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Organic–inorganic hybrid metallic conductors based on bis(ethylenedithio)tetrathiafulvalene cations and antiferromagnetic oxalate-bridged copper(II) dinuclear anions†

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The organic–inorganic hybrid β'' -(BEDT-TTF)₃[Cu₂(μ -C₂O₄)(C₂O₄)₂(CH₃OH)(H₂O)] (BEDT-TTF = bis(ethylenedithio)tetrathiafulvalene) composed of a BEDT-TTF donor and the oxalate-bridged binuclear anion [Cu₂(μ -C₂O₄)(C₂O₄)₂(CH₃OH)(H₂O)]²⁻ has been obtained by electrocrystallization. It crystallizes in the triclinic $P\bar{1}$ space group with cell parameters of $a = 7.4803(3)$ Å, $b = 9.3547(3)$ Å, $c = 18.6711(7)$ Å, $\alpha = 95.797(3)^\circ$, $\beta = 90.974(3)^\circ$, $\gamma = 93.508(3)^\circ$, $V = 1297.06(8)$ Å³, and $Z = 1$ at 100 K. The donor arrangement belongs to the β'' phase. From the TTF core bond lengths and Raman spectroscopy, the oxidation state of BEDT-TTF is assigned to $\sim +2/3$. CH₃OH or H₂O molecules bond to the metal atoms at the apical position of the square pyramid with an occupancy of 0.5. A supramolecular square lattice forms through hydrogen bonds between the antiferromagnetic binuclear anions in the anion sheet. From the band structure at 100 K, metallic conductivity is expected, which agrees with four-probe conductivity measurements: its conductivity is 11.5 S cm⁻¹ at room temperature, increases to 160 S cm⁻¹ at 7.6 K, and then decreases to 150 S cm⁻¹ at 2 K. From magnetic measurements, there is no long-range magnetic ordering, which is confirmed by specific heat measurements.

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Introduction

Organic–inorganic hybrid charge-transfer salts composed of organic donors and inorganic magnetic anions provide many dual-functional molecular crystals showing magnetism, such as

paramagnetism, long-range magnetic ordering, spin cross-over, and spin frustration, together with conductivity from insulating to semiconducting to metallic conductivity to even superconductivity. These materials have attracted considerable attention in materials science because of their potential uses in molecular spintronics and in the search for new superconductors.¹ The oxalate anion (C₂O₄²⁻) is one of the most commonly used coordination ligands to mediate magnetic interactions between transition metals. Research on oxalate-bridged zero-dimensional (0D) binuclear compounds started in the 1970s and it is still active for molecular magnets.² Research has been extended to two-dimensional (2D) honeycomb metal-oxalate compounds, and the long-range ordered (LRO) [(C₄H₉)₄N][CrMn(μ -C₂O₄)₃] molecular ferromagnet and [(C₄H₉)₄N][Fe^{II}Fe^{III}(μ -C₂O₄)₃] molecular ferrimagnet were discovered in the 1990s.³ Based on the oxalate-based molecular magnet with a honeycomb lattice, researchers have started to search for dual-functional molecular crystals composed of organic donors and inorganic metal-oxalate anions as charge-transfer salts. β'' -(BEDT-TTF)₄[(H₃O)Fe(C₂O₄)₃] \cdot C₆H₅CN (BEDT-TTF = bis(ethylenedithio)tetrathiafulvalene), the first superconductor containing a magnetic atom, was reported in 1995.⁴ The coexistence of ferromagnetic long-range order and metallic conductivity was observed in β -(BEDT-TTF)₃[CrMn(μ -C₂O₄)₃] \cdot xCH₂Cl₂, and α -(BETS)₃[CrMn(μ -C₂O₄)₃]

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$x\text{CH}_2\text{Cl}_2$ (BETS = bis(ethylenedithio)tetraselenafulvalene, $x = 2-3$) in the 2000s.⁵ Strong antiferromagnetic behaviour without long-range ordering was observed in the θ^{21} -(BEDT-TTF)₃[Cu₂(μ -C₂O₄)₃] \cdot 2CH₃OH (2) semiconductor and the θ^{21} -(BETS)₃[Cu₂(μ -C₂O₄)₃] \cdot 2CH₃OH (3) conductor in the 2010s.⁶ In organic-inorganic hybrid dual-functional molecular crystals with honeycomb metal-oxalate anions, organic molecules, such as C₆H₅CN, CH₂Cl₂, and CH₃OH, exist in the vacancies of the honeycomb lattice without bonding to the metal atoms.

