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## Advancing high-throughput combinatorial aging studies of hybrid perovskite thin films *via* precise automated characterization methods and machine learning assisted analysis†

Alexander Wieczorek, <sup>a</sup> Austin G. Kuba, <sup>b</sup> Jan Sommerhäuser, <sup>a</sup> Luis Nicklaus Caceres, <sup>a</sup> Christian M. Wolff <sup>\*b</sup> and Sebastian Siol <sup>\*a</sup>

To optimize material stability, automated high-throughput workflows are of increasing interest. However, many of those workflows either employ synthesis techniques not suitable for large-area depositions or are carried out in ambient conditions, which limits the transferability of the results. While combinatorial approaches based on vapour-based depositions are inherently scalable, their potential for controlled stability assessments has yet to be exploited. Based on  $\text{MAPbI}_3$  thin films as a prototypical system, we demonstrate a combinatorial inert-gas workflow to study intrinsic materials degradation, closely resembling conditions in encapsulated devices. Specifically, we probe the stability of  $\text{MAPbI}_3$  thin films with varying residual  $\text{PbI}_2$  content. A comprehensive set of automated characterization techniques is used to investigate the structure and phase constitution of pristine and aged thin films. A custom-designed *in situ* UV-Vis aging setup is used for real-time photospectroscopy measurements of the material libraries under relevant aging conditions, such as heat or light-bias exposure. These measurements are used to gain insights into the degradation kinetics, which can be linked to intrinsic degradation processes such as autocatalytic decomposition. Despite scattering effects, which complicate the conventional interpretation of *in situ* UV-Vis results, we demonstrate how a machine learning model trained on the comprehensive characterization data before and after the aging process can link changes in the optical spectra to phase changes during aging. Consequently, this approach does not only enable semi-quantitative comparisons of material stability but also provides detailed insights into the underlying degradation processes which are otherwise mostly reported for investigations on single samples.

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## 1 Introduction

The advancement of energy conversion and sensory technologies relies heavily on developing novel semiconducting materials with innovative functional properties. A class of hybrid organic-inorganic semiconductors that garnered attention for high-performance applications in light-emission,<sup>1</sup> photovoltaics,<sup>2</sup> and X-ray detection<sup>3</sup> are metal halide perovskites (MHP).

<sup>a</sup>Laboratory for Surface Science and Coating Technologies, Empa – Swiss Federal Laboratories for Materials Science and Technology, Switzerland. E-mail: sebastian.siol@empa.ch

<sup>b</sup>Institute of Electrical and Microengineering (IEM), Photovoltaic and Thin-Film Electronics Laboratory, EPFL –École Polytechnique Fédérale de Lausanne, Switzerland. E-mail: christian.wolff@epfl.ch

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MHPs crystallize in the perovskite structure and follow the general formula  $\text{ABX}_3$ , where typically  $\text{A} = \text{Cs}^+$ , methylammonium ( $\text{MA}^+$ ), formamidinium ( $\text{FA}^+$ ),  $\text{B} = \text{Sn}^{2+}$ ,  $\text{Pb}^{2+}$ , and  $\text{X} = \text{I}^-, \text{Br}^-, \text{Cl}^-$ .

A key challenge in the development of hybrid semiconductors is their limited long-term stability under operating conditions.<sup>4,5</sup> Although multiple concepts exist to predict the phase stability of hybrid perovskites, *e.g.* through the long-established Goldschmidt tolerance factor,<sup>6</sup> and octahedral tilting due to the A-site cation,<sup>7–10</sup> decomposition may occur *via* a multitude of pathways. These can include thermal instability,<sup>11</sup> photoinduced degradation,<sup>12</sup> as well as oxygen- and moisture-induced degradation.<sup>13</sup> Consequently, experimental assessments are necessary to measure and understand their degradation behaviour and unlock their true potential for the aforementioned applications on the commercial scale.

Nonetheless, performing aging studies can be especially challenging for multiple reasons:



(1) Aging is intrinsically time-intensive, even under accelerated application-specific aging conditions.<sup>14,15</sup>

(2) Assessing aging using only one material characterization method typically yields limited insights into the complex changes in the material chemistry which may result in wrong interpretations.

(3) Aging may result from multiple intrinsic and extrinsic factors at once, rendering the assignment to one specific factor difficult. At the same time, different extrinsic factors can limit the reproducibility across laboratories.

(4) The durability of a material can be highly dependent on the deposition technique.<sup>16</sup> Consequently, the results from stability screening may differ significantly for different synthesis routes.

So far, mainly the first issue has been targeted through various approaches. These include parallelized high-throughput experiments (HTE) on photovoltaic cells under accelerated aging conditions<sup>17</sup> as well as automated materials acceleration platforms for screening the optical stability of thin films.<sup>18</sup> While these strategies allow one to significantly reduce the time effort to probe material stability, there is a critical demand to address the remaining challenges.

A common approach for experimental thin-film screening using scalable deposition methods are vapour-based combinatorial techniques.<sup>19</sup> Using such techniques, material libraries can be prepared, which exhibit process parameter gradients or composition gradients across a single substrate. This way large areas of the synthesis phase space are accessed in one deposition run. Driven by advancements in experimental infrastructure and data processing,<sup>20</sup> automated characterization techniques have greatly improved in recent years,<sup>21</sup> yielding more thorough insights into the respective material properties across libraries. Although first reports have been made to use combinatorial approaches for the optimization of devices,<sup>22</sup> its potential for stability screening with respect to compositional stability has yet to be exploited.

In this work, we report a comprehensive workflow for combinatorial thin-film aging studies under highly controlled environmental conditions. As extrinsic factors are less relevant for the commercialization of MHPs due to modern encapsulation techniques,<sup>23,24</sup> we aim to exclude those extrinsic aging factors in our screening protocol. This is enabled through a comprehensive inert-gas workflow, in which the samples remain in nitrogen atmosphere during XRF, XRD and UV-Vis characterization, as well as the transfer between instruments. We note that aging under inert-gas to investigate intrinsic aging was recently recommended within a consensus statement for the investigation of perovskite optoelectronics.<sup>15</sup>

We demonstrate our workflow based on  $\text{MAPbI}_3$  as a prototypical system. Since its degradation mechanisms have already been investigated in detail, it serves as an ideal reference material. Furthermore, we employ a scalable deposition technique based on a two-step process that holds promise for the conformal deposition of these semiconductors on large areas. Using a novel custom-designed platform for combinatorial *in situ* UV-Vis aging experiments, we investigate how conversion gradients introduced during this process affect the resulting optical stability of the thin films. While initial analysis of the

UV-Vis data cannot be directly linked to the aging trends revealed from XRD measurements of the pristine and aged thin films, machine-learning assisted analysis of the UV-Vis spectra allows us to reveal common aging kinetics for the herein investigated thin films.

Overall, the results of this study highlight how a comprehensive stability-screening workflow performed under inert-gas atmosphere can generate valuable insights into the materials aging behaviour under application-relevant degradation conditions. While our investigation is focused on MHPs, the framework may also be used to accelerate the development of other promising material systems as stability assessments are crucial for a wide range of novel materials.

## 2 Experimental section

### 2.1 Materials

Lead iodide ( $\text{PbI}_2$ , 99.999%) beads were purchased from Sigma Aldrich. Methylammonium iodide (MAI, >98%) was purchased from Greatcell Solar Ltd. Isopropanol (iPrOH, 99.5%, anhydrous) was purchased from Sigma-Aldrich. All materials were used without further purification.

### 2.2 Sample preparation

The samples were deposited using a two-step technique which consists of the deposition of an inorganic  $\text{BX}_2$  template (here:  $\text{PbI}_2$ ) followed by coating of the organic  $\text{AX}$  salt (here: MAI) as further explained below.

