



Cite this: *J. Mater. Chem. C*, 2023, 11, 12356

## Controllable and tuneable growth of NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%)@Na(Yb/Y)F<sub>4</sub>-core@shell structures and the effect of their geometry on upconversion luminescence†

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Lanthanide doped upconversion nanoparticles (UCNPs) are a class of nanomaterials with excellent luminescence properties. The practical use of UCNPs, however, has been hindered by their relatively low upconversion (UC) quantum yields. Enhancing their efficiency, and therefore their brightness, is a critical goal for these emerging materials. In this study, a range of novel core@shell structures were synthesized using NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) as the core template to regulate the shell growth of Na(Yb/Y)F<sub>4</sub>. Here, we observed that the size and the morphology of the UCNPs can be fine-tuned by controlling the ratio of the Y<sup>3+</sup>/Yb<sup>3+</sup> ions within the shell material, resulting in small rod-like structures that form when using a high concentration of Y<sup>3+</sup> and larger hexagonal plate-like structures when using a high concentration of Yb<sup>3+</sup>. In terms of the optical properties, the UC luminescence and lifetime of the core-only and core@shell structures were measured. Overall, the emission intensity and lifetime of the UCNPs increased with the nanoparticle size. We observed that, under the same experimental conditions, larger core@shell nanostructures with a hexagonal shape showed brighter UC emission, compared to those with a small nanorod shape and core-only. Based on our findings, we propose a variety of energy transfer pathways arise through geometric alteration of the nanostructures upon varying the Y<sup>3+</sup>/Yb<sup>3+</sup> ion concentration ratios within the shell material. Our results could have important significance for understanding the relationship between the geometry of UCNPs and their UC optical properties.

Received 6th April 2023,  
Accepted 12th August 2023

DOI: 10.1039/d3tc01215h

rsc.li/materials-c

## Introduction

Lanthanide doped upconversion nanoparticles (UCNPs) are a class of luminescence nanomaterials with excellent chemical stability which convert low energy near infrared (NIR) photons into high energy visible (Vis) and ultraviolet (UV) light photons. UCNPs exhibit unique luminescent properties, including sharp excitation and emission bands, long luminescence lifetimes, large anti-Stokes shifts, low autofluorescence, no photobleaching, and high detection sensitivity.<sup>1–3</sup> For this reason, UCNPs have attracted considerable interest in recent years in many fields, including 3D display technology,<sup>4,5</sup> solar technology,<sup>6</sup> photovoltaics,<sup>7</sup> photocatalysis,<sup>8</sup> optical sensing,<sup>9–11</sup> and biomedical imaging applications.<sup>12,13</sup> Despite their potential, the practical use of UCNPs has been hindered by their relatively low upconversion (UC) quantum yield, so there is a drive to enhance this

efficiency to meet the growing demands of such applications. To this end, various approaches have been employed to tune the UC emission spectra. UCNPs are composed of a host material (typically NaYF<sub>4</sub>) in which lanthanide (Ln) ions are doped. Structural design is essential for high-efficiency UC emission, which involves the correct choice of the host material, the type of Ln ion used as dopants, their concentration, and their distribution across the UCNP.<sup>14,15</sup> Conventionally, in early UCNP development, activator dopant concentration was kept low (typically less than 3% for Er<sup>3+</sup> and less than 1% for Tm<sup>3+</sup>), while the number of sensitizer ions, Yb<sup>3+</sup>, was high to ensure an appropriate distance between the sensitizer and the activator ions to generate energy transfer and thus a bright emission.<sup>16</sup> However, with a high Yb<sup>3+</sup> sensitizer content, a concentration quenching effect process could occur, which would be detrimental to the luminescence efficiency.<sup>17–19</sup> Recently, various approaches have been developed to alleviate the threshold of concentration quenching in UCNPs. One common approach is to passivate the surface of UCNPs, creating a core@shell structure with one or more shells which can block the energy transfer path to surface quenchers and can also separate sensitizers and activators ions through the use of these multi shells.<sup>20–22</sup> Another strategy is to

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† Electronic supplementary information (ESI) available: TEM histograms, FTIR, XRD, luminescence and lifetime decay data. See DOI: <https://doi.org/10.1039/d3tc01215h>



irradiate the UCNPs with high excitation rates so that the majority of dopant ions are at excited states and the number of detrimental ground-state ions is reduced.<sup>23–25</sup> Choosing a host material with a large unit cell to maintain a uniform ion–ion distance between the crystal is probably the most common strategy to minimize ion segregation and avoids local concentration quenching.<sup>26–30</sup> Another approach is to incorporate another dopant such as Fe<sup>3+</sup>. Zhang *et al.* also observed that a high Yb<sup>3+</sup> concentration could be significantly enhanced through doping Fe<sup>3+</sup> into the host material. They claimed that in the new system, new energy levels form owing to the mixed electron wavefunctions of Yb<sup>3+</sup> and Fe<sup>3+</sup>. It was observed that the UV and blue emissions from the NaYbF<sub>4</sub>:Tm(0.5%) core doped with Fe<sup>3+</sup> (5%) were enhanced by up to 10 and 6 times, respectively, compared to those of the core sample alone.<sup>31</sup>