The charge-transfer salts of 0D oxalate-bridged dinuclear anions, antiferromagnetic insulator TTF₅[Fe₂(μ -C₂O₄)(C₂O₄)₄] \cdot 2C₆H₅CH₃ \cdot 2H₂O, antiferromagnetic semiconductor (BEDT-TTF)₄[Fe₂(μ -C₂O₄)(C₂O₄)₄] and TMTTF₄[Fe₂(μ -C₂O₄)(C₂O₄)₄] \cdot C₆H₅CN \cdot 4H₂O, were reported in the 2000s.⁷ They are different from the charge-transfer salts of 0D inorganic CuCl₄⁻, FeCl₄⁻ and FeBr₄⁻ with the π - d interaction between the donor and anion from the S/Se \cdots X interaction, such as the field-induced superconductivity in antiferromagnetic insulator λ -BETS₂FeCl₄, and the modulation of conductivity states between the insulator, metal, and superconductor by a magnetic field in λ -BETS₂Fe_{0.4}Ga_{0.6}Cl₄.^{1,8}

In 0D oxalate-bridged dinuclear copper compounds, the interaction between the two Jahn-Teller distorted metal ions

varies from ferromagnetic to antiferromagnetic.⁹ Magnetos-structural correlation plays an important role in compounds from classic molecular magnets to quantum spin liquids. For example, 0D oxalate-bridged [Cu₂(μ -C₂O₄)(C₂O₄)₂]²⁻ is one of the coordination isomers of the 2D honeycomb [Cu₂(μ -C₂O₄)₃]²⁻_n and 3D hyperhoneycomb [Cu₂(μ -C₂O₄)₃]²⁻_n lattices (Scheme 1).^{6,10} When the coordination environment of Cu²⁺ is modified from a square to an octahedron, 0D [Cu₂(μ -C₂O₄)(C₂O₄)₂]²⁻ is the building block of the 2D honeycomb [Cu₂(μ -C₂O₄)₃]²⁻_n and 3D hyperhoneycomb [Cu₂(μ -C₂O₄)₃]²⁻_n lattices. The magnetic structures of the 3D hyperhoneycomb [Cu₂(μ -C₂O₄)₃]²⁻_n lattice obtained from Jahn-Teller distorted 0D Cu₂(μ -C₂O₄)(C₂O₄)₄²⁻ orbital analysis show good agreement with theoretical analysis.^{11,12} Spin frustration with $f > 10$ ($f = |\theta|/Tc$) has been observed in the molecular magnet and organic-inorganic hybrid dual-functional molecular crystals with the 2D honeycomb anions: insulating [(C₃H₇)₃NH]₂[Cu₂(μ -C₂O₄)₃] \cdot 2.2H₂O, semiconductive θ^{21} -(BEDT-TTF)₃[Cu₂(μ -C₂O₄)₃] \cdot 2CH₃OH, conductive θ^{21} -(BETS)₃[Cu₂(μ -C₂O₄)₃] \cdot 2CH₃OH, and insulating [(C₂H₅)₃NH]₂[Cu₂(μ -C₂O₄)₃] with the 3D hyperhoneycomb anion.^{6,12,13} [(C₃H₇)₃NH]₂[Cu₂(μ -C₂O₄)₃] \cdot 2.2H₂O has been suggested to be a candidate



Scheme 1 Three coordination isomers of [Cu₂(C₂O₄)₃]²⁻: 0D (top left), 2D honeycomb (top right), and 3D hyperhoneycomb (bottom).

compounds. However, the crystal structures of **1** and **2** are different. Comparing the cell volumes of **1** and **2** at 290 K, the 200 Å³ difference is more than the volume of two CH₃ groups. This means that there is tight packing of the organic and inorganic building blocks inside the **1** crystal.

In **1**, the donor molecules stack in a quasi-face-to-face mode along the $[-110]$ direction with the sequence of $-A-A-B-$ (Fig. 2b). Although the stack is composed of a trimer unit, there is no displacement, and the stack is practically uniform. The in-plane displacement between the neighbouring donor molecules is almost half a molecule along the b axis. There are C–H···S interactions between the donor molecules.

The donor columns are arranged side-by-side along the $[1-10]$ direction, where there are S···S contacts and C–H···S interactions between the columns. The donor arrangement belongs to the β'' phase, as observed in the organic superconductors β'' -(BEDT-TTF)₂I₃, and β'' -(BEDT-TTF)₄[(H₃O)Fe(C₂O₄)₃]₂C₆H₅CN.^{4,18} This is different from **2** and θ^{21} -(BETS)₃[Cu₂(μ -C₂O₄)₃]₂CH₃OH, where the donor arrangement belongs to the θ^{21} phase (Fig. S3, ESI†).⁶ Considering the bond length of the TTF core with the standard deviation of 0.1 of δ , the oxidation state of BEDT-TTF is in basic agreement with the +2/3 average charge (Table 1), and it remained the same at 180 and 290 K (Table S2, ESI†).¹⁹

The C=C stretching frequency of the charge-transfer complexes of BEDT-TTF is a powerful way to determine the oxidation state. In the Raman spectrum of **1** (Fig. 3), the ν_2 mode is observed at 1488 cm⁻¹. This is the same as that of charge-transfer salts with BEDT-TTF^{+2/3}. The formal charge is deduced to be 0.66.²⁰

