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High performance *n*-type bismuth telluride based alloys for midtemperature power generation

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Currently more than 60% of primary energy used in industry or life is lost as waste heat in the temperature range of 400 \sim 900 K, and more and more attention are paid to mid-temperature thermoelectric (TE) power generation. Here we combine several strategies, i.e. alloying, doping and hot deformation, to improve the TE performance of n-type bismuth telluride based TE alloys for mid-temperature power generation. Se alloying was adopted to widen the band gap and suppress intrinsic conduction at elevated temperatures. When Se atoms completely substitute the Te⁽²⁾ atoms, the crystal structure of Bi₂Te₃ based alloys tends to be more ordered, resulting in the maximum value of band gap. And the induced alloying scattering significantly reduces the lattice thermal conductivity. Then Sbl₃ donor doping was used to increase the electron concentration to further suppress the detrimental effects of bipolar conduction. Finally we applied repetitive hot deformations to further improve the figure of merit *zT* and a peak *zT* of ~1.1 was obtained at about 600 K in the 0.1 at.% Sbl₃-Bi₂Te_{1.9}Se_{1.1} alloy, which was hot-deformed three times. The results demonstrated the great potential of the alloy for application in mid-temperature TE power generation .

1. Introduction

Currently more than 60% of primary energy used in industry or life is lost as waste heat in the temperature range of 400 ~ 900 K. And due to the difficulties in heat collection, high cost and some technological issues, much of the waste heat is rarely recovered^{1,2}. Thermoelectric (TE) technology, which can realize the direct conversion between heat and electricity, is promising to tackle these issues. The conversion efficiency of a TE device can be measured by the dimensionless figure of merit $zT = \alpha^2 \sigma T/\kappa$, where α is the Seebeck coefficient, σ the electrical conductivity, *T* the operating temperature, and κ the total thermal conductivity (including the lattice contribution $\kappa_{\rm ph}$, the carrier contribution $\kappa_{\rm el}$, and the ambipolar contribution $\kappa_{\rm amb}$), respectively^{3,4}.

Reviewing the primary TE materials that can be used for power generation, such as PbTe⁵⁻⁷, half-Heusler alloys^{8,9}, filled skutterudites¹⁰, Mg₂Si_{1-x}Sn_x¹¹⁻¹³, and SnSe¹⁴, their optimal service temperatures mainly reside in the range 650–900 K. However, in the temperature range of 500K-650K, few TE materials are reported to exhibit excellent performance¹⁵⁻¹⁷. Such an application gap has imposed a pressing demand for high-efficiency TE materials operating in the range 500–650K.

Bismuth telluride (Bi_2Te_3) based bulk materials have been the best commercial TE materials with high *zT* values of ~ 1 near

refrigeration¹⁸. With the demands on waste heat recovery, Bi₂Te₃ based compounds, (Bi,Sb)₂(Te,Se)₃, have recently attracted increasing interest in low temperature power generation applications¹⁹⁻²⁷. Hu *et al* reported a high *zT* of \sim 1.3 at 380K for *p*-type polycrystalline Bi_{0.3}Sb_{1.7}Te₃ alloys²⁰ and a figure of merit zT ~ 1.2 at 445 K for n-type polycrystalline Bi2Te2.3Se0.7 alloys applied for low-temperature power generation²¹, the similar results have been achieved by other research groups in recent years²⁸. Meanwhile, Wang et al reported that the highest zT value of *n*-type zone-melted Bi₂Te_{1.5}Se_{1.5} doped with I doping reached 0.86 at 600K, whose average zT between 400K and 600K is 0.8²⁹, Liu et al also reported a systematic study on the Bi₂Te₃-Bi₂Se₃-Bi₂S₃ system, found out that *n*-type Bi_2Te_2S showed a peak $zT \sim 0.8$ at 573K and Bi₂SeS₂~0.8 at 773K fabricated by high energy ball milling followed by hot pressing³⁰. However, in the temperature range of 500-650K, the present property of *n*-type is not high enough to match p-type counterparts which had a peak $zT \approx 0.92$ at 710K reported by Hu at al³¹. So it is in crying needs to develop a kind of n-type high-performance TE materials for midtemperature power generation.

