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## ARTICLE

## Effect of variations in manufacturing and material properties on the self-folding behaviors of hydrogel and elastomer bilayer structures

Jiayu Zhao <sup>a</sup>, Hesaneh Kazemi <sup>b</sup>, Hyunsun Alicia Kim <sup>b</sup>, Jinhye Bae <sup>a,c,d,e\*</sup>

The stimuli-responsive self-folding structure is ubiquitous in nature, for instance, the mimosa folds its leaves in response to external touch or heat, and the Venus flytrap snaps shut to trap the insect inside. Thus, modeling self-folding structures has been of great interest to predict the final configuration and understand the folding mechanism. Here, we apply a simple yet effective method to predict the folding angle of the temperature-responsive nanocomposite hydrogel/elastomer bilayer structure manufactured by 3D printing, which facilitates the study of the effect of the inevitable variations in manufacturing and material properties on folding angles by comparing the simulation results with the experimentally measured folding angles. The defining feature of our method is to use thermal expansion to model the temperature-responsive nanocomposite hydrogel rather than the nonlinear field theory of diffusion model that was previously applied. The resulted difference between the simulation and experimentally measured folding angle (*i.e.*, error) is around 5%. We anticipate that our method could provide insight into the design, control, and prediction of 3D printing of stimuli-responsive shape morphing (*i.e.*, 4D printing) that have potential applications in soft actuators, robots, and biomedical devices.

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Electronic Supplementary Information (ESI) available: See the supplementary material for obtaining the temperature-dependent thermal expansion coefficient of NC-PNIPAM (Fig. S1), the stress-strain curves for calculating E of NC-PNIPAM and PDMS (Fig. S2), the photographs for calculating Poisson's ratios of NC-PNIPAM at swelled and deswelled states, respectively (Fig. S3). The optical microscope images showing the dimension (Fig. S4) and folding angle (Fig.S5).

## 1 Introduction

2 Self-folding structures that are activated in response to  
3 external stimuli are of interest for their applications in self-  
4 assembly,<sup>1</sup> soft actuators,<sup>2</sup> biomedical devices,<sup>3</sup> and  
5 wearable devices.<sup>4</sup> Bilayer structure is one of the most  
6 commonly used designs to create a self-folding structure.<sup>5</sup>  
7 However, the entire bulk bilayer structure will undergo large  
8 deformation (*i.e.*, bending) when actuated, limiting the  
9 ability to form a more complex final configuration. Inspired  
10 by the ancient art of origami, hinge-based bilayer structures  
11 can greatly simplify the design space by localizing the  
12 deformation to hinges.<sup>6</sup> In these structures, the strain-  
13 mismatch generated between the active component of the  
14 hinge and passive component of the substrate in response  
15 to environmental cues, including temperature,<sup>7</sup> moisture,<sup>8</sup>  
16 light,<sup>9</sup> and electricity,<sup>10</sup> will result in folding of the structure.  
17 One of the most widely used active materials in self-folding  
18 structures is stimuli-responsive hydrogels, which are  
19 chemically or physically crosslinked hydrophilic polymers  
20 that can have volume expansion when immersed in water  
21 due to water absorption. This characteristic makes hydrogels  
22 a suitable choice for the active component of the hinge-  
23 based bilayer structure. Crosslinked poly(N-  
24 isopropylacrylamide) (PNIPAM) is a well-known thermo-  
25 responsive hydrogel that exhibits lower critical solution  
26 temperature (LCST) at around 32 °C,<sup>11</sup> which is close to the  
27 physiological temperature, making it a suitable material for  
28 biomedical applications. PNIPAM hydrogels can reversibly  
29 expand or shrink their volume by controlling the  
30 temperature below or above LCST, respectively.<sup>11</sup> Recently,  
31 we have reported the thermally responsive self-folding  
32 structure using the nanocomposite PNIPAM hydrogel as an  
33 active hinge and polydimethylsiloxane (PDMS) as a passive  
34 substrate.<sup>7</sup> Although we experimentally showed that the  
35 folding angle can be programmed with prescribed geometric  
36 parameters (*i.e.*, PDMS thickness and hinge width), their  
37 self-folding behavior has not yet been fully explored,  
38 especially in terms of the inevitable variations in  
39 manufacturing and measured material properties. Modeling  
40 the self-folding structures would allow us to understand and  
41 predict the folding process more accurately by providing  
42 insight into how the variations raised from material  
43 properties and the manufacturing process would influence  
44 the folding angles, therefore making it possible to precisely  
45 control the folding structure towards the programmed  
46 shape, enabling complex final configurations in various  
47 applications including soft robotics, biomedical devices, and  
48 aerospace. To date, Guo *et al.* demonstrated modeling of the  
49 programmable deformation of origami structures with  
50 temperature-sensitive hydrogels,<sup>12</sup> where the nonlinear field  
51 theory of coupled diffusion and deformation is used to  
52 model the hydrogel. However, the accuracy of their model  
53 remains unknown because the predicted shape deformation  
54 was not directly compared with the experimental results.  
55 Tang *et al.* adopted thermal expansion to model the shape  
56 morphing of the thermal responsive magnetic  
57 hydrogel/elastomer bilayer structures. Their simulation

58 results exhibited similar final configurations to the  
59 experimental results,<sup>13</sup> however, they didn't further  
60 examine the results quantitatively. Therefore, the  
61 quantitative accuracy of the simulations compared to the  
62 experimental results of the hydrogel/elastomer material  
63 systems has not been well investigated to our best  
64 knowledge.