The [Cu₂(μ -C₂O₄)(C₂O₄)₂(CH₃OH)(H₂O)²⁻] anion is a 0D oxalate-bridged binuclear anion (Fig. 2c). The oxalate-bridged 0D Cu₂(μ -C₂O₄)(C₂O₄)₂²⁻ unit is one of the coordination isomers of [Cu₂(μ -C₂O₄)₃²⁻], which have been reported to have 2D honeycomb and 3D hyperhoneycomb lattices (Scheme 1). In the anion, the Cu²⁺ atoms are square-pyramidally coordinated to two O atoms from the bisbidentate oxalate anion and two O atoms from the bidentate oxalate anion in the basal plane, and the O atom from H₂O or CH₃OH from the apical position. It is different from the 0D oxalate-bridged dinuclear anion [Fe₂(μ -C₂O₄)(C₂O₄)₄⁴⁻], and the Fe³⁺ atoms are octahedrally coordinated with two O atoms from the bisbidentate oxalate anions and four O atoms from the bidentate oxalate anions.⁷ In the basal plane, the Cu–O distances are 1.930(4)–2.001(3) Å. The metal atoms lie in the main plane of the bidentate oxalate anion. The planes of the two bidentate oxalate anions are staggered from the plane of the bisbidentate oxalate anion with an angle of 10.3(2)° in a “chair” conformation. The Cu–O distance from Cu to CH₃OH and H₂O is 2.231(4) Å. There are O–H···O hydrogen bonds between the anions, from H₂O/CH₃OH to uncoordinated O atoms on the bidentate oxalate anion (Fig. 2c). Because the methanol CH₃ group is oriented outside of the anion with occupation of 0.5, the occupation of the inside H atom is 1.0, and the occupation of another H atom of H₂O with the same orientation as the CH₃ group is 0.5. The hydrogen bond between the oxalate anion and apical ligand is along the a axis, as shown by thick dashed blue lines in Fig. 2c.

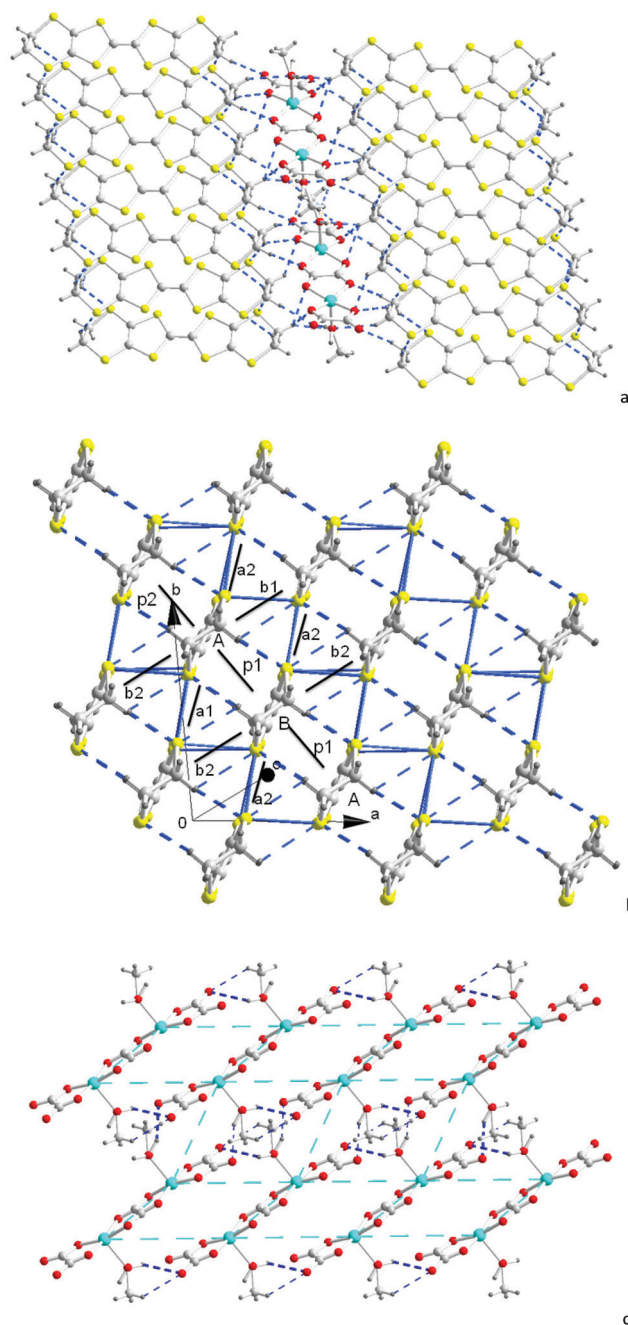


Fig. 2 Crystal structure of **1**: (a) packing diagram viewed along the a axis, (b) donor arrangement viewed along the c axis with the definition of the overlap integrals, and (c) anion arrangement viewed along the c axis. Colour code: Cu, cyan; O, red; S, yellow; C, grey; H, dark grey; hydrogen bonds, blue dash lines; S···S contacts, blue solid lines. The dashed cyan lines show the hydrogen-bonded 4×4 grid. The transfer integrals (black solid lines) are: p1, -43.8 ; p2, -29.8 ; a1, 175.8 ; a2, 176.4 ; b1, -82.8 ; b2, -96.1 meV.

