 wt% CNC suspension was measured using a Malvern Zetasizer Nano-ZS90 at 25 °C, with an equilibration time of 120 s. The same molar ratio of Fe to CNCs was used for the measurement of CNC-Fe³⁺ as in the preparation of CNC/ β -FeOOH, with a concentration of 0.1 wt%. The magnetic hysteresis (M–H) curves were collected on a superconducting quantum interferometer (SQUID, MPMS XL5) provided by Lake Shore Cryotronics, Inc., using a maximum applied magnetic field of $\pm 2 \text{ T}$ at 298 K. The ultraviolet-visible (UV-vis) and circular dichroism (CD) spectra of different dispersions were recorded with a Bio-Logic MOS-450 spectrophotometer. The film samples used for these measurements were 10 mm by 10 mm in size, whose surface was normal to the incident beam. The data were collected from 200 to 750 nm, with a data pitch of 1 nm, a slit width of 2 mm and an acquisition duration of 0.1 s. During the CD test, a weak magnetic field of approximately 50 mT was provided using a pair of annular neodymium magnets and was parallel to the direction of light propagation. The strength of the magnetic field was measured using an HT201 Gauss meter. Polarized optical microscopy (POM) was performed using a Leica DM4000M versatile upright microscope, with an exposure time of 250 ms. Diffusion reflectance circular dichroism (DRCD) spectra and corresponding UV-vis spectra were obtained in reflectance mode using a JASCO J-1700. The scan rate, data pitch, bandwidth and digital integration time were 200 nm min^{−1}, 1 nm, 1 nm, and 1 s, respectively.

Data availability

The data that support the findings of this study are available from the corresponding author, M. Y., upon reasonable request.

Author contributions

J. Tang: data curation, formal analysis, investigation, methodology, and writing – original draft; S. Feng: supervision; M.



Yang: conceptualization, funding acquisition, project administration, supervision, validation, and writing – review & editing.

Conflicts of interest

There are no conflicts to declare.