**2.2.1 Inorganic template preparation.** The inorganic templates were prepared *via* thermal evaporation in a Lesker Mini SPECTROS evaporation system. 1.1 mm thick borosilicate glass (EXG) substrates of  $50 \times 50 \text{ mm}^2$  size were ultrasonically cleaned in acetone, and ethanol, followed by drying in a flow of  $\text{N}_2$ . The substrates were then transferred into the evaporation system, substrate masks were used to generate a patterned template. With closed shutters, the  $\text{PbI}_2$  source was first heated until a stabilized rate of  $1 \text{ \AA s}^{-1}$  was acquired. Then, the shutters were opened and  $\text{PbI}_2$  was deposited at  $1 \text{ \AA s}^{-1}$  until a nominal film thickness of 200 nm was determined by a quartz crystal monitor (QCM). The substrate was intentionally not rotated during this process to induce thickness gradients resulting in conversion gradients from the  $\text{PbI}_2$  to the resulting perovskite.

**2.2.2 Deposition of organic precursors.** Previously prepared inorganic templates were transferred to another glovebox without breaking inert-gas conditions and all further steps were performed within this glovebox. The MAI solution was first prepared by dissolving 300 mg within 5 ml iPrOH. The mixture was first vortexed for 1 h and then filtered using a  $45 \mu\text{m}$  polytetrafluoroethylene (PTFE) syringe filter. All used solutions and solvents were then pre-heated at  $70^\circ\text{C}$  for 20 min. As a pre-wetting step, 1 ml preheated iPrOH was first deposited onto the inorganic templates and left resting for 1 min and finally spun at 1300 rpm for 10 s. 1 ml of the preheated MAI solution was then deposited and left resting on the template for 1 min. To promote the solution evaporation, the sample was then spun at 1300 rpm for 30 s. Films were then annealed at  $100^\circ\text{C}$  for



10 min. Residual organic salts were then removed through the dynamic addition of 500  $\mu$ l iPrOH at 2500 rpm. To reduce the amount of surface defects on the  $\text{MAPbI}_3$  surface without affecting the crystal lattice, a MAI-based surface treatment was applied as reported elsewhere.<sup>25</sup>

### 2.3 Characterization techniques

**2.3.1 X-Ray diffractometry (XRD).** XRD measurements were performed in a Bruker D8 Discovery diffractometer using  $\text{Cu K}\alpha$  radiation in Bragg-Brentano geometry. Within a glovebox, the sample was introduced in a homemade gas-tight XRD sample chamber equipped with a 1 mm thick polyether ether ketone (PEEK) dome. The dome material was chosen due to its high X-ray transmission properties compared to other widely available polymers.<sup>26</sup> Spectra for all samples on the library were acquired by moving the dome within the diffractometer using an  $x$ - $y$  stage. For each sample spectra, a reference spectra recorded using an empty dome was subtracted. Batch analysis of the XRD peaks for fitting was further performed using the open source package COMBIgor<sup>27</sup> within an Igor Pro 9 environment.  $2\theta$  scans were performed within the 5–50° range.  $\chi$  scans were performed by first setting the  $2\theta$  values to the maximum positions of the  $\text{PbI}_2$  (001) and  $\text{MAPbI}_3$  (110) features at 12.65° and 14.1°, respectively. The  $\chi$  angles were then scanned in the –5–90° range during the measurement.

For the semi-quantitative analysis, reference crystal structures were first retrieved from the Inorganic Crystal Structure Database (ICSD).<sup>28</sup> Using ICSD-68819 (ICSD release 2023.2) for  $\text{PbI}_2$  and ICSD-124919 (ICSD release 2023.2) for  $\text{MAPbI}_3$ , structure factors  $|F|$  were extracted with VESTA 3.5.8 and divided by their respective number of formula units  $z$ . The resulting values were  $|F_{\text{PbI}_2(001)}| = 63$ , and  $|F_{\text{MAPbI}_3(110)}| = 102$ . Furthermore, the multiplicities  $p$  of the reflection planes were  $p_{\text{PbI}_2(001)} = 2$ , and  $p_{\text{MAPbI}_3(110)} = 4$ .

To account for intensity variations resulting from different orientations of the sample specimen relative to the incident beam, the Lorentz factor  $L$  is further considered in the semi-quantitative analysis of the XRD data. It is defined as:<sup>29</sup>

$$L = \frac{1 + \cos^2(2\theta)}{\sin^2(\theta) \times \cos(\theta)}, \quad (1)$$

where  $\theta$  denotes the Bragg angle of the observed reflection.

**2.3.2 X-Ray fluorescence (XRF) spectroscopy.** XRF measurements were performed in a Fischerscope XDV-SDD X-ray fluorescence (XRF) system equipped with an Rh X-ray source to determine the atomic ratios of Cs, Pb, I and Br. Spectra for all samples on the library were acquired by moving the sample underneath the device optics using the integrated stage motors. The thickness of each thin film was further estimated from the attenuation of signals related to the sample substrate. To maintain all samples under inert-gas conditions, 3D-printed containers covered with a thin polyethylene (PE) foil were used. Measuring through this PE foil notably did not significantly affect the measured thickness due to the low sensitivity of XRF for 2nd row elements and minimal attenuation of the underlying signal.

**2.3.3 In situ UV-Vis spectroscopy.** *In situ* UV-Vis measurements were performed in a homemade setup. The combinatorial library, contained within the climate chamber, was mounted on a biaxial stepper motor stage (Thorlabs LTS300). During measurements, the sample was illuminated using an Deuterium-Halogen light source (Ocean Optics DH-2000-BAL). Two charge-coupled device (CCD) detectors (Ocean Optics HR4Pro XR-ES) were used to record the resulting transmission and reflection spectra in the 200 nm to 1100 nm range. A 50 : 50 polka dot beam splitter (Thorlabs BPD5254-G01) mounted at an 45°-angle was used to separate the reflected light from the incoming light beam. Polarization-resistant multi-mode fibres (Thorlabs M112L02) with a core diameter of 200  $\mu\text{m}$  were used to route the optical beams. All fibres were terminated with FOV-adjustable collimator lenses (Ocean Insight 74-ACR), mounted on biaxial kinematic mounts for aligning the beams. Two plano-convex focusing lenses (Thorlabs LA4052) were used to collimate light being reflected and transmitted through the sample. All components mentioned above are specified for use within the optical range from 200 nm to 1100 nm or better, allowing measurements across the entire UV-VIS spectrum. The beam path for the reflection measurement first passes straight through the beam splitter and a focusing lens, is reflected at the sample back through the lens and is then reflected by the beam splitter at an 90°-angle before entering the reflection spectrometer. The beam path for the transmission measurement passes straight through the beam splitter, the focusing lens, the sample and a second lens before entering the transmission spectrometer on the other side of the sample (Fig. S1†). The resulting measurement spot has diameter of  $<2\text{ mm}^2$ . Spectra for each sample on the library were acquired by moving the climate chamber using the aforementioned  $xy$  stage underneath the fixed optics. This ensured the stability of the optical setup while scanning the materials library. The entire measurement setup is enclosed in a box made from black anodized aluminium to block stray light.

Based on the obtained reflectance and transmission spectra, the reflection-corrected transmission  $T'$  was calculated according to:

$$T' = \frac{T}{1 - R}, \quad (2)$$

where  $T$  denotes the sample transmission and  $R$  denotes the measured sample reflection (Fig. S2†).

Based on this value, the absorption coefficient  $\alpha$  of the thin-film layer was calculated:

$$\alpha = \frac{-\ln\left(\frac{T'_{\text{Substrate+Film}}}{T'_{\text{Substrate}}}\right)}{d}, \quad (3)$$

where  $d$  denotes the thin-film thickness estimated from XRF measurements.

To ensure inert-gas conditions are maintained during the measurement, a homemade climate chamber was used. It was continuously held at an overpressure of approximately 20 mbar using a continuous  $\text{N}_2$  inflow and a one-way valve with a corresponding cracking pressure. To enable aging under stress-light



conditions, a  $70 \times 70 \text{ mm}^2$  white-light LED chip was used in combination with a light diffuser. Based on the characterized emission spectrum, the light intensity was calibrated to  $1 \text{ kW m}^{-2}$  using a Si photodiode with a known external quantum yield (EQY) (Fig. S3†). The spatial uniformity of our white light-stressing conditions were characterized to be better than  $\pm 1.4\%$  across the area of our combinatorial libraries in combination with a high temporal stability (Fig. S4†). The entire climate chamber body is heated to  $85^\circ\text{C}$  using two heating cartridges connected to a PID Controller (Red Lion PXU11A50). From infrared camera measurements of the thin films, a heating uniformity of  $\pm 1 \text{ K}$  was confirmed (Fig. S5†). As a result of these highly accurate aging conditions across the library area, reliable combinatorial aging is enabled. Overall, the aging conditions reflect the ISOS-LC3I standard.<sup>15</sup> Homemade code implemented in LabView 2021 was used to automate the measurements.