room temperature, and are mainly used in the solid state

When the Bismuth telluride (Bi_2Te_3) based compounds are applied to mid-temperature power generation, intrinsic excitation becomes the major limitation. The practical solution to suppressing bipolar conduction is to increase the band gap or majority concentration²⁰. To widen the band gap, Se alloying can be adopted. The band gap of Bi_2Te_3 is 0.15eV, and Bi_2Se_3 0.35eV, so replacing Te with Se can enlarge band gap, which have been proved by predecessors^{32,33}. In addition, Se alloying can introduce much lattice disorder to reduce the

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 κ_{ph} . On the other hand, Sbl₃ is generally used as a donor dopant to increase the majority concentration and improve the TE performance ³⁴.

In our previous work, Hu at al. demonstrated that repetitively hot-deformed $Bi_2Te_{2.3}Se_{0.7}$ alloys with engineered intrinsic point defects show the best zT in the temperature range 400-500 K^{21} . In order to further optimize the band gap of n-type Bi_2Te_3 based compounds for application in higher temperatures, in this paper, we fabricate a series of n-type hot-deformed $Bi_2Te_{3-x}Se_x$ (x=0.9-1.4) polycrystalline solid solutions with higher Se contents by a powder metallurgy process, including melting, ball milling (BM), hot pressing (HP) and subsequent hot deformation (HD). The effects of Se alloying on the suppression of intrinsic conduction and the increase of zT are systematically investigated. The peak zT is shifted above 500 K as a result of band gap widening. Further increasing the carrier concentration by Sbl₃ doping, the zT peak can be pushed to 600 K. A maximum zT~1.1 at 600 K is obtained by repetitive hot deformation, demonstrating the potential of Bi₂Te₃ based compounds with high Se contents for power generation in 500-650K.

2. Experimental

2.1 Melting

Commercial high-purity elemental chunks of 99.999% Bi, 99.999% Te, and 99.999% Se were used as raw materials. Appropriate quantities of each were weighed according to the nominal compositions of *n*-type $Bi_2Te_{3-x}Se_x$ (x = 0.9-1.4) and sealed into a quartz tube at 10^{-3} Pa. The elemental mixtures were subsequently melted at 1073K for 10 h in a rocking furnace to ensure the composition homogeneity, and then cooled in the furnace to room temperature.

2.2 Consolidation of powders, hot press and hot deformation

Ingots were ball-milled (MM 200, Retsch Gmbh, Haan, Germany) for 20 min at 20 Hz to yield fine powers. These powders were hot pressed into cylindrical shapes in a 10 mm graphite die at 673 K for 30 min under the uniaxial stress of 80 MPa, resulting in a cylinder bulk with a height of 2 mm. The initial hot-pressed bulk samples with different Se content x were named as HP-Bi₂Te_{3-x}Se_x(x=0.9-1.4). Subsequently, hot deformation (HD) was performed by repressing the HP samples in a larger graphite die with an inner diameter of 16, 20, and 25 mm each time at 823K for 30 min at the same pressure, named HD-Bi₂Te_{3-x}Se_x(x=0.9-1.4). The doping samples were called HD-ySbl₃+Bi₂Te_{3-x}Se_x (y=0.0005, 0.001, 0.0015, 0.002). This repetitive HD samples were finally named as HD_n-ySbl₃+Bi₂Te_{3-x}Se_x (*n*=1; 2; 3).