65 In recent years, self-folding structures fabricated by additive  
66 manufacturing (3D printing) provoke lots of interest,  
67 because it allows for fast prototyping of various kinds of  
68 materials with spatially programmed compositions and  
69 microstructures,<sup>14-16</sup> enabling functional materials with new  
70 properties that cannot be fabricated using conventional  
71 manufacturing techniques. Especially, 3D printing of active  
72 materials gives rise to "4D printing", with the 4<sup>th</sup> dimension  
73 being the time, the 3D printed object can have shape  
74 transformation over time in response to external stimuli.<sup>17,</sup>  
75 <sup>18</sup> Theoretical models have been developed for different  
76 material systems to guide the structural design and predict  
77 the final configurations.<sup>19, 20</sup> However, despite the recent  
78 advances, the understanding of 3D printing imperfection,  
79 specifically, the dimension difference between the printed  
80 and designed structures, on the shape transformation  
81 remains limited. Moreover, it has not been investigated how  
82 the folding angle would be influenced due to the inevitable  
83 variations in material properties of 3D printed samples.

84 In this work, we study the range of uncertainty observed in  
85 both manufacturing (*i.e.*, 3D printing) and sample-to-sample  
86 variation in material properties on the folding angle of the  
87 nanocomposite PNIPAM hydrogel/PDMS bilayer structures.  
88 We characterize the self-folding structures fabricated by  
89 extrusion-based 3D printing, quantify their responses by  
90 thermal actuation, model their self-folding behavior and  
91 quantify the error. We employ a thermal expansion model  
92 to predict the folding angle of the hinge-based bilayer  
93 structure of nanocomposite PNIPAM hydrogel/PDMS.  
94 Compared to the previously reported nonlinear field theory  
95 for modeling the thermal responsive hydrogels/PDMS  
96 bilayer structures,<sup>12</sup> where the energy function depends on  
97 the number of chains per polymer volume, the volume of a  
98 solvent molecule and the Boltzmann constant, our method  
99 is much simpler and computationally efficient while in good  
100 agreement with the experimental data (folding angle  
101 difference ~ 5%). As a result, the predicted folding angles  
102 using the average Young's modulus (E) of the  
103 nanocomposite hydrogel agree reasonably well (*i.e.*, error ~  
104 5%) with the experimentally obtained values, given the  
105 variabilities associated with the 3D printing process.  
106 Furthermore, the possible reasons causing the deviation  
107 between the computational and experimental results are  
108 discussed from both manufacturing and material aspects.  
109 Examining these factors is important in enabling the  
110 facilitation of self-folding structure design and providing a  
111 deeper insight into their folding mechanism. We anticipate  
112 that our work can contribute to the fundamental  
113 understanding to support the programming and  
114 manufacturing of shape transformations produced by  
115 thermal-responsive material systems.

1

**2 Experiments****3 Materials**

4 N-isopropylacrylamide (NIPAM, stabilized with 4-  
5 methoxyphenol,  $M_w = 113.16 \text{ g mol}^{-1}$ ) was purchased from  
6 Tokyo Chemical Industry (TCI) America. PDMS (Sylgard 184)  
7 was purchased from Dow Corning. N,N'-  
8 methylenebisacrylamide (BIS), Irgacure 2959, and  
9 benzophenone were purchased from Sigma-Aldrich (St  
10 Louis, MO, USA). Fumed silica nanoparticles (SiNPs, CAB-O-  
11 SIL EH-5) were purchased from Cabot Corporation. Nanoclay  
12 (NC, Laponite-RD) was obtained by BYK Additives &  
13 Instruments. All chemicals were used as received without  
14 further purification.

15

**16 Preparation of PDMS precursor inks**

17 The PDMS precursor inks were prepared by a simple one-pot  
18 mixing process containing PDMS base/crosslinker (10:1),  
19 benzophenone (1.8 wt% with respect to PDMS base), and  
20 SiNPs (15 wt% with respect to PDMS base) in a Thinky  
21 planetary mixer (Thinky U.S.A., Inc.) mixing at 2000 rpm for  
22 3 min and followed by degas process at 2000 rpm for 2 min  
23 to remove any air bubbles. Since benzophenone is in the  
24 solid state at room temperature, to achieve better mixing  
25 quality, it was heated in an oven at 70 °C for 10 min to melt  
26 before adding to the PDMS precursor. After mixing, the inks  
27 were loaded into a 10 mL syringe (Fisher Scientific) and  
28 centrifuged at 4000 rpm for 20 min to eliminate any air  
29 bubbles.

30

**31 Preparation of NC-PNIPAM precursor inks**

32 NIPAM solution (2 M) and BIS solution (0.13 M) were  
33 prepared by adding NIPAM and BIS to deionized (DI) water  
34 respectively and mixed in a vortex mixer until all chemicals  
35 were dissolved at room temperature. Next, NIPAM solution  
36 (10 mL, 2 M), BIS solution (120  $\mu\text{L}$ , 0.13 M), Irgacure 2959  
37 (0.04 g), and NC (1 g) were added into a 35 mL container  
38 (Thinky U.S.A., Inc.) and mixed at 2000 rpm for 5 min or  
39 longer until the solution was mixed well with no visible NC  
40 aggregates. Finally, the mixed ink was loaded into a 10 mL  
41 syringe (Fisher Scientific) and centrifuged at 2000 rpm for 10  
42 min to eliminate any visible air bubbles.

43

**44 Extrusion-based 3D printing and fabrication**

45 NC-PNIPAM/PDMS bilayer structures were fabricated by  
46 extrusion-based 3D printing using a 3D printer (Rokit,  
47 Invivo). The 10 mL syringe stored with PDMS precursor ink  
48 was placed in the extrusion carriage of the 3D printer and  
49 printed on the glass slides (75 mm  $\times$  50 mm  $\times$  1 mm) using a  
50 20-gauge blunt tip dispensing needle (0.6 mm inner  
51 diameter). The printed PDMS substrate with a hinge  
52 structure was cured in an oven at 80 °C for 30 min.  
53 Subsequently, NC-PNIPAM precursor ink was directly  
54 printed on the hinge section of the cured PDMS substrate,  
55 the printed bilayer structure was then transferred into a  
56 homemade transparent humid box to prevent NC-PNIPAM  
57 from drying out while UV irradiation (365 nm) with the

58 intensity of 253 mW  $\text{cm}^{-2}$  for 2 min and 22 s (UV source  
59 provided by Omnicure).