**2.3.4 Machine learning analysis.** The machine learning model was implemented using the open-source python packages XGBoost<sup>30</sup> and Scikit-Learn<sup>31</sup> within a Python 3.7.1 environment. The data set was first shuffled and split it into training and testing data sets with a ratio of 8 : 2. Then, hyperparameter optimization was performed using *k*-fold cross validation using root-mean-squared-error (RMSE) as the scoring metric. To prevent overfitting on our experimental data set, we limited the optimizable hyperparameter space to a low maximum tree-depth of 2–4, a low learning rate of  $10^{-3}$ – $10^{-1}$ , and a high amount of estimators of 200–1000. Furthermore, the regularization parameters Gamma, Alpha and Lambda have been optimized. A gamma value of 0.7 was found to be optimal, in agreement with the low maximum tree-depth. Lastly, the performance of the model was assessed through prediction on the earlier split test set using  $R^2$  and RMSE as metrics.

The code and data used for the training and evaluation of the machine-learning model is available at GitHub ([https://github.com/Empa-CT/PbI2\\_Residuals\\_From\\_UVVis](https://github.com/Empa-CT/PbI2_Residuals_From_UVVis)).

## 3 Results and discussion

### 3.1 Combinatorial workflow

Fig. 1 illustrates the combinatorial workflow used in this study. Central targets of this approach were the seamless integration of existing scalable deposition processes and the exclusion of external degradation effects during aging and characterization. To address the first aspect, we employ a two-step deposition process consisting of thermal evaporation of an inorganic  $\text{BX}_2$  template ( $\text{PbI}_2$  in this case), followed by solution processing and annealing of the organic salts AX (MAI in this case) with the goal to form the fully converted  $\text{ABX}_3$  perovskite phase. This approach was previously shown to be ideal for conformal coatings on micro-structured substrates.<sup>32</sup> Combined with its high potential for upscaling using slot-die coating and mitigation of toxic solvents such as DMF,<sup>33,34</sup> it is highly promising for an industrial scale deposition process, especially for other MHPs suitable for integration with Si-based optoelectronics (e.g. perovskite-Si tandem solar cells).<sup>35</sup> During this process, underconversion (*i.e.* incomplete formation of the MHP from the inorganic  $\text{BX}_2$  template) or overconversion (*i.e.* the

formation of an AX-rich phase) may easily occur.<sup>36,37</sup> Using an evaporation step for the inorganic precursor allows us to achieve thickness gradients on the substrate.<sup>21</sup> Different amounts of the inorganic template, in turn, result in areas of higher or lower conversion in the resulting MHP materials library.

In high-throughput aging experiments the effect of these conversion gradients on the materials stability is evaluated. To improve the validity of results, the libraries are patterned during synthesis, thereby preventing lateral aging effects across different samples (e.g. *via* lateral diffusion or local  $\text{PbI}_2$  outgassing).<sup>38,39</sup> Most importantly, to exclude external degradation the entire characterization process is performed under an inert-gas atmosphere. This is crucial to screen for intrinsic MHP stability,<sup>15</sup> which is important for implementation in encapsulated devices, and is typically not considered for stability-screening workflows. To this end, we developed custom solutions to perform X-ray and optical-based methods under inert-gas: first, a comprehensive thin-film characterization is performed based on XRD, XRF and UV-vis measurements, allowing to deduce structural information, thin-film thickness and the optical properties of the pristine thin films. Then, *in situ* UV-vis measurements are performed under highly-controlled aging conditions ( $85^\circ\text{C}$ ,  $1 \text{ kW m}^{-2}$  white light,  $\text{N}_2$  gas) to promote the intrinsic aging of the thin films while rapidly probing changes in the optical properties. Lastly, the comprehensive sample characterization initially performed on the pristine library is repeated on the aged library.

Based on these comprehensive experimental data sets of pristine and aged thin films, we are able to leverage data-driven approaches to enable a better understanding of the observed optical changes with respect to phase changes. This knowledge can then be used to provide a better analytical model for the *in situ* UV-Vis spectra and understand the underlying aging kinetics.

### 3.2 Effect of conversion ratio on pristine thin-film properties

Based on the characterization of the pristine thin films, the effect of varying  $\text{PbI}_2$  residuals across the library on the optical properties can be determined.

First, the amount of residual  $\text{PbI}_2$  within the MHP thin film was approximated using XRD measurements. This is a common approach where typically the most prominent peak intensities of  $\text{PbI}_2$  and the corresponding MHP are compared (Fig. 2a).<sup>40–43</sup> We further confirmed the absence of over-converted MHP phases through the lack of features at shallower  $2\Theta$  angles (Fig. S6†).

Thin films produced by PVD often exhibit pronounced texture. Especially in oblique-angle vapour-deposition setups, grains tend to grow towards the source.<sup>44,45</sup> Since the commonly employed Bragg–Brentano XRD measurements only probe features that align with the diffraction plane, preferential orientation can result in intensity variations, which are not correlated with the phase constitution of the material. To account for these effects, we performed azimuthal scans on  $\text{PbI}_2(001)$  and  $\text{MAPbI}_3(110)$  features (Fig. 2b) to probe the preferred orientation of each phase and account for geometry-specific intensity variations.<sup>46</sup> Compared to  $\text{PbI}_2$  which exhibited a clear out-of-plane orientation, the resulting  $\text{MAPbI}_3$  thin films were less textured. These findings are in agreement with

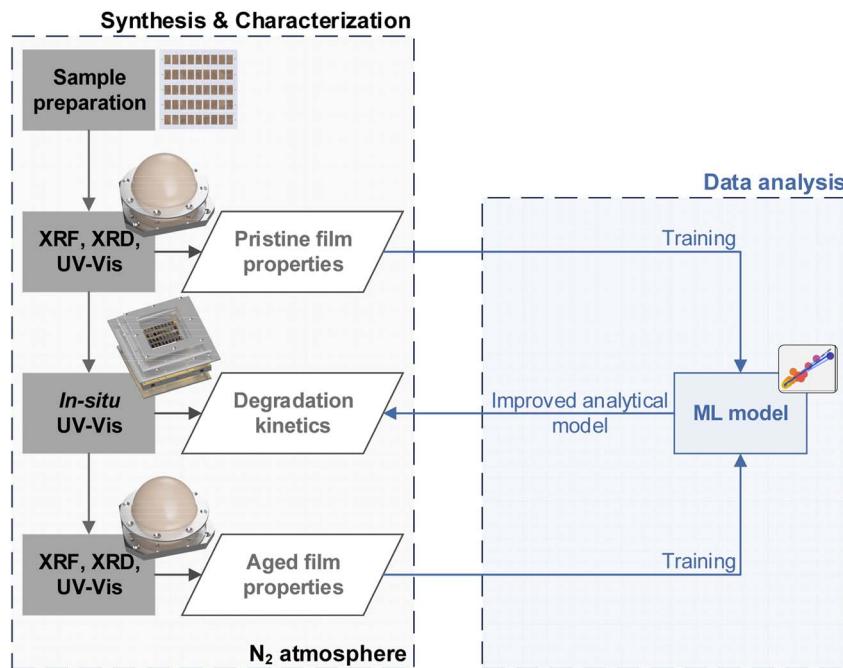


Fig. 1 Inert-gas synthesis & characterization workflow for combinatorial perovskite libraries. The rendered graphics within the inert-gas workflow depict a combinatorial perovskite library, the inert-gas XRD Dome, the climate chamber for UV-Vis aging, and repeatedly the inert-gas XRD Dome (top to bottom). The combined data sets obtained from the characterization of the pristine and aged thin films are then used to train a machine learning (ML) model for the analysis of the optical spectra obtained during the *in situ* UV-Vis aging characterization.

previous results from synchrotron-based Grazing-Incidence Wide-Angle X-ray Scattering (GIWAXS) measurements on two-step deposited MHPs grown from thermally evaporated templates.<sup>47,48</sup> As a result of the differences in texture between the two phases, determination of the phase fraction from the  $2\theta$  scans alone would result in an underestimation of the MHP phase ratio compared to  $\text{PbI}_2$ . To account for this effect, the normalized peak area  $A_{\text{norm}}$  is first calculated according to:

$$A_{\text{norm}} = A_{2\theta} \times \frac{A_\chi}{I_{\chi=0^\circ}} \quad (4)$$

where  $A_{2\theta}$  denotes the peak area obtained from the  $2\theta$  measurement,  $A_\chi$  denotes the area under curve for the  $\chi$  scan performed at the given reflection and  $I_{\chi=0^\circ}$  denotes the absolute intensity at  $\chi = 0^\circ$ . The latter factor effectively resembles the previously depicted normalized  $\chi$  scans (inset in Fig. 2b).

