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# Enhanced performance of perovskite solar cells via a bilateral electron-donating passivator as a molecule bridge†

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Defects at the buried interface are the primary factors contributing to recombination losses and instability in perovskite solar cells (PSCs) with n-i-p structure. Here, a molecule with bilateral electron-donating groups, 6-amino-1-hexanol (HAL), is introduced between SnO<sub>2</sub> and perovskite (PVK) to optimize the characteristics of the buried interfacial properties, as well as the PVK film quality. The surface defects of SnO<sub>2</sub> can be more effectively passivated, and its energy level structure can be tuned more appropriately. Meanwhile, the electron-donating groups in HAL can passivate the Pb<sup>2+</sup> defects and stabilize the [PbI<sub>6</sub>]<sup>4-</sup> octahedra at the buried interface. With the assistance of bilateral electron-donating groups, HAL can act as a molecular bridge to easily bridge SnO<sub>2</sub> and PVK to passivate the buried interfacial defects, thus enhancing device performance. As a result, the photovoltaic performance was significantly improved by the buried modification with HAL, ultimately achieving a champion PCE of 23.58%, far superior to the inferior PCE of 21.30% of the pristine device. Furthermore, after 35 days of storage at approximately 35% relative humidity (RH) and room temperature, the device based on HAL-modified SnO<sub>2</sub> demonstrates significantly better stability than the pristine device, as it still maintains over 90% of the initial efficiency. This study provides an effective molecular bridge strategy for further enhancing the performance of PSCs.

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## 1. Introduction

Since the excellent photovoltaic performance of perovskites (PVK) has been widely acknowledged, extensive research has been conducted on perovskite solar cells (PSCs) with the hope that they may serve as the next-generation replacement for silicon-based solar cells. In recent years, the power conversion efficiency (PCE) of PVK-based photovoltaic devices has significantly increased, from only 3.8% to over 25% certified PCE.<sup>1–12</sup>

For PSC devices with n-i-p structure, the quality of the electron transport layer (ETL) is a crucial factor in achieving stable and efficient normal PVK devices. This is because the incoming light must pass through the ETL before reaching the PVK layer, making it crucial to select the ETL to match the PVK

layer rationally.<sup>13</sup> Inorganic ETLs, with the advantages of low cost, high stability, and facile process technology, have been regarded as the most promising electron transport materials for preparing high-performance PSCs.<sup>14</sup> The most widely employed inorganic electron transport materials in normal PSCs include SnO<sub>2</sub>, TiO<sub>2</sub>, ZnO, and so on.<sup>14</sup> The PSCs based on TiO<sub>2</sub> and SnO<sub>2</sub> ETLs have both demonstrated high PCE exceeding 25%.<sup>12,15</sup> Notably, SnO<sub>2</sub> may be more promising than TiO<sub>2</sub> due to its non-photocatalytic activity and low-temperature processability.<sup>14,16</sup> Inevitably, SnO<sub>2</sub> also has certain disadvantages, such as oxygen vacancies and Sn interstitial defects on its surface, which serve as non-radiative recombination centers resulting in recombination losses.<sup>17,18</sup> Moreover, the construction of efficient and stable SnO<sub>2</sub>-based PSCs also requires high-quality PVK thin films. Simultaneously, solution-processed PVK films will generate high defect densities, which adversely affects film quality and surface contact and is a limiting factor in achieving further device performance improvements.<sup>19,20</sup> Generally, the deep-level defect concentration at the buried interface of PVK is much higher than that at the upper surface, which not only hampers electron transport but also severely impedes the crystal growth of the PVK layer.<sup>21,22</sup> Therefore, the buried interfacial defects will impede the enhancement of open-circuit voltage (*V*<sub>oc</sub>) and fill factor

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(FF) in PSCs.<sup>23</sup> Presently, the introduction of a functional molecule that can exert positive influences on both SnO<sub>2</sub> and PVK simultaneously remains a significant challenge. To tackle this challenge, many researchers have been inclined to investigate the optimization of the buried interface of PVK.<sup>17,24–27</sup> Clearly, the buried interfacial modification has emerged as an effective strategy for achieving high-efficiency PSCs with n–i–p structure.

Herein, we introduced a bilateral electron-donating groups passivator (6-amino-1-hexanol, HAL) as a molecular bridge at the buried interface of PSCs aiming to simultaneously regulate the surface defects of SnO<sub>2</sub> and the quality of PVK films. The HAL with the bilateral electron-donating groups can effectively mitigate surface defects on SnO<sub>2</sub> to suppress carrier recombination. Meanwhile, the presence of HAL on the SnO<sub>2</sub> surface is beneficial for improving energy-level structure and increasing electron donor species, which ultimately enhances electron extraction and transport. Furthermore, the improved surface wettability of SnO<sub>2</sub> films facilitates the formation of large PVK grains with low defect densities. Benefiting from the presence of electron-donating groups, the surface properties of SnO<sub>2</sub> are positively regulated, and the quality of PVK films is effectively improved. As a result, compared to the pristine device (21.30%), the n–i–p structured PSCs based on HAL-modified SnO<sub>2</sub> achieved a significant improvement, with a champion PCE of up to 23.58%. More importantly, the unencapsulated devices modified with HAL maintained over 90% of the initial PCE after aging for 35 days at approximately 35% relative humidity (RH). The present work demonstrates that molecules with bilateral electron-donating groups can act as molecular bridges at the buried interface of the SnO<sub>2</sub>-based PSCs, positively enhancing device performance.