60

**61 Characterization**

62 The rheological data were obtained using a rheometer (TA  
63 Instruments™ Discovery™ HR-30) using a 40 mm plate. The  
64 viscosity measurements were conducted using flow sweep  
65 mode with the shear rate ranging from 0.1 to 100  $\text{s}^{-1}$ . The  
66 storage and loss moduli were measured using oscillation  
67 mode at a frequency of 1 Hz with strain ranging from 1 to  
68 100%. Stress-strain data was obtained using a universal  
69 testing machine (Instron Corp., Instron 5982) with a strain  
70 rate of 10 mm  $\text{s}^{-1}$ . To prepare samples for the tensile test,  
71 the NC-PNIPAM were 3D printed to rectangular shapes (35  
72 mm  $\times$  10 mm  $\times$  0.6 mm), and after photo-crosslinking, they  
73 were swelled or de-swelled at 22 °C or 45 °C water bath,  
74 respectively. The average values of E for each condition were  
75 calculated based on the tensile test results of 5 samples. The  
76 PDMS inks (PDMS precursor + 15 wt% SiNPs) were 3D  
77 printed to rectangular shapes (35 mm  $\times$  10 mm  $\times$  0.6 mm)  
78 and thermally cured. The average E was obtained by the  
79 tensile test of 3 samples.

80 For calculation of the thermal expansion coefficient  $\alpha$  of NC-  
81 PNIPAM, the NC-PNIPAM hydrogel was fabricated into a  
82 rectangular rod-like shape (35 mm  $\times$  2 mm  $\times$  0.6 mm) using  
83 3D printing. After cross-linking, the NC-PNIPAM was first de-  
84 swelled in DI water at 45 °C for at least 48 hours to reach its  
85 equilibrium state, and the length was measured. It was then  
86 swelled in DI water at temperatures of 40.2, 38.1, 34.7, 28.6,  
87 and 22 °C, and the resulting lengths were measured,  
88 respectively. Optical micrographs were captured using an  
89 optical microscope (Keyence VHX1000).

90

**91 Simulation**

92 We applied a thermal expansion model to examine the  
93 folding angle of the hinge-based bilayer structure of NC-  
94 PNIPAM/PDMS due to its similarity with the isotropic  
95 swelling/deswelling of the temperature-responsive PNIPAM  
96 hydrogel. The thermal strain in the general form can be  
97 written as,

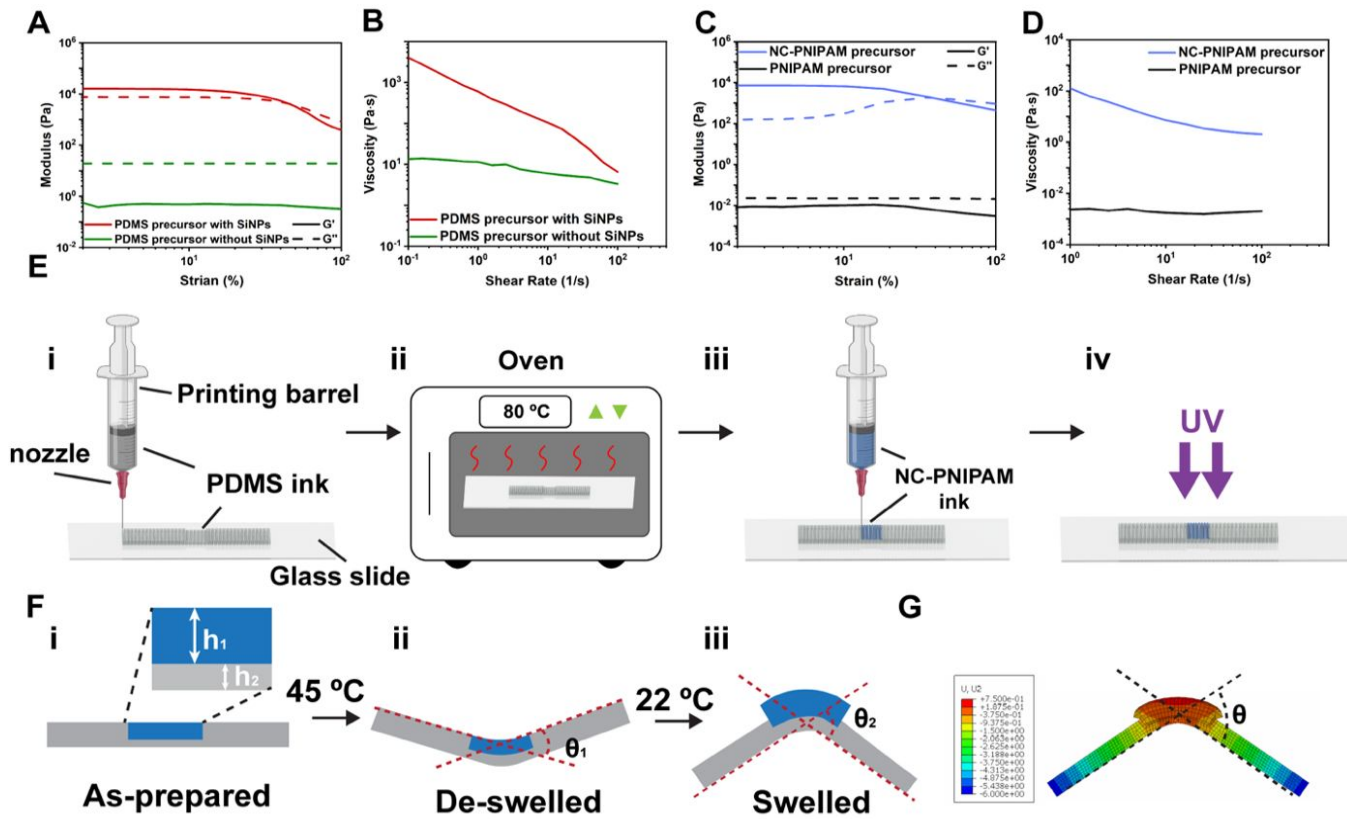
$$\epsilon_T = \alpha \Delta T = [\alpha_1 \quad \alpha_2 \quad \alpha_3] \Delta T \quad (1)$$

