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The preparation and phase diagrams of (⁷Li_{1-x}Fe_xOD)FeSe and (Li_{1-x}Fe_xOH)FeSe superconductors

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We report the phase diagram for the superconducting system (⁷Li_{1-x}Fe_xOD)FeSe and contrast it with that of (Li_{1-x}Fe_xOH)FeSe both in single crystal and powder forms. Samples were prepared via hydrothermal methods and characterized with laboratory and synchrotron X-ray diffraction, high-resolution neutron powder diffraction (NPD), and high intensity NPD. We find a correlation between the tetragonality of the unit cell parameters and the critical temperature, *T_c*, which is indicative of the effects of charge doping on the lattice and formation of iron vacancies in the FeSe layer. We observe no appreciable isotope effect on the maximum *T_c* in substituting H by by D. The NPD measurements definitively rule out an antiferromagnetic ordering in the non-superconducting (Li_{1-x}Fe_xOD)FeSe samples below 120 K, which has been reported in non-superconducting (Li_{1-x}Fe_xOH)FeSe.¹ A likely explanation for the observed antiferromagnetic transition in (Li_{1-x} *x*Fe_xOH)FeSe samples is the formation of impurities during their preparation such as Fe₃O₄ and Li_xFeO₂, which express a charge ordering transition known as the Verwey transition near 120 K. The concentration of these oxide impurities is found to be dependent on the concentration of the lithium hydroxide reagent and the use of H₂O vs. D₂O as the solvent during synthesis. We also describe the reaction conditions that lead to some of our superconducting samples to exhibit ferromagnetism below *T_c*.

Introduction

In the iron-based pnictide and chalcogenide superconductors, chemical doping and physical pressure are universal variables by which to tune the superconducting properties.^{2, 3} For example, the critical temperature, T_{cr} of 8 K in FeSe under ambient conditions⁴⁻⁷ can be raised to 38 K by externally applied pressure^{8, 9} or 44 K by intercalation of cationic species.¹⁰⁻¹³ The tetragonal (*P4/nmm*) structure of FeSe (Figure 1) consists of sheets of edge-sharing FeSe₄ tetrahedra held together by van der Waals interactions, which makes it an ideal host for intercalation chemistry. Negative pressure, or strain, has also been implicated as a parameter in the high T_c of 65 K – 100 K reported for single layered FeSe.

Given the propensity of the FeSe layered system for chemical and physical manipulation, FeSe is an ideal platform for understanding the superconductivity of the iron-based systems and for the preparation of new layered functional materials. The recently discovered ($Li_{1-x}Fe_xOH$)FeSe system,^{1, 17-21} which contains layers of LiOH alternating with layers of FeSe (Figure 1), offers such an opportunity. Iron may occupy the lithium site and therefore effectively charge dope the FeSe layer since ($Li_{1-x}Fe_xOH$) would be positively charged. Sun *et al.* have also reported that increased lithiation of the ($Li_{1-x}Fe_xOH$) layer would force iron to occupy vacancies in the FeSe layer, which can be detrimental to the superconducting properties.¹⁸



Figure 1. Crystal structures of three layered iron selenides: $K_{0.8}Fe_{1.6}Se_2$ (left) FeSe (centre) and ($Li_{1:x}Fe_xOH$)FeSe (right).

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Three outstanding issues in the (Li_{1-x}Fe_xOH)FeSe system and related phases are whether 1) an antiferromagnetic parent phase exists, 2) superconductivity coexists with ferromagnetism, and 3) any isotope effects on T_c exist. Critical to answering all three questions is the preparation of the deuteroxide version of (Li_{1-x}Fe_xOH)FeSe and constructing its phase diagram. Furthermore, hydrothermal synthesis under either H₂O or D₂O presents interesting differences in the purity of the resulting superconducting phases due to differences in the reaction kinetics. Thus, this study will help workers in the field understand the thermodynamic and kinetic factors in the preparation of phase pure and superconducting FeSe-based materials.

Experimental

Sample preparation

The preparation of powder samples were modified from a hydrothermal route reported in the literature.^{1, 22} Due to the high neutron absorption cross section of ⁶Li, our compounds were prepared doubly isotopically pure to form (⁷Li_{1-x}Fe_xOD)FeSe. For the synthesis of deuterated samples we first prepared the doubly isotopically pure ⁷LiOD as a precursor. ⁷LiOD was prepared by mixing a stoichiometric amount of ⁷LiCO₃ (Sigma Aldrich, 99 % for ⁷Li) and CaO (calcined from CaCO₃, Sigma Aldrich, 99 %) in D₂O. The CaCO₃ precipitate was filtered, and ⁷LiOD crystallized by evaporation of the solvent.

For a typical preparation of $(Li_{1,x}Fe_xOD)FeSe$, 5 mmol of Fe powder (Alfa Aesar, 99.9 %), 6 mmol of selenourea (Sigma Aldrich, 98 %) and 50 mmol of LiOD powder were suspended in 5 mL of distilled D₂O (Oxford Isotope, 99.9%).[§] The mixture was placed in a Teflon-lined stainless steel autoclave at 120-200 °C for 3-5 days. Afterwards, the autoclave was opened in an argon-filled glove bag, and the shiny black precipitate was washed with D₂O. The product was washed and centrifuged several times until the supernatant was clear. The remaining product was collected, vacuumed dried, and stored in a nitrogen-filled glove box. The yield of the product was usually between 50 % and 70 %.

