

RSC Advances

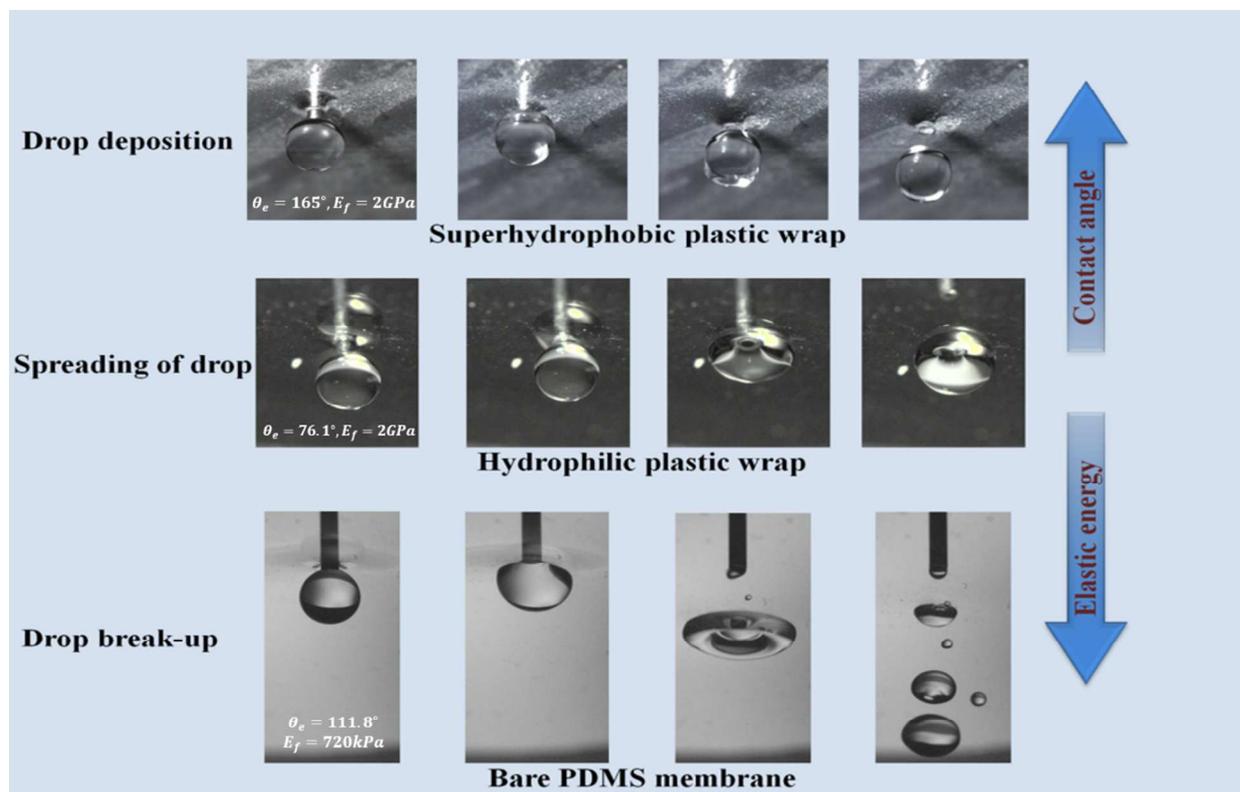


This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. This *Accepted Manuscript* will be replaced by the edited, formatted and paginated article as soon as this is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.





Cite this: DOI: 10.1039/xxxxxxxxxx

Needle-free drop deposition: Role of elastic membranes

Prashant Waghmare,^a Surjyasish Mitra,^b Naga Siva Kumar Gunda,^b and Sushanta K. Mitra^{*b}

Received Date

Accepted Date

DOI: 10.1039/xxxxxxxxxx

www.rsc.org/journalname

Contact angle measurement of low-energy surfaces (superhydrophobic, superoleophobic, etc.) with needle-drop assembly is critical towards characterizing such substrates. However, it is extremely difficult to detach the needle from the drop when it is brought in contact with a characterizing substrate and often one has to report contact angles with the needle attached with the drop itself. To overcome this challenge, here, we present a new technique to achieve a 'needle-free' drop by bringing the drop in contact with an additional elastic membrane, kept between the needle-drop assembly and the characterizing substrate. The detachment of the drop from the needle is achieved by retracting the needle-drop assembly at a finite speed and allowing the drop to receive the elastic energy of the soft flexible membrane. Such interaction of the drop with the elastic membrane allows the drop to get repelled from the elastic membrane and further gets deposited on a characterizing substrate. The repelling behavior of the drop can be controlled by appropriately selecting the mechanical and wetting properties of this additional elastic membrane. This technique not only provides a needle-free drop deposition that is independent of the physical properties of the liquid and the needle but it also allows achieving the drop size that is independent of the needle diameter. The experimental analysis and theoretical investigations suggest that the mechanical property of the elastic membrane, particularly its elasticity, plays an important role towards the success of the drop deposition technique.

