



Cite this: *Phys. Chem. Chem. Phys.*,
2022, 24, 1869

Formation of phenylacetylene and benzocyclobutadiene in the *ortho*-benzyne + acetylene reaction†

Morgan N. McCabe, ^a Patrick Hemberger, ^b Dario Campisi, ^{‡c}
Jeger C. Broxterman, ^a Engelbert Reusch,^d Andras Bodi ^b and
Jordy Bouwman  §¶||^{*a}

Ortho-benzyne is a potentially important precursor for polycyclic aromatic hydrocarbon formation, but much is still unknown about its chemistry. In this work, we report on a combined experimental and theoretical study of the *o*-benzyne + acetylene reaction and employ double imaging threshold photoelectron photoion coincidence spectroscopy to investigate the reaction products with isomer specificity. Based on photoion mass-selected threshold photoelectron spectra, Franck–Condon simulations, and ionization cross section calculations, we conclude that phenylacetylene and benzocyclobutadiene (**PA: BCBdiene**) are formed at a non-equilibrium ratio of 2:1, respectively, in a pyrolysis microreactor at a temperature of 1050 K and a pressure of ~20 mbar. The C₈H₆ potential energy surface (PES) is explored to rationalize the formation of the reaction products. Previously unidentified pathways have been found by considering the open-shell singlet (OSS) character of various C₈H₆ reactive intermediates. Based on the PES data, a kinetic model is constructed to estimate equilibrium abundances of the two products. New insights into the reaction mechanism – with a focus on the OSS intermediates – and the products formed in the *o*-benzyne + acetylene reaction provide a greater level of understanding of the *o*-benzyne reactivity during the formation of aromatic hydrocarbons in combustion environments as well as in outflows of carbon-rich stars.

Received 12th November 2021,
Accepted 23rd December 2021

DOI: 10.1039/d1cp05183k

rsc.li/pccp

1 Introduction

Polycyclic aromatic hydrocarbons (PAHs), organic molecules that contain two or more fused aromatic rings, are abundant on Earth and ubiquitous in the interstellar medium (ISM).

Terrestrially, these molecules have natural and anthropogenic sources (*via* combustion) and are of concern as their accumulation in the environment can lead to health issues.^{1,2} In the ISM, PAHs are detected as a family of molecules by the characteristic emission bands in the 3–20 μm mid-infrared (mid-IR) spectral region.^{3–5} These mid-IR emission bands have been reported towards many sources including H II regions, planetary nebula, and reflection nebula.⁶

Numerous studies have been conducted to simulate combustion environments and to reveal the PAH and soot formation mechanisms.^{1,7–12} These combustion studies also provide a basis for understanding PAH formation in the ISM, as it is assumed that PAHs are formed under the combustion-like conditions in the outflow of carbon-rich stars.^{5,13} However, the formation of the second aromatic ring, a critical step in the growth of aromatic molecules, has yet to be explained.¹⁴ New mechanisms, such as reactions involving *o*-benzyne, may contribute to growth of PAHs and the formation of a second ring.

The structure and reactivity of *o*-benzyne have been studied in some detail. Spectroscopic measurements characterized *o*-benzyne's structure, with IR stretching modes indicating the presence of a carbon–carbon triple bond, albeit with a lower

^a Laboratory for Astrophysics, Leiden Observatory, Leiden University, PO Box 9513, 2300 RA Leiden, The Netherlands. E-mail: jordy.bouwman@colorado.edu;

Fax: +31 7 1527 8432

^b Laboratory for Synchrotron Radiation and Femtochemistry, Paul Scherrer Institute, 5232 Villigen, Switzerland

^c Leiden Observatory, Leiden University, Niels Bohrweg 2, 2333 CA Leiden, The Netherlands

^d Institute of Physical and Theoretical Chemistry, University of Würzburg, Am Hubland, D-97074 Würzburg, Germany

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d1cp05183k

‡ Present address: Chemistry Department, The University of Chicago, 5735 S Ellis Ave, Chicago, IL 60637, USA.

§ Present address: Laboratory for Atmospheric and Space Physics, University of Colorado, Boulder, CO 80303, USA.

¶ Present address: Department of Chemistry, University of Colorado, Boulder, CO 80309, USA.

|| Present address: Institute for Modeling Plasma, Atmospheres, and Cosmic Dust (IMPACT), University of Colorado, Boulder, CO 80303, USA.

frequency than the average for the bond type.¹⁵ Further, computational studies confirmed that the species largely exhibits arylene character, due to the presence of a large singlet-triplet energy gap, with the singlet species being lower in energy. However, *o*-benzynes retain some small amount (*ca.* 10%) of biradical character, as confirmed by rotational spectroscopy.^{16,17} Recently, *o*-benzyne was detected toward the Taurus Molecular Cloud TMC-1,¹⁸ marking its first ever detection in the ISM. In addition, *o*-benzyne has recently awakened the interest of the combustion chemistry community, as it was shown that phenyl radicals can decompose to *o*-benzyne under combustion-relevant conditions.¹⁹ Moreover, *o*-benzyne self-reactions were shown to result in PAHs.²⁰ Due to its presence in the ISM and its potential role in combustion chemistry, *o*-benzyne may play a significant role in closing the gap between current PAH and soot formation models and the observed PAH abundance and rates of formation.^{14,21,22}

A few studies have aimed to understand gas-phase *o*-benzyne reactions with both closed-shell and radical reactants.^{23–26} A combined experimental and computational study by Friedrichs *et al.*²⁷ probed the products of *o*-benzyne's reaction with the small unsaturated hydrocarbons acetylene (C₂H₂), ethene (C₂H₄), and propene (C₃H₆). Computationally, they found 40–50 kJ mol^{−1} entrance barriers to the *o*-benzyne + hydrocarbon reactions. Furthermore, while their mass spectrometric approach did not allow for isomer-specific product assignment, they concluded based on computational results that both polycyclic and branched monocyclic reaction products are accessible. In addition, Friedrichs *et al.* proposed an “edge-on” and a “concerted” pathway for the *o*-benzyne + acetylene association reaction with the final product dependent on the favored angle of attack. In the absence of isomer-specific data, the dominant product formation mechanism remained unclear.