Single crystal (Li_{1-x}Fe_xOD)FeSe and (Li_{1-x}Fe_xOH)FeSe samples were prepared by replacing potassium cations with LiOD or LiOH from K_xFe_{2-v}Se₂ single crystals under hydrothermal conditions similar to those reported by Dong *et al.*²³ For the growth of the $K_x Fe_{2-y}Se_2$ single crystals, 1.8 g (13 mmol) of FeSe powder was mixed with 0.21 g (5.4 mmol) of potassium metal (Alfa Aesar, 99 %) to match the nominal composition of $K_{0.8}$ Fe₂Se₂.^{24, 25} The FeSe precursor was prepared by heating Fe (Alfa Aesar, 99.9 %) and Se (Alfa Aesar, 99 %) powders to 700 °C for 5 h followed by furnace cooling; the resulting phase does not need to be of the tetragonal β -FeSe form for the crystal growth. The FeSe/K mixtures were loaded in a guartz ampoule inside a nitrogen-filled glovebox, and the ampoules flame sealed under vacuum. In order to avoid oxidation of the samples from breaking of the ampoule due to potassiuminduced corrosion of the quartz walls, the sample container was sealed in a larger ampoule. For crystal growth of K_xFe₂₋

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 $_{\rm y}$ Se₂, the mixture was heated to 1030 °C over 10 h and held at 1030 °C for 3 hours to form a homogeneous melt. Subsequently, the melt was slowly cooled at a rate of 6 °C /hour to 650 °C to allow crystal growth. After cooling to room temperature, K_xFe_{2-y}Se₂ single crystals approximately 8 mm in diameter were recovered.

In order to compare the effect of D₂O to the reaction kinetics, single crystals of both $(Li_{1-x}Fe_xOH)FeSe$ and $(Li_{1-x}Fe_xOD)FeSe$ were prepared under identical hydrothermal conditions. For the preparation of $(Li_{1-x}Fe_xOH)FeSe$ single crystals, the K_xFe₂, _ySe₂ precursor (0.2 g - 0.4 g), 0.14 g (2.5 mmol) Fe powder, and 2 g (47 mmol) LiOH monohydrate were added to 5 mL water. For $(Li_{1-x}Fe_xOD)FeSe$ single crystals, to match the concentration of LiOH in water, 1.2 g (47 mmol) LiOD and 6 mL D₂O were used for reactions. The mixture was placed in a Teflon-lined stainless steel autoclave at 120-200 °C for 4-5 days. Silver colored single crystals were recovered by washing away excess powder with water and drying under vacuum overnight. The as-recovered single crystals retained similar shape to the starting K_xFe_{2-y}Se₂ single crystals.

Samples prepared in the absence of excess iron powder were not superconducting, which could be due to either oxidation of the iron or vacancy formation in the FeSe layer. To study the role of metal powders during the cation exchange reactions, experiments using Sn metal (Alfa Aesar, 99.9 %) instead of Fe powder for the preparation of $(Li_{1-x}Fe_xOH)FeSe$ and $(Li_{1-x}Fe_xOD)FeSe$ single crystals at 120 °C were carried out. Sn can react with hot concentrated bases to form soluble $[Sn(OH)_6]^{2+}$ species while evolving H₂ gas,²⁶ thus providing a stronger reducing environment than the hydrothermal reactions without the presence of metal powders.

Laboratory and synchrotron X-ray diffraction measurements

Room temperature powder X-ray diffraction (PXRD) data were collected on a Bruker D8 X-ray diffractometer (Cu K α radiation, $\lambda = 1.5418$ Å). Data were collected with a step size of 0.02° between 7° and 70° for Pawley fits to extract lattice constants and 7° and 120° for Rietveld fits to obtain more accurate lattice and structural parameters.

In order to find any possible crystallographic phase transitions that are coupled to either the superconducting or magnetization order parameters, temperature dependent (5-300 K) high-resolution synchrotron X-ray diffraction was carried out for powders of ground single crystals at Beamline 11-BM at the Advanced Photon Source (APS). In addition to finding subtle changes in symmetry due to peak splitting, the synchrotron measurements provide high-*Q* data and therefore more accurate structural parameters. Analysis of the high-*Q* reflections help determine any small changes in the iron occupancies in both the FeSe and LiOH layers, which could affect T_c 's of the sample as suggested by Sun *et al.*¹⁸ An Oxford helium cryostat (closed flow system) was used to reach a temperature that is close to liquid helium (≈ 4 K).

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Figure 2. (a) X-ray and (b) neutron powder diffraction for non-superconducting (⁷Li₁. _xFe_xOD)FeSe with x = 0.166. No magnetic phase could be indexed in the NPD, indicating lack of antiferromagnetic ordering. Weight percent fractions from structural refinements are as follows: 98% (⁷Li_{1:x}Fe_xOD)FeSe (gold ticks) and 2% Li₂CO₃ (green ticks). A few broad peaks corresponding to FeSe were fit by a LeBail routine (asterisk).