The relationship between the UCNP geometry with their optical properties is an area of investigation that has been drawing interest since the pathways of UC could be different for different dimensional systems depending on the ions' distribution across the host material. Therefore, there is still a lot of potential to be explored. Several groups have used synchrotron-based X-ray photoelectron spectroscopy (XPS) to analyse the distributions of Yb<sup>3+</sup> sensitizer ions within the NaLnF<sub>4</sub> core and have observed that the dopants of sensitizer and activator ions are not uniformly distributed within the host material.<sup>32,33</sup> Zhang *et al.*<sup>34</sup> also confirmed that this ‘‘heterogeneous doping’’ can further induce many negative effects on the optical properties of the UCNPs, especially the luminescent efficiency. These findings led us to consider that the sensitizer ions' distribution across the crystal lattice should play an important role.

The aim of this study was to evaluate a new tuneable growth mechanism and set of shapes for NaYbF<sub>4</sub>:Tm. To this end, a range of novel core@shell structures were synthesized by growing various NaYF<sub>4</sub>:Yb(*x*%) shell materials in which the Yb<sup>3+</sup> doping concentration was set at *x* = 0%, 20%, 40%, 60%, 80% and 100% around a core of NaYbF<sub>4</sub>:Tm(0.5%), with 5% Fe<sup>3+</sup> incorporated following the interesting findings of Zhang *et al.* Here, we have first observed that the nanoparticle size and the morphology can be tuned by precisely controlling the sensitizer concentration of Yb<sup>3+</sup> within the shell material. The resulting core@shell nanoparticles exhibited a pure hexagonal phase and increasing sizes with increasing Yb<sup>3+</sup> concentrations. Additionally, we correlated the structure with their optical properties for the range of the prepared UCNPs and have observed that the Yb<sup>3+</sup> concentration is not limited by the concentration quenching effect. We saw that a high dopant concentration of sensitizer ions enhanced the overall brightness of UCNPs. It was found that the UC intensities of the UV and blue emissions for the core@shell with 100% Yb<sup>3+</sup> were enhanced by ~75 and ~55-fold, respectively, compared to that of the core sample alone, showing also longer-lived emission lifetimes. Due to the core@shell formation, we hypothesize a new distribution for the pairs Yb<sup>3+</sup>–Tm<sup>3+</sup> and Tm<sup>3+</sup>–Tm<sup>3+</sup> within each UCNP from the core to the UCNP surface, forming distinct types of energy transfer network within the systems, depending on the introduction of Y<sup>3+</sup>/Yb<sup>3+</sup> ions. We believe that these results would complement and extend the findings on the distribution of

Y<sup>3+</sup>/Yb<sup>3+</sup> ion dopants around a core shell. We have investigated this using a NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) core, however, this could be further applied to other UCNP cores with different dopant ions.

## Experimental

### Materials

All the reagents for the synthesis of UCNPs were purchased from Sigma-Aldrich. These are: yttrium(III) chloride hexahydrate (YCl<sub>3</sub>·6H<sub>2</sub>O), ytterbium(III) chloride hexahydrate (YbCl<sub>3</sub>·6H<sub>2</sub>O), thulium(III) chloride hexahydrate (TmCl<sub>3</sub>·6H<sub>2</sub>O), iron(III) chloride hexahydrate (FeCl<sub>3</sub>·6H<sub>2</sub>O), oleic acid (OA, 90%), 1-octadecene (ODE, 90%), sodium hydroxide (NaOH, ≥98%) and ammonium fluoride (NH<sub>4</sub>F, ≥98%).

### Synthesis of NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) core UCNPs

NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) core UCNPs were synthesized using a Schlenk-line and following a previously reported protocol with slight modification to increase the scale.<sup>35</sup> In a 250 mL three-neck round-bottomed flask, YbCl<sub>3</sub>·6H<sub>2</sub>O (4.75 mmol, 1.83 g), TmCl<sub>3</sub>·6H<sub>2</sub>O (0.025 mmol, 9.60 mg), and FeCl<sub>3</sub>·6H<sub>2</sub>O (0.25 mmol, 67.58 mg) were mixed with 30 mL of OA (90%) and 60 mL of ODE (90%). The solution was then heated to 150 °C under the presence of N<sub>2</sub> gas for 2 h. After that, the solution was cooled down to room temperature, and a mixture of NaOH (12.5 mmol, 500 mg) and NH<sub>4</sub>F (20 mmol, 740.80 mg) dissolved in 30 mL of methanol was added dropwise under vigorous stirring. The mixture was stirred for 60 min at room temperature. After evaporating methanol, a vacuum was applied for 45 min at 100 °C. Finally, the solution was heated up to 305 °C and maintained for 30 min under a N<sub>2</sub> atmosphere. The solution was cooled down naturally, and UCNPs were precipitated by centrifugation and washed with ethanol three times (3000 rpm, 15 min). The as-prepared UCNPs were re-dispersed in hexane for further use.

