Faraday Discussions

Cite this: Faraday Discuss., 2022, 239, 51

PAPER

The kesterite–stannite structural transition as a way to avoid Cu/Zn disorder in kesterites: the exemplary case of the $Cu₂(Zn, Mn)SnSe₄$

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Received 14th February 2022, Accepted 14th April 2022 DOI: 10.1039/d2fd00042c

The solid solution series between $Cu₂ZnSnSe₄$, crystallizing in the kesterite type structure, and Cu₂MnSnSe₄, adopting the stannite type structure, i.e. Cu₂(Zn_{1-x}Mn_x)SnSe₄, was studied by a combination of neutron and X-ray powder diffraction. Powder samples with $0 \le x \le 1$ were synthesized by the solid state reaction of the pure elements and it was confirmed by wavelength-dispersive X-ray spectroscopy that each contained a homogeneous, off-stoichiometric quaternary phase. The lattice parameters and cation site occupancy factors were determined simultaneously by the Rietveld analysis of the neutron and X-ray powder diffraction data. The refined site occupancy factors were used to determine the average neutron scattering length of the cation sites in the crystal structure of the $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals, from which a cation$ distribution model was derived. For the end member $Cu₂ZnSnSe₄$, the disordered kesterite structure was confirmed and for Cu₂MnSnSe₄, the stannite structure was confirmed. The cross-over from the kesterite to stannite type structure by $\text{Zn}^{2+} \leftrightarrow$ Mn^{2+} substitution in the Cu₂Zn_{1-x}Mn_xSnSe₄ solid solution can be seen as a cation redistribution process among the positions (0, 0, 0), (0, $\frac{1}{2}$, $\frac{1}{4}$) and (0, $\frac{1}{4}$, $\frac{3}{4}$), including Cu⁺, Zn^{2+} and Mn²⁺. The Sn⁴⁺ cation does not take part in this process and remains on the 2b site. Moreover, the cross-over is also visible in the ratio of the lattice parameters c/ (2a), showing a characteristic dependence on the chemical composition. The order parameter Q, the quantitative measure of Cu/B^{II} disorder (B^{II} = Zn and Mn), shows a distinct dependence on the Mn/(Mn + Zn) ratio. In Zn-rich $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals, the order parameter $Q \sim 0.7$ and drops to $Q \sim 0$ (complete Cu/B^{II} disorder) in the compositional region $0.3 \ge x \ge 0.7$. In Mn-rich Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals, adopting the stannite type structure, the order parameter reaches almost $Q \sim 1$ (order). Thus, it can be concluded that only Mn-rich $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals do not$ **PAPER**
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show Cu/B^{II} disorder. A similar trend of the dependence on the chemical composition of both Cu/B^{II}-disorder and the band gap energy E_q in Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals was observed.

Introduction

In the last decade, quaternary chalcogenides, especially kesterite type semiconductors like $Cu₂ZnSn(S,Se)₄$ (CZTSSe), have gained a lot of attention as they can be used as absorbers in thin film solar cells. In contrast to the widely studied absorber material $Cu(In,Ga)Se₂$ (CIGS), kesterites consist of earth abundant and non-toxic elements. CZTSSe is applied as absorber layer in the only critical raw material free photovoltaic technology promising also low costs.

Recently, a power conversion efficiency record of 13.2% ¹ was reported for a CZTSSe-based thin film device in which, as usual, the polycrystalline CZTSSe absorber layer shows an off-stoichiometric composition.

Kesterite type compounds tolerate strong deviations from the stoichiometric composition, keeping the kesterite type structure but showing cation ratios Cu/ $(Zn + Sn)$ and Zn/Sn of lower or higher than one. This structural flexibility originates from the propensity of the kesterite type structure to stabilize copper vacancies, anti-sites, and interstitials. The charge balance is ensured by appropriate substitutions on the cation sites. Thus, the off-stoichiometric composition corresponds to certain point defects. This correlation is the basis for the offstoichiometry type model,² realized in off-stoichiometry types named A to $L³$ These intrinsic point defects (vacancies, anti-sites and interstitials) determine the electronic properties of the semiconductor significantly.⁴ It is agreed in the literature that the large band tailing observed in kesterite type semiconductors causes voltage losses, limiting the device efficiency.⁵ Faraday Discussions

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Additionally, a special kind of structural disorder, Cu/Zn disorder⁶ (Cu_{Zn} and Zn_{Cu} anti-sites in lattice planes perpendicular to the crystallographic \vec{c} -axis at z $=$ $\frac{1}{4}$ and $\frac{3}{4}$, $\{004\}$ lattice planes), is always present in kesterite type compounds^{3,7–12} and has been discussed as a possible reason for band tailing as well.⁵ But to separate the Cu_{Zn} and Zn_{Cu} anti-site point defects which cause the Cu/Zn disorder from point defects correlated with the off-stoichiometric composition of the kesterite type semiconductor according to the off-stoichiometry type is a very complex problem.