Ground powders of single crystals were packed in 0.4 mm Kapton capillaries tubes and sealed with epoxy. Diffraction data were collected between 0.5° and 46° with a step size of 0.0001° using a constant wavelength λ = 0.413964 Å (30 keV).

Magnetization measurements

Magnetic susceptibility measurements were carried out using a magnetic property measurement system (Quantum Design MPMS). Both field-cooled (FC) and zero-field-cooled (ZFC) magnetic susceptibility measurements were taken from 2K -300 K in direct current mode with an applied magnetic field of 1 or 3 mT.

Neutron powder diffraction measurements

All the neutron work was carried out with doubly isotopically pure samples (⁷Li_{1-x}Fe_xOD)FeSe at the NIST Center for Neutron Research (NCNR). The samples were loaded into He-filled vanadium cans and subsequently into a closed cycle refrigerator for low temperature measurements. Low temperature (4 K) diffraction data were collected on the BT-1 high-resolution NPD with the Cu(311) monochromator (λ = 1.540 Å). In addition to base temperature measurements, we performed NPD measurements at various temperatures 25 K, 75 K, 150 K, and room temperature to search for any crystallographic phase transitions.

High-intensity and coarse-resolution diffraction measurements were carried out on the BT-7 spectrometer (λ = 2.359 Å) using

the position sensitive detector (PSD) to search for magnetic Bragg peaks from base temperature up to 150 ${\rm K.}^{27}$

Results

Crystallography and phase diagram

Rietveld refinements with both XRD and NPD data were carried out with the TOPAS 4.2 software.²⁸ Representative fits to one of the deuteroxide samples are presented in Figure 2 for both laboratory X-rays and neutrons. Although the samples are mostly phase pure, some starting reagent Li₂CO₃ was found as an impurity in the neutron data, which is more of a bulk technique than X-ray diffraction. Furthermore, two very broad peaks could be indexed as close to the lattice parameters of the parent phase β -FeSe. Indeed, these peaks have also been observed in previous work.^{17, 19} The much broader peak width for the FeSe impurity is indicative of very small crystallite size and quantitatively fitting this phase is not possible given its nearly amorphous nature.

The temperature-dependent synchrotron diffraction data did not reveal any major crystallographic changes in the structure (Figure 3). Therefore, unlike the parent FeSe phase, which undergoes a tetragonal to orthorhombic phase transition near 75 K,²⁹ LiOHFeSe and LiODFeSe remain tetragonal down to base temperature (10 K). Rietveld refinements of one of the deuteroxide patterns at 7 K and 150 K are presented in Figure 3, and relevant structural parameters are in Table 1 for both superconducting and non-superconducting deuteroxide phases. Relevant bond distances and bond angles are also shown in Table 1. Only results from the synchrotron X-ray dataset are presented in Table 1, and structural parameters

Table 1. Rietveld refinement with synchrotron PXRD data collected at 7 K for a superconducting sample of (⁷Li_{1x}Fe_xOD)FeSe shown in Figure 3 and a non-superconducting sample. Both samples are fitted to a *P4/nmm* space group with origin choice 1. The tetrahedral angles α_1 and α_2 represent the Se–Fe–Se angles in and out of the basal plane, respectively.

a = 3.7725(1) Å, c = 9.1330(2) Å, R _{wp} = 12.83%, T _c = 37 K						
Atom	Wyckoff site	x	у	Z	Occ.	U _{iso} (Å ²)
Li/Fe1	2b	0	0	0.5	0.827/0.173(2)	0.0134
Fe2	2a	0.5	0.5	0	0.979(2)	0.0057
0	2c	0.5	0	0.4266(3)	1	0.0037(7)
Se	2c	0	0.5	0.1603(1)	1	0.0028(2
α_1 (°)	α_2 (°)	Fe-Fe (Å)		Fe-Se (Å)		
104.38(2)	112.07(1)	2.6675(1)		2.3875(4)		
a = 3.7820(1) Å, c = 9.0992(1) Å, R _{wp} = 10.66%, non-superconducting						
Atom	Wyckoff site	x	y	Z	Occ.	U _{iso} (Å ²)
Li/Fe1	2b	0	0	0.5	0.809/0.191(2)	0.0156
Fe2	2a	0.5	0.5	0	0.919(2)	0.0036
0	2c	0.5	0	0.4252(3)	1	0.0038(1
Se	2c	0	0.5	0.1609(1)	1	0.0019(6
α_1 (°)	<i>α</i> ₂ (°)	Fe-Fe (Å)		Fe-Se (Å)		
104.51(2)	112.01(1)	2.6743(1)		2.3914(3)		



Figure 3. Synchrotron powder X-ray diffraction pattern at (a) 7 K and (b) 150 K for a single crystal sample of (Li1-xFexOD)FeSe prepared at 120 °C for 5 days (T_c = 37 K). Rietveld refinement with the data collected at both temperatures did not reveal any lowering of symmetry from tetragonal *P4/nmm*. Tick marks representing the tetragonal phase are shown below the calculated, observed, and differences curves. The insets shown are a zoom in of the high-angle synchrotron data.

from the Rietveld refinements, including the neutron data, for the rest of the samples used to construct the full phase diagrams can be found in the ESI (Tables S1-S5).