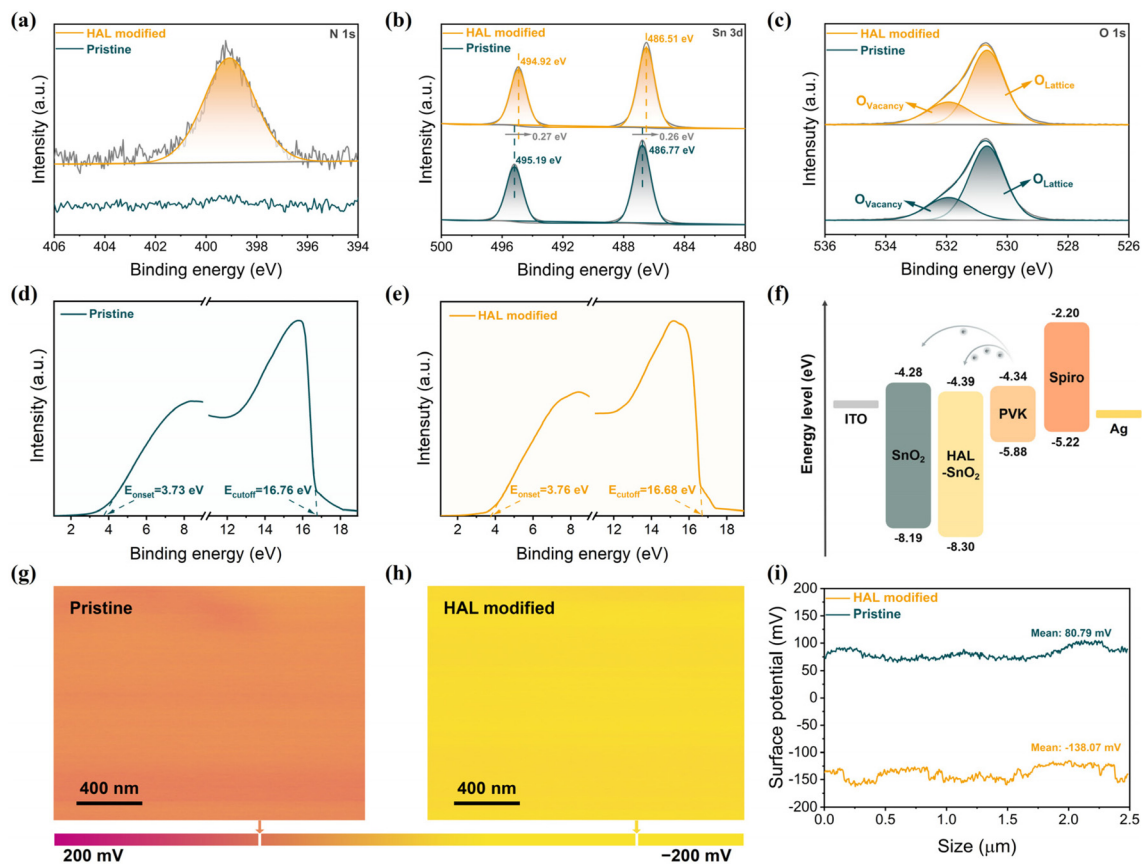
## 2. Results and discussion

The buried interfacial modification was used to introduce a linear organic molecule (6-amino-1-hexanol, HAL) between SnO<sub>2</sub> and PVK, whose chemical structure (Fig. S1†) has two functional groups, amino (–NH<sub>2</sub>) and hydrocarbon (–OH) groups. To confirm the possibility of HAL interacting with SnO<sub>2</sub> and PVK, we simulated the electrostatic potential (ESP) of HAL. From Fig. S2,† it can be observed that the high charge distribution of –NH<sub>2</sub> and –OH groups in the HAL molecule indicates that the –NH<sub>2</sub> and –OH groups of HAL are expected to passivate the oxygen vacancies of SnO<sub>2</sub>. More importantly, according to previous works, the uncoordinated Pb<sup>2+</sup> of the [PbI<sub>6</sub>]<sup>4–</sup> octahedral and the electron-donating groups (–NH<sub>2</sub>/–OH) could form Lewis adducts through the interaction between the lone-pair electrons in them, and the formation of Lewis acid–base pairs of Pb<sup>2+</sup> with electron-donating groups helps to stabilize the [PbI<sub>6</sub>]<sup>4–</sup> octahedra.<sup>28–30</sup> Furthermore, the –NH<sub>2</sub> and –OH groups of HAL can connect with the iodide from [PbI<sub>6</sub>]<sup>4–</sup> octahedra *via* H-bond formation, which is also beneficial for the stabilization of [PbI<sub>6</sub>]<sup>4–</sup> octahedra at the buried interface of PVK.<sup>31–33</sup>

To investigate the effect of introducing HAL as a modifier on SnO<sub>2</sub>, the analytical results for the characterization using X-ray photoelectron spectroscopy (XPS) can be depicted in Fig. 1a–c and Fig. S3.† From the N 1s spectra (Fig. 1a) of pristine SnO<sub>2</sub> and HAL-modified SnO<sub>2</sub>, the appearance of N 1s peaks in HAL-modified SnO<sub>2</sub> can confirm the presence of HAL on the SnO<sub>2</sub> surface, which is not present in the pristine films. The positions of the double peaks representing Sn 3d<sub>2/3</sub> and Sn 3d<sub>3/5</sub> in the Sn 3d patterns (Fig. 1b) of the HAL-modified SnO<sub>2</sub> films are shifted by 0.27 eV toward the low binding energy position relative to the pristine SnO<sub>2</sub>, which implies an enhancement of the electron cloud density. The change in electron cloud density may be attributed to the fact that the nitrogen in –NH<sub>2</sub> and oxygen in –OH with lone pairs of electrons can coordinate with the uncoordinated Sn<sup>4+</sup> in SnO<sub>2</sub>, which induces electron transfer from nitrogen and oxygen elements in HAL to SnO<sub>2</sub>.<sup>34,35</sup> The result of Sn 3d spectra confirms the existence of interaction between HAL and SnO<sub>2</sub>. From the O 1s spectra in Fig. 1c, an asymmetric peak profile appears in both pristine SnO<sub>2</sub> and HAL-modified SnO<sub>2</sub> films. The O 1s peaks of all samples can be deconvoluted into two peaks, the peak at 530.67 eV belongs to lattice oxygen (O<sub>Lattice</sub>), and the peak at 531.93 eV originates from vacancy oxygen (O<sub>Vacancy</sub>). Utilizing the  $O_{\text{Vacancy}}/(O_{\text{Lattice}} + O_{\text{Vacancy}})$  formula in conjunction with analysis of the O 1s spectra, we can calculate the peak area ratio for O<sub>Vacancy</sub>.<sup>34</sup> The calculated results showed that the peak area ratio of O<sub>Vacancy</sub> decreased from 0.43 to 0.38 for the HAL-modified SnO<sub>2</sub> sample compared with the pristine SnO<sub>2</sub> sample. This result proves that HAL can effectively passivate the oxygen vacancies, thus inhibiting defect-induced non-radiative recombination at the SnO<sub>2</sub>/PVK interface.

The interaction between HAL and SnO<sub>2</sub> motivated us to further investigate the influence of SnO<sub>2</sub> modification on the properties of SnO<sub>2</sub> films. From Fig. S4,† the UV-Vis transmission spectra demonstrate that the modification of HAL did not affect the high transmittance of the SnO<sub>2</sub> films. The UV-vis absorption spectroscopy (Fig. S5†) measurements of different SnO<sub>2</sub> films were also performed and converted into the Tauc plots. From the Tauc plots in Fig. S6,† the band gap ( $E_g$ ) of the SnO<sub>2</sub> film was not altered by the introduction of HAL, and the  $E_g$  of the HAL-modified SnO<sub>2</sub> film remained at 3.91 eV. Furthermore, possibly attributed to the presence of the electron-donating groups (–NH<sub>2</sub>/–OH) with lone pairs of electrons, the Tafel curves of different devices in Fig. S7† reveal that the HAL-modified sample exhibits higher electrical conductivity than the pristine sample.

