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# Insights into the oxygen evolution mechanism of transition metal-anchored holey graphyne

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Growing worldwide environmental concerns linked to the overuse of fossil fuels and rising energy demands are driving a thorough and comprehensive search for clean, sustainable sources of energy. Water electrolysis has recently become a highly appealing method for achieving optimal energy conversion and storage. This study employs density functional theory to investigate the catalytic performance along with the electronic properties of pristine and adatom-doped transition metal (TMs = Sc, Pt, Co, Cr, and Au)-anchored holey graphyne (HGY). Among all considered candidates, Pt-doped HGY gives the best oxygen evolution reaction (OER) with the overpotential equivalent to 0.74 V. This catalyst is deemed optimal for further exploration of the OER mechanism. Through molecular dynamics (MD) simulations, the structural along with thermal stability of Pt@HGY has been confirmed. The convincing results motivate the use of Pt-anchored HGY as an efficient OER mechanism catalyst.

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## Introduction

The pressing issue of meeting the increasing energy demand caused by rising populations and societal progress must be addressed without delay. The oxygen evolution reaction is hindered by its slow kinetics, which is caused by a complex multielectron reaction process. As a result, the reaction requires a greater overpotential in order to provide a sufficient current for the OER. Noble metal-based compounds such as IrO<sub>2</sub> and RuO<sub>2</sub> are commonly used as commercial electrocatalysts for the OER.<sup>1–3</sup> The limited availability and expensive cost severely restrict their use on a large commercial level. Significant efforts have been made in this case to determine which earth-copious components are suitable for OER catalysis.<sup>4,5</sup> The rate-determining step (RDS), which is the adsorption as well as the desorption of intermediates onto the active site, determines the kinetics of the OER; the effectiveness of this step is correlated with the active site's d-orbital electronic structure. Consequently, adjusting the electronic structure of the electrochemical reaction's active site at the atomic scale is a useful way to enhance the electrochemical performance of electrode materials. The elementary phase of oxygen evolution, which involves the formation of an O–O bond and the complex transfer of 4 protons along with 4 electrons, is considered to be the slowest

despite having the highest activation energy.<sup>6–8</sup> As a result, the reaction exhibits a relatively low level of efficiency, hence impeding its practical application. Nevertheless, due to the sluggish four-electron transfer, catalysts are indispensable for the vast majority of OER processes, and the prohibitively high cost of existing noble metal-based catalysts is a critical weakness to sustainable chemistry as well as industrialization.<sup>9,10</sup> Efficient electrocatalysts are necessary to control the overpotential and accelerate the process. Throughout the literature, it has been recognised that oxides of iridium and ruthenium are the most suitable options for catalysing the process of oxygen evolution,<sup>11–16</sup> as they are cost-efficient yet exhibit comparable performance. Hence, it is essential for researchers to possess a thorough comprehension of the reaction mechanisms implicated in every phase of the OER and strategically devise catalysts that can adequately meet market requirements, including water splitting reaction devices and metal air batteries. The OER is a vital reaction that plays a key part in a number of renewable energy technologies, such as metal air batteries and water splitting.<sup>17–20</sup>

Two-dimensional (2D) materials possess intriguing characteristics, including a substantial surface area, excellent thermal as well as mechanical stability, and adjustable electronic properties, along with versatility in various research domains and applications such as gas sensors, spintronics, energy storage, etc.<sup>16,21,22</sup> These attributes have contributed to their increasing popularity in recent times. Examples of these materials include porous carbon compounds like graphene, and transition metal dichalcogenides (TMDCs) such as MoS<sub>2</sub>, WS<sub>2</sub>, and VSe<sub>2</sub>, and hexagonal boron nitride, to name a few.<sup>14,15,17,20</sup> Because of the previously described qualities, high sensitivity, selectivity, and