Only a few experimental methods allow the determination of the Cu/Zn disorder directly. Conventional X-ray diffraction methods fail, because Cu^{1+} and Zn^{2+} are isoelectronic cations and thus have essentially the same scattering factor for X-rays. But anomalous X-ray diffraction and neutron diffraction overcome this problem, and especially neutron diffraction has been proven to be an extremely efficient experimental method to determine the distribution of isoelectronic cations in a crystal structure. Neutron diffraction has been applied to study the crystal structure, structural disorder and point defects in a number of kesterite type materials^{3,7-12}

In order to minimize or avoid Cu/Zn disorder, cation mutation strategies can be applied. Alloying the crystal structure of the material can change it from the kesterite to stannite type to completely avoid this disorder. Both crystal structures are based on tetrahedra in which each anion is surrounded by two Cu^+ , one

Table 1 Overview of the kesterite type and stannite type crystal structures with Cu $_2$ ZnSnSe4 (kesterite type)⁶ as exemplary cases. The colour code of the corresponding cations and anions in the crystal structure figures is given in the table

divalent and one Sn^{4+} cation $(Cu₂B^HSnX$ cation tetrahedron), but these tetrahedra are differently arranged. In doing so, the stacking sequence of the lattice planes perpendicular to the crystallographic \vec{c} -axis changes from Cu–Sn/Cu–Zn/Sn–Cu/ Zn–Cu/Cu–Sn (with disorder in the Cu–Zn planes) in the kesterite type structure to B^H –Sn/Cu–Cu/Sn– B^H /Cu–Cu/ B^H –Sn (with B^H = alloyed divalent cation) in the stannite type structure. Moreover, the crystal symmetry changes from $I\bar{4}$ (kesterite type structure) to $I\bar{4}2m$ (stannite type structure). Details of the two crystal structures are summarized in Table 1.

Such a cross-over from the kesterite to the stannite type structure was previously described for the $Cu_2(Zn_{1-x}Fe_x)SnS_4$ solid solution series.⁶ Whereas kesterite type $Cu₂ZnSnS₄$ shows Cu/Zn disorder, stannite type $Cu₂FeSnS₄$ does not show this type of cation disorder. Due to the presence of Cu/Zn disorder in $Cu₂ZnSnS₄$, the crystal structure of the material was named the disordered kesterite type structure.¹³

The quaternary compound semiconductor $Cu₂MnSnSe₄$ adopts the stannite type structure.¹⁴ By substituting Zn^{2+} by Mn^{2+} ($\text{Zn}^{2+} \leftrightarrow \text{Mn}^{2+}$) in kesterite type Cu₂ZnSnSe₄³, the Cu/Zn disorder should be suppressed. Thus, within the Cu₂(- $Zn_{1-x}Mn_x)SnSe_4$ solid solution series, a cross-over between these two crystal structures is expected and stannite type mixed crystals should not show Cu/Mn disorder.

In this work, we correlate the crystal structure, Cu/Zn disorder and intrinsic point defect scenario in $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals with the band gap energy of the quaternary semiconductors. Neutron diffraction enables us to differentiate the isoelectronic cations Cu^+ and Zn^{2+} , as well as the electronically similar Mn²⁺, in the crystal structure analysis.¹⁵ The neutron diffraction data have been refined by Rietveld analysis using the kesterite as well the stannite type crystal structures as starting models. The average neutron scattering length analysis method¹⁵ was applied to determine the cation distribution within the crystal structure, resulting in a cation distribution model for the $Cu_2(Zn_{1-x}Mn_x)$ SnSe₄ solid mixed crystals. These investigations enabled us to deduce the complex cation re-distribution process within the crystal structure which transforms the kesterite to the stannite type structure. Faraday Discussions Wew Article counter

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Experimental

Because neutron diffraction experiments require large sample volumes, powder samples covering the complete $\text{Cu}_2(\text{Zn}_x\text{Mn}_{1-x})\text{SnSe}_4$ solid solution series were synthesized by solid state reactions of the pure elements, which is a proven reliable method for a systematic approach.^{3,6-11} Weighted stoichiometric mixtures of elemental zinc (5N), copper (5N), tin (5N), manganese (5N) and selenium (6N) were placed in a pyrolytic graphite boat and sealed in evacuated quartz ampoules. Each tube was placed in a one-zone furnace and heated at a rate of 10 K $\mathrm{h}^{-1},$ with intermediate temperature steps of 250 \degree C and 450 \degree C held for 48 h each, up to a final temperature of 750 \degree C. After keeping the final temperature for 300 h, the silica tubes with the samples were cooled to room temperature dynamically by switching off the furnace. An annealing procedure of the samples followed to reach a uniform chemical composition of the quaternary phase. In doing so, the powder was homogenized (by grinding in an agate mortar and pressing pellets) and the sample pellets were annealed for 300 h in evacuated silica tubes in a one-

Fig. 1 BSE micrograph of the Cu₂(Zn_{1-x}Mn_x)SnSe₄ sample with $x = 0.9$. Grains of the quaternary phase, with the chemical composition $Cu_{2.02}Zn_{0.10}Mn_{0.90}Sn_{0.99}Se_4$, as well as grains of the secondary phase are visible (gray grains). The black part of the image is the epoxy in which the grains are embedded.

zone furnace at the same final temperature as before. The silica tubes underwent dynamic cooling. The last annealing step was performed at 600 $^{\circ}$ C (300 h), again followed by dynamic cooling.

Backscattered electron (BSE) micrographs (Fig. 1) revealed the presence of different phases in the Mn-rich samples.

Wavelength dispersive X-ray spectroscopy (WDX) was applied to determine the composition of the phases present using an electron microprobe analysis system. In order to obtain reliable results from the WDX measurements, the microprobe system was calibrated using NIST elemental standards. A high accuracy of the

Table 2 Overview of the synthesized $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals: cation ratios$ $Cu/(Zn + Mn + Sn)$ and $(Zn + Mn)/Sn$ of the quaternary phase obtained from the electron microprobe analysis (WDX spectroscopy, 2% error coming from the measurements should be considered), occurring secondary phases, final formulae of the quaternary phase and the corresponding off-stoichiometry types are given