2.3 Materials characterization

The phase structure of the samples was investigated by X-ray powder diffraction (XRD) on a Rigaku D/MAX-2550P diffract meter. The in-plane electrical conductivity σ and the Seebeek coefficient α were simultaneously measured on a commercial Linseis[®] LSR-3 system. The in-plane thermal diffusivity D measurement was performed on a Netzsch LFA 457 laser flash apparatus with a Pyroceram standard using the scheme

introduced by Xie et al³⁵. The specific heat $C_{\rm P}$ was measured on the Netzsch DSC 404C and the density $\rho_{\rm D}$ was estimated by an ordinary dimension and weight measurement procedure. The in-plane thermal conductivity was then calculated using the relation $\kappa = D\rho_{\rm D}C_{\rm P}$. The Hall coefficient $R_{\rm H}$ was determined at 300 K on a Quantum Design PPMS-9T instrument using a fourprobe configuration. The carrier concentration $n_{\rm H}$ and in-plane Hall mobility $\mu_{\rm H}$ were calculated according to $n_{\rm H} = 1/eR_{\rm H}$ and $\mu_{\rm H} = \sigma R_{\rm H}$, respectively, so that the absolute error is on the order of 5%.

3. Results and discussion

3.1 Effects of Se alloying on TE properties of Bi₂Te_{3-x}Se_x

3.1.1 Microstructure evolution

The phase structure of the $HD\text{-}Bi_2Te_{3\text{-}x}Se_x$ compounds was characterized by XRD, as shown in Fig. 1. It is clear that all the diffraction patterns can be indexed to be hexagonal Bi2Te2Se (JCPDS#29-0247). All the peaks shift towards higher degrees with the increase of Se content, indicating the formation of $Bi_2Te_{3-x}Se_x$ solid solutions. To investigate the texture degree, the orientation factor F of (001)-planes for the HP and repetitive HD samples was calculated according to the Lotgering method³⁶. The F value of 0.15 suggests that certain degree preferred orientation is formed in the HP samples, according to the result of *F* values found for n-type samples by Zhao at al. and Shen at al.^{28,37}. The higher F values of $0.3^{\circ}0.5$ for HD samples characterize their stronger anisotropy. The enhanced texture in HD samples can also be observed in the SEM results, as shown in Fig.2. It is shown that after hot deformation at 823K, the grain size grows and preferential orientation is formed in the HD samples. The anisotropy has to be taken into consideration for textured samples when zTvalues are calculated, as our previous studies indicated that an *F* value of 0.31 could overestimate the zT up to 60%^{24,38}.



Fig. 1 In-plane XRD patterns of the hot-deformed $Bi_2Te_{3-x}Se_x$ bulk samples (x =0.9-1.4), taken on the hot-deformed surfaces.



Fig. 2 SEM fractographs of the cross-sections parallel to the pressing direction for the bulk samples before (a) and after hot deformation for once (b), twice (c), and three times (d).



Fig. 3 (a) Se-content dependence of electron concentration (n_H) of hot pressed, hot-deformed $Bi_2Te_{3-x}Se_x$ bulk samples, together with the undoped single crystal³⁹ and I doped zone-melt samples²⁹. (b) Se-content dependence of carrier mobility (u_H) of hot-pressed, hot-deformed $Bi_2Te_{3-x}Se_x$ bulk samples, together with I doped zone-melt samples²⁹.

3.1.2 Carrier transport

Fig. 3a shows the room temperature carrier concentration n_H of polycrystalline $Bi_2Te_{3-x}Se_x$ samples, together with undoped single crystal³⁹ and I doped zone-melted crystal²⁹. Our samples were obtained by ball milling (BM), hot pressing (HP) and hot deformation (HD). Because of donor-like effect introduced by mechanical deformation⁴⁰, all the polycrystalline samples subject to BM (HP and HD samples) show high electron concentrations than single crystal and doped zone-melted crystal. Meanwhile, HD samples show the higher electron concentrations than HP counterparts because hot deformation process also induces donor-like effect⁴¹. The n_H of polycrystalline HD and HP samples firstly falls and then rises with increasing Se content, which is the result of donor-like effect combined with point defects. The similar phenomenon has also been observed in previous report, and shows a great consistency in the Se content of 0.5-1.5^{21,42,43}. While this work focuses on mid-temperature use rather than to achieve the highest zT^{21} of n-type Bi₂Te₃ based materials. So we further refine the content gradient to find out the content with largest band gap for higher temperature use. The electron concentration of the samples results from both inherent anon vacancies and deformation-induced donor-like effect. The inherent electron concentration keeps increasing with x, meanwhile the donor-like effect becomes weaker as a result of the lack of the antisite defects concentration when x < 1.