Computational studies to determine reaction pathways and accurate energies for reactions with *o*-benzyne are uniquely challenging. Specifically, the reactions between *o*-benzyne and hydrocarbons are likely to involve biradical intermediates with open-shell singlet (OSS) character, and as a result, care needs to be taken to describe these systems appropriately. Closed-shell single-determinant wave functions are insufficient to represent the wave function of these molecules because of the presence of independent radical centers. Potential energy surface exploration using multireference approaches with large enough basis sets is an intractable pursuit in these comparably large systems. Luckily, studies have shown that OSS energies can also be approximated based on OSS and triplet energies quite reliably, allowing for a straightforward density functional theory (DFT) description of these species.^{28–30}

The aim of this work is to determine the reaction mechanism and final isomeric products of the acetylene + *o*-benzyne reaction by double imaging photoelectron photoion spectroscopy³¹ (*i*²PEPICO) at the Swiss Light Source. Photoion mass-selected threshold photoelectron spectra offer an isomer-selective and sensitive detection tool to identify reactive intermediates in gas mixtures, and reveal reaction mechanisms, as also applied in catalysis and combustion environments.^{32,33} DFT computational

chemistry methods are employed to yield insights into the potential energy surface and the reaction mechanism at play taking into account the OSS character of the biradical species. Finally, statistical modelling is performed to reveal the isomerization reaction dynamics and put the measured isomer branching ratios into context. With the combined application of these methods, we identify the products of the acetylene + *o*-benzyne reaction and obtain a clear understanding of the reaction mechanism.

2 Experimental

The *o*-benzyne + acetylene reaction was studied using a pyrolysis microreactor connected to the CRF-PEPICO setup at the vacuum ultraviolet (VUV) beamline of the Swiss Light Source (SLS) at Paul Scherrer Institute (PSI) in Villigen, Switzerland. The VUV beamline, the *i*²PEPICO detection chamber^{34,35} and the pyrolysis reactor³⁶ have been described in much detail elsewhere and only a summary of the relevant experimental details is provided here.

Benzocyclobutenedione (**BCBdione**) was synthesized following the method by South and Liebeskind³⁷ and used as *o*-benzyne precursor. The solid **BCBdione** sample was placed in between a small aluminum pellet and glass wool in a 0.25 inch diameter sample container tube. The tube had a 100 μm pinhole at the end, situated at the entrance of the pyrolysis microreactor. The sample container tube was mounted in a copper block, of which the temperature was stabilized at 60 °C using water cooling and heating with a Huber minichiller. As shown in a previous study, the vapor pressure of **BCBdione** is sufficient at this temperature to create a stable *o*-benzyne signal, while being low enough to suppress bimolecular *o*-benzyne chemistry.²⁵ The reactant gas, neat acetylene (≥99.5% stabilized with acetone, Carbagas), flowed through the sample container and picked up the **BCBdione** vapor.

The resulting gas-phase mixture of acetylene and **BCBdione** expanded through the pinhole at the end of the sample container tube and into the pyrolysis microreactor. The microreactor consists of a ~4 cm long SiC tube with an inner diameter of 1 mm and is resistively heated over a length of ~2 cm using DC power. The temperature was estimated from the electrical power, based on previous surface temperature measurements as recorded using a thermocouple for a similar SiC reactor tube with equidistant electrodes. Temperatures were varied between 700–1200 K with an absolute uncertainty on the order of ±100 K, but with a small relative error between measurements.

The gas pressure in the sample tube containing the precursor was measured at 0.25–0.35 bar while flowing acetylene at a rate of 30 sccm. The gas expands through the pyrolysis reactor into the source chamber. The pressure in the source chamber was 5 × 10^{−4} mbar during gas flow. The pressure and residence time in the reactor tube can be inferred from these conditions based on computational fluid dynamics and are on the order of 20 mbar and 10–100 μs, respectively.^{38–40} The effusive beam containing reactants and products exiting the pyrolysis reactor was skimmed by a 2 mm skimmer, forming a molecular beam

in the i^2 PEPICO detection chamber. The background pressure of the detection chamber was $\sim 10^{-6}$ mbar during measurements with a baseline pressure of *ca.* one order of magnitude lower.

Synchrotron radiation was generated using a bending magnet and dispersed by a 150 grooves per mm grating. The dispersed light was focused at the 200 μm exit slit, situated in a differentially pumped rare gas filter and resulting in a resolution of 1:1500. The gas filter was filled with 10 mbar of a mixture of Kr, Ar, and Ne over an optical length of 10 cm to filter out high harmonic radiation of the grating above 14 eV. The monochromatized ionizing VUV synchrotron radiation intersects the molecular beam in the i^2 PEPICO detection chamber. The detection chamber is connected to the beamline, and the ionization region is 50 cm downstream from the exit slit. The electron-ion pairs formed in an ionization event are accelerated in opposite directions using a 218 V cm^{-1} field and are recorded in delayed coincidence. Electrons are velocity map imaged on a RoentDek delay line detector, and their arrival time serves as the start time for the time-of-flight (TOF) measurement of the associated cation. A second RoentDek detector on the opposite end records the position and arrival time of the space focused ions. Coincidence data were recorded over a photon energy range of 7.4–9.4 eV using a step size of 10 meV and an integration time of 60 s per energy point. These data were then used to construct the mass-selected threshold photoelectron spectra (ms-TPES) of the reaction products. Isomer-resolved product assignments are made based on the ms-TPES data in conjunction with reference and/or simulated spectra.