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compositional parameters was achieved by averaging over 10 local measured points within one grain and averaging over more than 30 grains of the quaternary phase showing the same compositional values. In cases of secondary phases being present, additional grains containing this phase were measured in the same way. Using the measured atomic percentages of the elements, the cation ratios Cu/(B^{II} + Sn) and B^{II}/Sn (B^{II} = Zn and Mn) were calculated. Based on these cation ratios, the chemical compositions of the $Cu_2(Zn_xMn_{1-x})SnSe_4$ mixed crystals were deduced by a procedure described elsewhere.⁸ An overview of the compositions of the quaternary mixed crystals, as well as the occurring secondary phases (when present), is given in Table 2. Faraday Discussions Wew Article on
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points within one grain and averaging over more than 30 grains of the quaternary
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The compositional and phase analysis revealed striking differences between Zn-rich and Mn-rich mixed crystals. Zn-rich mixed crystals form homogeneous, single phase $Cu_2(Zn_xMn_{1-x})SnSe_4$ compounds with an off-stoichiometric composition and no secondary phases are present. Mn-rich mixed crystals also form homogeneous, single phase $Cu_2(Zn_xMn_{1-x})SnSe_4$ compounds with an offstoichiometric composition but binary secondary phases are present. The offstoichiometry type of the $Cu_2(Zn_xMn_{1-x})SnSe_4$ mixed crystals was determined from the cation ratios according to a procedure reported previously.⁸ The cation ratio plot (Fig. 2) gives an overview of the synthesized $Cu_2(Zn_xMn_{1-x})SnSe_4$ mixed crystals.

X-ray diffraction (XRD) measurements of the polycrystalline powders were performed on a Panalytical X'Pert Pro MPD in the X-ray CoreLab of the Helmholtz-Zentrum Berlin für Materialien und Energie (HZB) using Bragg-Brentano geometry with a focusing beam and Cu- $K_{\alpha(1+2)}$ radiation. An instrumental setup with a sample spinner, a $\frac{1}{4}$ divergence slit, and a 5 mm mask were chosen. The measurements were performed over a 2θ range from 10 \degree to 130 \degree with steps of 0.013° and a measuring time of 100 s per step. The X-ray diffraction patterns of the

Fig. 2 Cation ratio plot showing the $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals synthesized. All of$ the Zn-rich mixed crystals are located between the G and F type lines. The Mn-rich $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals show a wider spread over the Cu-rich, B^{II}-rich area of$ the cation ratio plot. Each point refers to a homogeneous, single phase quaternary phase. The off-stoichiometry types refer to Table 4 and ref. 12.

polycrystalline powder samples within this study were analyzed by LeBail refinements¹⁶ using the FullProf Suite software package¹⁷ in order to obtain the lattice parameters of the $\text{Cu}_2(\text{Zn}_x\text{Mn}_{1-x})\text{SnSe}_4$ phase (see Table 3). Using the lattice parameters a and c, the tetragonal distortion $c/2a$ was calculated. The distinct change in the composition dependent $c/2a$ behaviour ("kink") is the first hint of the kesterite–stannite structural transition (Fig. 3). A similar type of "kink" was previously observed in the solid solution series $\text{Cu}_2(\text{Zn}_{1-x}\text{Fe}_x)\text{SnS}_4$.⁶

Neutron powder diffraction data were collected at the Berlin Research Reactor BER II at HZB using the fine resolution powder diffractometer FIREPOD (E9) ($\lambda =$ 1.798 \AA ; ambient temperature).¹⁸ Data analysis was done by full pattern Rietveld refinement¹⁹ using the FullProf Suite software package.¹⁷ The kesterite type structure (space group $I\bar{4}$) with Cu on Wyckoff positions 2a:(0, 0, 0) and 2c:(0, $\frac{1}{2}$, $\frac{1}{4}$), Zn on 2d: $(0, \frac{1}{2}, \frac{3}{4})$, Sn on 2b: $(\frac{1}{2}, \frac{1}{2}, 0)$ and Se on 8g: (x, y, z) , \circ as well as the stannite type structure (space group $I\bar{4}2m$)⁶ with Mn on Wyckoff position 2a:(0, 0, 0), Cu on 4d: $(0,\frac{1}{2},\frac{1}{4})$, Sn on 2b: $(\frac{1}{2},\frac{1}{2},0)$ and Se on 8i: (x,x,z) , were used as structural models in the Rietveld refinements. The kesterite type structure was applied for mixed crystals with $0 \leq Mn/(Mn + Zn) \leq 0.5$ and the stannite type structure for mixed crystals with $0.7 \leq Mn/(Mn + Zn) \leq 1.0$. The mixed crystal with $Mn/(Mn + Zn) =$ 0.6 was refined with both structure models separately. Global profile refinement parameters included a scale factor, a 2θ zero parameter, and parameters of a peak shape function (U, V, W) using a peak full width at half-maximum (FWHM) function of the form FWHM² = U tan 2θ + V tan θ + W. The peak shape function was a Thompson–Cox–Hastings pseudo-Voigt function for the neutron data and a pseudo-Voigt-peak shape for the X-ray data. Profile refinement parameters for the structure were unit cell parameters, metal site occupancies, anion coordinates and isotropic temperature factors. The refinement of the site occupancy factors (SOF) was done without any chemical constraints.

Examples of neutron diffraction patterns and the corresponding Rietveld analysis of the data are shown in Fig. 4.

