

# RSC Advances



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

*Accepted Manuscripts* are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. This *Accepted Manuscript* will be replaced by the edited, formatted and paginated article as soon as this is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.

Cite this: DOI: 10.1039/c0xx00000x

www.rsc.org/xxxxxx

ARTICLE TYPE

# A Review of Structure of Oxide Glasses by Raman Spectroscopy

Avadhesh Kumar Yadav,<sup>\*a</sup> and Prabhakar Singh<sup>a</sup><sup>a</sup>Department of Physics, Indian Institute of Technology (Banaras Hindu University), Varanasi-221005 (India)<sup>\*</sup>E-mail: yadav.av11@gmail.com

The family of oxide glasses is very wide and it is continuously developing. The rapid development of advanced and innovative glasses is under progress. The oxide glasses have variety of applications in the daily use articles as well as advanced technological fields like protection of X-ray, fibre, optical instruments or lab glasswares. The oxide glasses are basically consisted by network formers such as borate, silicate, phosphate, borosilicate, borophosphate and network modifiers like alkali, alkaline earth and transition metals. In present review article, Raman spectroscopic results for structures of borate, silicate, phosphate, borosilicate, borophosphate, aluminosilicate, phosphosilicate, alumino-borosilicate and tellurite glasses are summarized.

## 1. Introduction

### 1.1 Brief history of glass

The exact history of glass is not known but there are certain evidences which confirm that it is used about 3500 BCE in Mesopotamia. The available archaeological evidence gives the signature of presence glass articles in coastal north Syria, Mesopotamia or Ancient Egypt. The glass technology in South Asia was started with 1730 BCE [1]. The evidences of glass have also been found in Hastinapur and archaeological site in Takshashila, ancient India. The glass objects have been developed in Romanian Empire and Anglo-Saxon period in domestic, industrial and funerary contexts. In china, the development of glass was started late but nowadays, glass played a peripheral role in arts and crafts. The glasses from that period are made of barium oxide (BaO) and lead [2]. The tradition of use of the lead-barium glasses was disappeared at the end of the Han Dynasty (AD 220). The soda glasses with high alumina are rarely used around the Mediterranean area or in the Middle East. The soda-lime, borosilicate, glass fibre, lead, alkali-barium silicate, aluminosilicate, vitreous silica and pyrex glasses were investigated time to time. The large scale development glass technology was started mainly from 14<sup>th</sup> century [3]. The recent research on glasses are based on advanced glasses used for safety systems like X-ray protective glasses, fibre glasses, high refractive index glasses for optical instruments, photo-chromatic and high strength glasses.

Last few decades, the oxide glasses have attracted the attention of researchers and scientist due to their extensive properties which are very useful for many applications. They are easily shaped due to lack of underlying crystal structure. The covers of solar units and photovoltaic units improve the performance. In turbines, glass fibers reinforced composites are used for the purpose of storage of wind energy. For this purpose, boron-free glasses are commonly used [4]. The oxide glasses may be used in smaller power supplies, including dielectrics for super-capacitors,

sealants for high-temperature solid oxide fuel cells (SOFC) and electrolytes for electrochemical devices [5]. W. H. Zachariassen had investigated the silicate glasses and he found that such glasses are consist of silicon tetrahedrally coordinated to 4 oxygen atoms [6]. The phosphate glasses are investigated P. B. Price et al [7]. Borate glasses have been investigated in various glass systems B<sub>2</sub>O<sub>3</sub>-xLi<sub>2</sub>O<sub>3</sub> [8], xB<sub>2</sub>O<sub>3</sub>-(1-x)Li<sub>2</sub>O<sub>3</sub> [9-10], (30 - x)Li<sub>2</sub>O - xK<sub>2</sub>O - 10CdO/ZnO - 59B<sub>2</sub>O<sub>3</sub> [11], 60B<sub>2</sub>O<sub>3</sub>-(20-x)Na<sub>2</sub>O-10PbO-10Al<sub>2</sub>O<sub>3</sub>:xTiO<sub>2</sub>:yNd<sub>2</sub>O<sub>3</sub> [12], PbO-Bi<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub> [13], Bi<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-ZnO-Li<sub>2</sub>O [14], Gd<sub>2</sub>O<sub>3</sub>-MoO<sub>3</sub>-B<sub>2</sub>O<sub>3</sub> [15]. The silicate glasses were investigated in presence of Na<sub>2</sub>O, Li<sub>2</sub>O, K<sub>2</sub>O and CaO by various research groups [16-25]. The phosphate glasses are investigated for long time in many systems. The calcium, sodium, potassium, iron, gadolinium, zinc, lead, chromium, molybdenum, lanthanum, bismuth, tungsten, samarium, copper, cadmium, barium, silver contained phosphate glasses have been investigated for various application purpose [122-165]. Soon after single network former glasses, double network former glasses have also been investigated for composite network structure of borate, silicate and phosphate which are named as borosilicate, borophosphate etc glasses. In recent years, the structure properties of SnF<sub>2</sub>-SnO-P<sub>2</sub>O<sub>5</sub> glasses have been studied. In this study, the glass transition temperature was found to be very low [26]. CeO<sub>2</sub> in cerium iron borophosphate glasses increases the glass transition temperature and also increases thermal stability [27]. The thermal stability of Na<sub>2</sub>O-FeO-Fe<sub>2</sub>O<sub>3</sub>-P<sub>2</sub>O<sub>5</sub> was found to be poorer than sodium-free iron phosphate glasses [28]. TiO<sub>2</sub> is useful for increasing the thermal stability of BaO-Li<sub>2</sub>O-diborate glasses [29]. The low melting glasses were prepared in the glass system K<sub>2</sub>O-MgO-Al<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> [30]. For the purpose of immobilisation of nuclear waste, few borosilicate glass were also investigated in glass system 55 SiO<sub>2</sub>-15 B<sub>2</sub>O<sub>3</sub>-5 Al<sub>2</sub>O<sub>3</sub>-5 CaO-(20-x)Na<sub>2</sub>O-xCs<sub>2</sub>O [31].

Glass may be defined as inorganic solids, which are amorphous in nature. They are classified in several groups like

metallic alloys, ionic melts, aqueous solutions, molecular liquids and polymers. They have attracted the attention of researchers and scientists in recent years due to demand in many significant applications like fields of energy science and photonics. The substantial growth of glasses is still required.

## 1.2 Brief Chemistry and Physics of Glasses

Glasses are disordered materials i.e. lack of the periodicity in crystal structure [32] and are formed by cooling of melt of inorganic products to rigid condition without crystallization [33-34]. Several materials like organic polymers and metal alloys have also amorphous structures but all of them are not glasses [35]. The formation of glass requires critical cooling rates [36-39]. Thermodynamically, the glasses are non-equilibrium materials because their properties are the function of pressure, temperature and composition. They are always approached to its nearby metastable state [40-43]. Glass formation is possible if glass composition contains network formers. They are formed by the network formers such as borate, silicate, phosphate, borosilicate, borophosphate etc. It is found that the excellent glass formers have electro-negativity values of 1.7–2.1 on the Pauling scale [44]. The sufficient network modifiers are used in glass formation to modify the glass properties. The glass network is modified by alkali and alkaline earth metal modifiers like ZnO, PbO, TeO<sub>2</sub>, Bi<sub>2</sub>O<sub>3</sub>, MgO, CaO, SrO, BaO, CeO<sub>2</sub>, Cr<sub>2</sub>O<sub>3</sub>, La<sub>2</sub>O<sub>3</sub>, CdO, CuO, Li<sub>2</sub>O, Na<sub>2</sub>O, K<sub>2</sub>O, Al<sub>2</sub>O<sub>3</sub>, MgF<sub>2</sub> etc. The alkali ions are easily thermally activated which can move from one site to another within a glass. This type of movement of alkali ions within a glass structure enables to replace alkali ions near the surface of glass by other ions of the same valence [45-46]. Several local packing arrangements are available in glasses which enable the tuning of their properties through composition and processing. This permits creation of highly homogeneous materials on a macroscopic scale [47]. Zachariasen establish a theory which explains the criteria for vitreous structures with highly directed three-dimensional bond arrays. The glass structures are comprised of microcrystallites ~ 20 Å in size. The glass structures lie between a melt state and a glassy state [48-49]. The glassy network structure may be effectively understood by infrared and Raman spectroscopy. Certain viscosity models are used to explain the glass properties on temperature scale. The five temperatures named as standard points (associated with the viscous flow of glass), strain point (maximum temperature supported by glass), annealing point (relieve the internal stresses), softening point and working point (most glass forming operation acting at this temperature) [50].

### 1.2.1 Classification of Preparation Method of Glasses

The glasses are mainly synthesized by melt-quenched method, sol-gel method and chemical vapour deposition [51].

#### 1.2.1.1 Glass by Melt-Quench Technique

The melt-quench technique is the first discovered technique for preparation of glasses. In melt quench method, the inorganic raw materials are properly weighed according to composition. Then, the raw materials are mixed properly in mortar with pestle or ball mill in acetone media. The well mixed and dried ingredients are placed into high grade alumina or platinum crucible and crucible is placed in programmable temperature controlled furnace for melting. The homogenous melt is quenched in aluminium mould. In this way, the glass is formed by melt quench method. Formed

glasses are annealed at relatively low temperature to the glass casting temperature to remove the residual internal stresses which are produced due to temperature gradient during forming and subsequent cooling [52]. The melt quench method is advantageous for obtaining materials of large size compared to single crystal or polycrystalline ceramics. It has one more advantage that it provides large flexibility of composition over chemical vapour deposition or sol-gel method. This method has also certain disadvantages. For example, this method is not good for preparation of glasses of ultra-high purity which are used in optical communication. Certain contamination and impurity of crucible or furnace materials is also added in the glass in this method. Glasses, in which refractory materials like SiO<sub>2</sub>, TiO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub>, ZrO<sub>2</sub> etc are used, are difficult to be prepared by this method due to requirement of extremely high temperature.

#### 1.2.1.2 Glass by Sol-gel Method

In sol-gel method, the colloidal solution (sol) of raw materials is converted into gelatinous substance (gel). In this method, there is no requirement of any melting and quenching. The fine-grained gel bulks of various organic compounds are annealed at particular temperature and thereafter the annealed gel was sintered at higher temperature than the annealing temperature. During this process, there are two reactions take place, named as hydrolysis and condensation reaction [53].

#### 1.2.1.3 Glass by Chemical Vapour Deposition Method

The chemical vapour deposition (CVD) method for preparation of bulk glasses was developed in the early 1940. For the formation of glass material is followed by thermally activated homogeneous oxidation or hydrolysis of the initial metal halide vapour. The oxidation or hydrolysis reaction is activated by an oxy-hydrogen flame or oxygen plasma. In this method, the initial components are liquid at room temperature and their boiling temperature is very low compared with halides of alkali, alkaline earth, transition metals or rare-earth elements. The purification of raw materials is done by repeating the distillation below melting points. This method is advantageous for preparation of ultra-high purity glasses [54].

The structural study of glass networks is important for understanding the composition dependence of oxide glass properties. Various models are proposed for correlating the different properties with structure and composition of glassy system. Phillips and Thorpe investigated the compositional dependence of glass properties by analyzing the topology of the glassy network. This approach is based on the comparison of the number of atomic degrees of freedom with the number of interatomic force field constraints. The tendency for glass formation would be maximized if the number of degrees of freedom exactly equals the number of constraints in the glassy network [55]. Gupta and Mauro have demonstrated generalization of the Phillips and Thorpe approach in their study [56, 253].

This review article is focused on the structural study of various oxide glasses and how the structure of glassy network is modified by various dopants like alkali, alkaline and transition metals. The compositional dependent study of oxide glasses are also summarized in present review article. This review article is very useful for understanding the formation of glassy networks and their modification by various dopants.

## 1.3 Brief Introduction to Glassy Networks

A lot of work has been done on glassy materials till now. The glassy materials consist of various network former like borate, silicate, phosphate, borosilicate, borophosphate etc. The borate network consists of vibrations of isolated diborate ( $B_2O_3$ ) groups, breathing of O-atoms, loose diborate vibrations, chain or ring type meta- and pentaborate ( $B_5O_8$ ) groups, symmetric breathing vibrations of six member borate rings with one or two  $BO_3$  triangle replaced by a  $(BO_4)$  tetrahedral, symmetric breathing vibrations of boroxol rings, triborate ( $B_3O_5$ ), terta-borate ( $B_8O_{13}$ ), ortho-borate ( $BO_3$ ), pyroborate ( $B_2O_5$ ) groups, linked  $[BO_3]^{3-}$  triangles with one free vertex etc [45, 57]. The borate glasses have significant importance in several applications in thin amorphous films for battery application, bioactive glasses for tissue engineering, nuclear waste disposals, photonic applications, development of tuneable or short pulse lasers, optical fibre amplifiers and fibre lasers [58-61]. The borate networks are shown in Fig. 1.

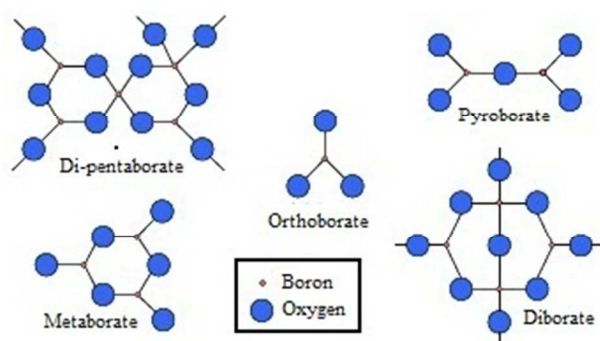


Fig. 1: The network units of borate glasses [107]. (Reproduced with permission of IOP Publishing)

The silicate network former is also significant for glass formation but the silica is not easily fusible and has not retained viscosity for long time. Thus, it cannot form the glass easily. Certain substances are used which help in the formation of glassy network, called fluxes [62]. The silicate glasses consist of tetrahedral 6-coordinated 4 oxygen atoms. 6-coordinated silicon is also found in few crystalline materials [63]. The silicate networks are found in the form of structural units of orthosilicate, Si - O stretching vibrations of tetrahedral silicate units, symmetric stretching vibrations of silicate tetrahedra, inter-tetrahedral Si - O - Si linkages and structural unit  $Q_{Si}^n$  ( $n = 0$  to 4), where,  $Q_{Si}$  represents the tetrahedral unit and  $n$  is the number of bridging oxygen (BO) per tetrahedron [64-66]. Silicate glasses are, generally, used for application of windows, lenses, mirror substrates, crucibles, trays and boats, UV transmitting optics (synthetic fused silica), IR transmitting optics and metrology components. The network units of silicate glasses are shown in Fig. 2. Alike to silicate network, phosphate network is formed by various structural units of phosphate. The phosphate glasses are formed at low temperature in comparison to silicate glasses. The vibrational units of phosphate networks are  $Q_p^n$  groups of  $PO_4$  tetrahedra ( $n =$  number of bridging oxygen per  $PO_4$  tetrahedron), symmetric stretching of P - O, asymmetric stretching vibration of P - O, P - O - P symmetric vibrations, bending vibrations of phosphate polyhedra and O - P - O bending modes. The phosphate glasses have numerous applications in photonics, biomedical applications, solid state electrolytes, fibres etc [67-

69]. The phosphate glasses are much sensitive in ionizing particles. This property of phosphate glasses enables the use for identification of heavy ions [70]. The phosphate networks are shown in Fig. 3.

Not only single network former glasses were performed but also double network former glasses were also investigated. The borosilicate glasses are one of the double network former glasses. The co-existence of borate and silicate networks is associated in borosilicate glasses. The interlinking of some borate and silicate networks is also possible in such glasses. The high frequency band  $> 850 \text{ cm}^{-1}$  are related with silicate networks,  $Q_{Si}^n$ , whereas mid frequency bands ( $400$  to  $850 \text{ cm}^{-1}$ ) are order superstructures, reedmergerite  $[BSi_3O_8]^-$  and danburite  $[B_2Si_2O_8]^{2-}$ . The bands ranges between  $300$ - $500 \text{ cm}^{-1}$  are the characteristic of mixed stretching and bending modes of Si-O-Si units while bands in the range  $550$ - $850 \text{ cm}^{-1}$  are caused by ring breathing modes. The bands at  $670$ ,  $770$  and  $808 \text{ cm}^{-1}$  are the signature of tetraborate groups, four- and three-coordinated boron in diborate and boroxol rings, respectively. The borosilicate glasses are mainly used in lab equipments, high quality medical devices (ampoules, dental cartridges, veterinary tracking devices etc), space exploration devices and electronics (microelectromechanical systems) [71-74].

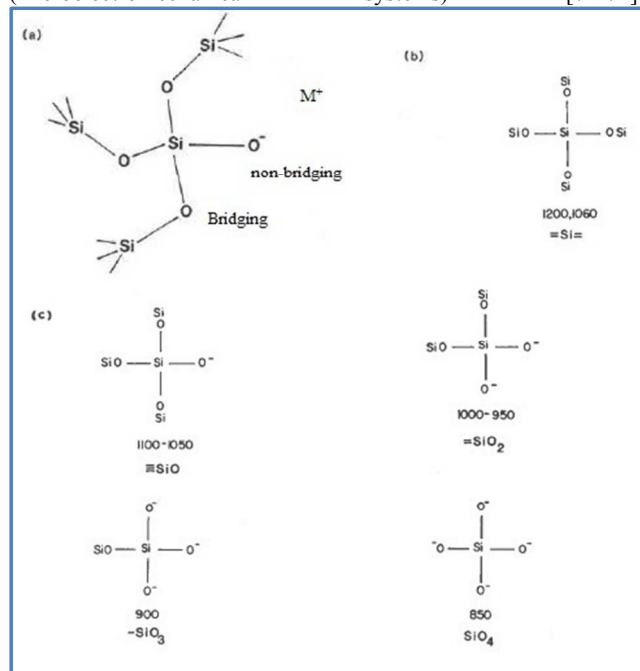


Fig. 2: The network units of silicate glasses [16]. (Reproduced with permission of Mineralogical Society of America)

Borate has an ability to change their properties and control them by a suitable modification of their chemical compositions within a relatively broad concentration region of their constituents in phosphate glasses. Addition of borate ( $B_2O_3$ ) to phosphate glasses improves their mechanical properties and chemical resistance against atmospheric moisture. The borophosphate glasses includes the networks of both borate and phosphate alongwith various linking vibrations. These glasses are very useful for potential applications in solid electrolytes, glass solders and glass seals [75-76]. Phosphosilicate glasses consist of networks of phosphate as well as silicate. The phosphosilicate glasses are used in forming the Raman fibres laser and amplifiers. Phosphosilicate

glasses are useful in developing high refractive index grating [77-78]. Aluminosilicate glasses contain the units of silicate and aluminate networks. These are very important for geological applications due to its highly refractory nature [79-80]. The aluminium based glasses have a great technological importance due to their strong and ductile nature. These play an important role where metallic glasses failed [81]. Almino-borosilicate glasses are formed by the networks of aluminate, borate and silicate. These glasses have wide range of applications in optical, thermal, electrical and mechanical properties. Along with these properties, the addition of  $\text{Al}_2\text{O}_3$  to borosilicate matrix improves the chemical durability, decreases thermal expansion coefficient, electrical even radiation resistant and widens the scope of applications of borosilicate glasses [50]. The doping of alkali and alkaline earth metals in these glasses modified the parent network. Rare-earth ion-doped glasses have importance in application like optical fibre lasers, amplifiers and scintillating glasses.  $\text{Sm}^{2+}$  doped glasses are advantageous for spectral hole burning [83-88].

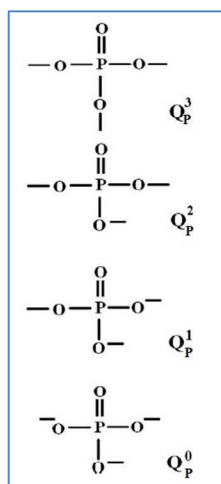


Fig. 3: Network units of phosphate glasses [164]. (Reproduced with permission of Elsevier)

## 2. Principle of Raman spectroscopy

Raman spectroscopy is a powerful tool for structural analysis and developed by Sir C. V. Raman and Kirishnan in 1928 [89]. They pointed in their study that Raman intensity is  $10^{-6}$  to  $10^{-9}$  times less to Rayleigh scattering. Such low intensity can be produced by lasers [90]. When a laser light is incident on vibrating molecules, then, the energy of photon may be changed. The excited molecule or atoms return to different states. The change in energy between the original state and the new state gives shift of emitted photon. The excited molecule goes to higher or lower frequency side to original state and referred as Stoke Raman scattering, (Stoke shift) or anti-Stoke Raman scattering (anti-Stoke shift), respectively [91-92]. The detailed description of working principle is shown in Fig. 4. The chemical composition and structure of molecules influence the modified scattering and no two of the spectra being exactly similar. Thus, Raman shift may be useful in distinguishing structure of different constituents and molecules [93-94].

Raman spectroscopy is used for determination of the structure, environment and dynamics of glassy materials. Furthermore, the

portability of the technique allows for its use in on-line process monitoring over other instruments like Infrared spectroscopy,

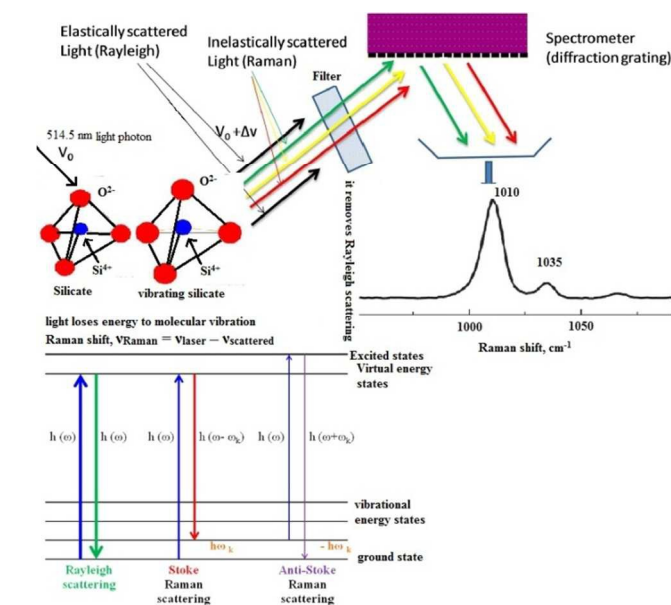


Fig. 4: Working principle of Raman spectroscopy

NMR and X-ray diffraction techniques [95]. Raman and infrared (IR) spectroscopy are used as complimentary techniques and both techniques differ slightly by their selection rule. IR spectra arise from a change in the dipole moment whereas Raman bands arise from a change in the polarizability. Few transitions are allowed in Raman spectroscopy but forbidden in IR spectra, so Raman spectroscopy gives the significant information in that situation. Similar to Raman spectroscopy, in FT-IR spectrometer of UV-is spectrometer, the characteristics of sample are analysed by the absorbance (or transmission) spectrum. In IR spectroscopy, molecules like water or acetone are strong IR absorbers whereas sodium chloride or potassium bromide is very less significant absorbers (window materials). Thus, the functional groups such as hydroxyl or carbonyl are easily detected whereas window materials are not easily detected. However, the situation is different for Raman spectroscopy. In Raman spectra, the ordinate axis normally has arbitrary rather than absorption or %T units because it is simply a measure of the number of scattered photon counts captured by the detector at any particular frequency. If the power of incident laser on the sample is varied, then, the intensity of the Raman spectrum will vary accordingly. Therefore, the peak height in a Raman spectrum is not simply a function of the sample thickness which is different for IR spectrum which depends on sample thickness.

In UV/Visible and IR spectroscopy, the spectral contribution from the instrument is removed by using a reference beam or by subtracting a background spectrum. However, in Raman spectroscopy, any contribution or variation due to the instrument is considered a single-beam is used in it. Due to this reason Raman spectroscopy is play a vital role for quantitative analyses. Thus, the unwanted contributions can be minimized by Raman spectroscopy.

X-ray fluorescence (XRF) is also useful for elemental analysis

of glassy network whereas the vibrational networks cannot understand significantly by XRF. One more thing, a very high energy is required for its operation. X-ray powder diffraction method is also used for understanding the nature of glassy materials but this method is not much useful for network analysis of glassy material [96-98].

### 3. Discussion

#### 3.1. Borate Glasses

A typical Raman spectrum of lithium borate glass is shown in Fig. 5 [99]. Raman study of lithium borate glasses depicts the networks of  $B_{(4)}-O$  vibration of tetraordinated boron at 520.5, 769 and 784  $\text{cm}^{-1}$ , the vibrations of  $B_5O_8(OH)_4$  units at the 555.5  $\text{cm}^{-1}$  and the vibrations of tricoordinated boron near 919  $\text{cm}^{-1}$ . The characteristics intense peaks at 500.5 and 884.5  $\text{cm}^{-1}$  correspond to the vibrations of  $B_{(3)}-O$  in  $B(OH)_3$ . The sol-derived borate glasses show the presence of boroxol ring as band is found at 807  $\text{cm}^{-1}$  [99].