2.1 Computational methods

Wherever possible, isomeric assignments are based on reference threshold photoelectron spectra. Otherwise, for species without a published photoelectron spectrum, adiabatic ionization energies (IEs) were calculated at the CBS-QB3 level of theory using the program Gaussian 16⁴¹ to support the interpretation of the experimental observations. Furthermore, TPE spectra were simulated by calculating Franck–Condon factors in the double harmonic approximation. For the Franck–Condon simulations, geometry optimizations and frequency calculations for the neutral and ionic ground state of the species were conducted with DFT using the B3LYP/6-311++G(d,p)^{42–46} level of theory. Franck–Condon factors were calculated from these vibrational normal modes using ezSpectrum.⁴⁷ The resulting stick spectra were convoluted with a Gaussian function with a full-width-at-half-maximum of 0.032 eV to simulate the rotational envelope and facilitate comparison with the experimental data.

In order to estimate the relative abundances of the isomers present in the ms-TPES, eZDyson was used to calculate absolute photoionization cross sections for each isomer.⁴⁸ Dyson orbitals were obtained using equation of motion for ionization potential with coupled cluster singles and doubles (EOM-IP-CCSD)^{49,50} with the correlation consistent polarized triple zeta basis set (cc-pVTZ)⁵¹ basis set, using Q-Chem 4.3.⁵² The continuum state

of the photoelectron was taken into account to compute photoionization cross sections, which were then combined with Franck–Condon simulations of the nuclear wave function overlap to predict relative TPES intensities of the isomers. Thus, we only rely on relative signal intensities to estimate the relative abundance of the isomers contributing to the photoelectron spectrum.

Potential energy surface (PES) calculations were performed in order to provide insight into the energy landscape of the acetylene + *o*-benzyne reaction. Relaxed coordinate scans were conducted at the M06-2X/6-311++G(d,p)^{46,53} level of theory and stationary points were located. M06-2X was chosen as it provides accurate energies when studying reactions involving large hydrocarbons.^{54–56} Due to the fact that the reaction starts from the closed-shell but in part biradical *o*-benzyne, some of the intermediates exhibit biradical character. To treat this properly using DFT, unrestricted singlet wave functions were calculated for the biradical intermediates by mixing the highest occupied and lowest unoccupied α orbitals in the initial guess of the wave function to destroy spatial symmetry of the α and β sets of orbitals. The wave function was thus allowed to converge to an open-shell singlet solution. The energy of OSS species was determined using the Ziegler–Cramer correction by calculating twice the OSS energy minus the triplet energy at the same geometry.^{29,30} This was done in instances where the mixing of the singlet and triplet states is comparable, as indicated by a spin expectation value $\langle S^2 \rangle \approx 1$. Lastly, Complete Active Space Self-Consistent Field (CASSCF)⁵⁷ calculations were performed along the entrance pathways (see ESI†) to confirm the entrance barriers and therewith obtain insights into the association mechanism at play.

The product distribution in the pyrolysis microreactor is determined by the radical precursor formation, the bimolecular kinetics of the association reaction, the competition between fast re-thermalization and the system exploring the adduct's potential energy surface, as well as the short reaction time, determined by the 10–100 μs residence time in the reactor. This means that non-equilibrium product distributions can result from fast cooling after exothermic reactions. To put the experimentally determined isomer distribution of the products into context, unimolecular microcanonical rate curves were calculated by Rice–Ramsperger–Kassel–Marcus (RRKM) theory using a computer program broadly validated in modeling unimolecular dissociative photoionization and developed by Sztáray *et al.*⁵⁸ The model provides microcanonical rates for unimolecular reactions based on the reactant density of states and the transition state numbers of states *via* the equation:

$$k(E) = \frac{\sigma N^\ddagger(E - E_0)}{h\rho(E)},$$

where $\rho(E)$ is the reactant density of states, $N(E - E_0)$ transition state number of states, h is Planck's constant, and σ is the degeneracy of the reaction channel. We are not addressing the kinetics of the association reaction but consider intermediates on the potential energy surface and calculate the associated isomerization rates.

3 Results

Mass spectra and mass-selected threshold photoelectron spectra for the reactions of *o*-benzynes with acetylene recorded at selected pyrolysis temperatures are presented first. Comparing the TPES data with reference data or with Franck–Condon simulated spectra allows us to identify the species with isomer specificity. Next, the potential energy surface, computed at the M06-2X/6-311++G(d,p) level of theory, is discussed to support the experimental findings.

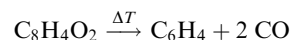
3.1 Mass spectrometry

Mass spectra are depicted in Fig. 1 comparing the products for (A) acetylene only and (B) acetylene + **BCBdione** recorded at a photon energy of 10.0 eV.

As can be seen from Fig. 1A as well as the inset, a number of peaks are present in the mass spectrum of pyrolyzed acetylene. In Fig. 1A, strong signals are seen at *m/z* 58 and *m/z* 43 and are attributed to the acetone ((CH₃)₂CO⁺) and the acetyl cation (CH₃CO⁺), the latter being a dissociative ionization fragment of acetone (see ESI† for further details).⁵⁹ Acetone is present in our gas flow as it is added to acetylene gas bottle as an inhibitor. Acetylene (*m/z* 26) is not seen in the mass spectrum, as it was recorded using a photon energy well below its

ionization potential of 11.4 eV.⁶⁰ In the inset, two peaks can be seen at *m/z* 83 and *m/z* 98 and are confirmed to be trace contaminants present and are not involved in the reactions occurring during pyrolysis (see ESI† for further details).