1 Introduction

Drop deposition on a given target substrate has a plethora of applications starting from characterizing the wetting properties of natural, bio-inspired, and artificial micro/nano-structured surfaces¹⁻⁵ to drop impact studies to determine the pre and post-impact dynamics⁶⁻¹⁰. In all of these examples, the critical step is to detach the drop reliably from the needle that is typically used in any drop deposition method - be it using the traditional pendant drop method⁶ or drop weight method^{4,11}. Often this critical step poses a challenge in itself when the surface energy of the needle-drop assembly is comparable with the surface energy of the drop and the characterizing substrate^{1,12,13} or when a significant height is required to dispense the drop reliably by its own weight on a given substrate or the needle diameter produces only finite range of drop sizes that can be deposited onto a given substrate¹⁴⁻²⁰. In the latter case, the volume of the drop

poses a limitation as one needs a significantly large drop volume to achieve deposition due to its own weight²¹.

Contact angle measurement has been one of the crucial tasks for accurate quantification of the wetting characteristics of a given substrate²²⁻²⁶. Recently, we have proposed a 'needle-free' drop deposition technique for a contact angle measurement of a substrate placed under-water, where a favorable spreading parameter between the drop, the surrounding liquid medium and the fluid at the interface of the surrounding medium and the needle has been appropriately configured to detach the drop from the needle¹. This study was recently been extended to achieve similar needle-free drop deposition, but in air medium²⁷, which is the common scenario for characterizing substrates or performing drop-impact studies. However, this technique²⁷ has a limitation in terms of minimum drop volume that could be detached from the needle, which has implication in terms of the drop radius in comparison to the capillary length scale for wetting studies²⁸, achievable Weber numbers and Reynolds numbers for drop impact studies^{17,18} etc. In our previous study²⁷, the drop detachment was achieved by bringing in an additional superhydrophobic rigid substrate and the needle with the drop at its end is withdrawn through such surface, the interaction of the drop with the

^a Department of Mechanical Engineering, University of Alberta, Edmonton, AB, T6G 2G8, Canada

^b Micro & Nano-scale Transport Laboratory, Department of Mechanical Engineering, Lassonde School of Engineering, York University, Toronto, Ontario M3J 1P3, Canada.

*E-mail: sushanta.mitra@lassonde.yorku.ca; Tel: +1 416 736 5924

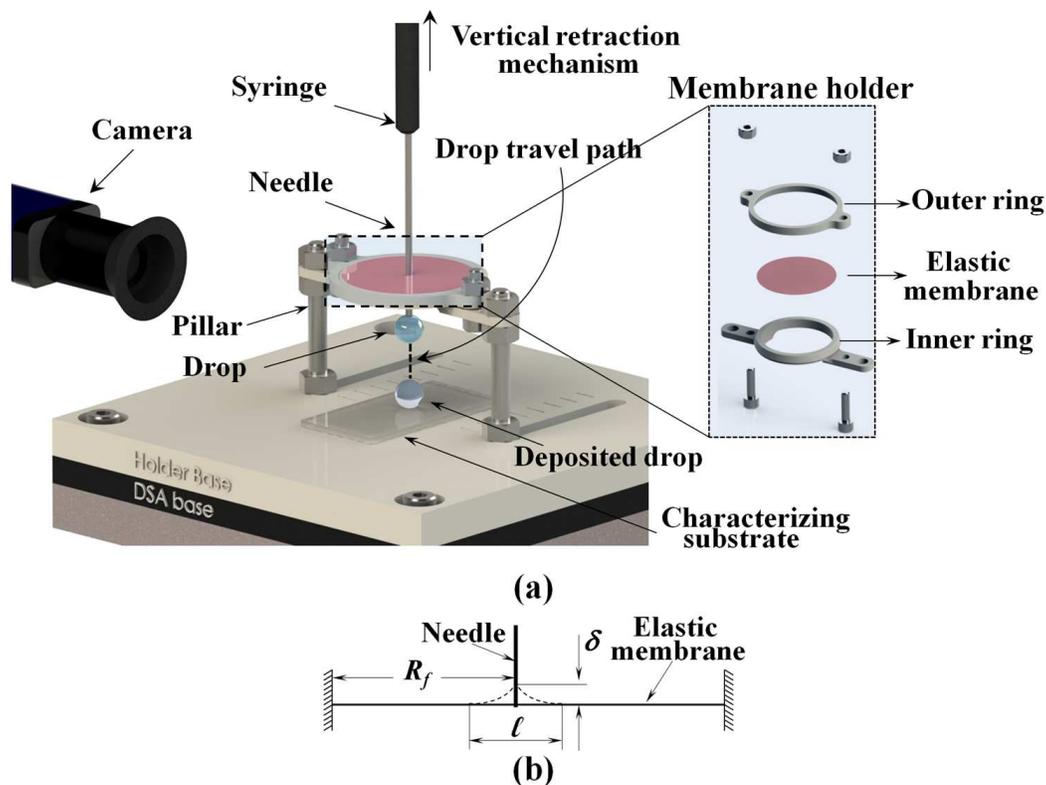


Fig. 1 (a) Schematic of a needle-free drop deposition system using elastic membrane. A specialized holder, shown in details in the inset diagram, is used to maintain the elastic membrane in a perfectly horizontal, wrinkle-free configuration. (b) Membrane configuration before (solid line) and after (dotted line) the impact with upward deflection over a length for a membrane of radius R_f .

superhydrophobic substrate creates the detachment of the drop from the needle and allows the drop to be deposited on the characterizing substrate kept below the superhydrophobic substrate. However, the extent of superhydrophobicity (attained by using special coating to achieve static contact angle $> 165^\circ$) alone could not produce reliable drop detachment for drop volume less than $3\mu\text{L}$ for needle diameter of 0.5 mm. Moreover, the argument of using a superhydrophobic coating on the needle itself to make the drop also has its own shortcomings. For that technique to work, prior knowledge of surface energies of the needle material and the characterizing substrate is required, which then itself defeat the very purpose of developing reliable tools to measure contact angle (or interpret the surface energies of unknown substrates).