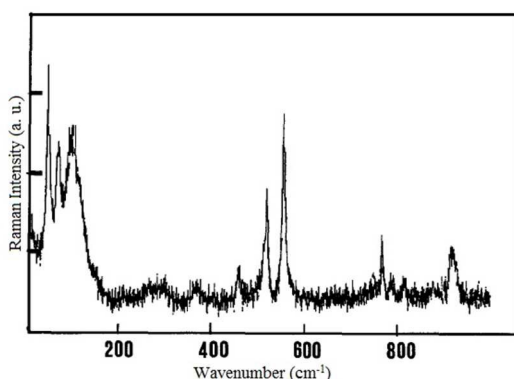


Fig. 5: Raman spectrum of lithium borate glass [99] (Reproduced with permission of Springer)

The Li cations are found in two distinct mode as non-bridging oxygen (NBO) type and bridging oxygen (BO) type sites. The primary network is formed by NBO and BO atoms in tetrahedral borate units, which reside close to NBO atoms. They modify the network by making cluster around NBO atoms [9]. The Raman spectra of  $xLi_2O \cdot (1-x)B_2O_3$  glasses consist of one component at 1030  $\text{cm}^{-1}$  corresponding to B-O bond stretching of  $BO_4$  units in diborate polycrystals and the others near 900, 940 and 1040  $\text{cm}^{-1}$  correspond to the same vibration but in tetraborate units. Increasing concentration of  $Li_2O$  leads to the formation of pentaborate, tetraborate, diborate, metaborate and pyroborate [10]. The spectra of  $Li_2O \cdot 2B_2O_3$  fluxure show the features of crystalline nature with strong characteristic band at 783, 1034 and 1174  $\text{cm}^{-1}$  and some weaker bands between 167 - 723  $\text{cm}^{-1}$ . The spectra of the lithium metaborate ( $Li_2O \cdot B_2O_3$ ) show the clear representation of 545 - 582, 765, 955, 1108 - 1128 and 1467  $\text{cm}^{-1}$ . When the sample is heat treated above melting temperature, then the bands around 640, 671, 724  $\text{cm}^{-1}$  and two less resolved bands in the range 1475 - 1495  $\text{cm}^{-1}$  were also observed [100]. Raman spectra of  $Li_2O \cdot 4B_2O_3$  at room temperature indicate the formation of boroxol rings containing one  $BO_4$  tetrahedron with increasing temperature, the stability of  $BO_3$  tetrahedron decreases and finally, destroyed [101]. The network structure of  $Li_2O-PbO-B_2O_3$  is attributed by  $[BO_{3/2}]^0$ ,  $[B_2O_{4/2}]^-$  and B-O-B bending

mode at 1485, 967 and 765  $\text{cm}^{-1}$ , respectively. The band due to diborate units at 535  $\text{cm}^{-1}$  is completely disappeared by the addition of PbO whereas a few additional bands in region 280 - 300  $\text{cm}^{-1}$  are incorporated. The emergence of peaks in the region of 1200  $\text{cm}^{-1}$  is associated with symmetric stretching of  $B_3$  units and is generated during the incorporation of Pb - O into the structure. The band at 640  $\text{cm}^{-1}$  region was attributed to a bending mode of the Pb - O - B links in the new structure [102]. The Raman spectra in range 150 - 1600  $\text{cm}^{-1}$  of  $(30-x)Li_2O \cdot xK_2O \cdot 10CdO/ZnO \cdot 59B_2O_3$  ( $x = 0, 10, 15, 20, \text{ and } 30$ ) doped with  $1MnO_2/1Fe_2O_3$  showed the four region of bands as: (i) 190 - 600  $\text{cm}^{-1}$ ; (ii) 600 - 820  $\text{cm}^{-1}$ ; (iii) 900 - 1200  $\text{cm}^{-1}$ ; (iv) 1200 - 1600  $\text{cm}^{-1}$ . The strong peaks around 775, 650, 500  $\text{cm}^{-1}$  and a weak peak  $\sim 805$   $\text{cm}^{-1}$  are present, which are attributed to localized breathing motions of oxygen atoms in the boroxol ring. The peaks around 950 and 1110  $\text{cm}^{-1}$  are occurred due to diborate groups in the structure. The peaks in the high frequency region are attributed to  $BO_2O^-$  triangles linked to  $BO_4$  units and  $BO_2O$  triangles linked to other triangular units [103].

The polarized Raman spectra of sodium borate glasses showed the short range structure of  $BO_3$  and  $BO_2O^-$  units. In HH (parallel to the inherent polarisation of the excitation laser) spectra, some strong Raman bands were found in frequency region 700 - 850  $\text{cm}^{-1}$ . The strongest band due to boroxol ring was also observed at 805  $\text{cm}^{-1}$  in the spectrum of glass. A broader and asymmetric band at lower frequency side was assigned with increase of  $Na_2O$ . A band near 1500  $\text{cm}^{-1}$  was found in the spectra and it shifts toward lower frequency side with the increase of  $Na_2O$  concentration whereas VH (perpendicular to the inherent polarisation of the excitation laser) spectra showed drastic difference from HH spectra. VH spectra showed some more bands at 773, 730 and 670  $\text{cm}^{-1}$  [104]. The spectra of  $B_2O_3-Na_2O$  glasses depict five major peaks at 480, 660, 770, 806 and 920  $\text{cm}^{-1}$ . The peaks due to the isolated diborate groups and the ring-type metaborate groups polymerized by  $BO_3$  and  $BO_4$  units remain unaffected with addition of  $Al_2O_3$  whereas peak at 770  $\text{cm}^{-1}$  was shifted towards low frequency side in the spectrum. This means that  $BO_4$  units were consumed and the boroxol ring ( $B_3O_6$ ) $^{3-}$  gradually increased. Another new vibrational peak at about 920  $\text{cm}^{-1}$  was appeared in 6 mol%  $Al_2O_3$  content. This peak is the characteristic of orthoborate type structure containing  $BO_3$  units, which can be transformed by the linking of pentaborate and tetraborate groups. The presence of orthoborate groups and absence of pentaborate groups proves the feasibility of the transformation pentaborate to the orthoborate groups with effect of  $Al_2O_3$  effect on the structure. This suggests that the addition of  $Al_2O_3$  destroyed the high polymeric borate units like pentaborate into the low polymeric groups such as boroxol rings, triborate and orthoborate groups [105]. There are four Raman bands in glassy matrix of  $60B_2O_3 \cdot (20-x)Na_2O \cdot 10PbO \cdot 10Al_2O_3 \cdot xTiO_2 \cdot yNd_2O_3$ . The bands at 755 and 772  $\text{cm}^{-1}$  were assigned due to chain type metaborate groups and the symmetric breathing vibration of six membered rings with one  $BO_4$  tetrahedron, respectively. The position of boroxol ring was found at 797  $\text{cm}^{-1}$  in this system. When  $TiO_2$  concentration was increased, then back conversion of  $BO_4-BO_3$  takes place. The incorporation of  $Ti^{4+}$  led to  $\sim 40\%$  smooth reduction of the  $BO_4$  groups due to the  $BO_4-BO_3$  back conversion effect [12].

The Raman spectrum of lead borate glasses in glass compositions PbO 80%, Bi<sub>2</sub>O<sub>3</sub> 20% and PbO 90%, B<sub>2</sub>O<sub>3</sub> 10% showed a very sharp band at 130 cm<sup>-1</sup> in first composition and shifted to 139 cm<sup>-1</sup> in second composition. Few broad bands centred at 710, 905, 1024, 1239 and 1308 cm<sup>-1</sup> are also observed in the pattern. In 1 mol% MoO<sub>3</sub> doped lead borate glass, two other peaks at 1927 and 3341 cm<sup>-1</sup> are introduced. On increasing the doping concentration of MoO<sub>3</sub> the band near 1255 cm<sup>-1</sup> become weaker. A sharp band at 133 cm<sup>-1</sup> in the 2 mol% Mo-containing lead borate glasses was found. The Raman peak in undoped lead borate glasses is assigned at 139 cm<sup>-1</sup> and this band was located at 144 cm<sup>-1</sup> in crystalline PbO. This suggests the presence of Pb<sup>2+</sup> in high lead borate glasses in the form of PbO<sub>4</sub> groups. The sharp band observed at 305 cm<sup>-1</sup> of lead borate glass containing 5 mol% of MoO<sub>3</sub> is not observed in un-doped lead borate glasses whereas it appeared at 296 cm<sup>-1</sup> in crystalline lead oxide and 282 cm<sup>-1</sup> in crystalline molybdenum trioxide MoO<sub>3</sub>. This can be predicted as Raman band at 305 cm<sup>-1</sup> is originated from different site of Pb<sup>2+</sup>. The weak broad Raman band in MoO<sub>3</sub>-lead borate glasses at 3311 - 3341 cm<sup>-1</sup> are assigned due to molecular water [106]. Sodium-doped lead borate glasses in glass system xNa<sub>2</sub>O·(100 - x)(10PbO·90B<sub>2</sub>O<sub>3</sub>), are made by structural units of boroxol ring resonance and growth of triborate [107]. The spectra of bismuth containing lead borate glasses showed Raman peaks at 400, 550, 710, 920 and 1220 cm<sup>-1</sup>. The peaks for heavy metal oxides are assigned in the range 380 - 580 and 650 - 950 cm<sup>-1</sup> due to the bridging anion modes and the non-bridging anion modes, respectively. The polarization behaviour of bismuth and lead cations have similar due to almost similar atomic weight. First two peaks correspond to the bridge-anion motion due to symmetric stretching of Bi - O - Bi and Pb - O - Pb combined with Bi - O - Pb [108]. The Raman spectra for the xFe<sub>2</sub>O<sub>3</sub>·(100 - x)[3B<sub>2</sub>O<sub>3</sub>·0.7PbO·0.3Ag<sub>2</sub>O] glasses (Fig. 6) confirm the presence of boroxol rings [B<sub>3</sub>O<sub>4.5</sub>], pyro- B<sub>2</sub>O<sub>5</sub><sup>4-</sup>, ditri- [B<sub>3</sub>O<sub>8</sub>] and dipenta-borate [B<sub>5</sub>O<sub>11</sub>] groups as Raman bands are present at ~770, ~800, ~1040 and ~1340 cm<sup>-1</sup>. A wide envelope appears around ~465 cm<sup>-1</sup> due to isolated diborate groups as well as Pb - O linking vibrations and the envelope at ~700 cm<sup>-1</sup> is the characteristic of symmetric breathing vibrations of metaborate rings. The silver ions shifted the peak at 806 to 800 cm<sup>-1</sup>. With increasing Fe<sub>2</sub>O<sub>3</sub> content in lead borate glasses containing silver decreases the intensity of peaks at 770 and 800 cm<sup>-1</sup> and grow the band at 1040 and 1340 cm<sup>-1</sup> [109]. The Raman bands in PbO - BaO - B<sub>2</sub>O<sub>3</sub> glasses are assigned at 770, 806, 1230, 1280 and 1450 cm<sup>-1</sup> [110]. The substitution of cadmium to lead in lead borate glasses was represented by Raman band in the range 200 - 300 cm<sup>-1</sup>. The band near 840 cm<sup>-1</sup> is not found and it suggests that CdO<sub>4</sub> is not formed in the glass network. Raman bands were found at 3691, 1275, 927, 720, 623 and 219 cm<sup>-1</sup> in spectra of glass with composition 30PbO<sub>2</sub>·20CdO·50B<sub>2</sub>O<sub>3</sub> [111].

Structure of barium borate glasses consist of mainly two band near 460 and 805 cm<sup>-1</sup>. With the addition of 0.5 mol% Mn in glassy matrix, one more band about 990 cm<sup>-1</sup> was assigned and the intensity of this band increases with increasing the concentration of Mn, i. e. MnO is helpful in formation of orthoborate group. The addition of Mn ions upto 3 mol% the band intensity of band ~ 460 cm<sup>-1</sup> increases and thereafter decreases. There are two new bands at ~ 630 and 700 cm<sup>-1</sup> are

also introduced with the addition. The band at ~ 630 cm<sup>-1</sup> was assigned due to vibration of ring and chain type of meta- and penta- borate groups whereas the band at 700 cm<sup>-1</sup> is attributed to vibrations of chain or ring- type metaborate groups [112]. There

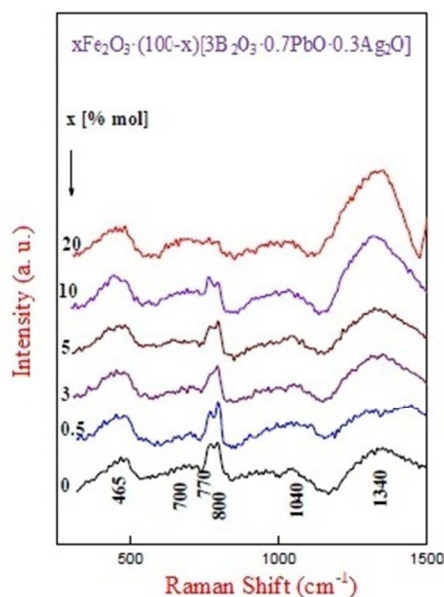


Fig. 6: Raman spectra for xFe<sub>2</sub>O<sub>3</sub>·(100 - x)[3B<sub>2</sub>O<sub>3</sub>·0.7PbO·0.3Ag<sub>2</sub>O] glasses [109]

are five Raman bands at ~ 490, ~ 690, ~ 800, ~ 875 and ~ 1250 cm<sup>-1</sup> that are present in 3B<sub>2</sub>O<sub>3</sub>-As<sub>2</sub>O<sub>3</sub> glass matrix. The band at ~ 490 cm<sup>-1</sup> is caused by the vibrations of isolated diborate groups and/or to vibrations of As - O bonds. The band near 875 cm<sup>-1</sup> decreases progressively with the addition of silver ions. The Raman spectra of xAg<sub>2</sub>O·(1 - x)[zB<sub>2</sub>O<sub>3</sub>·As<sub>2</sub>O<sub>3</sub>] (z = 1,2,3) glasses showed the bands at 490, 685, 803, 880, 960 and 1250 cm<sup>-1</sup>. The addition of manganese ions in xAg<sub>2</sub>O·(1 - x)[2B<sub>2</sub>O<sub>3</sub>·As<sub>2</sub>O<sub>3</sub>] vitreous matrix leads to structural modifications and to increase of disorder degree, particularly for higher manganese ions concentrations (y ≥ 20 %mol) in glass system x[(1 - y)Ag<sub>2</sub>O·yMnO]·(100 - x)[2B<sub>2</sub>O<sub>3</sub>·As<sub>2</sub>O<sub>3</sub>] [113-114].

The Raman studies of xSrO·(1 - x)B<sub>2</sub>O<sub>3</sub> (0.2 ≤ x ≤ 0.7) indicates the absence of band near 806 cm<sup>-1</sup> band whereas presence of bands at 989, 677 and 552 cm<sup>-1</sup> for x = 0.2. Raman spectra of 0.3SrO·0.7B<sub>2</sub>O<sub>3</sub> is dominated by the band at 671 cm<sup>-1</sup> and assigned to di-pentaborate groups whereas for 0.4SrO·0.6B<sub>2</sub>O<sub>3</sub>, a band at 799 cm<sup>-1</sup> is observed. In glass sample with composition 0.5SrO·0.5B<sub>2</sub>O<sub>3</sub>, the band near 744 cm<sup>-1</sup> is attributed to six membered rings with one BO<sub>4</sub> tetrahedra. The intensity of this band decreases with addition of SrO. Addition of metal cations in these glasses introduced a new band near 423 cm<sup>-1</sup>. The peak near 1413 cm<sup>-1</sup> in undoped sample was shifted towards higher wavenumber side with addition on Fe<sup>2+</sup>, Mn<sup>2+</sup> and Zn<sup>2+</sup> ions [115]. The structure of glasses in glass system 20MO.55Bi<sub>2</sub>O<sub>3</sub>.25B<sub>2</sub>O<sub>3</sub> (M = Sr, Ba) was determined by FT-Raman spectra. A broad Raman band at 124 cm<sup>-1</sup> was observed in both barium bismuth borate and strontium bismuth borate glasses. A few resolved Raman bands are also present in region 50 - 400 cm<sup>-1</sup>. These vibrations are assigned due to Bi atoms in the crystalline Bi<sub>24</sub>B<sub>2</sub>O<sub>39</sub>. The external vibration due to α-Bi<sub>2</sub>O<sub>3</sub> was appeared in the region 0 - 150 cm<sup>-1</sup> while the internal vibrations

are ranges 150 - 500  $\text{cm}^{-1}$ . The peak centred at  $\sim 128 \text{ cm}^{-1}$  in the Raman spectra indicates the presence of  $[\text{BiO}_3]$  and  $[\text{BiO}_6]$  units in the structure of glass and glass ceramics. Raman bands at 696 and 723  $\text{cm}^{-1}$  are associated with the symmetric bending vibration of  $[\text{BO}_3]^{3-}$  anion in  $\text{Bi}_{24}\text{B}_2\text{O}_{39}$  and indicate violation of the planarity of  $[\text{BO}_3]^{3-}$  anion. The band due to boroxol ring at 804  $\text{cm}^{-1}$  was disappeared during heat treatment of glasses. The bands at 773  $\text{cm}^{-1}$  is predominant with associated bands at 960, 663 and 487  $\text{cm}^{-1}$ . These bands are comparatively weak in barium bismuth borate glass ceramics as compared to strontium bismuth borate glass ceramics. The overtones of  $[\text{BO}_3]^{3-}$  anion are observed in range 1100 - 1500  $\text{cm}^{-1}$ . The strong band centred at  $\sim 1412 \text{ cm}^{-1}$  was accredited to linked  $[\text{BO}_3]^{3-}$  triangles with one free vertex as in  $\text{BiB}_3\text{O}_6$  heat treated samples. The heat treatment of glasses increases the diborate units as indicated by the increased intensity of the bands 1117, 1031 and 930  $\text{cm}^{-1}$  (strontium contained) and 1107, 1026 and 925  $\text{cm}^{-1}$  (barium contained) [116]. The Raman study of  $\text{Bi}_2\text{O}_3\text{-B}_2\text{O}_3\text{-ZnO-Li-O}$  showed that Bi can form  $[\text{BiO}_3]$  pyramidal or  $[\text{BiO}_6]$  octahedral units. The symmetric stretching anion motion in an angularly constrained Bi - O - Bi configuration was found in region 300 - 600  $\text{cm}^{-1}$ . The vibrations of Bi - O - Bi and  $[\text{BiO}_6]$  octahedral units were present near 390  $\text{cm}^{-1}$ . The position of this band is shifted towards lower wavenumber sides with increase of  $\text{Li}_2\text{O}$ . The shoulder around 555  $\text{cm}^{-1}$  was assigned due to stretching vibrations of Bi - O - / Bi - O - Zn and its intensity increases with the increase of ZnO content. The Raman peak at 220  $\text{cm}^{-1}$  indicates the presence of Zn - O tetrahedral bending vibrations of  $\text{ZnO}_4$  units in the present glass system [14].

The polarized Raman spectra of zinc borate glasses showed features of binary metaborate glasses. The polarized bands were assigned at 950, 840 and 1280  $\text{cm}^{-1}$ . The two weak bands at  $\sim 770$  and 800  $\text{cm}^{-1}$  were also observed Raman pattern. The unpolarized bands were found at 690 and 1420  $\text{cm}^{-1}$ . The strongly polarized shoulder localized at about 250  $\text{cm}^{-1}$  is ascribed to bending modes of  $\text{ZnO}_4$  units. The doping of  $\text{Eu}^{3+}$  in zinc borate increases intensity of Raman band. The band at 440  $\text{cm}^{-1}$  is assigned due to a  $\text{Eu}^{3+}$  - O stretching/ $\text{BO}_3^{3-}$  vibrational mode [117]. The Raman spectra of glasses in glass system  $60\text{B}_2\text{O}_3\text{-}10\text{TeO}_2\text{-}5\text{TiO}_2\text{-}24\text{R}_2\text{O:}1\text{CuO}$  (where R = Li, Na, K) depict the bands of boroxol ring, vibrations of ring / chain type meta borate units and stretching vibrations of B - O<sup>-</sup> bonds with non bridging oxygen. The networking modifying behaviour of  $\text{Na}_2\text{O}$  and  $\text{K}_2\text{O}$  is stronger than  $\text{Li}_2\text{O}$  [118].

Raman spectra of  $(1-x)[3\text{B}_2\text{O}_3\cdot\text{K}_2\text{O}]\cdot x\text{TiO}_2$  glasses consist of bands centred at 420, 475, 600, 670, 770, 800, 850, 930, 1230 and 1450  $\text{cm}^{-1}$ . The intensity of band at 770  $\text{cm}^{-1}$  is higher than that of band near 800  $\text{cm}^{-1}$ . The band at 770  $\text{cm}^{-1}$  is significant upto  $x = 0.2$  and at higher concentration of titanium oxide, the intensity of bands at 420, 475 and 670  $\text{cm}^{-1}$  were found to increase. For  $x = 0.5$ , two shoulders appear at 850 and 600  $\text{cm}^{-1}$ . Further addition of  $\text{TiO}_2$  introduces the appearance of ring type metaborate groups and loose  $\text{BO}_4^-$  tetrahedra. Therefore, the number of non-bridging oxygens increases with increasing titanium concentration and the glass structure becomes more randomized. The addition of  $\text{K}_2\text{O}$  in borate glasses changes the boron coordination number from 3 to 4 [119]. The calcium oxide modified the structure of borate glasses. The structure of borate glasses consist of boroxol groups,

a smaller number of pentaborate groups, diborate groups, chain type metaborate groups and pyroborate groups at lower concentration of calcium oxide whereas at higher concentration of calcium oxide, pentaborate, orthoborate and metaborate groups are also appeared. With increasing the content of calcium oxide, the number of non-bridging oxygen ions increases [120]. The yttrium oxide in calcium borate glasses further modified structure [121]. The glasses with the compositions of  $22.5\text{Gd}_2\text{O}_3\cdot x\text{WO}_3\cdot(47.5-x)\text{MoO}_3\cdot 30\text{B}_2\text{O}_3$  ( $x = 0 - 40$ ) consisted by structural units centred at 344, 840 and 944  $\text{cm}^{-1}$  as shown in Fig. 7. When these glasses were heat treated, then a few more sharp Raman bands were also introduced. The Raman peak at  $\sim 994 \text{ cm}^{-1}$  was assigned due to symmetric stretching vibrations in  $(\text{WO}_4)^{2-}$  tetrahedral unit in the glass composition with  $\text{Sm}_2\text{O}_3$  for  $x = 10$  [15]. Peak assignment of various vibrational units in borate glasses are listed in Table 1.

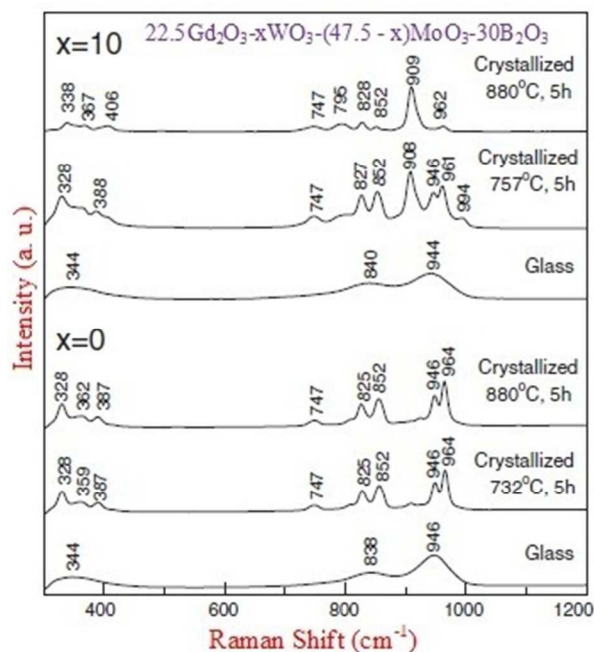


Fig. 7: Raman spectra of glasses in system  $22.5\text{Gd}_2\text{O}_3\cdot x\text{WO}_3\cdot(47.5-x)\text{MoO}_3\cdot 30\text{B}_2\text{O}_3$  ( $x = 0 - 40$ ) [15]. (Reproduced with permission of Elsevier)

Table 1: Assignment of main Raman bands in the spectra of borate glasses.