BCBdione was used as a precursor as it has previously been shown to be a clean source of *o*-benzynes *via*:^{25,37}



Mass spectra of **BCBdione** pyrolysis as a function of temperature show that **BCBdione** can be cleanly converted to *o*-benzynes at 1082 K (see Fig. S1 in ESI†). Fig. 1B shows the mass spectrum after pyrolysis of the *o*-benzynes precursor, **BCBdione**, in presence of acetylene. The inset of Fig. 1B shows two new peaks when **BCBdione** is added to the flow. The peak at *m/z* 76 corresponds to *o*-benzynes and the peak at *m/z* 102 corresponds to the association product formed from the acetylene + *o*-benzynes reaction:



3.2 Threshold photoelectron spectroscopy

The *m/z* 102 ms-TPES was recorded over the 7.4–9.4 eV range with a 10 meV step size to assign the reaction products and is shown in Fig. 2. Two distinct, vibrationally resolved bands are present, the first starting at 7.74 eV and the second at 8.82 eV. Literature reference data were used to assign the vibronic structure starting at 8.82 eV. A clear match is apparent with the previously recorded ms-TPES of phenylacetylene (**PA**) taken from Hemberger *et al.*⁶¹ and shown in purple. The onset also matches well with the previously reported ionization energy of **PA** of 8.825 eV.⁶² Benzocyclobutadiene (**BCBdiene**) may be responsible for the signal at 7.74 eV, with a previously reported IE of 7.87 eV⁶³ obtained by photoelectron spectroscopy. The IE found by Koenig *et al.*⁶³ is notably higher in energy than the onset seen here; the difference in energy is likely due to the lower resolution of the photoelectron spectrum taken at the time.

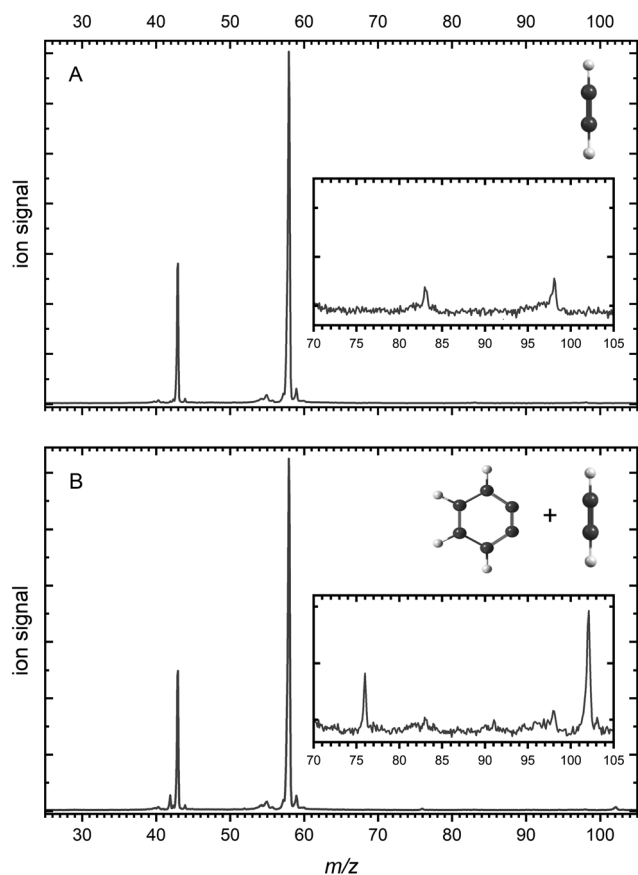


Fig. 1 Mass spectra recorded at a photon energy of 10.0 eV and a microreactor temperature of 1026 K and 996 K (respectively) for (A) neat acetylene and (B) neat acetylene + **BCBdione** (1%).

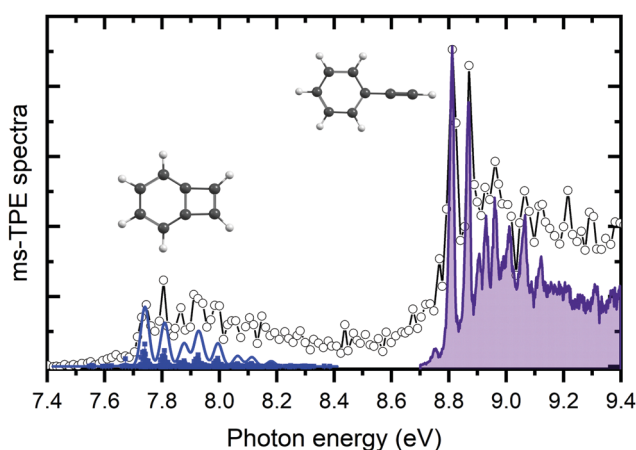


Fig. 2 The *m/z* 102 ms-TPES in the 7.4–9.4 eV photon energy range for the *o*-benzynes (1%) + acetylene reaction at a reactor temperature of 1050 K. The experimental spectrum is fit using a Franck–Condon simulated spectrum for **BCBdiene** and reference data for **PA**.

In order to determine if **BCBdiene** is indeed responsible for the m/z 102 signal, we turn to quantum chemical computations and Franck–Condon simulations to assign the spectrum. Using CBS–QB3 and G4 levels of theory, the adiabatic IE of **BCBdiene** is predicted at 7.75 eV and 7.76 eV, respectively, in very good agreement with the onset of the ms-TPES band. The B3LYP/6-311++G(d,p) optimized ground state structure of the neutral and the cation were subsequently used in a Franck–Condon simulation. The resulting stick spectrum was then convoluted with a 32 meV FWHM Gaussian profile to account for the rotational envelope and the experimental energy resolution to facilitate comparison with the recorded spectrum. From Fig. 2, it is clear that the simulated spectrum of **BCBdiene** matches very well with the vibronic structure seen in the experiment, leading us to conclude that **BCBdiene** is indeed the product from the reaction. A zoom in of the **BCBdiene** ms-TPE spectrum is shown in the ESI,[†] where further information is provided on the vibrational normal modes responsible for the observed progression. From our data we conclude that the ionization energy is 7.74 ± 0.04 eV, which is significantly lower than the value of Koenig *et al.*,⁶³ also cited in the NIST Chemistry Webbook.⁶⁴

Photoionization cross section calculations with ezDyson based on the EOM-IP-CCSD/cc-pVTZ Dyson orbitals show that the ionization cross sections to the ground cation states of **BCBdiene** and **PA** agree to within 3% for electrons with less

than 300 meV kinetic energy. Thus, the relative TPES peak intensities will be determined by the nuclear wave function overlap between the neutral and the final cation rovibrational state and the relative abundance of the two products in the sample stream. Therefore, based on the predicted intensities from the Franck–Condon simulations, we conclude that the products form approximately in a ratio of 2:1 of **PA** to **BCBdiene**, respectively, under our experimental conditions. As will be shown later, the non-equilibrium product distribution indicates that the less stable **BCBdiene** is collisionally stabilized in the microreactor faster than it is converted to the more stable **PA**.