The ‘needle-free’ drop deposition technique presented in this work is motivated by studies conducted in literature for understanding the drop impact on elastic membranes^{29,30}, where ethanol drops from a certain height were imparted on a stretchable polymer film kept below and it was found that drop impact energy gets transferred to the elastic (soft) membrane thereby reducing the drop splashing effects. Hence, it can be inferred that the additional energy transfer mechanism between the drop and a substrate can be obtained by utilizing a deformable solid substrate, as opposed to a rigid superhydrophobic surface, as done in our earlier study²⁷. In this present work, we take this added advantage provided by such soft membranes, where we coat such membranes with superhydrophobic coating and utilize it as a top low-energy substrate to obtain ‘needle-free’ drop deposition in air

medium.

2 Experimental Methods

A schematic of the experimental setup used in this study is shown in Fig. 1. The experimental set-up consists of a regular contact angle measurement system, DSA 100 (Krüss, Germany) with two key modifications. First, we have designed a special holder (details provide in inset diagram (a) in Fig. 1), which allows an elastic (soft) deformable membrane to be attached in a particular fashion to avoid any wrinkling in the elastic membrane and further undesirable defects due to non-uniform tension³¹. To achieve this condition, the elastic membrane was secured between two rings; both rings were fabricated using the prototype machine Eden 350V, Stratasys Ltd., USA. The one end of the inner ring was clamped to the outer ring (with the membrane held in between), whereas the other end of the inner ring was fixed to the pillar-stand of this specially designed membrane holder. This holder with the elastic membrane was attached firmly to the characterizing base-table of the DSA system with nuts and bolts. This arrangement not only allowed to place a membrane with negligible initial tension or stretching but also provided a perfectly horizontal membrane configuration, which remain parallel with the bottom characterizing substrate placed at the base of the DSA system. The second modification is done by attaching the needle-drop assembly to a traversing mechanism (Zaber, Canada) which allows the needle to be retracted vertically at a speed varying between $0.1\text{mm/s} - 100\text{mm/s}$. This retraction mechanism is sim-

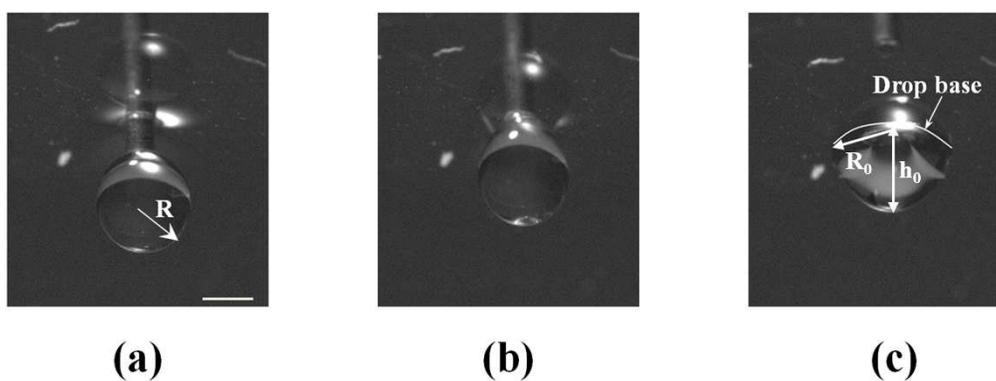


Fig. 2 Snap shots of the interaction of the needle-drop assembly with the hydrophilic elastic membrane (plastic wrap) $\theta \approx 76.1^\circ$ and needle retraction speed of 30mm/s. (a) Drop-needle assembly before impacting the elastic membrane (b) Impact of the drop on the elastic membrane. (c) The equilibrium configuration of the drop after spreading on to the elastic membrane. The scale-bar at the bottom of Fig. 2(a) represents 1mm.