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
465 - 500	isolated diborate groups	105, 113-115
535	diborate units	102
600 - 650	symmetric breathing vibrations of metaborate rings	100, 103, 111, 115, 119
650-660	pentaborate groups	103, 105
700 - 735	symmetric breathing vibrations of metaborate chains	10, 100, 104, 12, 106, 108, 111, 116
740-775	Symmetric breathing vibrations of six membered rings with one $\text{BO}_4^-$ tetrahedra units	10, 100, 104, 105, 12, 109, 115, 119
765	B-O-B bending mode	102
800-808	boroxol ring	99, 103, 105, 12, 109, 112, 112-115, 117, 119
835-840	Pyroborate vibrations	115, 126



875 - 1000	ortho-borate groups	10, 105, 106, 108, 111, 112-114, 117, 119
1000-1110	diborate groups	10, 100, 102, 103, 106, 112-115
1200	symmetric stretching of B <sub>3</sub> units	102
1216 - 1260	Pyro-borate groups	102, 106, 108, 109, 119
1300 -1600	B-O <sup>-</sup> stretching in metaborate rings and chains	102, 106, 109, 115, 116
3300 - 3500	molecular water	102
0 - 150	external vibration of α-Bi <sub>2</sub> O <sub>3</sub>	116
128	presence of [BiO <sub>3</sub> ] and [BiO <sub>6</sub> ] units	116
150 - 500	internal vibrations of α-Bi <sub>2</sub> O <sub>3</sub>	116
220	Zn - O tetrahedral bending vibrations of ZnO <sub>4</sub> units	117
250	bending modes of ZnO <sub>4</sub> units	118
390	[BiO <sub>6</sub> ] octahedral units	117
300 - 600	angularly constrained Bi - O - Bi configuration	117
555	stretching vibrations of Bi - O - / Bi - O - Zn	117
440	Eu <sup>3+</sup> -O stretching / BO <sub>3</sub> <sup>3-</sup> vibrational mode	118

### 3.2. Silicate Glasses

The networks of silicate glass are shown in Fig. 8. Raman spectroscopic studies of silicate glasses showed the formation of structural units of orthosilicate, silicon-oxygen stretching vibrations of tetrahedral silicate units, symmetric stretching vibrations of silicate tetrahedra, inter-tetrahedral Si - O - Si linkages and structural unit Q<sub>Si</sub><sup>n</sup>, where Q represents the tetrahedral unit and n the number of bridging oxygen (BO) per tetrahedron. For silicon compounds, n varies between 0 and 4, where Si is a central tetrahedral atom ranging from O<sup>0</sup> which represents orthosilicates SiO<sub>4</sub><sup>4-</sup>, Q<sub>Si</sub><sup>4</sup> (tectosilicates), Q<sub>Si</sub><sup>3</sup>, Q<sub>Si</sub><sup>2</sup> and Q<sub>Si</sub><sup>1</sup> representing intermediate silicate structures.

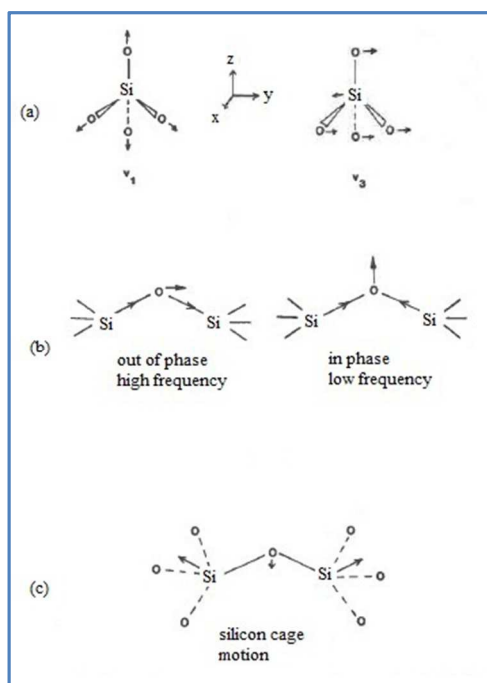


Fig. 8: The selected networks of silicate glasses [16]. (Reproduced with permission of Mineralogical Society of America)

These networks can be modified by adding certain network modifiers such as alkali and alkali earth atoms. The high frequency bands are described in mainly polarized Raman bands centred in the range 1050 - 1100, 950 - 1000, 900 and 850 cm<sup>-1</sup>, they are designated as disilicate, metasilicate, pyrosilicate and orthosilicate, respectively. The low frequency bands are mainly associated with Si - O - Si linkage [16]. The presence of water group in SiO<sub>2</sub> glasses was confirmed by Raman band at 3598 cm<sup>-1</sup>. The weak band near 2350 cm<sup>-1</sup> was arisen due to Si - OH groups involved in intratetrahedral hydrogen bonding across an edge of the SiO<sub>4</sub> tetrahedron. The Si - OH stretching mode at 970 cm<sup>-1</sup> is more prominent than other bands and showing that more silanol groups are present in comparison to suprasil. The bands at 430, 800, 1060 and 1200 cm<sup>-1</sup> arises due to fundamental vibrations of the dry SiO<sub>2</sub> glass. The sharp peaks at 490 and 600 cm<sup>-1</sup> are the characteristic of defect bands. Raman bands at 541 and 1100 cm<sup>-1</sup> in the spectra of dry and wet sodium silicate glasses indicate the presence of SiO units (Q<sub>Si</sub>, species: SiO<sub>4</sub> units with one nonbridging oxygen) and vibrations of the bridging oxygen in the Si - O - Si linkage, respectively. The shoulders peaks near 850, 960 and 1000 cm<sup>-1</sup> suggest that the three components are present in the high-frequency envelope whereas four components near 970, 1060, 1100, and 1150 cm<sup>-1</sup> are similar to hydrous sodium silicate glasses [17].

The effect of cations on symmetric vibrational wavenumber of NBO of Si - O - T in high frequency range is small, whereas the intensity of Raman peak increases with increasing the radius of cation. The scattering cross section increases in order of Li, Na, K, Rb and Cs [18]. The Raman spectrum of sulphur induced sodium calcium silicate glasses indicates the band due to sulphur ions is located at 990 cm<sup>-1</sup>. Not only this band, but some other bands due to silicate network are also present in the pattern [19]. The presence of rare earth oxides in soda-lime-silicate glasses led to a shift of the peak positions at 1100, 790 and 550 cm<sup>-1</sup> and attributed to Si - O - Si asymmetric stretching, Si - O - Si symmetric stretching and bending vibration, respectively. It was observed that peak position of bands at 1100 and 790 cm<sup>-1</sup> are shifted towards low frequency side and peak at 550 cm<sup>-1</sup> was shifted towards high frequency side with doping of lanthanide elements (La<sub>2</sub>O<sub>3</sub>, CeO<sub>2</sub>, Nd<sub>2</sub>O<sub>3</sub> and Gd<sub>2</sub>O<sub>3</sub>) into soda-lime-silicate glass. Y<sub>2</sub>O<sub>3</sub> in soda-lime-silicate glass led the frequencies of peak position at 1100, 790 and 550 cm<sup>-1</sup> shift towards high frequency side. The O<sup>n</sup> species in soda-lime-silicate glasses changes as Si<sub>2</sub>O<sub>7</sub><sup>6-</sup> + 2Si<sub>2</sub>O<sub>5</sub><sup>2-</sup> = 5SiO<sub>3</sub><sup>2-</sup> + SiO<sub>2</sub> with addition of lanthanide elements [20]. The band near 980 cm<sup>-1</sup> was observed in both bulk (suprasil) and sol-gel silica glass. The defect mode bands at 493 and 606 cm<sup>-1</sup> are also assigned in the matrix. Raman bands at ~356, ~264, ~207 and ~128 cm<sup>-1</sup> in quartz crystal are attributed to lattice modes whereas the strong peak at ~465 cm<sup>-1</sup> is the characteristic of symmetric stretching vibration. In the high-frequencies, the bands at ~807 and ~1083 cm<sup>-1</sup> are assigned to Si - O - Si bending and SiO<sub>4</sub> asymmetric stretching vibration, respectively. The addition of CaO in silica glass changes the relative intensities of bands as indicated in Fig. 9. The intensity of band at 1050 cm<sup>-1</sup> was increased due to addition of CaO. With addition of modifier content of MgO, peak position of 805 cm<sup>-1</sup> is shifted towards lower wavenumber side. Comparatively, the CaO changes the structure more significantly than MgO [21-22].

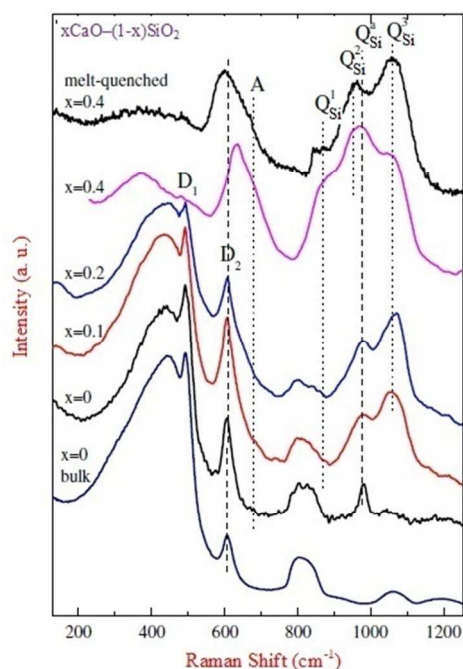


Fig. 9: Raman spectra of  $x\text{CaO}-(1-x)\text{SiO}_2$  ( $x = 0, 0.1, 0.2, 0.4$ ) glasses [21]. (Reproduced with permission of Springer)

The main silicate structure is not much changed at smaller content of metal oxide, but two additional peaks are assigned at 500 and 600  $\text{cm}^{-1}$ . The modifying nature of Cs, Rb, K, Na and Li in silica glass is in increasing order because the intensity of peak at 950  $\text{cm}^{-1}$  increases in same order. The silicate bands are broaden in order  $\text{Cs} < \text{Rb} < \text{K} < \text{Na} < \text{Li} < \text{Ca} < \text{Mg}$  which may suggest that the perturbation of silicate units are increased with increasing cation strength [23].

Raman spectra of potash - lime - silica glasses depict the bands near 300 and 450  $\text{cm}^{-1}$ . The addition of sulphur in glass matrix increases Raman cross-section of vibrational modes due to S - O bands in comparison to less polarized Si - O bonds. The intense Raman features in such glasses at 1000 - 1030  $\text{cm}^{-1}$  indicates the mixture of three calcium sulphates: gypsum ( $\text{CaSO}_4 \cdot 2\text{H}_2\text{O}$ ), (1008  $\text{cm}^{-1}$ ), bassanite ( $\text{CaSO}_4 \cdot 1/2\text{H}_2\text{O}$ ) (1016  $\text{cm}^{-1}$ ) and anhydrite ( $\text{CaSO}_4$ ), (1026  $\text{cm}^{-1}$ ). The bands in Raman spectra in the range of 608 - 679  $\text{cm}^{-1}$  and 1110 - 1167  $\text{cm}^{-1}$  are the feature of silicate networks whereas the peak in the range 412 - 494  $\text{cm}^{-1}$  is the characteristic of cation-oxygen vibrational modes [24]. In Raman spectra of lead silicate glasses, a band near 1000  $\text{cm}^{-1}$  was shifted towards longer wavelength side and becomes broader with increasing the content of lead. The intensity of Raman peak at 140  $\text{cm}^{-1}$  was increased by increasing content of lead. Other peaks at 450, 950 and 1060  $\text{cm}^{-1}$  are also modified by PbO. After UV irradiation with energy density 150  $\text{mJ}/\text{cm}^2$ , the significant changes in the spectra were observed. The intensity of Pb - O band at 140  $\text{cm}^{-1}$  was decreased after UV irradiation and no new band appears in the Raman spectra. The decrease of intensity after UV irradiation is caused by the broken Pb - O bond in lead silicate glasses and the broken Pb - O bond is related to the energy density of UV beam [25]. The lead silicate glasses provide a simple absorption spectrum gamma irradiation [122]. Band assignment of various vibrational units in silicate glasses are listed in Table 2.

Table 2: Assignment of main Raman bands in the spectra of silicate glasses.

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
580	Si-O <sup>0</sup> rocking motions in fully polymerized $\text{SiO}_2$ ( $\text{Q}^4$ ) units	20
600	Si-O-Si bending vibration in depolymerized structural units	20
700-850	Si-O-Si symmetric stretching of bridging oxygen between tetrahedra	20
1050 - 1100	disilicate	16
950 - 1000	metasilicate	16
790	Si - O - Si symmetric stretching	20
807	Si - O - Si bending	21-22
850	Orthosilicate, $\text{SiO}_4^{4-}$	16
900	pyrosilicate	16
970	Si - OH stretching mode	17
1083	$\text{SiO}_4$ asymmetric stretching vibration	21-22
1100	Si - O - Si asymmetric stretching	20
2350	Si-OH groups involved in intratetrahedral hydrogen bonding across an edge of the $\text{SiO}_4$ tetrahedron	17

### 3.3. Phosphate Glasses

The typical Raman spectra of phosphate glass are shown in Fig. 10. The structure of the phosphate glasses are mainly consisted by  $\text{Q}_n^p$  groups of  $\text{PO}_4$  tetrahedra ( $n =$  number of bridging oxygen per  $\text{PO}_4$  tetrahedron). The symmetric stretching of P - O is assigned at 1260  $\text{cm}^{-1}$  of an asymmetric profile (620  $\text{cm}^{-1}$ ), symmetric stretching vibration of P - O (1380  $\text{cm}^{-1}$ ), the symmetric stretching of a non bridging oxygen on a  $\text{Q}_n^p$  tetrahedron at 1170  $\text{cm}^{-1}$ , P - O - P symmetric vibrations at 690  $\text{cm}^{-1}$ , the symmetric stretching of the orthophosphate groups  $\text{PO}_4^{2-}$  near 960  $\text{cm}^{-1}$ . These bands are modified by dopants like earth and alkaline-earth atoms. Addition of  $\text{TiO}_2$  upto 10 mol % reveals the depolymerization of the phosphate glass network by systematic conversion of  $\text{Q}_n^p$  structural units into  $\text{Q}_m^p$  and finally into  $\text{Q}_0^p$  structural units. Even though  $\text{Q}_n^p$  to  $\text{Q}_m^p$  conversion is taking place due to breaking of P - O - P linkages, formation of P - O - Ti and P - O - Al linkages provide cross linking between short P-structural units. Above 10 mol %  $\text{TiO}_2$  content, network is highly depolymerized due to the formation of  $\text{Q}_0^p$  structural units and cross-linking becomes poor. The band near 700  $\text{cm}^{-1}$  is assigned due to the symmetric stretching of P - O - P linkages in  $\text{Q}_n^p$  and  $\text{Q}_m^p$  structural unit. On addition of  $\text{TiO}_2$  in phosphate glass, the intensity of the bands near 1280, 1170 and 700  $\text{cm}^{-1}$  is decreased and a few new bands at 1210, 1090, 900, 750 and 625  $\text{cm}^{-1}$  are introduced. The bands at 1210 and 1090  $\text{cm}^{-1}$  are assigned due to asymmetric and symmetric stretching vibrations of  $\text{PO}_3$  groups ( $\text{Q}^1$  structural units), respectively. The coordination of  $\text{Ti}^{4+}$  ions by oxygen atoms is found in Raman spectra and band near 625  $\text{cm}^{-1}$  is attributed to the vibration of octahedral titanate units  $\text{TiO}_6$ , as its relative intensity increases with increasing  $\text{TiO}_2$ . The band centred at 930  $\text{cm}^{-1}$  is associated with shorter Ti - O bond than for octahedral titanium and attributed to the titanyl bonds (Ti - O or -Ti - O - Ti-) associated with a five coordinated Ti. The band at 900 and 525  $\text{cm}^{-1}$  are related with asymmetric stretching of P - O - P bridges and bending vibrations of P - O bonds, respectively [123]. The

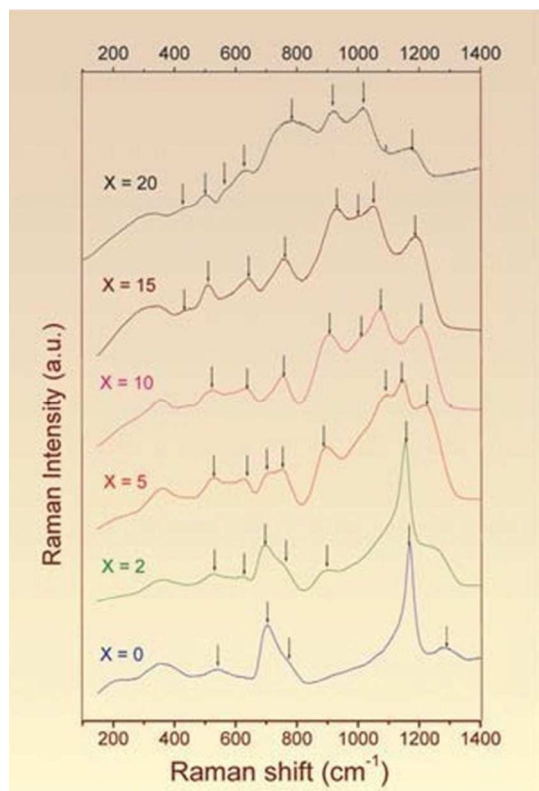


Fig. 10: Raman spectra of  $40\text{Na}_2\text{O}-10\text{Al}_2\text{O}_3-x\text{TiO}_2-(50-x)\text{P}_2\text{O}_5$  glasses [123]

bands in Raman spectra of  $39\text{AlF}_3 \cdot 11\text{NaF} \cdot 10\text{LiF} \cdot (40-x)\text{CaF}_2 \cdot \text{MgF}_2 \cdot \text{SrF}_2 \cdot \text{BaF}_2 \cdot x\text{NaPO}_3$  system at 515, 570, and 620  $\text{cm}^{-1}$  is assigned to octahedral  $[\text{AlF}_6]$ , five-coordinated  $[\text{AlF}_5]$  and tetrahedral  $[\text{AlF}_4]$  structural units, respectively. The intensity of band at 1330  $\text{cm}^{-1}$  decreases whereas, the intensity of bands at 1070, 860 and 775  $\text{cm}^{-1}$  increases with increase of phosphate content. The low frequency band at 70  $\text{cm}^{-1}$  is boson band. Raman peaks at 250 and near range 350 - 420  $\text{cm}^{-1}$  are associated with symmetric stretching vibrations of fluorine species bonded by the modifying cations [124]. The bands at 1194, 1294 and 1349  $\text{cm}^{-1}$  are prominent in Raman spectra of  $\text{Na}_2\text{O}-\text{Al}_2\text{O}_3-\text{P}_2\text{O}_5$ . With increasing the content of  $\text{Al}_2\text{O}_3$ , intensity of bands near 1200  $\text{cm}^{-1}$  is enhanced whereas the band at 1161  $\text{cm}^{-1}$  is disappeared. The band at 676  $\text{cm}^{-1}$  is shifted towards higher wavenumber side due to formation of aluminium metaphosphate structure. Due to introduction of  $\text{Al}_2\text{O}_3$ , the band located at 1349  $\text{cm}^{-1}$  is disappeared. The weak vibration due to  $\text{PO}_2$  bending and in chains of O - P - O bending motions is assigned at 300 and 360  $\text{cm}^{-1}$ , respectively. With addition of  $\text{Fe}_2\text{O}_3$  in this glass system, the intensity of the bands at 1257, 1185 and 706  $\text{cm}^{-1}$  is increased whereas position of band of at 1349  $\text{cm}^{-1}$  is shifted to 1309  $\text{cm}^{-1}$  due to depolymerization. The ferric oxide in sodium phosphate glasses forms the structure of metaphosphate and connected by P - O - Fe bonds [125]. The bands in Raman spectrum of zinc-alumino-sodium-phosphate glasses with addition of  $\text{Pr}^{3+}$  and  $\text{Nd}^{3+}$  are located at 545, 698, 1042 and 1162  $\text{cm}^{-1}$  [126]. Raman spectra of  $55\text{P}_2\text{O}_5 \cdot 30\text{CaO} \cdot (25-x)\text{Na}_2\text{O} \cdot x\text{TiO}_2$  ( $0 \leq x \leq 5$ ) system have bands near 695, 705, 930, 1180, 1280 and 1370  $\text{cm}^{-1}$ . The band at 930  $\text{cm}^{-1}$  is attributed to Ti-O band and the intensity of band at 1370  $\text{cm}^{-1}$  is increased by addition of  $\text{TiO}_2$  upto 5 mol% [127]. Raman spectra of  $\text{Li}_2\text{O}-\text{BaO}-\text{Al}_2\text{O}_3-\text{La}_2\text{O}_3-\text{P}_2\text{O}_5$  glasses showed the bands are found at 1850, 1700, 1180 and 700  $\text{cm}^{-1}$  [128]. The spectra of phosphate glasses containing  $\text{TiO}_2$  consist of bands which are assigned at 340 - 540, 693 - 697, 1010 - 1020, 1160 - 1173, 1248 - 1271  $\text{cm}^{-1}$ . The band in range 1010 - 1020  $\text{cm}^{-1}$  is

disappeared at higher content of  $\text{TiO}_2$  and peak at 905  $\text{cm}^{-1}$  is assigned to asymmetric  $\text{TiO}_5$  [129].

Raman investigations of  $x\text{MoO}_3 \cdot (100-x)[\text{P}_2\text{O}_5 \cdot \text{CaO}]$ ,  $0 \leq x \leq 50$  mol% glass system showed that Raman bands are centred at 316, 392, 536, 693, 1175 and 1274  $\text{cm}^{-1}$ . The addition of  $\text{MoO}_3$  about  $0.3 \leq x \leq 1$  mol%, a new band at 1032  $\text{cm}^{-1}$  is introduced. The intensity of band at 683  $\text{cm}^{-1}$  was decreased at  $5 \leq x \leq 10$  mol% of  $\text{MoO}_3$  whereas peak at 436 is shifted towards higher wavenumber side. The band located at 700  $\text{cm}^{-1}$  becomes weaker and the peak at 1270  $\text{cm}^{-1}$  is disappeared.  $\text{MoO}_3$  is strong modifier and forms the vibration at 600  $\text{cm}^{-1}$  due to asymmetric vibration of (Mo - O - Mo). The band at 894  $\text{cm}^{-1}$  is caused to orthomolybdate,  $(\text{MoO}_4)^{2-}$  units. The band at 963  $\text{cm}^{-1}$  is assigned due to stretching vibrations of partially isolated Mo - O bonds in deformed  $(\text{MoO}_6)$  units [130]. The structure of  $x\text{CaO} \cdot (100-x)(0.4\text{Fe}_2\text{O}_3 \cdot 0.6\text{P}_2\text{O}_5)$  ( $x = 0, 10, 20, 30, 40, 50$  mol%) glasses are made by structural units of phosphate and modified by calcium as well as iron. Their Raman bands are observed at 626, 772, 950, 1071 and 1241  $\text{cm}^{-1}$  [131]. Raman spectra of calcium coacervate show the peaks, which are found at 1168 and 698  $\text{cm}^{-1}$ . The Raman spectrum of the croconate ion consist of bands located at 564, 649, 1605 and 1722  $\text{cm}^{-1}$  [132]. In Raman spectra of glasses of glass system  $x\text{Ga}_2\text{O}_3 \cdot (1-x)\text{P}_2\text{O}_5$ , there are five Raman bands are observed at 1320, 1210, 715, 620 and 350  $\text{cm}^{-1}$ . The depolarization by Na, Ca, Zn, Mg, Al and Be is found at 350  $\text{cm}^{-1}$ . The low frequency band is assigned due to Ga - O - M bridges. The band near 640  $\text{cm}^{-1}$  is assigned in gallate system with four fold coordinated Ga atoms [133-134].

The structure of  $\text{PbO}-\text{MoO}_3-\text{P}_2\text{O}_5$  glasses is formed by networks of phosphate units and modified by  $\text{MoO}_3$ . The incorporation of  $\text{MoO}_3$  results the transformation of octahedral  $\text{MoO}_6$  units into tetrahedral  $\text{MoO}_4$  units. The band at 930  $\text{cm}^{-1}$  is shifted towards high frequency side with increasing content of  $\text{MoO}_3$ . The intense band near 1158  $\text{cm}^{-1}$  is shifted towards lower frequency side with increasing content of  $\text{MoO}_3$ . The bands in region 850 - 970  $\text{cm}^{-1}$  are attributed to Mo - O and Mo - O bonds in  $\text{MoO}_6$  tetrahedra [135]. The structure of  $\text{PbO}-\text{Ga}_2\text{O}_3-\text{P}_2\text{O}_5$  glasses doped with  $\text{Cr}_2\text{O}_3$  is modified by forming tetrahedral  $\text{CrO}_4^{2-}$  structural units. The phosphate network is modified by all Pb, Ga and Cr ions. The band at 398  $\text{cm}^{-1}$  is assigned due to Ga - O - P linkages and peak at 368  $\text{cm}^{-1}$  is attributed to  $\text{GaO}_6$  vibrational groups. The band at 270  $\text{cm}^{-1}$  is associated with Cr-O vibrations of  $\text{CrO}_4^{2-}$  units [136]. Lead phosphate network can also be modified by Zn and Cu ions. The phosphate network after modification is assigned at positions 690, 747, 904, 1026 and 1155  $\text{cm}^{-1}$ . The intensity of band 1155  $\text{cm}^{-1}$  is decreased with addition of  $\text{BaO}$  [137]. High ZnO content in lead phosphate glasses forms the network of P - O - Zn bonds [138]. The network of lead-gallium phosphate glasses can be modified by rare earth ions ( $\text{Eu}^{3+}$ ,  $\text{Dy}^{3+}$ ,  $\text{Tb}^{3+}$ ,  $\text{Er}^{3+}$ ). The spectrum contains two shoulders centred at 1050 and 1120  $\text{cm}^{-1}$  corresponds to symmetric stretching vibration of diphosphates in  $\text{Q}^1$  units. The other two shoulders near 900 and 1250  $\text{cm}^{-1}$  are attributed to symmetric stretching vibration of the  $\text{PO}_4$  in the  $\text{Q}_p^0$  units and asymmetric stretching vibration of the  $\text{PO}_2$  in  $\text{Q}_p^2$  units, respectively. The high frequency band about 1120  $\text{cm}^{-1}$  is related to phonon energy of the lead phosphate glass and its intensity decreases, when the  $\text{PbO}$  is substituted by the  $\text{PbF}_2$  components

[139]. The bands near 1172, 1096, 1032, 973, 930, 742, 627 and 474  $\text{cm}^{-1}$  are found in the glasses in the system  $\text{PbO-Fe}_2\text{O}_3\text{-P}_2\text{O}_5$ . The bands at 1080 and 966  $\text{cm}^{-1}$  are related to  $[\text{P}_2\text{O}_7]^{2-}$  and terminal  $[\text{PO}_3]^{2-}$  symmetric groups in  $\text{Q}_p^1$  pyrophosphates [140].  $\text{In}_2\text{O}_3$  in  $\text{PbO-P}_2\text{O}_5\text{-As}_2\text{O}_3$  plays an important role in network modification and form the structural unit of  $\text{InO}_6$ . The band centred at 200  $\text{cm}^{-1}$  is related with In - O vibrations of  $\text{InO}_6$  structural units. The band near 110 and 250 - 350  $\text{cm}^{-1}$  is attributed to Pb - O ionic bond vibrations and  $\text{PbO}_4$  structural vibrations, respectively. With increasing the concentration of  $\text{In}_2\text{O}_3$ , the intensity of the bands due to symmetric vibrations of the phosphate groups is increased whereas, the intensity of bands due to torsional vibrations of  $\text{PO}_4$  structural units and  $\text{As}_2\text{O}_3$  structural vibrations is decreased [141]. The fraction of orthophosphate  $\text{Q}_p^0$  units in lead iron phosphate glasses increases with increasing  $\text{Cr}_2\text{O}_3$  content. The band at 600  $\text{cm}^{-1}$  is assigned due to the different P - O, Fe - O and Pb - O stretching and bending modes. The intense peak is assigned at 1062  $\text{cm}^{-1}$  in Cr free glass. The intensity of band at 979  $\text{cm}^{-1}$  is enhanced with Cr addition. The ferric oxide also forms the network of Fe - O stretching vibration in  $\text{FeO}_6$  octahedra in  $\alpha\text{-Fe}_2\text{O}_3$  and it is indicated by band near 614  $\text{cm}^{-1}$  [142-143].

