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ARTICLE

Effect of strand molecular length on mechanochemical transduction in elastomers probed with uniform force sensors

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The mechanical properties of a polymer network reflect the collective behavior of all of the constituent strands within the network. These strands comprise a distribution of states, and a central question is how the deformation and tension experienced by a strand is influenced by strand length. Here, we address this question through the use of mechanophore force probes with discrete molecular weights. Probe strands, each bearing a mechanochromic spiropyran (SP), were prepared through an iterative synthetic strategy, providing uniform PDMS-functionalized SP force probes with molecular weights of 578, 1170, and 2356 g/mol. The probes were each doped (9 mM) into the same silicone elastomer matrix. Upon stretching, the materials change color, consistent with the expected conversion of SP to merocyanine (MC). The critical strain at which measurable mechanochromism is observed is correlated with the strain hardening of the matrix, but it is independent of the molecular length of the probe strand. When a network with activated strands is relaxed, the color dissipates, and the rate of decoloration varies as a function of the relaxing strain (ϵ_r); faster decoloration occurs at lower ϵ_r . The dependence of decoloration rate on ϵ_r is taken to reflect the effect of residual tension in the once-activated strands on the reversion reaction of MC to SP, and the effect of that residual tension is indistinguishable across the three molecular lengths examined. The combination of discrete strand synthesis and mechanochromism provides a foundation to further test and develop molecular-based theories of elasticity and fracture in polymer networks.

Introduction

The collective behavior of strands dictates the mechanical properties of a polymer network.^{1–4} Thanks in part to the advent of single-molecule force spectroscopy/AFM,^{5–8} the mechanical behaviors of single molecules have been experimentally and theoretically well-characterized. Often, the various influences of the specific characteristics of a given strand (e.g., composition, length, orientation relative to applied strain, connectivity) within the collective, however, are hidden within the behavior of the ensemble. Elucidating the contributions of such molecular microstates will test and inform the evolution of theories of polymer elasticity and fracture, which might further lead to molecularly optimized mechanical properties. In addition, a better understanding of molecular structure-activity relationships in this area is likely to benefit emerging approaches to mechanically adaptive systems, especially in the field of mechanochemistry, where mechanophores induce changes in material properties such as stress-induced color

changes,^{9–15} stress-strengthening,^{16–19} triggered degradation,^{20–23} and release of chemical cargo.^{24–29} For example, the elastically active strands within a polymer network generally comprise a range of molecular lengths between cross-links or entanglements (M_x , typically on the order of 10^2 to 10^3 g/mol), and an intuitive expectation is that shorter strands should serve as sites of focused tension, as a result of reaching their finite extensibility at smaller strains. The relationship between strand molecular length and mechanical response within a given network, therefore, represents an important factor that might guide the molecular design of mechanically active networks.³⁰

Mechanophores represent powerful and promising tools for investigating this structure-activity relationship.^{9,31,32} Numerous examples of using mechanophores to probe fundamental questions of polymer network physical behavior have been reported. For example, Moore, Sottos, and co-workers applied spiropyran-based mechanophores in polyacrylate-based^{11,33–35} and polyurethane-based^{36,37} networks to extensively investigate how the force transduction from macroscopic stress to the molecular tension felt by mechanophore probes is affected by various structural parameters, such as the positions of the mechanophore probes inside a polymer/filler composite^{38–42} (i.e., matrix or interface) and orientation of chains.³⁶ For rubbery networks, Creton and co-workers incorporated 1,2-dioxetane-based mechanophore probes¹² into tough multi-networks to visualize where molecular chain scission occurs and the number of bonds that are broken during fracture.⁴³ Clough et al. used similar 1,2-

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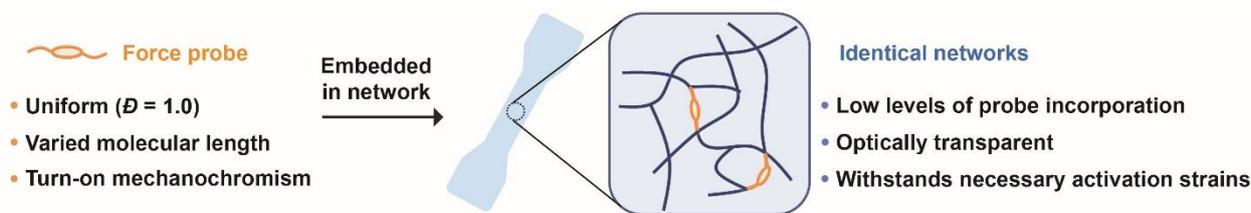


Fig. 1 The approach to analyze molecular behavior inside a polymer network

dioxetane-based mechanophores to elucidate the contribution of chain scission to the Mullins effect in a silicone elastomer.⁴⁴ Recently, Lin et al. incorporated azobenzene-based mechanophores into a silicone elastomer and experimentally determined the average force that the azobenzene probe experiences as a function of the macroscopic strain of the network.⁴⁵ Building off of these and related advances, we sought to take advantage of the ability to probe strand state within a network to address a central question: to what extent does the tension felt by a strand with a rubbery, elastomeric network depend on the length of that particular strand, vs. the continuum behavior of the network to which it is tethered?