However, the antisite defects concentration begins to increase when x>1, which causes the re-enhanced donor-like effect leading the increase of the electron concentration²¹.

Fig. 3b shows the room temperature carrier mobility u_H of polycrystalline $Bi_2Te_{3-x}Se_x$ samples, together with I doped zonemelted ingot. We can see that the u_H of polycrystalline samples are much smaller than the zone-melted one, as a result of that polycrystalline samples have weaker texture and finer grain sizes that scatter carriers more seriously. Compared with HP samples, the u_H of HD counterparts is larger, as a result of texture enhancement and grain growth, although the latter has the higher carrier concentration^{21, 22, 28}.



Fig. 4 (a) Temperature dependence of Seebeck coefficient of the hot deformed Bi_2Te_3 , se_x samples. (b) Room temperature and 500 K Seebeck coefficients of the single crystals, HP and HD Bi_2Te_3 , se_x samples. (c) Se-content dependence of band gap of the hot-deformed Bi_2Te_3 , se_x bulk samples. (d) Band gap and electron concentration dependence of temperature when the peak occurs of the hot deformed Bi_2Te_3 , se_x bulk samples.

3.1.3 Seebeck coefficients and band gap

Fig. 4a and b show that the absolute α of polycrystalline HP and HD samples first slightly rises and then falls with increasing x (0.9-1.4), corresponding to the change in n_H in Fig. 3a. Meanwhile, absolute α of HD samples are smaller than HP counterparts both at 300 K and 500 K as a result of donor-like effect induced carrier concentration increase^{21,22}.

In order to examine the effects of Se alloying on the band gap of Bi₂Te_{3-x}Se_x samples, the band gap of both HP and HD samples was estimated by the relationship $Eg = 2e\alpha_{max}T_{max}$ where α_{max} is the peak Seebeck coefficient and T_{max} the corresponding temperature²⁰. As displayed in Fig. 4c, it is shown the band gap first increases and then falls down with increasing Se content x between 0.9-1.4 and it reaches maximum when x=1.1. It is reported that the presence of a maximum and a minimum in energy gap results from the two different types of Te layers in the Bi₂Te₃ structure^{44,45}. For single crystals, band gap of Bi₂Te₃ alloyed with Se is found to enlarge with the increase of Se content x, and when x reaches around 1, the band gap goes to its maximum value.⁴⁶ Then the

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band gap begins to decrease with increasing x from 1. In current work, band gap shows a peak at around x=1.1 for the hot-deformed polycrystalline materials and the variation of band gap with Se content agrees well with the results in literature⁴⁷.

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In Fig.4d, it is worth noting that the temperature of maximum α is pushed to higher temperature with Se alloying although the band gap decreases at x>1.1, mainly due to the increased concentration of majority. As above mentioned, the majority increases with increasing Se content, which can offset the influence of the decrease of band gap at x>1.1.



Fig. 5 (a) Temperature dependence of in-plane electrical conductivity of the hot deformed $Bi_2Te_{3-x}Se_x$ samples.(b) Room temperature and 500 K electrical conductivity of the single crystals, HP and HD Bi2Te3-xSex samples. (c) Temperature dependence of power factor of the hot deformed Bi2Te3-xSex samples. (d) Room temperature and 500 K power factor of the single crystals, HP and HD Bi2Te3-xSex samples.

3.1.4 Electrical transport properties

Fig.5a and b show temperature and Se content dependence of in-plane electrical conductivity of $\text{Bi}_2\text{Te}_{3-x}\text{Se}_x$ samples, respectively. The in-plane electrical conductivity σ of HD samples decreases with measurement temperature, indicative of metal-like conduction behavior. The σ firstly decreases then increases with increasing *x*, consistent with the Se content dependence of carrier concentration. The HD samples exhibit the higher σ than HP counterparts for the studied range of Se content due to the higher μ_H and n_H of the former(Fig.3a, b).