The Raman spectra of iron phosphate glasses investigated by Milankovic et al. showed that network of phosphate glasses are modified by incorporation of  $\text{MoO}_3$ . Addition of  $\text{MoO}_3$  in glass matrix replaced the stronger P - O - P bonds with a weaker Mo - O - P bonds as the bond covalence decreases from P - O to Mo - O and the bands located at 850 and 980  $\text{cm}^{-1}$  become stronger. The iron works as network modifier for phosphate network and band assigned at 400  $\text{cm}^{-1}$  due to  $\text{PO}_4$  polyhedra was modified by it. The Raman band near 864  $\text{cm}^{-1}$  is associated with the stretching mode of isolated orthomolybdate  $(\text{MoO}_4)^{2-}$ . Due to increase of NBO and increase of  $(\text{PO}_4)^{3-}$ , some bands show the splitting in the spectra. The asymmetric vibration of Mo-O-Mo is assigned at 620  $\text{cm}^{-1}$ , whereas symmetric Mo - O - Mo vibrational mode is appeared at 520  $\text{cm}^{-1}$  [144]. The addition of  $\text{La}_2\text{O}_3$  in iron phosphate glasses lowers the peak position of band at 1072  $\text{cm}^{-1}$  and increases the relative area of band at 1231  $\text{cm}^{-1}$  whereas decreases for band at 1072  $\text{cm}^{-1}$ . This results is related with depolymerization of phosphate network  $2\text{Q}_p^1 = \text{Q}_p^2 + \text{Q}_p^0$  [145]. The most intense peak in Raman spectra of cesium doped iron phosphate glasses is assigned at 1085  $\text{cm}^{-1}$  due to  $\text{Q}_p^1$ . This band is shifted towards lower wavenumber side with increase of Cs content. The intensity of band at 555  $\text{cm}^{-1}$  is enhanced due to disproportion of pyrophosphate linkage. The intensity of this band increases with increase of cesium loading and indicates the increasing concentration of orthophosphate units. The bands at 270 and 200  $\text{cm}^{-1}$  are related with bending vibrations of P - O - P and Cs - O vibrations, respectively [146]. The P - O - P network is depolymerized with increasing  $\text{CeO}_2$  content in glass compositions and results the conversion of  $\text{Q}_p^1$  to  $\text{Q}_p^2$  structural units. The addition of  $\text{CeO}_2$  in iron phosphate glasses introduced some new bands near 1000, 1150 and 1279  $\text{cm}^{-1}$ . The band at 1037  $\text{cm}^{-1}$  is shifted towards higher wavenumber side with the substitution of  $\text{FeO}_3$  by  $\text{CeO}_2$ . The band at 522  $\text{cm}^{-1}$  is related to Fe - O bonds and a new peak at 460  $\text{cm}^{-1}$  is attributed to cerium oxygen polyhedron vibration of  $[\text{CeO}_8]$  unit symmetrical stretching mode at higher content of  $\text{CO}_2$  [147]. The

depolymerization of the phosphate network, in the form of conversion of  $\text{Q}_p^1$  to  $\text{Q}_p^0$  groups ( $2\text{Q}_p^1 + \text{Na}_2\text{O} = 2\text{Q}_p^0$ ), is take place with the addition of  $\text{Na}_2\text{SO}_4$  content and results the conversion of P - O - P bridging oxygens to P - O -  $\text{Na}^+$  non-bridging oxygens. The band at 1039  $\text{cm}^{-1}$  is shifted to higher frequency region with addition of modifier  $\text{Na}_2\text{O}$  ( $\text{Na}_2\text{SO}_4 = \text{Na}_2\text{O} + \text{SO}_3$ ) content. The intensity of the band at 931  $\text{cm}^{-1}$  is also increased with this addition [148]. The vibrations at 170, 306 and 420  $\text{cm}^{-1}$  arise due to network vibrations and librations of Fe - O polyhedral and phosphate network units, respectively. The peak position of 1090  $\text{cm}^{-1}$  is shifted towards low wavenumber side with heat treatment at transition temperature of glass [149].

The  $\text{MoO}_3$  plays an important role for modification in structural units of glass system  $\text{ZnF}_2\text{-Bi}_2\text{O}_3\text{-P}_2\text{O}_5$ . With addition of  $\text{MoO}_3$  in glass matrix form the structural units of  $\text{MoO}_4$  and  $\text{MoO}_6$ . The intensity of band due to asymmetric vibrations of Mo - O - Mo linkages with increase in dopant concentration is observed to increase with a slight shift in the position towards lower wavenumber. In Raman spectra of motioned system, bands due to  $\text{BiO}_6$  and  $\text{BiO}_3$  pyramidal structural units are assigned at 547 and 230  $\text{cm}^{-1}$ , respectively. The modifier Zn takes the interstitial positions in the form of dangling bonds along with non-bridging oxygens as Zn - O - P linkages and fluorine breaks the P - O network [150-151]. The doping of  $\text{WO}_3$  in place of  $\text{MoO}_3$  forms the network of bending vibrations of O - W - O with corner shared octahedra. The noticeable decrease in the intensity of W - O stretching was also observed which is associated  $\text{WO}_4$  tetrahedral units as given in Fig. 11 [152]. Trivalent samarium doped K - Mg - Al zinc fluorophosphate glasses are formed by networks of phosphate and modifiers by K, Mg, Al, Zn and F. The stretching vibration of Al - O linkages in  $\text{AlO}_4$  structural units was observed [153-154]. The structure of  $2\text{P}_2\text{O}_5\text{-CaO-0.05ZnO}$  glass matrix contains Raman bands near 345, 460, 903, 959, 1011, 1137, 1168, 1297 and 1400  $\text{cm}^{-1}$ . The content of  $\text{Ag}_2\text{O}$  in the glass system modified the network of ring structures. The slight depolymerization was also taken place in glass matrix. Raman spectra of  $x\text{Ag}_2\text{O}\cdot(100-x)[2\text{P}_2\text{O}_5\text{-CaO-0.05ZnO}]$  confirm the modifying nature of Ca, Zn and Ag. With increasing content of  $\text{Ag}_2\text{O}$ , the intensity of band near 554  $\text{cm}^{-1}$  is increased and the peak position of band at 690  $\text{cm}^{-1}$  is shifted towards lower frequency side due to change in the chain P - O - P bond angles. The silver oxide favours the presence of orthophosphate, pyrophosphate and small metaphosphate chains, these informations are revealed by the increase in intensity of the bands at 554, 1100 and 670  $\text{cm}^{-1}$  [155].

Raman spectroscopic investigations of  $\text{TiO}_2\text{-P}_2\text{O}_5\text{-CaO}$  glasses showed that the structure consists of a distorted Ti octahedral linked to pyrophosphate unit through P - O - Ti bonds at higher content of  $\text{TiO}_2$  and thus, the depolymerization was take place in P - O - P network. In Raman spectra of  $\text{TiO}_2$  free calcium phosphate glasses, the strong bands near 680, 1175 and 1290  $\text{cm}^{-1}$  are assigned in the glassy network whereas with addition of  $\text{TiO}_2$ , the existed bands are disappeared. The new bands at 765 and 920  $\text{cm}^{-1}$  are also introduced with  $\text{TiO}_2$  addition [156].  $\text{Li}_2\text{O}$  further modified the network of  $\text{TiO}_2\text{-P}_2\text{O}_5\text{-CaO}$  glasses [157]. Raman spectra of  $(\text{P}_2\text{O}_5\text{-K}_2\text{O-Al}_2\text{O}_3\text{-CaF}_2)$  glasses depict the bands at 335, 543, 625, 728, 1075, 1135 and 1284  $\text{cm}^{-1}$ . The band centred at 728  $\text{cm}^{-1}$  is related with vibrations of P - F chain formation,

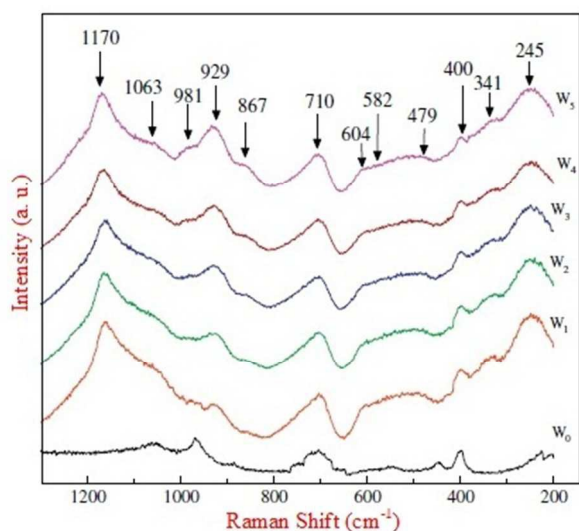


Fig. 11: Raman spectra of  $\text{ZnF}_2 \cdot \text{Bi}_2\text{O}_3 \cdot \text{P}_2\text{O}_5$  glass-ceramics doped with different concentrations of  $\text{WO}_3$  [152]. (Reproduced with permission of Elsevier)

whereas less intense band at  $625 \text{ cm}^{-1}$  is assigned due to  $\text{Al}_2\text{O}_3$  vibrations. The band at  $335 \text{ cm}^{-1}$  is attributed to  $\text{K}_2\text{O}$  [158-159].  $\text{CuO}$  and  $\text{V}_2\text{O}_5$  work as network modifier for calcium phosphate glasses. At higher concentration of  $\text{CuO}$  and  $\text{V}_2\text{O}_5$ , the glass matrix forms the network of  $\text{V} - \text{O} - \text{P}$  and  $\text{Cu} - \text{O} - \text{P}$ . The band assigned at  $1175 \text{ cm}^{-1}$  indicates the depolymerization of phosphate network. The band at  $980 \text{ cm}^{-1}$  is related with  $\text{V} - \text{O}$  vibration in  $\text{VO}_5$  tetragonal pyramid and stretching vibrations of  $\text{O} - \text{V} - \text{O}$  in metavanadate chains are assigned at  $970\text{-}905 \text{ cm}^{-1}$  [160]. Raman spectra of  $(50 - y)\text{Na}_2\text{O} \cdot y\text{CuO} \cdot 10\text{Bi}_2\text{O}_3 \cdot 40\text{P}_2\text{O}_5$  ( $0 \leq y \leq 25 \text{ mol}\%$ ) glasses depict the formation of  $\text{P} - \text{O} - \text{Bi}$  and  $\text{P} - \text{O} - \text{Cu}$  bonds.  $\text{Bi}_2\text{O}_3$  acts as network modifier and is incorporated in the form of  $\text{BiO}_6$  units. No depolymerization was taken place with the replacement of  $\text{Na}_2\text{O}$  by  $\text{CuO}$  [161]. As  $\text{CuO}$  content increases in  $x\text{CuO} \cdot (1 - x)\text{P}_2\text{O}_5$  glasses, the intensity of the  $\text{P} - \text{O}$  band associated with the  $\text{Q}^3$   $\text{Cu}$  tetrahedra increases to  $x < 0.33$  and decreases with a concomitant increase of the intensity of the band at  $1265 \text{ cm}^{-1}$  due to the asymmetric vibration of the  $\text{PO}_2$  groups on  $\text{Q}^2$  tetrahedra. When  $x > 0.33$ , the isolated  $\text{Cu}$ -octahedra begin to share common oxygens to form a sub-network in the phosphate matrix. The intensity of band at  $640 \text{ cm}^{-1}$  increases with addition of  $\text{CuO}$ .  $\text{CuO}$  in phosphate glasses promotes the depolymerization effects [162].

Raman spectra of cadmium phosphate glasses reveal the formation of  $\text{P} - \text{O} - \text{V}$  and  $\text{V} - \text{O} - \text{V}$  bonds instead of  $\text{P} - \text{O} - \text{P}$  vibrations after addition of  $\text{V}_2\text{O}_5$ . The presence of band at  $760 \text{ cm}^{-1}$  is attributed to  $\text{V} - \text{O} - \text{V}$  vibration. The intensity of bands at  $615$  and  $974 \text{ cm}^{-1}$  is increased with addition of  $\text{V}_2\text{O}_5$  about  $20 \text{ mol}\%$ . These bands are related with  $\text{P} - \text{O} - \text{V}$  and  $\text{V} - \text{O}$  vibration  $\text{VO}_5$  trigonal bipyramids, respectively [163]. The  $\text{Li}^+$  and  $\text{Cd}^{2+}$  cations in the phosphate glass network are responsible for shifting the bands as well as intensity modification. The band at  $1120 \text{ cm}^{-1}$  is shifted towards high frequency side with increase of  $\text{Li}_2\text{O}$  content [164]. With increasing the content of  $\text{BaO}$  in lithium barium phosphate glasses, the intensity of all bands is enhanced. The bands at  $712$  and  $1163 \text{ cm}^{-1}$  are shifted slightly towards low wavenumber side with increase of  $\text{BaO}$ . The decrease in wavenumber is related with an increase of average length of the  $\text{P}$

-  $\text{O}$  bond resulting from  $\pi$ -bond delocalization on the  $\text{Q}_p^3$  tetrahedra [165]. Raman spectra of phosphate glasses with addition of  $\text{Er}^{3+} / \text{Yb}^{3+}$  ions depict  $312, 380, 685, 1165$  and  $1275 \text{ cm}^{-1}$  [166]. The assignment of main Raman bands in the spectra of phosphate glasses are listed in Table 3.

Table 3: Assignment of main Raman bands in the spectra of phosphate glasses.

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
1300-1350	$\text{P}=\text{O}$ stretching of terminal oxygen	123
1267-1294	$\text{PO}_2$ asymmetric stretching of non-bridging oxygen, in $\text{Q}^2$ units	123
1260, 1380	symmetric stretching of $\text{P} - \text{O}$	123
1210	asymmetric stretching vibrations of $\text{PO}_3$ groups	123
1170	symmetric stretching of a non bridging oxygen on a $\text{Q}_p^2$ tetrahedron	123
1161	$\text{PO}_2$ symmetric stretching of non-bridging oxygen, in $\text{Q}^2$ units	123
1065-1097	$\text{PO}_2$ symmetric stretching of non-bridging oxygen, in $\text{Q}^1$ units	123
1090	symmetric stretching vibrations of $\text{PO}_3$ groups	123
1080	$[\text{P}_2\text{O}_7]^{2-}$ symmetric groups in $\text{Q}_p^1$ pyrophosphates	140
966	terminal $[\text{PO}_3]^-$ symmetric groups in $\text{Q}_p^1$ pyrophosphates	140
960	symmetric stretching of the orthophosphate groups $\text{PO}_4^{3-}$	123
940	$\text{PO}_2$ symmetric stretching of non-bridging oxygen, in $\text{Q}^0$ units	123
900	asymmetric stretching of $\text{P} - \text{O} - \text{P}$ bridges	123
700	the symmetric stretching of $\text{P} - \text{O} - \text{P}$ linkages in $\text{Q}_p^2$ and $\text{Q}_p^1$ structural unit	123
690	$\text{P} - \text{O} - \text{P}$ symmetric vibrations	123
676-708	$\text{P} - \text{O} - \text{P}$ symmetric stretching of bridging oxygen, in $\text{Q}^2$ units	123
525	bending vibrations of $\text{P} - \text{O}$ bonds	123
390-410	$\text{O} - \text{P} - \text{O}$ bending modes	123
398	$\text{Ga} - \text{O} - \text{P}$ linkages	136
360	chains of $\text{O} - \text{P} - \text{O}$ bending motions	125
300	$\text{PO}_2$ bending $\text{O} - \text{P} - \text{O}$ bending motions	125

### 3.4. Borosilicate Glasses

Typical Raman spectra of alkali borosilicate glasses are shown in Fig. 12. The network structure of borosilicate glasses is formed by borate as well as silicate structural units. The structure of lithium silicate glasses consists of Raman bands at  $850 \text{ cm}^{-1}$  ( $\text{Q}_{\text{Si}}^2$ ),  $950 \text{ cm}^{-1}$  ( $\text{Q}_{\text{Si}}^2$ ),  $1000 \text{ cm}^{-1}$  ( $\text{Q}_{\text{Si}}^3$ ),  $1050 \text{ cm}^{-1}$  ( $\text{Q}_{\text{Si}}^3$ ),  $1080 \text{ cm}^{-1}$  ( $\text{Q}_{\text{Si}}^3$ ) and  $1150 \text{ cm}^{-1}$  ( $\text{Q}_{\text{Si}}^4$ ). The borosilicate structure is associated with the vibrational modes of the two medium-range ( $400$  to  $850 \text{ cm}^{-1}$ ) order superstructures, reedmergnerite  $[\text{BSi}_3\text{O}_8]^-$  and danburite  $[\text{B}_2\text{Si}_2\text{O}_8]^{2-}$  and high frequency region ( $850 - 1250 \text{ cm}^{-1}$ ) vibration related with silicon  $\text{Q}_{\text{Si}}^n$  vibrational modes [167]. Raman spectra of sodium borosilicate glasses depict the bands centred at  $350, 500, 630, 670, 770$  and  $808 \text{ cm}^{-1}$ . The bands between  $300$  and  $500 \text{ cm}^{-1}$  are attributed to mixed stretching and

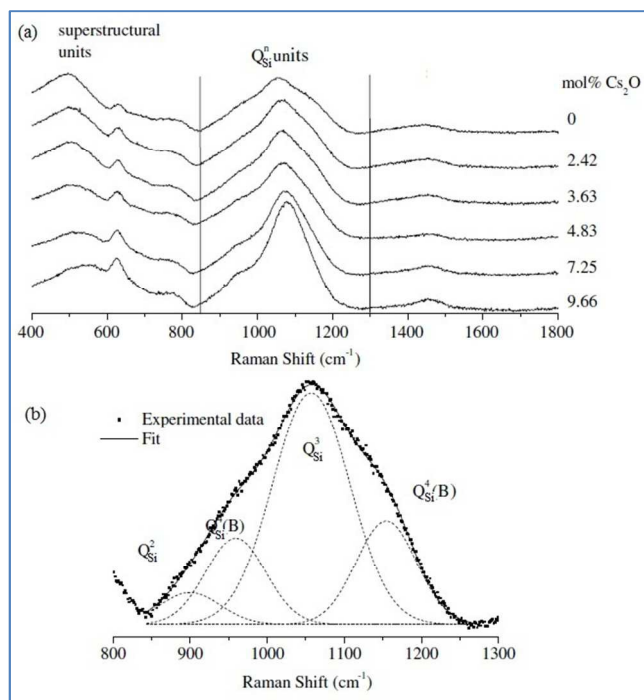


Fig. 12: Raman spectra of alkali borosilicate glasses [167]. (Reproduced with permission of Elsevier)

bending modes of Si-O-Si units. The bands in the range 550 - 850  $\text{cm}^{-1}$  are caused by ring breathing modes. The bands at 670, 770 and 808  $\text{cm}^{-1}$  are the signature of tetraborate groups, four- and three-coordinated boron in diborate and boroxol rings, respectively. The intensity of peak at 808  $\text{cm}^{-1}$  decreases with increase of sodium content whereas the band near 770  $\text{cm}^{-1}$  shows opposite behaviour. This means sodium ions favour the formation of  $\text{BO}_4$  units at the expense of  $\text{BO}_3$  units. The polymerization of the silicate network first increases with  $\text{Na}_2\text{O}$  content upto  $\text{Na}_2\text{O}/\text{B}_2\text{O}_3 = 0.6$  and thereafter decreases with the increase of  $\text{Na}_2\text{O}$  content [168]. Sulphate in borosilicate glass showed its presence in the form of vibrational modes of  $\text{SO}_4$  tetrahedra. The major bands of sulphate, assigned at 1100, 1000, 990, 620 and 460  $\text{cm}^{-1}$  are related to asymmetric S - O stretch modes, symmetric S - O stretching, symmetric S - O stretch vibrations of tetrahedral  $\text{SO}_4^{2-}$ , asymmetric O - S - O bend modes and symmetric O - S - O bend modes, respectively [169]. When the  $\text{B}_2\text{O}_3$  content is increased, the broad bands near 1500 and 800  $\text{cm}^{-1}$  are developed. The shoulder peak at 590  $\text{cm}^{-1}$  is assigned due to symmetric oxygen breathing vibrations of three-membered siloxane rings. The network modifier ions  $\text{Na}^+$  and  $\text{Ba}^{2+}$  are used for conversion of  $\text{BO}_3$  units into  $\text{BO}_4$  units before forming non bridging oxygen (NBO) in the network. As the silica content increases in the glassy matrix and the ratio Na/Si (or Ba/Si) decreases, which could depolymerise the silicate network after forming  $\text{BO}_4$  units increases. It can be concluded that the silicate network in  $\text{SiO}_2$ -rich borosilicate glasses remains more polymerized than in boron-rich glasses [170]. The bands in the range 1200 - 1100  $\text{cm}^{-1}$  are disappeared as the ratio  $\text{B}_2\text{O}_3/\text{Na}_2\text{O}$  becomes zero. The intensity of the band at 1077  $\text{cm}^{-1}$  is also increased. With increasing amount of  $\text{Na}_2\text{O}$  in the glass, the intensity of band around 630  $\text{cm}^{-1}$  was increased and the band in the range 550 - 450  $\text{cm}^{-1}$  begins to dominate, which results from

formation of  $[\text{BO}_4]$  tetrahedral of small degree of polymerization. The intensities of bands at 385, 360, 287, 1850 and 1975  $\text{cm}^{-1}$  are increased with increasing ratio  $\text{B}_2\text{O}_3/(\text{Na}_2\text{O} + 3\text{La}_2\text{F}_6)$  [171]. When barium is substituted for the sodium cation, the position of the peak at 540  $\text{cm}^{-1}$  is shifted to low frequency region and the intensity of band at 635  $\text{cm}^{-1}$  was decreased. The high frequency band around 1100  $\text{cm}^{-1}$  becomes strongest with sodium introduction. The substitution of Ba and Ca cations enhances the intensity of 960  $\text{cm}^{-1}$  and shifted towards the low frequency [172]. The Raman spectra of  $\text{Na}_2\text{O}-\text{B}_2\text{O}_3-\text{SiO}_2$  exhibit the bands at  $\sim 430, 485, 600, 800, 1060$  and  $1200 \text{ cm}^{-1}$ . The bands around 485 and 600  $\text{cm}^{-1}$  (known as defect lines  $\text{D}_1$  and  $\text{D}_2$ ) are assigned due to the ring breathing mode of four and three-membered silicate rings, respectively. Raman peaks at 770 and 805  $\text{cm}^{-1}$  become more significant for glasses with higher borate content [173]. In Raman spectra of  $30\text{Na}_2\text{O} \cdot 5\text{SiO}_2 \cdot 65[(1-x)\text{P}_2\text{O}_5 \cdot x\text{B}_2\text{O}_3]$  glasses, the intensity of bands at 680 and 770  $\text{cm}^{-1}$  decreases with increase of  $\text{B}_2\text{O}_3$ , which are attributed to the (P - O - P)<sub>sym</sub> stretching vibration and  $\text{BO}_4$  stretching in borate species, respectively. The bands at 1200 and 1350  $\text{cm}^{-1}$  are disappeared for  $x = 0.250$ . The intensity at 620  $\text{cm}^{-1}$  increases with an increase in  $\text{B}_2\text{O}_3$  content. The bands at 500, 780 and 1100  $\text{cm}^{-1}$  are appeared in  $\text{B}_2\text{O}_3$  rich glasses. The Raman band at 1160  $\text{cm}^{-1}$  is attributed to the symmetric stretching mode of the terminal oxygen on each tetrahedron. The intensity of band at 780  $\text{cm}^{-1}$  increases and assigned to  $\text{BO}_3$  stretching vibration [174]. The intensity of Raman bands at 810  $\text{cm}^{-1}$  is increased whereas decreased for band at 795  $\text{cm}^{-1}$  by heat treatment of glasses [175]. The S-O symmetrical stretching modes near 1000  $\text{cm}^{-1}$  from tetrahedral  $\text{SO}_4$  environments was observed. The peak near 300  $\text{cm}^{-1}$  is assigned to S - S stretch in the dithionate anion  $\text{S}_2\text{O}_6^{2-}$  [176-177].

