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Optimization of the nanotwin-induced zigzag surface of copper by electromigration

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By adding nanotwins to Cu, the surface electromigration (EM) slows down. The atomic mobility of the surface step-edges is retarded by the triple points where a twin meets a free surface to form a zigzag-type surface. We observed that EM can alter the zigzag surface structure to optimize the reduction of EM, according to Le Chatelier's principle. Statistically, the optimal alternation is to change an arbitrary (111)/(hkl) zigzag pair to a pair having a very low index (hkl) plane, especially the (200) plane. Using in-situ ultrahigh vacuum and high-resolution transmission electron microscopy, we examined the effects of different zigzag surfaces on the rate of EM. The calculated rate of surface EM can be decreased by a factor of ten.

1 Introduction

30 In Al interconnect technology, electromigration (EM) occurs by 2 grain boundary diffusion at the device operation temperature? 3 4 To slow down EM in Al, the alloying of approximately 1 at. $\frac{3}{3}$ 5 Cu solute has been found to be effective^[1]. In Cu interconnera technology, EM occurs by surface diffusion, and solutes, such 6 7 as Sn, Si, P, Al, and Mn, have been proposed to reduce the EMG 8 which follows the same alloying approach taken in Al 9 technology^[2-4]. The main problem with the use of alloying is37 10 performance penalty owing to the increase in Cu resistivity 11 which is not desirable, as the device dimensions continue **39** 12 shrink. We report here that instead of solute alloying, a high 13 density of nano-twins can serve the same purpose without sacrificing the electrical conductivity of pure Cu^[5-10]. Mode 14 15 importantly, we found that EM can alter the nanotwin-inducad 16 zigzag surface structure. The normal evolution of the zigzag 17 surface suggests a statistically optimal orientation of nanotwi45 18 in a polycrystalline Cu interconnect to further slow EM. What 19 introducing nanotwins into Cu, the free surface where a twin 20 intersects is transformed to a zigzag surface because a vallage and a ridge are formed by the two triple points at the 21 22 intersections of the pair of twin planes with the free surface 23 Typically, the zigzag surface consists of two sets of alternating 24 surface planes, for example, (111)/(422) /(111) /(422)... It has 25 been reported that the triple points can slowdown the net rate of surface diffusion when the EM-induced surface step migrates 26 27 from the (111) surface plane to the (422) surface plane on the 53

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zigzag surface^[11]. This is because of the necessity of the change of the surface step on the (111) surface to that on the (422) surface and the long incubation time that is needed to nucleate a new step. However, in addition to the delay of step migration across triple points, local surface structures are also believed to significantly affect the EM behaviors of surface steps^[12-16]. In turn, how EM can alter the local surface structure and how feedback based on Le Chatelier's principle can reduce EM further are addressed in this paper.

On different orientations of Cu grains, the zigzag surface is different, and a different reduction rate of EM is found. Generally speaking, we found that on a zigzag surface, while one set of the alternating surface plane tends to be (111), the other set can be a low-index (hkl) plane. The selection is based on the surface energy from thermodynamics, yet under EM, the selection based on kinetics is more favorable. For comparison, we have also measured the rate of EM on a (111) surface that has no nano-twins; the rate is orders of magnitude faster than that on a zigzag surface having nano-twins. In this report, the effect of the configuration of the surface structures of nanotwin-modified Cu on EM has been observed by an *in-situ* TEM-EM technique. We correlate the magnitude of the EM-induced mass flux to the specific geometry of the surface step-edges on the zigzag surfaces.

Results and discussion

We performed a direct comparison between the EM-induced migration of step-edges in a $(0\overline{1}1)$ -oriented Cu grain with and without nanotwins. A detailed description of the experimental procedure is provided in the supplementary materials. The two regions of examination are indicated in the schematic diagram of a Cu grain in Figure 1a. A void with a new free surface is formed in this grain when atomic layers are gradually removed from the right side of the Cu grain under EM. The new free surface has two regions of different surface planes with step edges, as seen in Figure 1b and 1c, depending on whether there are nanotwins. Figure 1b is the region of a free surface of the



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stressing time (Figs. 3a-c).

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(111) plane without nanotwins, and there are some surfa**48** step-edges on the (111) plane that form a curved surface.

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 $\begin{array}{l} \mbox{5} & \mbox{Figure 1. EM-induced Cu surface planes. (a) A schematic diagram of a (011)- \\ \mbox{oriented Cu grain without and with nanotwin-structured regions. (b) and (c) \\ \mbox{Corresponding HRTEM images of the evolved step edges at unnano-twinned and \\ \mbox{nano-twinned surfaces under EM.} \end{array}$