Magnetization results and the phase diagrams

The SQUID magnetic susceptibility measurements for the series of hydroxide samples prepared through the powder routes are presented in Figure 4a. The deuteroxide samples, which were all derived from the single crystal route, are shown in Figure 4b. Only one sample within the hydroxide series expressed a ferromagnetic signal within the superconducting regime. Although some studies have found no evidence of ferromagnetism below T_c in their samples,^{1, 20} under the right synthesis conditions, we have observed a ferromagnetic signal in the superconducting regime as first reported by Pachmayr *et al.*¹⁷ A similar plot (Figure S1) for the hydroxide system prepared via the single crystal route can be found in the ESI.

We have constructed superconducting phase diagrams in Figure 5 that relate the critical temperatures T_c to the lattice constants compiled from the SQUID data and diffraction results from all the samples. The lattice parameters of the tetragonal unit cell found at room temperature were used as the x-axis versus T_c in the phase diagrams. More specifically, we found the best correlation to T_c is that of the tetragonality parameter, which is the simple c/a ratio. The corresponding superconducting volume fractions ($4\pi\chi$) were also established by SQUID magnetometry (Figure 4).

We found that T_c and its volume fraction increased with the lattice constant c and decreased with lattice constant a.

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Therefore, those with the highest tetragonality gave the maximum T_c and superconducting volume fractions. For samples to exhibit significant superconducting volume fractions ($4\pi\chi > 10$ %), the lattice constant c must be larger than ~9.20 Å and a smaller than ~3.80 Å. These trends in the lattice parameters are consistent with the findings of Sun *et al.* on their hydroxide analogues,¹⁸ where they attribute a large a lattice constant to iron vacancies in the FeSe layers and therefore iron slightly oxidized above 2+.

Our combined diffraction experiments did indeed find variations on the iron occupancies, whether superconducting or non-superconducting. As Table 1 shows, when the occupancy of the Fe2 site falls from near full to 91.9(2) %, superconductivity is lost. The differences overall between the hydroxide samples prepared by powder routes and the single crystal ones could be due the accommodation of iron vacancies during the syntheses. As Sun et al. found in their samples,²⁸ when FeSe in its tetragonal β -phase is used as the host for intercalation via hydrothermal synthesis, iron powder is necessary as a co-reagent in order to fill in the resulting vacancies. When starting with Fe powder and selenourea, we found more variability in the amount of iron vacancies and therefore a larger spread in the tetragonality parameter (Figure 5). Our powder preparation therefore can be described as in-situ growth of alternating FeSe and Li_{1-x}Fe_xOD layers whereas our single crystal preparation is a post-synthetic modification method (also known as soft chemistry).



Figure 4. Magnetic susceptibility of samples prepared by (a) powder routes and (b) single crystal routes of deuteroxide series ($Li_{1.x}Fe_xOD$)FeSe. For $T_c = 32$ K, a ferromagnetic transition can be noted at $T_f = 10$ K. Zero-field cooled data of powder samples and single crystal samples collected with applied fields of 1 mT and 3 mT, respectively, are shown.



Figure 5. Superconducting phase diagrams of (a) ($Li_{1:x}Fe_xOH$)FeSe as comparison for that of (b) single crystal ($^7Li_{1:x}Fe_xOD$)FeSe and (c) powder samples and ($Li_{1:x}Fe_xOD$)FeSe. The critical temperatures T_c are related to the tetragonality parameter, which is the simple ratio of the lattice parameters c/a. Note that for the single crystal samples, the crystals exhibited a superconducting phase fraction typically larger than 100 %, likely due to the geometry of the crystals.

Neutron results

То verifv whether any of the samples exhibit antiferromagnetism, we searched for any superlattice peaks in the NPD patterns that could arise below 120 K, the AFM transition in the parent phase of (Li_{1-x}Fe_xOD)FeSe reported by Dong *et al.*¹ No superlattice reflections were observed in the BT-1 high-resolution NPD patterns, and our deuterated samples allowed for a low background in case of a small Fe signal. Indeed, in the arsenide systems the iron moment can be small in the parent phases such as 0.36(5) $\mu_{\rm B}/{\rm Fe}$ in LaOFeAs³⁰ and 0.25(7) μ_B /Fe in NdOFeAs,³¹ and any hydrogen incoherent background would easily overwhelm such a small signal from long-range magnetic ordering in the NPD.

The samples were measured up to 150 K on BT-1, and no long range magnetic ordering was observed. NPD patterns were taken with a PSD on BT-7, which has a much higher flux than BT-1 at low angles, also revealed no antiferromagnetic peaks in the non-superconducting samples (Figure 6). Difference patterns between 150 K and 4 K are shown in Figure 6, revealing no residual intensity and only differences arising from thermal expansion and thus peak positions.