At the same time, the ultraviolet photoelectron spectroscopy (UPS) technique was also used to investigate whether the interaction of HAL with SnO<sub>2</sub> changed the energy-level structure of SnO<sub>2</sub>/HAL. Fig. 1d and e exhibit the UPS spectra of pristine and HAL-modified SnO<sub>2</sub> films, respectively. Based on the data of the onset ( $E_{\text{onset}}$ ) and the secondary electron cutoff ( $E_{\text{cutoff}}$ ) of the UPS spectra, the conduction band minimum ( $E_{\text{CBM}}$ ), work function ( $W_F$ ), and valence band maximum ( $E_{\text{VBM}}$ ) of the samples can be arithmetically derived.<sup>36</sup> The detailed para-



**Fig. 1** (a) N 1s, (b) Sn 3d, and (c) O 1s XPS spectra of different SnO<sub>2</sub> films. UPS spectra of (d) SnO<sub>2</sub> and (e) HAL-modified SnO<sub>2</sub> films. (f) Energy band diagram of the materials employed in PSCs. KPFM images of (g) SnO<sub>2</sub> and (h) HAL-modified SnO<sub>2</sub> films. (i) Surface potential of different SnO<sub>2</sub> films.

eters of the energy level for the samples are displayed in Table S1.† Combining the data obtained from UPS analysis and the energy level distribution from previous literature, we plotted the corresponding energy band diagrams in Fig. 1f.<sup>36</sup> It is noteworthy from Fig. 1f that the  $E_{\text{CBM}}$  of the pristine and HAL-modified SnO<sub>2</sub> films are  $-4.28$  eV and  $-4.39$  eV, respectively. Compared to the  $E_{\text{CBM}}$  ( $-4.34$  eV) of PVK, the pristine SnO<sub>2</sub> film with a higher  $E_{\text{CBM}}$  ( $-4.28$  eV) would hinder electron transport and induce carrier accumulation at the buried interface. In contrast, the HAL-modified SnO<sub>2</sub> film with a lower  $E_{\text{CBM}}$  ( $-4.39$  eV) than that of PVK would be more favourable for electron transport and extraction between SnO<sub>2</sub> and PVK. The analytical results based on UPS measurements illustrate that the modification of HAL can construct a smoother energy gradient between SnO<sub>2</sub> and PVK, which facilitates the reduction of energy loss and the enhancement of photovoltage. More importantly, the HAL-modified SnO<sub>2</sub> film with a deeper  $E_{\text{VBM}}$  can create a higher energy barrier for hole injection at the SnO<sub>2</sub>/PVK interface, further suppressing the loss of charge recombination. Additionally, surface electronic properties of SnO<sub>2</sub> films were assessed for various samples by measuring their contact potential difference (CPD) using Kelvin probe force microscopy (KPFM) techniques. From Fig. 1g–i, it can be noted that the HAL-modified SnO<sub>2</sub> film has

a lower average CPD of  $-138.07$  mV than the pristine SnO<sub>2</sub> film ( $80.79$  mV). The reduction in CPD may be attributed to the increased number of electron donor species associated with the introduction of HAL, which would result in an increase in the film work function, in agreement with the findings from UPS analysis.

For devices with n–i–p structures, the surface properties of SnO<sub>2</sub> films profoundly affect the quality of the deposited PVK films. To determine the influence of HAL modification on the morphology of SnO<sub>2</sub> films, we measured the films using scanning electron microscopy (SEM) and atomic force microscopy (AFM). From Fig. S8,† it can be observed that there is significant agglomeration on the pristine SnO<sub>2</sub> film. In contrast, there is no significant agglomeration on the HAL-modified SnO<sub>2</sub> film. As seen from the AFM images of Fig. S9,† the surface of HAL-modified SnO<sub>2</sub> is flatter and more uniform, and the roughness of the modified SnO<sub>2</sub> is significantly reduced relative to the pristine SnO<sub>2</sub>, which contributes to the interfacial contact between SnO<sub>2</sub> and PVK.<sup>20</sup> Besides, the contact angle tests were performed on the pristine and HAL-modified SnO<sub>2</sub> films using the DMF/DMSO mixture solution, as illustrated in Fig. S10.† Compared to ITO/SnO<sub>2</sub> ( $20^\circ$ ), the DMF/DMSO mixture solution had a smaller contact angle ( $8^\circ$ ) on ITO/SnO<sub>2</sub>/HAL. The smaller the contact angle between the

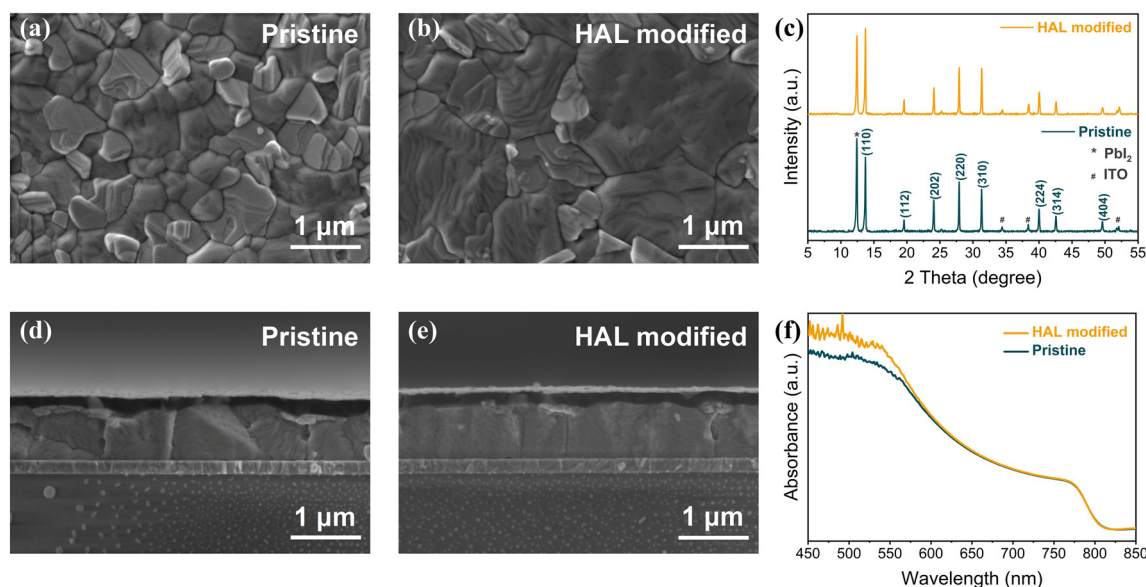
DMF/DMSO mixture solution and the substrate, the smaller the Gibbs free energy required for nucleation.<sup>37</sup> The above analyses demonstrate that the modification of HAL is more favorable to the nucleation of PVK and the improvement of the interfacial contact between PVK and SnO<sub>2</sub>.