Our broad approach is summarized in Fig. 1. We envisioned mechanochromic probe strands of similar composition but different molecular lengths incorporated at low levels within the same polymer network continuum. By maintaining the same mechanophore reporter in all of the probe strands, any observed differences in probe strand response from one network to another could be attributed to the differences in the molecular length of the probe. We established the following criteria for an initial investigation: (i) a distinct, strain-dependent signal that can be detected at a level of probe strand incorporation that is too low to have a significant effect on network mechanical properties and topology; (ii) miscibility of the probe strand and other starting materials for network synthesis, to avoid challenges associated with the aggregation of mechanochromic molecules; (iii) truly uniform probe strands across a relevant range of strand length ($M_x = 10^2$ - 10^3 g/mol, $D = 1.0$). The first criterion is satisfied through the use of the well-established mechanochromic force probe spiropyran, but the latter two criteria required synthetic method development.

The synthetic challenge began with the desire for uniform probes. A distribution of molecular weights is a necessary consequence of most abiological polymer synthesis methods, arising from the intrinsically stochastic nature of conventional step growth and addition polymerizations. This challenge has been creatively addressed through the development of solid-phase synthesis,⁴⁶⁻⁴⁸ iterative exponential growth,⁴⁹⁻⁵² and related approaches to discrete polymers.⁵³⁻⁵⁶ To the best of our knowledge, however, these methods have yet to be applied to the synthesis of a series of mechanophore-incorporated strands of precise molecular weight. We therefore set out to synthesize discrete spiropyran-based probe strands that satisfy the above criteria. We chose the poly(dimethyl siloxane) (PDMS) elastomer Sylgard 184 as a testing matrix, since it is an industrially important and widely used elastomer that has also proven to be a robust platform for polymer mechanochemistry,^{10,13,57} including mechanochromism. We

then used these SP-doped polymer networks to analyze the effects of strand molecular length on the mechanochemical reactivities and the tension experienced by the strand inside a common elastomeric matrix.

Results and Discussion

Bulk network effects and mechanochemical reactivities

Because we ultimately must compare networks that are made separately, we first evaluated how small variations in curing might influence molecular force transduction. For this purpose, we used commercially available silicone Sylgard 184 as the network matrix. Whereas many common polymer networks (e.g., butyl acrylate networks made via free radical polymerization) have difficulty in achieving quantifiable levels of mechanophore activation,⁹ this widely used silicone elastomer is a proven platform for mechanophore activation and quantification due to its combination of optical transparency and mechanical robustness.^{10,13,57,58} We chose **1c** (533 g/mol) for the force probe due to its proven activation inside the silicone elastomer.¹³ Elastomer films with **1c** (9 mM) were obtained through a platinum-catalyzed hydrosilylation reaction (Fig. 2), and differences in cured network structure and mechanical properties were created by intentional variation of the base to curing agent ratio of the two-component Sylgard kit (base:curing agent from 8:1 to 12:1, see Table S1). There is no SP aggregation observed in the X-ray scattering measurements (See SAXS, MAXS, and WAXS data in Fig. S2). We note that: (i) the filler content of each elastomer varies slightly as the base to curing agent ratio changes; and, (ii) it is possible that SP inhibits the Pt catalyst, lowering local crosslinking density around SP. Successful data normalization (see below) suggests that contributions from these effects are minimal, but they cannot be ruled out entirely. The films are mechanically robust, and at large strains they exhibit the expected color change to blue that is associated with the mechanically coupled conversion of SP to MC.¹³

The mechanochemical response of the probe was quantified as a function of uniaxial tension (Fig. 3a). As expected, the initial modulus and onset strain of strain hardening (ϵ_{SH} , defined as the intersection of linear fits to the stress/strain curve before and after the nonlinear transition in the curve) vary slightly across replicates of the same 10:1 mixing ratio, and they vary even more when the component ratio is altered intentionally, with the 8:1 mixture being stiffer (with lower ϵ_{SH}) and 12:1 being more compliant (higher ϵ_{SH}). Simultaneously, the onset of mechanochromism (ϵ_c) was obtained from in-situ digital image color (RGB) analysis (Fig. 3a). There is a distribution of initial

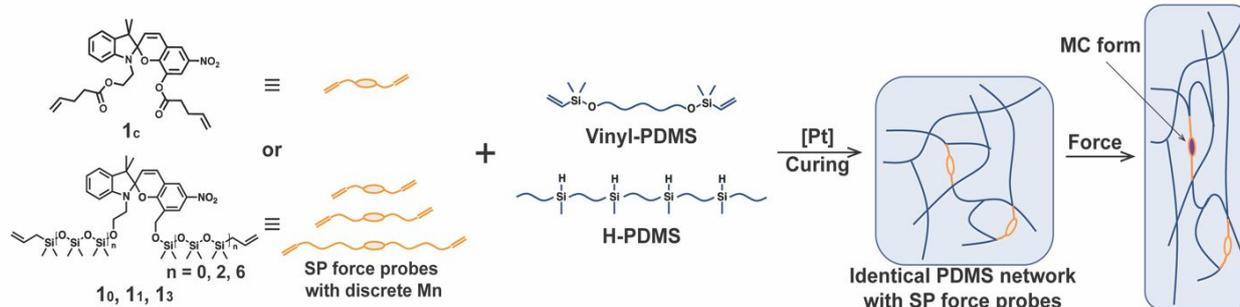


Fig. 2 Schematic of a PDMS elastomer incorporating SP probes.