For measurements performed at room temperature, the fraction of chemical formal units of  $\text{MAPbI}_3$  against  $\text{PbI}_2$  may be approximated as the conversion ratio  $\varphi$ <sup>29</sup>

$$\varphi = \frac{\frac{A_{\text{norm, MAPbI}_3(110)}}{|F_{\text{MAPbI}_3(110)}|^2 p_{\text{MAPbI}_3(110)} L_{\text{MAPbI}_3(110)}}}{\frac{A_{\text{norm, PbI}_2(001)}}{|F_{\text{PbI}_2(001)}|^2 p_{\text{PbI}_2(001)} L_{\text{PbI}_2(001)}}} = \frac{A_{\text{norm, MAPbI}_3(110)}}{A_{\text{norm, PbI}_2(001)}} \frac{|F_{\text{PbI}_2(001)}|^2 p_{\text{PbI}_2(001)} L_{\text{PbI}_2(001)}}{|F_{\text{MAPbI}_3(110)}|^2 p_{\text{MAPbI}_3(110)} L_{\text{MAPbI}_3(110)}} \quad (5)$$

where  $F$  further denotes the structure factor per formal unit for the given reflection plane,  $p$  denotes the multiplicity of the given

reflection plane within each crystal structure, and  $L$  defines the Lorentz-polarization factor, a trigonometric term considering intensity changes resulting from the orientation of the sample towards the incident beam (see method section). Errors in this approximation may arise from amorphous phases within the thin film or the significant presence of a layered structure within the thin film. However, as these factors would be expected to be similar for all samples, they result in a systematic error of  $\varphi$ . Consequently, it serves as a suitable semiquantitative descriptor for the phase ratio within the studied thin films.

The investigated combinatorial library exhibits a large range of conversion ratios ranging from  $\varphi \approx 0.2$  to  $\varphi \approx 5$  (Fig. 2c), rendering it ideal for the combinatorial analysis of  $\text{PbI}_2$  residuals within  $\text{MAPbI}_3$ . From repeated measurements of the same sample after extended periods of time within the inert-gas XRD solution, we further verified the phase stability of the library during this characterization (Fig. S7†).

Although not clearly visible by eye (Fig. S8†), subtle changes in optical properties across the library can be successfully resolved from automated UV-Vis measurements. A clear trend with respect to the phase ratio was observed: Samples with higher MHP phase contents exhibit an overall reduced absorption below and above the bandgap of 1.6 eV (Fig. 3a) compared to those with lower MHP phase fractions. The  $\text{PbI}_2$  template deposited by thermal evaporation does not exhibit notable optical absorption within this range (Fig. S9†), indicating that optical scattering from mixed phases may cause this effect. Notably, the intensity of the first derivative of the reflection-correct sample transmission  $T'$  (Fig. 3b), and hence the steepness of the optical



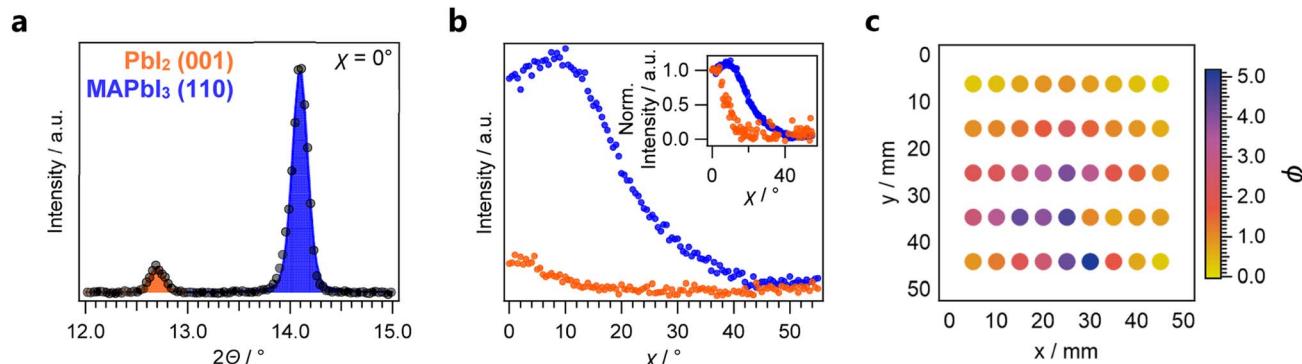


Fig. 2 XRD characterization of the combinatorial  $\text{MAPbI}_3$  sample. (a) 2-Theta scan for a selected sample and (b) Azimuthal  $\chi$ -scan for the same sample. The inset depicts traces that have been normalized at  $\chi = 0^\circ$ , according to the term defined in eqn (3). Based on these individual results, the conversion ratio is calculated for each sample. (c) Conversion ratio  $\varphi$  as function of the position on the combinatorial library.

absorption, is highly correlated to  $\varphi$ . In contrast, the maximum exhibited shifts of less than 0.1 eV and is therefore less sensitive to the changes of  $\varphi$ . Likewise, manually extracted bandgaps by Tauc plot analysis vary by less than 0.2 eV across the complete data set (Fig. S10†), which lies below inaccuracies typically resulting from automatic Tauc-plot analysis algorithms.<sup>49</sup> Hence, observation of the changed steepness of the absorption onset appears more robust than extracted bandgaps to track phase-composition changes within the thin film. Furthermore, these results highlight how a comprehensive automated thin-film analysis allows to resolve trends otherwise not obtainable from either experiments on few samples or HTE workflows relying on optical measurements alone.

### 3.3 Effect of conversion ratio on aging

To experimentally screen hybrid perovskite thin films with respect to their phase stability, RGB-camera-based optical characterization has previously been successfully employed.<sup>50</sup> Despite first successful demonstrations of image-based spectra-prediction for semiconductors by Stein *et al.*,<sup>51</sup> reliably resolving finer spectral changes within the optical absorption from this approach remains elusive.

Indeed, despite the volatile nature of the  $\text{MA}^+$  cation, we observed minute spectral changes during aging for thin films with higher  $\varphi$  under inert conditions (Fig. 4a). We observed a reduction in the absorption-onset slope, reduced optical absorption above the  $\text{MAPbI}_3$  bandgap, and an increase in the sub-bandgap absorption (Fig. 4b). These trends are aligned with *in situ* UV-Vis aging studies performed for single MHP samples.<sup>52</sup> Furthermore, the resemblance of the spectra from aged samples to spectra from as-deposited thin films with the same corresponding phase ratio  $\varphi$  suggests a back conversion of  $\text{MAPbI}_3$  into  $\text{PbI}_2$  upon outgassing of MAI as the primary aging mechanism.

This mechanism was further confirmed through post-mortem XRD measurements: Remarkably, we determined the loss in conversion to be equal for all samples as evident from the linear relationship between MHP/ $\text{PbI}_2$  phase ratios  $\varphi$  before and after the aging process (Fig. 4b). Through observation of absorption-coefficient changes  $\Delta\alpha$  both below (at 1.4 eV, Fig. 4d) and above the bandgap (at 1.7 eV, Fig. 4e), the evolution of optical properties as a function of aging time can be visualized across the library.