To further verify that the introduction of HAL can ameliorate the quality of PVK films, we also measured the top-view and cross-sectional SEM images of samples with ITO/SnO<sub>2</sub> (without and with HAL)/PVK structures. From Fig. 2a and b, it can be seen that the PVK film of the HAL-modified sample has a flatter surface, more blurred grain boundaries, and larger grain size relative to the pristine sample. The cross-sectional SEM images (Fig. 2d and e) also directly verify the results of the top-view SEM images with better quality for the PVK films with HAL-modified SnO<sub>2</sub>. The surface roughness of PVK films based on different ETL substrates was analyzed using AFM measurements. It is known from Fig. S11† that the roughness (28.5 nm) of PVK film on the HAL-modified SnO<sub>2</sub> is lower compared to the pristine PVK film (38.3 nm). The smaller the roughness of the PVK surface, the more favorable the contact at the PVK/HTL interface.<sup>20</sup> This phenomenon contributes to the extraction and transfer of holes from PVK to HTL.

From the XRD patterns in Fig. 2c, it can be noted that the modification of HAL did not alter the crystal structure of PVK. However, the (110) crystalline plane of the PVK films with HAL-modified SnO<sub>2</sub> has a stronger diffraction peak intensity, which indicates that the modification of HAL improved the crystalline intensity of PVK. The XRD results are consistent with the SEM analysis. Additionally, UV-Vis absorption spectroscopy measurements were also used to determine the absorption of PVK films based on different substrates. Due to the improved quality of the PVK films, it can be noted from Fig. 2f that the PVK film with HAL-modified SnO<sub>2</sub> displays a

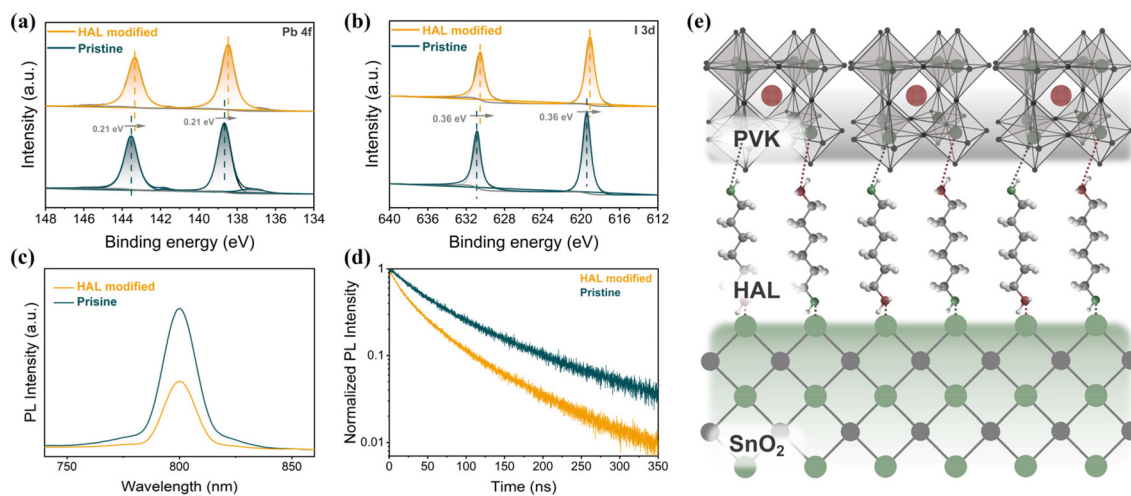
higher absorption intensity in the range of 400–600 nm, which indicates that the buried modification of HAL has a positive effect on the PVK films. As shown in the Tauc plots of Fig. S12† obtained from the conversion of Fig. 2f, we find that the  $E_g$  of the PVK film does not change due to the modification of HAL. Further, we obtained the corresponding plots (Fig. S13†) of Urbach energy ( $E_U$ ) from the UV-Vis absorption spectra. The  $E_U$  is a specific physical parameter that is related to the crystallinity, trap states, and electronic disorder of the material, from which we can infer the PVK film quality. From Fig. S13,† the obtained data reflect that the PVK film with the HAL-modified SnO<sub>2</sub> has a lower  $E_U$  (37.74 meV) relative to the pristine samples (46.95 meV). The smaller the  $E_U$ , the better the PVK film quality.<sup>38</sup> The result of  $E_U$  is consistent with our previous findings, which prove that the modification of HAL can effectually enhance the quality of PVK films and suppress the defect states.

The improved quality of PVK films leads us to consider whether there is also an interaction between HAL and PVK. The XPS spectra based on the PbI<sub>2</sub> films without and with HAL are shown in Fig. 3a and b. It can be observed from the Pb 4f spectra in Fig. 3a that after the introduction of HAL, the double peaks on behalf of Pb 4f<sub>7/2</sub> and Pb 4f<sub>5/2</sub> are overall shifted towards the lower binding energy by 0.21 eV compared to the pristine films. The variation of the binding energy in the Pb 4f caused by the HAL modification demonstrates the presence of an interaction between HAL and PbI<sub>2</sub>. Furthermore, the two shoulder peaks representing Pb<sup>0</sup> are more obvious on the pristine PVK film than on the PVK film based on the HAL-modified SnO<sub>2</sub>, implying that the presence of HAL can inhibit the Pb<sup>0</sup> defects. From the I 3d spectra in Fig. 3b, the I 3d<sub>5/2</sub> and 3d<sub>3/2</sub> peaks are shifted by 0.36 eV overall toward the lower binding energy relative to the pristine



**Fig. 2** Top-view SEM images of PVK films with (a) pristine and (b) HAL-modified SnO<sub>2</sub> films. (c) XRD patterns and (f) UV-Vis absorption spectra of PVK films with different SnO<sub>2</sub> films. Cross-sectional SEM images of PSC devices with (d) pristine and (e) HAL-modified SnO<sub>2</sub> films.





**Fig. 3** (a) Pb 4f and (b) I 3d XPS spectra of PbI<sub>2</sub> films based on different SnO<sub>2</sub> films. (c) SSPL and (d) TRPL spectra of PVK films based on different SnO<sub>2</sub>. (e) Schematic diagram illustrating potential interaction mechanisms for the buried interfacial modification of HAL.