The disordered hydrogen bonds in the anion sheet form a supramolecular square lattice of metal atoms. The Cu···Cu distances are 5.215 Å within the binuclear anion, and 6.690 and 7.480 Å between the Cu atoms connected through an O–H···O hydrogen bond. When the temperature is decreased from 300

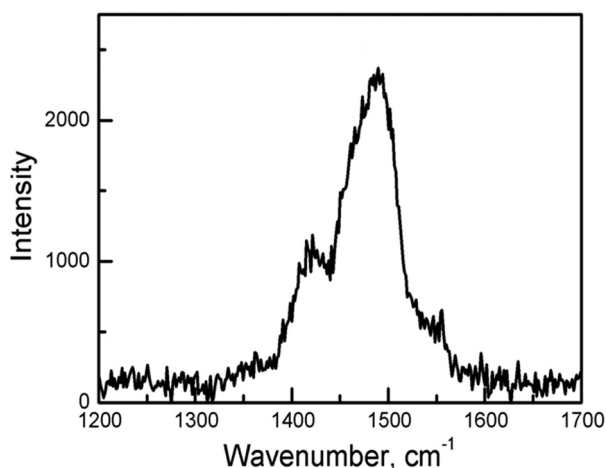
Table 1 Formal charges of the BEDT-TTF molecules at 100 K



$$\delta = (b+c)-(a+d)$$

$$Q = 6.347 - 7.463\delta$$

	<i>a</i>	<i>B</i>	<i>c</i>	<i>d</i>	δ	<i>Q</i>
A	1.368	1.733	1.750	1.355	0.756	0.705
		1.735	1.748	1.358		
		1.733	1.746			
		1.731	1.746			
B	1.372	1.730	1.744	1.359	0.745	0.787
		1.730	1.748			
		1.730	1.748			
Total						2.2

Fig. 3 Raman spectra of the best developed surfaces of single crystals of **1** with $\lambda = 514.5$ nm at room temperature.

to 100 K, the Cu–O distances on the square plane elongate, the apical Cu–O distance contracts, and the Jahn–Teller distortion remains.

There are C–H...O hydrogen bonds between the donor and the anion (Table 2). The S...X and Se...X contacts between the donor and anion, which are observed in charge-transfer salts composed of an organic donor and inorganic anion 0D mononuclear CuCl_4^{2-} , FeCl_4^- , FeBr_4^- , dinuclear $[\text{Cu}_2(\mu\text{-Cl})_2\text{Cl}_4^{2-}]$ and 1D $[\text{Fe}(\mu\text{-C}_2\text{O}_4)\text{Cl}_2]_n$, do not exist.^{1,8,21}

The room-temperature conductivity of **1** is 11.5 S cm^{-1} . When the temperature decreases, the conductivity increases, similar to that of a metal, and reaches 160 S cm^{-1} at 7.6 K, and it then slowly decreases to 150 S cm^{-1} at 2 K (Fig. 4). This is different from charge-transfer salts with the same donor arrangement: $\beta''\text{-(BEDT-TTF)}_3(\text{HSO}_4)_2$, $\beta''\text{-(BEDT-TTF)}_3(\text{ReO}_4)_2$, and $\beta''\text{-(BEDT-TTF)}_3(\text{ClO}_4)_2$, which show a metal–insulator transition at around 150, 100, and 170 K, respectively, accompanied by charge ordering.²² This is different from the reported charge-transfer salts of 0D oxalate-bridged iron dinuclear anions, insulator $\text{TTF}_5[\text{Fe}_2(\mu\text{-C}_2\text{O}_4)(\text{C}_2\text{O}_4)_4] \cdot 2\text{C}_6\text{H}_5\text{CH}_3 \cdot 2\text{H}_2\text{O}$,

Table 2 C–H...O hydrogen bonds between the donor and anion

D–H...A	H...A, Å	D...A, Å	$\angle \text{D–H...A}, ^\circ$
C7–H7A...O	2.66	3.333(5)	125.6
C8–H8B...O5	2.65	3.311(5)	124
C9–H9A...O3	2.38	3.241(5)	145.7
C10–H10A...O2	2.54	3.234(5)	127.2
C10–H10A...O6	2.57	3.333(5)	133.6
C14–H14B...O1	2.43	3.218(5)	136.6
C15–H15B...O4	2.40	3.388(5)	172.8

Fig. 4 Temperature dependence of the conductivity of **1**. The inset shows the resistance in the low-temperature range.

semiconductor $(\text{BEDT-TTF})_4[\text{Fe}_2(\mu\text{-C}_2\text{O}_4)(\text{C}_2\text{O}_4)_4]$ and $\text{TMTTF}_4[\text{Fe}_2(\mu\text{-C}_2\text{O}_4)(\text{C}_2\text{O}_4)_4]$, and $\text{C}_6\text{H}_5\text{CN} \cdot 4\text{H}_2\text{O}$ and 2D oxalate-bridged anions, semiconductor **2** and conductor $\theta^{21}\text{-(BETS)}_3[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3] \cdot 2\text{CH}_3\text{OH}$ with a metal–semiconductor transition at 150 K.⁶