ilar to the one used in our previous study²⁷. The interaction of the needle-drop assembly with the elastic membrane causes the membrane to be deflected with an amount δ in the vertical upward direction, as shown in Fig. 1(b). The elastic (soft) deformable membrane, used in this study, is a plastic wrap (NoName brand) purchased from a local store (Edmonton, Canada). The wetting properties of this plastic wrap were determined by performing static equilibrium contact angle measurements ($\theta = 76.1^\circ \pm 6.6$), using DSA 100 (Krüss, Germany) system. This measured value is based on Young-Laplace equation for equilibrium contact angle. A more accurate measurement can be performed by using the ‘Tadmor method’³³, however, in case of flexible substrates, like the plastic wrap used here, the definition of a static contact angle is still in an evolving phase³⁴. A commercially available coating spray (NeverWet) was used to alter the wetting properties of this elastic membrane, thereby achieving equilibrium contact angle $\sim 165^\circ$. Deionized water (PURELAB Ultra, ELGA) was used as a working fluid to generate drops with the help of stainless steel needles (NE44, Krüss, Germany) and the glass syringe (SY20, Krüss, Germany). To perform the needle-free drop deposition, the elastic membrane was punctured using a stainless steel needle (needle diameter 0.5mm) approximately at the center of the membrane to create a hole. Then the needle attached with the water filled syringe was inserted through this hole. After that, a drop of a certain volume ($2\mu\text{L}$) was generated at the tip of the needle. Finally, the drop-needle assembly was retracted with a finite velocity to allow the drop to impact on the bottom coated surface of the elastic membrane. The interaction (impact) of the drop with the elastic membrane was recorded using a Phantom V711 (ViSion Research, USA).

3 Results and discussions

The ‘needle-free’ drop deposition was first attempted with the uncoated elastic membrane (hydrophilic deformable substrate) attached to the holder, as shown in Fig. 2. A drop volume of $2\mu\text{L}$ was produced with the needle of outside diameter 0.5mm and was withdrawn through the elastic membrane. The needle retraction speed in this case was maintained at 30 mm/s. Fig-

ure 2 shows different time snaps of the interaction of the water drop with the uncoated elastic membrane. As evident from Fig. 2, the drop detaches from the needle but spreads on the elastic membrane. It is to be noted that the retraction of the needle with a finite speed along with the drop attached at its end imparts energy to the membrane during the contact of the drop with the membrane causing a finite deflection in the membrane. Together with the inherent background stress, due to the special holding arrangement of the elastic membrane, the membrane vibrates post impact (see Supplementary Video S1). It is evident from this video and Fig. 2(c) that in the case of an uncoated elastic membrane, the drop does not detach upon impact. Rather, it spreads on the elastic membrane. This is due to the more dominant drop-membrane adhesion force by the virtue of its hydrophilic nature that dominates over the elastic response of the membrane, hence the cause for drop spreading. Therefore, to obtain the detachment of the drop from the membrane, thereby achieving the desired ‘needle-free’ drop deposition, firstly the wettability of the membrane should be such that it opposes spreading and ensures minimal drop-membrane adhesion, i.e., a hydrophobic/superhydrophobic one. Secondly, the elasticity of the membrane should be such that the post impact elastic response dominates over drop-membrane adhesion and aids drop detachment. Hence, we present both the scenarios: we changed the wettability of the membrane significantly - the hydrophilic membrane was converted into a superhydrophobic membrane using the specialized spray, discussed earlier; Secondly, we varied the elasticity of the membrane by using a different material for the deformable membrane, viz., Polydimethylsiloxane (PDMS), details of which are discussed later.

To restrict the spreading of the drop on the elastic membrane, first we have applied superhydrophobic coating on the elastic membrane and then the coated membrane is attached to the holder in a similar manner as done for the uncoated elastic membrane case. Here, all other experimental parameters are kept the same as shown for Fig. 2. Hence, as expected, it was found that the superhydrophobicity (low energy) of the membrane reduces force of adhesion between the drop and the elastic membrane,

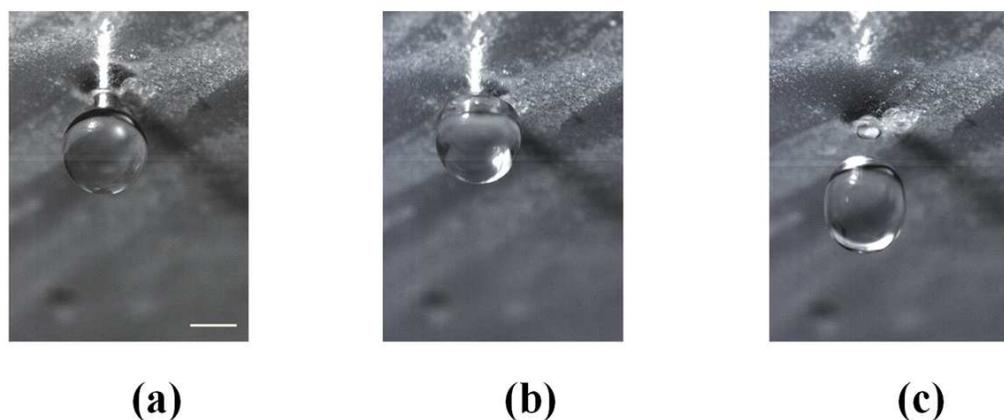


Fig. 3 Needle-drop interaction with a coated superhydrophobic elastic membrane (plastic wrap) ($\theta \approx 165^\circ$) and needle retraction speed of 30mm/s (a) Drop before impacting the elastic membrane (b) Drop impact on the elastic membrane (c) The deflection of membrane due to drop impact and the needle-free drop descending due to gravity. The scale-bar at the bottom of Fig. 3(a) represents 1mm.