Below the bandgap, a linear increase in the absorption coefficient of up to  $\Delta\alpha = 3 \times 10^3 \text{ cm}^{-1}$  was observed over the

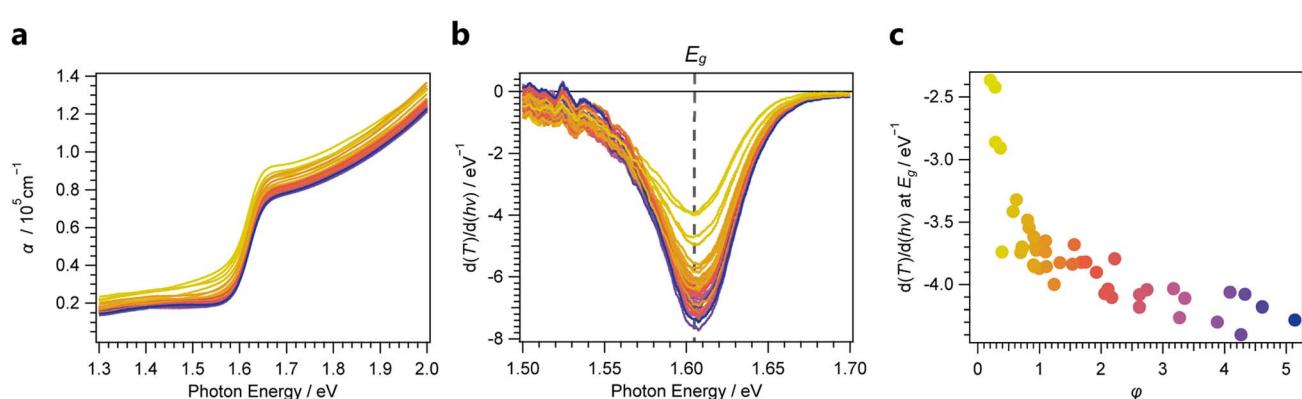


Fig. 3 (a) Absorption coefficients  $\alpha$  with varying  $\varphi$ . (b) First derivative of the reflection-corrected transmission  $T'$ , showing the steepness of the transmission onset near the bandgap  $E_g$ . (c) Extracted values of subplot b, indicating a clear correlation between the MHP/ $\text{PbI}_2$  and the steepness of the transmission onset near the bandgap.



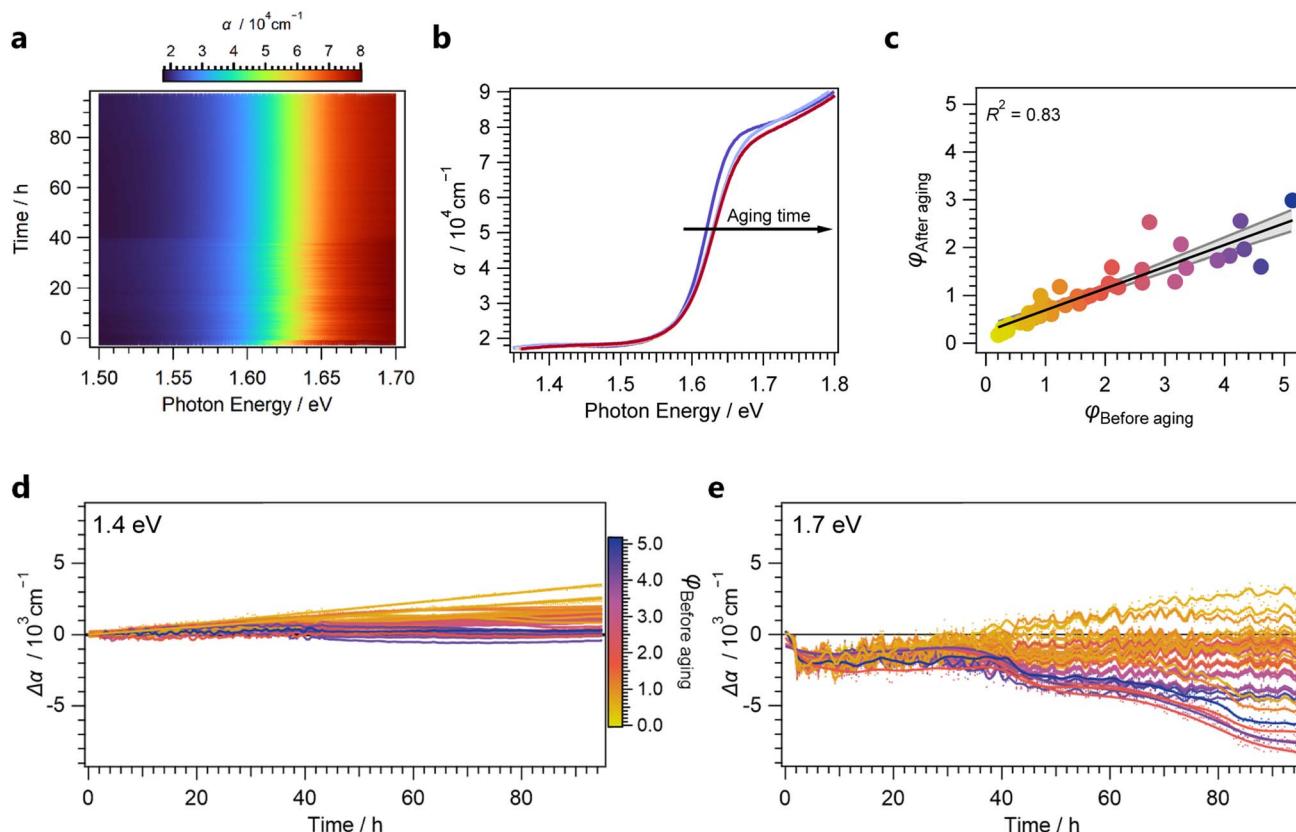


Fig. 4 (a) Change of optical properties over time for a single sample. (b) Selected spectra for a single sample at equidistantly spaced aging-time intervals. (c) Results from phase ratio determination before and after aging process, indicating a linear loss across the library. (d) Changes in absorption below the bandgap (at 1.4 eV). (e) Changes in optical absorption above the bandgap (at 1.7 eV).

aging time. This increase is particularly profound for samples exhibiting low MHP/PbI<sub>2</sub> phase ratios, indicating an increase in optical light scattering. Scattering effects in thin films can have multiple origins, such as interfacial roughness, as well as interference effects, and bulk scattering.<sup>53</sup> Thus, a quantitative understanding would require further optical modelling which is beyond the scope of the present work.

Still, these trends enable a qualitative understanding of the complex changes occurring above the bandgap: initially, a decrease in the optical absorption within 10 h of aging was observed for all samples. Afterwards, while low-converted samples exhibit a clear increase in absorption from likewise increased scattering, highly-converted samples exhibit a decrease in the absorption coefficient, resulting in a range of  $\Delta\alpha = 11 \times 10^3 \text{ cm}^{-1}$ . These effects can likely be assigned to an overlay of scattering effects as well as a loss of optical absorption from outgassing of the organohalides. Outgassing of the organohalides would induce the aforementioned back conversion into PbI<sub>2</sub> which exhibits a higher absorption onset. Interestingly, the observations indicate a non-linear trend for the outgassing with aging time. These kinetics are most pronounced for the highly-converted samples where the increase of the optical scattering during the aging time was observed to be negligible below the bandgap. The S-shaped feature in the optical absorption changes above the bandgap

at 40 h can be linked to autocatalytic degradation resulting from the local formation of I<sub>2</sub> from HI vapour during aging.<sup>54</sup> This resembles recent reports where similar trends were determined from XRD measurements with lower temporal resolution.<sup>42</sup>

Overall, these findings showcase the complex and nonlinear changes in optical spectra that occur during the degradation process of MHP thin films and result in implications for their optical stability screening: the apparent optical absorption of MHP thin films above the bandgap during aging can be heavily influenced by scattering effects, complicating the quantitative analysis. While previous approaches using RGB-cameras were able to deduce similar kinetic results from red colour channels alone,<sup>55</sup> these effects are critical for samples with overall lower optical absorption than the herein investigated thin films with PbI<sub>2</sub> residuals. Furthermore, the autocatalytic nature of the herein observed degradation mechanism induces nonlinear changes of the optical properties during aging. As a result, linear extrapolation of trends observed only before and after aging may give a falsified indication on the actual materials stability.