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Discussion

Relation between structural parameters and superconductivity

In preparing our deuteroxide samples, we found the reaction temperature to influence the lattice constants. Mild hydrothermal reaction temperatures (120 °C) led to samples expressing a higher T_c , while the reaction temperature above 180 °C led exclusively to either non-superconducting samples or ones with very low volume fractions ($4\pi\chi < 1$ %). Reaction times also affected the lattice constants. Longer reaction times (> 3 days) yielded samples with slightly larger *a* and smaller *c* (*i.e.* smaller tetragonality parameters).

While all the deuterated samples followed the trend shown in the phase diagram (Figure 5c), similarly prepared hydrated samples deviated in their behaviour. Indeed, some hydroxide samples with lattice parameters matching those in the phase diagram from earlier literature¹ did not exhibit superconductivity (Figure 4a).

Interestingly, samples prepared at lower temperatures with the described mixing ratio and longer reaction times expressed coexistence between superconductivity and ferromagnetism (Figure 3a). Thus, while longer reaction times above 180 °C led to lower T_c 's or non-superconducting samples, longer reaction times at lower temperatures (120 °C) produced the ferromagnetic signal in superconducting (Li_{1-x}Fe_xOD)FeSe. As described in the next section, this might be a kinetic effect from the increasing amount of oxidized iron in water from longer reaction times.

Relation between magnetism and superconductivity

In several of our non-superconducting samples, we have observed an antiferromagnetic transition close to 120 K. Dong et al. have claimed that the hydroxide samples in the nonsuperconducting dome were antiferromagnetic parent phases with a T_N close to 120 K, and therefore that the selenides and arsenides have the same underlying physics with respect to the superconducting mechanism. This is a very important claim that could have large implications in the field of ironbased superconductors. None of our non-superconducting deuteroxide samples, however, exhibited this antiferromagnetic signal in the SQUID measurements, which led us to believe that the antiferromagnetism may not be intrinsic to this system.

Our findings in the preparation of hydroxide and deuteroxide samples revealed the strong possibility that the 120 K transition observed in the SQUID magnetization measurements arise from iron oxide impurities. The so-called Verwey transition, which corresponds to a Fe²⁺/Fe³⁺ charge ordering transition in Fe₃O₄ also occurs near 120 K.^{32, 33} Furthermore, structurally related Li_xFeO₂ phases can express T_N from 100 K to 300 K according to amount of intercalated lithium cations.^{34, 35}

In order to study the formation of iron oxide impurities, a sample was prepared under similar hydrothermal conditions but without the addition of selenourea. As pointed out by Sun et al. in their extensive study of the formation of (Li₁. _xFe_xOH)FeSe, the strongly basic conditions (pH > 14) of the



Figure 6. (a) The NPD pattern of non-superconducting phase of $(^7Li_{1:x}Fe_xOD)FeSe$ at 150 K and (b) 4 K. The difference between the two patterns in (c) reveals no antiferromagnetic peaks.

synthesis strongly favour the formation of Fe³⁺ species according to the electrochemical-pH phase equilibrium diagram (i.e. Pourbaix) of iron.³⁶ Therefore, without the selenourea reagant to stabilize Fe²⁺ species, a large amount of mixed valent iron oxides are produced from hydrothermal and highly basic conditions.

The XRD pattern of the as-recovered sample from hydrothermal synthesis without selenourea was fit by the Fe₃O₄ structure, and its magnetic susceptibility measurement was in very good agreement with typical Verwey transition at 120 K (Figure 7a,b). As shown in Figure 7c, a (Li_{1-x}Fe_xOH)FeSe sample with lattice constants in the supposed superconducting region displayed no super-conductivity, but a transition similar to charge ordering in Fe₃O₄. Peaks in the XRD pattern (indicated by * in Figure 7b) that cannot be matched with the (Li1., Fe, OH)FeSe phase was indexed well with the strongest peaks of Fe₃O₄. In addition, our synchrotron XRD data for a non-superconducting (Li_{1-x}Fe_xOH)FeSe single crystal sample revealed small amounts of Fe₃O₄ impurity (Figure S2), which was not observed by laboratory X-ray measurements. Therefore, it is likely that the 120 K transition in nonsuperconducting (Li_{1-x}Fe_xOH)FeSe samples is extrinsic and due to magnetic impurities Fe_3O_4 or structurally related Li_xFeO_2 .

Hydrothermal synthesis of samples with D_2O under similar conditions as those with H_2O did not lead to appreciable oxide impurity. We therefore conclude that the observed differences in the acid-base chemistry of H_2O and D_2O lead to different products for similar reaction conditions. Indeed, a hydrothermal treatment of iron powder with D_2O and with LiOD did not lead to complete conversion to Fe₃O₄ but left unreacted iron powder (approximately 50%). Since under highly basic conditions, Fe³⁺ should be favoured thermodynamically, we believe the kinetics for the oxidation of iron with D_2O is slower than in H_2O due to the stronger D–O bond than the H–O bond.

All deuteroxide single crystal samples prepared in D_2O showed noticeably higher T_c 's than their hydroxide counterparts prepared at identical conditions (Table S2). The T_c of hydroxide samples can be improved by reducing the reaction time, which likely minimizes the extent of Fe oxidation.