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quick recovery, 2D materials are now widely used as efficient catalysts to develop electrocatalytic processes.<sup>23</sup> Graphene is an exceptionally desirable nanomaterial on account of its remarkable mechanical strength, and phononic and captivating electro-optical characteristics, as it is composed of carbon atoms arranged in the most stable 2D configuration. The characteristics exhibited by graphene have resulted in significant progress within the domains of contemporary chemistry and physics. The electrical structure of graphene, characterised by a zero-bandgap, imposes limitations on its use within the domain of semiconducting materials. Thereupon, it is crucial to identify avant-garde types of two-dimensional carbon allotropes that demonstrate peculiar semiconducting properties, for instance an appropriate energy bandgap and enhanced mobility.<sup>24–27</sup> Graphene possesses a honeycomb lattice structure, while graphyne exhibits several geometrical configurations. At present, scholars have predominantly examined four unique geometries of graphynes, *viz.* a-, b-, g-, and 6,6,12-graphyne.<sup>28–30</sup> The remarkable attributes of graphyne make them strong competitors to graphene, especially in specific applications. The design and coalescence of the anecdote carbon allotrope, namely holey graphyne (HGY), was inspired by this intriguing molecule, as reported. The atomic structure of HGY is characterised by the presence of uniform pores arranged in 6, 8, and 24 vertex rings. The rings are surrounded by acetylene connections of zero, two, and six, which are produced through sp–sp bonds. Two C atoms from the benzene ring are joined to 2 C atoms from the following octagon to provide a 50% sp to sp<sup>2</sup> C ratio.<sup>31</sup> It was previously found that HGY is composed of a p-type semiconductor that has elevated electron and Hall mobilities. At ambient temperature, the system exhibits a hole mobility of  $2.13 \times 10^9 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ , an electron mobility of about  $10^4 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ , and a band gap of 1.0 eV.<sup>32</sup> Then, in 2021, Gao *et al.* determined that the overall gravimetric density of Li-coated HGY could surge to 12.8% when they investigated its H<sub>2</sub> storage capacity.<sup>33</sup> An investigation conducted in 2022 by Xinghui *et al.* examined the stability, vibrational, and optical properties of HGY. The researchers discovered that HGY sheets synthesized with AB stacking are exceptionally crystalline. The semiconducting nature, non-linear sp bonding, and unique pi-conjugated structure of HGY make it very suitable for a wide range of applications.<sup>34</sup> In their research, Juhee *et al.* investigated the adsorption and desorption properties of hydrogen (H) in the Ti-functionalized HGY system. It was also observed that the titanium atom exhibits an adsorption energy of  $-0.38$  electron Volt per hydrogen atom, which allows it to gather seven H<sub>2</sub> molecules. The outcome of this process yields a hydrogen gravimetric density of 10.52 weight percent.<sup>35</sup> Subsequently, the H storage potential of HGY anchored with Sc as well as Y was examined by Mukesh *et al.* They discovered that the gravimetric densities could reach 9.80% and 9.34%.<sup>36</sup> In a study conducted in 2023, Seetha *et al.* examined the detection of ammonia on HGY, decorated with Cu, Sc, and Pd. The results showed that among all the decorations, Sc-doped HGY was identified as the best NH<sub>3</sub> sensor. At a temperature of 600 K, the sample demonstrated an adsorption

energy of 1.49 eV, a notable charge transfer of  $0.113e$ , and a recoverable time of 3.2 s.<sup>37</sup> Darshil *et al.* conducted a study which showed that the HGY catalysts, doped with different TMs like Fe, Cr, and Co, exhibit a high level of activity for the HER with  $\Delta G$  values as low as  $-0.21$ ,  $-0.14$ , and  $-0.05$  electron Volt.<sup>31</sup> Manman *et al.* found that SACs buoyed by graphydyne and holey graphyne are a potential option for the CO<sub>2</sub>RR after studying them using machine learning. By considering the complete elemental miscellany of metal sites across the periodic table, it was discovered that 25 catalysts could both successfully activate CO<sub>2</sub> along with the HER.<sup>38</sup> Muhammad *et al.* examined the potential of HGY for battery applications, and it was observed that compared to other carbonaceous materials (200–500 mAh g<sup>-1</sup> for Na and 450–750 mAh g<sup>-1</sup> for Li) and traditional graphite anodes (372 mAh g<sup>-1</sup> for Li and < 35 mAh g<sup>-1</sup> for Na), Na (558 mAh g<sup>-1</sup>) was revealed to have much larger storage capacities.<sup>39</sup> Bohayra examined the mechanical, electrical, and thermal characteristics of C and BN-HGY and discovered that the semiconducting C-HGY along with insulating NB-HGY system had moderate lattice thermal conductivity and ultrahigh negative thermal expansion coefficients.<sup>40</sup> Yong *et al.* have examined a higher-order topological invariant connected to the material's C2 symmetry and proved that the ensuing corner modes materialize in nano-flakes that match the hitherto disclosed precursor structure.<sup>41</sup> While initial theoretical reckoning of HGY have been documented for applications such as hydrogen storage, battery technology, and catechol detection; HGY's potential as a general water splitting mechanism involving certain transition metals with the OER remains unexplored.

In this study, we report using DFT simulations, an unprecedentedly potent way of oxygen evolution *via* pristine and transition metal (TM = Au, Co, Cr, Pt, and Sc) HGY monolayers. The utilisation of theoretical simulations aids in the clarification of the adsorption mechanism at the orbital level. These insights include favourable adsorption configurations, adsorption energy, bonding angle and distance, charge transfer, and binding energy. In light of theoretical considerations such as adsorption energy, Löwdin charge analysis and partial density of states, the phenomenon of charge transfer has been examined. Additionally, it has been observed that even at extremely high temperatures up to 500 K, the Pt element maintains its binding with HGY. These *ab initio* molecular dynamics simulations provide additional substantiation for the thermal stability of the system.