### 3.4 Machine learning assisted analysis of aging kinetics

To gain a more detailed understanding of the degradation kinetics across the complete library, we continued our investigation with data-driven approaches to extract kinetic



information from the recorded *in situ* UV-Vis spectra. The high correlation of trends observed from XRD and UV-Vis measurements motivated us to develop a model that can predict  $\varphi$  from UV-Vis spectra obtained during the aging process, effectively yielding insights about the phase composition evolution from our *in situ* UV-Vis investigation.

To ensure generalization while still being able to account for non-linear trends occurring in optical absorption spectra, we focused on ensemble methods, striking a balance between simple linear regression and data-intensive deep learning approaches. Based on a gradient boosting regressor, optimized for smaller data sets (see method section),<sup>30</sup> we obtained a model which generalized well on the testing data set (Fig. 5), yielding a root-mean-square-error (RMSE) of 0.37 for the prediction of  $\varphi$ .

The feature selection was performed based on the previous empirical trends observed for MHP thin films with varied  $\varphi$ : To learn from the complete spectral shape, we first chose the ratio of optical absorption below and above the bandgap as our input feature. To this end, absorption coefficient ratios at 1.7 eV vs. 1.5 eV as well as at 2.0 eV vs. 1.4 eV were used. Inclusion of these features allows to consider sub-bandgap absorption and hence scattering effects in our analysis. Additionally, the initially introduced maximum from the first derivative of the optical transmission around the MHP bandgap (Fig. 3b) was included. Analysis of the feature importance revealed the latter feature to be the most relevant, while the two absorption coefficient ratios used to learn from the spectral shape were similarly important (Fig. S11†).

When applying this method on the *in situ* UV-vis data set, we were able to unravel the kinetics of less converted samples, previously inaccessible through more classic observations of the optical spectra alone (Fig. 6b):

A continuous decrease of  $\varphi$  and therefore the loss of the  $\text{MAPbI}_3$  phase during aging was observed, resulting in a rapid loss of at least 20% from the initial value after 20 h for all studied samples. This is in stark contrast to the smaller decrease in the optical absorption observed for all investigated thin films above the bandgap and lack of any changes below the bandgap (Fig. 4e).

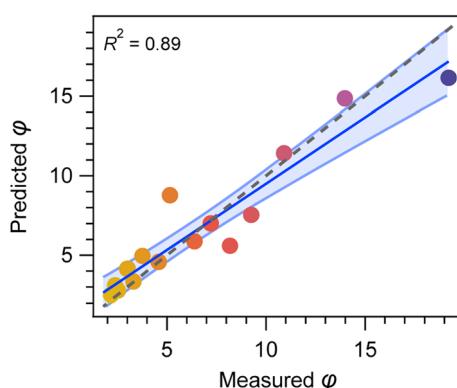


Fig. 5 Prediction of the MHP/PbI<sub>2</sub> phase ratio  $\varphi$  from UV-Vis spectra by the ensemble regression model, achieving a high  $R^2$  score of 0.89 and a RMSE of 0.37.

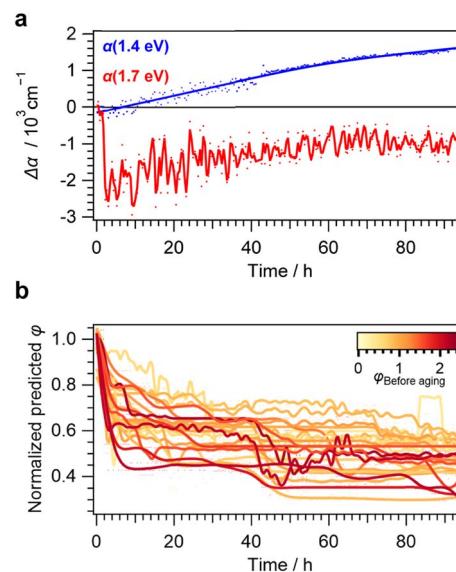


Fig. 6 (a) Observation of the change in optical absorption  $\Delta\alpha$  for a representative sample with  $\varphi_{\text{Before aging}} = 0.90$  and thus large amounts of scattering effects as evident from the absorption increase below the bandgap. (b) Normalized predicted kinetic traces expressed as changes in  $\varphi$  for samples with low  $\text{MAPbI}_3$  phase ratios ( $\varphi < 2.5$ ), showing a continuous decrease.

Even more strikingly, an S-shaped feature at approx. 40 h of aging time was observed for most samples. This is in agreement with the autocatalytic degradation at this point and the qualitative observations made from thin films without significant scattering effects. Fluctuations overlaying these trends can be seen for extremely low-converted samples presented in yellow hues. Most likely, this is a result of the prediction error as samples exhibiting low  $\varphi$  values before aging had lower absolute changes of  $\varphi$  during aging (Fig. 4c). Beyond the inflection point of this degradation effect, further decreases stayed within 10% of the initial  $\varphi$  value, even when accounting for fluctuations within the predictions. This again highlights the nonlinear nature of the degradation behavior detected from *in situ* UV-Vis characterization with high temporal resolution.

Interestingly, applying the regressor model on the higher converted samples gives less steady trends (Fig. S12†). This is both apparent from higher fluctuations and predictions beyond  $\varphi \geq 1$  at 50 h of aging time. These shortcomings can be explained by the training bias as entries with  $\varphi > 2.5$  were only found in unaged samples and therefore comprised only 19% of the training data (Fig. S13†). Furthermore, despite their high potential for data-driven experimental research,<sup>56–58</sup> tree-based regression models are typically not well-suited for predictions beyond the median or even extrapolation outside of the training data set.<sup>59</sup>

For the herein employed workflow however, these shortcomings are not crucial since applying the machine-learning analysis is of most use for degraded and low-converted samples that are typically more prevalent within the training data. Thus, our findings highlight both the utility and limitations of applying these models for the quantitative analysis of UV-Vis data for material degradation analysis.



## 4 Conclusions & outlook

A comprehensive automated inert-gas workflow was developed and employed to investigate the aging of MHP thin films with various amounts of  $\text{PbI}_2$  residuals using a combination of XRF, XRD and UV-Vis characterization methods. From combinatorial *in situ* UV-Vis aging studies, we obtained a detailed understanding of the underlying degradation kinetics. While the overall mechanism is in alignment with an autocatalytic degradation mechanism as reported elsewhere for experiments not performed under-inert gas,<sup>42,55</sup> we were able to detect the independence of the degradation kinetics from the phase composition for underconverted MHP samples. We further explored how machine-learning assisted analysis of UV-Vis data with featurization based on a detailed materials characterization before and after the aging experiments can be used to link results from optical analysis to conversion rate. This enables deeper insights into the degradation kinetics beyond qualitative discussions alone. The plateauing effect observed for all investigated thin films within this study suggests a common self-limiting mechanism for their aging beyond the inflection point which could be of interest for further investigations.

Overall, beyond the aforementioned increased understanding of the MHP thin films aging kinetics, our results highlight the potential, but also important considerations of accelerated aging studies based on optical spectroscopy. Relying on optical changes alone, especially at one wavelength only, may provide limited insights of the actual phase stability within the thin films. These effects would be especially pronounced for material systems exhibiting lower optical absorption coefficient maxima, *e.g.* charge-transport-layers, transparent conductive films, or antireflective coatings, but can be minimized by analysis of the complete spectral shape. For selected material systems, the simultaneous measurement at few selected wavelengths, *e.g.* below, around and above the bandgap, may be sufficient. In addition, we show that including complementary characterization techniques within the workflow can provide crucial validity checks and enable a more detailed scientific understanding of the underlying aging mechanisms. The usage of machine-learning assisted analysis techniques on the other hand can be used to extract actionable intelligence from the large amounts of data generated by these workflows. Lastly, the herein presented inert-gas workflow is ideal for integrating other techniques such as X-ray photoelectron spectroscopy (XPS) for surface analysis mapping, which would provide an even deeper understanding of materials aging behavior.<sup>60,61</sup>

## Author contributions

Alexander Wieczorek: conceptualization, formal analysis, investigation, methodology, software, validation, visualization, writing – original draft, writing – review and editing. Austin G. Kuba: investigation, methodology, supervision, writing – review and editing. Jan Sommerhäuser: software, writing – review and editing. Luis Nicklaus Caceres: software, writing – review and editing. Christian Wolff: funding acquisition, resources,

supervision, writing – review and editing. Sebastian Siol: conceptualization, investigation, methodology, funding acquisition, project administration, resources, supervision, writing – review and editing.