Figure 7. (a) Powder XRD and Rietveld analysis of phase pure Fe_3O_4 prepared under similar hydrothermal conditions to that of $(Li_{1,x}Fe_xOH)FeSe$ in the absence of selenourea. (b) The corresponding magnetization data of the Fe_3O_4 sample indicating the Verwey transition near 125 K. (c) The powder XRD of $(Li_{1,x}Fe_xOH)FeSe$ with the Fe_3O_4 impurity marked along its strongest reflection. (d) The corresponding magnetization data of this non-superconducting $(Li_{1,x}Fe_xOH)FeSe$ sample, exhibiting the 125K transition similar to that of Fe^{2+}/Fe^{3+} charge ordering seen in Fe_3O_4 (inset).

Although we have established here that the (Li_{1-x}Fe_xOH)FeSe system does not likely have a parent antiferromagnetic phase, we do not suggest that the chalcogenide-based superconductors are not linked to the arsenide-based ones. The Fe_{1+x}TeyCh_{1-y} for *Ch* = Se and S phases, for example, exhibit a rich magnetic phase diagram^{37, 38} before superconductivity sets in with chemical substitution.³⁹ The ordered vacancy phase of K_{0.8}Fe_{1.6}Se₂ has also shown an antiferromagnetic transition at high temperatures (~559 K)⁴⁰ while the disordered vacancy phase exhibits a *T_c* close to 30 K.^{41.45} What distinguishes those two systems, however, from the present compound and FeSe, is the lack of a large magnetic moment on iron in the latter compounds.⁴⁶⁻⁵¹ In Fe_{1+x}Te it can be as large as 2 μ_B^{52} and in K_{0.8}Fe_{1.6}Se₂ as large as 3.3(1) μ_B .⁴⁰

As to the ferromagnetic transition observed at ~10 K in the sample with a T_c of 34 K, several authors have also observed it in in the hydroxide analogues. Pachmeyer *et al.* attribute the long range magnetic order at 18 K to the iron cations partially substituted on the Li site,¹⁷ while Lu *et al.* assign this transition as being antiferromagnetic (~12 K) according to their NMR studies.¹⁹ Our recent small angle neutron scattering study illustrates the formation of long-range magnetic order below 12.5 K, but with a moment too small to see with diffraction.²² This observation likely arises from the dilute occupancy by the iron cation on the Li site.

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Conclusions

In conclusion, we have successfully mapped out a phase 6. diagram for (Li_{1-x}Fe_xOD)FeSe and have found that the highest T_c for deuterated samples is 42 K, and the T_c for both 7. deuterated and hydroxide samples correlate with lattice constants. From synchrotron and neutron diffraction studies, iron vacancies in the FeSe layer seem to lead to the low 8. tetragonality ratio. The phase diagram of our samples therefore shows a dome-like appearance for the superconducting regime, which had not been observed 9. before.^{1,18} Since the presence of more vacancies implies a higher oxidation state for iron, the samples with optimal Tc does suggest that an optimal electron doping is also necessary for the selenide system. Since the highest T_c observed for the hydroxide sample was also approximately 42 K, we conclude that there is no isotope effect on the superconducting properties in substituting H by D. Mild hydrothermal preparation (120 °C) for long reaction times can lead to the coexistence of ferromagnetism and superconductivity. Finally, any claims of anti-ferromagnetism in the parent phase of this system should be re-evaluated in light of the easy preparation of oxide impurities with transition temperatures near the vicinity of 120 K in H_2O .

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Notes and references

§ The identification of any commercial product does not imply endorsement or recommendation by the National Institute of Standards and Technology.

- X. Dong, H. Zhou, H. Yang, J. Yuan, K. Jin, F. Zhou, D. Yuan, L. Wei, J. Li and X. Wang, J. Am. Chem. Soc., 2014, 137, 66-69.
- 2. J. Paglione and R. L. Greene, *Nat. Phys.*, 2010, **6**, 645-658.
- 3. D. C. Johnston, Adv. Phys., 2010, **59**, 803-1061.
- T. M. McQueen, Q. Huang, V. Ksenofontov, C. Felser, Q. Xu, H. Zandbergen, Y. S. Hor, J. Allred, A. J. Williams, D. Qu, J. Checkelsky, N. P. Ong and R. J. Cava, *Phys. Rev. B*, 2009, **79**, 014522.
- F.-C. Hsu, J.-Y. Luo, K.-W. Yeh, T.-K. Chen, T.-W. Huang, P. M. Wu, Y.-C. Lee, Y.-L. Huang, Y.-Y. Chu, D.-C. Yan and M.-