## Computational methodology

Quantum ESPRESSO was used to perform calculations in accordance with density functional theory (DFT).<sup>42,43</sup> The generalized gradient approximation (GGA) was used to represent the exchange–correlation energies, and the projector augmented wave (PAW) potentials were used to analyse the electron–ion interactions.<sup>44</sup> Integration over the Brillouin zone was conducted utilising a *k*-point mesh of  $5 \times 5 \times 1$  and a plane wave



kinetic energy cutoff of 50 Ry. Both the atomic positions and the structural parameters were relaxed until the energy reached a convergence of  $10^{-6}$  eV, respectively. A 20 Å vacuum was used to eliminate the imperceptible forces between neighbouring layers in the z-direction. The utilisation of Grimme's DFT-D2 strategies was implemented in order to tackle the extended van der Waals interactions that encompass heterostructures and bilayers.<sup>45</sup> The thermal and structural stability have been verified by molecular dynamics simulations, utilizing the VASP<sup>46</sup> code.

## Results and discussion

### Structural, binding energy and electronic properties, and the OER mechanism of pristine HGY and TM-anchored HGY

The lattice constants for pristine HGY were determined to be 10.796 Å based on the hexagonal unit cell's structure. Fig. 1(a and b) show the 24 C atoms that make up the HGY supercell. By applying an adamant vacuum level of 20 Å across the aperiodic orientation and fully relaxing the atomic position with regard to cell size, the pristine HGY supercells were optimized.

The scrutiny of pristine HGY and its function in the OER process of water splitting will be covered. O adsorption alone is not a reliable signal for forecasting OER reactions, according to a literature review. The entire activity of the OER, comprising OH, O, and OOH intermediates, has been distinctly surveyed in this research work. The adsorption geometries of HGY with the potential adsorption sites are shown in Fig. 1(c).

Fig. 1 shows the optimised geometry of the HGY monolayer from both (a) top view and (b) side perspectives. The optimal bond lengths (Å) along with the bond angles (°) of OER intermediates (O, OH, and OOH) adsorbed at various potential adsorption sites (P1–P9) on pristine HGY are summarized in Table S1. The lengths of the C–C bonds proximal to the binding site for O adsorption range from 1.21 to 1.32 Å, whereas those

next to the binding site extend slightly up to 1.47 Å. The site-dependent lattice distortions caused by O anchoring are indicated by the large fluctuation in the C–C–O bond angles, which range from 61.9° at P4 to as high as 143.3° at P2 and P5. In accordance with molecular stability, the O–O ( $\approx 1.36$ – $1.38$  Å) and O–H ( $\approx 0.98$  Å) bond lengths in the event of OOH adsorption stay almost constant across all locations. Disparities in adsorption strength and contact distance with the HGY surface are reflected in the O–C bond lengths, which range greatly from 2.47 Å (P1) to 3.10 Å (P9). The angular parameters likewise exhibit a substantial dependency on the adsorption site: the O–O–H angle is typically about 103°, the C–C–O angle ranges from 83.5° (P4) to 114.3° (P7), and the C–O–O angle ranges from 107.6° (P1) to 173.2° (P7). According to these findings, the inherent OOH geometry is maintained, although the binding orientation and HGY interaction differ considerably depending on the adsorption location.

Overall, the structural study shows that O adsorption results in localized bond distortions, while OOH adsorption, particularly at sites P7 and P9, produces greater fluctuations in adsorption distance along with angular distortion. The sensitivity of pristine HGY to intermediate binding is highlighted by its site-specific structural flexibility, which is essential for comprehending its catalytic action in the OER. Although stronger binding sites like P1 and P2 run the risk of over binding and impairing the reaction kinetics, they may effectively stabilize intermediates. On the other hand, OOH, the crucial intermediate in the OER, is frequently destabilized by weaker binding sites like P7 and P9, which lessens its catalytic significance.

In order to provide a detailed elucidation of the adsorption reaction mechanism, the HGY transition state was calculated. Initially, the reaction included a proton being transferred to an adsorbed OH molecule *via* a surface-adsorbed H<sub>2</sub>O molecule. A proton is released by the adsorbed hydroxide (OH), creating an adsorbed atomic oxygen. Adsorbed hydroperoxyl (OOH) is the result of this atomic oxygen's subsequent reaction with OH<sup>−</sup>



Fig. 1 (a) Top and (b) side view of HGY, and (c) pictorial depiction of HGY for OER activity with possible adsorption sites represented by pink circles. Colour scheme: silver C; pink OER activity.



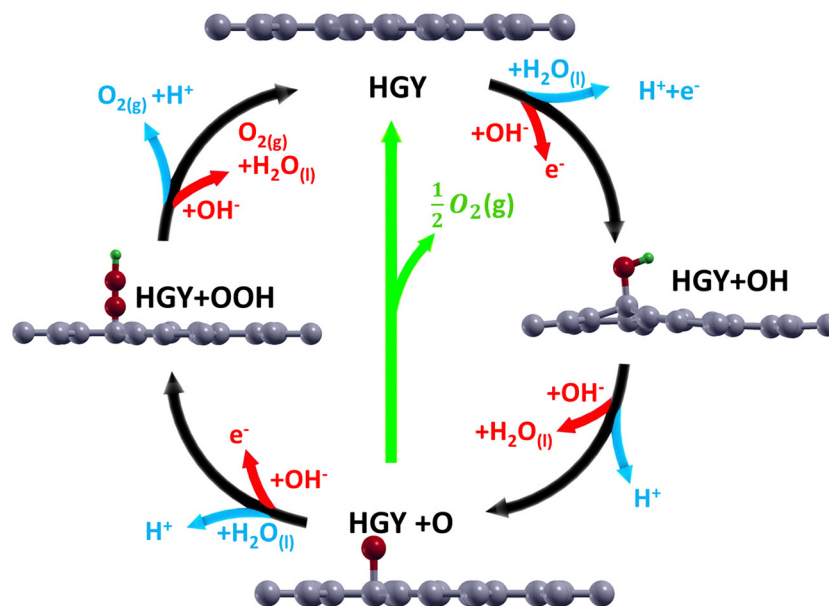


Fig. 2 Schematic representation of the typical OER pathway for pristine HG Y. Colour scheme: silver C; red O; green H.