The wettability of the diamond by the borosilicate glass melt at temperatures above 700  $^\circ\text{C}$  was enhanced due to the oxidation of the diamonds. The irregular pores become distinct in the diamond-borosilicate glass composites sintered above 800  $^\circ\text{C}$  [178]. The depolarization of silicate network takes place with increase of zirconia and causes to reduction of intensity of band due to Si - O - Si links. Some of the highly polymerized  $\text{Q}_3^{\text{Si}}$  and  $\text{Q}_4^{\text{Si}}$  units are replaced by the less polymerized  $\text{Q}_2^{\text{Si}}$  units in the glass structure. The most borate units within the glass structure are probably linked to silicate tetrahedral in range 650 - 800  $\text{cm}^{-1}$  and it is indicated by the lack of any prominent peaks in this range. The band near 703  $\text{cm}^{-1}$  is related to the zirconia network (Zr - O stretch) [179]. The structure of  $\text{Eu}^{3+}$  ions doped  $\text{MgO}-\text{PbO}-\text{B}_2\text{O}_3-\text{SiO}_2-\text{Nd}_2\text{O}_3$  glasses consist of bands centred at 1510, 1385, 955, 842, 732, 645, 487 and 298  $\text{cm}^{-1}$  (Fig. 13). The last band near 298  $\text{cm}^{-1}$  was attributed to Ln - O - Ln (Nd, Eu) stretching vibrations. The position of first band was shifted towards higher wavenumber side with increase of doping concentration of Eu. The bands at 605 and 527  $\text{cm}^{-1}$  are related to B - O - Si linkages along with  $\text{PbO}_4$  units and Si - O - Pb, respectively. The polymerization of the glass network take place due to the replacement of bonds B - O - B, Si - O - Si with more resistant B - O - Si, Si - O - Pb bonds with addition of  $\text{Eu}^{3+}$  ions [180]. The structural units of  $\text{TiO}_4$  exist in  $\text{CaO}-\text{B}_2\text{O}_3-\text{SiO}_2-\text{TiO}_2$  glasses and the degree of depolarization decreases with increase of  $\text{CaO}/\text{SiO}_2$  ratio. The increase in  $\text{B}_2\text{O}_3$  leads to increase of  $\text{BO}_4$

tetrahedral units in glass matrix. The addition of  $\text{TiO}_2$  in this glass system introduces two peaks near  $840$  and  $726 \text{ cm}^{-1}$  whereas the existed bands at  $500 - 800 \text{ cm}^{-1}$  are disappeared. The bands near  $948$ ,  $1040$  and  $1300 - 1500 \text{ cm}^{-1}$  are shifted towards low frequency side. The bands at  $840$  and  $726 \text{ cm}^{-1}$  are the characteristic of  $\text{Ti} - \text{O} - \text{Si}$  or  $\text{Ti} - \text{O} - \text{Ti}$  vibrations and deformation of  $\text{O} - \text{Ti} - \text{O}$  or  $\text{O} - (\text{Si}, \text{Ti}) - \text{O}$  in chain or sheet units or both [181].

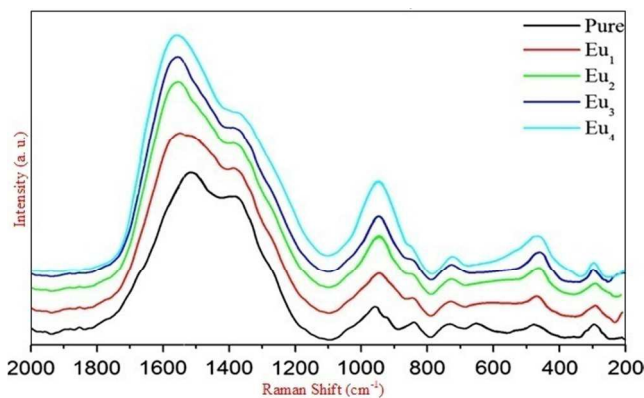


Fig. 13: Raman spectra of  $\text{MgO-PbO-B}_2\text{O}_3\text{-SiO}_2\text{-Nd}_2\text{O}_3\text{-Eu}_2\text{O}_3$  glasses [180]. (Reproduced with permission of Elsevier)

The structure of borosilicate glass in presence of modifier  $\text{BaO}$  also contains the network of borate and silicate. The addition of  $\text{BaO}$  decreases the intensity of band near  $816 \text{ cm}^{-1}$ . A new band at  $780 \text{ cm}^{-1}$  is incorporated at  $42 \text{ mol} \%$  of  $\text{BaO}$ . The bands at  $\sim 843$ ,  $877$ ,  $934$ ,  $1000$  and  $1040 \text{ cm}^{-1}$  are shifted to the lower frequency region and cause to low network polymerization after formation of  $\text{BO}_4$  units due to the increment of  $\text{BaO}$  content instead of  $\text{SiO}_2$  [182]. The position of peak at  $675 \text{ cm}^{-1}$  is shifted towards higher wavenumber side with increasing ratio of  $\text{BaO/SrO}$ . This is associated with variation of ionic radii of  $\text{Ba}$  and  $\text{Sr}$ .  $\text{La}_2\text{O}_3$  and  $\text{Fe}_2\text{O}_3$  in barium strontium titanate (BST) borosilicate glasses changes the peak positions of  $675$  and  $820 \text{ cm}^{-1}$  towards higher wavenumber side [183-185]. Raman band about  $735 \text{ cm}^{-1}$  occurs due to metaborate groups and symmetric breathing vibrations of  $\text{BO}_3$  triangles replaced by boron tetrahedral ( $\text{BO}_4$ ). Raman shift has been shifted towards high wavenumber side with increasing the doping concentration of  $\text{Fe}_2\text{O}_3$  [52]. Raman spectra of  $\text{AO-SiO}_2\text{-B}_2\text{O}_3\text{-La}_2\text{O}_3$ , ( $\text{A} = \text{Mg}, \text{Ca}, \text{Sr}, \text{Ba}$ ) glasses showed bands near  $770 - 779$  and  $1274 - 1280 \text{ cm}^{-1}$ , which are related to borate network. The bands near  $796 - 800$  and  $1060 - 1070 \text{ cm}^{-1}$  caused to silicate network. Tetrahedral coordination occurs in both silicon and boron in this glass system. Lanthanum is occupied interstitial octahedral co-ordination in the local network. The modifiers, such as  $\text{Ba}$ ,  $\text{Ca}$ ,  $\text{Mg}$ ,  $\text{Sr}$  and  $\text{La}$ , break the  $\text{B} - \text{O}$  or  $\text{Si} - \text{O}$  bonds and occupy interstitial or substitutional position in glass network. As the ionic radii of modifier ion increase in order  $\text{Mg}^{2+} > \text{Ca}^{2+} > \text{Sr}^{2+} > \text{Ba}^{2+}$ , then polarization follow the Fajan's rule, i.e. weak polarization by  $\text{Mg}$  and strong polarization by  $\text{Ba}$  [186]. Raman spectra of lead bismuth borosilicate glasses show the bands at  $1582$  and  $1382 \text{ cm}^{-1}$  in high frequency region [187]. After irradiation, the peak at  $480 \text{ cm}^{-1}$  is shifted toward higher wave number side. This suggests that the decrease of the  $\text{Si} - \text{O} - \text{Si}$  bond angle is decreased. A new weak peak at  $605 \text{ cm}^{-1}$  is introduced by irradiation in glasses and assigned due to the  $\text{D}_2$

defect peak involving the three-membered rings of  $\text{SiO}_4$  tetrahedra. The weak band at  $630 \text{ cm}^{-1}$  is associated either with the hindered bending vibration of ring-type metaborate groups or the breathing mode of danburite-like rings. The  $\text{Q}_n^s$  species of silicon remain almost unchanged after Ar-irradiation [188]. The Raman band assignments of borosilicate glasses are listed in Table 4.

Table 4: Assignment of main Raman bands in the spectra of borosilicate glasses

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
1510-1570	B-O stretching modes involving with one NBO of $[\text{BO}_3]$ triangles and molecular oxygen stretching vibration modes	183-185
1385-1397	Stretching vibrations of B-O bond in $\text{BO}_4$ units from different borate groups	183-185
1160	symmetric stretching mode of the terminal oxygen on each tetrahedron	174
947-955	Stretching vibrations of Si-O bond and due some presence of $\text{BO}_4$ units	183-185
840-842	Pyroborate unit along with ortho-silicate ( $\text{SiO}_4$ ) <sup>4-</sup> composition or vibrations of Ti-O-Si or Ti-O-Ti structural units	181
808	boroxol rings	168
722-732	Chain-type metaborate groups containing NBO or deformation of O-Ti-O or O-(Si, Ti)-O in chain or sheet units or both	181
550-850	Ring breathing modes	167, 168
400 - 850	reedmergnerite $[\text{BSi}_3\text{O}_8]^-$ and danburite $[\text{B}_2\text{Si}_2\text{O}_8]^{2-}$	167
605	B-O-Si linkages along with $\text{PbO}_4$ units	180
590	symmetric oxygen breathing vibrations of three-membered siloxane rings	170
300-500	Mixed stretching and bending modes of Si-O-Si units	167
475-487	Si-O-Si, Si-O-B isolated vibrations	180

### 3.5. Borophosphate Glasses

Typical Raman spectra tin borophosphate glasses are shown in Fig. 14. In this study, it was found that the Raman bands centred at  $1330$ ,  $1270$ ,  $1050$ ,  $1000$ ,  $950$ ,  $760$ ,  $700 \text{ cm}^{-1}$  and some bands near  $500 \text{ cm}^{-1}$ . The vibrations at  $1330$  and  $700 \text{ cm}^{-1}$  are attributed to metaborate groups whereas the band near  $1000 \text{ cm}^{-1}$  is attributed to pyrophosphate or orthophosphate units. The presence of metaborate chains in glass matrix was found by the vibration at  $760 \text{ cm}^{-1}$ . The evidence of pyrophosphate units was also confirmed by peak located at  $1050 \text{ cm}^{-1}$ . The vibration at  $1270 \text{ cm}^{-1}$  is caused by orthoborate/pyroborate units [189]. Network of lithium borophosphate glasses, in presence of germanium oxide, consists of stretching vibrations of phosphate structural units in high ( $1000 - 1300 \text{ cm}^{-1}$ ) and mid frequency regions ( $500 - 700 \text{ cm}^{-1}$ ) and these units are modified by  $\text{Ge}$  ions.  $\text{Ge}$  free glasses showed a band at  $1159 \text{ cm}^{-1}$ , which is caused by metaphosphate structure. The vibration at  $1260 \text{ cm}^{-1}$  is related with asymmetric,  $\nu_{\text{as}}(\text{PO}_2)$ , stretching mode in  $\text{Q}_p^2$  units whereas the weaker band centred at  $1093 \text{ cm}^{-1}$  is attributed to  $\nu_s(\text{PO}_3)$ , in  $\text{Q}_p^1$  units. The bands at  $703$  and  $665 \text{ cm}^{-1}$  are assigned due to the symmetric stretching mode of P-O-P bridging bond in  $\text{Q}_p^1$  and  $\text{Q}_p^2$  structure, respectively. A broad band around  $320 \text{ cm}^{-1}$  is attributed to the cage vibrational frequencies of  $\text{Li}^+$  ions. The addition of  $5 \text{ mol.}\%$   $\text{GeO}_2$  shifted the band around  $1159$  to  $1106$

$\text{cm}^{-1}$  and also showed depolymerization of the  $\text{Q}_p^2$  structure by the formation of more  $\text{Q}_p^1$  phosphate units. The intensity of band at  $665 \text{ cm}^{-1}$  is reduced with increasing  $\text{GeO}_2$  and it is completely

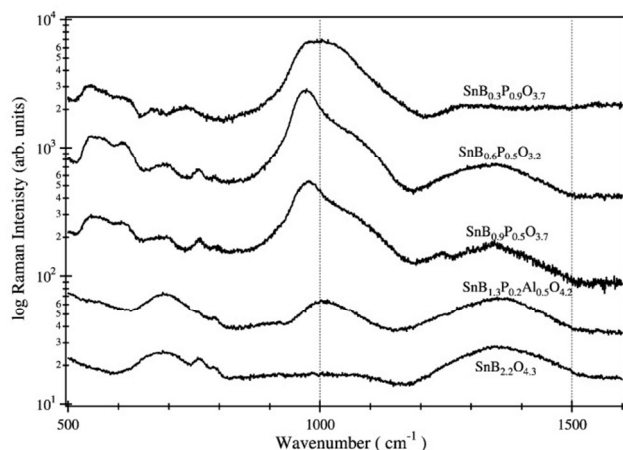


Fig. 14: Raman spectra of various tin borophosphate glasses [189]. (Reproduced with permission of Elsevier)

disappeared for 15 to 25 mol% of  $\text{GeO}_2$  whereas the intensity of bands at  $700$  and  $751 \text{ cm}^{-1}$  are enhanced due to more depolymerization of  $\text{Q}_p^1$  units by this addition. A new band at  $600 \text{ cm}^{-1}$  is introduced due to Ge - O - Ge bending modes or Ge - O - P stretching modes [190]. The four intense bands at  $650$ ,  $696$ ,  $880$  and  $1112 \text{ cm}^{-1}$  are found in Raman spectrum of  $75\text{P}_2\text{O}_5\text{-}20\text{B}_2\text{O}_3\text{-}4.9\text{Na}_2\text{O}\text{-}0.1\text{Er}_2\text{O}_3$  glass [191].

The Raman bands at  $636$  and  $350 \text{ cm}^{-1}$  in  $\text{Na}_2\text{O}\text{-Ga}_2\text{O}_3\text{-P}_2\text{O}_5$  glasses are attributed to  $\text{GaO}_4$  and  $\text{GaO}_6$  groups, respectively. The band centred at  $515$  and  $540 \text{ cm}^{-1}$  are caused by Ga - O - Ga bonds between  $\text{GaO}_4$  tetrahedra and the band at  $675 \text{ cm}^{-1}$  is associated with the vibrations involved with non-bridging oxygens from the  $\text{GaO}_4$  tetrahedra. The band at  $\sim 540 \text{ cm}^{-1}$  was attributed to  $[\text{GaO}_4]^-$ . In Ti containing glasses, a intense band near  $900 \text{ cm}^{-1}$  is caused by  $[\text{TiO}_4]$  groups [192]. Raman spectra of calcium borophosphate consisted by networks of ortho  $\text{Q}_p^0$  and pyro  $\text{Q}_p^1$  phosphate units. With addition of 20 mol% of  $\text{B}_2\text{O}_3$  in calcium phosphate glass, the band  $1045 \text{ cm}^{-1}$  is shifted to  $1038 \text{ cm}^{-1}$  and new bands at  $1000$  and  $675 \text{ cm}^{-1}$  are also introduced [193]. In Raman spectra of  $(1-x)[0.5\text{K}_2\text{O}\cdot 0.1\text{B}_2\text{O}_3\cdot 0.4\text{P}_2\text{O}_5]\cdot x\text{Nb}_2\text{O}_5$  glasses, a strong band in region  $907 - 911 \text{ cm}^{-1}$  was observed and the band of phosphate groups are assigned at  $1070 - 1148 \text{ cm}^{-1}$  and  $1212 - 1244 \text{ cm}^{-1}$ . The new bands at  $811 - 817$  and  $638 \text{ cm}^{-1}$  are introduced in glasses for more than 16.7 mol%  $\text{Nb}_2\text{O}_5$ . The vibrations centred at  $907 \text{ cm}^{-1}$  are the characteristic of symmetric stretching vibration of isolated  $\text{NbO}_6$  octahedra. The dominating bands at  $817$  and  $638 - 648 \text{ cm}^{-1}$  are assigned due to Nb - O - Nb bonds in glass with 37.5 mol% of  $\text{Nb}_2\text{O}_5$ . The bands at  $820 - 840 \text{ cm}^{-1}$  in barium-potassium niobato-phosphate glasses are attributed to the deformation vibrational modes of Nb - O - Nb bridges between  $\text{NbO}_6$  octahedra [194].

Raman spectroscopic studies of  $50\text{PbO}\cdot 10\text{B}_2\text{O}_3\cdot 40\text{P}_2\text{O}_5\cdot x\text{TiO}_2$  glasses showed that major bands are located at  $329$ ,  $729$ ,  $936$ ,  $1100$  and  $1210 \text{ cm}^{-1}$  (Fig. 15). The strong peak at  $1100$  and  $1210 \text{ cm}^{-1}$  are shifted towards lower wavenumber side whereas the peak at  $729$  and  $936 \text{ cm}^{-1}$  are shifted towards higher wavenumber

side with increasing the content of  $\text{TiO}_2$  [195].  $\text{TeO}_2$  in lead borophosphate glasses form the structural units of  $\text{TeO}_3$  and  $\text{TeO}_4$  at its low and high concentrations, respectively. The intensity of band at  $1092 \text{ cm}^{-1}$  decreases with  $\text{TeO}_2$  and also shifted towards low frequency side. This change is associated with depolymerisation of the phosphate network. The bands at  $550$  and  $850 \text{ cm}^{-1}$  are attributed to the vibration of oxygen atoms in P - O - P bridges and are gradually replaced by the band characteristic of different Te - O vibrations. New bands in the region  $453 - 458 \text{ cm}^{-1}$  are ascribed due to bending vibrations of O - Te - O and Te - O - Te linkages. The bands lying at  $620 - 630 \text{ cm}^{-1}$ ,  $660 - 670 \text{ cm}^{-1}$  are related to stretching vibrations of  $\text{TeO}_4$  trigonal bipyramids, whereas, the vibrations at  $720 - 750 \text{ cm}^{-1}$  and  $760 - 800 \text{ cm}^{-1}$  are associated with stretching vibrations of  $\text{TeO}_3$  trigonal pyramids [196].

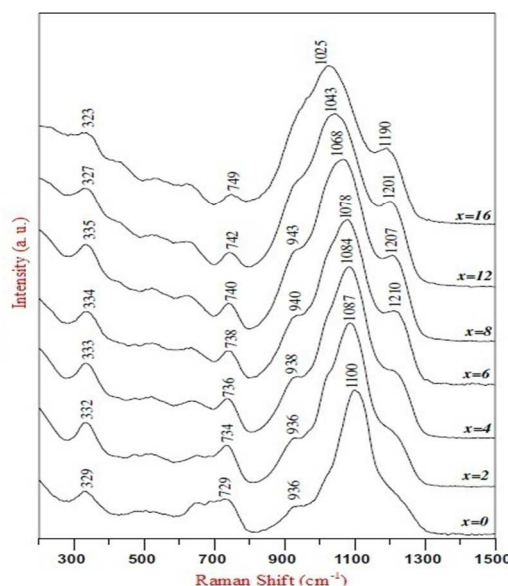


Fig. 15: Raman spectra of the glass series  $50\text{PbO}\text{-}10\text{B}_2\text{O}_3\text{-}40\text{P}_2\text{O}_5 + x\text{TiO}_2$  [196]. (Reproduced with permission of Elsevier)

Zinc oxide plays an important role in modification of borate and phosphate structural units. In the spectra, a sharp band assigned at  $808 \text{ cm}^{-1}$  is the characteristic of  $\text{B}_2\text{O}_3$  and the band at  $968 \text{ cm}^{-1}$  is caused by vibrations of isolated  $\text{PO}_4$  units in the structural network of borophosphate glasses [197]. If  $\text{TiO}_2$  was added in zinc borophosphate glasses, then, the intense bands at  $1162$  and  $666 \text{ cm}^{-1}$  are shifted towards to lower frequency side whereas, the band at  $747 \text{ cm}^{-1}$  is shifted to  $774 \text{ cm}^{-1}$  successively. The band at  $506 \text{ cm}^{-1}$  was disappeared above 4 mol% of  $\text{TiO}_2$ . The shifts of bands are associated with distortion of borate and phosphate units. The shift in sodium contained glasses is lesser as compared to Zn modified glasses due to difference in ionic field strength [198]. All the Raman bands are shifted towards lower wavenumber side with substitution of  $\text{Sb}_2\text{O}_3$  instead of  $\text{TiO}_2$  in above glass system. The vibrations of structural units containing Sb - O are assigned in the range of  $350 - 700 \text{ cm}^{-1}$ .  $\text{Sb}_2\text{O}_3$  modifies the network in the form of  $\text{SbO}_3$  units with single atom of Sb. The low concentration of  $\text{Sb}_2\text{O}_3$  forms the isolated  $\text{SbO}_3$  unit. The existence of three bands in place of two is associated with the splitting of the symmetry of  $\text{SbO}_3$  units in the glass network.  $\text{Sb}_2\text{O}_3$  content replicate the depolymerisation of phosphate chains. Raman bands in the range of  $520 - 690 \text{ cm}^{-1}$  are



attributed to  $\text{SbO}_3$  pyramids.  $\text{SbO}_3$  units linked into chains and clusters with Sb - O - Sb bridges at higher  $\text{Sb}_2\text{O}_3$  content and are confined by bands in the region  $380 - 520 \text{ cm}^{-1}$  [199].  $\text{MoO}_3$  weakened the bond strength for zinc borophosphate glasses. The presence of the polarized vibrational band at  $976 \text{ cm}^{-1}$  is associated with  $\text{MO}_x$  symmetric stretching vibrations and the depolarized band at  $878 \text{ cm}^{-1}$  is ascribed due to the Mo - O - Mo stretching vibration. The incorporation of molybdate units in glass matrix is responsible for depolymerization of phosphate chains and the formation of P - O - Mo bonds [200]. The tellurium oxide in above glass system forms structural units of  $\text{TeO}_3$ ,  $\text{TeO}_{3+1}$  and  $\text{TeO}_4$ . The phosphate units are depolymerized by incorporation of  $\text{TeO}_2$  and also boron coordination is modified. When the ratio of  $\text{B}_2\text{O}_3/\text{P}_2\text{O}_5$  increases, then  $\text{TeO}_4$  units are replaced by  $\text{TeO}_3$  units as the number of oxygen atoms in the glass is decreased [201-204]. The peak assignment of Raman spectra of borophosphate glasses are listed in Table 5.

Table 5: Assignment of Raman bands in the spectra of borophosphate glasses.