### Synthesis of NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%)@NaYF<sub>4</sub>-core@shell UCNPs

The as-synthesised UCNPs served as a core to grow an epitaxial shell of NaYF<sub>4</sub>. The synthesis was again performed using a standard Schlenk-line following a method described in the literature with some adjustments.<sup>36,37</sup> In a 100 mL three-neck round bottom flask, YCl<sub>3</sub>·6H<sub>2</sub>O (0.5 mmol, 151.68 mg) was mixed with 6 mL of OA and 15 mL of ODE under a N<sub>2</sub> gas flow. The solution was heated to 150 °C to form a homogenous mixture. After that, the solution was cooled down to 80 °C under a steady flow of N<sub>2</sub>, and the previously prepared core UCNPs (125 mg) dispersed in hexane (20 mg mL<sup>−1</sup>) were added. After removal of the hexane, the temperature of the solution was reduced to room temperature and a mixture of NaOH (1.25 mmol, 50 mg) and NH<sub>4</sub>F (2 mmol, 74.08 mg) dissolved in 5 mL of methanol was added dropwise under vigorous stirring. The mixture was stirred for 45 min at room temperature. After evaporation of the methanol, the solution was heated up to 305 °C and maintained for 1 h and 30 min under a N<sub>2</sub> atmosphere. The as-prepared UCNPs were



cooled down naturally and were precipitated by centrifugation, and washed with ethanol three times (8000 rpm, 15 min). Finally, UCNPs were re-dispersed in hexane, and stored for further use.

### Synthesis of NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%)@NaYF<sub>4</sub>:Yb (x%)-core@shell

NaYF<sub>4</sub>:Yb (x%), where x = 20%, 40%, 60%, 80% and 100% were created around the core structure by using a similar procedure as described above. (Note, where x = 0%, the shell is denoted NaYF<sub>4</sub> and where x = 100% the shell is referred to as NaYbF<sub>4</sub>). Here, both YCl<sub>3</sub>·6H<sub>2</sub>O and YbCl<sub>3</sub>·6H<sub>2</sub>O were incorporated into the synthesis at various ratios, as stated in Table 1.

### Characterization

Transmission electron microscopy (TEM) was performed on a FEI Tecnai with an operating voltage of 120 kV. TEM grids were

**Table 1** Proportions of YCl<sub>3</sub>·6H<sub>2</sub>O and YbCl<sub>3</sub>·6H<sub>2</sub>O employed for shell growth synthesis where Yb<sup>3+</sup> doping concentration were set at 0%, 20%, 40%, 60%, 80% and 100%

Shell material	YCl <sub>3</sub> ·6H <sub>2</sub> O (mmol)	YCl <sub>3</sub> ·6H <sub>2</sub> O (mg)	YbCl <sub>3</sub> ·6H <sub>2</sub> O (mmol)	YbCl <sub>3</sub> ·6H <sub>2</sub> O (mg)
NaYF <sub>4</sub>	0.5	151.68	—	—
NaYF <sub>4</sub> :Yb(20%)	0.4	121.34	0.1	38.75
NaYF <sub>4</sub> :Yb(40%)	0.3	91.01	0.2	77.50
NaYF <sub>4</sub> :Yb(60%)	0.2	60.67	0.3	116.25
NaYF <sub>4</sub> :Yb(80%)	0.1	30.34	0.4	155.00
NaYbF <sub>4</sub>	—	—	0.5	193.75

Carbon Films on 300 Mesh Grids Copper (from Agar Scientific). The nanoparticle size distribution was analysed with ImageJ (National Institutes of Health, USA) software considering over 150 nanoparticles for the statistical analysis. XRD analysis was performed using a Bruker X-ray diffractometer equipped with a CuKα source ( $\lambda = 1.54059$  nm) in transmission mode and using a stainless-steel plate. Fourier transform infrared (FTIR) spectra were recorded in the spectral range of 525 and 4000 cm<sup>-1</sup> on a Vertex 70 Fourier transform infrared spectrometer (FT-IR) spectrometer at room temperature. For the sample preparation, nanoparticles were dried and deposited directly on the glass (approximately 1 mg required for analysis). The UC luminescence spectra and lifetime decay were obtained using a FLS1000 Photoluminescence Spectrometer equipped with an external 980 nm CW laser with 0.001 to 2 W power. The UCNPs were dispersed in hexane to the concentration of 10 mg mL<sup>-1</sup>. All samples were measured using a quartz cuvette with a path length of 10 mm. Best fit for the life curve decay was obtained using the Advanced Analysis of Fluorescence Kinetics software FAST.