9 Figure 1c is the region of a free surface with nanotwins that 10 form a zigzag shape. The zigzagged surface consists of 11 alternating (111) and (200) planes, and surface step-edges are 12 inherent in the (200) plane to maintain the zigzag surface. The 13 evolution of the surface structure, morphology, and EM rate on 14 these two regions were examined. In Figure 1b, we observe that 15 under EM, the upper edge of the Cu without nanotwins has a 16 single step-edge and developed into an array of surface step-17 edges with (111) terraces in changing a smooth surface to a 18 slightly curved (111) surface, see Figs. 2a-d. When electron 19 carriers are scattered by the surface step-edges, they transfer 20 their momentum and provide the energy needed for the change 21 of the surface morphology in Figure 2e-g. Later, the array of 22 surface step-edges fades away and becomes a flat(111) surface 23 again, as seen at the top of Figure 2h. The curved surface at the 24 bottom of the image in Figure 2h is connected to a twin 25 boundary. As a consequence, the free surface recedes deeper into the remaining grain, as shown by the schematic diagram in 26 27 Figure 1a. It is reported that (111) surface planes have the lowest activation energy and are the preferred diffusion pla**54** over any other surface plane^[17,18]. In addition, the electro**55** 28 29 30 flowing in <110> directions have less resistance than those 56 31 <100> or <111> directions, indicating that the current density in <110> directions is higher when an electrical field is appli**58** 32 ^[21]. Repeating step-edge diffusion on the (111) surface along 33 34 the <110> directions observed in Figure 2 leads to a unique 35 stepped atomic structure in crystalline Cu. After long periods 60 36 surface EM, there was a substantial change in surface 37 morphology of a Cu grain. Movie S1 shows the real tinge 38 movement of the step-edges on the nanotwin-free region during 39 the first 24 seconds in Figure 2. Additional evidence θ_1 concurrent EM-induced multiple step-edge migration in a twig5 40 41 free $(0\overline{1}1)$ -oriented Cu grain is shown in movie S2. 66

42 In contrast to the nanotwin-free surface, the lower part of the 43 surface of this $(0\overline{1}1)$ -oriented Cu grain, which has a pair 66 44 twin boundaries, has evolved into a surface with a zigzag shapped 45 as shown in Figure 3. The zigzag surface was found to hay 46 alternating ($\overline{1}11$) and (200) planes, see Figure 3a. Under EM2 47 we observe that the migration of step-edges in the regions



between triple points is similar to that in the twin-free region.

The surface step-edges indicated on the (200) plane (Figure

3b) are locations to see the initiated migration of the surface

step-edge and the removal of atomic layers with the current

Figure 2. Step-edge migrations with no nanotwins. (a-h) HRTEM images of the $(0\overline{1}1)$ -oriented Cu grain under current stressing as a function of time. The time of the image capture is given in the rectangular box at the upper-right corner. The first two numbers are in units of minutes, the second two numbers are in units of seconds, and the following two smaller numbers are in units of 1/30 seconds.

Under EM, the zigzag surface and the triple points tend to persist. As shown in Figs. 3c-d, the zigzag surface has evolved into ($\overline{111}$) and ($\overline{311}$)planes. From Figs. 3c to 3d, the change of alternating ($\overline{111}$) /(200) to alternating ($\overline{111}$) /($\overline{311}$) planes seems to have been affected by the advancing sunken (111) planes in the upper untwinned region (Figure 1b) due to the faster rate of EM. It led to a change in the relative orientation of the normal grain surface with respect to the flow direction of electrons. The change indicates that the geometry of the triple point or the three surface forces at the triple point may not be in thermodynamic equilibrium as given by the classical Young's equation. It is a kinetic effect of EM to alter the surface planes of the zigzag. Furthermore, it is assumed that a free surface has no normal stress. However, whether there is a normal strain in a

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- 1 ridge-type triple point is uncertain, and it is a subject that nee25
- 2 further study. Movies S3 and S4 show the direct observation **26** 3 the step-edge migration on a meta-stable zigzag surface **OT**
- 4 alternating $(\overline{1}11)$ and (200) planes and on alternating $(\overline{1}128)$
- 5 and $(\overline{3}11)$ planes, respectively.



Figure 3. Step-edge migrations and pinning of triple points with nanotwins. (a-66
HRTEM images of the (011)-oriented Cu grain with nanotwins under current
stressing as a function of time. The time of the image capture is given in the
rectangular box at the lower-left corner. The first two numbers are in units of
minutes, the second two numbers are in units of seconds, and the following two 70
smaller numbers are in units of 1/30 seconds.

13 Figure 4a shows the number of removed {111} atomic layers as a function of the current stressing for different Cu surface 14 15 planes. It quantitatively indicates EM-induced step-edge migrations at the surfaces with and without nanotwins. The 16 17 effect of the nanotwinned surface on the EM rate demonstrated by the significant change of the slope of the 18 19 measured curves compared to those of the untwinned surface see insets in Figure 4a. The EM rate to remove Cu atomic 20 layers at the free surface without nanotwins has a very sharp slope (Line a) compared to the gradual slopes of the nanotwig 21 22 modified zigzag surfaces. Because the nanotwinned region has a zigzag surface, Line b-1 and Line b-2 are the numbers of 23 24

removed Cu atomic layers counted from two ($\overline{1}11$) planes separated by $a(\overline{3}11)$ plane, and similarly, Line c-1 and Line c-2 are for the $(\overline{1}11)/(200)/(\overline{1}11)$ zigzag surface. In each of these step-wise curves, we can break each step into two periods: the very narrow width of the vertical line and the very wide width of the horizontal line. The very short period of the width of the vertical line represents the time of the propagation of an atomic layer, which shows no significant difference for all curves and it is less than 0.1 sec per atomic layer. However, the long period of the width of the horizontal line represents a stagnation time or incubation time between two vertical steps (the time it takes for the electromigration of a new free surface layer to occur); it is this long waiting time that greatly reduces the rate of EM. More twin boundaries can increase numbers of wide widths of the horizontal line and therefore are expected to slow down the EM-induced voiding.