molecular orientations and conformations of the force probes, as indicated in previous studies, which means that even though all mechanophores are themselves identical, some mechanophores activate at lower strains than others because the local tension is not uniform.^{59–61} In this study, we consider the critical strain for activation, ε_c , which we define as the intersection of linear fits to the RGB ratio curve before and after the nonlinear transition in the curve and represents the strain where major activation events start. The mechanochromism also changes from one film to another, but it tracks with the shifts in ε_{SH} (Fig. 3a), a correlation that is borne out by plotting color as a function of normalized strain, $\bar{\varepsilon} = \varepsilon/\varepsilon_{SH}$. As shown in Fig. 3b, the mechanochromism across the five samples collapses to a common onset strain of approximately $\bar{\varepsilon} = 1.1$. The length of the probe stays the same, but the onset strain for mechanochromism varies, and that variation is correlated with the mechanical behavior of the surrounding network. The surrounding polymer network, and in particular its strain hardening behavior, is one of the determining factors of force transduction.

Local strand effects and mechanochemical activation

We next sought to isolate the effects of strand length on force transduction. For a randomly coiled single macromolecule, the initial end-to-end distance (R_i) is proportional to the square root of the number of Kuhn lengths along its backbone ($N^{1/2}$), while the maximum length (R_{max}) is directly proportional to N . As such, the maximum extension to which an individual chain can be stretched scales as $\lambda_t = R_{max}/R_i = N^{1/2}$.⁶² If individual strands inside a polymer network deform affinely, strand length becomes an important factor in determining the onset of mechanochromism.

To analyze the effects of strand molecular length, we began by synthesizing discrete force probes that satisfy our criteria (mechanochromic, uniform, miscible) in the context of silicone elastomers. We adopted a strategy based on the iterative synthesis⁵³ of uniform poly(dimethyl siloxane)⁵¹ from an SP core (Fig. 4a). The SP core was coupled to an α,ω -chlorohydroxy oligosiloxane chain extender through nucleophilic substitution of an SP hydroxyl with the chloride on the chain extender. Catalytic oxidation of the oligosilane hydride regenerated terminal hydroxyl groups for subsequent extension on both sides through a second and/or third cycle. This pathway allows for siloxane extension and is expected to help improve the

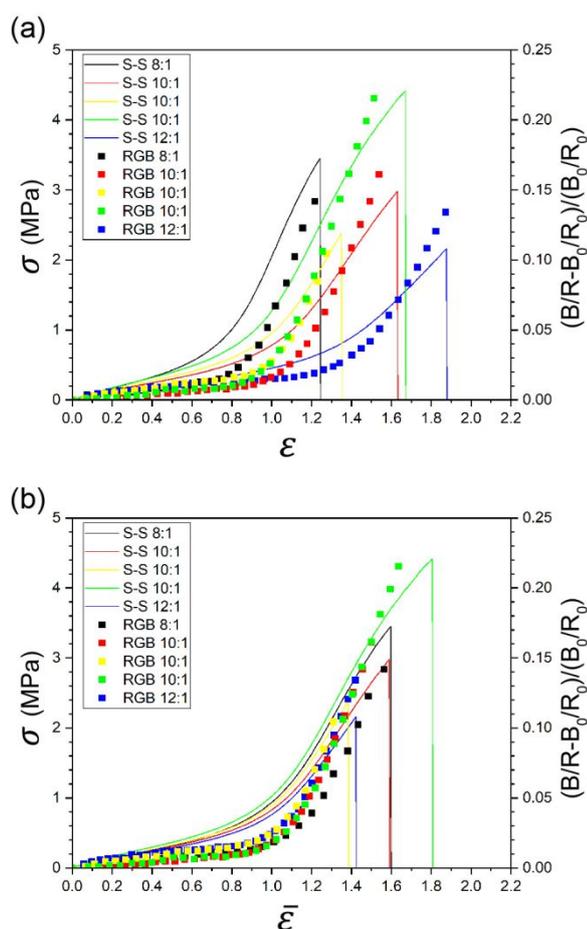
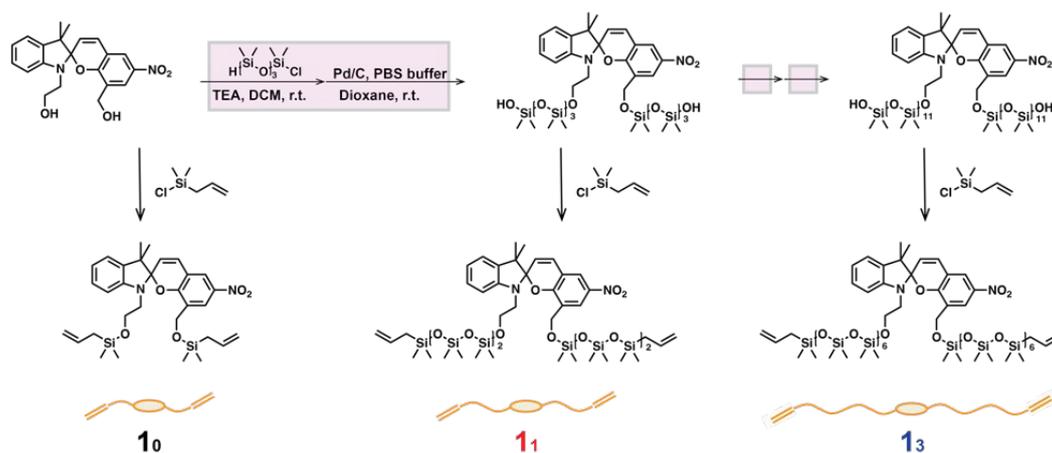