thereby preventing spreading, and together with the favorable elastic response of the membrane, a successful drop detachment is achieved. This ‘needle-free’ drop descends due to gravity and subsequently gets deposited on the desired characterizing substrate, kept at the bottom of this elastic membrane (Please refer Supplementary Video S2). Figure 3 shows the snapshots of the needle-free drop deposition using the coated superhydrophobic plastic membrane. As mentioned earlier, the second approach to detach the drop would be to change the elastic energy of the membrane. To achieve this, we replace the elastic membrane used earlier (i.e., the plastic wrap) by $\sim 1\text{mm}$ thick PDMS membrane. The PDMS membrane was produced using the spin-coating technique and appropriate operating parameters³² were used to obtain the desired thickness of $\sim 1\text{mm}$ to allow the flexible property of the PDMS membrane. We observed that a bare PDMS membrane is hydrophobic with equilibrium contact angle $\theta = 111.8^\circ \pm 3.7$. It should be noted here that contact angle measurement on soft PDMS is challenging since the applicability of Young’s law for soft solids is debatable due to deformation of such surface at the three phase contact line³⁴. Upon retracting the drop-needle assembly through it, it was found that the bare(uncoated) PDMS membrane proved to be successful in detaching the drop by the virtue of its significant elasticity, as depicted in Fig. 4. However, the energy imparted due to its elasticity and vibrating modes are significantly high and hence we observe that the drop does not retain its original spherical shape (as opposed to the case shown for coated plastic wrap membrane in Fig. 3 where the drop retains its spherical shape), but rather breaks up into number of satellite drops. It is to be noted that such a high magnitude of elastic energy, as observed here, is similar to effects (i.e., flattening of the drop, recoiling, splashing and break-up³⁵) that we witness during drop-impact on a solid substrate with a very high impact velocity. Recoiling and break-up of the drop can be observed in Figs. 4 (c)-(d). The wetting property of the membrane allows the drop to spread on the membrane and we speculate that this spreading of the drop further initiates such unwanted phenomena (See Supplementary Video S3). It was also observed that the certain

amount of liquid volume always remains attached to the membrane. Therefore, to avoid these undesirable effects and also to achieve the detachment of an intact drop from membrane (as oppose to satellite drop formation), it is necessary to restrict spreading of the drop on the membrane. Therefore, we have changed the wetting properties of the PDMS membrane in a similar way in which we had earlier changed the wetting properties of the plastic wrap (elastic membrane used in Fig. 3). We noticed that when we apply coating to the PDMS surface, due to the superhydrophobic nature of the coated PDMS membrane (contact angle $\theta \approx 165^\circ$), the entire drop gets detached from the membrane and the drop break-up was arrested (See supplementary VideoS4). It should be noted here that post impact, the drop, though undergoes detachment, it does not retain its original spherical shape changing from a donut like shape to a more elongated spherical shape at later times. Therefore, we have demonstrated that by tuning the elastic or wetting properties of elastic membrane, one can easily obtain the needle-free drop deposition In the upcoming section, we have provided theoretical analysis to this technique in terms of the wetting and elastic properties of the membrane to achieve the desired ‘needle-free’ drop deposition.

4 Theoretical Analysis

To explain the four different scenarios observed here, a theoretical analysis is presented here. For ease of analysis, we have applied conservation of energy at three different stages of the process: firstly, before the impact of the drop with the elastic membrane, secondly at the moment of maximum deflection of the membrane due to impact and finally, after the impact where the drop detaches or spreads on the membrane. For all these stages, the system in consideration is the drop and the elastic membrane. The drop attached to the needle, before impact, corresponds to a spherical non-deformed drop of radius R , retracting with a velocity V_i . The total energy of the system at this initial stage can be expressed as,

$$K.E._1 + S.E._1 = \frac{2}{3}\pi R^3 \rho V_i^2 + 4\pi\gamma_{da}R^2 \quad (1)$$

where ρ (1000 kg/m³) and γ_{da} (0.072 N/m) are the density and the surface tension of the liquid drop, respectively. At stage 2, where the membrane is at its maximum deflection δ , the kinetic energy of the drop gets stored in as elastic energy of the membrane while its surface energy depends on its configuration at that point. It should be noted that the drop in stage 2 is in contact with the membrane and tends to spread due to inertia. There are two contributions to the elastic energy stored in the membrane: one due to the inherent elastic nature of the membrane (E_{els}^m) i.e., its own elasticity while the background stress (σ_b) due to the tension in the membrane is the other contribution (E_{σ}^m). The energy of the drop-membrane system at this stage can be expressed as,

$$K.E._2 + S.E._2 = E_{els}^m + E_{\sigma}^m + S.E._2 = c_1 E_f h_f l^2 \left(\frac{\delta}{l}\right)^4 + c_2 \sigma_b h_f l^2 \left(\frac{\delta}{l}\right)^2 + \gamma_{da} A_{da} + \gamma_{dm} A_{dm} + E_{diss}$$