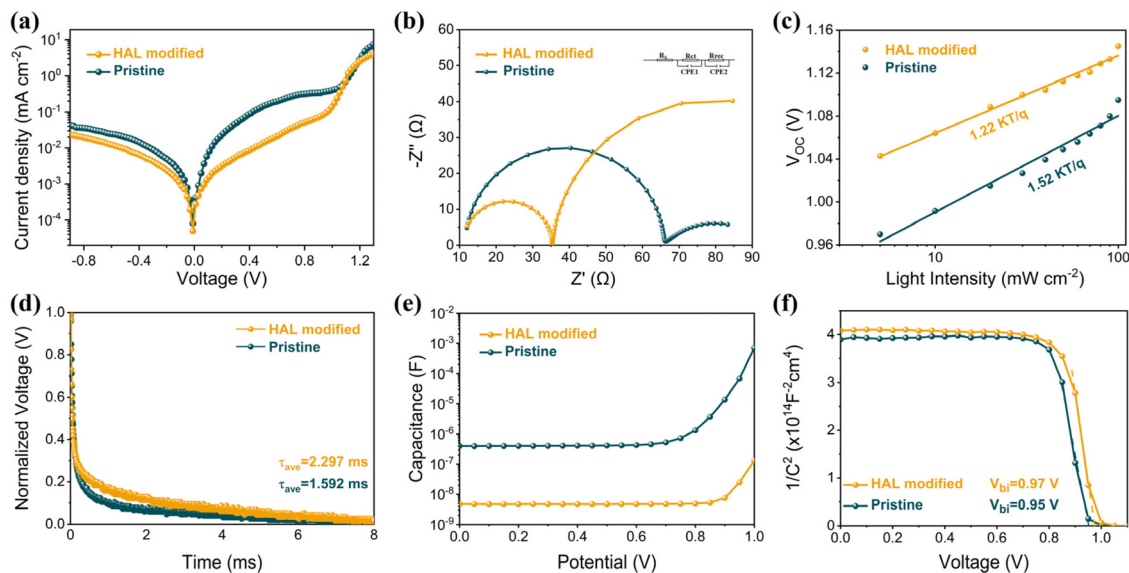
film. The observed shift in the I 3d XPS spectra suggests that the interaction between HAL and PbI<sub>2</sub> may have resulted in a decrease in the oxidation state of iodide or weakening of the Pb–I bond.<sup>39</sup> The interaction of HAL and PbI<sub>2</sub> may be attributed to the existence of the electron-donating groups (–NH<sub>2</sub>/–OH) with lone pairs of electrons in HAL. As illustrated in Fig. 3f, we have drawn the relevant schematic diagrams to demonstrate the buried modification of HAL more intuitively. Due to the electron-donating groups with lone pairs of electrons on HAL, the introduction of HAL can effectively passivate the surface defects on SnO<sub>2</sub> and improve the contact at the SnO<sub>2</sub>/PVK interface. Meanwhile, the electron-donating groups of HAL can modulate the uncoordinated Pb<sup>2+</sup> defects and stabilize [PbI<sub>6</sub>]<sup>4–</sup> octahedra at the bottom interface of PVK through this coordination and hydrogen bonding.

To clarify the effect of the buried modifications of HAL on the carrier dynamics of the SnO<sub>2</sub>/PVK interface, the steady-state photoluminescence (SSPL) and time-resolved photoluminescence (TRPL) measurements of the samples with ITO/SnO<sub>2</sub> (without and with HAL)/PVK structures were revealed in Fig. 3c and d, respectively. In the SSPL spectra in Fig. 3c, it can be found that the sample with the introduction of HAL for the modification obviously exhibits a lower steady-state PL intensity than the samples without the introduction of HAL, which suggests a more efficient extraction and transport of photo-generated electrons at the SnO<sub>2</sub>/PVK interface. The carrier extraction lifetimes of different samples were determined by fitting the TRPL plots shown in Fig. 3d with a bi-exponential decay equation, and the detailed fitted data are displayed in Table S2.†<sup>34</sup> Compared with the pristine sample, the average carrier lifetime of the HAL-modified sample is lower, proving that the HAL modification favors the transport of photo-generated electrons from PVK to SnO<sub>2</sub>. The reduction in the average carrier lifetime of the sample with the ITO/SnO<sub>2</sub>/HAL/PVK structure corresponds to the reduced steady-state PL intensity. These findings strongly confirm that the introduction of HAL

can effectively reduce the defects of SnO<sub>2</sub> and PVK and improve electron extraction and transport by suppressing the non-radiative recombination, ultimately improving  $V_{OC}$  and FF in the device.

Measuring the impact of buried modifications on improving performance using a complete device with n–i–p structure is necessary. Fig. 4a shows that the leakage current of the PSC with HAL is smaller than the pristine PSC. The decrease in dark current suggests that the introduction of HAL can effectively mitigate current leakage, improving the  $J_{SC}$  and FF of the PSCs.<sup>40</sup> Electrochemical impedance spectroscopy (EIS) was employed to acquire additional information on carrier transport behavior and interfacial properties. The Nyquist plots of the PSC devices at 1.0 V bias voltage were acquired under AM 1.5 G solar irradiation conditions. The corresponding equivalent circuit diagram is marked in Fig. 4b, and the fitted measurement parameters are listed in Table S3.† The Nyquist plot is made up of two semicircles that are located in different frequency ranges. In impedance measurements, the high-frequency component usually indicates the charge transfer resistance ( $R_{ct}$ ), while the low-frequency component is typically associated with the complex resistance ( $R_{rec}$ ).<sup>39</sup> Relative to the pristine device, the HAL-modified device exhibits a significant decrease in  $R_{ct}$  coupled with a marked increase in  $R_{rec}$ , which demonstrated that the introduction of HAL can effectively enhance carrier extraction and transfer, as well as inhibit carrier recombination at the ETL/PVK interface.

From the light-intensity-dependent  $V_{OC}$  measurement in Fig. 4c, the device with HAL displays a smaller slope (1.22 kT q<sup>–1</sup>) compared to the pristine device (1.52 kT q<sup>–1</sup>). The deviation of the slope from unity (kT q<sup>–1</sup>) is a well-established metric that reflects the presence of trap-assisted non-radiative recombination in PSCs. The lower value for the HAL-modified device indicates a reduction in non-radiative recombination.<sup>18</sup> This result implies that the HAL modification strategy may lead to improved performance and stability of PSCs by redu-



**Fig. 4** (a) The dark-state  $J$ - $V$  curves, (b) Nyquist plots, (c) Light-intensity-dependent  $V_{OC}$  curves, (d) TPV curves, (e)  $C$ - $V$  curves, and (f)  $M$ - $S$  curves for PSCs based different  $\text{SnO}_2$  films.