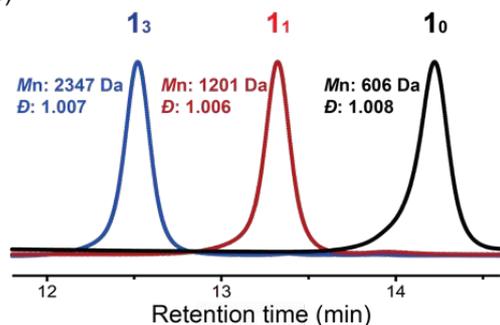
Fig. 3 Stress-Strain curves and RGB ratio analysis of PDMS with different base and curing ratios: (a) before normalization and (b) after normalization. The ratios indicate base to curing agent ratios, and the three 10:1 samples are different batches.

miscibility of the probe strands and the silicone matrix. As this process was iterated with full control over the end-group functionality, the molecular weight grew linearly in a defined manner. Following zero, one, or three iterations, the strands were capped at the hydroxyl groups with allyl groups to give probe strands **1₀**, **1₁**, and **1₃**, respectively, where the subscript *n* refers to the number of extension iterations involved in the synthesis of **1_n**. Each of the three PDMS-based spiropyran (SP) force probe strands possess a single SP per chain, and the probes span discrete molecular weights of 578, 1170, and 2356 g/mol ($\bar{D} = 1.0$) for **1₀**, **1₂**, and **1₃**, respectively.

(a)



(b)



(c)

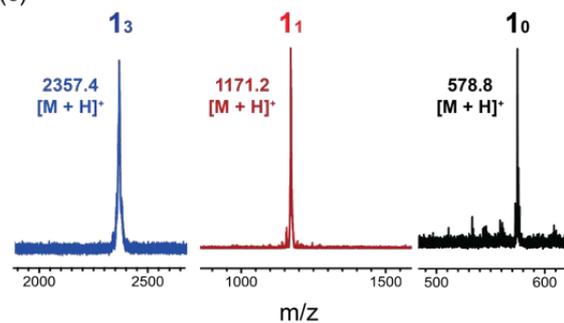


Fig. 4 (a) synthetic scheme for the SP probe strands with discrete molecular weights, (b) chloroform SEC traces of the SP probe strands, and (c) MALDI-TOF of the SP probe strands

The probe strands were characterized by NMR spectroscopy (Fig. S3-8), size exclusion chromatography (SEC, Fig. 4b), and MALDI-TOF mass spectrometry (Fig. 4c). ^1H and ^{13}C NMR show the signature peaks of SP in each probe strand, and peak integration in ^1H NMR confirms the incorporation of one SP mechanophore per strand (Fig. S3-8). The SEC traces show sharp peaks with decreased retention times as the number of synthetic iterations increases, indicating the corresponding increase in molecular weight. Furthermore, MALDI-TOF mass spectra are dominated by a single peak at the expected mass/charge ratio of the strands, thus supporting their uniformity. We estimate the number of Kuhn lengths in the three strands to be $N = 1.1$, 2.6 , and 5.7 for **1₀**, **1₂**, and **1₃**, respectively (see Section G in ESI).⁶²

We incorporated these discrete SP sensors into Sylgard 184 silicone elastomers (average molecular weight of 3,000 to 3,500

g/mol per elastically active strand with the functionality of **4**, corresponding to the number of Kuhn units, $N_{\text{ave}} \sim 8 - 9$) as shown in Fig. 2. Each of the SP probes **1_n** were incorporated into the silicone matrix under identical conditions (9 mM probe, 10:1 base to curing agent) to produce elastomeric films through a platinum-catalyzed hydrosilylation. Here, the probe strands are miscible with the Sylgard kit, a feature we attribute to the use of siloxane extenders in the iterative synthetic approach. There is no SP aggregation observed in the X-ray scattering measurements (See Fig. S2 for SAXS, MAXS and WAXS data). Indeed, when polar triazoles were used to link the SP to PDMS

handles, the probe strand was immiscible with the Sylgard matrix and good films were not obtained (results not shown).

The mechanical properties of the various films under uniaxial tension are consistent with well-formed silicone networks (Fig. 5a). The Young's moduli (E), ~ 1.0 - 1.2 MPa of the films with the various replicates fall within the range previously observed for replicates of the network doped with **1c**. There might be some differences in the number of effective probe strands due to the different propensities of the three probes to form elastically inactive loops.⁶³ The presence of loops, however, only changes the magnitude of the mechanochromism and not the strain-dependent behavior examined here. Additionally, there is no evidence that the chain length of the SP probes has a significant effect on the mechanical properties of the network at this level of incorporation. When stretched, all films develop a blue color

that is clearly visible to the naked eye in the stretched regions, demonstrating successful mechanophore activation of SP to MC as shown previously.¹³