98 where  $\Delta T$  is the temperature change, and  $\alpha$  is the thermal  
99 expansion coefficient. We assume isotropic thermal  
100 expansion.

101 We also assume that the volume change of NC-PNIPAM,  
102 which is initially at 45 °C and then placed in water of 22 °C,  
103 is the result of thermal expansion only. We obtained the  
104 thermal expansion coefficients for NC-PNIPAM at multiple  
105 temperatures based on the experiment described in the  
106 previous section and by performing a curve fitting. Since NC-  
107 PNIPAM expands once cooled, all of these values are  
108 negative. Because the PDMS does not elongate once put in  
109 cooler water (*i.e.*, 22 °C), we use  $\alpha = 0$  for the PDMS  
110 substrate.

111 We use the neo-Hookean hyperelastic model available in  
112 Abaqus,<sup>21</sup> where the strain energy function is given by

$$W = \frac{\mu}{2} (\bar{I}_1 - 3) + \frac{k}{2} (J - 1)^2 \quad (2)$$



**FIG. 1.** Log-log plots of (A) modulus as a function of stress and (B) viscosity as a function of shear rate for PDMS precursor with and without SiNPs. Log-log plots of (C) modulus as a function of stress and (D) viscosity as a function of shear rate for PNIPAM and NC-PNIPAM precursor. (E) Schematic illustration of the fabrication process. (i) 3D printing of PDMS ink into a cuboid ( $10 \times 20 \times 0.8 \text{ mm}^3$ ) with a hinge structure in the middle, (ii) and the printed structure was then transferred to an oven to cure at  $80 \text{ }^\circ\text{C}$  for 30 min. (iii) The NC-PNIPAM ink was directly printed onto the hinge structure of the PDMS substrate. (iv) The NC-PNIPAM precursor was photo-crosslinked using UV irradiation. (F) Schematic illustration of the cross-sectional view of a hinge-based bilayer structure of NC-PNIPAM (blue)/PDMS (gray) (i) as prepared with the thickness of NC-PNIPAM and PDMS denoted as  $h_1$  and  $h_2$ , respectively; (ii) de-swelled at  $45 \text{ }^\circ\text{C}$  water bath with a negative folding angle  $-\theta_1$  and (iii) swelled at  $22 \text{ }^\circ\text{C}$  with a positive folding angle  $\theta_2$ , respectively. (G) Finite element analysis (FEA) of the hinge-based bilayer structure with a folding angle  $\theta$ .

1 In the above expression,  $\mu$  is the shear modulus,  $k$  is the bulk  
2 modulus, and  $I_1$  is the first strain invariant, defined as

$$I_1 = \text{trace}(\bar{\mathbf{B}}) \quad (3)$$

3 where  $\bar{\mathbf{B}} = \bar{\mathbf{F}}\bar{\mathbf{F}}^T$  denotes the deviatoric stretch matrix,  $\bar{\mathbf{F}} =$

4  $J^{-\frac{1}{3}}\mathbf{F}$  is the distortional component of the deformation

5 gradient defined as  $\mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}}$ , and  $J = \det(\mathbf{F})$  is the

6 determinant of the deformation gradient. The second Piola–

7 Kirchhoff stress  $\mathbf{S}$  can be computed as

$$\mathbf{S} = \mathbf{F}^{-1} \frac{\partial W}{\partial \mathbf{F}} \quad (4)$$

8 The static equilibrium of the unit cell under finite  
9 deformation is given by

$$\mathbf{R} = \mathbf{F}_{ext} - \mathbf{F}_{int} = \mathbf{0} \quad (5)$$

10 where  $\mathbf{R}$  is the residual force,  $\mathbf{F}_{ext}$  is the external force, and

11  $\mathbf{F}_{int}$  is the internal force. This equation can be discretized  
12 using the finite element method and be written as

$$\mathbf{r} = \mathbf{f}_{ext} - \mathbf{f}_{int}(\mathbf{u}) = \mathbf{0} \quad (6)$$

13 where  $\mathbf{r}$  is the residual nodal force vector,  $\mathbf{f}_{ext}$  is the

14 external nodal force vector,  $\mathbf{f}_{int}$  is the internal nodal force

15 vector that depends on the nodal displacement vector,  $\mathbf{u}$ .

16 The internal nodal force vector is defined by

$$\mathbf{f}_{int}(\mathbf{u}) = \frac{\partial (\int_V W(\mathbf{u}) dV)}{\partial \mathbf{u}} \quad (7)$$

17 and can be solved iteratively for the displacement using the

18 Newton–Raphson method.<sup>22</sup> We use E and Poisson’s ratios

19 of NC-PNIPAM and PDMS obtained from the tensile test. We

20 performed the simulation in Abaqus. We first created the

21 structure as a 3D deformable solid, based on the geometry

22 of the printed structure. We then assigned material

23 properties of PDMS and NC-PNIPAM to corresponding

24 sections, using elastic isotropic materials with the

25 experimentally measured E, Poisson’s ratio, and  $\alpha$ . When

26 assigning material properties for the NC-PNIPAM, we note

27 the values of E and Poisson’s ratios that we measured in the

28 initial and final temperatures (*i.e.*,  $45$  and  $22 \text{ }^\circ\text{C}$ , respectively)

29 are different. Therefore, we use temperature-dependent

30 material properties, which assume a linear relationship

31 between the E and temperature, and between Poisson’s

32 ratio and temperature. We also use a temperature-

33 dependent thermal expansion coefficient. For the final step,

34 we use a predefined temperature field of  $22 \text{ }^\circ\text{C}$ . We used the

35 encastre boundary conditions to fix the symmetry plane of

36 the structure in the hydrogel section. We created a mesh

37 using hexahedral elements.