## Conflicts of interest

The authors declare no conflict of interest.

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## References

- 1 B. Guo, R. Lai, S. Jiang, L. Zhou, Z. Ren, Y. Lian, P. Li, X. Cao, S. Xing, Y. Wang, W. Li, C. Zou, M. Chen, Z. Hong, C. Li, B. Zhao and D. Di, *Nat. Photonics*, 2022, **16**, 637–643.
- 2 H. Lai, J. Luo, Y. Zwirner, S. Olthof, A. Wieczorek, F. Ye, Q. Jeangros, X. Yin, F. Akhundova, T. Ma, R. He, R. K. Kothandaraman, X. Chin, E. Gilshtein, A. Müller, C. Wang, J. Thiesbrummel, S. Siol, J. M. Prieto, T. Unold, M. Stolterfoht, C. Chen, A. N. Tiwari, D. Zhao and F. Fu, *Adv. Energy Mater.*, 2022, **12**, 2202438.
- 3 K. Sakhatskyi, B. Turedi, G. J. Matt, E. Wu, A. Sakhatska, V. Bartosh, M. N. Lintangpradipto, R. Naphade, I. Shorubalko, O. F. Mohammed, S. Yakunin, O. M. Bakr and M. V. Kovalenko, *Nat. Photonics*, 2023, **17**, 510–517.
- 4 W. Fu, A. G. Ricciardulli, Q. A. Akkerman, R. A. John, M. M. Tavakoli, S. Essig, M. V. Kovalenko and M. Saliba, *Mater. Today*, 2022, **58**, 275–296.
- 5 T. D. Siegler, A. Dawson, P. Lobaccaro, D. Ung, M. E. Beck, G. Nilsen and L. L. Tinker, *ACS Energy Lett.*, 2022, **7**, 1728–1734.
- 6 V. M. Goldschmidt, *Naturwissenschaften*, 1926, **14**, 477–485.



- 7 A. M. Glazer and H. D. Megaw, *Philos. Mag.*, 1972, **25**, 1119–1135.
- 8 P. M. Woodward, *Acta Crystallogr., Sect. B: Struct. Sci.*, 1997, **53**, 44–66.
- 9 M. W. Lufaso and P. M. Woodward, *Acta Crystallogr., Sect. B: Struct. Sci.*, 2004, **60**, 10–20.
- 10 J.-H. Lee, N. C. Bristowe, J. H. Lee, S.-H. Lee, P. D. Bristowe, A. K. Cheetham and H. M. Jang, *Chem. Mater.*, 2016, **28**, 4259–4266.
- 11 B. Conings, J. Drijkoningen, N. Gauquelin, A. Babayigit, J. D'Haen, L. D'Olieslaeger, A. Ethirajan, J. Verbeeck, J. Manca, E. Mosconi, F. De Angelis and H.-G. Boyen, *Adv. Energy Mater.*, 2015, **5**, 1500477.
- 12 H. Wei, S. Chen, J. Zhao, Z. Yu and J. Huang, *Chem. Mater.*, 2020, **32**, 2501–2507.
- 13 L. Lanzetta, T. Webb, N. Zibouche, X. Liang, D. Ding, G. Min, R. J. E. Westbrook, B. Gaggio, T. J. Macdonald, M. S. Islam and S. A. Haque, *Nat. Commun.*, 2021, **12**, 1–11.
- 14 M. Saliba, M. Stolterfoht, C. M. Wolff, D. Neher and A. Abate, *Joule*, 2018, **2**, 1019–1024.
- 15 M. V. Khenkin, E. A. Katz, A. Abate, G. Bardizza, J. J. Berry, C. Brabec, F. Brunetti, V. Bulović, Q. Burlingame, A. Di Carlo, R. Cheacharoen, Y.-B. Cheng, A. Colsmann, S. Cros, K. Domanski, M. Dusza, C. J. Fell, S. R. Forrest, Y. Galagan, D. Di Girolamo, M. Grätzel, A. Hagfeldt, E. von Hauff, H. Hoppe, J. Kettle, H. Köbler, M. S. Leite, S. Liu, Y.-L. Loo, J. M. Luther, C.-Q. Ma, M. Madsen, M. Manceau, M. Matheron, M. McGehee, R. Meitzner, M. K. Nazeeruddin, A. F. Nogueira, Ç. Odabaşı, A. Osheroval, N.-G. Park, M. O. Reese, F. De Rossi, M. Saliba, U. S. Schubert, H. J. Snaith, S. D. Stranks, W. Tress, P. A. Troshin, V. Turkovic, S. Veenstra, I. Visoly-Fisher, A. Walsh, T. Watson, H. Xie, R. Yıldırım, S. M. Zakeeruddin, K. Zhu and M. Lira-Cantu, *Nat. Energy*, 2020, **5**, 35–49.
- 16 Y. Zhao, J. Zhang, Z. Xu, S. Sun, S. Langner, N. T. P. Hartono, T. Heumueller, Y. Hou, J. Elia, N. Li, G. J. Matt, X. Du, W. Meng, A. Osvet, K. Zhang, T. Stubhan, Y. Feng, J. Hauch, E. H. Sargent, T. Buonassisi and C. J. Brabec, *Nat. Commun.*, 2021, **12**, 2191.
- 17 N. T. P. Hartono, H. Köbler, P. Graniero, M. Khenkin, R. Schlattmann, C. Ulbrich and A. Abate, *Nat. Commun.*, 2023, **14**, 4869.
- 18 S. Sun, A. Tiihonen, F. Oviedo, Z. Liu, J. Thapa, Y. Zhao, N. T. P. Hartono, A. Goyal, T. Heumueller, C. Batali, A. Encinas, J. J. Yoo, R. Li, Z. Ren, I. M. Peters, C. J. Brabec, M. G. Bawendi, V. Stevanovic, J. Fisher and T. Buonassisi, *Matter*, 2021, **4**, 1305–1322.
- 19 H. Koinuma and I. Takeuchi, *Nat. Mater.*, 2004, **3**, 429–438.
- 20 A. Zakutayev, N. Wunder, M. Schwarting, J. D. Perkins, R. White, K. Munch, W. Tumas and C. Phillips, *Sci. Data*, 2018, **5**, 180053.
- 21 J. M. Gregoire, L. Zhou and J. A. Haber, *Nat. Synth.*, 2023, **2**, 493–504.
- 22 I. Susic, A. Kama, L. Gil-Escríg, C. Dreessen, F. Palazon, D. Cahen, M. Sessolo and H. J. Bolink, *Adv. Mater. Interfaces*, 2023, **10**, 2202271.
- 23 Y. Cheng and L. Ding, *Energy Environ. Sci.*, 2021, **14**, 3233–3255.
- 24 Q.-Q. Chu, Z. Sun, D. Wang, B. Cheng, H. Wang, C.-P. Wong and B. Fang, *Matter*, 2023, **6**, 3838–3863.
- 25 X. J. She, C. Chen, G. Divitini, B. Zhao, Y. Li, J. Wang, J. F. Orri, L. Cui, W. Xu, J. Peng, S. Wang, A. Sadhanala and H. Sirringhaus, *Nat. Electron.*, 2020, **3**, 694–703.
- 26 M. Nentwich, T. Weigel, C. Richter, H. Stöcker, E. Mehner, S. Jachalke, D. V. Novikov, M. Zschornak and D. C. Meyer, *J. Synchrotron Radiat.*, 2021, **28**, 158–168.
- 27 K. R. Talley, S. R. Bauers, C. L. Melamed, M. C. Papac, K. N. Heinselman, I. Khan, D. M. Roberts, V. Jacobson, A. Mis, G. L. Brennecka, J. D. Perkins and A. Zakutayev, *ACS Comb. Sci.*, 2019, **21**, 537–547.
- 28 D. Zagorac, H. Müller, S. Ruehl, J. Zagorac and S. Rehme, *J. Appl. Crystallogr.*, 2019, **52**, 918–925.
- 29 D. Beaudoin, T. Maris and J. D. Wuest, *Nat. Chem.*, 2013, **5**, 830–834.
- 30 T. Chen and C. Guestrin, in *Proceedings of the 22nd ACM SIGKDD International Conference on Knowledge Discovery and Data Mining*, ACM, New York, NY, USA, 2016, vol. 13–17, pp. 785–794.
- 31 F. Pedregosa, G. Varoquaux, A. Gramfort, V. Michel, B. Thirion, O. Grisel, M. Blondel, P. Prettenhofer, R. Weiss, V. Dubourg, J. Vanderplas, A. Passos, D. Cournapeau, M. Brucher, M. Perrot and E. Duchesnay, *J. Mach. Learn. Res.*, 2011, **12**, 2825–2830.
- 32 F. Sahli, J. Werner, B. A. Kamino, M. Bräuninger, R. Monnard, B. Paviet-Salomon, L. Barraud, L. Ding, J. J. Diaz Leon, D. Sacchetto, G. Cattaneo, M. Despeisse, M. Boccard, S. Nicolay, Q. Jeangros, B. Niesen and C. Ballif, *Nat. Mater.*, 2018, **17**, 820–826.
- 33 Z. Wu, E. Bi, C. Li, L. Chen, Z. Song and Y. Yan, *Sol. RRL*, 2023, **7**, 2200571.
- 34 T. H. Kim and S. G. Kim, *Saf. Health Work*, 2011, **2**, 97–104.
- 35 X. Y. Chin, D. Turkay, J. A. Steele, S. Tabean, S. Eswara, M. Mensi, P. Fiala, C. M. Wolff, A. Paracchino, K. Artuk, D. Jacobs, Q. Guesnay, F. Sahli, G. Andreatta, M. Boccard, Q. Jeangros and C. Ballif, *Science*, 2023, **381**, 59–63.
- 36 T. Moser, K. Artuk, Y. Jiang, T. Feurer, E. Gilshtein, A. N. Tiwari and F. Fu, *J. Mater. Chem. A*, 2020, **8**, 21973–21982.
- 37 Q. Guesnay, F. Sahli, C. Ballif and Q. Jeangros, *APL Mater.*, 2021, **9**, 100703.
- 38 S. Wang, Y. Jiang, E. J. Juarez-Perez, L. K. Ono and Y. Qi, *Nat. Energy*, 2017, **2**, 1–8.
- 39 F. Fu, S. Pisoni, Q. Jeangros, J. Sastre-Pellicer, M. Kawecki, A. Paracchino, T. Moser, J. Werner, C. Andres, L. Duchêne, P. Fiala, M. Rawlence, S. Nicolay, C. Ballif, A. N. Tiwari and S. Buecheler, *Energy Environ. Sci.*, 2019, **12**, 3074–3088.
- 40 J. Zhao, Y. Deng, H. Wei, X. Zheng, Z. Yu, Y. Shao, J. E. Shield and J. Huang, *Sci. Adv.*, 2017, **3**, eaao5616.
- 41 A. G. Kuba, A. J. Harding, R. J. Richardson, B. E. McCandless, U. K. Das, K. D. Dobson and W. N. Shafarman, *ACS Appl. Energy Mater.*, 2022, **5**, 10731–10741.
- 42 C. L. C. Ellis, H. Javaid, E. C. Smith and D. Venkataraman, *Inorg. Chem.*, 2020, **59**, 12176–12186.