3.3 Potential energy surface calculations

The acetylene + *o*-benzyne (C_8H_6) potential energy surface (PES) was explored using DFT and the results are shown in Fig. 3. An entrance barrier of 55 kJ mol^{-1} (**T1**) was found to form the biradical intermediate (**INT1**) at -71 kJ mol^{-1} . From this initial adduct, the reaction path branches with the lower energy biradical pathway proceeding through a series of small rotational barriers; **T3** ($\langle S^2 \rangle = 1$) at -49 kJ mol^{-1} , **INT2** ($\langle S^2 \rangle = 1$) at -76 kJ mol^{-1} , and finally through **T4** ($\langle S^2 \rangle = 1$) at -59 kJ mol^{-1} for ring closure, forming **BCBdiene** at -288 kJ mol^{-1} . Alternatively, we have also located a closed-shell singlet pathway, initiated by the formation of a $[2 + 1]$ closed-shell intermediate (**INT3**) at 16 kJ mol^{-1} by crossing a barrier of 56 kJ mol^{-1} (**T6**), which also proceeds to **BCBdiene**. However, this pathway is

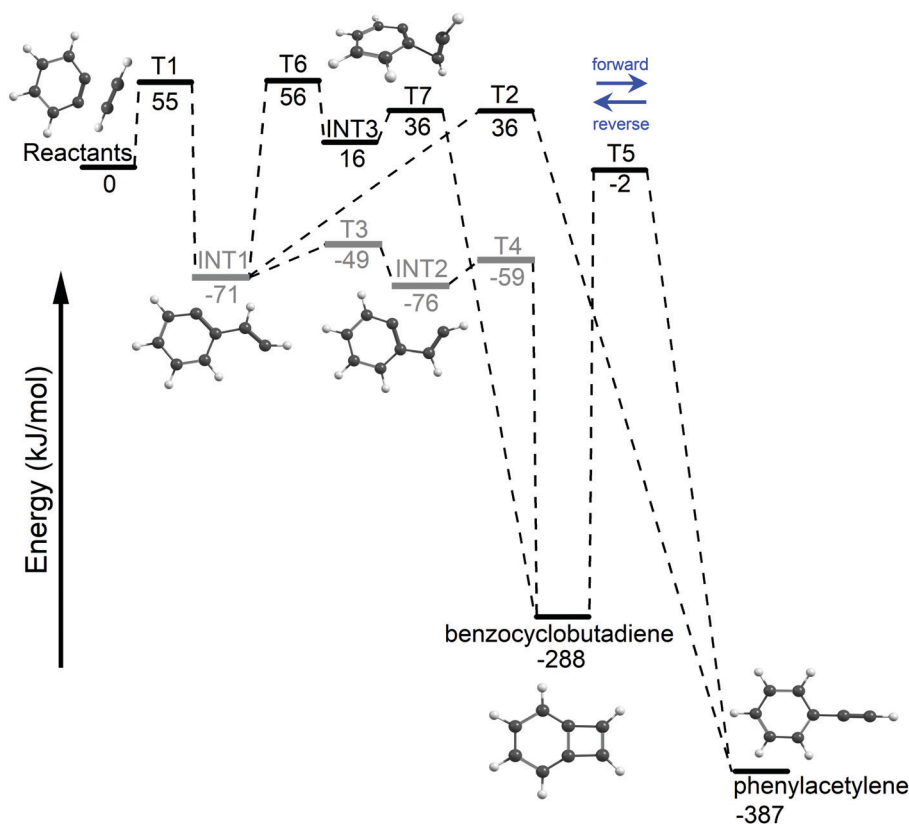


Fig. 3 The potential energy surface for the acetylene + *o*-benzyne reaction using M06-2X/6-311++G(d,p). Open-shell intermediates and transition states are indicated in gray.

unlikely to compete with the lower energy biradical pathway. The third pathway from the initial adduct (**INT1**) proceeds *via* a hydrogen transfer to the radical site on the ring, forming **PA** (-387 kJ mol^{-1}). This route also requires overcoming a transition state energy of 36 kJ mol^{-1} . **BCBdiene** may also isomerize to **PA** *via* a hydrogen transfer over **T5** at -2 kJ mol^{-1} .

In addition to density functional calculations at using the M06-2X functional, the G4 composite method⁶⁵ was also applied to refine the potential energy surface and obtain improved energetics. For closed-shell species, the agreement between the DFT and G4 results was generally good (with energies varying by no more than 10 kJ mol^{-1}), but G4 could not be applied consistently to describe open-shell singlet intermediates and transition states. Consequently, all energies are reported at the M06-2X/6-311++G(d,p) level of theory.

3.4 Kinetic modeling

RRKM theory was used to provide insight into the product branching for the two reactions. To this end, microcanonical rates were calculated as a function of the internal energy for the interconversion of the **BCBdiene** and **PA** products of the acetylene + *o*-benzyne reaction (Fig. 4).