From Fig.5c and d, it is found that HD samples have higher *PF* for all the compositions compared to HP counterparts due to the remarkably increased σ , consistent with our previous report^{21, 22, and 28}. The maximum *PF* of ~ 2.4×10⁻³ Wm⁻¹K⁻² is obtained for HD sample at *x*=1.1 at 300K and stays ~2.0 at 500K.

3.1.5 Thermal conductivity and zT

Fig. 6a, b shows that Se alloying also has an obvious influence on κ . The κ of polycrystalline Bi₂Te_{3-x}Se_x samples firstly drops and then rises with increasing x in accordance with the variation of σ (Fig.5a, b). It can be understood because the

contribution of κ_{el} is higher when σ is larger. In the same way, the HD counterparts show higher κ compared with the HP samples. The κ_{ph} of HD samples in Fig. 6c, d was estimated by subtracting the κ_{el} from the total κ , where the κ_{el} was calculated using the Wiedemann-Franz law with $L_0 = 2.0 \times 10^{-8}$ $V^2 K^{-2}$. The κ_{ph} of HD samples keep decreased with increasing Se content, and are lower than that of the HP counterparts. As previously studied^{20, 21}, multiscale microstructure effects play important roles in the reduced κ_{ph} in polycrystalline HD Bi₂Te₃. xSe_x samples. Both antisite defects and vacancies behave as phonon scattering centers and vacancies show stronger scattering on phonon because of the larger mass differences and stress field fluctuations between the occupied sites and the vacancies. During deformation process, a large number of

vacancies V_{Bi}^{-} and V_{Te}^{\bullet} (or V_{Se}^{\bullet}), and the high-density lattice distortions and dislocations are generated. All of these multiscale defects can effectively scatter phonon to reduce $\kappa_{ph}^{21,31}$.



Fig. 6 (a) Temperature dependence of in-plane thermal conductivity of the hot deformed $Bi_2Te_{3-x}Se_x$ samples, (b) Room temperature and 500 K thermal conductivity of the single crystals, HP and HD $Bi_2Te_{3-x}Se_x$ samples, (c) Temperature dependence of in-plane lattice thermal conductivity of the hot deformed $Bi_2Te_{3-x}Se_x$ samples, and (d) Room temperature and 500 K lattice thermal conductivity of the single crystals, HP and HD $Bi_2Te_{3-x}Se_x$ samples. (e) Temperature dependence of in-plane zT of the hot deformed $Bi_2Te_{3-x}Se_x$ samples. (f) Room temperature and 500 K zT of the single crystals, HP and HD $Bi_2Te_{3-x}Se_x$ samples.

The dimensionless figure of merit zT of the HD Bi₂Te_{3-x}Se_x samples is presented in Fig.6e, f with both electrical and thermal properties measured along the in-plane direction. All

the samples with x>1 exhibit a considerable enhancement in zT at higher temperatures compared to the x<1 compounds, especially above 500 K. The highest zT ~ 0.9 at 510K was obtained in Bi₂Te_{1.9}Se_{1.1} HD sample. Although the PF of doped zone-melt samples are higher than polycrystalline counterparts (Fig.5d), the final zT at 500K of HD samples are higher than that of doped zone-melt counterparts, as showed in Fig.6f, as a result of reduced κ_{ph} (Fig.6b). So the HD samples are a better choice used for mid-temperature than doped zone-melt counterparts.



Fig. 7 Temperature dependence of in-plane electrical conductivity (a), Seebeck coefficient (b), power factor (c), in-plane thermal conductivity (d), in-plane ambipolar thermal conductivity (e), and the in-plane zT (f) of the hot-deformed $ySbl_3-Bi_2Te_{1.9}Se_{1.1}$ samples.