Diffuse reflectance spectroscopy (DRS) measurements were carried out in air at room temperature by a spectrophotometer equipped with an integrating sphere (PerkinElmer UV/Vis-spectrometer Lambda 750S). The wavelength range of the

Fig. 3 Lattice parameters a and c (top) and tetragonal distortion $c/2a$ (bottom) of Cu₂(- Zn_xMn_{1-x})SnSe₄ mixed crystals depending on the chemical composition. The lines are guides for the eye.

measurement was adjusted to 800–1800 nm with a step size of 1 nm. Tauc plots were obtained by plotting $(F(R) \times hv)^2$ versus the photon energy.²⁰ The linear part of the curve was extrapolated to the baseline, and the optical band gap was extracted from the value of the intersection, the same way it was previously reported for $Cu₂ZnSnS₄$, $Cu₂ZnSnSe₄$ and $Cu₂ZnGeSe₄$ powder samples.⁴

Results

The cation distribution model

The average neutron scattering length analysis method $15,21$ was applied to determine the distribution of the cations Cu^+ , Zn^{2+} , Mn²⁺ and Sn⁴⁺ on the four cation sites of the kesterite type structure (2a, 2b, 2c and 2d), as well as on the three cation sites of the stannite type structure (2a, 2b and 4d). The experimental average neutron scattering lengths \bar{b}_j (exp) were calculated according to eqn (1)

Fig. 4 Rietveld refinement of the neutron diffraction data for $Cu_{2.01}Zn_{0.62}Mn_{0.42}Sn_{0.98}Se_4$ (left) and Cu_{2.05}Mn_{0.98}Sn_{0.99}Se₄ (right). The red dots are experimental data, the black line is the obtained fit, the blue ticks are the Bragg peak positions of the kesterite and stannite type structures, respectively, the red ticks are the Bragg peak positions of the MnSe phase, and the blue line is the difference between the experimental and calculated data.

(kesterite structure) and eqn (2) (stannite structure) using the cation site occupancy factors (SOFs) determined by the Rietveld refinement of neutron diffraction data and the neutron scattering length b of the cations ($b_{Cu} = 7.718$ fm, $b_{Zn} = 5680$ fm, $b_{\text{Mn}} = -3.73$ fm, $b_{\text{Sn}} = 6.225$ fm).²²

Each SOF is indexed by the Wyckoff letter of the corresponding site $(SOF_i with$ $j = 2a$, 2b, 2c and 2d for the kesterite type structure and $j = 2a$, 4d and 2b for the stannite type structure).

$$
\begin{array}{l}\n\overline{b}_{2a}(\exp) = \text{SOF}_{2a} \times b_{Cu} \\
\overline{b}_{2c}(\exp) = \text{SOF}_{2c} \times b_{Cu} \\
\overline{b}_{2d}(\exp) = \text{SOF}_{2d} \times b_{Zn} \\
\overline{b}_{2b}(\exp) = \text{SOF}_{2b} \times b_{Sn}\n\end{array} \tag{1}
$$

$$
\begin{array}{l}\n\overline{b}_{2a}(\exp) = \text{SOF}_{2a} \times b_{\text{Mn}} \\
\overline{b}_{4d}(\exp) = \text{SOF}_{4d} \times b_{\text{Cu}} \\
\overline{b}_{2b}(\exp) = \text{SOF}_{2b} \times b_{\text{Sn}}\n\end{array} \tag{2}
$$

The experimental average neutron scattering lengths calculated by eqn (1) and (2) are compared with the neutron scattering lengths of the cations copper, zinc, manganese and tin ($b_{\text{Cu}}, b_{\text{Zn}}, b_{\text{Mn}}, b_{\text{Sn}}$). At a first glance, it becomes obvious that the experimental average neutron scattering length of site 2b $(\bar{b}_{\rm 2b}(\rm exp))$ is almost equal to the neutron scattering length of tin (within an experimental uncertainty) in all the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals studied (see Fig. 5). This is a strong indication that site 2b is occupied almost exclusively by tin; the resulting cation distribution on site 2b can be seen in Fig. 6.

But the situation is significantly different for the other cation sites. The value of the experimental average neutron scattering length of site 2a $(\bar{b}_{2\text{a}}(\text{exp}))$ is almost constant at values of b_{Cu} within $0 \le x \le 0.6$, and decreases drastically afterwards (Fig. 5). This is a strong indication that the site 2a is almost exclusively occupied by Cu in Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals with $0 \le x \le 0.6$. The strong decrease of $\bar{b}_{2a}(\exp)$ for $x > 0.6$ can be explained by the substitution of copper on 2a by manganese. Mn has a negative scattering length, so even a small amount of manganese on the copper site 2a has an enormous influence on the average neutron scattering length of this site. The values of the average neutron scattering lengths $\bar{b}_{2c}(\exp)$ and $\bar{b}_{2d}(\exp)$ for the cation 2c and 2d sites decrease constantly, as long as $\bar b_{\rm 2a}(\rm exp)$ remains constant. This behaviour can be explained by the incorporation of manganese on these cation sites (2c and 2d). As soon as the average neutron scattering length of site 2a $(\bar{b}_{\rm 2a} ({\rm exp}))$ decreases, the values of $\bar{b}_{\rm 2c}({\rm exp})$ and $\bar{b}_{\rm 2d}({\rm exp})$ (and later $\bar{b}_{\rm 4d}({\rm exp}))$ increase (see Fig. 5). Faraday Discussions Wew Article on 14 compared average neutron scattering lengths calculated by eqn (1) and compares and tin $(b_{\rm Cu}, b_{\rm Zn}, b_{\rm Mn}, b_{\rm M})$. At a first glance, it becomes obvious that the experimental averag

Due to the Cu/Zn disorder present in kesterite type materials, the situations at the 2c and 2d sites are even more complex. Considering the values of the experimental average neutron scattering lengths \bar{b}_{2c} (exp) and \bar{b}_{2d} (exp), it can be concluded that the copper 2c site is partly occupied by zinc and/or manganese, respectively, decreasing the experimental average neutron scattering length with increasing Mn content in the $\text{Cu}_2(\text{Zn}_{1-x}\text{Mn}_x)\text{SnSe}_4$ mixed crystals. On the other hand, the experimental average neutron scattering length of the zinc 2d site