(hydroxide). The adsorbed OOH and its original surface were transformed into molecular  $O_2$  in the fourth step. Fig. 2 illustrates the adsorption of OH on the slab and the genesis of \*OOH by the adsorption of \*O on OH, releasing energy in accordance with chemical principles. According to the Gibbs free energy diagram, spontaneity was not seen, however, since an external electric field was necessary to perturb the initial electric field balance and promote the movement of the adsorption group. The anaerobic HG Y adsorbs species including \*O and \*OH, as well as \*OOH, as depicted in Fig. 2. To ascertain the adsorption energies of the species \*OH, \*O, and \*OOH, we individually attach each species to a different site, such as the top, hollow, and bridge of the carbon atom. Subsequently, we relax the geometry of the system, as illustrated in Fig. 1(c). The hollow sites are denoted by P6, P7, P8, and P9, while P3, P4, and P5 represent the bridge sites between the two carbon atoms, and P1 and P2 represent the C top sites.

The OER is accomplished using procedures (5) to (9) from the SI. The free energy of the reaction was calculated for each individual step.

$$\Delta G = \Delta E + \Delta E_{ZPE} - T\Delta S + \Delta G_U + \Delta G_{pH} \quad (1)$$

The symbol  $\Delta E$  signifies the total energy change between the initial reactants and final products in a chemical process. The zero-point energy corrections are denoted by  $\Delta ZPE$ .  $\Delta S$  denotes the alteration in vibrational entropy with respect to a particular temperature  $T$ .  $\Delta G_U$  is equivalent to  $-eU$ , where  $U$  denotes the electrode potential and  $e^-$  signifies the fundamental charge. The correction for  $H^+$  free energy is expressed as  $\Delta G_{pH}$ .<sup>10</sup>

The overpotential, symbolised as  $\eta$ , is calculated by comparing the Gibbs free energy differences at each stage.

$$\eta_{OER} = \max\{\Delta G_1, \Delta G_2, \Delta G_3, \Delta G_4\}/e - 1.23 \quad (2)$$

where  $\Delta G_n$  are the free energy of reactions (5)–(8) (from the SI).<sup>47,48</sup>

Fig. 3 illustrates the computed reaction energies associated with each OER process taking place at the active sites, as well as the excess potential of pristine HG Y. For the purpose of assessment, the utmost reaction energies for the HG Y sites were computed as follows: 3.09, 4.12, 2.64, 3.14, 4.12, 3.22, 3.80, 4.03, and 3.93 eV for P1, P2, P3, P4, P5, P6, P8, and P9 (step 3 from OH\* to OOH\*), and P7 (step 1 from \* to OH\*). The overpotential values for P1 to P9 are 1.86, 2.89, 1.41, 1.91, 2.89, 1.99, 2.57, 2.80, and 2.70 V, respectively. Following this, the OER activity declines in the subsequent order: P2 > P5 > P8 > P9 > P7 > P6 > P4 > P1 > P3. Inferring from the data from Fig. 3, the activity of the P3 site (bridge site 1) for the OER surpasses those of the other sites of HG Y.

The octagon void shown in Fig. 4(a) has been used to anchor the TMs (Au, Pt, Co, Cr, and Sc). The prediction of a structural modification in the HG Y is evident from Table S2, which indicates that the interaction between the HG Y support and TMs induces a modification in the HG Y upon their introduction. The average distances between transition metals (Pt, Sc, Au, Cr, and Co) and the carbon atoms of the HG Y support are 2.07, 2.77, 3.06, 2.07, and 1.97 Å, for transition metals attached to HG Y.

According to the following relationship, the binding energy is a crucial factor in the recurrent usage of single atom catalysts (SACs):

$$E_b = E(\text{TMs}/\text{HG Y}) - E(\text{HG Y}) - E(\text{TM}) \quad (3)$$

The aggregate energy of the HG Y system doped by TMs is designated as  $E(\text{TMs}/\text{HG Y})$ , the isolated TM atom's total energy is epitomized by  $E(\text{TM})$ . The  $E_b$  values of the TMs (Au, Co, Cr, Pt, and Sc) on HG Y were determined in the current work to be





Fig. 3 Free energy diagram of the OER for pristine HGy for all possible sites at zero potential. The black line is the rate determining step of OER activity. (a)–(i) represents sites P1–P9.