To understand the conductivity behaviour, the electronic structure was analyzed. Because there are no strong interactions between the donor and anion as observed in $\lambda\text{-(BETS)}_2\text{FeCl}_4$, $\kappa\text{-BETS}_2\text{FeBr}_4$ and $\text{TTF}_2\text{Fe}(\mu\text{-C}_2\text{O}_4)\text{Cl}_2$, we can ignore the $\pi\text{-d}$ interaction.^{1,21} The electronic structure of **1** is determined by the organic donor layer, similar to $\beta\text{-(BEDT-TTF)}_3[\text{CrMn}(\mu\text{-C}_2\text{O}_4)_3] \cdot x\text{CH}_2\text{Cl}_2$, **2** and $\theta^{21}\text{-(BETS)}_3[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3] \cdot 2\text{CH}_3\text{OH}$.^{5,6} The band structure was calculated from the transfer integrals between the donor molecules using the tight-binding approach (Fig. 2b).²³ The resulting band structure shows strong dispersion (Fig. 5). The Fermi surface consists of a hole pocket around the Y point and an electron pocket around the Γ point, but the global feature is derived from a large elliptical surface associated with the very small trimerization. This energy band agrees with the metallic conductivity, as expected from the conductivity measurements above 7.6 K.

In charge-transfer salts with metal-oxalate anions, the magnetic property is always controlled by the inorganic anion, and it varies from paramagnetic 0D $\text{Fe}(\text{C}_2\text{O}_4)_3^{3-}$, antiferromagnetic 0D $[\text{Fe}_2(\mu\text{-C}_2\text{O}_4)(\text{C}_2\text{O}_4)_4]^{4-}$, and antiferromagnetic LRO for 1D $[\text{Fe}(\mu\text{-C}_2\text{O}_4)\text{Cl}_2]_n$ to ferromagnetic LRO for 2D $[\text{CrMn}$



Fig. 5 Band structure and Fermi surface of **1**.

$(\mu\text{-C}_2\text{O}_4)_3^-]_n$ to spin frustration for 2D $[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3^{2-}]_n$.^{4-7,21} By empirical orbital analysis, the interaction between the Cu spins within the dinuclear anions can be analyzed.⁹ The magnetic orbital ($d_{x^2-y^2}$) on two metal atoms is coplanar, so a strong antiferromagnetic interaction is expected. The temperature-dependent susceptibility was measured under 10 000 Oe. The Pauli paramagnetism from the conducting unit in **1** should be in the same range as observed in charge-transfer complexes with ET_2^{3+} . For example, in organic conductors $\beta''\text{-(BEDT-TTF)}_3(\text{ClO}_4)_2$ and $\beta''\text{-(BEDT-TTF)}_3(\text{HSO}_4)_2$, the Pauli paramagnetism is $6.5\text{--}7.5 \times 10^{-4} \text{ cm}^3 \text{ mol}^{-1}$.²⁴ At 300 K, χT is $0.408 \text{ cm}^3 \text{ K mol}^{-1}$ after Pauli paramagnetism was subtracted (Fig. 6). It is similar to the previously reported value in honeycomb copper-oxalate-framework compounds with Jahn–Teller distortion at 300 K [**2**, $0.43 \text{ cm}^3 \text{ K mol}^{-1}$; $[(\text{C}_3\text{H}_7)_3\text{NH}]_2[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3] \cdot 2.2\text{H}_2\text{O}$, $0.465 \text{ cm}^3 \text{ K mol}^{-1}$], and a square copper-format-framework compound at 300 K [$\theta\text{-(BEDT-TTF)}_2\text{Cu}_2(\mu\text{-HCOO})_5$, $0.494 \text{ cm}^3 \text{ K mol}^{-1}$].^{6,11,25}

With a decreasing temperature, the χ value smoothly increases with a broad shoulder at around 190 K. Below 190 K, it decreases and reaches a minimum at 55 K. The χ value then increases and reaches $0.008 \text{ cm}^3 \text{ mol}^{-1}$ at 2 K. From the broad maximum around 190 K, it is an antiferromagnetic system, and the antiferromagnetic interaction J could be estimated from the empirical function $|J|/kT_{\text{max}} = 1.599$ as -304 K .