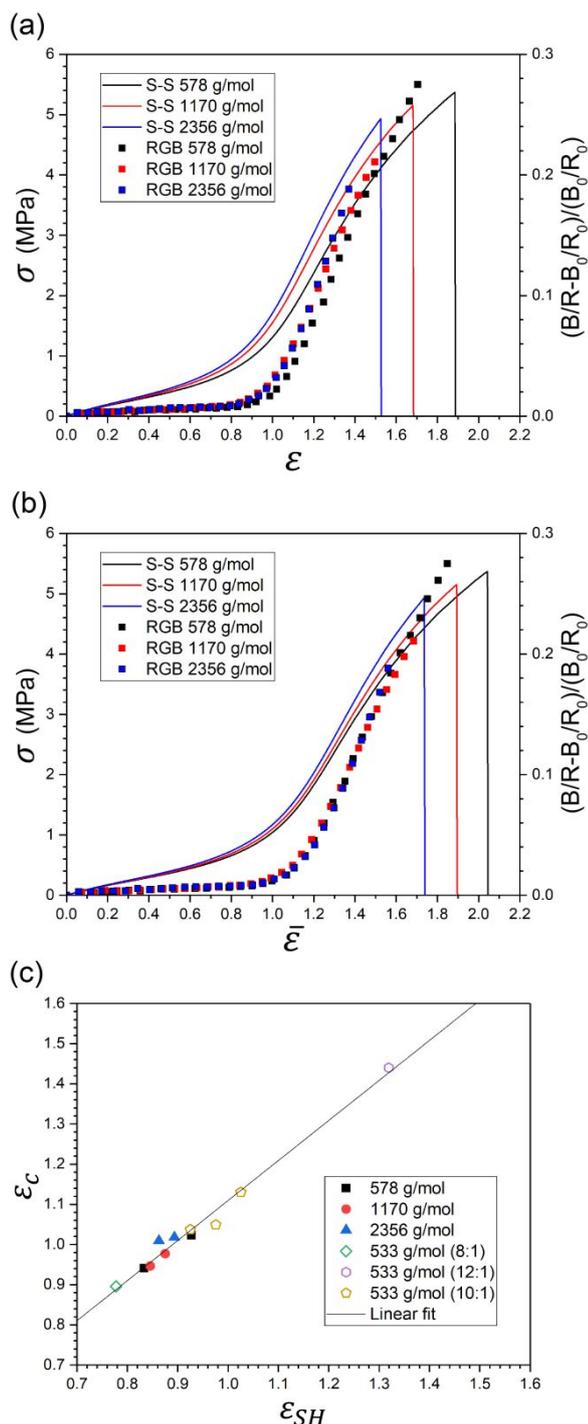


Fig. 5 Stress-Strain curves and RGB ratio analysis: (a) before normalization and (b) after normalization, and (c) the onset strains for mechanochromism against those for strain hardening. R^2 of the linear fit is 0.98.

The mechanochromism of each film was tracked (Fig. 5a), and it again correlates with strain hardening behavior (Fig. 5b and c) in a manner that is indistinguishable from that observed in Fig. 3. A correlation between macroscopic strain hardening and mechanochromism has been observed in previous studies using SP probes.^{10,59,60} Chen et al. fabricated multiple networks with SP probes in the first networks that bear the majority of the stress.^{59,60} Using a fixed probe, they changed crosslinking density (i.e., strand length between crosslinkers) and observed that the onset of mechanochromism correlated with the onset of strain hardening as the latter shifted in response to the change in crosslinking density. Similarly, Lin et al. incorporated SP probes as cross-linkers into silicone elastomers, maintaining a constant cross-linker length but changing the critical force for activation by varying the regiochemistry of the cross-linking attachments on the SP.¹⁰ More reactive/better coupled SP probes resulted in greater mechanochromism; the onset of the mechanochromism, however, remained fixed and located at the onset of strain hardening.

The general correlation of strain hardening with mechanochromism makes sense, as the >100 pN forces required for SP activation⁶⁴ at the single strand level occur in the regime that is the single-molecule equivalent of strain hardening, where conformational degrees of freedom have been exhausted to the extent that force starts to increase more dramatically as a function of chain extension. Here, we ask the question as to how varying the critical extension of the probe strand (by changing its contour length) impacts the delivery of tension to the mechanophore embedded within it. A strong strand length dependency in the onset of the mechanochemical reaction (ϵ_c) would be observed if the deformation of individual probe strands were affine, as would be expected from single molecule extension behavior. However, there is no obvious strand length dependency. The experimentally observed ϵ_c are 1.02, 0.98, and 1.01 for PDMS with $\mathbf{1}_0$, $\mathbf{1}_1$, and $\mathbf{1}_3$, respectively, whereas an $N^{1/2}$ dependence on ϵ_c would result in activation strains of ~ 2.2 for $\mathbf{1}_1$ and 3.7 for $\mathbf{1}_3$, based on the $\mathbf{1}_0$ benchmark. This approximation using $N^{1/2}$ dependence is not rigorous and probably overestimates ϵ_c to some degree, as the real initial end-to-end distance (R_i) becomes smaller than $N^{1/2}b$ for a shorter chain (b : the length of Kuhn unit).⁶² The key point, however, is that the added flexible chain length in $\mathbf{1}_1$ and, especially, in $\mathbf{1}_3$ is sufficient to change the critical single-molecule extension required for activation in a way that, were it to translate to the bulk, would be experimentally observable. Therefore, the results indicate non-affine deformation inside the elastomers. We speculate that this non-affine behavior is partially attributed to the force balance inside the networks, as concluded previously by Chen et al.^{59,60} Simply speaking, for a given chain orientation, a shorter chain must deform less than a longer chain to generate the same amount of force, which prevents individual strands from affine deformation.