The rate-limiting transition state, **T5**, connects the two species (see Fig. 3). The forward rate curve shown in Fig. 4 depicts the **PA** formation rate constant starting from the **BCBdiene** isomer, while the reverse rate curve shows the rate constant in the opposite direction. The vertical dotted line in this figure indicates the minimum energy needed for the *o*-benzyne + C_2H_2 association reaction to proceed, *i.e.*, the energy of the association transition state **T1** to form the initial adduct. The horizontal line in the figure shows an estimate of the inverse residence time of molecules in the pyrolysis microreactor.³⁸

The lowest energy pathway on the potential energy surface leads to **BCBdiene**. However, there is enough energy in the system to cross the barrier **T5** leading to **PA**. The RRKM rates provide us insights into the equilibration between the two

products on the time scale of the reactor. The reaction from **BCBdiene** leading to **PA** is several orders of magnitude faster than the reverse. Hence, on the time scale on which reactions occur in our reactor, the excess energy drives the reaction forward to form **PA** in larger abundance. However, even at the highest internal energies, the forward rate is still commensurate with the residence time in the microreactor. Based on the isomerization free energy, the **BCBdiene** abundance should be negligible in equilibrium. Combined with the experimentally observed 2 : 1 ratio of **PA** : **BCBdiene**, this implies that a fraction of **BCBdiene** is stabilized upon rethermalization in the microreactor, while a slightly larger fraction proceeds to form **PA**.

4 Discussion and conclusion

The results presented here can be compared with the study by Friedrichs *et al.*,²⁷ who suggested that the formation of phenylacetylene was dependent on a single “concerted” route, with an estimated barrier of over 70 kJ mol^{-1} . The reported energy of the transition state involved in this pathway was deemed “less reliable” as the optimization did not fully converge, yet the suggested route was assumed to exist based on similar reactions between *o*-benzyne and other unsaturated hydrocarbons. It was concluded that **BCBdiene** is the primary product of the reaction, although the authors proposed that subsequent isomerization may take place. This could, however, not be confirmed as the study was limited to recording mass spectrometric data only. Moreover, the authors had not presented a mechanism for the isomerization of **BCBdiene** to **PA**.

By applying i^2 PEPICO spectroscopy, we are now able to provide isomer-specific detection of the reaction products, allowing us to elucidate these key questions posed earlier. Based on the estimated photoionization cross sections combined with Franck–Condon factors, we have shown that, under our experimental conditions, **PA** is the favored reaction product with an estimated 2 : 1 branching ratio with respect to **BCBdiene**. Our PES calculations accounted for open-shell singlets explicitly, and allowed us to locate two additional routes for the formation of **PA** in the title reaction. Furthermore, a subsequent isomerization step was located, which accounts for the conversion of **BCBdiene** to **PA**. Kinetic modeling was employed to show that the dynamic competition between this isomerization step and the re-thermalization and, thus, stabilization of **BCBdiene** can explain the limited, 67% conversion of **BCBdiene** to the preferred **PA** product in the SiC microreactor.

The results of this study yield insights into the role of *o*-benzyne in the growth of aromatic molecules. The association reaction of *o*-benzyne with acetylene first has to overcome a sizeable entrance barrier of 55 kJ mol^{-1} . Afterwards, this bimolecular reaction produces a mix of a polycyclic species and an open-chain isomer of C_8H_6 composition, favoring the open-chain species **PA**. It may proceed in hot environments, such as in combustion and in outflows of carbon-rich stars, but, due to the high entrance barrier, it is highly unlikely to contribute to growth of aromatic molecules in cold regions of

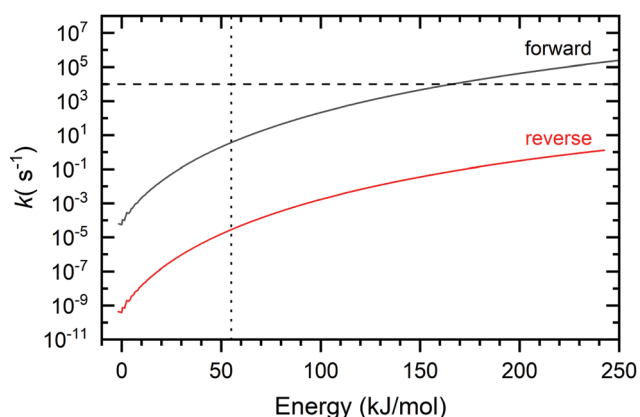


Fig. 4 Forward and reverse rate constants for crossing **T5** between **PA** and **BCBdiene** (see Fig. 3 for reference). The dotted vertical line indicates the minimum energy present needed for the reaction to proceed. The horizontal line indicates the average residence time of molecules inside the pyrolysis microreactor.

the ISM. Moreover, the polycyclic **BCBdiene** is formed by (collisional) stabilization, and the system is expected to shift more towards equilibrium, *i.e.*, towards **PA**, if the reaction time is not constrained. Although the system possesses sufficient internal energy to dissociate, collisional stabilization may not be essential for the survival of the adduct: the dissociation will proceed slower than the reverse isomerization step, which may lengthen the lifetime of the adduct to allow for IR-fluorescence, which could contribute to the stabilization of the C₈H₆ products. Further collisional stabilization may not occur in low-density interstellar molecular clouds. On the whole, this suggests reactions of *o*-benzyne with radicals are more likely the driving force behind the formation of larger poly-ringed species, at least in the ISM, as they are barrierless and proceed *via* addition–elimination mechanisms.^{12,25}

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

JB acknowledges the Netherlands Organisation for Scientific Research (Nederlandse Organisatie voor Wetenschappelijk Onderzoek, NWO) for a Vidi grant (grant number 723.016.006) which was used to support this work. This work was carried out on the Dutch national e-infrastructure with the support of SURF Cooperative (e-infra 46011 and EINF-997). The *i*²PEPICO experiments were performed at the VUV beamline at the Swiss Light Source (Paul Scherrer Institute, 5232 Villigen, Switzerland). PH and AB gratefully acknowledge funding by the Swiss Federal Office of Energy (BFE Contract Number SI/501269-01).

