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1 First evidence of the dramatic enhancement of the
2 reactivity of methyl formate (HC(O)OCH_3) with
3 OH at temperatures of the interstellar medium: A
4 gas-phase kinetic study between 22 K and 64 K

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17 **ABSTRACT.** Gas-phase chemistry of neutral-neutral reactions of interest in the interstellar
18 medium (ISM) is poorly-understood. The rate coefficients (k_{OH}) for the majority of the
19 reactions of hydroxyl (OH) radicals with interstellar oxygenated species are unknown at
20 temperatures of the ISM. In this work, we present the first determination of k_{OH} for
21 HC(O)OCH₃ between (22.4±1.4) and (64.2±1.7)K. The CRESU (French acronym for
22 *Cinétique de Réaction en Ecoulement Supersonique Uniforme* or Reaction Kinetics in a
23 Uniform Supersonic Flow) technique was used to create a chemical reactor with a uniform
24 temperature and gas density and the pulsed laser photolysis/laser induced fluorescence
25 technique was used to generate OH radicals and to monitor their temporal profile. It was
26 observed that $k_{\text{OH}}(T)$ increases one order of magnitude in only ~40K
27 ($k_{\text{OH}}(T=22\text{K})=(1.19\pm 0.36)\times 10^{-10}\text{ cm}^3\text{s}^{-1}$ and $k_{\text{OH}}(T=64\text{K})=(1.16\pm 0.12)\times 10^{-11}\text{ cm}^3\text{s}^{-1}$), and ~3
28 orders of magnitude with respect to $k_{\text{OH}}(T=298\text{K})$. This reaction is a very efficient loss
29 route for HC(O)OCH₃ in the gas phase and may have a great impact on the interpretation
30 by astrophysical models of the HC(O)OCH₃ abundances in cold regions of the ISM.

31

32 **1 Introduction**

33 Observations of interstellar material started in the last century and nowadays more than 180
34 species, besides the most abundant H₂ and He, are confirmed to be present in the interstellar
35 medium (ISM). Among them, radicals and neutral molecular species containing C, H, and
36 O, such as alcohols, aldehydes, acids and esters, have been detected in dense molecular
37 clouds ($T = 10\text{-}100$ K) and star-forming regions ($T > 100$ K). For instance, hydroxyl (OH)
38 radicals and methyl formate (HC(O)OCH₃) were detected in the ISM molecular clouds.¹⁻⁸
39 The observed abundances relative to H₂ in hot cores and hot corinos were reported to be
40 3×10^{-8} and 1.4×10^{-8} (in Orion), 4×10^{-7} (in IRAS 16293-2422) and 7×10^{-7} (IRAS 4A). The
41 diverse abundances can be attributed to differences in physical conditions (for instance, the
42 temperature evolution during the warm-up timescale of the object) that affect the chemistry.
43 Recent detections of methyl formate in colder objects, such as protostars Barnard 1b,⁹ in
44 cold environment of IRAS 16293-2422¹⁰ and prestellar cores¹¹ demonstrate that
45 understanding the formation of complex organic molecules (COMs) in these environments
46 is one of the key issues to model properly the observed abundances of these molecules in
47 the ISM.

48 The formation mechanism of HC(O)OCH₃ has been the subject of several
49 studies,¹²⁻¹⁹ but it is not yet well-understood. In hot cores and corinos, it is believed that
50 some organic molecules are formed in grain-surfaces and subsequently injected into the gas
51 phase during the warm-up phase. Once in the gas-phase these molecules are ionized being
52 the precursors of other COMs, such as HC(O)OCH₃. Some models exclusively use gas
53 phase reactions, for example, Horn et al.¹² proposed several formation routes for
54 HC(O)OCH₃ via the recombination reaction of a protonated and a neutral species followed