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
1260, 1380	symmetric stretching of P - O	123
1210	asymmetric stretching vibrations of $\text{PO}_3$ groups	123
1170	symmetric stretching of a non bridging oxygen on a $\text{Q}_p^2$ tetrahedron	123
1080	$[\text{P}_2\text{O}_7]^{2-}$ symmetric groups in $\text{Q}_p^1$ pyrophosphates	140
966	terminal $[\text{PO}_3]^-$ symmetric groups in $\text{Q}_p^1$ pyrophosphates	140
960	symmetric stretching of the orthophosphate groups $\text{PO}_4^{3-}$	123
600	Ge - O - Ge bending modes or Ge - O - P stretching modes	190
550, 850	vibration of oxygen atoms in P - O - P bridges	196
968	vibrations of isolated $\text{PO}_4$ units	199

### 3.6 Aluminosilicate Glasses

Raman spectra of glasses for glass system  $5\text{Na}_2\text{O} \cdot 20\text{CaO} \cdot 5\text{Al}_2\text{O}_3 \cdot (60 - x)\text{SiO}_2 \cdot x\text{ZnO}$  (where  $x = 0, 4, 7$  and  $10 \text{ mol}\%$ ) showed the bands by deconvolution, which are centred at  $868, 956, 1020$  and  $1075 \text{ cm}^{-1}$  and attributed to the  $\text{Si}\backslash\text{NBO}$  stretching modes of mainly silicate units,  $\text{Q}_{\text{Si}}^n$  (Fig. 16). The band near  $868 \text{ cm}^{-1}$  is the characteristic of  $\text{Q}_{\text{Si}}^1$ . The peak at  $956 \text{ cm}^{-1}$  is assigned due to  $\text{Q}_{\text{Si}}^2$  whereas peak at  $1019 \text{ cm}^{-1}$  is the cause of asymmetric stretching of BO and related to  $\text{Q}_{\text{Si}}^1 - \text{Q}_{\text{Si}}^3$  units. The band at  $1020 \text{ cm}^{-1}$  is associated with  $\text{Q}_{\text{Si}}^3$ . Raman bands in wavenumber range  $400 - 800 \text{ cm}^{-1}$  are assigned at  $412, 544, 622$  and  $779 \text{ cm}^{-1}$ . The vibrations in the range  $400 - 650 \text{ cm}^{-1}$  are caused by bending vibrations of the bridging oxygen (BO) bonds of  $\text{SiO}_4$  and band near  $779 \text{ cm}^{-1}$  is caused by the Si - O - Si network,  $\text{ZnO}_4$  tetrahedra and  $\text{AlO}_4$  units with three BOs and one NBO. The low frequency bands located at  $222, 279$  and  $338 \text{ cm}^{-1}$  are related with cations of modifying networks. The vibration of cations  $\text{Na}^+$  and  $\text{Ca}^{2+}$  are centred at  $222$  and  $279 \text{ cm}^{-1}$ , respectively. The coupled vibration of Ca - O and  $\text{SiO}_4$  network are found at  $338 \text{ cm}^{-1}$ . The addition of ZnO in glass matrix shifts

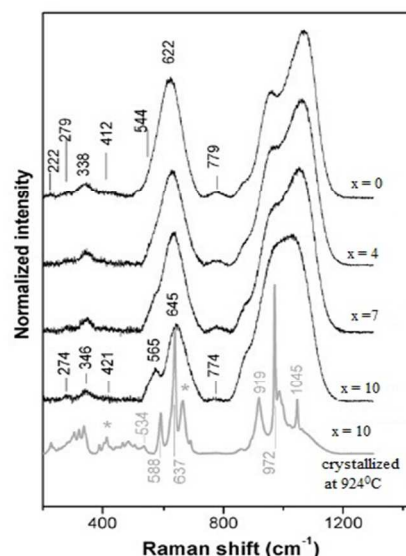


Fig. 16: Raman spectra of glasses for glass system  $5\text{Na}_2\text{O} \cdot 20\text{CaO} \cdot 5\text{Al}_2\text{O}_3 \cdot (60 - x)\text{SiO}_2 \cdot x\text{ZnO}$  [203]. (Reproduced with permission of Elsevier)

the bands at  $279$  and  $338 \text{ cm}^{-1}$  to position  $274$  and  $346 \text{ cm}^{-1}$ , respectively. With increasing ZnO/ $\text{SiO}_2$  molar ratio, the band at  $622 \text{ cm}^{-1}$  is shifted towards  $645 \text{ cm}^{-1}$  and finally, split into two components for  $10 \text{ mol}\%$  ZnO content glasses [203].

Sodium silicate and aluminosilicate can be analyzed by Raman spectroscopy in low frequency region of  $20 - 250 \text{ cm}^{-1}$ . In this region, the peak at  $69 \text{ cm}^{-1}$  is the signature of Raman scattering involving in rotational motions of interconnected tetrahedral units. The position of this peak is shifted to  $78 \text{ cm}^{-1}$  with addition of  $\text{Al}_2\text{O}_3$  upto  $6 \text{ mol}\%$ . Raman peak at  $170 \text{ cm}^{-1}$  is assigned due to narrower distribution of T - O - T angles, or by a structural localization of those vibrations in specific clusters of molecules in the glass network. The band centred at  $490 \text{ cm}^{-1}$  is shifted to  $482 \text{ cm}^{-1}$ . The band at  $597 \text{ cm}^{-1}$  becomes broader with increase of content of  $\text{Al}/(\text{Al} + \text{Na})$ . This result suggests the presence of three-membered rings in the glass network. Raman band located at  $781 \text{ cm}^{-1}$  is shifted to  $799 \text{ cm}^{-1}$  with replacement of  $\text{Na}_2\text{O}$  by  $\text{Al}_2\text{O}_3$ . A shoulder near  $810 - 820 \text{ cm}^{-1}$  is the feature of Si - O stretching involving oxygen motions in the Si - O - Si plane or to the motion of the Si atom in its oxygen cage or threefold - degenerate rigid cage vibrational mode of  $\text{TO}_2$  units. The high frequency bands at  $960, 1070, 1100, 1150$  and  $1200 \text{ cm}^{-1}$  are the characteristic of silicate network and vibrational modes of  $\text{TO}_4$  tetrahedra [204]. A Raman band at  $495 \text{ cm}^{-1}$  was observed in albite glass and showed the presence of polymeric structure in such glasses. The band assigned at  $595 \text{ cm}^{-1}$  is the characteristic of structural defect in the glass involving T - O (non-bridging) vibrations (Fig. 17). The defect line,  $\text{D}_2$ , of  $\text{SiO}_2$  is assigned at  $606 \text{ cm}^{-1}$ . The Al in borosilicate glass shifted the band positioned at  $950 \text{ cm}^{-1}$  to  $900 \text{ cm}^{-1}$  while peaks at  $980$  and  $1110 \text{ cm}^{-1}$  showed opposite trend with previous result. The polymerization become more significant with increasing Al/Na ratio [205].

The structure of aluminosilicate glasses consists of silicate and aluminate networks. The band centred at  $1060$  and  $1200 \text{ cm}^{-1}$  are attributed to Si - O stretching vibrations of  $\text{SiO}_4$  tetrahedra. The intensities of bands at  $1000$  and  $1100 \text{ cm}^{-1}$  are increased with increase of  $\text{Al}_2\text{O}_3$ . This increase in the intensity of these bands is

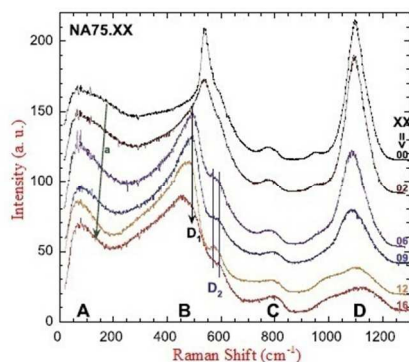


Fig. 17: Raman spectra of  $75\text{SiO}_2-(25-x)\text{Na}_2\text{O}-x\text{Al}_2\text{O}_3$  glasses [204]. (Reproduced with permission of Elsevier)

due to interlinking of Si-O-Al network. The intensity of band at  $435\text{ cm}^{-1}$  was decreased with increase of  $\text{Al}_2\text{O}_3$  and suggested to decrease of Si - O - Si linkage in tetrahedral units [206]. Aluminosilicate glasses along with silica-calcium, silica sodium aluminate and silica-potassium aluminate showed the appearance of bands centred at  $1120$ ,  $1100$ ,  $930$  and  $890\text{ cm}^{-1}$ . The degree of metal cations induced network clustering increases in order of  $\text{K} < \text{Na} < \text{Li} < \text{Ca} < \text{Mg}$  within the  $\text{SiO}_2\text{-MAlO}_2$  glasses ( $M = \text{K}, \text{Na}, \text{Li}, \text{Ca}$  and  $\text{Mg}$ ) [207-208]. The low frequency bands at  $450$ ,  $500$  and  $600\text{ cm}^{-1}$  are associated with motions of bridged oxygen in T - O - T linkages and band at  $560\text{ cm}^{-1}$  in Al rich sample is caused by presence of Al - O - Al bridges in  $\text{CaO-Al}_2\text{O}_3\text{-SiO}_2$  system. The intensity of band at  $800\text{ cm}^{-1}$  was decreased with decrease of  $\text{SiO}_2$  and shifted towards lower wavenumber side. The high frequency Raman spectra consist of peaks at  $1000$  and  $1080\text{ cm}^{-1}$ . Al coordination exists mainly in four-fold type and varies between  $\text{Q}_{\text{Si}}^2$  and  $\text{Q}_{\text{Si}}^4$  species as a function of  $\text{MO}/\text{Al}_2\text{O}_3$  ratio ( $M = \text{Ca}, \text{Sr}, \text{Ba}$ ). At very higher concentration of  $\text{CaO}$  (more than  $62\text{ mol}\%$ ) a Raman band between  $300 - 400\text{ cm}^{-1}$  was also present. A Raman band at  $170\text{ cm}^{-1}$  and a shoulder around  $300\text{ cm}^{-1}$  was observed in Ba content glasses. The intensity of peak at  $170\text{ cm}^{-1}$  increases with Ba content. The peak position of Raman bands due to metal cation decreases with size of cation [209-212]. Raman band located at  $450\text{ cm}^{-1}$  is the characteristic of motion of bridging oxygen in the plane perpendicular to the Si - O - Si (Al) bond. The vibration at  $710\text{ cm}^{-1}$  gives the evidence of  $\text{AlO}_4$  tetrahedra and the band at  $760\text{ cm}^{-1}$  corresponds to the stretching mode of the Al - O bond with aluminium in four-fold coordination. The band at  $960\text{ cm}^{-1}$  in the spectra is caused by stretching vibration of silicon-oxygen tetrahedra with two corners shared with aluminium-oxygen or calcium-oxygen polyhedral [213]. The depolarization ratio in calcium aluminosilicate glasses is  $0.007$  and  $0.485$  for intense bands at  $540$  and  $96\text{ cm}^{-1}$ , respectively. This result showed that the vibration at  $540\text{ cm}^{-1}$  is highly symmetric in nature. The polarized spectra for this glass system showed peaks centred at  $96$ ,  $540$ ,  $767$  and  $867\text{ cm}^{-1}$ . Similar bands are also observed in other glass system with a slight variation in their peak positions [214]. The main peak assignment in the Raman spectra of Aluminosilicate glasses are listed in Table 6.

Table 6: Assignment of Raman bands in the spectra of Aluminosilicate glasses

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
1060, 1200	Si - O stretching vibrations of $\text{SiO}_4$ tetrahedra	206
960	stretching vibration of silicon-oxygen tetrahedra with two corners shared with aluminium-oxygen or calcium-oxygen polyhedral	213
810 - 820	Si - O stretching involving oxygen motions in the Si - O - Si plane or to the motion of the Si atom in its oxygen cage or threefold-degenerate rigid cage vibrational mode of $\text{TO}_2$ units	204
779	Si - O - Si network, $\text{ZnO}_4$ tetrahedra and $\text{AlO}_4$ units with three BOs and one NBO	203
450, 500, 600	motions of bridged oxygen in T - O - T linkages	209-212
400 - 650	bending vibrations of the bridging oxygen (BO) bonds of SiO	203

### 3.7 Phosphosilicate Glasses

A typical Raman spectrum of phosphosilicate glass is shown in Fig. 18 [215]. Raman spectra of phosphosilicate glasses consist of networks of phosphate as well as silicate. Raman peak at  $1320\text{ cm}^{-1}$  is attributed to stretching vibrations of O - P double bonds and the bands below  $860\text{ cm}^{-1}$  are assigned due to Si - O and P - O bonds in Si - O - Si, P - O - Si and P - O - P linkages. The peaks at  $1200$  and  $1020\text{ cm}^{-1}$  are caused by P - O - P linkages and band near  $1145\text{ cm}^{-1}$  is due to P - O - Si linkages [215-216]. The bands are located at  $340$ ,  $426$ ,  $580$ ,  $648$ ,  $720$ ,  $780$ ,  $882$ ,  $956$ ,  $1040$  and  $1450\text{ cm}^{-1}$  in the Raman spectra of various compositions of  $\text{CaO-MgO-SiO}_2$  system doped with  $\text{B}_2\text{O}_3$ ,  $\text{P}_2\text{O}_5$ ,  $\text{Na}_2\text{O}$  and  $\text{CaF}_2$  (Fig. 19).

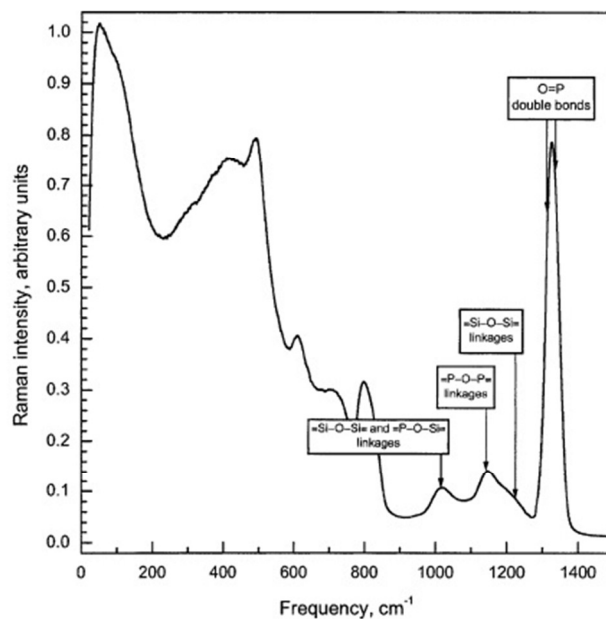


Fig. 18: Raman spectrum of phosphosilicate glass [215]. (Reproduced with permission of Elsevier)

The vibrations in the regions  $800 - 1300\text{ cm}^{-1}$  are the characteristic of the asymmetric vibration of  $\text{SiO}_4$  tetrahedra whereas the bands

at 952, 590 and 425  $\text{cm}^{-1}$  are assigned due to the symmetric

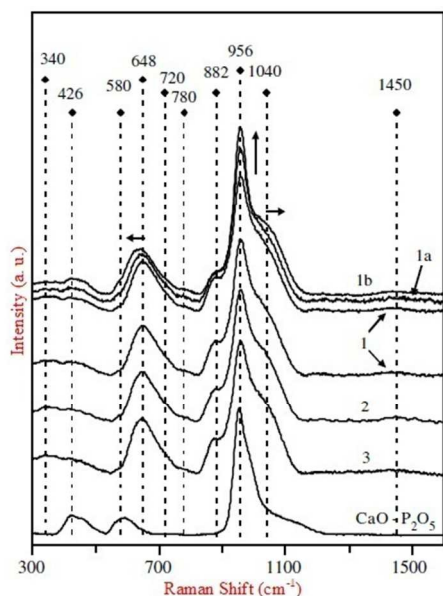


Fig. 19: Raman spectra of calcium magnesium silicate glasses doped with  $\text{B}_2\text{O}_3$ ,  $\text{P}_2\text{O}_5$ ,  $\text{Na}_2\text{O}$  and  $\text{CaF}_2$  [217]. (Reproduced with permission of Elsevier)

stretching of the P - O bonds and O - P - O bending modes of the orthophosphate  $\text{PO}_4^{3-}$  unit ( $\text{Q}_2^0$ ), respectively [217]. The silicate networks Si - O - Si and intensity of the Si - O - NBO stretching mode are modified by cations Na, Mg and Ca [218-219]. The addition of silver in calcium phosphosilicate glasses changes some networks such as the contribution of  $\text{Q}_2^0$ ,  $\text{Q}_3^0$  and  $\text{Q}_4^0$  is progressively vanishing by increasing the concentration of silver and a weak band around 1025  $\text{cm}^{-1}$  was also assigned. The position of band at 956  $\text{cm}^{-1}$  was shifted at 967  $\text{cm}^{-1}$  with this addition. The spectra also show the occurrence of a depolymerization process of  $\text{SiO}_4$  network [220]. The average distance between silicates, phosphates, inter silicate-phosphate and aluminium-phosphate chains was increased by increasing the concentrations of  $\text{AlO}_6$  structural units and it leads to the increase in the average bond length of Tm - O and Er - O due to weaker field around Tm - O and Er - O ions [221]. The main peak assignment in the Raman spectra of Phosphosilicate glasses are listed in Table 7.

Table 7: Assignment of Raman bands in the spectra of Phosphosilicate glasses

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
1320	stretching vibrations of O - P double bond	215-216
1200, 1020	P - O - P linkages	215-216
1145	P - O - Si linkages	215-216
860	Si - O and P - O bonds in Si - O - Si, P - O - Si and P - O - P linkages	215-216

### 3.8 Alumino-borosilicate Glasses

Raman spectra of NaF alumino-borosilicate glass consist of vibrations located at 1075, 947, 802, 723 and 480  $\text{cm}^{-1}$  (Fig. 20). The Raman band located at 802  $\text{cm}^{-1}$  is the characteristic of Si - O - Si vibration and symmetric breathing vibration of the boroxol

rings. Sodium fluoride (NaF) in alumino-borosilicate glass increases the stability of boroxol ring. The metaborate group is assigned at 723  $\text{cm}^{-1}$  whereas lower intensity bands centered at 1075, 947 and 480  $\text{cm}^{-1}$  are attributed to the stretching vibration of Si - O - Si and B - O - Si. Instead of above bands, few more bands at 1114, 1000, 885, 770, 657, 598, 520, 491 and 447  $\text{cm}^{-1}$  are also assigned in Raman spectrum. The band located at 1114  $\text{cm}^{-1}$  is attributed to diborate groups. The band at 885  $\text{cm}^{-1}$  indicates the presence of pyroborate groups. The existence of symmetric breathing vibration of six-membered rings  $\text{BO}_4$  tetrahedron was confirmed by band located at 770  $\text{cm}^{-1}$ . The bands at 1000 and 520  $\text{cm}^{-1}$  indicate the availability of the vibration of Si - O - B - O - B, B - O - Si and Si - O - Si bending. The rocking vibration is also found in network of glass at 475  $\text{cm}^{-1}$ . Raman band at 951  $\text{cm}^{-1}$  is attributed to B - O - Si stretching vibration. The shoulder peak at 447  $\text{cm}^{-1}$  is assigned due to the bending or rocking vibrations of the B - O - Si linkages, B - O - B, B - O - Si and Si - O - Si [222].

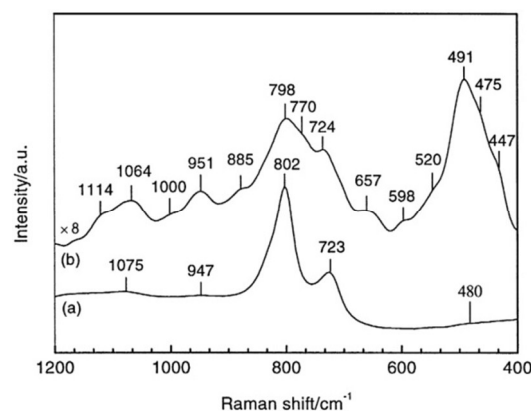


Fig. 20: Raman spectra of NaF alumino-borosilicate glasses [222]. (Reproduced with permission of Elsevier)

Raman spectroscopy of sodium-alumino-borosilicate glasses in presence of  $\text{Gd}_2\text{O}_3$  showed that the band at 1050  $\text{cm}^{-1}$  is shifted towards lower wavenumbers side with addition of  $\text{Gd}_2\text{O}_3$  and intensity of band centred at 450  $\text{cm}^{-1}$  was increased. One more band at 300  $\text{cm}^{-1}$  was ascribed for the glasses with 7.06 and 13.58 mol% of  $\text{Gd}_2\text{O}_3$ . The borate network was much influence by network modifying cations as compared to silicate network. Gd cations are more effective than Na cations for modifications of network [223]. Raman spectra of strontium alumino-borosilicate glasses exhibit vibrations of borate as well as silicate network along with some linking networks. The asymmetric ring breathing vibrations is assigned at 430  $\text{cm}^{-1}$ , symmetric stretching vibrations of Si - O - Si at 800  $\text{cm}^{-1}$ , stretching vibrations associated to  $\text{SiO}_4$  composed of sites with three and four-fold rings ranging 1000 - 1200  $\text{cm}^{-1}$  and symmetric stretching vibrations at 1200  $\text{cm}^{-1}$ . The borate networks are pure  $\text{B}_2\text{O}_3$  exhibit a strong band at 805  $\text{cm}^{-1}$  due to symmetric beating vibrations of boroxol ring oxygen. The boroxol ring structure is splitted into different ring structures (i.e. triborate, tetraborate or pentaborate, etc.) in the presence of suitable modifiers in the glass network (Fig. 21). After the irradiation of  $\gamma$ -ray, the intensity of bands due to asymmetric vibrations was increased and shifted to higher wavenumber side. The effect of radiation was lesser in  $\text{Al}_2\text{O}_3$  added glasses [224]. Not only  $\gamma$ -ray irradiation but also  $\beta$ -

ray changes the network of the glasses. The bending vibration of Si - O - Si at  $460\text{ cm}^{-1}$  is shifted in the glasses irradiated with  $10^9$  Gy.  $\beta$ -ray radiation decreases the  $Q_{\text{Si}}^2$  species as compared to  $Q_{\text{Si}}^3$  ones [225-226]. The  $\beta$ -ray irradiation in rare earth aluminoborosilicate glasses reduces the Cr and Mn ions from the matrix. The Raman band at  $460\text{ cm}^{-1}$  is shifted to  $480\text{ cm}^{-1}$  with adding the Cr instead of Sm and Gd. Silicate networks in the range  $900\text{--}1200\text{ cm}^{-1}$  are not much affected by Cr whereas new bands centered at  $845$  and  $900\text{ cm}^{-1}$  are introduced in the matrix. The behaviour of Raman band at  $\sim 1550\text{ cm}^{-1}$  due to molecular oxygen could not be analyzed under irradiation due to polishing of glass surface [226-229]. Similar to  $\beta$ -ray irradiation, the  $\gamma$ -irradiation is effective for reducing  $\text{Sm}^{3+}$  to  $\text{Sm}^{2+}$  ions from  $10\text{ kGy}$  dose [230]. The Raman band assignment in Aluminoborosilicate glasses are listed in Table 8.

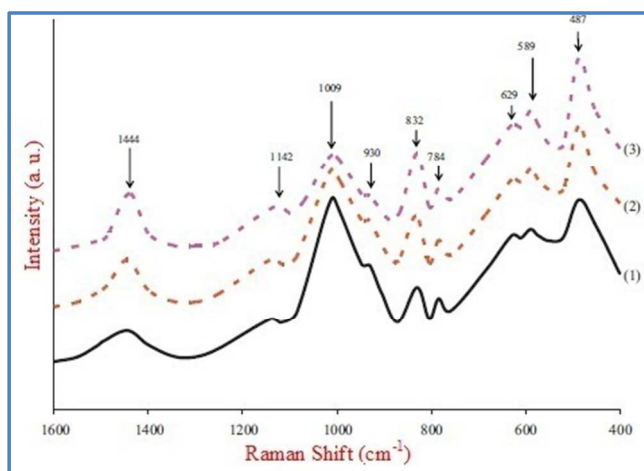


Fig. 21: Raman spectra of glass  $40\text{SrO}-5\text{Al}_2\text{O}_3-15\text{B}_2\text{O}_3-40\text{SiO}_2$  irradiated with doses of (1)  $0\text{ kGy}$ ; (2)  $10\text{ kGy}$ ; and (3)  $30\text{ kGy}$  [224]. (Reproduced with permission of Elsevier)

Table 8: Assignment of Raman bands in the spectra of Aluminoborosilicate glasses

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
1114	diborate groups	222
800 -1300	asymmetric vibration of $\text{SiO}_4$ tetrahedra	222
1000, 520	vibration of Si - O - B - O - B, B - O - Si and Si - O - Si bending	222
952	symmetric stretching of the P - O - bonds	222
951	B - O - Si stretching vibration	222
885	pyroborate groups	222
770	symmetric breathing vibration of six-membered rings one $\text{BO}_4$ tetrahedron	222
590, 425	O - P - O bending modes of the orthophosphate $\text{PO}_4^{3-}$ unit	222
447	bending or rocking vibrations of the B - O - Si linkages, B - O - B, B - O - Si and Si - O - Si	222

### 3.9 Tellurite Glasses

The Raman spectra of various tellurite glasses are shown in Fig. 22. In tellurite glasses, the tellurium oxide works as glass former. The tellurite oxide has mainly three basic structures,

namely  $\text{TeO}_4$  (trigonal bipyramid,  $\text{tbp}$ ),  $\text{TeO}_3$  (pyramid) and an intermediate with  $\text{TeO}_{3+\delta}$  polyhedral. The four oxygen atoms are covalently bonded with central tellurium atom in  $\text{TeO}_4$ . The bipyramid structure is formed by two equatorial and two apex oxygen sites. The trigonal pyramid structure consisted of two bridging oxygen sites and one non-bridging oxygen atom [231-232]. The Raman bands near  $667\text{ cm}^{-1}$  are related with combined vibrations of asymmetric stretching of  $\text{Te}_{\text{eq}}\text{O}_{\text{ax}}-\text{Te}$  bonds and symmetric stretching of  $\text{TeO}_4$   $\text{tbp}$ s. The addition of  $\text{Na}_2\text{O}$  in glassy matrix changes the intensity of bands due to stretching vibrations of non-bridging  $\text{Te}-\text{O}^-$  bonds in  $\text{TeO}_3$   $\text{tp}$ s ( $753$  and  $792\text{ cm}^{-1}$ ). The position of band assigned to the symmetric bending vibration of  $\text{TeO}_4$   $\text{tbp}$ s is shifted towards low wavenumber side. The intense band at  $40\text{ cm}^{-1}$  is associated with Boson peak. The increase of  $\text{Na}_2\text{O}$  content in tellurite glasses results the transformation of  $\text{TeO}_4$   $\text{tbp}$ s to  $\text{TeO}_3$   $\text{tp}$ s due to increased number of NBOs [233-234]. Assignment of main Raman bands in the spectra of tellurite glasses are listed in Table 9.