Fig. 5 Average neutron scattering length \bar{b}_j of the cation sites $j = 2a$, 2c and 2d for Cu₂(Zn_xMn_{1-x})SnSe₄ mixed crystals with $0 \le x \le 0.6$, as well as $j = 2a$ and 4d for Cu₂(- $\text{Zn}_x\text{Mn}_{1-x}$)SnSe₄ mixed crystals with 0.7 $\leq x \leq 1$ (full symbols – experimental values $(\bar{b}_i$ (exp)); star – calculated values (\bar{b}_i (calc)) according to the chosen cation distribution model (Fig. 6); lines - neutron scattering length of the elements Cu, Zn and Mn²²).

indicates a partial occupation by copper first and manganese later, first increasing and later decreasing the average neutron scattering length with increasing Mn content in the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals.

Considering the experimentally determined chemical composition of the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals as obtained by WDX analysis, and assuming that all four cation sites of the kesterite type structure and all three cation sites of the stannite type structure are fully occupied by the cations, the respective cation occupancy of each of the crystallographic sites can be derived with high reliability. Vacancies are not considered here, because the respective off-stoichiometry types of the quaternary mixed crystals (D, F, G, K, I) are not expected to contain vacancies.¹²

Based on a cation distribution model which is derived from the considerations above and taking into account the according off-stoichiometry type that the mixed crystal belongs to (Table 4), an average neutron scattering length can be calculated according to eqn (3):

$$
\overline{b}_j(\text{calc}) = \sum_k (k)_j^{\text{model}} \times b_k \tag{3}
$$

where $(k)_{j}^{\text{model}}$ is the fraction of the cation k on the crystallographic site j (2a, 2b, 2c, 2d for the kesterite type structure and 2a, 4d, 2b for the stannite type structure) and b_k is the neutron scattering length of the cation k.

In this calculation, the requirement that the sum of a cation species in the cation distribution model is the cation species amount measured by wavelength

Fig. 6 Cation distribution model for the end members $Cu₂ZnSnSe₄$ and $Cu₂MnSnSe₄$, as well as the Cu₂(Zn_xMn_{1-x})SnSe₄ mixed crystals, as determined experimentally by neutron diffraction. The cation sites 2a, 2b, 2c and 2d of the kesterite type structure are shown.

dispersive spectroscopy (WDX) in the chemical analysis of the sample has to be taken into account.

$$
\sum_{j} (k)^{\text{model}} = k^{\text{WDX}} \tag{4}
$$

In order to establish a reasonable cation distribution model, the experimental average neutron scattering lengths \bar{b}_j (exp) and the calculated average neutron scattering lengths \bar{b}_j (calc) are compared, minimizing their difference by varying the cation fractions $(k)_{j}^{\text{model}}$ (agreement within the experimental error). The final cation distribution model for the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals is presented in Fig. 6.

The cation re-distribution process in $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals$

On the basis of the cation distribution model, it was possible to deduce the cation re-distribution process which transforms kesterite type $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals to stannite type mixed crystals when increasing the Mn content.

In detail, this transformation goes in the following way: when starting to substitute Zn by Mn in Cu₂ZnSnSe₄ (Zn²⁺ \leftrightarrow Mn²⁺), the substituted Mn occupies the 2c and 2d sites (with a strong preference for the 2d site), whereas the 2a site is fully occupied by copper. When increasing the Mn content above $x = 0.4$ in the

Fig. 7 Comparison of the cation distribution in the Cu₂(Zn_xMn_{1-x})SnSe₄ (above) and $Cu₂(Zn_xMn_{1-x})SnSe₄$ (below) solid solution series. For simplicity, the 2c and 2d sites of the kesterite type structure are transformed into the equivalent 4d site of the stannite type structure (right figures). The cation distribution on the 2a sites of the kesterite and stannite type structures are shown in the left figures.

 $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals, the substituted manganese occupies the 2a$ site as well, in addition to the 2c and 2d sites. When a manganese content of $x >$ 0.7 is reached in the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals, the substituted manganese occupies the 2a site, and the 2c and 2d sites are (nearly) only occupied by copper. It has to be taken into account that the 2c and 2d sites of the kesterite type structure are equivalent to the 4d site of the stannite type structure. Thus, an occupation of the 4d site by copper only and an occupation of the 2a site by the divalent cations of zinc and manganese indicate that the cation re-distribution process is finished and the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals with $x > 0.7$ adopt the stannite type crystal structure. It should be noted that the strongest change of the lattice parameter ratio $c/2a$ occurs also in this compositional range. **Paper**

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The cation re-distribution process deduced for $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ is quite complex and different from the re-distribution process in the $Cu_2(Zn_xFe_{1-x})SnS_4$ series described in the literature;⁶ a comparison is shown in Fig. 7. In Cu₂(Zn₋ xFe_{1-x} SnS₄, the substituted iron (according to $\text{Zn}^{2+} \leftrightarrow \text{Fe}^{2+}$) is never present at the 2c and 2d sites (or the corresponding 4d site). The substituted iron always occupies the 2a site. This difference between the cation re-distribution processes may be related to the different divalent cations of iron and manganese, as well the different anions in the two solid solution series. But there is also a common feature of the cation re-distribution process in both series: for $x \ge 0.7$, the 4d site is occupied only by Cu, indicating the finished cross-over from the kesterite type to the stannite type crystal structure.

Intrinsic cation point defects in $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals

In addition to the cross-over from the kesterite to stannite type structure, which influences the cation distribution in $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals strongly, changes in the intrinsic cation point defects can be observed.