1

## 2 Results and Discussion

3 The fabrication and synthesis of the precursor inks in detail  
 4 were reported in our previous paper.<sup>7</sup> Here, we briefly  
 5 summarize this process. The hinge-based bilayer structures  
 6 of the NC-PNIPAM/PDMS are fabricated by extrusion-based  
 7 3D printing, which is one of the most commonly used in  
 8 additive manufacturing.<sup>23</sup> Extrusion-based 3D printing  
 9 greatly enlarges the design space for patterning viscous  
 10 material (ink) into a 3D structure in a layer-by-layer manner.  
 11 The printable ink should possess shear-thinning behavior to  
 12 facilitate the extrusion, and solid-like behavior with storage  
 13 modulus ( $G'$ ) > loss modulus ( $G''$ ) to maintain the shape  
 14 retention after deposition.<sup>16</sup> We formulated both the PDMS  
 15 and NC-PNIPAM precursor inks using the compositions from  
 16 our previous paper,<sup>7</sup> so that they not only meet the  
 17 rheological requirements to allow the extrusion-based 3D  
 18 printing, but also ensure strong adhesion between the  
 19 hydrophobic PDMS and hydrophilic NC-PNIPAM with the  
 20 adhesion strength greater than the fracture strength of the  
 21 NC-PNIPAM hydrogel, which was 14.8 kPa.<sup>7</sup> Specifically, the  
 22 PDMS precursor ink is composed of PDMS precursor (10:1  
 23 base: crosslinker), benzophenone to create covalent  
 24 bonding between PDMS and NC-PNIPAM, and SiNPs as the  
 25 rheological modifier. It is noted that the PDMS precursor  
 26 without SiNPs possesses liquid-like behavior ( $G'' > G'$ ) (Fig.  
 27 1A) and low viscosity  $\eta = 13$  at  $0.1 \text{ s}^{-1}$  (Fig. 1B). After the  
 28 addition of SiNPs, the network formed between the silanol  
 29 groups on the surface of SiNPs<sup>24</sup> endows the PDMS ink with  
 30 solid-like behavior as  $G'$  (16710 Pa) >  $G''$  (8013 Pa) (Fig. 1A)  
 31 and shear-thinning properties as  $\eta = 4052 \text{ Pa}\cdot\text{s}$  at  $0.1 \text{ s}^{-1}$  and  
 32  $6.6 \text{ Pa}\cdot\text{s}$  at  $100 \text{ s}^{-1}$  (Fig. 1B). The NC-PNIPAM precursor ink is  
 33 composed of NIPAM as a monomer, BIS as a crosslinker,  
 34 Irgacure 2959 as a photoinitiator, and NC as a rheological  
 35 modifier. NC is known to form a so-called “house-of-cards”  
 36 structure driven by the electrostatic forces between its  
 37 positively charged surfaces and negatively charged edges.<sup>25</sup>  
 38 The addition of NC transforms the PNIPAM precursor from a  
 39 liquid-like fluid with low viscosity to a solid-like paste with  
 40 shear-thinning properties, where  $G'$  (7263 Pa) >  $G''$  (158 Pa)  
 41 (Fig. 1C) and  $\eta = 7063 \text{ Pa}\cdot\text{s}$  at  $0.1 \text{ s}^{-1}$  and  $6.9 \text{ Pa}\cdot\text{s}$  at  $100 \text{ s}^{-1}$   
 42 (Fig. 1D).

43 The PDMS precursor ink was first printed into a cuboid ( $10 \times$   
 44  $20 \times 0.8 \text{ mm}^3$ ) with a hinge structure in the middle (Fig. 1E-i)  
 45 and cured in an oven heated to  $80 \text{ }^\circ\text{C}$  (Fig. 1E-ii). In the  
 46 following text, we denote the cured PDMS precursor ink as  
 47 PDMS. The NC-PNIPAM precursor ink was then printed onto  
 48 the hinge structure of the PDMS substrate (Fig. 1E- iii) and  
 49 photo-crosslinked by UV light ( $365 \text{ nm}$ ,  $253 \text{ mW cm}^{-2}$ ) for 142  
 50 s (Fig. 1E-iv). The resulting NC-PNIPAM is a temperature-  
 51 responsive nanocomposite hydrogel with reversible  
 52 expansion and collapse of the PNIPAM network due to  
 53 swelling and deswelling by water diffusion,<sup>26</sup> respectively.  
 54 The strain-mismatch generated between active NC-PNIPAM  
 55 and passive PDMS at high and low temperatures (*i.e.*,  $45$  and  
 56  $22 \text{ }^\circ\text{C}$ , respectively) will result in the folding of the structure.  
 57 We note that the folding directions at  $45$  and  $22 \text{ }^\circ\text{C}$  are  
 58 opposite due to the deswelling and swelling of the NC-

59 PNIPAM hinge. In the as-prepared state, the bilayer  
 60 structure of NC-PNIPAM/PDMS is flat (Fig. 1F-i). The bilayer