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References

- W. Ma, L. Xu, A. F. De Moura, X. Wu, H. Kuang, C. Xu and N. A. Kotov, *Chem. Rev.*, 2017, **117**, 8041–8093.
- C. Li, J. Zhao, X. Gao, C. Hao, S. Hu, A. Qu, M. Sun, H. Kuang, C. Xu and L. Xu, *Adv. Mater.*, 2023, **35**, 2308198.
- J. Lu, Y. Xue, K. Bernardino, N. N. Zhang, W. R. Gomes, N. S. Ramesar, S. Liu, Z. Hu, T. Sun, A. F. de Moura, N. A. Kotov and K. Liu, *Science*, 2021, **371**, 1368–1374.
- H. E. Lee, H. Y. Ahn, J. Mun, Y. Y. Lee, M. Kim, N. H. Cho, K. Chang, W. S. Kim, J. Rho and K. T. Nam, *Nature*, 2018, **556**, 360–365.
- Y. Li, J. Cheng, J. Li, X. Zhu, T. He, R. Chen and Z. Tang, *Angew. Chem., Int. Ed.*, 2018, **57**, 10236–10240.
- J. Yeom, U. S. Santos, M. Chekini, M. Cha, A. F. De Moura and N. A. Kotov, *Science*, 2018, **359**, 309–314.
- Z. Li, Q. Fan, Z. Ye, C. Wu, Z. Wang and Y. Yin, *Science*, 2023, **380**, 1384–1390.
- T. Bai, J. Ai, J. Ma, Y. Duan, L. Han, J. Jiang and S. Che, *Angew. Chem., Int. Ed.*, 2021, **60**, 20036–20041.
- J. Cai, J. Zhao, X. Gao, W. Ma, D. Meng, H. Zhang, C. Hao, M. Sun, H. Kuang, C. Xu and L. Xu, *ACS Nano*, 2022, **16**, 11066–11075.
- T. Bai, J. Ai, Y. Duan, L. Han and S. Che, *Small*, 2022, **18**, 2104509.
- J. L. Greenfield, J. Wade, J. R. Brandt, X. Shi, T. J. Penfold and M. J. Fuchter, *Chem. Sci.*, 2021, **12**, 8589–8602.
- S. F. C. S. Chan, G. De Stasio, S. A. Welch, M. Girasole, B. H. Frazer, M. V. Nesterova and J. F. Banfield, *Science*, 2004, **1656**, 1656–1658.
- E. D. Sone, S. Weiner and L. Addadi, *J. Struct. Biol.*, 2007, **158**, 428–444.
- W. Huang, D. Montroni, T. Wang, S. Murata, A. Arakaki, M. Nemoto and D. Kisailus, *Acc. Chem. Res.*, 2022, **55**, 1360–1371.
- A. Querejeta-Fernández, G. Chauve, M. Methot, J. Bouchard and E. Kumacheva, *J. Am. Chem. Soc.*, 2014, **136**, 4788–4793.
- A. Querejeta-Fernández, B. Kopera, K. S. Prado, A. Klinkova, M. Methot, G. Chauve, J. Bouchard, A. S. Helmy and E. Kumacheva, *ACS Nano*, 2015, **9**, 10377–10385.
- X. Gao, B. Han, X. Yang and Z. Tang, *J. Am. Chem. Soc.*, 2019, **141**, 13700–13707.
- W. J. Kauzmann, J. E. Walter and H. Eyring, *Chem. Rev.*, 1940, **26**, 339–407.
- G. González-Rubio, J. Mosquera, V. Kumar, A. Pedrazo-Tardajos, P. Lombart, D. M. Solís, I. Lobato, E. G. Noya, A. Guerrero-Martínez, J. M. Taboada, F. Obelleiro, L. G. MacDowell, S. Bals and L. M. Liz-Marzán, *Science*, 2020, **368**, 1472–1477.
- S. Elazzouzi-Hafraoui, Y. Nishiyama, J. L. Putaux, L. Heux, F. Dubreuil and C. Rochas, *Biomacromolecules*, 2008, **9**, 57–65.
- M. Giese, L. K. Blusch, M. K. Khan and M. J. MacLachlan, *Angew. Chem., Int. Ed.*, 2015, **54**, 2888–2910.
- Z. Cheng, Y. Ma, L. Yang, F. Cheng, Z. Huang, A. Natan, H. Li, Y. Chen, D. Cao, Z. Huang, Y. H. Wang, Y. Liu, R. Yang and H. Zhu, *Adv. Opt. Mater.*, 2019, **7**, 1801816.
- T. Chen, Q. Zhao, X. Meng, Y. Li, H. Peng, A. K. Whittaker and S. Zhu, *ACS Nano*, 2020, **14**, 9440–9448.
- S. Kang, G. M. Biesold, H. Lee, D. Bukharina, Z. Lin and V. V. Tsukruk, *Adv. Funct. Mater.*, 2021, **31**, 2104596.
- S. Zhao, H. Zhu, J. Lu, M. Li, L. Zhao, L. Zhou and L. Gao, *Adv. Funct. Mater.*, 2022, **32**, 2201927.
- C. Wang, S. R. Mouchet, O. Deparis, J. Li, E. Paineau, D. Dragoe, H. Remita and M. N. Ghazzal, *Small*, 2024, **20**, 2402211.
- W. Zhang, X. Cheng, S. H. Chen and M. Anthamatten, *Langmuir*, 2023, **39**, 9180–9185.
- C. Wang, J. Li, E. Paineau, A. Laachachi, C. Colbeau-Justin, H. Remita and M. N. Ghazzal, *J. Mater. Chem. A*, 2020, **8**, 10779–10786.
- T. D. Nguyen, E. Lizundia, M. Niederberger, W. Y. Hamad and M. J. MacLachlan, *Chem. Mater.*, 2019, **31**, 2174–2181.
- J. Majoinen, J. Hassinen, J. S. Haataja, H. T. Rekola, E. Kontturi, M. A. Kostianen, R. H. A. Ras, P. Törmä and O. Ikkala, *Adv. Mater.*, 2016, **28**, 5262–5267.
- A. Chakraborty, Nonappa, B. Mondal, K. Chaudhari, H. Rekola, V. Hynninen, M. A. Kostianen, R. H. A. Ras and T. Pradeep, *J. Phys. Chem. C*, 2021, **125**, 3256–3267.
- H. Yu, H. Huang, J. Liang and J. Deng, *Nanoscale*, 2018, **10**, 12163–12168.
- K. Ståhl, K. Nielsen, J. Jiang, B. Lebech, J. C. Hanson, P. Norby and J. van Lanschot, *Corros. Sci.*, 2003, **45**, 2563–2575.
- C. Luna, M. Ilyn, V. Vega, V. M. Prida, J. González and R. Mendoza-Reséndez, *J. Phys. Chem. C*, 2014, **118**, 21128–21139.
- L. B. Mao, H. L. Gao, H. Bin Yao, L. Liu, H. Cölfen, G. Liu, S. M. Chen, S. K. Li, Y. X. Yan, Y. Y. Liu and S. H. Yu, *Science*, 2016, **354**, 107–110.
- S. Llàcer Navarro, K. Nakayama, A. Idström, L. Evenäs, A. Ström and T. Nypelö, *Cellulose*, 2021, **28**, 9633–9644.
- S. Mann, C. C. Perry and J. Webb, *Proc. R. Soc. London, Ser. B*, 1986, **227**, 179–190.
- R. M. H. Rumney, S. C. Robson, A. P. Kao, E. Barbu, L. Bozycki, J. R. Smith, S. M. Cragg, F. Couceiro, R. Parwani, G. Tozzi, M. Stuer, A. H. Barber, A. T. Ford and D. C. Górecki, *Nat. Commun.*, 2022, **13**, 3753.