Fig. 8 Concentration of Cu_{2d} and B^{II}_{2c} point defects (related to Cu/B^{II} disorder) and Cu_{Sn}, Cu_i, II_{Cu}, II_{Sn} and II_i point defects (related to the off-stoichiometry type), as well as the order parameter Q calculated from the Cu and B^{II} distribution on the Wyckoff positions 2c and 2d (according to egn (5)) dependent on the Mn/(Mn + Zn) ratio.

Based on the analysis of the WDX data, most of the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals can be located in the G–F area of the cation ratio plot (see Fig. 2), whereas the Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystal with $x = 0.8$, as well as the end member of the series, Cu₂MnSnSe₄, are in the K–D region. The Cu₂($\text{Zn}_{1-x}\text{Mn}_x\text{)SnSe}_4$ mixed crystal with $x = 0.9$ is in the F–I area of the cation ratio plot. The obtained intrinsic point defect scenario in the $Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals is summarized in$ Table 5. According to the cation distribution model derived from the average neutron scattering length analysis, all expected off-stoichiometry type related intrinsic cation point defects have been revealed (see Fig. 8). Thus, it can be concluded that the cation re-distribution process does not affect the offstoichiometry type related intrinsic point defect scenario.

Cu/Zn disorder in Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals

Since the discovery of the Cu/Zn disorder during the investigation of the Cu₂($\text{Zn}_{1-x}\text{Fe}_x\text{SnS}_4$ solid solution series by neutron diffraction,⁶ it is well known that Cu and Zn are disordered in the 2c and 2d sites of the kesterite type structure, forming Cu_{Zn} and Zn_{Cu} anti-sites. In case of the Cu₂(Zn_{1-x}Mn_x)SnSe₄ solid solution studied here, Cu/B^{II} disorder has to be discussed ($B^H = Zn$ and Mn), formed by Cu_{2d} and B_{2c}^{II} anti-site defects. It has to be noted, that the concentration of these anti-site point defects is much higher than the concentration of the offstoichiometry type related defects (see Fig. 8).

The presence of Cu/Zn disorder in the end member $Cu₂ZnSnSe₄$ is very visible by the significant difference between the experimental average neutron scattering lengths of the 2c and 2d sites and the neutron scattering length of the respective cations, b_{Cu} and b_{Zn} . It can be found that $\bar{b}_{2c}(\exp) < b_{\text{Cu}}$ and $\bar{b}_{2d}(\exp) > b_{\text{Zn}}$. These scattering length differences can be explained by the formation of intrinsic point defects. In this case, zinc on copper anti-site defects (Zn_{Cu}) would decrease the average neutron scattering length of the 2c site because $b_{\text{Zn}} = 5.680 \text{ fm} < b_{\text{Cu}} =$ 7.718 fm; likewise, the copper on zinc anti-site defect (Cu_{Zn}) would lead to an

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increase of the average neutron scattering length of the 2d site, in this way forming the Cu/Zn disorder. The resulting cation distribution on the 2c and 2d Wyckoff sites can be seen in Fig. 6 (the equivalent 4d site is shown in Fig. 7). Similar behaviour of the experimental average neutron scattering lengths $\bar{b}_{\text{2c}}(\text{exp})$ and $\bar{b}_{2d}(\exp)$ can be recognized for the Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystal with x = 0.1 (see Fig. 6 and 7).

The cation re-distribution process due to the cross-over from the kesterite to the stannite type crystal structure in the $Cu_2(Zn_{1-r}Mn_x)SnSe_4$ mixed crystals with $0.1 < x < 0.7$ introduces the necessity of extending the concept of Cu/Zn disorder to consider both divalent cations (B^H with B=Zn and Mn). Cu/ B^H disorder would be formed by equal amounts of Cu_{2d} and $\text{B}^\text{II}_{\text{2c}}$ anti-site defects, respectively. The fact that in $\text{Cu}_2(\text{Zn}_{1-x}\text{Mn}_x)\text{SnSe}_4$ mixed crystals with 0.4 < x < 0.7, the divalent cations also partly occupy the 2a site (besides copper), is related to the cation redistribution process, because the 2a site in the stannite type structure is occupied by the divalent cation exclusively. **Paper**

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Using the derived cation distribution and taking into account the cation redistribution process, the order parameter Q of the Cu/B^{II} disorder in kesterite type $Cu₂(Zn,Mn)SnSe₄ mixed crystals was calculated according to$

$$
Q = \frac{[Cu_{2c} + B_{2d}^{II}] - [B_{2c}^{II} + Cu_{2d}]}{[Cu_{2c} + B_{2d}^{II}] + [B_{2c}^{II} + Cu_{2d}]}.
$$
\n(5a)

Eqn (5a) was modified from the definition of the order parameter in the literature²³ to introduce the two divalent cations present in $\text{Cu}_2(\text{Zn}_{1-x}\text{Mn}_x)\text{SnSe}_4$ mixed crystals. The order parameter is defined so that $Q = 0$ for complete disorder and $Q = 1$ for order.

The order parameter Q , calculated by eqn $(5a)$, is shown in Fig. 9. In the end member Cu₂ZnSnSe₄ and the mixed crystal with $x = 0.1$, significant Cu/B^{II} disorder is present. Kesterite type $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals with 0.1 < x \leq 0.7 (compositional region of the cation re-distribution process) show full Cu/B^{II} disorder ($Q = 0$). In the compositional region where the Cu₂(Zn_{1-x}Mn_x)SnSe₄

Fig. 9 Band gap energy E_q and order parameter Q of Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals dependent on the Mn/(Mn + Zn) ratio.