The incubation time depends on the triple point configuration; the (111)/(200) pair has the longest incubation time and, in turn, the slowest EM rate. This is because it is harder to nucleate a step-edge on a low-index plane than on a high-index plane.

By taking the reciprocal of the slope of the curves, the average time required to remove one {111} Cu atomic layer in the nanotwinned regions can range from 3 times (3.85 sec) per atomic layer for the $(\overline{1}11) / (\overline{3}11)$ zigzag pairs to one order of magnitude larger (12.2 sec) per atomic layer for the $(\overline{1}11)$ /(200) zigzag pairs, compared to that (1.34 sec) per atomic layer on the $(\overline{1}11)$ planes without nanotwins. By computing the incubation period in the different stepwise curves captured in Figure 4a, frequency distribution curves of the incubation period for the step-edge migration are plotted in Figure 4b. The zigzag surfaces lead to an increase as well as a wider frequency distribution in the incubation period and consequently, a decrease in the rate of EM. Clearly, there is a strong dependence of the EM on the two alternating atomic planes of the zigzag surface. We found that the EM induces a feedback selection of the zigzag surface. The $(\overline{1}11)/(200)$ combination appears to be one of the most meta-stable zigzag pairs under EM, as is further demonstrated in the process of the refilling of a void of nanotwinned Cu (supplementary materials Figure S2).

i In the classical theory of EM, the atomic flux equation is

$$J_{EM} = CMF = C\frac{D}{kT}Z^*e\rho j$$

where C is the Cu concentration (C=1/ Ω_{Cu} atomic volume of Cu, Ω_{Cu} =1.182×10⁻²³ cm³), M is the atomic mobility and F is the EM force. From the step-edge displacement in Figure 4a, the atomic flux J can be calculated by the total atoms crossing a

unit area per unit time (Crossing area A=thickness $\times \delta$, where

 $\delta \approx 0.3$ nm is the effective diameter of the Cu atom). The divided time is the average period to remove one Cu atomic layer, which is the reciprocal of the slope of the stepwise curves (Figure 4a). The Cu step-edge diffusivity D can be obtained by giving the measured atomic flux and the defined driving force F. Here, we assume ρ_{Cu} (resistivity of Cu) =1.68×10⁻⁸ ohm·m, Z^* (effective charge number) is -5, and an average current density of 2×10⁶ A/cm² for computing the driving force. Our calculated step-edge diffusivity for the free surface without a nanotwin-modified structure is 3.8×10^{-10} cm²/sec, whereas for the nanotwin-modified surface, it ranges from 1.61×10^{-10}

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cm²/sec to 5.07×10^{-11} cm²/sec for the $(\overline{1}11)/(\overline{3}11)$ and the 1 $(\overline{1}11)/(200)$ pairs of the zigzag surface, respectively. Such a lowered diffusivity indicates that the step-edge diffusion 292 3 4 much more difficult on the nanotwin-modified surface. TBO diffusivities are 1.8×10^{-10} cm²/sec and 4.1×10^{-12} cm²/sec for tBe 5 high-symmetry (straight) and low-symmetry (rough) step-edgg2 reported in the literature, respectively^[16,19,20]. In the case of 6 7 partial symmetry^[16], the effective diffusivity can be calculated 8

9 as 1.4×10^{-11} cm²/sec.



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64 65 11 Figure 4. Measurement of EM rate in three surface planes. (a) Plot of the number removed {111} Cu atomic layers as a function of current stressing time in regions without and with nanotwins. (b) Frequency distribution curves of incubation period 204 12 13 step-edge migration computed from (a). The insets are the three corresponding HRTEA 14 15 images of the $(0\overline{1}1)$ -oriented Cu grain edges under EM. 69

71 16 It can be seen that our calculated results agree well with the 17 literature data, indicating that the assumptions used in this

18 calculation are reasonable.

19 Conclusions

20 In conclusion, we have demonstrated that the nanotwin-induced zigzag surface structures of Cu can be altered by EM. The 21 22 alternation by the feedback of EM can produce a high density 23 of the (111)/(200) pairs of zigzag surfaces on a polycrystalline 24 Cu interconnect. The correlation between the specific zigzag 25 surfaces and the magnitude of the mass flux driven by EM has 26 been studied by direct observations using in-situ ultrahigh

27 vacuum and high-resolution TEM.

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