- 39 S. Li, C. Huang, J. Hao and C. Wang, *J. Comput. Chem.*, 2014, **35**, 415–426.
- 40 G. Zhu, M. L. Sushko, J. S. Loring, B. A. Legg, M. Song, J. A. Soltis, X. Huang, K. M. Rosso and J. J. De Yoreo, *Nature*, 2021, **590**, 416–422.
- 41 J. F. Banfield, S. A. Welch, H. Zhang, T. T. Ebert and R. L. Penn, *Science*, 2000, **289**, 751–754.
- 42 B. C. Park, M. J. Ko, Y. K. Kim, G. W. Kim, M. S. Kim, T. M. Koo, H. E. Fu and Y. K. Kim, *Nat. Commun.*, 2022, **13**, 1144.
- 43 S. Dong and M. Roman, *J. Am. Chem. Soc.*, 2007, **129**, 13810–13811.
- 44 U. A. Agú, M. I. Oliva, S. G. Marchetti, A. C. Heredia, S. G. Casuscelli and M. E. Crivello, *J. Magn. Magn. Mater.*, 2014, **369**, 249–259.
- 45 D. M. Sherman and T. D. Waite, *Am. Mineral.*, 1985, **70**, 1262–1269.
- 46 H. Yao and Y. Ishikawa, *J. Phys. Chem. C*, 2015, **119**, 13224–13230.
- 47 B. Zhu, F. Wang, C. Wang, Y. Cao, L. Guo, J. Zhang and Y. Gu, *J. Phys.:Condens. Matter*, 2016, **28**, 285801.
- 48 C. Adamo and D. Jacquemin, *Chem. Soc. Rev.*, 2013, **42**, 845–856.
- 49 D. Wang, S. Feng and M. Yang, *ACS Nano*, 2024, **18**, 29636–29647.
- 50 A. Kumar and S. K. Gupta, *J. Nanopart. Res.*, 2013, **15**, 1466.
- 51 H. Zhang, X. Lv, Y. Li, Y. Wang and J. Li, *ACS Nano*, 2010, **4**, 380–386.
- 52 J. Hu, S. Li, J. Chu, S. Niu, J. Wang, Y. Du, Z. Li, X. Han and P. Xu, *ACS Catal.*, 2019, **9**, 10705–10711.
- 53 J. Zhao, B. Wu, X. Huang, Y. Sun, Z. Zhao and M. Ye, *Adv. Sci.*, 2022, **9**, 2201678.
- 54 X. Chen, Y. Zeng, Z. Chen, S. Wang and C. Xin, *Front. Chem.*, 2020, **8**, 328.
- 55 H. Qi, L. Cao, J. Li, J. Huang, Z. Xu, Y. Cheng, X. Kong and K. Yanagisawa, *ACS Appl. Mater. Interfaces*, 2016, **8**, 35253–35263.
- 56 J. Liu, M. Zheng, X. Shi, H. Zeng and H. Xia, *Adv. Funct. Mater.*, 2016, **26**, 919–930.
- 57 T. Lu and F. Chen, *J. Comput. Chem.*, 2012, **33**, 580–592.
- 58 Z. Liu, T. Lu and Q. Chen, *Carbon*, 2020, **165**, 461–467.
- 59 X. Gao, X. Zhang, K. Deng, B. Han, L. Zhao, M. Wu, L. Shi, J. Lv and Z. Tang, *J. Am. Chem. Soc.*, 2017, **139**, 8734–8739.
- 60 H. Chen, J. Zhou and J. Deng, *Polym. Chem.*, 2016, **7**, 125–134.
- 61 A. K. Vishneratina, F. Purcell-Milton, R. Serrano-García, V. A. Kuznetsova, A. O. Orlova, A. V. Fedorov, A. V. Baranov and Y. K. Gun'ko, *J. Mater. Chem. C*, 2017, **5**, 1692–1698.
- 62 T. T. Zhuang, Y. Li, X. Gao, M. Wei, F. P. García de Arquer, P. Todorović, J. Tian, G. Li, C. Zhang, X. Li, L. Dong, Y. Song, Y. Lu, X. Yang, L. Zhang, F. Fan, S. O. Kelley, S. H. Yu, Z. Tang and E. H. Sargent, *Nat. Nanotechnol.*, 2020, **15**, 192–197.
- 63 J. S. Kamal, R. Gomes, Z. Hens, M. Karvar, K. Neyts, S. Compennolle and F. Vanhaecke, *Phys. Rev. B:Condens. Matter Mater. Phys.*, 2012, **85**, 035126.