K. Wu, Proc. Natl. Acad. Sci. U.S.A., 2008, **105**, 14262-14264.

- H. Kotegawa, S. Masaki, Y. Awai, H. Tou, Y. Mizuguchi and Y. Takano, *J. Phys. Soc. Jpn.,* 2008, **77**, 113703.
- S. Margadonna, Y. Takabayashi, M. T. McDonald, K. Kasperkiewicz, Y. Mizuguchi, Y. Takano, A. N. Fitch, E. Suard and K. Prassides, *Chem. Comm.*, 2008, DOI: 10.1039/B813076K, 5607-5609.
- S. Medvedev, T. McQueen, I. Troyan, T. Palasyuk, M. Eremets, R. Cava, S. Naghavi, F. Casper, V. Ksenofontov and G. Wortmann, *Nat. Mater.*, 2009, **8**, 630-633.
- T. Imai, K. Ahilan, F. L. Ning, T. M. McQueen and R. J. Cava, Phys. Rev. Lett., 2009, **102**, 177005.
- E.-W. Scheidt, V. Hathwar, D. Schmitz, A. Dunbar, W. Scherer, F. Mayr, V. Tsurkan, J. Deisenhofer and A. Loidl, *Eur. Phys. J. B*, 2012, **85**, 1-5.
- M. Burrard-Lucas, D. G. Free, S. J. Sedlmaier, J. D. Wright,
 S. J. Cassidy, Y. Hara, A. J. Corkett, T. Lancaster, P. J. Baker,
 S. J. Blundell and S. J. Clarke, *Nat. Mater.*, 2013, **12**, 15-19.
- T. Ying, X. Chen, G. Wang, S. Jin, X. Lai, T. Zhou, H. Zhang, S. Shen and W. Wang, J. Am. Chem. Soc., 2013, 135, 2951-2954.
- S. J. Sedlmaier, S. J. Cassidy, R. G. Morris, M. Drakopoulos, C. Reinhard, S. J. Moorhouse, D. O'Hare, P. Manuel, D. Khalyavin and S. J. Clarke, J. Am. Chem. Soc., 2014, 136, 630-633.
- 14. J.-F. Ge, Z.-L. Liu, C. Liu, C.-L. Gao, D. Qian, Q.-K. Xue, Y. Liu and J.-F. Jia, *Nat. Mater.*, 2015, **14**, 285-289.
- 15. S. He, J. He, W. Zhang, L. Zhao, D. Liu, X. Liu, D. Mou, Y.-B. Ou, Q.-Y. Wang and Z. Li, *Nat. Mater.*, 2013, **12**, 605-610.
- 16. S. Tan, Y. Zhang, M. Xia, Z. Ye, F. Chen, X. Xie, R. Peng, D. Xu, Q. Fan and H. Xu, *Nat. Mater.*, 2013, **12**, 634-640.
- U. Pachmayr, F. Nitsche, H. Luetkens, S. Kamusella, F. Brückner, R. Sarkar, H.-H. Klauss and D. Johrendt, *Angew. Chem. Int. Ed.*, 2015, **54**, 293-297.
- H. Sun, D. N. Woodruff, S. J. Cassidy, G. M. Allcroft, S. J. Sedlmaier, A. L. Thompson, P. A. Bingham, S. D. Forder, S. Cartenet, N. Mary, S. Ramos, F. R. Foronda, B. H. Williams, X. Li, S. J. Blundell and S. J. Clarke, *Inorg. Chem.*, 2015, 54, 1958-1964.
- X. F. Lu, N. Z. Wang, H. Wu, Y. P. Wu, D. Zhao, X. Z. Zeng, X. G. Luo, T. Wu, W. Bao, G. H. Zhang, F. Q. Huang, Q. Z. Huang and X. H. Chen, *Nat. Mater.*, 2015, 14, 325-329.
- X. F. Lu, N. Z. Wang, G. H. Zhang, X. G. Luo, Z. M. Ma, B. Lei, F. Q. Huang and X. H. Chen, *Phys. Rev. B*, 2014, **89**, 020507.
- 21. U. Pachmayr and D. Johrendt, *Chem. Comm.*, 2015, **51**, 4689-4692.
- J. W. Lynn, X. Zhou, C. K. H. Borg, S. R. Saha, J. Paglione and E. E. Rodriguez, *Physical Review B*, 2015, **92**, 060510.
- X. Dong, K. Jin, D. Yuan, H. Zhou, J. Yuan, Y. Huang, W. Hua, J. Sun, P. Zheng and W. Hu, *Phys. Rev. B*, 2015, **92**, 064515.
- 24. Y. Mizuguchi, H. Takeya, Y. Kawasaki, T. Ozaki, S. Tsuda, T. Yamaguchi and Y. Takano, *Appl. Phys. Lett.*, 2010, **98**, 042511.
- J. Ying, X. Wang, X. Luo, A. Wang, M. Zhang, Y. Yan, Z. Xiang, R. Liu, P. Cheng and G. Ye, *Phys. Rev. B*, 2011, 83, 212502.
- 26. P. G. Harrison, *Chemistry of tin*, Blackie, 1989.
- 27. J. Lynn, Y. Chen, S. Chang, Y. Zhao, S. Chi, W. Ratcliff, B. Ueland and R. Erwin, *J. Res. NIST*, 2012, **117**, 61-79.