The temperature dependent susceptibility data were analyzed by a “singlet–triplet” model with $H = -JS_{\text{A}}S_{\text{B}}$, with

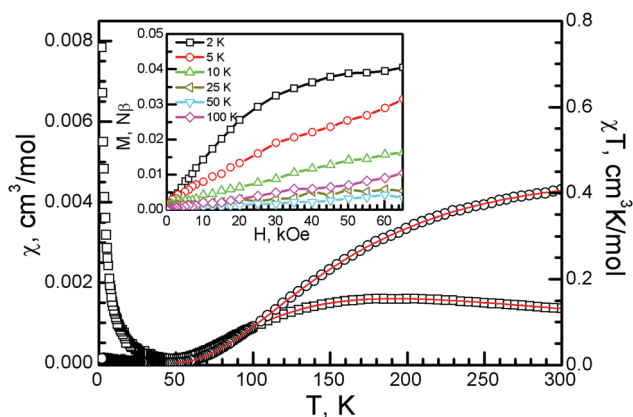


Fig. 6 Temperature dependence of the magnetic susceptibility, χ (squares) and χT (circles), together with the dimer model (red lines). The inset shows the isothermal magnetization at 2, 5, 10, 50 and 100 K.

intermolecular interaction between dinuclear units, which was used for oxalate-bridged copper dinuclear compounds.²⁶ The equation was used to fit the results above 50 K, giving $g = 2.01$ (1), $J = -323$ (1) K, and $R = 1.56 \times 10^{-7}$ ($R = \sum(\chi_0 - \chi_c)^2 / \sum \chi_0^2$). This shows that the singlet–triplet energy gap of **1** is above 200 cm^{-1} , which is in the same range as that of the reported oxalate-bridged dinuclear Cu(II) compounds.²⁶

The isothermal magnetization was measured at temperatures from 2 to 300 K (Fig. 6, inset). The isothermal magnetization smoothly increased and reached $0.0395 \text{ N}\beta$ at 65 kOe and 2 K, which is much smaller than the value of $1 \text{ N}\beta$ for isolated, spin only Cu^{2+} with $S = 1/2$ and $g = 2.0$. The magnetization linearly increased with increasing magnetic field and reached $0.0033 \text{ N}\beta$ at 50 K and $0.0104 \text{ N}\beta$ at 100 K. These values confirm the strong antiferromagnetic interaction between the oxalate-bridged metal atoms. They are in the same range as those of the 3D hyperhoneycomb compound $[(\text{C}_2\text{H}_5)_3\text{NH}]_2[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3]$ ($0.0351 \text{ N}\beta$ at 2 K, $0.0304 \text{ N}\beta$ at 5 K, and $0.0145 \text{ N}\beta$ at 300 K and 65 kOe) and the 2D honeycomb compound $[(\text{C}_3\text{H}_7)_3\text{NH}]_2[\text{Cu}(\mu\text{-C}_2\text{O}_4)_3] \cdot 2.2\text{H}_2\text{O}$ ($0.044 \text{ N}\beta$ at 2 K and 65 kOe).¹² They are also comparable with those of the oxalate-bridged binuclear copper compound $\{\text{tmen}(2\text{-MeIm})\text{Cu}\}_2(\text{C}_2\text{O}_4)(\text{PF}_6)_2$, for which the isothermal magnetization has been estimated to be about $0.05 \text{ N}\beta$ at 70 kOe.²⁷ Zero-field-cooled, field-cooled, and remnant magnetization measurements from 2 to 100 K under 100 Oe showed no bifurcation (Fig. S3, ESI†).

Specific heat experiments were performed from 2 to 120 K under 0 and 5 T (Fig. S4, ESI†). No λ -peak was observed between 2 and 120 K. Combining with the X-ray diffraction experiment from room temperature to 100 K, there is neither magnetic ordering nor a structural transition between 2 and 290 K.

Because the susceptibility measurements were performed with selected high-quality crystals, the antiferromagnetic interaction should be the intrinsic behaviour of the crystal. The observed magnetic properties are determined not only by the oxalate-bridged binuclear anion, but also by the hydrogen-bonded square lattice anion. Because the magnetic interaction through hydrogen bonds is around $1\text{--}10 \text{ cm}^{-1}$ and much weaker than the oxalate-bridge, the magnetic interactions between the anions should be much weaker than the intradimer interaction.²⁷ The magnetic properties of **1** are controlled by the oxalate-bridged binuclear anion sheet.

Conclusions

$\beta''\text{-(BEDT-TTF)}_3[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)(\text{C}_2\text{O}_4)_2(\text{CH}_3\text{OH})(\text{H}_2\text{O})]$ is the first magnetic organic conductor incorporating oxalate-bridged binuclear Cu(II) anions. It shows metallic conductivity down to 2 K and proves that modification of the donor arrangement improves the conductivity from its polymorphic compound $\theta^{21}\text{-(BEDT-TTF)}_3[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3] \cdot 2$ solvent rather than by replacement of S by Se on the TTF core to obtain the isostructural compound $\theta^{21}\text{-(BETS)}_3[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3] \cdot 2$ solvent. The antiferromagnetic interaction in the hydrogen-bonded oxalate-bridged

Cu(II) binuclear anion square lattice is stronger than that in the $[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)_3]^{2-}_n$ honeycomb lattice. No long range order is observed above 2 K. $\beta''\text{-(BEDT-TTF)}_3[\text{Cu}_2(\mu\text{-C}_2\text{O}_4)(\text{C}_2\text{O}_4)_2(\text{CH}_3\text{OH})(\text{H}_2\text{O})]$ is a new member of magnetic conductors.