## Results and discussion

### Effect of epitaxial shelling on NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) UCNPs

Fig. 1a shows the transmission electron microscopy (TEM) image of the synthesized NaYbF<sub>4</sub>:Tm(0.5%) Fe(5%) core UCNPs showing uniformity in size with a mean diameter of 11.5 nm ± 3.6 nm. Here, we used 5 mol% of Fe<sup>3+</sup> ions which, according to



**Fig. 1** TEM images of core and core@shell structures using different Yb<sup>3+</sup> concentrations during shell growth. (a) NaYbF<sub>4</sub>:Tm(0.5%) Fe(5%) core and (b)–(g) NaYbF<sub>4</sub>:Tm(0.5%) Fe(5%)@NaYF<sub>4</sub>:Yb(x%)-core@shell where x = 0%, 20%, 40%, 60%, 80%, and 100%. Size histograms are shown in Fig. S1 (ESI†). All scale bars show 100 nm.





previous reports, is optimal for the UC enhancement in UCNPs.<sup>31</sup> Core@shell structures were achieved through epitaxial growth, in which NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) served as the core template for the shell growth.<sup>38–40</sup> Different lanthanide (Ln) doped shell materials of NaYF<sub>4</sub>:Yb (x%) were used, where *x* refers to the Yb<sup>3+</sup> doping concentration incorporated into the shell material. For this study, six samples were prepared: NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%)@NaYF<sub>4</sub>:Yb(*x*%)–core@shell structures, where *x* was set at 0%, 20%, 40%, 60%, 80% and 100%. Note, where *x* = 0%, the shell is denoted NaYF<sub>4</sub> and where *x* = 100% the shell is referred to as NaYbF<sub>4</sub>.

Overall, it was observed that the deposition of the shell material on the core led to highly unbalanced epitaxy in different crystallographic directions when different shell materials were used. As one can see from Fig. 1(b)–(g) core@shell nanostructure synthesized with a low concentration of Yb<sup>3+</sup> tended to form small rod-shaped structures while a higher concentration of Yb<sup>3+</sup> favoured the larger hexagonal plate nanostructures. This result is expected due to the smaller ionic radius of Yb<sup>3+</sup> vs. Y<sup>3+</sup> and its higher relative electronegativity allowing for a greater binding energy to F<sup>−</sup> ions to occur, which results in the formation of a larger structure.<sup>41</sup> The size and morphology are summarized in Table 2.

Fourier-transform infrared (FTIR) spectroscopy was used to characterize the functional groups on the surface of the nanoparticles (see Fig. S2, ESI†). The spectra show two main vibrations at 2924 cm<sup>−1</sup> and 2853 cm<sup>−1</sup>, which were attributed respectively to the asymmetric and symmetric stretching modes of the methylene (−CH<sub>2</sub>) groups of the OA molecule. The characteristic bands at ~1700 cm<sup>−1</sup> corresponding to the stretching vibration of C=O from OA, were practically absent in all the spectra. Instead, two bands at 1559 cm<sup>−1</sup> and 1459 cm<sup>−1</sup> appeared and were attributed to the asymmetric and symmetric stretching vibrations of the carboxylate group [−COO<sup>−</sup>]<sub>3</sub>Ln<sup>3+</sup> of the NaLnF<sub>4</sub>, respectively.

To obtain information about the crystal structures, we carried out X-ray diffraction (XRD) measurements (see Fig. S3, ESI†). It shows the XRD patterns of the core and the different

core@shell structures. Our experiments revealed that the initial core had a pure cubic phase ( $\alpha$ ) while all core@shell structures had a pure hexagonal phase ( $\beta$ ), without other impurity peaks. In addition, we saw that some peaks, such as the (101) and (201) reflections, were shifted to lower 2 theta angles when NaYF<sub>4</sub> was grown as the shell material and slightly shifted to higher values when NaYbF<sub>4</sub> was used instead (see Fig. S4, ESI†). This is consistent with the Bragg equation  $\lambda = 2d\sin\theta$ ; Y<sup>3+</sup> ions, which have a larger ionic radius compared to Yb<sup>3+</sup> ions, will form unit cells with larger *d*-spacing, and therefore they will need to be compressed to fit around the NaYbF<sub>4</sub> core.<sup>42,43</sup>