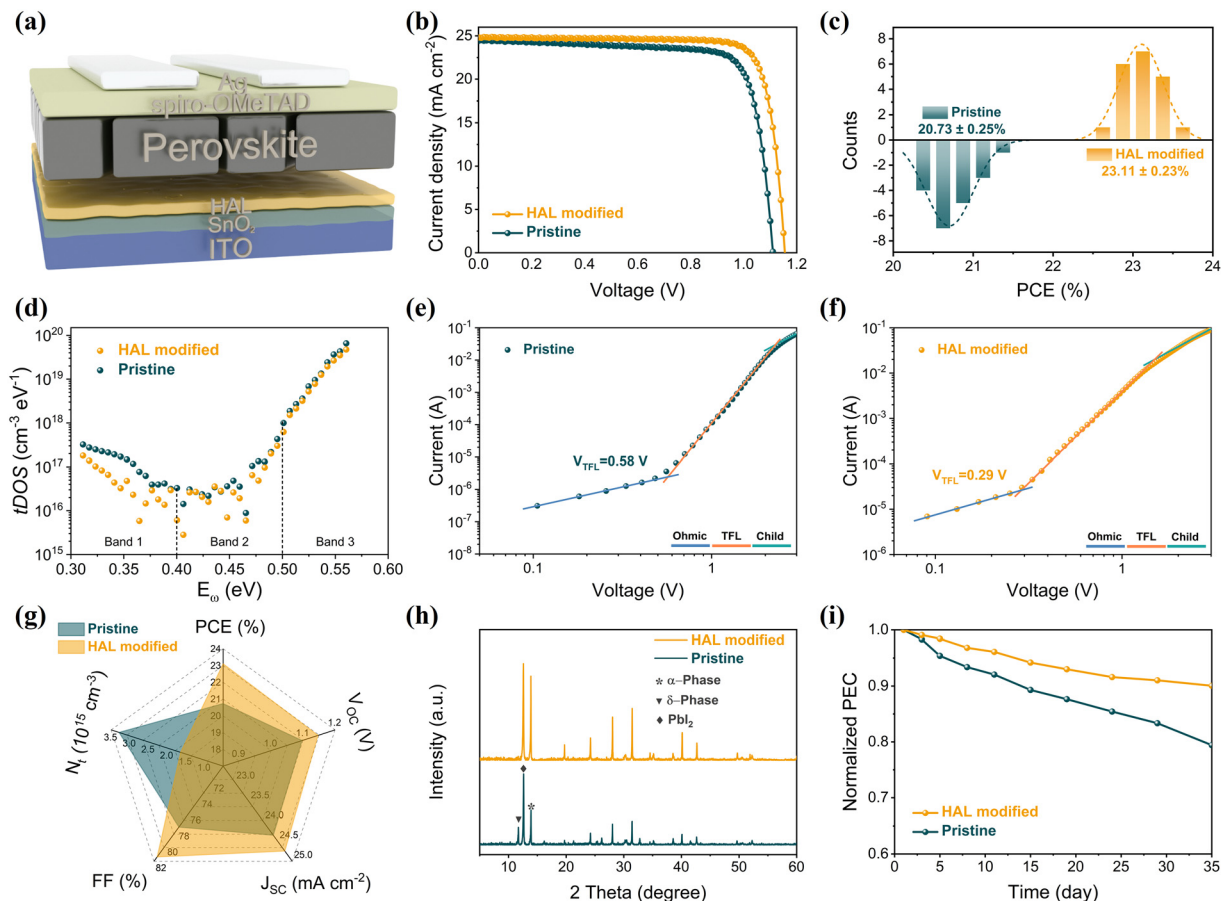
cing the energy losses attributed to non-radiative recombination. In addition, the transient photovoltage (TPV) measurements of the device further confirm that the introduction of HAL promotes the carrier transfer and suppresses the recombination, and the results are shown in Fig. 4d. After introducing HAL, the decay rate of TPV slows down, indicating a longer carrier lifetime, which is associated with reduced non-radiative interfacial recombination.<sup>21</sup> In other words, the introduction of HAL can effectively passivate defect-induced recombination to achieve faster carrier transport and longer carrier lifetime.

Fig. 4e shows the Capacitance-Voltage ( $C$ - $V$ ) plots of the devices with different ETLs. It can be clearly seen that the capacitance of the device with HAL is much smaller than that of the pristine, which proves that the presence of HAL can reduce the charge accumulation at the buried interface.<sup>36</sup> The Mott-Schottky ( $M$ - $S$ ) curves presented in Fig. 4f clearly demonstrate that the built-in potential ( $V_{bi}$ ) of the device with HAL is significantly greater than that of the original device. The increase of  $V_{bi}$  means a stronger electric field between ETL and PVK, which is beneficial for enhanced carrier separation and extraction.<sup>40</sup> These properties promote the charge collection, increasing the FF and contributing to higher  $V_{OC}$ .

Fig. 5a illustrates the device structure of PSCs with n-i-p structure in this work, where HAL is inserted between  $\text{SnO}_2$  and PVK layers by the buried modification. The positive effects of HAL on the passivation of  $\text{SnO}_2$  defects, the tuning of energy levels, and the promotion of PVK crystallization are bound to have an impact on the photovoltaic performance of PSCs. To further explore the optimal concentration of HAL for efficient PSCs, we set up concentration gradients of 0, 0.5, 1.0 and 1.5  $\text{mg mL}^{-1}$  and assembled the complete devices. The  $J$ - $V$  curves revealing the impact of different HAL concentrations on the performance of devices are displayed in Fig. S14,<sup>†</sup> and the detailed corresponding parameters are

listed in Table S4.<sup>†</sup> From the data collected from the  $J$ - $V$  curves, we can find that the optimal HAL concentration is 1.0  $\text{mg mL}^{-1}$ . Fig. 5b illustrates a comparison of the PCE between the champion and pristine devices, revealing the superior performance of the champion device. The champion device can achieve a PCE of 23.58% (with a  $V_{OC}$  of 1.155 V, FF of 0.8227, and  $J_{SC}$  of 24.82  $\text{mA cm}^{-2}$ ), which is significantly higher than the pristine device with a PCE of 21.30%,  $V_{OC}$  of 1.110 V, FF of 0.7861, and  $J_{SC}$  of 24.41  $\text{mA cm}^{-2}$ . The EQE spectrum (Fig. S15<sup>†</sup>) of the device with HAL demonstrates an integrated current density of 24.40  $\text{mA cm}^{-2}$ , which matches the measurement results obtained from the  $J$ - $V$  characterization. To determine the reproducibility of the devices, we selected 20 devices from each of the HAL-modified devices and pristine devices, and further evaluated the statistical distribution of the collected photoelectric data, which corresponding statistical data are listed in Table S5.<sup>†</sup> As shown in Fig. 5c and Fig. S16,<sup>†</sup> HAL-modified devices exhibit superior photoelectric performance over the pristine devices, especially higher PCE and more concentrated PCE distribution, which is primarily benefiting from the improvement in  $V_{OC}$  and FF.