We considered the observed onsets of strain hardening and mechanochromism in the context of simple expectations based on network structure. The limiting extension ratio of a network is related to N_{ave} by $\lambda_{lim} = N_{ave}^{1/2} = \epsilon_{lim} + 1$ and is estimated to be about 2.9 (i.e., $N_{ave} = 8 - 9$). Since λ_{lim} represents the extension

ratio, where stress/force diverge, it is larger than the observed onset extension ratios of the strain hardening ($\lambda \sim 1.9$, $\varepsilon \sim 190$ %strain) and the mechanochromism ($\lambda \sim 2.0$, $\varepsilon \sim 200$ %strain). We fitted the initial portion of each curve in Fig. 5a with an empirical constitutive model by Gent that takes into account strain hardening^{65,66} and extracted a characteristic extension for strain hardening λ_m of 2.3 (fitting details are provided in Section J of the ESI). The λ_m extracted from the experimental data is similar to, but smaller than, the value of $\lambda_{lim} = 2.9$ obtained from N_{ave} . Empirically, strain hardening behavior depends on multiple factors (e.g., strand length, entanglements, presence of phase separation).^{66–69} Some systems have similar λ_{lim} and λ_m ,⁶⁶ but other systems exhibit significant differences between the two values.^{68,69} One possibility for the difference in our case might be contributions from nanofillers in the surrounding matrix.

Strand length and mechanophore deactivation

Finally, we consider the mechanophores that are “turned on” when the strain exceeds ε_c and ask how those strands relax (i.e., how the mechanophores “turn back off” by reverting to SP from MC) as the strain applied to the network is reduced to values below ε_c . Generally, a polar environment favors the more polar MC form, making SP easier to open and more difficult to close;¹⁴ here, the non-polar nature of the silicone has the opposite effect and leads to decoloration on reasonable time scales. Specifically, the decoloration kinetics of the elastomers were analyzed as a function of time and normalized strain ($\bar{\varepsilon} = \varepsilon/\varepsilon_{SH}$) to distinguish bulk network effects from local strand effects. Silicone elastomers were first stretched to $\bar{\varepsilon} = 1.1$ to induce force-activated MC states. After equilibrating for 30 min, the strain was lowered to a normalized relaxation strain (ε_r), and the decoloration of the force-activated MC states back to SP was monitored as a function of the extent of relaxation (i.e., from $\bar{\varepsilon} = 1.1$ to $\varepsilon_r = 1.0 - 0.1$).

To properly account for the effect of residual tension on decoloration rates, we first quantified the rates of MC reversion in unstrained films, by photochemically converting SP to MC and observing the rate at which color fades under negligible tension in the unstretched films, where the local environment of the low- T_g polymer is liquid-like. The longest probe strands fade about 2 times more quickly than the shortest strands, which we attribute the difference to small variations in local polarity that result from the tethers (Table S2). We therefore focus our analysis on the relative decoloration kinetics of a given probe as a function of ε_r .

Two trends emerge from the data. First, lower ε_r corresponds to faster decoloration (Fig. 6); just as the high tensions at large strains accelerate the conversion of SP to MC, the reversion of MC to SP is accelerated at lower strains. This strain-coupled effect is more obvious at ε_r between 1.0 and 0.6. Below $\varepsilon_r = 0.6$, the decoloration is effectively indistinguishable, indicating that the activated MCs experience negligible tension at these low strains. Second, as with the activation studies, the relative effect of tension on MC-to-SP reversion kinetics is indistinguishable across the series of probes. We quantified the strain dependence by normalizing the time axis of the