–1.40, –2.44, –5.69, –2.45, and –3.60 eV, respectively, using eqn 3. HGy doped with a distinct TM support exhibits a negative magnitude of  $E_b$ , which implies a firm binding combined with suitable appropriateness for the substrate of the catalyst (refer to Fig. 4(b)). For instance, Darshil *et al.* found that Zr, Y, V, Sc, Mn, Co, and Fe had binding energies of –4.72, –3.57, –5.23, –2.50, –2.44, –2.27, and –5.69 eV on HGy, respectively.<sup>31</sup> Likewise, the binding energies of –4.09 and –0.69 eV for Sc- and Pd-doped HGy were reported by Seetha *et al.*<sup>49</sup> In comparison, our calculated binding energy for Sc is –4.05 eV, which falls well within the range of previously reported values. A direct comparison, particularly for Sc and other metals, indicates that the slight variations observed among the studies mainly results from differences in pseudo-potentials and basis sets used in computational packages such as QE and VASP. Overall, our findings confirm consistency and reliability, as the obtained binding energies align with established literature values.

Simulations were performed by keeping oxygen atom at 2 Å atop the TMs, to investigate the OER process of adatom doping of TMs on HGy. Next, as shown in Fig. 4(c–f [top view] and g–j [side view]), we proceeded to analyse the intermediates of the OER activity, O, OH, and OOH. The adsorption of intermediates causes a modest or minimal change in structural

characteristics like bond length and angle in HGy, SACs anchored by TMs. Table S2 makes it evident that this is because of the liaison that takes place between TMs and intermediates. To facilitate comparison, Tables S1 and S2 detail alterations in structural parameters, such as bond length and bond angle, between the HGy-anchored intermediate and TMs HGy SACs. The distance between intermediates such as C–O, O–H, and O–O is between 0.97 and 3.10 Å. The bond lengths spanning from 0.96 to 1.99 Å exist between intermediates such as TM–O, O–H, and O–O, as well as O–H.

In the context of the OER, 4.92 eV is the total energy required to break down water. For the reaction to continue efficiently, the  $\Delta G$  for each elementary stage in the OER should ideally be minimized. Specifically, the overpotential of the OER must be at least 1.23 eV. Hence, foremost efficient OER catalysts' intermediate Gibbs free energy values will position the total reaction energy around or beneath this threshold. Gibbs free energy profiles and over potential are shown in Fig. 6(a–e) to illustrate the OER process that occurs at the loftiest TMs, comprising Au, Co, Cr, Pt, and Sc, throughout the reaction route.

The maximum reaction energies for the HGy system are found to be as follows: 2.40 eV for Au@HGy (step 2 involves the transformation from OH\* to O\*), 2.54 eV for Co@HGy [(OOH\* to O<sub>2</sub>) in step 4], 3.19 eV for Cr@HGy [(O\* to OOH\*) in step 3],





Fig. 4 (a) Energetic position for TM-anchored HGY and (b) binding energy of TM (Au, Co, Cr, Pt, and Sc)-anchored HGY. Top view of (c) TM-HGY (TM = Pt, Au, Co, Cr, and Sc), (d) TM-HGY/O, (e) TM-HGY/OH, and (f) TM-HGY/OOH and side view of (g) TM-HGY (TM = Pt, Au, Co, Cr, and Sc), (h) TM-HGY/O, (i) TM-HGY/OH, and (j) TM-HGY/OOH. Colour scheme: silver C; blue S; navy blue TM; red O; green H.

1.97 eV for Pt@HG Y [( \*O to \*OOH) in step 3], and 3.32 eV for Sc@HG Y [( \*OOH to O<sub>2</sub>) in step 4]. It is found that the maximum reaction energy for Pt@HG Y is 1.97 eV, which is significantly closer to the ideal value of 1.23 eV. To boost the OER process, Pt@HG Y helps to strengthen the link that forms between the catalyst's surface and the \*O intermediate.

The d-band centre is widely recognized as a valuable descriptor for quantifying the strength of the interaction amid adsorbates and catalysts on various surfaces. The following equation is used to determine the d-band centre:

$$\varepsilon_d = \frac{\int_{-\infty}^{\infty} p(\varepsilon) \varepsilon d\varepsilon}{\int_{-\infty}^{\infty} p(\varepsilon) d\varepsilon} \quad (4)$$

The energy width of the d orbital is represented by  $\varepsilon$ , while the density of states is represented by  $\rho$ .<sup>50</sup> The d orbital of the metal atom has a major impact on the surface electronic structure of transition metals. The projected density of states (PDOS) along with the d band centre  $\varepsilon_d$  of OH, O, and OOH adsorption on Pt@HG Y were investigated in order to better understand the binding properties of the species involved in

the associative mechanism (Fig. 7(a–c)). After the intermediates are adsorbed, the d-band centre of each model shows a direction irreconcilable to the Fermi level, with values ranging from  $-1.15$  eV. The metal forfeits excess electrons in order to form an equitable bond with the intermediate due to electron exchange hybridization between the d orbits of Pt@HG Y and the sp orbits of the intermediate; additionally, a larger and less appropriate d band centre indicates an enhanced occupancy of antibonding orbitals after the d band of the metal catalyst hybridizes with the s/p band of the adsorbate, leading to the dissociation barrier of the reactants along with enabling the desorption of intermediates across the catalytic reaction system. Moreover, the d band centre boosts by 4.14 electron Volt due to the connection that links H and \*OH creating \*O. This result supports the theory that electrons can occupy the antibonding orbitals when the H atom is present. Likewise, when the \*O is bonded with the –OH, it experiences greater electron loss due to orbital hybridization; as a result, the \*OOH exhibits a greater density of states in proximity to the Fermi level, which further reduces the d band centre. According to these findings, the electronic field as well as metal properties have a direct