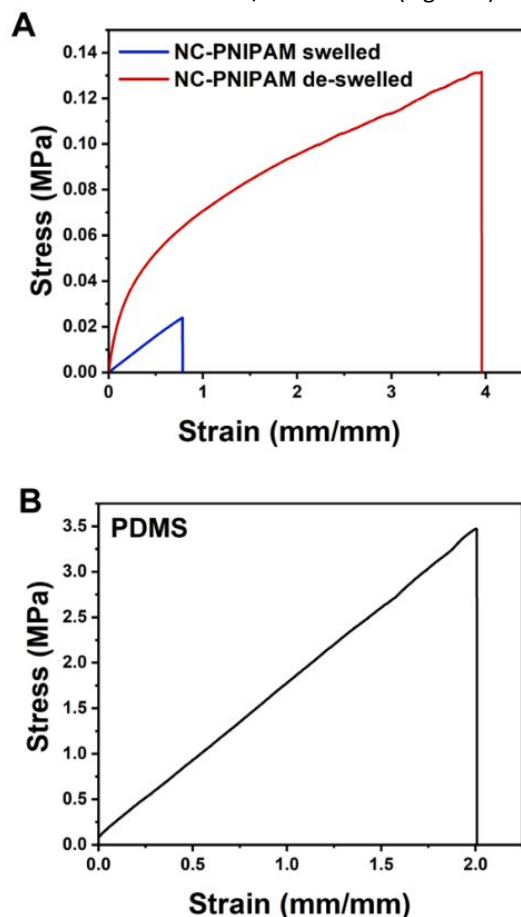
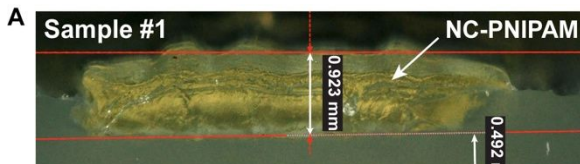


FIG. 2. (A) Stress-strain curves of the NC-PNIPAM at swelled state (blue) and de-swelled state (red). (B) Stress-strain curve of PDMS substrate.

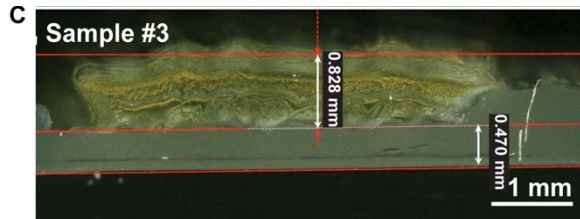
61 structure was then transferred to a  $45 \text{ }^\circ\text{C}$  water bath for 12  
 62 hours, which refers to the initial condition, to release the  
 63 residual stress generated during the fabrication or curing  
 64 process, resulting in a negative folding angle  $-\theta_1$  at the  
 65 equilibrium state due to the de-swelling of the NC-PNIPAM  
 66 (Fig. 1F-ii). After this step, the bilayer structure of NC-  
 67 PNIPAM/PDMS was transferred to a  $22 \text{ }^\circ\text{C}$  water bath for 12  
 68 hours to allow the NC-PNIPAM to reach the equilibrium state  
 69 by swelling, resulting in a positive folding angle  $\theta_2$  (Fig. 1F-  
 70 iii). The simulated folding angle  $\theta$  is compared with the  
 71 experimentally obtained total angle change  $\theta_1 + \theta_2$  to  
 72 evaluate the accuracy of the model (Fig. 1G).





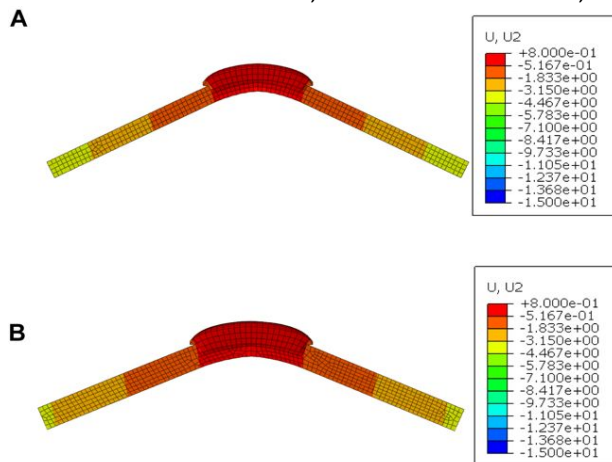
**Table 1.** Avg., min., and max. E of NC-PNIPAM de-swelled and swelled at 45 and 22 °C, respectively.

Temperature	45 °C	22 °C
Avg. E (kPa)	324	22
Min. E (kPa)	230	11
Max. E (kPa)	418	33



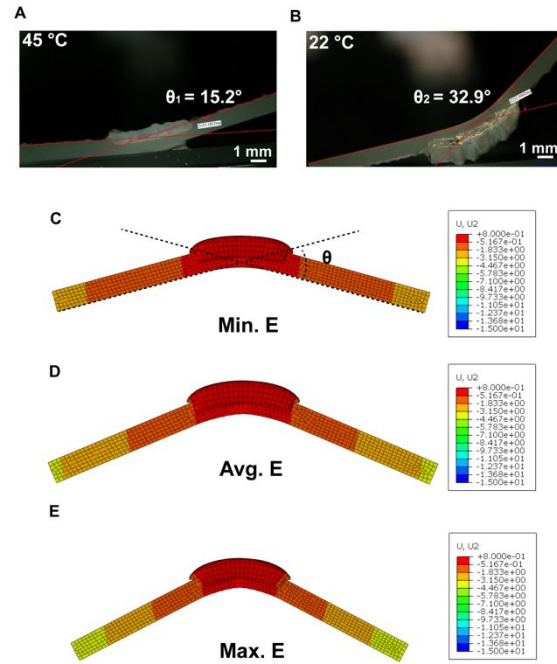
**FIG. 3.** Optical microscope photographs of the cross-sectional view of the hinge-based bilayer structure of NC-PNIPAM/PDMS printed with the target thickness ( $h_1 = 0.6$  mm,  $h_2 = 0.4$  mm). (A) Sample #1 with  $h_1 = 0.923$  mm,  $h_2 = 0.492$  mm; (B) sample #2 with  $h_1 = 0.779$  mm,  $h_2 = 0.473$  mm; (C) sample #3 with  $h_1 = 0.828$  mm,  $h_2 = 0.470$  mm.