- R. W. Cheary and A. Coelho, J. Appl. Crystallogr., 1992, 25, 51. 109.
- T. McQueen, A. Williams, P. Stephens, J. Tao, Y. Zhu, V. 52. Ksenofontov, F. Casper, C. Felser and R. Cava, *Phys. Rev. Lett.*, 2009, **103**, 057002.
- C. de La Cruz, Q. Huang, J. Lynn, J. Li, W. Ratcliff II, J. L. Zarestky, H. Mook, G. Chen, J. Luo and N. Wang, *Nature*, 2008, 453, 899-902.
- Y. Chen, J. Lynn, J. Li, G. Li, G. Chen, J. Luo, N. Wang, P. Dai, C. dela Cruz and H. Mook, *Phys. Rev. B*, 2008, **78**, 064515.
- 32. F. Walz, J Phys.: Condens. Matter, 2002, 14, R285.
- 33. Z. Zhang and S. Satpathy, *Phys. Rev. B*, 1991, **44**, 13319.
- M. Tabuchi, K. Ado, H. Kobayashi, I. Matsubara, H. Kageyama, M. Wakita, S. Tsutsui, S. Nasu, Y. Takeda and C. Masquelier, *J. Solid State Chem.*, 1998, 141, 554-561.
- M. Tabuchi, S. Tsutsui, C. Masquelier, R. Kanno, K. Ado, I. Matsubara, S. Nasu and H. Kageyama, J. Solid State Chem., 1998, 140, 159-167.
- J.-M. R. Génin, G. Bourrié, F. Trolard, M. Abdelmoula, A. Jaffrezic, P. Refait, V. Maitre, B. Humbert and A. Herbillon, *Environ. Sci. Technol.*, 1998, **32**, 1058-1068.
- E. E. Rodriguez, C. Stock, P. Zajdel, K. L. Krycka, C. F. Majkrzak, P. Zavalij and M. A. Green, *Phys. Rev. B*, 2011, 84, 064403.
- P. Zajdel, P.-Y. Hsieh, E. E. Rodriguez, N. P. Butch, J. D. Magill, J. Paglione, P. Zavalij, M. R. Suchomel and M. A. Green, J. Am. Chem. Soc., 2010, 132, 13000-13007.
- Y. Mizuguchi and Y. Takano, J. Phys. Soc. Jpn., 2010, 79, 102001.
- W. Bao, Q.-Z. Huang, G.-F. Chen, M. A. Green, D.-M. Wang, J.-B. He and Y.-M. Qiu, *Chin. Phys. Lett.*, 2011, 28, 086104.
- 41. J. Guo, S. Jin, G. Wang, S. Wang, K. Zhu, T. Zhou, M. He and X. Chen, *Phys. Rev. B*, 2010, **82**, 180520.
- W. Hang-Dong, D. Chi-Heng, L. Zu-Juan, M. Qian-Hui, Z. Sha-Sha, F. Chun-Mu, H. Q. Yuan and F. Ming-Hu, *Eur. Phys. Lett.*, 2011, **93**, 47004.
- X. G. Luo, X. F. Wang, J. J. Ying, Y. J. Yan, Z. Y. Li, M. Zhang,
 A. F. Wang, P. Cheng, Z. J. Xiang, G. J. Ye, R. H. Liu and X.
 H. Chen, *New J. Phys.*, 2011, **13**, 053011.
- F. Ming-Hu, W. Hang-Dong, D. Chi-Heng, L. Zu-Juan, F. Chun-Mu, C. Jian and H. Q. Yuan, *Eur. Phys. Lett.*, 2011, 94, 27009.
- A. F. Wang, J. J. Ying, Y. J. Yan, R. H. Liu, X. G. Luo, Z. Y. Li, X. F. Wang, M. Zhang, G. J. Ye, P. Cheng, Z. J. Xiang and X. H. Chen, *Phys. Rev. B*, 2011, **83**, 060512.
- V. Y. Pomjakushin, D. Sheptyakov, E. Pomjakushina, A. Krzton-Maziopa, K. Conder, D. Chernyshov, V. Svitlyk and Z. Shermadini, *Phys. Rev. B*, 2011, 83, 144410.
- D. Wang, J. He, T.-L. Xia and G. Chen, *Phys. Rev. B*, 2011, 83, 132502.
- P. Zavalij, W. Bao, X. Wang, J. Ying, X. Chen, D. Wang, J. He, X. Wang, G. Chen and P.-Y. Hsieh, *Phys. Rev. B*, 2011, 83, 132509.
- D. P. Shoemaker, D. Y. Chung, H. Claus, M. C. Francisco, S. Avci, A. Llobet and M. G. Kanatzidis, *Phys. Rev. B*, 2012, 86, 184511.
- X. Ding, D. Fang, Z. Wang, H. Yang, J. Liu, Q. Deng, G. Ma, C. Meng, Y. Hu and H.-H. Wen, *Nat. Commun.*, 2013, 4, 1897.

- D. Guterding, H. O. Jeschke, P. Hirschfeld and R. Valenti, *Phys. Rev. B*, 2015, **91**, 041112.
- W. Bao, Y. Qiu, Q. Huang, M. Green, P. Zajdel, M. Fitzsimmons, M. Zhernenkov, S. Chang, M. Fang and B. Qian, *Phys. Rev. Lett.*, 2009, **102**, 247001.