Fig. 5 Representation of the  $4e^-$  OER process of the 4H site over TM-HGY with the optimized configurations for O, OH, and OOH (intermediates). Colour scheme: silver C; navy blue TM (TM = Au, Pt, Co, Cr, and Sc); red O; green H.



Fig. 6 Gibbs free energy diagram of the OER for (a) Au, (b) Co, (c) Cr, (d) Pt, and (e) Sc in HGY at zero potential.

bearing on the OER of the Pt@HGY system. The electronic properties of O, H, C, and Pt atoms are compared, and it is found that O exhibits the optimum electronegativity. Additionally, the \*O and \*OOH formation phases were chosen because the incarnation of \*O or \*OOH was the crucial step in pursuing the redistribution of charge, and comprehending that the d block of the Pt atoms that exist in Pt@HGY for the OER mechanism are extensive *versus* a single atom of O as well as H. In all phases comprising the intermediates, Fig. 7 shows that

the interaction between Pt and C is almost the same. It is found that the interaction between O and Pt is negligible for both \*O and \*OOH when the Löwdin charge for transition state 3 amid the \*O and \*OOH is critiquing. Fig. 7(f) makes it clear that when OOH is attached, the charges of oxygen atoms only slightly change, yet the charges of H atoms drop relative to OH [Fig. 7e]. Table S2, where the bond length represents the electron cloud of the two elements, offers an additional comprehension method. According to Table S2, Pt@HGY + OH has