1 To calculate the thermal expansion coefficient, we measured  
 2 the length of the rectangular rod-like shape at its equilibrium  
 3 state in the water of 45, 40.2, 38.1, 34.7, 28.6, and 22 °C and  
 4 plotted  $\frac{\Delta L}{L_0}$  vs.  $T$ , in which  $\Delta L$ ,  $L_0$ , and  $T$  refer to length change,  
 5 initial length, and temperature, respectively. We performed  
 6 a cubic curve fitting (norm of residuals = 0.02325, Fig. S1) and  
 7 obtained the thermal expansion coefficient of NC-PNIPAM  
 8 hydrogel for temperatures of 45, 40, 35, and 30 °C as -  
 9 0.0213, -0.0192, -0.0251, and -0.0392, respectively.  
 10 We performed the tensile tests to obtain E and Poisson's  
 11 ratios for NC-PNIPAM and PDMS. For the temperature-  
 12 responsive NC-PNIPAM, the tensile tests were conducted  
 13 using samples de-swelled at 45 °C and swelled at 22 °C to  
 14 match with the initial and final conditions set in the  
 15 simulation, respectively. The E can be calculated from the  
 16 initial slopes (0 ~ 0.1 mm/mm strain) of the stress-strain  
 17 curves (Fig. S2), which yield  $22 \pm 11$ ,  $324 \pm 94$ , and  $2000 \pm 188$   
 18 kPa for swelled NC-PNIPAM, de-swelled NC-PNIPAM, and



**FIG. 4.** The simulated folding structure using the (A) target thickness ( $h_1 = 0.6$  mm,  $h_2 = 0.2$  mm) and (B) actual thickness ( $h_1 = 0.779$  mm,  $h_2 = 0.473$  mm). Color bars shown on the right indicate the simulated displacement in the vertical direction.

19 PDMS, respectively. The representative plots of NC-PNIPAM  
 20 and PDMS are shown in Fig. 2A and Fig. 2B, respectively.



**FIG. 5.** Optical microscope photographs of the hinge-based bilayer structure of NC-PNIPAM/PDMS (A) de-swelled at 45 °C and (B) swelled at 22 °C. The simulated folding structure using the (C) min., (D) avg., and (E) max. E of the NC-PNIPAM. Color bars shown on the right indicate the simulated displacement in the vertical direction.

21 The calculated minimum (min.), average (avg.), and  
 22 maximum (max.) E of NC-PNIPAM was summarized in Table  
 23 1, in which the min. and max. E were calculated by  
 24 subtracting and adding the standard deviation (SD) to the  
 25 avg. E, respectively. We note that the E of NC-PNIPAM at 45 °C  
 26 is higher than the one at 22 °C, which can be attributed to  
 27 the collapsed network of PNIPAM due to de-swelling at a  
 28 higher temperature. On the other hand, Poisson's ratio  $\nu$  by  
 29 definition is the negative of the ratio of transverse strain  
 30 ( $\epsilon_{\text{trans}}$ ) to axial strain ( $\epsilon_{\text{axial}}$ ), which can be calculated using the  
 31 initial and final dimensions of the tensile tested samples (Fig.  
 32 S3),  $\nu = -\frac{(w_f - w_0)/w_0}{\epsilon_{\text{axial}}}$ , in which  $w_0$  and  $w_f$  are the initial  
 33 final width, respectively, and  $\epsilon_{\text{axial}}$  is recorded by the tensile  
 34 test machine. The Poisson's ratio was calculated based on 3  
 35 samples for each condition, yielding the value of  $0.14 \pm 0.017$   
 36 and  $0.28 \pm 0.015$  at 45 and 22 °C for NC-PNIPAM,  
 37 respectively. As for PDMS, we directly adopt Poisson's ratio  
 38 of 0.49 from the literature since it is a common material.<sup>27</sup>  
 39 We performed the Finite Element Analysis (FEA) for five  
 40 printed bilayer structures of NC-PNIPAM/PDMS with the  
 41 same programmed dimensions, NC-PNIPAM thickness ( $h_1$ ) of  
 42 0.6 mm and PDMS substrate thickness ( $h_2$ ) of 0.4 mm.  
 43 However, it turns out that each printed sample has a slightly  
 44 different thickness with  $h_1 = 0.923, 0.779, 0.828, 0.586,$  and  
 45 0.64 mm;  $h_2 = 0.492, 0.473, 0.470, 0.369,$  and 0.394 mm,  
 46 respectively (Fig. 3A-C, Fig. S4). The differences between the  
 47 target and actual thickness can be calculated as (actual  
 48 thickness - target thickness)/actual thickness, yielding values  
 49 ranging from -2 ~ 35%. This thickness variation is caused by

**Table 2.** Folding angle errors (%) using the min., avg., and max. E for sample #1, #2, and #3, respectively. Min. and Max. E are calculated as avg.E -SD and avg.E +SD, respectively.