The improvement of photoelectric performance affirms the positive effect of the HAL buried modification, and this enhancement may be due to the effective defect passivation of HAL. To further investigate the impact of HAL on the defects of devices, we employed the trap density of states (tDOS) of different devices. As shown in Fig. 5d, the tDOS spectra contain three distinct types of trap states, corresponding to the three energy bands (labeled as band 1, band 2, and band 3).<sup>41</sup> It is evident that the overall tDOS of the HAL-modified device is lower than that of the pristine device, especially the tDOS of the HAL-modified device at shallow energy levels (band 1) is considerably lower than that of the pristine device. The traps belonging to band 1 at the low-energy region are usually



**Fig. 5** (a) Schematic device structure of PSCs with HAL-modified  $\text{SnO}_2$ . (b) The  $J$ - $V$  curves, (c) PCE histograms, and (d) tDOS spectra of PSCs with different  $\text{SnO}_2$  films. The SCLC curves for electron-only devices with (e) pristine and (f) HAL-modified  $\text{SnO}_2$ . (g) Radar graph of average photoelectric parameters of PSCs with different  $\text{SnO}_2$  films. (h) XRD stability and (i) PCE stability of PSCs with different  $\text{SnO}_2$  films in an air environment with a RH of about 35%.

associated with iodide vacancies at grain boundaries, indicating that the introduction of HAL can effectively reduce the trap states associated with vacancies at grain boundaries.<sup>9,41</sup> The reduction of the defects at the shallow energy level may be attributed to the fact that the bilateral electron-donating groups of HAL can coordinate with uncoordinated  $\text{Pb}^{2+}$  to stabilize the  $[\text{PbI}_6]^{4-}$  octahedra at the buried interface.

Moreover, we further evaluated the defect states of PSCs before and after modification using the space-charge-limited current (SCLC) technique for electron-only devices with ITO/ $\text{SnO}_2$ /PVK/PCBM/Ag structures. It is well known that the trap-filled limit voltage ( $V_{\text{TFL}}$ ) is positively related to the trap density ( $N_{\text{trap}}$ ). The smaller the  $V_{\text{TFL}}$ , the smaller the  $N_{\text{trap}}$ .<sup>12,40</sup> From the dark-state  $J$ - $V$  curves of Fig. 5e and f, the  $V_{\text{TFL}}$  values for the HAL-modified and pristine electron-only devices are 0.29 V and 0.58 V, respectively. By calculation, it is known that the  $N_{\text{trap}}$  of the HAL-modified device ( $1.64 \times 10^{15} \text{ cm}^{-3}$ ) is smaller than that of the pristine device ( $3.29 \times 10^{15} \text{ cm}^{-3}$ ), which also proves that HAL is effective for defect passivation. The Radar graph (Fig. 5g) is used to provide a more visual analysis of the average photovoltaic parameters of the HAL-modified and pristine devices, and it is clear that the HAL-modified

device has a lower  $N_{\text{trap}}$  than the pristine devices, resulting in better photovoltaic performance.

The positive role of HAL in photovoltaic performance prompted us further to investigate its influence on the stability of PSCs. As shown in Fig. 5h, the XRD technique was utilized to monitor the degradation state of the PVK films. After 35 days of storage in air at 35% RH, the (110) peak of PVK in the pristine film has significantly degraded, while the  $\text{PbI}_2$  and  $\delta$ -phase peaks have increased noticeably. In contrast, the phase transition of the PVK film with HAL-modified  $\text{SnO}_2$  has hardly increased with time. The degradation in the pristine devices can be ascribed to the large amounts of defects in the PVK films, which induce the degradation of the PVK films. In addition, we assess the long-term stability of HAL-modified and pristine devices in an air storage environment (RH of  $\sim 35\%$ , temperature of  $\sim 25^\circ \text{C}$ ). From Fig. 5i, the HAL-modified device has a lower decay rate than the pristine device during the 35-day stability test. After 35 days, the HAL-modified device still retained more than 90% of the initial PCE, whereas the pristine device only retained 79% of the initial PCE. The improved stability may be benefit from the fact that the buried modification of HAL not only effectually passivates the defects

at the SnO<sub>2</sub>/PVK interface but also stabilizes [PbI<sub>6</sub>]<sup>4-</sup> octahedra by the interactions to stabilize the perovskite structure.

### 3. Conclusion

In summary, an effective molecular bridge strategy is developed to positively modulate the buried interface in SnO<sub>2</sub>-based PSCs. Owing to the bilateral electron-donating groups (-NH<sub>2</sub>/-OH), HAL can act as a molecular bridge to passivate the defects on the SnO<sub>2</sub> surface and simultaneously improves the quality of PVK films. Modification of SnO<sub>2</sub> with HAL results in lower defect density and improved energy level matching, which is beneficial for electron extraction and transport. Furthermore, the presence of electron-donating groups on HAL-modified SnO<sub>2</sub> promotes the formation of low-defect PVK films and stabilizes the buried interface of PVK through interactions, which facilitates the suppression of interfacial defects and the reduction of voltage loss. Consequently, PSCs with HAL-modified SnO<sub>2</sub> achieve a high PCE of 23.58%, which is significantly higher than that of the pristine PSCs. Benefiting from the reduction of defects, PSCs based on HAL-modified SnO<sub>2</sub> have exhibited superior stability compared to the pristine PSCs. After storing under approximately 35% RH for 35 days, the HAL-modified SnO<sub>2</sub>-based PSCs maintain more than 90% of the initial efficiency. This work demonstrates an effective molecular bridge strategy to optimize the buried interface in n-i-p PSCs, which further improves the performance and stability of PSCs.

## 4. Experimental sections

### 4.1 Materials

SnO<sub>2</sub> hydrocolloid solution (15 wt%) was purchased from Alfa Aesar. Indium tin oxide (ITO, 14 Ω sq<sup>-1</sup>), methylamine hydrochloride (MACl, 99.5%), methylammonium iodide (MAI, 99.5%), and formamidinium iodide (FAI, 99.9%) were purchased from Advanced Election Technology Co., Ltd. Lead iodide (PbI<sub>2</sub>, 99.99%) was purchased from Xi'an Polymer Light Technology Corp. 2,2',7,7'-Tetrakis(*N,N*-di-*p*-methoxyphenylamino)-9,9'-spirobifluorene (Spiro-OMeTAD, 99.8%) was purchased from Shenzhen Feiming Science and Technology Co., Ltd. Chlorobenzene (CB, 99%), acetonitrile (ACN, 99%), dimethyl sulfoxide (DMSO, 99.7%), isopropanol (IPA, 99.5%), 4-*tert*-butylpyridine (*t*-BP, 98%), dimethylformamide (DMF, 99.8%), and bis(trifluoromethane) sulfonimide lithium salt (Li-TFSI, 99.95%) were purchased from Sigma-Aldrich. All chemicals and solvents are used directly without further purification.