Table 4 Cation substitutions leading to the off-stoichiometry types F, G, K, D and I. Given is the cation substitution reaction, the expected defect complex and the corresponding chemical formula for $Cu_{2}(Zn_{1-x}Mn_{x})SnSe_{4}$. Here, II stands for B^{II}

mixed crystals adopt the stannite type structure $(0.7 < x < 1)$, another type of cation disorder is present. The Cu/B^{II} disorder occurs in the $\{010\}$ lattice planes (between the 2a and 4d sites). The order parameter Q of the Cu/B^{II} disorder in stannite type mixed crystals was calculated according to

$$
Q = \frac{[Cu_{4d} + B_{2a}^{II}] - [B_{4d}^{II} + Cu_{2a}]}{[Cu_{2a} + B_{4d}^{II}] + [B_{2a}^{II} + Cu_{4d}]}.
$$
 (5b)

It increases drastically to values of $Q > 0.8$, reaching a $Q \sim 1$, for the end member $Cu₂MnSnSe₄$. This compositional dependent behaviour of the order parameter Q shows that the Cu/B^{II} disorder in Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals is influenced by the cation re-distribution process. The detailed cation distributions, as well as the Q values, are presented in Table 5. It should be noted that the concentration of anti-site point defects correlated to Cu/B^H disorder (Fig. 8) is almost an order of magnitude higher than the concentration of each of the offstoichiometry type related defects.

Band gap energy $E_{\rm g}$

In an attempt to correlate the obtained intrinsic cation point defect concentrations and the Cu/B^{II} disorder expressed by the order parameter *O* with the optoelectronic properties of $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals, the optical band gap energy $E_{\rm g}$ was determined by analyzing the DRS data. The results are summarized in Table 5.

The dependence of the band gap energy E_g on the Mn/(Zn + Mn) ratio in $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals is shown in Fig. 9. It is found that the band gap energy E_g varies between 0.90 eV for Cu₂ZnSnSe₄ and 1.21 eV for Cu₂MnSnSe₄ with a non-linear behaviour.

A comparison of the band gap energy E_{φ} dependence on the Mn/(Zn + Mn) ratio to the dependence of the order parameter Q on the Mn/(Zn + Mn) ratio (Fig. 9) highlights a very similar trend for both dependences. In both cases, the values (E_{φ}) and Q) slightly increase with the beginning of the $\text{Zn}^{2+} \leftrightarrow \text{Mn}^{2+}$ substitution. In the compositional region of the cation re-distribution process, both E_g and Q drop and stay almost constant. As soon as the $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals adopt the stannite type structure, $E_{\rm g}$ and Q increase.

Conclusions

Wavelength dispersive X-ray spectroscopy, X-ray and neutron diffraction, and diffuse reflectance spectroscopy were used to investigate the $Cu_2(Zn_xMn_{1-x})SnSe_4$ solid solution series, grown by solid state reaction.

The site occupancy factors, obtained by simultaneous Rietveld analysis of the neutron and X-ray diffraction data were used to determine the average neutron scattering length of the cation sites in kesterite type $Cu₂ZnSnSe₄$, stannite type $Cu₂MnSnSe₄$ and $Cu₂Zn_{1-x}Mn_xSnSe₄$ mixed crystals, giving new insights into the cation distribution.

The cross-over from the kesterite to the stannite type structure by $\text{Zn}^{2+} \leftrightarrow \text{Mn}^{2+}$ substitution in the $Cu_2Zn_{1-x}Mn_xSnSe_4$ solid solution series can be seen as a cation re-distribution process among the positions at $(0, 0, 0)$, $(0, \frac{1}{2}, \frac{1}{4})$ and $(0, \frac{1}{4}, \frac{3}{4})$, which includes Cu $^{\dagger},$ Zn $^{2+}$ and Mn $^{2+}.$ The Sn $^{4+}$ cation does not take part in this process and remains on the 2b site. Moreover, the cross-over is visible in the ratio of the lattice parameter ratio $c/(2a)$, showing a characteristic dependence on the chemical composition. A similar behaviour of the lattice parameter ratio was observed for the $Cu_2(Zn_{1-r}Fe_r)SnS_4$ solid solution series.⁶ **Proper**
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Starting the $\text{Zn}^{2+} \leftrightarrow \text{Mn}^{2+}$ substitution in the $\text{Cu}_2\text{Zn}_{1-x}\text{Mn}_x\text{SnSe}_4$ solid solution, substituted manganese occupies the 2c & 2d sites, whereas the 2a site is fully occupied by copper. When the substituted manganese in $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals exceeds 40% ($x \ge 0.4$), it also occupies the 2a site, besides the 2c and 2d sites. When reaching a manganese content of 70% ($x = 0.7$), the substituted manganese overwhelmingly occupies the cation 2a site, whereas the amounts of manganese on the 2c and 2d sites strongly decrease. Thus, these two sites (2c and 2d), which are equivalent to the 4d site in the stannite type structure, are nearly only occupied by copper. It can be concluded that $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals with a manganese content of more than 70% ($x \ge 0.7$) adopt the stannite type structure, finishing the structural cross-over.

The off-stoichiometry type related intrinsic point defects have been deduced from the cation distribution model, as well the Cu/B^{II} disorder (B^{II} = Zn and Mn). Interestingly, the order parameter Q, the quantitative measure of Cu/B^{II} disorder, shows a distinct dependence on the Mn/(Mn + Zn) ratio. In Zn-rich $Cu₂(Zn₁$ $_{\rm -x}$ Mn_x)SnSe₄ mixed crystals, the order parameter shows Q \sim 0.7 and drops to Q \sim 0 (complete Cu/B^{II} disorder) in the compositional region 0.3 $\geq x \geq 0.7$. In Mn-rich $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals, adopting the stannite type structure, the order parameter reaches $Q \sim 1$ (order). Thus, it can be concluded that only Mnrich Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals do not show Cu/B^{II} disorder.