3.2 Further optimization by Sbl₃ doping

Above discussion has demonstrated that Se alloying can widen the band gap and suppress the intrinsic conduction. The temperature for peak *zT* has been raised to 530K (*x*=1.1) from 445K (*x*=0.7) However, the ambipolar diffusion is still serious above 500 K (Fig. 6a and c). We aim at improving the service temperature to 600 K. In the following part, the bipolar conduction above 500 K is expected to be suppressed by increasing the carrier concentration through Sbl₃ doping in HD-Bi₂Te_{1.9}Se_{1.1} sample. The carrier concentration increases with the Sbl₃ doping content (Fig.S2) The electrical conductivity increases with the doping content, as shown in Fig.7a. Although the Seebeck show the certain decrease with doping, the temperature that α_{max} occurs is increased from 540 K to Journal Name

590 K at higher doping concentration in Fig 7b. From Fig 7c, the *PF* of doped samples are higher above about 420K.

In view of the thermal conductivity, it is clear that the minimum values for each samples has been pushed to higher temperatures with increasing Sbl₃ content. We evaluated the ambipolar contribution κ_{amb} for the alloys of different Sbl₃ content using the method introduced in Ref. 20. It is shown that the contribution of κ_{amb} to the total thermal conductivity is quite large when the carrier concentration is low, especially at high temperatures. With increasing carrier concentration, the ambipolar effects have been greatly reduced. κ_{amb} of pristine sample takes over 50 % of κ at 600 K while the sample y=0.001 only holds 8.3%. Hence, the suppression of ambipolar diffusion significantly contributes to the high *zT* of ~0.95 at 580 K for HD-0.001Sbl₃+Bi₂Te_{1.9}Se_{1.1} sample, compared to HD-Bi₂Te_{1.9}Se_{1.1} counterpart of 0.9 at 510 K.



Fig. 8 (a) Temperature dependence of in-plane zT of the repetitive hot-deformed 0.001Sbl₃-Bi₂Te_{1.9}Se_{1.1}. (b) Thermal transport properties for the Bi₂Te₂Se_{1.1} bulk samples before and after different hot-deformation numbers: κ_R/κ_A and *F* samples.



Fig. 9 (a) zT advances of the mid-temperature n-type $Bi_2Te_{3-x}Se_x$ -based alloys, (b) Temperature dependence of zT values for the re-prepared HD3-0.001Sbl₃- $Bi_2Te_{1.9}Se_{1.1}$ bulk samples, (c) Temperature dependence of zT values of the n-type TE materials applied in mid-temperature ranges, (d) Temperature dependence of zT values of the n-type bismuth telluride compounds applied in different temperature ranges.

3.3 Enhancing the zT by repetitive hot deformation

It has been reported that repetitive hot deformation can effectively improve the TE performance for *n*-type bismuth

telluride based alloys through enhanced texture and increased carrier concentration^{21,22}. In this work Fig 8a shows that the optimal 0.001Sbl₃+Bi₂Te_{1.9}Se_{1.1} sample hot deformed triple has highest zT values of ~1.1 at 590K. Fig 8b shows that the texture degree was greatly enhanced by repetitive hot deformations, resulting that the F and (κ_R/κ_A) value are getting larger with the increase HD times, which must be taken into consideration in measuring TE parameters for the zT calculation to avoid a overestimation, as demonstrated in previous work^{28,37}.

Fig. 9a shows the zT improvement of *n*-type Bi₂Te_{3-x}Se_x-based alloys in this work. The zT values have been successfully increased from 0.45 to 1.1 and the optimal service temperature has been improved from 480 K to 590 K through Se alloying, optimizing doping and repetitive hot deformation.

The HD3-0.001Sbl₃+Bi₂Te_{1.9}Se_{1.1} bulk samples have been reprepared triple by the same procedure. The high zT values have been reproduced; meanwhile all the samples have an average value of zT ~0.9 in the temperature range from 300 K to 650 K. (Fig. 9b). Compared with other n-type medtemperature TE materials such as Mg₂Si¹¹, Bi₂Te_{1.5}Se_{1.5}²⁹, $Bi_2Te_2S,\ Bi_2SeS_2^{\ 30},\ our \ work \ shows \ the \ best \ TE \ performance$ during whole temperature range, especially in the range of 500-650 K (Fig. 9c).