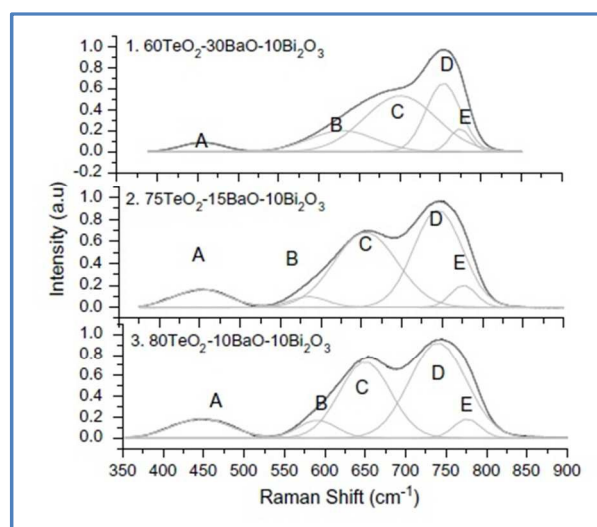


Fig. 22: Raman spectra of various tellurite glasses [231]. (Reproduced with permission of Elsevier)

Table 9: Assignment of main Raman bands in the spectra of tellurite glasses.

Wave number ( $\text{cm}^{-1}$ )	Raman assignments	Reference
40	Boson peak	233-234
60-80	$\beta\text{-TeO}_2$	239
110-143	$\gamma\text{-TeO}_2$	239
300	$\text{TeO}_3$	236
440-478	stretching and bending vibration of $\text{Te}-\text{O}-\text{Te}$ linkages in $\text{TeO}_4$ ( $\text{tbp}$ 's), $\text{TeO}_{3+\delta}$ polyhedra and $\text{TeO}_3$ ( $\text{tp}$ 's)	231, 235, 236, 243, 246
581-623	vibration of the continuous network comprised of $\text{TeO}_4$ $\text{tbp}$ 's	231
648-700	antisymmetric vibrations of $\text{Te}-\text{O}-\text{Te}$ in $\text{TeO}_{3+\delta}$ , $\text{TeO}_4$ , $\text{TeO}_3$ and $\text{TeO}_4$ networks	231, 233-234, 239, 250
716-753	stretching vibrations between tellurium and non-bridging oxygen (NBO's) atoms	231, 233-234, 239, 250
773-793	continuous network vibration of $\text{TeO}_4$ $\text{tbp}$ 's and a $\text{TeO}^-$ stretching vibration of $\text{TeO}_{3+\delta}$ polyhedra or $\text{TeO}_3$	231, 233-234, 239, 250

The study of  $x\text{Na}_2\text{O} \cdot (35-x)\text{V}_2\text{O}_5 \cdot 65\text{TeO}_2$  glass showed the increase in intensities of  $\text{TeO}_4$  tbp (443-478  $\text{cm}^{-1}$  and 671-675  $\text{cm}^{-1}$ ) and  $\text{TeO}_3$  tp (783-793  $\text{cm}^{-1}$ ) with increase of  $\text{Na}_2\text{O}$  content. This result is also consistent with above result of transformation of  $\text{TeO}_4$  tbps to  $\text{TeO}_3$  tps [235]. The network  $\text{Te-O-Te}$  bridges are broken by addition network modifiers  $\text{Na}_2\text{O}$ ,  $\text{ZnO}$  and  $\text{ZnF}_2$  in tellurite and forms the non-bridging oxygens. The intensity of vibrations  $\text{TeO}_3$  (300  $\text{cm}^{-1}$ ),  $\text{Te-O-Te}$  (470  $\text{cm}^{-1}$ )/ $\text{TeO}_4$  (675  $\text{cm}^{-1}$ ) are reduced by replacement of  $\text{ZnO}$  with  $\text{ZnF}_2$  in glassy matrix. The tellurium chain was broken by introduction of  $\text{P}_2\text{O}_5$  in tellurite glasses and suggests the formation of  $\text{Te-O-P}$  bonds [236]. Raman peaks of crystallized sodium tellurite glasses are very sharp and intense. The peak positions were shifted towards lower wavenumber side with increase of crystallization temperature [237]. The network of  $\text{TeO}_4$  bipyramids are broken to networks of  $\text{TeO}_{3+1}$  and  $\text{TeO}_3$  polyhedra at high pressures [238]. Raman spectra of zinc tellurite glasses confirms the presence  $\beta\text{-TeO}_2$  (60-80  $\text{cm}^{-1}$ ),  $\gamma\text{-TeO}_2$  (110-143  $\text{cm}^{-1}$ ). The three dimension network of asymmetrical vibration of  $\text{TeO}_4$  unit with one lone pair of electron at equilateral position is linked with  $\text{Te}_{\text{ax}}\text{-O}_{\text{eq}}\text{-Te}$  linkages. The formation of  $\text{Zn}_2\text{Te}_3\text{O}_8$  units changes the bands positioned at 295, 343 and 700-800  $\text{cm}^{-1}$  [239-241]. The peak positions of bands at 423  $\text{cm}^{-1}$  is found to shift towards lower wave numbers with increasing  $\text{ZnO}$  content whereas bands at 661 and 720  $\text{cm}^{-1}$  are shifted towards larger wavenumber side with increasing  $\text{ZnO}$  concentration. The intensities of bands at 661  $\text{cm}^{-1}$  (decrease) and 720  $\text{cm}^{-1}$  (increase) are reversed order with increase of  $\text{ZnO}$  content from 18 to 35 mol%. This indicates the decrease in  $\text{Te-O}$  coordination due to conversion of  $\text{TeO}_4$  into  $\text{TeO}_3$  structural units with increase of  $\text{ZnO}$  content [242]. Raman study of  $\text{TeO}_2\text{-La}_2\text{O}_3\text{-TiO}_2$  glasses shows the chain formation of  $\text{Te-O-Ti-O-Te-}$  at lower content of  $\text{TiO}_2$  whereas the chain formation becomes saturated and some  $\text{TiO}_4$  polyhedron are also formed in glassy matrix [243]. The structural change from  $\text{TeO}_4$  trigonal bipyramid (tbps) to  $\text{TeO}_3$  trigonal pyramid (tps) via  $[\text{TeO}_{3+1}]$  is also observed with increasing  $\text{Ta}_2\text{O}_5$  content in glass [244]. The Raman spectra (89-x)  $\text{TeO}_2\text{-10TiO}_2\text{-1Nd}_2\text{O}_3\text{-xWO}_3$  indicates a depolymerisation of tellurite glass as the  $\text{Te-O-Te}$  inter-chain linkages are progressively substituted by stronger  $\text{Te-O-W}$  bridges by the addition of  $\text{WO}_3$  [245]. The addition of  $\text{Nb}_2\text{O}_5$  in tellurite glasses decreases the connectivity by the deformation of  $\text{Te-O-Te}$  linkages [246]. The network modifying nature of  $\text{K}^+$  ions is greater than that of  $\text{Li}^+$  ions in the alkali metal tellurite glasses due to larger ionic radius of potassium than that of lithium [247]. The thallium ions in tellurite glasses induces the island-like structure of  $[\text{TeO}_3]^{2-}$  [248]. The alkali ions ( $\text{Li}$  and  $\text{Na}$ ) transports drive the secondary event as NBO migration through  $\text{BO-NBO}$  switching [249]. Khanna et al studied the Raman spectra of boro-tellurite and aluminoboro-tellurite glasses. Raman spectra of boro-tellurite glass depicts the bands at 450, 503, 615, 665, 718, 762, and 902  $\text{cm}^{-1}$  whereas alumino-boro-tellurite glasses have Raman bands located at 337, 450, 505, 567, 606, 666, 724, 754 and 864  $\text{cm}^{-1}$  [250]. The structural units  $\text{TeO}_4$  tbp's and  $\text{TeO}_{3+1}$  polyhedra changes into  $\text{TeO}_3$  tp's with an increase of  $\text{B}_2\text{O}_3$  content [251]. The diamond in tellurite glasses is the good medium for quantum information. The first order diamond Raman band is assigned at 1350  $\text{cm}^{-1}$  in nano-diamond induced tellurite glasses [252].

## 4. Conclusions

A detailed study of Raman spectroscopic data on alkali and alkaline earth metal doped various oxide glasses is reviewed. The various models were helpful for the truthful correlation of data. A general agreement was found that the structure of glassy network is compositional dependent. The vibrational properties of structural units are better understood in their structure and bonding as the vibrational spectroscopy cannot be useful as a tool for analyzing their detailed molecular structure. Raman spectroscopic studies of rectify the modification in vibration networks of borate by alkali and alkaline earth metals. The basic borate networks are formed by NBO and BO atoms in tetrahedral borate unit [9]. The formation of pentaborate, tetraborate, diborate, metaborate and pyroborate units [10], disappears of diborate units [102], boroxol ring [116] and high polymeric units such as pentaborate [105] by certain dopants was clearly noted by Raman spectroscopic studies of various borate glasses. It was found in Raman spectroscopy of oxide glasses that the modifying nature of alkali and alkaline metal are in the order of  $\text{Cs} < \text{K} < \text{Na} < \text{Li} < \text{Ca} < \text{Mg}$  [23]. The alkaline metals are effective in the order of  $\text{Ba} < \text{Sr} < \text{Ca} < \text{Mg}$  whereas the alkalis are in the order of  $\text{K} < \text{Na} < \text{Li}$ . The significant information glass forming ability of various glassy networks was obtained Raman spectroscopic studies in the present review article. On the basis of above investigations about oxide glasses, dopants  $\text{MoO}_3$ ,  $\text{CeO}_2$ ,  $\text{Bi}_2\text{O}_3$ ,  $\text{Sb}_2\text{O}_3$  and  $\text{TeO}_2$  form mainly  $\text{MoO}_6$  or  $\text{MoO}_4$ ,  $\text{CeO}_8$ ,  $\text{BiO}_3$  or  $\text{BiO}_6$ ,  $\text{SbO}_3$  and  $\text{TeO}_3$  or  $\text{TeO}_4$  units, respectively. Raman study is also providing significant UV,  $\beta$  and  $\gamma$ -ray irradiations. It was found that these radiations decrease the  $\text{Q}_{\text{Si}}^2$  species in comparison to  $\text{Q}_{\text{Si}}^3$  and it reduces the oxidation state of cations [230]. The tellurium oxide works as network former as well as network modifier. With the addition of alkali oxides, the  $\text{TeO}_4$  tbps units transform into  $\text{TeO}_3$  tps due to increased number of NBOs.

The great advances have been seen in the structural studies of glasses in past few decades still few open questions remain to study. The presence or absence of superstructural units in the glass network is the most important question. The mesounits in an intermediate- or medium-range order in the oxide structure are affecting a variety of the properties of the glass and are constituted by 7-20 or more atoms. The primary information about mesounits and their possible arrangements can be derived from the study of crystal structures isocompositional with the glasses. Such types of informations are questions for Raman spectroscopy. The quantitative information of structural units of glasses is a question for Raman spectroscopic analysis. It is always a question that informations obtained by Raman spectroscopy are either complete information or not. The vibrations, which are antisymmetrical with respect centre, are Raman inactive. In this case, it is a question that how can antisymmetric vibrations be analysed by Raman spectroscopy. Such vibrations are analysed by infrared spectroscopy. What do the analogies with organic high polymers teach us for the glass? In what respect is given glass to be called a polymer?

The absence of change of P-coordination in phosphate glasses as found in borate are unanswered by structural study of phosphate glasses by Raman spectroscopy. There is no particular relationship seen between distribution of tetrahedral linkages and glass properties. A question is also arising that are there any

differences between the non-bridging oxygens on individual  $Q_2^+$  and  $Q_1^+$  tetrahedra.

### Acknowledgements

5 Authors are thankful to the publishers Elsevier, Mineralogical Society of America, Springer, IOP science and EBSCO Publishing for giving permission for reproduction of cited contributions for academic purpose as review article. Authors acknowledge to the authors of cited articles. The funding agency 10 CSIR, New Delhi, (India) is gratefully acknowledged for financial support under project No. 03(1300)/13/EMR-II.

### References:

- 15 [1] L. Dussubieux, B. Gratuze and M. Blet-Lemarquand, *J. Archaeol. Sci.*, 2010, **37**, 1646.
- [2] A. Macfarlane and G. Martin, *Glass: A World History*. University of Chicago Press, 2002.
- [3] R. K. Brow and M. L. Schmitt, *J. Euro. Ceram.Soc.*, 2009, **29**, 1193.
- 20 [4] B. Wei, H. Cao and S. Song, *Mater. Des.*, 2010, **31**, 4244.
- [5] A. Goel, M. J. Pascual and J. M. F. Ferreira, *Int J Hydrogen Energy*, 2010, **35**, 6911.
- [6] W. H. Zachariassen, *J. Am. Chem. Soc.*, 1932, **54**, 3841.
- [7] P. B. Price, L. M. Cook and A. Marker, *Nature*, 1987, **325**, 137.
- 25 [8] T. Lopez, E. Haro-Poniatowski, P. Bosh, M. Asomoza And R. Gomez, M. Massot and M. Balkanski, *J. Sol-Gel Science and Technology*, 1994, **2**, 891.
- [9] Ch.-P. E. Varsamis, A. Vegiri and E. I. Kamitsos, *Cond. Matter Phys*, 2001, **4**, 119.
- 30 [10] K. S. Oo, M. Lwin, P. Kaung and S. Htoon, *J. Myan. Acad. Arts & Sc.*, 2006, **4**, 277.
- [11] G. Padmaja and P. Kistaiah, *J. Phys. Chem. A*, 2009, **113**, 2397.
- [12] N. C. A. de Souza, C. C. Santos, I. Guedes, N. O. Dantas and M. V. D. Vermelho, *Optical Mater.*, 2013, **35**, 2544.
- 35 [13] Q. Chen, Q. Chen, S. Wang, *Open J. Inorg. Non-met. Mater.*, 2011, **1**, 1.
- [14] S. Bale and S. Rahman, *ISRN Spectroscopy*, 2012, **2012**, 1.
- [15] Y. Wang, T. Honma and T. Komatsu, *J. Non-Cryst. Solids*, 2014, **383**, 86.
- 40 [16] P. McMillan, *Am. Mineralogist*, 1984, **69**, 622.
- [17] P. F. McMillan and R. L. Remmele, Jr., *Am. Mineralogist*, 1986, **71**, 772.
- [18] Y. Jinglin, J. Guochang, C. Hui and X. Kuangdi, *Rare Metals*, 2006, **25**, 431.
- 45 [19] R. P. Hapanowicz and R. A. Condrate, Sr., *J. Solid State Chem.*, 1996, **123**, 183.
- [20] M. Wang, J. Cheng, M. Li and F. He, *Physica B*, 2011, **406**, 3865.
- [21] A. G. Kalampounias, *Bull. Mater. Sci.*, 2011, **34**, 299.
- [22] B. O. Mysen, L. W. Finger, D. Vigro and F. A. Seifert, *Am. Mineralogist*, 1982, **67**, 686.
- 50 [23] P. McMillan, *Am. Mineralogist*, 1984, **69**, 645459.
- [24] L. D. Ferri, P. P. Lottici and G. Vezzalini, *Corrosion Sci.*, 2014, **80**, 434.
- [25] H. Jia, G. Chen and W. Wang, *Optical Mater.*, 2006, **29**, 445.
- 55 [26] H. Liu, J. Ma, J. Gong and J. Xu, *J. Non-Cryst. Solids*, 2015, **419**, 92.
- [27] M. Karabulut, H. Ertap and M. Yüsek, *J. Non-Cryst. Solids*, 2015, **417–418**, 39.
- [28] L. Ma, R. K. Brow, L. Ghussn and M. E. Schlesinger, *J. Non-Cryst. Solids*, 2015, **409**, 131.
- 60 [29] M. Farouk, *J. Non-Cryst. Solids*, 2014, **402**, 74.
- [30] L. Song, J. Wu, Z. Li, X. Hao and Y. Yu, *J. Non-Cryst. Solids*, 2015, **419**, 16.
- [31] P. Stoch and A. Stoch, *J. Non-Cryst. Solids*, 2015, **411**, 106.
- [32] G. Pablo, Debenedetti and F. H. Stillinger, *Nature*, 2001, **410**, 259.
- 65 [33] M. Yamane and Y. Asahara, *Glasses for Photonics*, Cambridge University Press, Cambridge, 2000.
- [34] A. L. Greer, *Nature*, 1999, **402**, 132.
- [35] E. Axinte, *Materials and Design*, 2011, **32**, 1717.
- [36] J. Schroers, *Nature*, 2014, **512**, 142.
- 70 [37] N. Karpukhina, R. G. Hilla and R. V. Law, *Chem. Soc. Rev.*, 2014, **43**, 2174.
- [38] G. W. Morey, *J. Industrial and Engineering Chemistry*, 1922, 823.
- [39] J. Kieffer, *J. Phys. Chem. B*, 1999, **103**, 4153.
- [40] A.K. Varshneya and J. C. Mauro, *European J. Glass Science and Technology Part A*, 2010, **51**, 28.
- 75 [41] L. Wondraczek, S. Sen, H. Behrens and R. E. Youngman, *Phys. Rev. B*, 2007, **76**, 014202.
- [42] J.C. Mauro, D.C. Allan and M. Potuzak, *Phys. Rev. B*, 2009, **80**, 094204.
- 80 [43] J. C. Mauro and M. M. Smedskjaer, *J. Non-Cryst. Solids*, 2014, **396–397**, 41.
- [44] J. E. Stanworth, *Nature*, 1952, **169**, 581.
- [45] C. R. Gautam and A. K. Yadav, *ISRN Ceramics*, 2012, **2012**, 1.
- [46] M. C. Eckersley, P. H. Gaskell, A. C. Barnes and P. Chieux, *Nature*, 1988, **335**, 525.
- 85 [47] S. Singh, M. D. Ediger and J. J. de Pablo, *Nature Materials*, 2013, **12**, 139.
- [48] Zhong-Hong Jiang and Qin-Yuan Zhang, *Prog. Mater. Sci.*, 2014, **61**, 144.
- 90 [49] M. Poulain, *Nature*, 1981, **293**, 279.
- [50] A. Fluegel, *Glass Technol: Eur J Glass Sci Technol A*, 2007, **48**, 13.
- [51] A. K. Yadav and C. R. Gautam, *Adv. Appl. Ceram.*, 2014, **113**, 193.
- [52] A. K. Yadav and C. R. Gautam, *Spectroscopy Lett.*, 2015, **48**, 514.
- [53] A. K. Yadav and C. R. Gautam, *Adv. Appl. Ceram.*, 2014, **113**, 193.
- 95 [54] M. Yamane and Y. Asahara, *Glasses for photonics*, Cambridge University, Cambridge, UK, 2000.
- [55] J. C. Phillips and M. F. Thorpe, *Solid State Commun.* 1985, **53**, 699.
- [56] P. K. Gupta and J. C. Mauro, *The Journal of Chemical Physics*, 2009, **130**, 094503.
- 100 [57] P. G. Pavani, K. Sadhana and V. C. Mouli, *Physica B*, 2011, **406**, 1242.
- [58] K. Cho, J. Oh, T. Lee and D. Shin, *J. Power Sources*, 2008, **183**, 431.
- [59] S. Baccaro, N. Catallo, A. Cemmi and G. Sharma, *Nuclear Instr. & Meth. Phys. Res. Section B-Beam Interactions with Mater. and Atoms*, 2011, **269**, 167.
- 105 [60] Y. J. Chen, X. H. Gong, Y. F. Lin, Z. D. Luo and Y. D. Huang, *Optical Materials*, 2010, **33**, 71.
- [61] S. P. Singh, R. P. S. Chakradhara, J. L. Rao and B. Karmakar, *J. Alloys and Comp.*, 2010, **493**, 256.
- 110 [62] R. Gehring, *Scientific American Inc*, 1890, **737**, 11776.
- [63] R. Dupree, D. Holland and M.G. Mortuza, *Nature*, 1987, **328**, 416.
- [64] L. Cormier, G. Calas and B. Beuneu, *J. Non-Cryst. Solids*, 2011, **357**, 926.
- [65] A. M. Efimov, *J. Non-Cryst. Solids*, 1999, **253**, 95.
- 115 [66] E. M. Rabinovich, *J. Mater. Sci.*, 1976, **11**, 925.
- [67] M. Karabulut, E. Melnik, R. Stefan, G. K. Marasinghe, C. S. Ray, C. R. Kurkjin and D. E. Day, *J. Non-Cryst. Solids*, 2001, **288**, 8.
- [68] K. El-Egili, H. Doweidar, Y. M. Moustafa and I. Abbas, *Physica B*, 2003, **339**, 237.
- 120 [69] G. Guimbretière, D. Bégué, M. Dussauze and V. Rodriguez, *Vibrational Spect.*, 2012, **63**, 426.
- [70] P.B. Price, L.M. Cook and A. Marker, *Nature*, 1987, **325**, 137.
- [71] A. A. Osipov, L. M. Osipova and V. E. Eremyashev, *Glass Phys. Chem.*, 2013, **39**, 105.
- 125 [72] O. N. Koroleva, L. A. Shabunina and V. N. Bykov, *Glass and Ceramics*, 2011, **67**, 340.
- [73] A. Goel, J. S. McCloy, C. F. W. Jr, B. J. Riley, M. J. Schweiger and C. P. Rodriguez and J. M. F. Ferreira, *Int. J. Appl. Glass Sci.*, 2012, **1**.
- 130 [74] J. Blizard, *Industrial and Engineering Chemistry*, 1947, **39**, 1215.
- [75] L. Koudelka, P. Mosner and J. Šubčík, *IOP Conf. Ser.: Mater. Sci. Eng.* 2009, **2**, 012015.
- [76] D. Kim, C. Hwang, D. Gwoo, T. Kim, N. Kim and B.-K. Ryu, *Electr. Mater. Lett.*, 2011, **7**, 343.
- 135 [77] A. Naumaan and J. T. Boyd, *J. Electrochem. Soc.*, 1980, **127**, 1414.
- [78] A. Tilocca, *Phys. Rev. B*, 2007, **76**, 224202.
- [79] S. W. Lee, K. S. Ryoo, J. E. Kim, J. H. Lee, C. D. Kim and K. S. Hong, *J. Mater. Sci.*, 1994, **29**, 4577.