It is known that  $\beta$ -NaLnF<sub>4</sub> generally features the (0001) crystal planes for the top and the bottom and six equivalents  $\pm(10-10)/\pm(1-100)/\pm(01-10)$  axial planes, usually simplified as (10-10).<sup>38,44,45</sup> The rate of overall crystal growth, according to general principles, is related to the relative growth rate of the individual crystal planes within the crystal structure.<sup>46,47</sup> The major cause of anisotropy during shell growth is the lattice mismatch between the core and shell material. Another important factor is the competing effect of the selective adsorption of the OA ligand vs. the deposition of the shell material on specific faces of the core.<sup>48–50</sup> Fig. 2 shows the proposed mechanism of the epitaxial growth of the NaLnF<sub>4</sub> (Ln = Y and Yb) shell around the NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) core.

When the core@shell structure was created with 0% and 20% Yb<sup>3+</sup>, it is postulated that OA adsorbed preferentially on the axial planes, thereby slowing the growth in this direction. Consequently, the shell grew preferentially on the [0001] lateral direction, and thus, produced structures with a rod-like shape and high AR. By increasing the Yb<sup>3+</sup> content to 40%, the shell material could adhere to the (0001) facets as well as to the less preferred (10-10) axial facets, and as a result, the growth rates would be comparable along both directions. Here, the ratio between the length and width was very close to 1, more precisely AR = 0.86. When the concentration of Yb<sup>3+</sup> was above 60%, we reasoned that the material would adhere selectively more on the prismatic axial facets of the core. This way the shell

**Table 2** Composition, shape, sizes, and aspect ratios (ARs) of the resulting NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) @ NaYF<sub>4</sub>:Yb – core@shell nanostructures. Sizes were determined through TEM measurement of >100 UCNPs for the statistical analysis (see Fig. S1, ESI)

Shell material	Core@shell length ( <i>L</i> ) (nm)	Core@shell width ( <i>W</i> ) (nm)	AR ( <i>L</i> / <i>W</i> )	Shape
NaYF <sub>4</sub>	34.3 ± 1.7	26 ± 1	1.32	
NaYF <sub>4</sub> :Yb(20%)	43.6 ± 1.7	35.8 ± 2.8	1.22	
NaYF <sub>4</sub> :Yb(40%)	52.4 ± 1.9	60.7 ± 3.7	0.86	
NaYF <sub>4</sub> :Yb(60%)	69.1 ± 5.7	115.2 ± 6.8	0.60	
NaYF <sub>4</sub> :Yb(80%)	103.0 ± 8.1	200.7 ± 11.8	0.51	
NaYbF <sub>4</sub>	119.1 ± 9.7	255.9 ± 9.6	0.47	



**Fig. 2** Proposed mechanism of epitaxial growth of NaLnF<sub>4</sub> shells around the NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) core.



grew faster in the axial direction, leading to the formation of hexagonal shaped core@shell. When the core@shell structure was created with 80% and 100%  $\text{Yb}^{3+}$ , the shell grew preferentially on the (10–10) axial facets, and thus, produced mostly flattened hexagonal plate-shaped structures and low AR (0.51 and 0.47). These arguments can explain the change of morphology we observed in our experiments.

### Upconversion luminescence properties of nanoparticles with different core@shell structures as a function of $\text{Yb}^{3+}$ content in the shell

To systematically correlate the optical properties to the structure of both the core and the different core@shell structures, the UC luminescence properties were investigated under 980 nm irradiation.

One can observe from the luminescent spectra several characteristic enhanced UC emissions, where the UV emissions at 344 nm and 362 nm can be assigned to the  $^1\text{I}_6 \rightarrow ^3\text{F}_4$  and  $^1\text{D}_2 \rightarrow ^3\text{H}_6$  transitions, and the blue emissions at 450 and 475 nm correspond to the  $^1\text{D}_2 \rightarrow ^3\text{F}_4$  and the  $^1\text{G}_4 \rightarrow ^3\text{H}_6$  transitions of the  $\text{Tm}^{3+}$  ion, respectively (see Fig. S5, ESI† for the full spectra).