- 64 J. Hu, L. S. Li, W. Yang, L. Manna, L. W. Wang and A. P. Alivisatos, *Science*, 2001, **292**, 2060–2063.
- 65 E. U. Condon, W. Altar and H. Eyring, *J. Chem. Phys.*, 1937, **5**, 339–407.
- 66 S. Ma, Y. K. Jung, J. Ahn, J. Kyhm, J. Tan, H. Lee, G. Jang, C. U. Lee, A. Walsh and J. Moon, *Nat. Commun.*, 2022, **13**, 3259.
- 67 H. Lu, C. Xiao, R. Song, T. Li, A. E. Maughan, A. Levin, R. Brunecky, J. J. Berry, D. B. Mitzi, V. Blum and M. C. Beard, *J. Am. Chem. Soc.*, 2020, **142**, 13030–13040.
- 68 M. Kepenekian, R. Robles, C. Katan, D. Saporì, L. Pedesseau and J. Even, *ACS Nano*, 2015, **9**, 11557–11567.
- 69 R. L. Penn and J. F. Banfield, *Science*, 1998, **281**, 969–971.
- 70 M. H. Tsai, S. Y. Chen and P. Shen, *Nano Lett.*, 2004, **4**, 1197–1201.
- 71 D. Qu, H. Zheng, H. Jiang, Y. Xu and Z. Tang, *Adv. Opt. Mater.*, 2019, **7**, 1801395.
- 72 S. N. Fernandes, P. L. Almeida, N. Monge, L. E. Aguirre, D. Reis, C. L. P. de Oliveira, A. M. F. Neto, P. Pieranski and M. H. Godinho, *Adv. Mater.*, 2017, **29**, 1603560.
- 73 P. Li, L. Li, K. J. Jeong, X. Yu, X. Yu and Y. Xu, *Adv. Opt. Mater.*, 2022, **10**, 2102616.
- 74 X. Wang, Q. Wang, X. Zhang, J. Miao, J. Cheng, T. He, Y. Li, Z. Tang and R. Chen, *Adv. Opt. Mater.*, 2022, **10**, 2200761.
- 75 A. Salij, R. H. Goldsmith and R. Tempelaar, *J. Am. Chem. Soc.*, 2021, **143**, 21519–21531.
- 76 H. Maeda and Y. Maeda, *Langmuir*, 2011, **27**, 2895–2903.
- 77 D. Lu, X. Yu, S. Zhu and Y. Xu, *Cellulose*, 2023, **30**, 8325–8338.
- 78 S. Beck, J. Bouchard, G. Chauve and R. Berry, *Cellulose*, 2013, **20**, 1401–1411.
- 79 B. Frka-Petesic, G. Guidetti, G. Kamita and S. Vignolini, *Adv. Mater.*, 2017, **29**, 1701469.
- 80 T. Nypelö, *J. Mater. Chem. C*, 2022, **10**, 805–818.
- 81 F. Kimura, T. Kimura, M. Tamura, A. Hirai, M. Ikuno and F. Horii, *Langmuir*, 2005, **21**, 2034–2037.
- 82 X. Zhang, S. Kang, K. Adstedt, M. Kim, R. Xiong, J. Yu, X. Chen, X. Zhao, C. Ye and V. V. Tsukruk, *Nat. Commun.*, 2022, **13**, 5804.
- 83 X. Liu, Y. Du, S. Mourdikoudis, G. Zheng and K. Y. Wong, *Adv. Opt. Mater.*, 2023, **11**, 2202859.
- 84 K. J. Jeong, D. K. Lee, V. T. Tran, C. Wang, J. Lv, J. Park, Z. Tang and J. Lee, *ACS Nano*, 2020, **14**, 7152–7160.
- 85 B. Han, X. Gao, J. Lv and Z. Tang, *Adv. Mater.*, 2020, **32**, 1801491.
- 86 T. Ji, S. Che and Y. Duan, *Acc. Mater. Res.*, 2023, **4**, 1057–1067.
- 87 X. Wu, C. Hao, L. Xu, H. Kuang and C. Xu, *Small*, 2020, **16**, 1905734.
- 88 Y. Kitagawa, H. Segawa and K. Ishii, *Angew. Chem., Int. Ed.*, 2011, **50**, 9133–9136.
- 89 G. L. J. A. Rikken and E. Raupach, *Nature*, 1997, **390**, 493–494.
- 90 B. J. Lemaire, P. Davidson, J. Ferré, J. P. Jamet, P. Panine, I. Dozov and J. P. Jolivet, *Phys. Rev. Lett.*, 2002, **88**, 125507.