In previous work, our group applied hot deformation directly to commercialized $Bi_2Te_{2.79}Se_{0.21}$ ingots, whose Se content is low and suitable for use in room temperature range (300-400 K), and $zT \sim 1.2$ at 350K was achieved⁴⁸. Meanwhile, we also further increased Se content and demonstrated that repetitive hot-deformed $Bi_2Te_{2.3}Se_{0.7}$ alloys show the best zT in the temperature range 400–500 K^{21} . In this paper, we increase Se content to 1.1 and fabricate high zΤ HD3-0.001Sbl₃+Bi₂Te_{1.9}Se_{1.1} material applied in the temperature range of 500-650 K. These results demonstrate the promising application of bismuth telluride alloys for power generation (Fig. 9d).

Conclusions

In summary, we successfully shifted the maximum zT values of n-type bismuth telluride-based alloys to relatively high temperatures (approximately 600K) by Se-alloying and Sbl₃ doping. The detrimental effects of minority carriers on the Seebeck coefficient and thermal conductivity were suppressed as a result of increases in both the band gap and electron concentration. At last, high-performance bismuth telluridebased alloys were fabricated utilizing the repetitive HD method. The Se alloying gives out a bigger band gap among all samples to suppress detrimental intrinsic conduction. A donor-like effect and Sbl₃ doping increase the electron concentration simultaneous, inhabiting the ambipolar effects for higher temperature use. Aside from this, the lattice thermal conductivity was considerably reduced by the presence of recrystallization-induced in situ nanostructures and high density lattice defects. As a consequence of these factors, the HD3-0.001Sbl₃+Bi₂Te_{1.9}Se_{1.1} materials, showed a maximum zTof ~1.1 at approximately 600K and the largest average zT_{av} of

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0.9 in the range of 300–650K. These data indicate significant promise for these materials in mid-temperature power generation. As these results proved extremely reproducible, this is a promising approach for the mass production of highperformance TE materials.

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Notes and references

1 Kim SI, Lee KH, Mun HA, Kim HS, Hwang SW, Roh JW, Science 2015;348:109-14.

2 Poudel B, Hao Q, Ma Y, Lan Y, Minnich A, Yu B, Science 2008;320:634-8.

3 Tritt TM, Science 1999;283:804-5.

4 Dresselhaus MS, Chen G, Tang MY, Yang RG, Lee H, Wang DZ, Advanced Materials 2007;19:1043-53.

5 Pei Y, Lensch-Falk J, Toberer ES, Medlin DL, Snyder GJ, Advanced Functional Materials 2011;21:241-9.

6 Biswas K, He J, Blum ID, Wu Cl, Hogan TP, Seidman DN, Nature 2012;489:414-8.

7 Hsu KF, Loo S, Guo F, Chen W, Dyck JS, Uher C, Science 2004;303:818-21.

8 Xie H, Wang H, Pei Y, Fu C, Liu X, Snyder GJ, Advanced Functional Materials 2013;23:5123-30.

9 Yu C, Zhu T-J, Shi R-Z, Zhang Y, Zhao X-B, He J, Acta Materialia 2009;57:2757-64.

10 Shi X, Yang J, Salvador JR, Chi M, Cho JY, Wang H, Journal of the American Chemical Society 2011;133:7837-46.

11 Du Z, Zhu T, Chen Y, He J, Gao H, Jiang G, Journal of Materials Chemistry 2012;22:6838.

12 Liu W, Tan X, Yin K, Liu H, Tang X, Shi J, Physical Review Letters 2012;108.

13 Liu X, Zhu T, Wang H, Hu L, Xie H, Jiang G, Advanced Energy Materials 2013;3:1238-44.

14 Zhao L-D, Lo S-H, Zhang Y, Sun H, Tan G, Uher C, Nature 2014;508:373-7.

15 Hsu C-T, Huang G-Y, Chu H-S, Yu B, Yao D-J, Applied Energy 2011;88:1291-7.

- 16 Wu C, Applied Thermal Engineering 1996;16:63-9.
- 17 T. D. Barr and F. Dahlen, J. Geophys. Res., 1989, 94, 3923.