Author contributions

The manuscript was written through contributions from all authors. B. Zhang and D. Zhu managed the project. B. Zhang and Y. Zhang synthesised the samples. Y. Zhang performed magnetic measurements. Z. Wang, Z. Gao and G. Chang carried out single crystal X-ray diffraction experiments. Y. Zhang and D. Wang carried out conductivity and heat capacity experiments. D. Yang performed the EDS measurements. G. Yan recorded the Raman spectra. Z. Zhao and F. Liu carried out XPS experiments. T. Mori carried out ESR experiments and calculated the band structure. Q. Li carried out elemental analysis experiments. B. Zhang conducted experiments and analyzed the data. B. Zhang, T. Mori, Y. Zhang and Z. Wang wrote the main manuscript text. All authors have given approval to the final version of the manuscript.

Abbreviations

BEDT-TTF bis(ethylenedithio)tetrathiafulvalene.
BETS bisethylene(tetraselenafulvalene).

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- P. Day, M. Kurmoo, T. Mallah, I. Marsden, R. Friend, F. Pratt, W. Hayes, D. Chasseau, J. Gaultier, G. Bravis and L. Ducasse, *J. Am. Chem. Soc.*, 1992, **114**, 10722; P. Cassoux, *Science*, 1997, **272**, 1277; H. Kobayashi, H. Tomita, T. Naito, A. Kobayashi, F. Sakai, T. Watanabe and P. Cassoux, *J. Am. Chem. Soc.*, 1996, **118**, 368; F. Palacio and J. S. Miller, *Nature*, 2000, **408**, 421; H. Kobayashi, H. Cui and A. Kobayashi, *Chem. Rev.*, 2004, **104**, 5265.
- T. R. Felthouse, E. J. Laskowski and D. N. Hendrickson, *Inorg. Chem.*, 1977, **16**, 1077; M. Julve, M. Verdaguer, A. Gleizes, M. Philoche-Levisalles and O. Kahn, *Inorg. Chem.*, 1984, **23**, 3808; O. Kahn, *Angew. Chem., Int. Ed. Engl.*, 1985, **24**, 834; M. Julve, A. Gleizes, L. M. Chamoreau, E. Ruiz and M. Verdaguer, *Eur. J. Inorg. Chem.*, 2018, 509 and references cited therein.
- H. Tamaki, Z. J. Zhong, N. Matsumoto, S. Kida, M. Koikawa, N. Achiwa, Y. Hashimoto and H. Okawa, *J. Am. Chem. Soc.*, 1992, **114**, 6974; C. Mathoniere, S. G. Carling, D. Yusheng and P. Day, *J. Chem. Soc., Chem. Commun.*, 1994, 1551.
- A. W. Graham, M. Kurmoo and P. Day, *J. Chem. Soc., Chem. Commun.*, 1995, 2061; P. Cassoux, *Science*, 1997, **272**, 1277; S. Benmansour and C. J. Gomez-Garcia, *Magnetochemistry*, 2021, **7**, 93.
- E. Coronado, J. R. Galan-Mascaros, C. J. Gomez-Garcia and V. Laukhin, *Nature*, 2000, **408**, 447; A. Albertola, E. Coronado, J. R. Galan-Mascaros, C. Gimenez-Saiz and C. J. Gomez-Garcia, *J. Am. Chem. Soc.*, 2003, **125**, 10774.
- B. Zhang, Y. Zhang and D. Zhu, *Chem. Commun.*, 2012, **48**, 197; B. Zhang, Y. Zhang, Z. Wang, S. Gao, Y. Guo, F. Liu and D. Zhu, *CrystEngComm.*, 2013, **15**, 3529.
- M. Clemente-Leon, E. Coronado, J. Galan-Mascaros, C. Gimenez-Saiz, C. Gomez-Garcia and J. Fabre, *Synth. Met.*, 1999, **103**, 2279; S. Rashid, S. Turner, P. Day, M. Light and M. Hursthouse, *Inorg. Chem.*, 2000, **39**, 2426; E. Coronado, J. Galan-Mascaros and C. Gomez-Garcia, *J. Chem. Soc., Dalton Trans.*, 2000, **2**, 205.
- P. Day, M. Kurmoo, T. Mallah, I. Marsden, R. Friend, F. Pratt, W. Hayes, D. Chasseau, J. Gaultier, G. Bravis and L. Ducasse, *J. Am. Chem. Soc.*, 1992, **114**, 10772; H. Kujiwara, E. Fujiwara, Y. Nakazawa, B. Narymbetov, K. Kato, H. Kobayashi, A. Kobayashi, M. Tokumoto and P. Cassoux, *J. Am. Chem. Soc.*, 2001, **123**, 306; B. Zhang, H. Tanaka, H. Fujiwara, H. Kobayashi, E. Fujiwara and A. Kobayashi, *J. Am. Chem. Soc.*, 2002, **124**, 9982.
- O. Kahn, *Angew. Chem.*, 1985, **24**, 834; O. Kahn, O. Kahn, *Molecular Magnetism*, VCH Publisher, Inc., 1993, p. 103; J. Cano, P. Alemany, S. Alvarez, M. Verdaguer and E. Ruiz, *Chem. – Eur. J.*, 1998, **4**, 476.
- B. Zhang, Y. Zhang, Z. Wang, D. Wang, P. Baker, F. Pratt and D. Zhu, *Sci. Rep.*, 2014, **4**, 06451; B. Zhang, P. Baker, Y. Zhang, D. Wang, Z. Wang, S. Su, D. Zhu and F. Pratt, *J. Am. Chem. Soc.*, 2018, **140**, 122.
- A. C. Jacko and B. J. Powell, *arXiv*: 1805.05495.
- B. Zhang, Y. Zhang and D. Zhu, *Dalton Trans.*, 2012, **14**, 8509.
- L. Balents, *Nature*, 2010, **464**, 199.
- P. W. Anderson, *Science*, 1987, **235**, 1196.
- CrysAlisPro 1.171.39.9f*, Rigaku OD, 2015.
- G. M. Sheldrick, *SHELX97*, University of Göttingen, Göttingen, Germany, 1997; G. M. Sheldrick, *SHELXL-2014/7*, University of Göttingen, Göttingen, Germany, 2014.
- O. Kahn, *Molecular Magnetism*, VCH Publisher, Inc., 1993, p. 3.
- T. Mori, *Bull. Chem. Soc. Jpn.*, 1998, **71**, 2509; T. Mori, *Electronic Properties of Organic Conductors*, Springer, 2016.
- P. Guionneau, C. J. Kepert, G. Bravis, D. Chasseau, M. R. Truter, M. Kurmoo and P. Day, *Syn. Met.*, 1997, **86**, 1973.