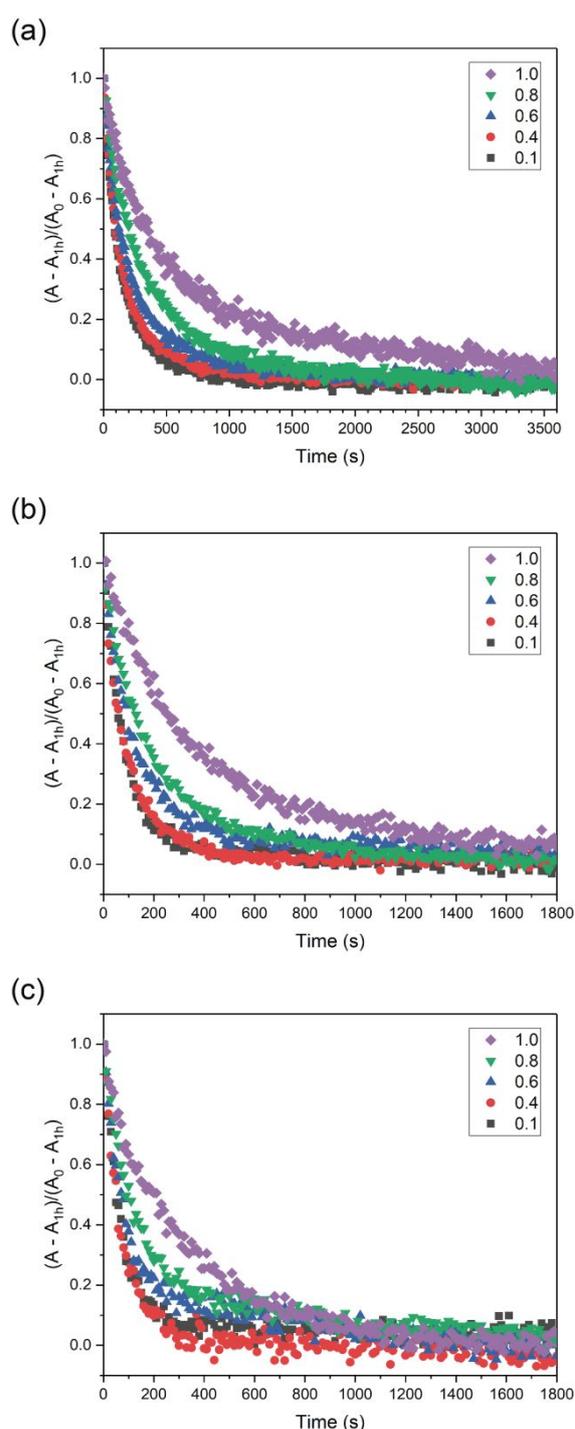


Fig. 6 The relaxation kinetics of force activated MC back to SP states at each ε_r : (a) 578 g/mol, (b) 1170 g/mol, and (c) 2356 g/mol. The absorbance values are normalized as $(A - A_{1h})/(A_0 - A_{1h})$ for comparison.

decoloration at a given strain by the dimensionless scaling factor t_{rel} . Specifically, at each ε_r we choose t_{rel} so that plots of decoloration vs. t/t_{rel} collapse onto the reference curve obtained at $\varepsilon_r = 0.1$ (where, as described below, strain has a negligible impact on the kinetics; details of the scaling are found in Section I of the ESI). The values of t_{rel} therefore signify how much the decoloration is slowed at a given strain relative to

decoloration at negligible strain. This normalization is phenomenological: the decoloration curves are not perfectly described by a single exponential decay, and this shift enables us to capture the cumulative effect across the distribution of relaxation modes. We can then express the change in kinetics as a relative rate constant k_{rel} given by equation (1), where k_{rel} and k_0 are the reaction constants at a given ε_r and at the force free state, respectively:

$$k_{\text{rel}}/k_0 = 1/t_{\text{rel}} \quad (1)$$

The normalized rate constants for decoloration are shown in Fig. 7a, confirming that residual tension affects the MC-to-SP reversion with no discernible influence from the probe strand length.

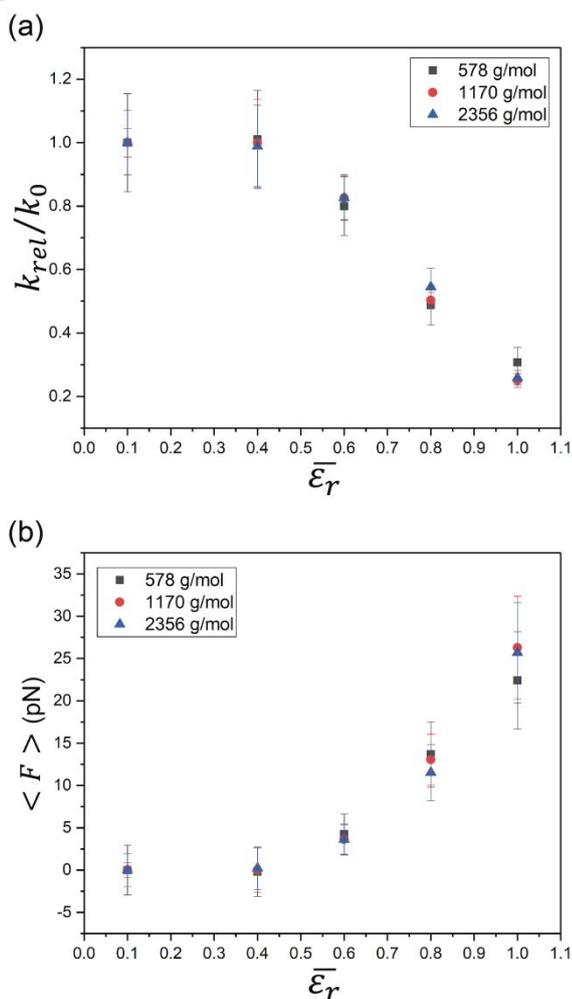


Fig. 7 (a) Ratios of reaction constants and (b) calculated average forces of pre-activated probe strands as a function of relaxation strain and strand molecular weight.

The magnitude of the average tension experienced in the strands can be estimated from the strain-dependent decoloration rates. At the molecular level, the relationship between rate and force of tension is often complex,^{70–72} but can often be estimated by assuming the simplest model of mechanochemical coupling (eq. 2), where k_B , T , and Δx^\ddagger are Boltzmann constant, temperature, and the change in end-to-end distance between MC and transition states, respectively.

$$\langle F \rangle = (k_B T / \Delta x^\ddagger) \ln(k_{\text{rel}}/k_0) \quad (2)$$

Within the constraints of this model, prior literature suggests a value of $\Delta x^\ddagger = -2.13 \text{ \AA}$.⁶⁴ Below $\varepsilon_r = 0.6$, the average residual force calculated from eq. 1 and 2 is negligible, but increases from a few pN at $\varepsilon_r = 0.6$ to ~ 25 pN at $\varepsilon_r = 1.0$ (Fig. 7b). Most importantly, the change in activated strand tension upon relaxing the network, like the strain-induced activation itself, is independent of strand length.