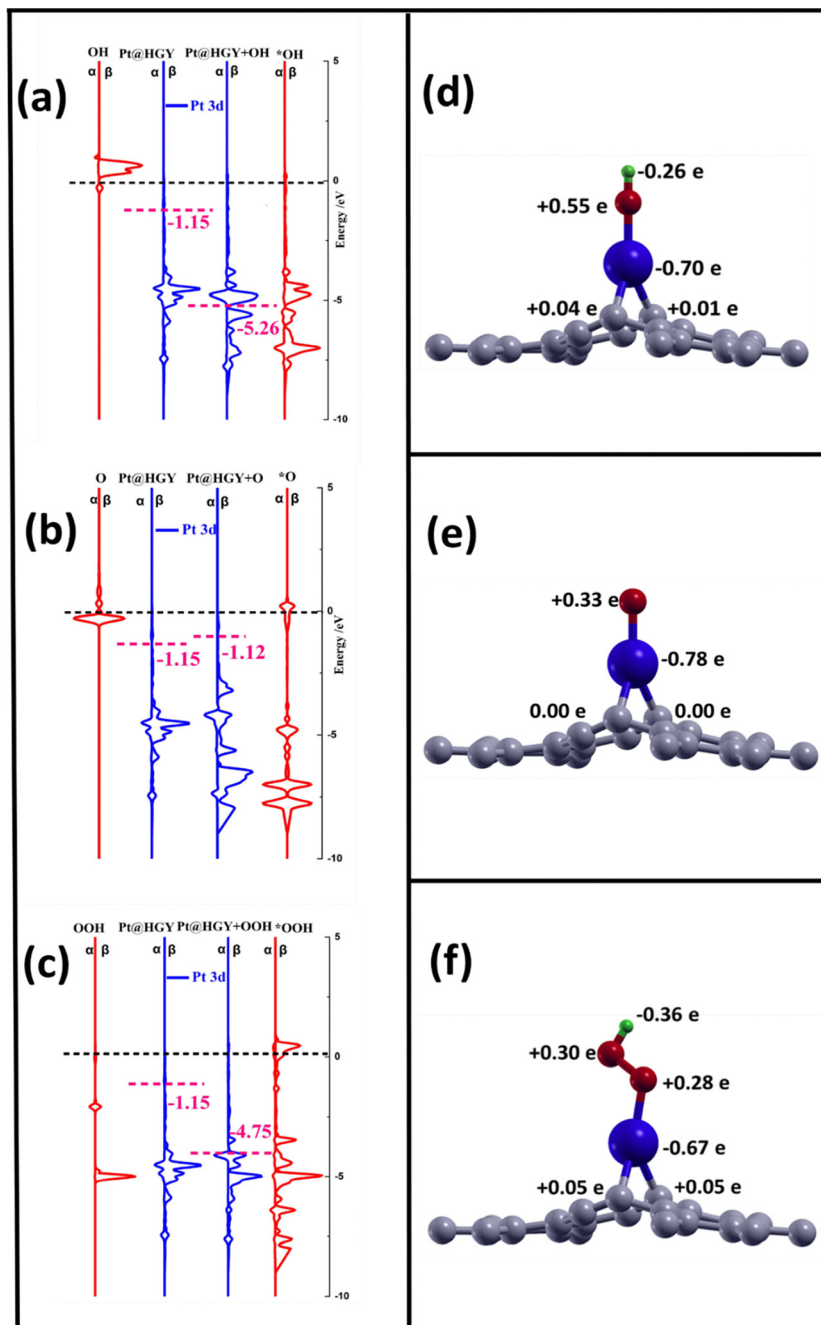


Fig. 7 (a)–(c) PDOS and diagrammatic representation of 3d orbitals of Pt@HGY, where the solid blue line embodies the Pt 3d orbitals. 2p orbitals of intermediates prior and subsequent to the adsorption on the Pt@HGY substrate. The pink dotted line marks the position of the d band centre of the Pt@HGY substrate prior and subsequent to the adsorption of O, OH, and OOH. Löwdin charge analysis (d)–(f) of the corresponding adsorbed OH, O, and OOH. Colour scheme: silver C; navy blue Pt; green H; red O.

a bond length of 1.82 Å between Pt and O. Likewise, Pt@HGY + OOH has a total bond length of 4.4 Å, with the bond lengths among Pt–O, O–O, along with O–H being 1.98, 1.45, and 0.97 angstrom. An increase in bond length, which results in a decreased occurrence of charge transfer, is responsible for the H atom's charge decreasing from  $-0.26$  to  $-0.36e$ . It is clear that the total charges show an upward trend in transition state III, which is betwixt the \*O and \*OOH functional groups. This observation suggests that the overall direction of electron

transmission during the hydroxyl adsorption process is from the bottom to the top. A new charge equilibrium is ultimately established by the structure following the adsorption process.

The significance of the d-band centre as a crucial electronic structural descriptor controlling the OER catalytic activity has been demonstrated by a number of earlier investigations. In order to facilitate charge transfer and catalytic performance, Guoyu *et al.* reported that the FeCu/NC catalyst has an optimally positioned d-band centre at about  $-1.806$  eV, which



correlates with its low OER overpotential and enhanced density of states (DOS) near the Fermi level.<sup>48</sup> Similar to this, Sun *et al.* showed that the OER overpotential of Ir@NiN<sub>2</sub> (0.28 V) is linearly related to the location of the metal d-band centre, indicating its crucial function in regulating the catalytic efficiency and adsorption strength.<sup>51</sup>

### Volcano plot and scaling relationship of pristine HGY and TM@HGY

The overpotential and the strength of the interaction between intermediates and active atoms are typically characterized in the OER process by examining the relationship among the Gibbs free energy of three intermediates ( $\Delta G_{\text{OH}}^*$ ,  $\Delta G_{\text{O}}^*$ , and  $\Delta G_{\text{OOH}}^*$ ). The scaling relationship between the  $\Delta G_{\text{OH}}^*$  and  $\Delta G_{\text{OOH}}^*$  of the analyzed HGY and TM-doped HGY are shown in Fig. 8(a). Among the four elementary stages, the deprotonation process ( $\Delta G_{\text{O}}^* - \Delta G_{\text{OH}}^*$ ), or the creation of OOH species ( $\Delta G_{\text{OOH}}^* - \Delta G_{\text{O}}^*$ ) typically exhibits the highest free energy shift, serving as the PDS. Additionally, the difference between  $\Delta G_{\text{O}}^*$  and  $\Delta G_{\text{OH}}^*$  can be employed as a descriptor of overpotential ( $\eta$ ) in the volcano map based on the scaling relationship between  $\Delta G_{\text{OH}}^*$  and  $\Delta G_{\text{OOH}}^*$  ( $\Delta G_{\text{OH}}^* = \Delta G_{\text{OOH}}^* + 3.2$  eV). The scaling relation, where  $\Delta G_{\text{OH}}^* = 0.80\Delta G_{\text{OOH}}^* + 3.64$  eV, is usually met, as Fig. 8(a) illustrates. A good match is demonstrated by the intercept and slope, which have modest deviations within a respectable range ( $R^2 = 0.85$ ).

The overpotential  $\eta_{\text{OER}}$  can be computed by comparing the potentials of  $\Delta G_{\text{OOH}}^*$  and  $\Delta G_{\text{O}}^*$ , as the potential determining stages of the OER for both pristine as well as TM-doped HGY predominantly occur during the  $\Delta G_3$  (OOH\* to O\*) transition. The  $\eta_{\text{OER}}$  was characterized using the  $\Delta G_{\text{OOH}}^* - \Delta G_{\text{O}}^*$  descriptor, and a graph of the volcano curve can be produced (refer to Fig. 8(b)). The matching overpotentials for every type are additionally presented in the volcano plot for easy viewing. Therefore, the discovery of the Pt top in HGY at the volcano's summit attests to the system's better electrocatalytic activities relative to other pristine sites and TM-anchored HGY.