Notes and references

- H. Richter and J. B. Howard, *Prog. Energy Combust. Sci.*, 2000, **26**, 565–608.
- C. A. Menzie, B. B. Potocki and J. Santodonato, *Environ. Sci. Technol.*, 1992, **26**, 1278–1284.
- L. J. Allamandola, A. G. G. M. Tielens and J. R. Barker, *Astrophys. J.*, 1985, **290**, L25–L28.
- A. Leger and J. L. Puget, *Astron. Astrophys.*, 1984, **137**, L5–L8.
- A. G. G. M. Tielens, *Annu. Rev. Astron. Astrophys.*, 2008, **46**, 289–337.
- E. Peeters, S. Hony, C. Van Kerckhoven, A. Tielens, L. J. Allamandola, D. M. Hudgins and C. W. Bauschlicher, *Astron. Astrophys.*, 2002, **390**, 1089–1113.
- M. Frenklach and E. D. Feigelson, *Astrophys. J.*, 1989, **341**, 372.
- Z. Mansurov, *Combust., Explos. Shock Waves*, 2005, **41**, 727–744.
- M. Schenk, N. Hansen, H. Vieker, A. Beyer, A. Götzhäuser and K. Kohse-Höinghaus, *Proc. Combust. Inst.*, 2015, **35**, 1761–1769.
- T. Yang, T. P. Troy, B. Xu, O. Kostko, M. Ahmed, A. M. Mebel and R. I. Kaiser, *Angew. Chem., Int. Ed.*, 2016, **55**, 14983–14987.
- L. Zhao, R. I. Kaiser, W. Lu, B. Xu, M. Ahmed, A. N. Morozov, A. M. Mebel, A. Hasan Howlader and S. F. Wnuk, *Nat. Commun.*, 2019, **10**, 3689.
- D. E. Couch, A. J. Zhang, C. A. Taatjes and N. Hansen, *Angew. Chem., Int. Ed.*, 2021, **60**, 27230–27235.
- I. Cherchneff, J. R. Barker and A. G. G. M. Tielens, *Astrophys. J.*, 1992, **401**, 269.
- K. O. Johansson, M. P. Head-Gordon, P. E. Schrader, K. R. Wilson and H. A. Michelsen, *Science*, 2018, **361**, 997–1000.
- J. G. Radziszewski, B. A. Hess and R. Zahradnik, *J. Am. Chem. Soc.*, 1992, **114**, 52–57.
- A. C. Scheiner, H. F. Schaefer and B. Liu, *J. Am. Chem. Soc.*, 1989, **111**, 3118–3124.
- R. D. Brown, P. D. Godfrey and M. Rodler, *J. Am. Chem. Soc.*, 1986, **108**, 1296–1297.
- J. Cernicharo, M. Agúndez, R. I. Kaiser, C. Cabezas, B. Tercero, N. Marcelino, J. R. Pardo and P. de Vicente, *Astron. Astrophys.*, 2021, **652**, L9.
- A. Comandini, S. Abid and N. Chaumeix, *J. Phys. Chem. A*, 2017, **121**(31), 5921–5931.
- F. Hirsch, E. Reusch, P. Constantinidis, I. Fischer, S. Bakels, A. M. Rijs and P. Hemberger, *J. Phys. Chem. A*, 2018, **122**, 9563–9571.
- A. Keller, R. Kovacs and K.-H. Homann, *Phys. Chem. Chem. Phys.*, 2000, **2**, 1667–1675.
- A. D'Alessio, A. D'Anna, P. Minutolo, L. Sgro and A. Violi, *Proc. Combust. Inst.*, 2000, **28**, 2547–2554.
- A. Comandini and K. Brezinsky, *J. Phys. Chem. A*, 2012, **116**, 1183–1190.
- A. Comandini, S. Abid and N. Chaumeix, *J. Phys. Chem. A*, 2017, **121**, 5921–5931.
- M. N. McCabe, P. Hemberger, E. Reusch, A. Bodi and J. Bouwman, *J. Phys. Chem. Lett.*, 2020, **11**, 2859–2863.
- A. Matsugi and A. Miyoshi, *Phys. Chem. Chem. Phys.*, 2012, **14**, 9722–9728.
- G. Friedrichs, E. Goos, J. Gripp, H. Nicken, J.-B. Schönborn, H. Vogel and F. Temps, *Z. Phys. Chem.*, 2009, **223**, 387–407.
- C. Wentrup, *Aust. J. Chem.*, 2010, **63**, 979–986.
- W. T. G. Johnson, M. B. Sullivan and C. J. Cramer, *Int. J. Quantum Chem.*, 2001, **85**, 492–508.
- D. Kvaskoff, H. Lüerssen, P. Bednarek and C. Wentrup, *J. Am. Chem. Soc.*, 2014, **136**, 15203–15214.
- T. Baer and R. P. Tuckett, *Phys. Chem. Chem. Phys.*, 2017, **19**, 9698–9723.
- P. Hemberger, J. A. van Bokhoven, J. Pérez-Ramírez and A. Bodi, *Catal. Sci. Technol.*, 2020, **10**, 1975–1990.
- P. Hemberger, A. Bodi, T. Bierkandt, M. Köhler, D. Kaczmarek and T. Kasper, *Energy Fuels*, 2021, **35**, 16265–16302.
- B. Sztáray, K. Voronova, K. G. Torma, K. J. Covert, A. Bodi, P. Hemberger, T. Gerber and D. L. Osborn, *J. Chem. Phys.*, 2017, **147**, 013944.
- A. Bodi, P. Hemberger, T. Gerber and B. Sztáray, *Rev. Sci. Instrum.*, 2012, **83**, 083105.
- D. W. Kohn, H. Clauberg and P. Chen, *Rev. Sci. Instrum.*, 1992, **63**, 4003–4005.