Here, E_f is the elasticity of the membrane. h_f , δ and l are respectively the membrane thickness, its maximum deformation and the length over which the deformation is significant (see Table. 1 for detailed values of these parameters). The background stress can be estimated as $\sigma_b \sim M_{ring} g / 2\pi R_f h_f$, where M_{ring} (0.6g) is the mass of the ring and R_f (15mm) is the radius of the membrane. In Eq. 2, A_{da} and A_{dm} are the drop-air and drop-membrane surface area, respectively, and γ_{da} and γ_{dm} are the respective interfacial tensions. E_{diss} is the viscous dissipation, which is negligible here. In Eq. 2, c_1 and c_2 are numerical constants. Unfortunately, the exact magnitudes of c_1 and c_2 are unknown due to lack of knowledge on nonlinear elastic properties of these materials. This stage could be crucial for the eventual outcome of the process. Since most of the stored elastic energy in the membrane will be eventually transferred back to the drop as kinetic energy, whether the drop will detach or not will depend on the competition between the stored elastic energy and surface energy of the drop at this stage. Hence, for $E_{els}^m + E_{\sigma}^m > \gamma_{da} A_{da} + \gamma_{dm} A_{dm}$, drop detachment will be favored while for the opposite case, the drop will spread further on the membrane. Hence, as mentioned previously, depending on the wetting property and elastic nature of the membrane, four probable scenarios arise. For a hydrophilic plastic membrane, the drop completely spreads on the membrane due to dominant drop-membrane interfacial energy compared to the elastic response of the membrane, and together with the deformed membrane performs oscillations with diminishing amplitude. Hence, majority of the total energy of the system is now in the form of surface energy of the drop in its final spreading configuration, while a certain portion of the total energy is lost due to the vibrations (E_{vib}). The energy for stage 3 corresponding to this outcome can be written as,

$$E_{3,I} = S.E._3 + E_{vib} = \left[\gamma_{da} \pi (R_0^2 + h_0^2) + \gamma_{dm} \pi R_0^2 \right] + \rho_f h_f l^2 (\delta/\tau)^2 \quad (3)$$

where ρ_f is the density of the membrane and τ is the contact time of impact. R_0 and h_0 are respectively the radius and height of the drop in its maximum spreading configuration (see Fig.

2(c)). As reported by Courbin et al.³⁰, for wide range of impact velocities (0 to 1.4m/s), the contact time can be written as $\tau \approx \delta/0.2V_i$. For the superhydrophobic plastic membrane, the drop does not spread on the membrane by the virtue of unfavorable drop-membrane surface energy and aided by the more dominant elastic response of the membrane, the drop smoothly detaches (see Supplementary Video S3). For this case, the drop practically regains most of its initial kinetic energy, while a small amount of energy is lost due to vibrations of the membrane. The energy for stage 3 corresponding to this is,

$$E_{3,II} = K.E._3 + E_{vib} = \frac{2}{3} \pi R^3 \rho V_f^2 + \rho_f h_f l^2 (\delta/\tau)^2 \quad (4)$$

where V_f is slightly less than V_i .

The situation is different for the more flexible PDMS membrane. For the un-coated hydrophobic one, even though the drop spreads initially due to inertia, but the elastic response of the (2) more flexible PDMS membrane is such that vibrations of larger amplitudes eventually breaks up the drop into satellite drops. This can be attributed to the fact that for the more elastic PDMS membrane, less energy is available to the drop for spreading and there is more residual kinetic energy in the membrane, resulting in oscillations of larger amplitudes. So, in addition to stage 3, there is another stage where the drop in its maximum spreading configuration deforms and actually breaks up, and finally falls due to gravity. Also, for the superhydrophobic (coated) PDMS membrane, the drop detaches from the membrane, albeit with a deformed shape, and falls due to gravity. It should be noted that for all the scenarios discussed above, the energies E_1 , E_2 and E_3 are same order of magnitude by the virtue of elastic collision. Unfortunately, a more detailed energy analysis is beyond the scope of the present work. On the other hand, a force analysis of the drop impact and subsequent detachment can also be performed to account for the four scenarios described above. In that regard, the competition between the force generated due to the drop impact and the drop-surface adhesion force^{1,36} will dictate the eventual outcome of the process. The authors would like to point out that drop impact and break up is a field in its own, while in the present work, the focus is entirely on needle-free drop deposition. Compared to our previous study²⁷, using a rigid hydrophobic top substrate to facilitate drop detachment, drop deposition using an elastic membrane minimizes drop deformation and energy loss upon impact and subsequent detachment due to perfectly elastic interaction, and can be considered as a lower bound of any such post impact drop deformation.

It should be noted here that the retraction velocity is important for successful drop detachment upon impact with the elastic membrane. The needle retraction velocity of 30 mm/s corresponds to a Weber number ($We = \rho V_i^2 R / \gamma$) much lower than unity, which is perfectly suited for successful drop detachment upon impact. For $We > 1$, unwanted phenomenon, such as splashing upon impact or self-detachment from the needle during retraction is expected, as reported in our previous study²⁷. On the other hand, a very low retraction velocity will not produce the sufficient elastic response from the membrane to enable drop detachment. Instead, for such a scenario, the drop will adhere to the membrane