As a reminder to the reader, all the samples were prepared from the same core material used throughout, *i.e.* with 94.5%  $\text{Yb}^{3+}$  as the sensitizer and 0.5%  $\text{Tm}^{3+}$  as the activator. Therefore, the only synthetic modification was the incorporation of  $\text{Y}^{3+}$  and  $\text{Yb}^{3+}$ , in different concentrations for the core@shell structures. With the core-only structures, nanoparticles had a high surface area to volume ( $S/V$ ) ratio, therefore, most of the  $\text{Ln}^{3+}$  dopants were exposed to surface deactivations arising from surface defects (*i.e.*, disorder, vacancy, and interstitial defects).<sup>20</sup> Overall, it was observed that upon introduction of any shell material, the recovery from the core's surface defects enhanced the UC luminescence. Compared to the UC emission of the core, the emission from the core@shell prepared with 100% of  $\text{Y}^{3+}$  ions was enhanced by  $\sim 4$  fold and  $\sim 2.6$  fold in the UV and blue range, respectively, as can be seen in Fig. 3a. As the same core material is used throughout, we can assume the same distribution and concentration of  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  ions in both the core and core@ $\text{NaYF}_4$  shell structures. As such, we attribute the change in the emission intensity to the surface quenching effect, which depends primarily on the particle size. When the shell was created using 100% of  $\text{Yb}^{3+}$  ions as the material, the UC emission was enhanced by  $\sim 75$  and  $\sim 55$ -fold with respect to the core-only structure, shown in Fig. 3b. In general, for a thin surface coating in the 0–2.5 nm range, a linear dependence of UC emission intensity and shell thickness has been observed.<sup>51</sup> However, it has been demonstrated that a steady enhancement in UC emission is generally unexpected by using very thick shells and the UC emission intensity reaches a plateau at a shell thickness of  $\sim 3$  nm.<sup>40,52</sup> Based on this finding, we propose that the variation in the UC efficiency could depend on the final size but also on the different pathways of UC in the new geometry for the  $\text{NaYbF}_4$ @- $\text{NaYbF}_4$  core@ shell.

As known, when UCNPs are excited at 980 nm there is an absorption by sensitizers  $\text{Yb}^{3+}$ , producing an efficient energy transfer to the activator  $\text{Tm}^{3+}$  ions, which occurs because each

$\text{Tm}^{3+}$  ion is surrounded by many excited  $\text{Yb}^{3+}$  ions.<sup>53</sup> In general, the energy transfer depends on the strength of the ion–ion interaction and on the concentration of  $\text{Yb}^{3+}$ – $\text{Tm}^{3+}$  and  $\text{Tm}^{3+}$ – $\text{Tm}^{3+}$  ions within the crystal.<sup>54</sup> Therefore, as the relative concentration of  $\text{Yb}^{3+}$  and  $\text{Tm}^{3+}$  is expected to be constant in the core of all samples, we can assume that the gradual incorporation of  $\text{Yb}^{3+}$  ions and their new arrangement within the shell crystal structure will favour this energy transfer process, to a greater or lesser extent.

When  $\text{Yb}^{3+}$  was incorporated into the shell material, as shown in Fig. 3c (Scheme II and III), we can assume that the variation in particle shapes leads to a different pathway for energy transfer for the pair  $\text{Yb}^{3+}$ – $\text{Tm}^{3+}$  in the crystal structure. The  $\text{Yb}^{3+}$  sensitizer located at the shell would act as a bridge and allow for energy migration between  $\text{Yb}^{3+}$  ions across the shell, resulting in the excitation of core  $\text{Tm}^{3+}$  ions that are ultimately responsible for emitting UC light. This is more prominent where the shell is composed of 100%  $\text{Yb}^{3+}$  (Fig. 3c (Scheme III)) *vs.* shells containing  $\text{Y}^{3+}$  (Fig. 3c (Scheme II)). The rod like structures that form when  $\text{Y}^{3+}$  is present in the core encourage more 2D energy transfer networks through the length of the structure, while the hexagonal plate like nanostructures that form upon less to no  $\text{Y}^{3+}$  present in the shell allow for greater opportunities for 3D energy transfer networks across the overall crystal structure. Emission enhancement is further demonstrated when looking at the integrated intensities, which shows the enhancement of UC luminescence as a function of  $\text{Yb}^{3+}$  concentration incorporated in the shell growth (see Fig. S6, ESI†).