- 43 N. Pant, A. Kulkarni, M. Yanagida, Y. Shirai, T. Miyasaka and K. Miyano, *ACS Appl. Energy Mater.*, 2020, **3**, 6215–6221.
- 44 J. Patidar, A. Sharma, S. Zhuk, G. Lorenzin, C. Cancellieri, M. F. Sarott, M. Trassin, K. Thorwarth, J. Michler and S. Siol, *Surf. Coat. Technol.*, 2023, 129719.
- 45 J. M. Nieuwenhuizen and H. B. Haanstra, *Philips Tech. Rev.*, 1966, **27**, 87–91.
- 46 J. L. Baker, L. H. Jimison, S. Mannsfeld, S. Volkman, S. Yin, V. Subramanian, A. Salleo, A. P. Alivisatos and M. F. Toney, *Langmuir*, 2010, **26**, 9146–9151.
- 47 Q. Guesnay, C. J. McMonagle, D. Chernyshov, W. Zia, A. Wieczorek, S. Siol, M. Saliba, C. Ballif and C. M. Wolff, *ACS Photonics*, 2023, **10**, 3087–3094.
- 48 H. Li, J. Zhou, L. Tan, M. Li, C. Jiang, S. Wang, X. Zhao, Y. Liu, Y. Zhang, Y. Ye, W. Tress and C. Yi, *Sci. Adv.*, 2022, **8**, 1–9.
- 49 M. Schwarting, S. Siol, K. Talley, A. Zakutayev and C. Phillips, *Mater. Discov.*, 2017, **10**, 43–52.
- 50 R. Keesey, A. Tiihonen, A. E. Siemann, T. W. Colburn, S. Sun, N. T. Putri Hartono, J. Serdy, M. Zeile, K. He, C. A. Gurtner, A. C. Flick, C. Batali, A. Encinas, R. R. Naik, Z. Liu, F. Oviedo, I. M. Peters, J. Thapa, S. I. Parker Tian, R. H. Dauskardt, A. J. Norquist and T. Buonassisi, *Digital Discovery*, 2023, **2**, 422–440.
- 51 H. S. Stein, D. Guevarra, P. F. Newhouse, E. Soedarmadji and J. M. Gregoire, *Chem. Sci.*, 2019, **10**, 47–55.
- 52 J. Yang, B. D. Siempelkamp, D. Liu and T. L. Kelly, *ACS Nano*, 2015, **9**, 1955–1963.
- 53 S. Schröder, M. Trost, M. Garrick, A. Duparré, X. Cheng, J. Zhang and Z. Wang, *Thin Solid Films*, 2015, **592**, 248–255.
- 54 A. Latini, G. Gigli and A. Ciccioli, *Sustainable Energy Fuels*, 2017, **1**, 1351–1357.
- 55 R. R. Naik, A. Tiihonen, J. Thapa, C. Batali, Z. Liu, S. Sun and T. Buonassisi, *npj Comput. Mater.*, 2022, **8**, 72.
- 56 B. Selvaratnam and R. T. Koodali, *Catal. Today*, 2021, **371**, 77–84.
- 57 R. E. Kumar, A. Tiihonen, S. Sun, D. P. Fenning, Z. Liu and T. Buonassisi, *Matter*, 2022, **5**, 1353–1366.
- 58 K. Wieczorek, A. Groetsch, K. Pajor, M. Jain, A. M. Müller, C. Vockenhuber, J. Schwiedrzik, A. Sharma, F. F. Klimashin and J. Michler, *Adv. Sci.*, 2023, 2302997.
- 59 H. Zhang, D. Nettleton and Z. Zhu, *arXiv*, 2019, preprint, 1904.10416 DOI: [10.48550/arXiv.1904.10416](https://doi.org/10.48550/arXiv.1904.10416).
- 60 A. Wieczorek, H. Lai, J. Pious, F. Fu and S. Siol, *Adv. Mater. Interfaces*, 2023, **10**, 2370024.
- 61 S. Zhuk, A. Wieczorek, A. Sharma, J. Patidar, K. Thorwarth, J. Michler and S. Siol, *Chem. Mater.*, 2023, **35**, 7069–7078.