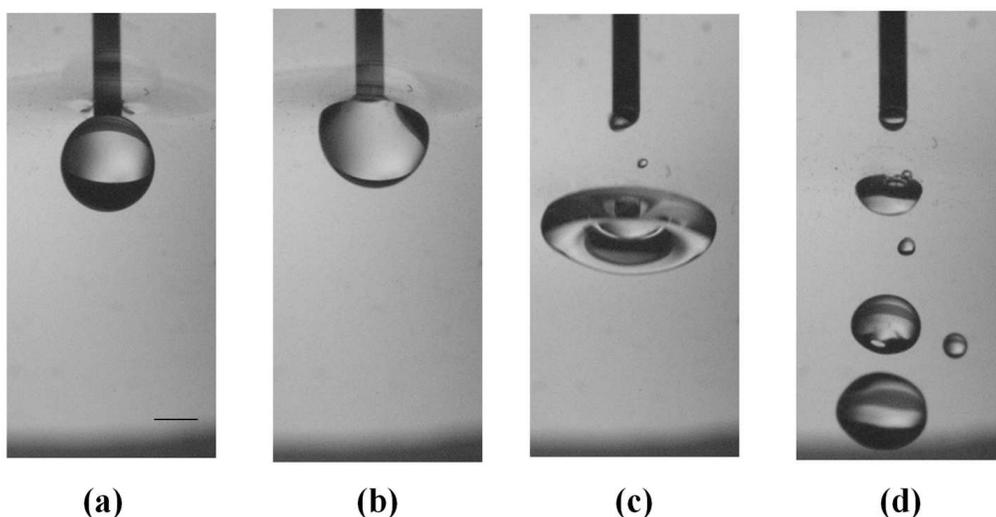


Fig. 4 Needle-drop interaction with a bare PDMS membrane ($\theta \approx 111.8^\circ$) and needle retraction speed of 30mm/s. (a) Drop before impacting the PDMS membrane. (b) Maximum spreading of the drop after impact. (c) Recoiling of the drop and the maximum deflection in the membrane during the vibration in the membrane. (d) Breakup of the drop into number of satellite drops due to multiple vibrations in the membrane. The scale-bar at the bottom of Fig. 4(a) represents 1mm.

surface.

Table 1 Different parameters of plastic and PDMS membranes

Parameters	Plastic Wrap	PDMS
ρ	1190kg·m ⁻³	965kg·m ⁻³
E_f	2 GPa	720 kPa
h_f	20 μ m	1mm
δ	2mm	5mm
l	2mm	5mm
σ_b	3.12kPa	0.624kPa
τ	0.133sec	0.665sec

5 Conclusions

In this work, a ‘needle-free’ drop deposition in air medium is obtained that is independent of the mechanical, geometrical and wetting properties of the needle and the characterizing substrate. This is achieved by bringing in an additional flexible substrate (soft elastic membrane) between the needle-drop assembly and the bottom characterizing substrate. The ‘needle-free’ technique works when the needle-drop assembly is retracted upward at a finite speed, away from the characterizing substrate, and the interaction of the drop with the flexible membrane ensures negligible energy loss and drop deformation upon impact thereby allowing the drop to be detached from the needle and settle on the bottom characterizing substrate by gravity for accurate contact angle measurement of the characterizing substrate. We have considered two different kinds of membrane substrates - plastic wrap and PDMS membrane with varying elastic property. We also varied the wetting characteristics of each substrate by providing specialized coating to make them superhydrophobic surfaces. It was found that in case of plastic wrap, whose elasticity is significantly smaller than the PDMS membrane, the elasticity of such surface alone is not good enough to detach the drop from the needle. Hence, for such low elastic energy membranes, one need to alter the wetting characteristics of the elastic membrane so that

the drop spreading on to such surfaces can be minimized and thereby achieving the desired ‘needle-free’ drop deposition in air. On the other hand, large elasticity of the PDMS membrane creates unwanted drop splashing, break-up, etc. which to a certain extent can be arrested by again altering the wetting characteristics of such highly deformable membrane. Moreover, in our previous study²⁷, we observed that the failure of the needle-free drop deposition for drop volume less than 3 μ L corresponding to the needle diameter of 0.5mm, which is of same dimension as that used in our present study. This is because, in case of a rigid solid substrate, there exists a finite gap between the needle and the inner surface of the drilled hole, which allows capillary imbibition of drops (smaller than 3 μ L volume) to occur inside the gap. Such problem can be avoided in case of an elastic membrane, where there exists practically negligible gap between the needle and the flexible elastic membrane. Hence, the proposed technique eliminates the dependence of the drop diameter, to be used for characterizing a given substrate, on the needle diameter, which can be of great advantage to studies related to drop impact. Therefore, the ‘needle-free’ drop deposition technique, presented here, has an universal appeal and not only helps in accurately determining the contact angle of low-energy surfaces (superhydrophobic, superoleophobic, etc.) but at the same time opens up a new avenue to revisit the drop impact studies with controlled drop diameters and over a wide range of impact Weber number.

6 Acknowledgment

The financial assistance from NSERC, through Grant No. RGPIN-2014-05236, is acknowledged here. The authors would like to thank Aleksey Baldygin, Ph.D. student in Department of Mechanical Engineering, University of Alberta, for his help towards designing and setting up the experimental system. The authors would also like to thank Dr. Siddhartha Das, who was an Assistant Professor in Department of Mechanical Engineering at the

University of Alberta, and currently at the University of Maryland, for fruitful discussions regarding drop impact study.