The room temperature DRS measurements revealed a strong effect of the order parameter Q on the electronic properties of the studied $Cu_2(Zn_{1-x}Mn_x)SnSe_4$ mixed crystals. A similar trend of the dependence on the chemical composition

for both Cu/B^{II}-disorder and band gap energy E_g in Cu₂(Zn_{1-x}Mn_x)SnSe₄ mixed crystals was observed.

Author contributions

Galina Gurieva: design of the study, ND measurements and data analysis, conceptualization & writing - original draft, writing - review & editing. Sara Niedenzu: solid state synthesis, WDX & XRD measurements and data analysis. Nikita Siminel: DRS methodology, DRS measurements and data analysis. Alexandra Franz: support in ND measurements (instrument scientist). Susan Schorr: conceptualization, supervision, methodology, writing – review & editing.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

We thank HZB for the allocation of neutron diffraction beamtime. We thank the Freie Universitaet Berlin for the allocation of WDX measuring time. The research leading to the presented results has been partially supported by the CUSTOM-ART project funded from the European Union's Horizon 2020 Research and Innovation Programme under the grant agreement no. 952982. Also, the research leading to the presented results has been partially supported by the INFINITE-CELL projects funded from the European Union's Horizon 2020 Research and Innovation Programme under the Marie Skłodowska-Curie grant agreement no. 777968. Faraday Discussions

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 Author contributions

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Notes and references

- 1 J. Zhou, X. Xu, B. Duan, H. Wu, J. Shi, Y. Luo, D. Li and Q. Meng, Nano Energy, 2021, 89, 106405.
- 2 A. Lafond, L. Choubrac, C. Guillot-Deudon, P. Deniard and S. Jobic, Z. Anorg. Allg. Chem., 2012, 638, 2571–2577.
- 3 G. Gurieva, L. E. Valle Rios, A. Franz, P. Whitfield and S. Schorr, J. Appl. Phys., 2018, 123, 161519.
- 4 S. Schorr and G. Gurieva, in Crystallography in Materials Science, ed. S. Schorr and C. Weidenthaler, De Gruyter, 2021.
- 5 T. Gokmen, O. Gunawan, T. K. Todorov and D. B. Mitzi, Appl. Phys. Lett., 2013, 103, 103506.
- 6 S. Schorr, H.-J. Hoebler and M. Tovar, Eur. J. Mineral., 2007, 19, 65–73.
- 7 L. E. Valle Rios, K. Neldner, G. Gurieva and S. Schorr, J. Alloys Compd., 2016, 657, 408.
- 8 G. Gurieva, D. M. Többens, S. Levcenco, T. Unold and S. Schorr, J. Alloys Compd., 2020, 846, 156304.
- 9 G. Gurieva, D. M. Toebbens, M. Ya. Valakh and S. Schorr, J. Phys. Chem. Solids, 2016, 99, 100–104.
- 10 R. Gunder, J. A. Marquez-Prietro, G. Gurieva, T. Unold and S. Schorr, CrystEngComm, 2018, 20, 1491.

- 11 S. Schorr, Sol. Energy Mater. Sol. Cells, 2011, 95(6), 1482.
- 12 S. Schorr, G. Gurieva, M. Guc, M. Dimitievska, A. Perez-Rodriguez, V. Izquierdo-Roca, C. Schnohr, J. Kim, W. Jo and J. M. Merino, J. Phys.: Energy, 2020, 2, 012002.
- 13 M. Ya Valakh, V. M. Dzaghan, I. S. Babichuk, X. Fontane, A. Perez-Rodriguez and S. Schorr, JETP Lett., 2013, 98, 255–258.
- 14 V. Sachanyuk, I. Olekseyuk and O. Parasyuk, Phys. Status Solidi A, 2006, 203(3), 459–465.
- 15 S. Schorr, in Neutron Applications in Materials for Energy, ed. G. J. Kearley and V. Peterson, Springer, 2015.
- 16 A. Le Bail, Powder Diffr., 2005, 20(4), 316–326.
- 17 J. Rodriguez-Carvajal and T. Roisnel, <https://www.ill.eu/sites/fullprof/>.
- 18 Helmholtz-Zentrum Berlin fur Materialien und Energie, E9: The Fine ¨ Resolution Powder Diffractometer (FIREPOD) at BER II, Journal of large-scale research facilities JLSRF, 2017, 3, A103.
- 19 H. M. Rietveld, J. Appl. Crystallogr., 1969, 2, 65–71.
- 20 J. Tauc, R. Grigorovici and A. Vancu, Phys. Status Solidi B, 1966, 15, 627–637.
- 21 S. Schorr et al., X-Ray and Neutron Diffraction on Materials for Thin-Film Solar Cells, in Advanced Characterization Techniques for Thin Film Solar Cells, ed. D. Abou-Ras, T. Kirchartz and U. Rau, Wiley-VCH Verlag GmbH & Co. KGaA, 2011, p. 347. Franch Schott, Sol. Energy Mater. Sol. Cells, 2011, 95(6), 1492.

15 S. Schorn, G. Gurieva, M. Gue, M. Dimitievska, A. Perez-Rodriguez, V. Izquierdo-Roca, C. Schnohr, J. Kim, W. Jo and J. M. Merino, *J. Phys.*

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	- 22 A.-J. Dianoux and G. Lander, Neutron Data Booklet, Old City, Philadelphia, 2003.
	- 23 D. M. Többens, G. Gurieva, S. Levcenko, T. Unold and S. Schorr, Phys. Status Solidi B, 2016, 253, 1890–1897.