- [80] J. L. Rygel, Y. Chen, C. G. Pantano, T. Shibata, J. Du, L. Kokou, R. Woodman and J. Belcher, *J. Am. Ceram. Soc.*, 2011, **94**, 2442.
- [81] R. W. Cahn, *Nature*, 1989, **341**, 183.
- [82] M. Handke, M. Sitarz, M. Rokita and E. Galuskin, *J. Molecular Struct.*, 2003, **651**, 39.
- [83] L. Cormier and D. R. Neuville, *Chem. Geo.*, 2004, **213**, 103.
- [84] S. L. S. Rao, G. Ramadevudu, Md. Shareefuddin, A. Hameed, M. N. Chary and M. L. Rao, *Int. J. Eng., Sci. Tech.*, 2012, **4**, 25.
- [85] D. P. Singh and G. P. Singh, *J. Alloys Comp.*, 2013, **546**, 224.
- [86] G. Calas, L. Cormier, L. Galoisy and P. Jollivet, *C. R. Chimie*, 2002, **5**, 831.
- [87] J. Qiu, X. Jiang, C. Zhu, H. Inouye, J. Si and K. Hirao, *Opt Lett.*, 2004, **29**, 370.
- [88] Q. Yu, F. Chen, T. Xu, S. Dai and Q. Zhang, *J. Non-Cryst. Solids*, 2013, **378**, 254.
- [89] M. Diem (Ed.), *Introduction to Modern Vibrational Spectroscopy*, New York: Wiley-Interscience, 1993.
- [90] C. V. Raman and K. S. Krishnan, *Nature*, 1928, **121**, 501.
- [91] C. V. Raman, *Indian J. Phys.*, 1928, **2**, 387-398.
- [92] C. C. Lin, M. T. Kuo and H. C. Chang, *J. Med. Biolog. Engg.*, 2010, **30**, 343-354.
- [93] C. V. Raman and K. S. Krishnan, *Indian J. Phys.*, 1928, **2**, 399.
- [94] C. V. Raman, K. S. Krishnan, *Nature (London)*, 1928, **121**, 619.
- [95] L. A. Lyon, C. D. Keating, A. P. Fox, B. E. Baker, H. Lin, S. R. Nicewarner, S. P. Mulvaney and M. J. Natan, *Analytical Chemistry*, 1998, **70**, 341R.
- [96] H. H. Willard, Jr. L. L. Merritt and J. A. Dean, *Raman Spectroscopy*. In *Instrumental Methods of Analysis*, 5<sup>th</sup> ed, D. Van Nostrand Co., New York, 1974, 189.
- [97] B. H. Stuart, *Infrared Spectroscopy, Fundamentals and Applications*. J. Wiley, Chichester, West Essex, England, 2004, 2.
- [98] M. S. Shackley, *X-Ray Fluorescence Spectrometry (XRF) in Geoarchaeology*. Springer, New York, 2010, 16
- [99] T. Lopez, E. H. Poniatowski, P. Bosh, M. Asomoza, R. Gomez, M. Massot and M. Balkanski, *J. Sol-Gel Sci. Tech.*, 1994, **2**, 891.
- [100] M. Loubser, *Chemical and Physical Aspects of Lithium Borate Fusion*, Distertation M. Sc. University of Pretoria, Pretoria, 2009.
- [101] W. Di, W. Song-ming, Z. Qing-li, S. Dun-lu, Y. Shao-tang, Y. Jing-lin and W. Yuan-yuan, *Chin. J. Quan. Elect.*, 2011, **28**, 241.
- [102] M. Ganguli and K. J. Rao, *J. Solid State Chem.*, 1999, **145**, 65.
- [103] G. Padmaja and P. Kistaiah, *J. Phys. Chem. A*, 2009, **113**, 2397.
- [104] T. Yano, N. Kunimine, S. Shibata and M. Yamane, *J. Non-Cryst. Solids*, 2003, **321**, 137.
- [105] Q. Jiao, X. Yu, X. Xu, D. Zhou and J. Qiu, *J. Solid State Chem.*, 2013, **202**, 65.
- [106] A. A. A. Shama and F. H. El-Batal, *Egypt. J. Solids*, 2006, **29**, 49.
- [107] M. Affatigato, S. Feller, A. K. Schue, S. Blair, D. Stentz, G. B. Smith, D. Liss, M. J. Kelley, C. Goater and R. Leelasagar, *J. Phys.: Condens. Matter*, 2003, **15**, S2323.
- [108] Q. Chen, Q. Chen, S. Wang, *Open J. Inorg. Non-met. Mater.*, 2011, **1**, 1.
- [109] V. Timar, R. Lucăcel-Ciceo and I. Ardelean, *Quan. Electr. Optoelectronics*, 2008, **11**, 221.
- [110] J. Schwarz and H. Tichá, *J. Optoelectr. Adv. Mater.*, 2003, **5**, 69.
- [111] A. A. Alemi, H. Sedghi, A. R. Mirmohseni and V. Golsanamlu, *Bull. Mater. Sci.*, 2006, **29**, 55.
- [112] A. Toderas, M. Toderas, S. Filip and I. Ardelean, *Optoelectr. Adv. Mater. Rapid Comm.*, 2012, **6**, 331.
- [113] B. S. Cristian, *Study of some structural and physical properties of some B<sub>2</sub>O<sub>3</sub> based glasses containing Ag<sub>2</sub>O and 3d ions*, Ph.D. thesis, Babeş-Bolyai University, 2011.
- [114] S. C. Baidoc and I. Ardelean, *J. Optoelectr. Adv. Mater.*, 2008, **10**, 3205.
- [115] R. Uning, W. M. Hua and R. Hussin, *Structural studies of strontium borate glass doped with Zn<sup>2+</sup>, Mn<sup>2+</sup> and Fe<sup>2+</sup> ions*, UMTAS 2011.
- [116] S.R. Rejisha, P.S. Anjana, N. Gopakumar and N. Santha, *J. Non-Cryst. Solids*, 2014, **388**, 68.
- [117] M. Bettinelli, A. Speghini, M. Ferrari and M. Montagna, *J. Non-Cryst. Solids*, 1996, **201**, 211.
- [118] S. Suresh, M. Prasad, G. Upender, V. Kamalaker and V. C. Mouli, *Ind. J. Pure and Appl. Phys.*, 2009, **47**, 163.
- [119] D. Maniu, I. Ardelean, T. Iliescu, S. Cinta, V. Nagel and W. Kiefer, *J. Molecular Struct.*, 1999, **480**, 657.
- [120] D. Maniu, T. Iliescu, I. Ardelean, S. Cinta-Pinzaru, N. Tarcea and W. Kiefer, *J. Molecular Struct.*, 2003, **651**, 485.
- [121] C. N. Santos, D. D. S. Meneses, P. Echegut, D. R. Neuville, A. C. Hernandez and A. Ibanez, *Appl. Phys. Lett.*, 2009, **94**, 151901.
- [122] R. S. Barker, D. A. Richardson, E. A. G. Mcconkey and R. Rimmer, *Nature*, 1960, **187**, 135.
- [123] B. Tiwari, A. Dixit, G. P. Kothiyal, M. Pandey and S. K. Deb, *Preparation and characterization of phosphate glasses containing titanium*, National Symposium on Science & Technology of Glass and Glass-Ceramics (NSGC-06) held during September 15-16, at BARC, Mumbai, 2006,
- [124] L. F. Santos, R. M. Aldeida, V. K. Tikhomirov and A. Jha, *J. Non-Cryst. Solids*, 2001, **284**, 43.
- [125] A. Mogus-Milankovic, A. Gajovic, A. Santic and D. E. Day, *J. Non-Cryst. Solids*, 2001, **289**, 204.
- [126] K. Brahmachary, D. Rajesh, S. Babu and Y. C. Ratnakaram, *J. Molecular Struct.* 2014, **1064**, 6.
- [127] A. Kiani, J. V. Hanna, S. P. King, G. J. Rees, M. E. Smith, N. Roohpour, V. Salih and J. C. Knowles, *Acta Biomaterialia*, 2012, **8**, 333.
- [128] M. Elisa, I. C. Vasiliu, C. E. A. Grigorescu, B. Grigoras, H. Niciu, D. Niciu, A. Meghea, N. Ifimie, M. Giurginca, H. J. Trodahl and M. Dalley, *Optical Mater.*, 2006, **28**, 621.
- [129] F. H. El-Batal, *Ind. J. Pure and Appl. Phys.*, 2009, **47**, 631.
- [130] D. Toloman, R. Ciceo-Lucacel, D. A. Magdas, A. Regos, A. R. Biris, C. Leostean, I. Ardelean, *J. Alloys Comp.*, 2013, **556**, 67.
- [131] B. Qian, X. Liang, C. Wang and S. Yang, *J. Nuclear Mater.*, 2013, **443**, 140.
- [132] C. I. R. D. Oliveira, L. F. C. D. Oliveira, F. A. D. Filho, Y. Messaddeq and S. J. L. Ribeiro, *Spectrochimica Acta Part A*, 2005, **61**, 2023.
- [133] D. Ilieva, B. Jivov, G. Bogachev, C. Petkov, I. Penkov and Y. Dimitriev, *J. Non-Cryst. Solids*, 2001, **283**, 195.
- [134] D. Ilieva, B. Jivov, D. Kovacheva, Ts. Tsacheva, Y. Dimitriev, G. Bogachev and C. Petkov, *J. Non-Cryst. Solids*, 2001, **293-295**, 562.
- [135] L. Koudelka, I. Rösslerová, J. Holubová, P. Mošner, L. Montagne and B. Revel, *J. Non-Cryst. Solids*, 2011, **357**, 2816.
- [136] G. L. Flower, M. S. Reddy, G. S. Baskaran and N. Veeraiiah, *Optical Mater.*, 2007, **30**, 357.
- [137] S. S. Sastry and B. R. V. Rao, *Physica B*, 2014, **434**, 159.
- [138] G. Le Saout, F. Fayon, C. Bessada, P. Simon, Annie Blin and Y. Vaills, *J. Non-Cryst. Solids*, 2001, **293-295**, 657.
- [139] W. A. Pisarski, L. Zur, T. Goryczka, M. Sołtys and J. Pisarska, *J. Alloys Comp.*, 2014, **587**, 90.
- [140] A. M. Milankovic, A. Santic, S. T. Reis, K. Furic and D. E. Day, *J. Non-Cryst. Solids*, 2005, **351**, 3246.
- [141] G. S. Baskaran, G. L. Flower, D. K. Rao and N. Veeraiiah, *J. Alloys Comp.*, 2007, **431**, 303.
- [142] A. Šantić, Ž. Skoko, A. Gajović, S. T. Reis, D. E. Day and A. M. Milanković, *J. Non-Cryst. Solids*, 2011, **357**, 3578.
- [143] A. Santic, A. M. Milankovic, K. Furic, V. Bermanec, C. W. Kim and D. E. Day, *J. Non-Cryst. Solids*, 2007, **353**, 1070.
- [144] A. M. Milankovic, A. Santic, A. Gajovic and D. E. Day, *J. Non-Cryst. Solids*, 2003, **325**, 76.
- [145] B. Qian, X. Liang, S. Yang, S. He and L. Gao, *J. Molecular Struct.*, (2012)**1027**, 31.
- [146] K. Joseph, M. Premila, G. Amarendra, K. V. G. Kutty, C. S. Sundar and P. R. V. Rao, *J. Nuclear Mater.*, 2012, **420**, 49.
- [147] Y. M. Lai, X. F. Liang, S. Y. Yang, J. X. Wang, L. H. Cao and B. Dai, *J. Molecular Struct.*, 2011, **992**, 84.
- [148] Y. M. Lai, X. F. Liang, S. Y. Yang, J. X. Wang and B. T. Zhang, *J. Molecular Struct.*, 2012, **1013**, 134.
- [149] S. Chakraborty and A. K. Arora, *Vibrational Spectr.*, 2012, **61**, 99.
- [150] P. S. Rao, P. M. V. Teja, A. R. Babu, C. Rajyasree and D. K. Rao, *J. Non-Cryst. Solids*, 2012, **358**, 3372.
- [151] N. Kitamura, K. Fukumi, J. Nakamura, T. Hidaka, T. Ikeda, H. Hashima and J. Nishii, *J. Non-Cryst. Solids*, 2011, **357**, 1188.

- [152] P. S. Rao, S. B. M. Krishna, S. Yusub, P. R. Babu, C. Tirupataiah and D. K. Rao, *J. Molecular Struct.*, 2013, **1036**, 452.
- [153] S. Thomas, R. George, M. Rathaiah, V. Venkatramu, S. N. Rasool and N. V. Unnikrishnan, *Physica B*, 2013, **431**, 69.
- [154] S. H. Santagneli, G. Poirier, M. T. Rinke, S. J. L. Ribeiro, Y. Messaddeq and H. Eckert, *J. Non-Cryst. Solids*, 2011, **357**, 2126.
- [155] R. C. Lucacel, A. O. Hulpus, V. Simonn and I. Ardelean, *J. Non-Cryst. Solids*, 2009, **355**, 425.
- [156] A. M. B. Silva, R. N. Correia, J. M. M. Oliveira and M. H. V. Fernandes, *J. Euro. Ceram. Soc.*, 2010, **30**, 1253.
- [157] I. D. F. Gimenez, I. O. Mazali and O. L. Alves, *J. Phys. Chem. Solids*, 2001, **62**, 1251.
- [158] S. N. Rasool, L. R. Moorthy and C. K. Jayasankar, *Optics Comm.*, 2013, **311** 156.
- [159] M. Elisa, B. A. Sava, I. C. Vasiliu, R. C. C. Monteiro, J. P. Veiga, L. Ghervase, I. Feraru and R. Iordanescu, *J. Non-Cryst. Solids*, 2013, **369**, 55.
- [160] N. Vedeau, O. Cozar, I. Ardelean, B. Lendl and D. A. Magdas, *Vibrational Spectr.*, 2008, **48**, 259.
- [161] A. Chahine, M. Et-tabirou and J. L. Pascal, *Mater. Lett.*, 2004, **58**, 2776.
- [162] E. Metwalli, M. Karabulut, D. L. Sidebottom, M. M. Morsi and R. K. Brow, *J. Non-Cryst. Solids*, 2004, **344**, 128.
- [163] N. Kerkouri, M. Haddad, M. Et-Tabirou, A. Chahine and L. Laanab, *Physica B*, 2011, **406**, 3142.
- [164] C. Ivascu, I. B. Cozar, L. Daraban and G. Damian, *J. Non-Cryst. Solids*, 2013, **359**, 60.
- [165] C. Ivascu, A. Timar Gabor, O. Cozar, L. Daraban and I. Ardelean, *J. Molecular Struct.*, 2011, **993**, 249.
- [166] S. Hraiech, M. Ferid, Y. Guyot and G. Boulon, *J. Rare Earths*, 2013, **31**, 685.
- [167] B. G. Parkinson, D. Holland, M. E. Smith, C. Larson, J. Doerr, M. Affatigato, S. A. Feller, A. P. Howes and C. R. Scales, *J. Non-Cryst. Solids*, 2008, **354**, 1936.
- [168] D. Manara, A. Grandjean and D. R. Neuville, *Am. Mineralogist*, 2009, **94**, 777.
- [169] M. Lenoir, A. Grandjean, S. Poissonnet and D. R. Neuville, *J. Non-Cryst. Solids*, 2009, **355**, 1468.
- [170] M. Lenoir, A. Grandjean, Y. Linard, B. Cochain and D. R. Neuville, *Chem. Geo.*, 2008, **256**, 316.
- [171] M. Jroda and C. Paluszkiwicz, *J. Molecular Struct.*, 2007, **834–836**, 302.
- [172] V. E. Eremyashev, A. A. Osipov and L. M. Osipov, *Glass and Ceramics*, 2011, **68**, 205.
- [173] A. Winterstein-Beckmann, D. Möncke, D. Palles, E. I. Kamitsos and L. Wondraczek, *J. Non-Cryst. Solids*, 2014, **401**, 110.
- [174] H. Yamashita, K. Nagata, H. Yoshino, K. Ono and T. Maekawa, *J. Non-Cryst. Solids*, 1999, **248**, 115.
- [175] D. Chen, H. Miyoshi, H. Masui, T. Akai and T. Yazawa, *J. Non-Cryst. Solids*, 2004, **345–346**, 104.
- [176] D. A. McKeown, I. S. Muller, Hao Gan, I. L. Pegg and C. A. Kendziora, *J. Non-Cryst. Solids*, 2001, **288**, 191.
- [177] D. A. McKeown, I. S. Muller, H. Gan, I. L. Pegg and W. C. Stolte, *J. Non-Cryst. Solids*, 2004, **333**, 74.
- [178] X. Zhang, Y. Wang, J. Lu, J. Zang, J. Zhang, E. Ge, *Int. Journal of Refractory Metals & Hard Mater.*, 2010, **28**, 260.
- [179] D. A. McKeown, I. S. Muller, A. C. Buechele, I. L. Pegg and C. A. Kendziora, *J. Non-Cryst. Solids*, 2000, **262**, 126.
- [180] T. G. V. M. Rao, A. R. Kumar, K. Neeraja, N. Veeraiah and M. R. Reddy, *J. Alloys Comp.*, 2013, **557**, 209.
- [181] Z. Wang, Q. Shu and K. Chou, *ISIJ International*, 2011, **51**, 1021.
- [182] L. Rezazadeh, S. Baghshahi, A. Nozad Golikand and Z. Hamnabard, *Ionics*, 2013, 1.
- [183] A. K. Yadav, C. R. Gautam, A. Gautam, V. K. Mishra, *Phase Transitions*, 2013, **86**, 1000.
- [184] A. K. Yadav and C. R. Gautam, *Optics & Photonics Journal*, 2013, **3**, 1.
- [185] C. R. Gautam, A. K. Yadav, V. K. Mishra and K. Vikram, *Open J. Non-met. Mater.*, 2012, **2**, 47
- [186] G. Kaur, O. P. Pandey, K. Singh, *J. Non-Cryst. Solids*, 2012, **358**, 2589.
- [177] C. R. Gautam, A. K. Singh and A. K. Yadav, *Int. J. Appl. Natural Sci.*, 2012, **1**, 69.
- [188] G. F. Zhang, T. S. Wang, K. J. Yang, L. Chen, L. M. Zhang, H. B. Peng, W. Yuan and F. Tian, *Nuclear Instr. Meth. Phys. Res. B*, 2013, **316**, 218.
- [189] C. Gejke, E. Zanghellini, J. Swenson and L. Borjesson, *J. Power Sources*, 2003, **119–121**, 576.
- [190] P. Mošner, M. Vorokhta, L. Koudelka, L. Montagne, B. Revel, K. Sklepić and A. M. Milanković, *J. Non-Cryst. Solids*, 2013, **375**, 1.
- [191] N. Sdiri, H. Elhouichet and M. Ferid, *J. Non-Cryst. Solids*, 2014, **389**, 38.
- [192] K. J. Rao, K. C. Sobha and S. Kumar, *Proc. Indian Acad. Sci. (Chem. Sci.)*, 2001, **113**, 497.
- [193] A. Saranti, I. Koutselas and M. A. Karakassides, *J. Non-Cryst. Solids*, 2006, **352**, 390.
- [194] L. Koudelka, J. Pospisil, P. Mosner, L. Montagne and L. Delevoeye, *J. Non-Cryst. Solids*, 2008, **354**, 129.
- [195] L. Koudelka, P. Mosner, M. Zeyer and C. Jager, *J. Non-Cryst. Solids*, 2003, **326–327**, 72.
- [196] L. Koudelka, I. Rösslerová, Z. Černošek, P. Mošner, L. Montagne and B. Revel, *J. Non-Cryst. Solids*, 2014, **401**, 124.
- [197] L. Koudelka and P. Mosner, *Mater. Lett.*, 2000, **42**, 194.
- [198] L. Koudelka, P. Mosner, J. Pospisila, L. Montagne and G. Palavit, *J. Solid State Chem.*, 2005, **178**, 1837.
- [199] L. Koudelka, J. Šubcik, P. Mosner, L. Montagne and L. Delevoeye, *J. Non-Cryst. Solids*, 2007, **353**, 1828.
- [200] J. Šubcik, L. Koudelka, P. Mošner, L. Montagne, B. Revel and I. Gregora, *J. Non-Cryst. Solids*, 2009, **355**, 970.
- [201] P. Mosner, K. Vosejpková, L. Koudelka, L. Montagne and B. Revel, *Mater. Chem. Phys.* 2010, **124**, 732.
- [202] K. Vosejpkova, L. Koudelka, Z. Cernosek, P. Mosner, L. Montagne and B. Revel, *J. Phys. Chem. of Solids*, 2012, **73**, 324.
- [203] S. Petrescu, M. Constantinescu, E. M. Anghel, I. Atkinson, M. Olteanu and M. Zaharescu, *J. Non-Cryst. Solids*, 2012, **358**, 3280.
- [204] C. L. Losq, D. R. Neuville, P. Florian, G. S. Henderson and D. Massiot, *Geochimica et Cosmochimica Acta*, 2014, **126**, 495.
- [205] D. A. McKeown, F. L. Galeener and G. E. Brown Jr, *J. Non-Cryst. Solids*, 1984, **68**, 361.
- [206] M. Okuno, N. Zotov, M. Schmucker, H. Schneider, *J. Non-Cryst. Solids*, 2005, **351**, 1032.
- [207] P. McMillan, B. Piriou and A. Navrotsky, *Geochimica et Cosmochimica Acta*, 1982, **46**, 2021.
- [208] P. McMillan and B. Piriou, *J. Non-Cryst. Solids*, 1983, **55**, 221.
- [209] D. R. Neuville, L. Cormier and D. Massiot, *Geochimica et Cosmochimica Acta*, 2004, **68**, 5071.
- [210] D. R. Neuville, L. Cormier and D. Massiot, *Chem. Geo.*, 2006, **229**, 173.
- [211] D. R. Neuville, L. Cormier, A. M. Flank, V. Briois and D. Massiot, *Chem. Geo.* 2004, **213**, 153.
- [212] M. Licheron, V. Montouillout, F. Millot and D. R. Neuville, *J. Non-Cryst. Solids*, 2011, **357**, 2796.
- [213] C. Huang and E. C. Behrman, *J. Non-Cryst. Solids*, 1991, **128**, 310.
- [214] J. F. Stebbins, E. V. Dubinsky, K. Kanehashi and K. E. Kelsey, *Geochimica et Cosmochimica Acta*, 2008, **72**, 910.
- [215] V. G. Plotnichenko, V. O. Sokolov, V. V. Koltashev and E. M. Dianov, *J. Non-Cryst. Solids*, 2002, **306**, 209.
- [216] A. Trukhin and B. Capoen, *J. Non-Cryst. Solids*, 2005, **351**, 3640.
- [217] S. Agathopoulos, D. U. Tulyaganov, J. M. G. Ventura, S. Kannan, A. Saranti, M. A. Karakassides and J. M. F. Ferreira, *J. Non-Cryst. Solids*, 2006, **352**, 322.
- [218] H. Aguiar, E. L. Solla, J. Serra, P. González, B. León, F. Malz and C. Jäger, *J. Non-Cryst. Solids*, 2008, **354**, 5004.
- [219] H. Aguiar, J. Serra, P. González and B. León, *J. Non-Cryst. Solids*, 2009, **355**, 475.
- [220] A. Vulpoi, L. Baia, S. Simon, V. Simon, *Mater. Sci. Eng. C*, 2012, **32**, 178.
- [221] Y. Gandhi, M. V. R.C. Rao, C. S. Rao, I. V. Kityk and N. Veeraiah, *J. Luminescence*, 2011, **131**, 1443.
- [222] K. Cheng, J. Wan and K. Liang, *Mater. Lett.*, 2001, **47**, 1.
- [223] H. Li, Y. Su, L. Li and D. M. Strachan, *J. Non-Cryst. Solids*, 2001, **292**, 167.



- [224] A. R. Kumar, T. G. V. M. Rao, K. Neeraja, M. R. Reddy and N. Veeraiah, *Vibrational Spectr.*, 2013, **69**, 49.
- [225] E. Malchukova, B. Boizot, G. Petite and D. Ghaleb, *J. Non-Cryst. Solids*, 2008, **354**, 3592.
- 5 [226] E. Malchukova and B. Boizot, *Phys. Solid State*, 2008, **50**, 1687.
- [227] E. Malchukova, B. Boizot, G. Petite and D. Ghaleb, *The Eur. Phys. J. Appl. Phys.*, 2009, **45**, 10701.
- [228] E. Malchukova, B. Boizot, D. Ghaleb and G. Petite, *J. Non-Cryst. Solids*, 2006, **352**, 297.
- 10 [229] E. Malchukova, B. Boizot, G. Petite and D. Ghaleb, *J. Non-Cryst. Solids*, 2007, **353**, 2397.
- [230] E. Malchukovaa, B. Boizot, D. Ghaleb, G. Petite, *Nuclear Instr. Meth. Phys. Res. A*, 2005, **537**, 411
- [231] A. Jha, B. Richards, G. Jose, T. T. Fernandez, P. Joshi, X. Jiang and  
15 J. Lousteau, *Progress mater. sci.*, 2012, **57**, 1426.
- [232] F. Pietrucci, S. Caravati and M. Bernasconi, *Phys. Rev. B*, 2008, **78**, 064203.
- [233] A. Santic, A. M. Milankovic, K. Furic, M. R. Linaric, C. S. Ray and D. E. Dayd, *Croatica Chemica Acta*, 2008, **81**, 559.
- 20 [234] J. Heo, D. Lam, G. H. Sigel, Jr. E. A. Mendoza and D. A. Hensley, *J. Am. Ceram. Soc.*, 1992, **75**, 277.
- [235] M.M. Umair and A. K. Yahya, *Mater. Chem. Phys.*, 2013, **142**, 549.
- [236] M. Irannejad, G. Jose, A. Jha and P. Steenson, *Optics Communications*, 2012, **285**, 2646.
- 25 [237] A. Jha, P. Joshi and S. Shen, *Optics Express*, 2008, **16**, 13526.
- [238] E. Stavrou, C. Raptis and K. Syassen, *Phys. Rev. B*, 2010, **81**, 174202.
- [239] D. Dutta, M. P. F. Graca, M. A. Valente and S.K. mendiratta, *Solid state ionics*, 2013, **230**, 66.
- 30 [240] M. Ceriotti, F. Pietrucci and M. Bernasconi, *Phys. Rev. B*, 2006, **73**, 104304-1 - 104304,
- [241] S. Manning, H. Ebendorff-Heidepriem and Tanya M. Monro, *Optical Mat. Express*, 2012, **2**, 140.
- [242] A. Kaur, A. Khanna, C. Pesquera, F. González and V. Sathe, *J. Non-Cryst. Solids*, 2010, **356**, 864.
- 35 [243] W. Stambouli, H. Elhouichet and M. Ferid, *J. Molecular Structure*, 2012, **1028**, 39.
- [244] I. Yakine, A. Chagraoui, A. Moussaoui and A. Tairi, *J. Mater. Environ. Sci.*, 2012, **3**, 776.
- 40 [245] Hssen Fares, Ifa Jlassi, Habib Elhouichet and Mokhtar Férid, *J. Non-Cryst. Solids*, 2014, **396–397**, 1.
- [246] A. Kaur, A. Khanna, V. G. Sathe, F. Gonzalez and B. Ortiz, *Phase Transitions*, 2013, **86**, 598.
- [247] R. Akagi, K. Handa, N. Ohtori, A.C. Hannon, M. Tatsumisago and N. Umesaki, *J. Non-Cryst. Solids*, 1999, **256 & 257**, 111.
- 45 [248] T. Hayakawa, M. Koduka, M. Nogami, J. R. Duclere, A. P. Mirgorodsky and P. Thomas, *Scripta Materialia*, 2010, **62**, 806.
- [249] M. H. Bhat, M. Ganguli and K. J. Rao, *Current Science*, 2004, **86**, 676.
- 50 [250] N.Kaur and A. Khanna, *J. Non-Cryst. solids*, 2014, **404**, 116.
- [251] P. Damas, J. Coelho, G. Hungerford and N. S. Hussain, *Mater. Res. Bull.*, 2012, **47**, 3489.
- [252] M. R. Henderson, B. C. Gibson, H. Ebendorff-Heidepriem, K. Kuan, S. Afshar V., J. O. Orwa, I. Aharonovich , S. Tomljenovic-Hanic, A. D. Greentree, S. Praver and T. M. Monro, *Adv. Mater.*, 2011, **23**, 2806.
- 55 [253] M. M. Smedskjaer, J. C. Mauro, S. Sen and Y. Yue, *Chem. Mater.* 2010, **22**, 5358.

60