References

- 1 P. R. Waghmare, S. Das and S. K. Mitra, *Soft Matter*, 2013, **9**, 7437–7447.
- 2 D. Vadhilo, A. Soucemarianadin, C. Delattre and D. Roux, *Physics of Fluids (1994-present)*, 2009, **21**, 122002.
- 3 X. Deng, L. Mammen, H.-J. Butt and D. Vollmer, *Science*, 2012, **335**, 67–70.
- 4 O. E. Yildirim, Q. Xu and O. A. Basaran, *Physics of Fluids*, 2005, **17**, 062107.
- 5 B.-B. Lee, P. Ravindra and E.-S. Chan, *Chemical Engineering Communications*, 2008, **195**, 889–924.
- 6 X. Deng, F. Schellenberger, P. Papadopoulos, D. Vollmer and H.-J. Butt, *Langmuir*, 2013, **29**, 7847–7856.
- 7 P. Tsai, S. Pacheco, C. Pirat, L. Lefferts and D. Lohse, *Langmuir*, 2009, **25**, 12293–12298.
- 8 D. Richard, C. Clanet and D. Quéré, *Nature*, 2002, **417**, 811–811.
- 9 B. Qian, M. Loureiro, D. A. Gagnon, A. Tripathi and K. S. Breuer, *Physical Review Letters*, 2009, **102**, 164502.
- 10 C. Antonini, A. Amirfazli and M. Marengo, *Physics of Fluids*, 2012, **24**, 102104.
- 11 B.-B. Lee, P. Ravindra and E.-S. Chan, *Colloids and Surfaces A: Physicochemical and Engineering Aspects*, 2009, **332**, 112–120.
- 12 A. Tuteja, W. Choi, M. Ma, J. M. Mabry, S. A. Mazzella, G. C. Rutledge, G. H. McKinley and R. E. Cohen, *Science*, 2007, **318**, 1618–1622.
- 13 D. Bartolo, C. Josserand and D. Bonn, *Journal of Fluid Mechanics*, 2005, **545**, 329–338.
- 14 J. C. Bird, R. Dhiman, H.-M. Kwon and K. K. Varanasi, *Nature*, 2013, **503**, 385–388.
- 15 A. Yarin, *Annu. Rev. Fluid Mech.*, 2006, **38**, 159–192.
- 16 D. Richard and D. Quéré, *EPL (Europhysics Letters)*, 2000, **50**, 769.
- 17 M. Pasandideh-Fard, Y. Qiao, S. Chandra and J. Mostaghimi, *Physics of Fluids*, 1996, **8**, 650–659.
- 18 D. Roux and J. Cooper-White, *Journal of Colloid and Interface Science*, 2004, **277**, 424–436.
- 19 Y. C. Jung and B. Bhushan, *Langmuir*, 2008, **24**, 6262–6269.
- 20 K. Okumura, F. Chevy, D. Richard, D. Quéré and C. Clanet, *EPL (Europhysics Letters)*, 2003, **62**, 237.
- 21 J. Zou, P. F. Wang, T. R. Zhang, X. Fu and X. Ruan, *Physics of Fluids*, 2011, **23**, 044101.
- 22 E. Atefi, J. A. Mann Jr and H. Tavana, *Langmuir*, 2013, **29**, 5677–5688.
- 23 R. Dufour, M. Harnois, V. Thomy, R. Boukherroub and V. Senez, *Soft Matter*, 2011, **7**, 9380–9387.
- 24 C. Fuentes, K. Beckers, H. Pfeiffer, L. Tran, C. Dupont-Gillain, I. Verpoest and A. Van Vuure, *Colloids and Surfaces A: Physicochemical and Engineering Aspects*, 2014, **455**, 164–173.
- 25 L. Gao and T. J. McCarthy, *Langmuir*, 2009, **25**, 14105–14115.
- 26 L. Jenkins and A. Donald, *Langmuir*, 1999, **15**, 7829–7835.
- 27 P. R. Waghmare and S. K. Mitra, *Journal of Applied Physics*, 2014, **116**, 114903.
- 28 P.-G. De Gennes, F. Brochard-Wyart and D. Quéré, *Capillarity and wetting phenomena: drops, bubbles, pearls, waves*, Springer Science & Business Media, 2013.
- 29 R. E. Pepper, L. Courbin and H. A. Stone, *Physics of Fluids (1994-present)*, 2008, **20**, 082103.
- 30 L. Courbin, A. Marchand, A. Vaziri, A. Ajdari and H. A. Stone, *Physical Review Letters*, 2006, **97**, 244301.
- 31 E. Cerda, K. Ravi-Chandar and L. Mahadevan, *Nature*, 2002, **419**, 579–580.
- 32 M. Liu, J. Sun, Y. Sun, C. Bock and Q. Chen, *Journal of Micromechanics and Microengineering*, 2009, **19**, 035028.
- 33 R. Tadmor, *Langmuir*, 2004, **20**, 7659–7664.
- 34 A. Marchand, S. Das, J. H. Snoeijer and B. Andreotti, *Physical Review Letters*, 2012, **109**, 236101.
- 35 R. Rioboo, C. Tropea and M. Marengo, *Atomization and Sprays*, 2001, **11**, 155–165.
- 36 R. Tadmor, P. Bahadur, A. Leh, H. E. N'guessan, R. Jaini and L. Dang, *Physical Review Letters*, 2009, **103**, 266101.