18 DiSalvo FJ, Science 1999;285:703-6.

19 Hu L, Gao H, Liu X, Xie H, Shen J, Zhu T, Journal of Materials Chemistry 2012;22:16484.

20 Hu L-P, Zhu T-J, Wang Y-G, Xie H-H, Xu Z-J, Zhao X-B, NPG Asia Materials 2014;6:e88.

21 Hu L, Zhu T, Liu X, Zhao X, Advanced Functional Materials 2014;24:5211-8.

Article 22 Hu LP, Liu XH, Xie HH, Shen JJ, Zhu TJ, Zhao XB, Acta Materialia 2012:60:4431-7. 23 Zhu T, Xu Z, He J, Shen J, Zhu S, Hu L, Journal of Materials Chemistry A 2013;1:11589. 24 Yan X, Poudel B, Ma Y, Liu WS, Joshi G, Wang H, Nano letters 2010;10:3373-8. 25 Shen J-J, Zhu T-J, Zhao X-B, Zhang S-N, Yang S-H, Yin Z-Z., Energy & Environmental Science 2010;3:1519-23. 26 Li J, Tan Q, Li J-F, Liu D-W, Li F, Li Z-Y, Advanced Functional Materials 2013;23:4317-23. 27 Mehta RJ, Zhang Y, Zhu H, Parker DS, Belley M, Singh DJ, Nano letters 2012;12:4523-9. 28 Zhao LD, Zhang BP, Li JF, Zhang HL, Liu WS, Solid State Sciences 2008;10:651-8. 29 Wang S, Tan G, Xie W, Zheng G, Li H, Yang J, Journal of Materials Chemistry 2012;22:20943. 30 Liu W, Lukas KC, McEnaney K, Lee S, Zhang Q, Opeil CP, Energy Environ Sci 2013;6:552-60. 31 Hu LP, Zhu TJ, Yue XQ, Liu XH, Wang YG, Xu ZJ, Acta Materialia 2015;85:270-8. 32 Imamuddin M, Dupre A, physica status solidi (a) 1972;10:415-24. 33 Teramoto I, Takayanagi S, J Phys Chem Solids 1961;19:124-9. 34 Delves R, Bowley A, Hazelden D, Goldsmid H, Proceedings of the Physical Society 1961;78:838. 35 Xie W, He J, Zhu S, Holgate T, Wang S, Tang X, Journal of Materials Research 2011;26:1791-9. 36 Lotgering FK, Journal of Inorganic and Nuclear Chemistry 1959;9:113-23. 37 Shen JJ, Hu LP, Zhu TJ, Zhao XB, Applied Physics Letters 2011;99:124102. 38 Ben-Yehuda O, Shuker R, Gelbstein Y, Dashevsky Z, Dariel MP, Journal of Applied Physics 2007;101:113707. 39 Birkholz U, Astrophysik Physik Und Physikalische Chemie 1958;13:780-92. 40 Jiang J, Chen L, Bai S, Yao Q, Wang Q, Scripta materialia 2005;52:347-51. 41 Navrátil J, Starý Z, Plecháček T, Materials Research Bulletin 1996;31:1559-66. 42 West D, Sun YY, Wang H, Bang J, Zhang SB, Physical Review B 2012;86. 43 Stary Z, Horak J, Stordeur M, Stolzer M, J Phys Chem Solids 1988:49:29-34. 44 I. G. Austin and A. Sheard, J. Electron. Control 1957;3, 236. 45 M. J. Smith, E. S. Kirk, and C. W. Spencer, J. Appl. Phys. 1960; 31, 1504 . 46 [46] Miller GR, Li C-Y, Spencer CW, Journal of Applied Physics 1963;34:1398. 47 J. Black, E. Conwell, L. Seigle, and C. W. Spencer, J. Phys. Chem. Solids 1957;2, 240. 48 Drabble JR, Goodman CH, J Phys Chem Solids 1958;5:142-4.

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Title: "High performance n-type bismuth telluride based alloys for mid-temperature power generation" by Zhenglong Tang, Lipeng Hu, and Tiejun Zhu, Xiaohua Liu, Xinbing Zhao.



We combine Se alloying, SbI₃ doping and repeated hot deformation to obtain high-performance n-type Bi₂Te₃ based med-temperature thermoelectric materials for power generation.