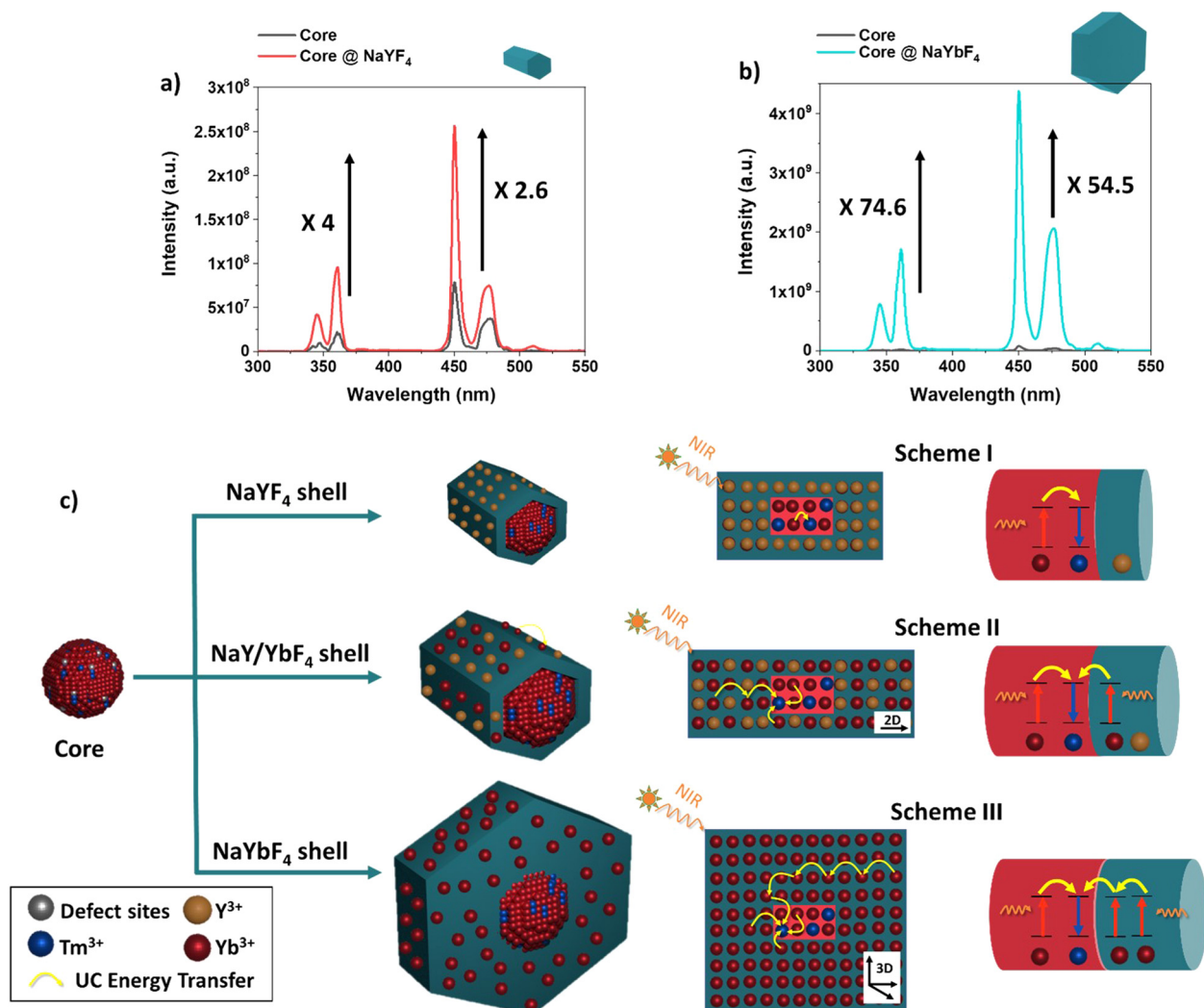
The pump power dependence of the UV and blue emission for the samples was also studied, with the results displayed on a double logarithmic plot (see Fig. S7, ESI†). It is known that the relationship between the UC emission intensity ( $I_{\text{UC}}$ ) and laser power ( $P$ ) can be approximately expressed as:  $I_{\text{UC}} \sim P^n$ , where ( $n$ ) refers to the number of photons absorbed per photon emitted, and its value can be derived from the slope of the fitted line of the plot  $\log(I_{\text{UC}})$  *vs.*  $\log(P)$ .<sup>55</sup> As expected, the UC emission changed gradually when increasing the pump power. The values of the exponent  $n$  were extracted from a linear fit of the experimental data. In detail, the values of the parameter  $n$  (photons absorbed per photon emitted) for 345 nm, 362 nm, 450 nm and 475 nm were found to be 2.00, 1.85, 1.94 and 1.69 for the core, 2.18, 1.95, 1.88 and 1.98 for the core@shell with 0%  $\text{Yb}^{3+}$ , and 2.21, 1.94, 1.901 and 1.85 for the core@shell with 100%  $\text{Yb}^{3+}$ . Overall, the  $n$ -values of the 4 bands are much lower than the theoretical value ( $n = 5$  for 345 nm and  $n = 4$  for 362 nm, 450 nm, and 475 nm) that could be strongly ascribed to the saturation effect, which has been seen previously by other groups.<sup>56–60</sup>

### Lifetime measurements

The upconversion luminescent (UCL) decay behaviours for the pair  $\text{Yb}^{3+}$ – $\text{Tm}^{3+}$  were subsequently studied. As shown in (Fig. S8(a)–(c), ESI†) it was found that the emission of  $\text{Tm}^{3+}$  ( $^1\text{D}_2 \rightarrow ^3\text{H}_6$  transition 362 nm) had a decay lifetime of  $\tau = 49 \mu\text{s}$  for the core, while the decay of  $\text{Tm}^{3+}$  emission when coated with a layer of  $\text{NaYF}_4$  and  $\text{NaYbF}_4$  registered a lifetime of 198  $\mu\text{s}$  and 234  $\mu\text{s}$ , respectively.