In addition to the primary observation regarding the (lack of) dependence on strand length, it is worth comparing the average tension per strand inferred here to that reported previously by Lin et al. for the cis-to-trans isomerization of a short azobenzene (AB) derivative ($M_n \sim 400$ g/mol, comparable to **1a**) similarly embedded in Sylgard 184. In the AB system, there is no mechanical pre-activation; instead the rate of AB isomerization was fit to a two-state model that treated the AB derivatives as being either under negligible tension or being part of a subset of strands that experience measurable effects of tension akin to those explored here. The characteristic residual forces acting on the SP probes at high normalized strain (e.g., ~ 25 pN at $\varepsilon_r = 1.0$ and $\varepsilon \sim 0.9$) are similar to, but slightly lower than, those inferred previously from the AB kinetics at similar strain (~ 40 pN). Differences in the magnitude of tension are expected, as the populations of strands being probed in the two studies are not identical. For example, the design of the present study does not include the small subset of SP strands under highest tension that do not revert back to MC; no such exclusion is intrinsic to the design in the previous work of Lin et al.

Conclusions

The use of discrete force probes of varied contour length provides insight into the behavior of specific sets of initial microstates within the canonical ensemble of a silicone elastomer. Polymer networks comprise a complex mixture of local strand lengths and topological connections, the details of which might vary greatly from one network to another in ways that influence the transduction of force among strands. The methods employed here using discrete force probes represent what we perceive to be a rich set of opportunities to use the increased precision afforded by contemporary synthetic methods to tease apart the contributions of specific structural microstates to the ensemble behavior of networks. We expect that these types of studies will help inform future molecularly based structure activity relationships and comprehensive physical models of network properties, including descriptions of tension distribution and fracture behavior that have so far been the sole province of simulation.^{73,74}

For the specific components studied here, the observations indicate that the bulk behavior of the network dominates mechanophore response to the exclusion of local strand length. On a fundamental level, the results cannot be interpreted through the lens of affine deformation of an individual strand. Rather, the macroscopic stress-strain curve, and especially the onset of strain-hardening, is the key correlant of molecular tension in overstressed strands. On a practical level, the results

suggest a reasonable tolerance for strand length when embedding mechanophores in a rubbery matrix, as compared to cross-linker length effects observed in glassy polymers.³³ Such freedom in design might be useful in improving the compatibility of mechanophore and matrix, or in addressing the challenges of low mechanophore activation^{9,44,75} by incorporating multiple mechanophores (including multiple types of mechanophores) into a single strand.⁷⁶ To that end, the Sylgard 184 matrix employed here is a commercially important and widely used elastomer. Unlike “conventional” elastomers that have difficulty in activating detectable amounts of mechanophores,⁹ the silicone elastomer has proven to be a robust platform for covalent mechanochemistry.^{10,13,37} As such, the findings should be applicable to other mechanically responsive systems.

Finally, we speculate as to the generality and limits of the present observations. At the level of the polymer matrix, the silicone network in this study is a randomly crosslinked network, with a distribution of strand molecular weight that is similar to that of many common elastomers (e.g., those made through vulcanization). Future work exploring the influence of various network architectures, such as end-linked networks of low-dispersity components (e.g., tetra-PEG gels⁷⁷), is an attractive target, but will require matrices that accommodate sufficient levels of mechanophore activation. At the level of the probe strand, the range of discrete probe strand lengths explored here ($N \sim 1-6$) is sufficient to show that large variations in the degree of individual strand extension does not have an impact of macroscopic strain threshold, but we note that the probe strands employed here are shorter than the average strand in the network ($N_{\text{avg}} \sim 8-9$). We hypothesize that different behavior would be observed if the probe strand is (much) longer than the effective strand. For example, in the extreme, a strand with infinitely long molecular length should feel almost negligible force, even as the network around it is stretched to the point of strain hardening. Due to synthetic feasibility, this regime is not presently accessible to us, but this limit is a technical and not a theoretical one. These and other questions of molecular structure in polymer networks, including the effect of junction functionality and network topology, provide rich opportunities for future investigation.

Data availability

The experimental procedures and supporting data for this article have been included in the ESI.†

Author contributions

T.O., W.W., J.A.J., and S.L.C. conceived the project. T.O., W.W., J.A.J., and S.L.C. designed the experiments. T.O. and B.E.S. performed the syntheses and structural characterization of the SP diol derivatives. W.W. conducted the synthesis and structural characterization of the SP probes with uniform molecular weights. T.O. synthesized PDMS elastomers with the SP probes and characterized their structures. T.O. conducted the tensile tests, image capturing, and reaction

kinetics tests, and analyzed the data. T.O., W.W., J.A.J., and S.L.C. discussed the results. J.A.J. and S.L.C. provided the funding. T.O., W.W., J.A.J., and S.L.C. wrote the manuscript.

‡ T. Ouchi and W. Wang contributed equally to this work.

Conflicts of interest

There are no conflicts to declare.

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