### Structural along with thermal stability by MD simulations of Pt@HGY

DFT simulations are conducted at 0 K, which may not be indicative of the structure at ambient temperature or at higher temperatures. Specifically, in the case of the OER, the substrate needs to maintain its structural integrity both at normal temperatures, then during the process of desorption, ensuring that the substrate may be used again in the subsequent cycle. At the outset, the temperature of Pt@HGY was increased to 500 K in a microcanonical way in a duration of 5 ps with 1 fs time step. During the subsequent phase, the system was brought to a state of equilibrium at a temperature of 500 K by applying heat through a Nose-Hoover thermostat in a canonical way. This process was carried out using the VASP code.

Using MD simulations at an equilibrium temperature of 500 K, the bond length variation of C-C (1.32 Å) and Pt-C (1.98 Å) atomic pairs has been shown over a period of 3500 fs (Fig. 9(a)). According to the research, these atomic pairs obligated a net bond length fluctuation of 1.18% and 2.59%. The slight variation seen could be attributed to the thermal excitations causing the mobility of the Pt atom. Considering the thermal vibrations of the Pt atom, the higher temperature does not have a notable impact on the stability of the adorned structure. Therefore, the doping of Pt on the HGY system is likely to remain very stable even at elevated temperatures, which makes the system well-suited for adsorbing intermediates intricate in the OER mechanism. Fig. 9(b) displays the variations in temperature and total energy as a function of the simulated time steps. No signs of deterioration were observed in the Pt@HGY framework, confirming the thermodynamic stability of the structure at high temperatures.

A comprehensive comparison between the present work, HGY and TM-anchored HGY, and previously reported nanomaterials is presented in Table 1. Based on the results presented in Fig. 5, it can be concluded that Pt's OER activity in HGY surpasses that of both pristine and other doped TMs, such as Au, Co, Cr, and Sc.



Fig. 8 (a) Scaling relations between  $\Delta G_{\text{OH}}^*$  and  $\Delta G_{\text{OOH}}^*$ ; (b) volcano curve between negative OER overpotential ( $-\eta_{\text{OER}}$ ) and  $\Delta G_{\text{OOH}}^* - \Delta G_{\text{O}}^*$  for HGY and TM@HGY.





Fig. 9 (a) Bond length fluctuation of Pt–C and C–C atomic pairs of Pt@HGY at 500 K and (b) temperature along with total energy fluctuation with respect to the MD simulation step at 500 K for Pt@HGY.

Table 1 The computed  $\eta$  of TM@HGY is compared to previously published nanomaterials, respectively

System	$\eta$ (V)	Ref.
Co@GY	0.77	52
Ni@3B-GY	0.53	52
graN- $\gamma$ -GY@Rh	0.27	53
Co@R graphyne	0.74	54
Ni@R graphyne	0.31	54
Ni@H4,4,4-GY	0.34	55
Co@GY/Mn@GY	Co = 0.92 Mn = 2.96	56
Co@GY/Fe@GY	Co = 0.98 Fe = 1.62	56
Co@GY/Ni@GY	Co = 0.38 Ni = 0.86	56
Co@GY/Cu@GY	Co = 0.78 Cu = 1.21	56
Co@GY/Co@GY	0.55	56
Co@GY/GY	0.53	56
Ni@S-GPY	0.43	57
Pt@S-GPY	0.59	57
Ir@S-GPY	0.36	57
Pd@S-GPY	0.40	57
HGY	1.41	Present work
Pt@HGY	0.74	Present work

## Conclusions

The electronic structural properties of HGY and TM-doped HGY, as well as their respective OER performance, are analysed using DFT. The catalytic reaction mechanism is then analysed in detail for Pt@HGY, which exhibits the greatest catalytic performance. The findings indicate a strong correlation between the adsorption characteristics of the material and the position of the d band centre. The analysis of the overpotential along with d band centre illustration reveals that the activation energy barrier of the OER can be diminished by modulating

the d band centre in a direction that deviates from the Fermi level. The suitability of the d band centre position at  $-1.15$  eV for achieving a balance between dissociation as well as adsorption of OER intermediates has been demonstrated. The partial density of states and Löwdin charge transfer findings indicate that the presence of unoccupied antibonding orbitals in Pt@HGY facilitates the adsorption process with the intermediate, leading to surface delocalization. Additionally, the interaction between the metal d orbital with the middle sp orbital leads to a reduction of the d band centre and a favourable range for bond length adsorption. Moreover, it has been observed that the Pt element continues to be attached to HGY even when exposed to elevated temperatures of up to 500 K. The aforementioned outcome confirms the structure's thermal stability, which is supported by *ab initio* MD simulations. This study offers valuable insights into the catalytic performance of Pt-anchored HGY for the oxygen evolution reaction, as well as the design strategy for generating high OER activity in transition metal-anchored holey graphyne.

## Conflicts of interest

There are no conflicts to declare.

## Data availability

All data supporting this study are included within the article and the SI, and are available from the corresponding author upon reasonable request. The authors will make the data publicly available in an appropriate repository upon publication. Supplementary information is available. See DOI: <https://doi.org/10.1039/d5cp02097b>.



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