Error	Sample #1	Sample #2	Sample #3	Sample #4	Sample #5
	$h_1=0.923\text{ mm}, h_2=0.492\text{ mm}$	$h_1=0.779\text{ mm}, h_2=0.473\text{ mm}$	$h_1=0.828\text{ mm}, h_2=0.470\text{ mm}$	$h_1=0.586\text{ mm}, h_2=0.369\text{ mm}$	$h_1=0.640\text{ mm}, h_2=0.394\text{ mm}$
Min. E	-38.93%	-37.50%	-38.10%	-28.13%	-31.74%
Avg. E	-12.21%	-4.16%	-7.16%	3.13%	2.40%
Max. E	6.87%	16.67%	12.19%	21.88%	19.45%

1 the limited printing precision when the nozzle size (0.6 mm)  
 2 is larger or close to the target dimension.<sup>28</sup> To study the  
 3 effect of 3D printing imperfection (*i.e.*, inaccurate printed  
 4 thickness) on the folding behavior, we compare the  
 5 predicted folding angles from the FEA model created using  
 6 the target thickness ( $h_1 = 0.6\text{ mm}$ ,  $h_2 = 0.4\text{ mm}$ , Fig.4A),  
 7 defined as  $\theta_t$ , with the predicted folding angles from the FEA  
 8 models created using the actual thickness measured for the  
 9 printed samples, defined as  $\theta_a$ . Fig. 4B shows an example of  
 10 the simulated sample #2 with  $\theta_a = 46^\circ$ . We note that the  
 11 materials' properties (*i.e.*, averaged E and Poisson's ratio)  
 12 were kept the same when running the simulations. The  
 13 simulation results show that  $\theta_t = 52^\circ$  and  $\theta_a$  of sample #1, #3,  
 14 #4, and #5 are  $46^\circ$ ,  $48^\circ$ ,  $66^\circ$ ,  $60^\circ$ , and  $62^\circ$ , respectively. We  
 15 denote the average of  $\theta_a$  as  $\theta_a$ . Therefore, the error caused  
 16 by the manufacturing process can be calculated as  $(\theta_t - \theta_a)/$   
 17  $\theta_a$ , which is 4.9%. Note that the printing precision can be  
 18 improved by carefully tuning the ink viscosity as well as  
 19 printing parameters including printing speed, nozzle size,  
 20 and layer height.

21 Next, we examine the effect of sample-to-sample variation  
 22 in E on the folding angle. We created five models for each of  
 23 these NC-PNIPAM/PDMS bilayer structures based on their  
 24 actual dimensions after 3D printing and curing (shown in Fig.  
 25 3 and Fig. S4). For each sample, we run the simulation using  
 26 the avg., min., and max. E of NC-PNIPAM and compare the  
 27 predicted folding angles with the experimental results.

28 The full profiles showing the folding angles of all five samples  
 29 can be found in Fig. S5. Here, we show sample #2 as one  
 30 example where the experimental folding angle for the  
 31 bilayer structure of NC-PNIPAM/PDMS is  $-15.2^\circ$  and  $32.9^\circ$  at  
 32 equilibrium de-swelled state at  $45^\circ\text{C}$  (Fig. 5A) and swelled  
 33 state at  $22^\circ\text{C}$  (Fig. 5B), respectively. Thus, the experimentally  
 34 measured folding angle is  $\theta_e = \theta_1 + \theta_2 = 15.2^\circ + 32.9^\circ = 48.1^\circ$ .  
 35 A structure with the same geometry was created in Abaqus  
 36 using the min., avg., and max. E of NC-PNIPAM (Fig. 5C-E),  
 37 and the predicted folding angles of  $30^\circ$ ,  $46^\circ$ , and  $56^\circ$  are  
 38 obtained, respectively. Therefore, the errors were calculated  
 39 as the  $(\theta_s - \theta_e)/\theta_e$ , where  $\theta_s$  is the angle from simulation,  
 40 which are  $-37.5\%$ ,  $-4.16\%$ , and  $16.67\%$  for these three  
 41 conditions, respectively. The errors for the folding angles of  
 42 all five samples are summarized in Table 2, which indicates  
 43 that the prediction using the average E of NC-PNIPAM gives  
 44 the most accurate result, with the smallest average error of  
 45 5.8%.

46 In addition to the sample-to-sample variation, another  
 47 possible reason for the discrepancy between the  
 48 experimental and computational folding could be  
 49 inconsistent environmental conditions (*i.e.*, temperature

50 and humidity) while measuring E using the tensile test  
 51 method. The samples are tested at ambient conditions  
 52 without temperature and humidity control, thus the  
 53 temperature and water content of NC-PNIPAM may vary  
 54 continuously during testing. This can lead to variations in the  
 55 degree of swelling/deswelling thus mechanical properties.  
 56 This effect could be more severe for the tensile test of NC-  
 57 PNIPAM at the de-swelled state of  $45^\circ\text{C}$ , as the temperature  
 58 will drop from  $45^\circ\text{C}$  to room temperature as soon as the  
 59 samples are taken out from the hot water bath. We  
 60 anticipate that the prediction error can be further minimized  
 61 if the error range in E can be minimized from the  
 62 measurements by better environmental control.

## 63 Conclusions

64 In summary, we have applied a simple yet effective method  
 65 by utilizing the thermal expansion model to predict the  
 66 folding angle of the temperature-responsive hinge-based  
 67 bilayer structure of NC-PNIPAM/PDMS fabricated by 3D  
 68 printing. The effect of the accuracy of 3D printed dimensions  
 69 was investigated on the folding angle. The properties of the  
 70 materials including thermal expansion coefficients, E, and  
 71 Poisson's ratios were measured experimentally and assigned  
 72 to the materials in the simulation. The simulations were  
 73 conducted using the min., avg., and max. E of NC-PNIPAM,  
 74 and the errors of the simulations conducted using the avg. E  
 75 yield around 5%. Given the variations in the printing process  
 76 and the material properties, we believe our work can lead to  
 77 new perspectives on modeling shape morphing systems of  
 78 temperature-responsive material-based structures, and  
 79 such modeling can facilitate the design, optimization, and  
 80 manufacturing of these structures that may find applications  
 81 in soft actuators/robots, biomedical devices, and drug  
 82 delivery systems.

## 83 Conflicts of interest

84 The authors declare no conflict of interest.

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