- 37 M. S. South and L. S. Liebeskind, *J. Org. Chem.*, 1982, **47**, 3815–3821.
- 38 Q. Guan, K. N. Urness, T. K. Ormond, D. E. David, G. B. Ellison and J. W. Daily, *Int. Rev. Phys. Chem.*, 2014, **33**, 447–487.
- 39 S. Grimm, S.-J. Baik, P. Hemberger, A. Bodi, A. M. Kempf, T. Kasper and B. Atakan, *Phys. Chem. Chem. Phys.*, 2021, **23**, 15059–15075.
- 40 M. V. Zagidullin, R. I. Kaiser, D. P. Porfiriev, I. P. Zavershinskiy, M. Ahmed, V. N. Azyazov and A. M. Mebel, *J. Phys. Chem. A*, 2018, **122**, 8819–8827.
- 41 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery, Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 16, Revision A.03*, Gaussian Inc., Wallingford CT, 2016.
- 42 A. D. Becke, *J. Chem. Phys.*, 1993, **98**, 1372–1377.
- 43 C. Lee, W. Yang and R. G. Parr, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1988, **37**, 785–789.
- 44 S. H. Vosko, L. Wilk and M. Nusair, *Can. J. Phys.*, 1980, **59**, 1200.
- 45 P. J. Stephens, F. J. Devlin, C. F. Chabalowski and M. J. Frisch, *J. Phys. Chem.*, 1994, **98**, 11623–11627.
- 46 R. Krishnan, J. S. Binkley, R. Seeger and J. A. Pople, *J. Chem. Phys.*, 1980, **72**, 650–654.
- 47 S. Gozem and A. I. Krylov, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2021, e1546.
- 48 S. Gozem, A. O. Gunina, T. Ichino, D. L. Osborn, J. F. Stanton and A. I. Krylov, *J. Phys. Chem. Lett.*, 2015, **6**, 4532–4540.
- 49 J. Paldus and X. Li, in *A Critical Assessment of Coupled Cluster Method in Quantum Chemistry*, John Wiley & Sons, Ltd, 1999, pp. 1–175.
- 50 R. J. Bartlett and M. Musiał, *Rev. Mod. Phys.*, 2007, **79**, 291–352.
- 51 T. H. Dunning, *J. Chem. Phys.*, 1989, **90**, 1007–1023.
- 52 Y. Shao, L. F. Molnar, Y. Jung, J. Kussmann, C. Ochsenfeld, S. T. Brown, A. T. B. Gilbert, L. V. Slipchenko, S. V. Levchenko, D. P. O'Neill, R. A. DiStasio, R. C. Lochan, T. Wang, G. J. O. Beran, N. A. Besley, J. M. Herbert, C. Y. Lin, T. Van Voorhis, S. H. Chien, A. Sodt, R. P. Steele, V. A. Rassolov, P. E. Maslen, P. P. Korambath, R. D. Adamson, B. Austin, J. Baker, E. F. C. Byrd, H. Dachsel, R. J. Doerksen, A. Dreuw, B. D. Dunietz, A. D. Dutoi, T. R. Furlani, S. R. Gwaltney, A. Heyden, S. Hirata, C.-P. Hsu, G. Kedziora, R. Z. Khalliulin, P. Klunzinger, A. M. Lee, M. S. Lee, W. Liang, I. Lotan, N. Nair, B. Peters, E. I. Proynov, P. A. Pieniazek, Y. M. Rhee, J. Ritchie, E. Rosta, C. D. Sherrill, A. C. Simmonett, J. E. Subotnik, H. L. Woodcock, W. Zhang, A. T. Bell, A. K. Chakraborty, D. M. Chipman, F. J. Keil, A. Warshel, W. J. Hehre, H. F. Schaefer, J. Kong, A. I. Krylov, P. M. W. Gill and M. Head-Gordon, *Phys. Chem. Chem. Phys.*, 2006, **8**, 3172–3191.
- 53 Y. Zhao and D. G. Truhlar, *Theor. Chem. Acc.*, 2008, **120**, 215–241.
- 54 P. A. Jensen, M. Leccese, F. D. S. Simonsen, A. W. Skov, M. Bonfanti, J. D. Thrower, R. Martinazzo and L. Hornekær, *Mon. Not. R. Astron. Soc.*, 2019, **486**, 5492–5498.
- 55 D. Campisi, F. D. S. Simonsen, J. D. Thrower, R. Jaganathan, L. Hornekær, R. Martinazzo and A. G. G. M. Tielens, *Phys. Chem. Chem. Phys.*, 2020, **22**, 1557–1565.
- 56 D. Campisi and A. Candian, *Phys. Chem. Chem. Phys.*, 2020, **22**, 6738–6748.
- 57 B. O. Roos, P. R. Taylor and P. E. Sigbahn, *Chem. Phys.*, 1980, **48**, 157–173.
- 58 B. Sztáray, A. Bodi and T. Baer, *J. Mass Spectrom.*, 2010, **45**, 1233–1245.
- 59 J. Bouwman, A. Bodi, J. Oomens and P. Hemberger, *Phys. Chem. Chem. Phys.*, 2015, **17**, 20508–20514.
- 60 K. Kimura, S. Katsumata and Y. Achiba, *Handbook of HeI photoelectron spectra of fundamental organic molecules ionization energies, Ab Initio assignments, and valence electronic structure for 200 molecules*, Japan scientific societies Press; Halsted Press, Tokyo; New York, 1981.
- 61 P. Hemberger, A. J. Trevitt, T. Gerber, E. Ross and G. da Silva, *J. Phys. Chem. A*, 2014, **118**, 3593–3604.
- 62 J. M. Dyke, H. Ozeki, M. Takahashi, M. C. R. Cockett and K. Kimura, *J. Chem. Phys.*, 1992, **97**, 8926–8933.
- 63 T. Koenig, D. Imre and J. A. Hoobler, *J. Am. Chem. Soc.*, 1979, **101**, 6446–6447.
- 64 NIST, NIST Chemistry WebBook, National Institute of Standards and Technology, Gaithersburg MD, 20899, 2009, vol. NIST Standard Reference Database Number 69.
- 65 L. A. Curtiss, P. C. Redfern and K. Raghavachari, *J. Chem. Phys.*, 2007, **126**, 084108.