- 20 H. H. Wang, J. R. Ferraro, J. M. Williams, U. Geiser and J. A. Schlueter, *J. Chem. Soc., Chem. Comm.*, 1994, 1893; T. Yamamoto, M. Uruichi, K. Yamamoto, K. Yakushi, A. Kawamoto and H. Taniguchi, *J. Phys. Chem. B.*, 2005, **109**, 15226.
- 21 A. Kobayashi, A. Sato, A. Arai, H. Kobayashi, C. Faulmann, N. Kushch and P. Cassoux, *Solid State Commun.*, 1997, **103**, 371; B. Zhang, Z. Wang, H. Fujiwara, H. Kobayashi, M. Kurmoo, K. Inoue, T. Mori, S. Gao, Y. Zhang and D. Zhu, *Adv. Mater.*, 2005, **17**, 1988; B. Zhang, Z. Wang, Y. Zhang, K. Takahashi, Y. Okano, H. B. Cui, H. Kobayashi, K. Inoue, M. Kurmoo, F. L. Pratt and D. Zhu, *Inorg. Chem.*, 2006, **45**, 3275; B. Zhang, Y. Zhang, Z. Gao, G. Chang, S. Su, D. Wang, Y. Guo and D. Zhu, *Eur. J. Inorg. Chem.*, 2014, 4028; B. Zhang, Y. Zhang, Z. Wang, D. Yang, Z. Gao, D. Wang, Y. Guo, D. Zhu and T. Mori, *Dalton Trans.*, 2016, 45, 16561.
- 22 T. Mori, *Chem. Rev.*, 2004, **104**, 4947; N. Takubo, N. Tajima, H. Yamamoto, H. Cui and R. Kato, *Phys. Rev. Lett.*, 2013, **110**, 227401.
- 23 T. Mori, A. Kobayashi, Y. Sasaki, H. Kobayashi, G. Saito and H. Inokuchi, *Bull. Chem. Soc. Jpn.*, 1984, **57**, 627.
- 24 B. Rothaemel, L. Forro, J. R. Cooper, J. S. Schilling, M. Weger, P. Bele, H. Brunner, D. Schweitzer and H. J. Keller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1986, **34**, 704; A. Miyazaki, T. Enoki, H. Uekusa, Y. Ohashi and G. Saito, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1997, **55**, 6847.
- 25 B. Zhang, Y. Zhang, Z. Wang, Z. Gao, D. Yang, D. Wang, Y. Guo and D. Zhu, *ChemistryOpen*, 2017, **6**, 320.
- 26 B. Bleaney and K. Bowers, *Proc. R. Soc. London*, 1952, **214**, 451; Y. Journaux, J. Sletten and O. Kahn, *Inorg. Chem.*, 1985, **24**, 4063; L. Soto, J. Garcia, E. Escriba, J. Legros, J. Tuchages, F. Dahan and A. Fuertes, *Inorg. Chem.*, 1989, **28**, 3378; O. Kahn, *Molecular Magnetism*, VCH Publisher, Inc., 1993, p. 104; P. Bergerat, O. Kahn, P. Legoll, M. Drillon and M. Guillot, *Inorg. Chem.*, 1994, **33**, 2049.
- 27 F. M. Romero, R. Ziessel, M. Bonnet, Y. Pontillon, E. Ressouche, J. Schweizer, B. Delley, A. Grand and C. Paulsen, *J. Am. Chem. Soc.*, 2000, **122**, 1298.