When the content of  $\text{Yb}^{3+}$  in the shell increased from 0% to 100%, the emission lifetime of  $\text{Tm}^{3+}$  prolonged noticeably by 36





**Fig. 3** Comparison of UC emission spectra of (a) NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) core vs. NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%)@NaYF<sub>4</sub>-core shell, and (b) NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%) core vs. NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%)@NaYbF<sub>4</sub>-core shell. For the measurements, nanoparticles were dispersed in hexane and measured with a 980 nm CW laser excitation. (c) Schematic illustration of the UC luminescence mechanism for the core@NaYF<sub>4</sub>, core@NaY/YbF<sub>4</sub> and core@NaYbF<sub>4</sub> structures, when measured with a 980 nm CW laser excitation. Scheme I represent a core composed of a sensitizer and activator coated with an optical inert outer shell. Scheme II illustrates the mechanism when the core is coated with a partially active shell of Y<sup>3+</sup> and Yb<sup>3+</sup>, demonstrating the suggested 2D energy transfer network along the length of the rod like nanostructure. Scheme III illustrates the mechanism when the core is coated with an active shell of 100 Yb<sup>3+</sup>, demonstrating the suggested 3D energy transfer network throughout the entirety of the hexagonal platelike structure.

$\mu\text{s}$ , indicating that the Yb<sup>3+</sup> concentration is a crucial factor affecting the energy transfer from the Yb<sup>3+</sup> to Tm<sup>3+</sup>. As a result, the longer lifetime at the 362 nm emission peak observed in the core@shell structured UCNPs indicates the lower non-radiative relaxation rate which agrees with the result shown in Fig. 3. For the  $^1\text{D}_2 \rightarrow ^3\text{F}_4$  transition 450 nm (see Fig. S8(d)–(f), ESI<sup>†</sup>) a similar trend was observed; a decay lifetime of  $\tau = 60 \mu\text{s}$  for the core, 201  $\mu\text{s}$  when coated with a NaYF<sub>4</sub> and 205  $\mu\text{s}$  when coated with NaYbF<sub>4</sub>.

## Conclusion

In this study we have developed novel NaYbF<sub>4</sub>:Tm(0.5%)Fe(5%)@Na(Y/Yb)F<sub>4</sub> core@shell UCNP structures, possessing

varied Y<sup>3+</sup>/Yb<sup>3+</sup> concentration ratios. We have demonstrated that the incorporation of Yb<sup>3+</sup> ions into the shell material further enhanced the emission in both the UV and visible regions, while also offering the ability to incrementally tune particle size and morphology. Based on our XRD observations and aligned with known theory, we propose a shift in the shell growth mechanism from a preferential growth in the [0001] lateral direction when low Yb<sup>3+</sup> is present in the shell material, producing structures with a rod-like shape, to a preferential growth on the (10–10) axial facets when higher Yb<sup>3+</sup> is present in the shell material, producing flattened hexagonal plate-shaped structures. An increase in the emission intensity was also observed at higher Yb<sup>3+</sup> concentrations in the shell material. We hypothesize that the enhancement in the UC efficiency is attributed to a combination of an increase in particle size,



but also variation of the energy transfer pathways arising through geometric alteration.

We propose that a reduction in the  $Y^{3+}/Yb^{3+}$  ratio encourages a 3D energy transfer network throughout the entirety of the plate-like structures, providing greater opportunities for energy transfer where the  $Yb^{3+}$  concentration is high within the shell. In contrast, at higher  $Y^{3+}/Yb^{3+}$  ratios fewer opportunities for energy transfer are present, resulting in more 2D energy transfer networks along the length of the formed rod-like structures where the  $Y^{3+}$  concentration is high. These findings are useful in designing core@shell nanostructures, which could be exploited further in creating accurate models of core@shell designs expanded to other cores, highlighting new findings for multifunctionality and optimal performance in luminescent applications.

## Author contributions

EUH and MG designed the experiments and wrote the manuscript. EUH and KL conducted all laboratory work, including the synthesis, characterisation and data interpretation.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

The authors would like to acknowledge funding from the UKRI Future Leaders Fellowship (MR/T021306/1), from the EPSRC for a vacation bursary (KL). The authors also thank the UoL Materials Innovation Factory (MIF) for access to laboratories, the UoL Biomedical Electron Microscopy Unit (EM Unit) for access to TEM and Amy Lunt of the Chong group (UoL) for access to analytical equipment (XRD).

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