### 4.2 Device fabrication

To attain a substrate with optimal cleanliness, the ITO glass substrate ought to undergo sequential sonication using IPA, acetone, and ethanol, each for 15 min. After the washing process, the glass should be dried under a stream of nitrogen gas and subsequently subjected to a combined UV-ozone and plasma treatment for 30 minutes. The SnO<sub>2</sub> hydrocolloid solu-

tion needs to be further diluted with deionized water at a volume ratio of 1:4 and then sonicated for 15 min to ensure homogeneous mixing. The diluted SnO<sub>2</sub> hydrocolloid solution was spin-coated onto the ITO substrate at 3500 rpm for 25 s and subsequently heat-treated at 150 °C for 30 min to gain ETL. Different concentrations of HAL aqueous solutions (0.5 mg mL<sup>-1</sup>, 1.0 mg mL<sup>-1</sup>, and 1.5 mg mL<sup>-1</sup>) were equipped as needed. The HAL modification solution was spin-coated onto the prepared SnO<sub>2</sub> ETLs at 4000 rpm for 25 s, followed by heat treatment at 130 °C for 25 min to gain HAL-modified SnO<sub>2</sub> ETLs. The FA<sub>0.92</sub>MA<sub>0.08</sub>PbI<sub>3</sub> light-absorbing layer was formed by a simple two-step method. Under a nitrogen atmosphere, a PbI<sub>2</sub> solution (1:9 volume ratio of DMSO to DMF) with a concentration of 1.5 M was spin-coated on ETL/ITO at 1500 rpm for 30 s, followed by annealing at 70 °C for 60 s. After that, the ammonium salt solution (90.0 mg FAI, 9.0 mg MACl, and 6.4 mg MAI dissolved in 1.0 mL IPA) was then deposited on the PbI<sub>2</sub> films at 2000 rpm for 30 s and subsequently heat-treated at 150 °C for 15 min to gain PVK films. To 1.0 mL of CB, 75.0 mg of spiro-OMeTAD, 28.8 μL of *t*-BP, and 17.5 μL of Li-TFSI solution (520.0 mg of Li-TFSI in 1.0 mL of ACN) were added sequentially and mixed homogeneously to gain the HTL solution. The HTL solution was spin-coated onto the PVK layers at 3500 rpm for 30 s to gain the HTL layers, followed by oxidation at 10% humidity for 12 hours. Finally, a silver electrode with a thickness of 80.0 nm was deposited by the ultra-vacuum thermal evaporation.

### 4.3 Film characterization

The atomic states and energy level structure changes on the surface of the samples were investigated using the photoelectron spectrometer systems (ESCLAB 250xi, Thermo Scientific) equipped with XPS and UPS. The interactions between substances were also analyzed using nuclear magnetic resonance (AVANCE III 500 MHz, Bruker). The surface properties of the samples were determined by SEM (SU-8000, Hitachi) and AFM (Multimode-8j, Bruker). The crystallinity of the samples was characterized by XRD using Cu Kα radiation (SmartLab, Rigaku). The SSPL intensity was collected using a highly sensitive fluorescence spectrophotometer (LF-1303003, Thermo Scientific). The UV-Vis absorption of the samples was measured with a spectrometer (Lambda 950UV/VIS/NIR, PerkinElmer). The physical parameter  $E_U$  can be utilized to quantify the film quality. The correlation between  $E_U$  and absorption coefficient ( $\alpha$ ) can be denoted by the following formula:<sup>38</sup>

$$\alpha = \alpha_0 \exp \frac{H\nu}{E_U}$$

where  $H\nu$  and  $\alpha_0$  correspond to the photon energy and a constant, respectively. A time-dependent single photon counting device (Pico harp 300) was applied to record the TRPL decay on an Omin-k Monochromator/Splitter (Zolix). To analyze the TRPL decay curve, a generally used method is to fit it to a bi-exponential equation, expressed as follows:<sup>24</sup>

$$f(t) = A_1 e^{-\frac{t}{\tau_1}} + A_2 e^{-\frac{t}{\tau_2}} + B$$



where  $\tau_1$  corresponds to the fast decay lifetime,  $\tau_2$  corresponds to the slow decay lifetime, while  $A_1$  and  $A_2$  are relative amplitudes. The average decay lifetime ( $\tau_{\text{avg}}$ ) can be calculated using the following equation:<sup>34</sup>

$$\tau_{\text{avg}} = \frac{A_1\tau_1^2 + A_2\tau_2^2}{A_1\tau_1 + A_2\tau_2}$$

#### 4.4 Device characterization

For the acquisition of the  $J$ - $V$  curve of the PSCs, the data were recorded using a Keithley 2420 source meter under AM 1.5G illumination in a solar simulator (Oriel Sol 3A, Newport). The light intensity of simulated sunlight was calibrated using a Si solar cell that is NREL-certified. Meanwhile, the dark  $J$ - $V$  curves were obtained by measuring the data in a dark environment using the aforementioned method. After obtaining the dark  $J$ - $V$  curves for electron-only devices, their electron defect density can be analyzed using the SCLC technique. Once the value of  $V_{\text{TFL}}$  has been determined, the value of  $N_{\text{trap}}$  can be obtained by employing the following formula:<sup>40</sup>

$$N_{\text{trap}} = \frac{2\epsilon\epsilon_0 V_{\text{TFL}}}{eL^2}$$

where  $e$  signifies the electric charge,  $L$  signifies the PVK's thickness,  $\epsilon$  signifies the PVK's relative permittivity, and  $\epsilon_0$  signifies the vacuum permittivity. Based on the  $J$ - $V$  curves data recorded under different light intensities ( $I$ ), the curves of light-intensity-dependent  $V_{\text{OC}}$  can be plotted. The correlation between  $I$  and  $V_{\text{OC}}$  adheres to the following formula:<sup>40</sup>

$$V_{\text{oc}} = \frac{nKT}{e} \ln(I) + \text{constant}$$

where  $n$  signifies the ideal coefficient,  $K$  signifies the Boltzmann constant, and  $T$  signifies the absolute temperature. TPV,  $C$ - $V$ ,  $M$ - $S$ , and EIS measurements were all performed using an electrochemical workstation (IM6e, Zahner). The built-in electric field can be employed to assess whether charge extraction is difficult or not. The relationship between capacitance ( $C$ ) and  $V_{\text{bi}}$  can be expressed by the following equation:<sup>15</sup>

$$\frac{1}{C^2} = \frac{2}{\epsilon\epsilon_0 e N_D S^2} (V_{\text{bi}} - V)$$

where  $S$  and  $N_D$  signify the active area and the doping density, respectively. The curves of tDOS can be obtained through angular frequency related capacitance, using the following formula:<sup>42</sup>

$$N_t(E_\omega) = -\frac{\omega V_{\text{bi}}}{KT e W} \frac{dC}{d\omega}$$

where  $N_t$  signifies the trap density,  $E_\omega$  signifies the demarcation energy,  $\omega$  signifies the angle frequency, and  $W$  signifies the depletion width. Based on the applied  $\omega$ ,  $E_\omega$  can be determined from the following equation:<sup>42</sup>

$$E_\omega = KT \ln\left(\frac{\omega_0}{\omega}\right)$$

where  $\omega_0$  signifies the attempt-to-escape frequency.

## Data availability

The data that support the findings of this study are available from the corresponding author upon reasonable request.

## Conflicts of interest

The authors declare no competing financial interest.

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