55 by dissociative recombination with electrons. Nevertheless, this gas-phase chemical model
56 underpredicts the observed abundances of HC(O)OCH₃ in hot cores at $T \sim 100$ K. More
57 recently, Balucani *et al.*¹⁸ proposed a new chemical model for the gas-phase formation of
58 HC(O)OCH₃ in colder environments, based on the fact that ejection of CH₃OH takes place
59 from the outer shell of the cold prestellar core ($T \sim 10$ K) contrary to previous assumptions.
60 These authors concluded that although the gas-phase formation of methyl formate cannot
61 be neglected, grain-surface formation is an important formation route for this COM. In gas-
62 grain models, it is assumed that the formation of HC(O)OCH₃ occurs via the reaction of
63 methoxy (CH₃O) and formyl (HCO) at temperatures around 30-40 K, afterwards it is
64 gradually evaporated into the gas phase.^{13,14} The evidence that the CH₃O+HCO reaction
65 can occur in interstellar granular ices was experimentally proven by Bennett and Kaiser¹⁵
66 and Modica and Palumbo.¹⁹ What is clear is that gas-phase formation of HC(O)OCH₃ is
67 interconnected with grain-surface formation, but what about the removal processes of
68 HC(O)OCH₃ in the gas phase at temperatures of cold objects? Evidently, the gas-phase loss
69 processes of this COM need to be included as well in chemical models to properly interpret
70 the observed abundances. Among them, ultraviolet (UV) photodissociation processes and
71 radical-COM reactions have to be taken into account in the models. UV photodissociation
72 of methyl formate has been investigated at room temperature (yielding mainly CH₃O and
73 HCO radicals),²⁰⁻²² while the chemical kinetics for the gas-phase reaction of OH radicals
74 with methyl formate,



76 has been investigated at temperatures between 233 and 1500 K and pressures between 0.4
77 and 1.63 atm.²³⁻²⁹ Up to date no chemical kinetic studies of HC(O)OCH₃ have been

78 reported at lower temperatures. Therefore, the aim of this work is to report for the first time
79 k_{OH} for reaction (1) between 22 and 64 K, temperatures of interest in the ISM.

80 **2 Experimental**

81 **2.1 General Aspects**

82 The CRESU system employed in this work is schematically shown in **Figure 1**. This setup,
83 which was exhaustively described by Jiménez *et al.*³⁰, couples the uniform supersonic gas
84 expansion technique, to achieve the ultra-low temperatures needed for determining $k_{\text{OH}}(T)$
85 between 22 and 64 K, and the pulsed laser photolysis-laser induced fluorescence (PLP-LIF)
86 kinetic technique.

87 Ultra-low temperatures are achieved by the isentropic expansion of a gas mixture,
88 consisting in the buffer gas, the OH-precursor (H_2O_2), and diluted $\text{HC}(\text{O})\text{OCH}_3$, through a
89 Laval nozzle from a high pressure reservoir (P_{res}) to a vacuum chamber (P_{cham}). In this
90 work, the Laval nozzle described in reference 30 was used for getting all temperatures. This
91 nozzle was initially designed to obtain 22 K in helium (conditions **A** in **Table 1**). However,
92 calculations performed with a home-made program, which take into account the existing
93 profile of the divergent-convergent nozzle, indicate that other uniform supersonic flows
94 could be obtained by varying the pressure and temperature conditions as well as the mass
95 flow rates (F) for different buffer gases (helium, nitrogen or a mixture of both). The
96 optimal P_{res} and P_{cham} and ranges of the calibrated mass flows for the buffer gas (F_{buffer}), the
97 buffer gas through the H_2O_2 bubbler ($F_{\text{H}_2\text{O}_2}$) and diluted $\text{HC}(\text{O})\text{OCH}_3$ (F_{R}) are summarized
98 in **Table 1**. In the case of conditions **B** (see **Table 1**), the buffer gas is a mixture of He
99 (40%) and N_2 (60%) independently introduced in the reservoir. As described elsewhere,³⁰⁻³²

100 the buffer gas is flowing through a glass bubbler containing a pre-concentrated aqueous
101 solution of H₂O₂. In case **B**, N₂ was only flowed through the bubbler. Diluted methyl
102 formate was prepared in two 20-L storage bulbs connected in series. Mixing ratios of
103 methyl formate (*f*) in the 40-L storage bulb ranged from *ca.* 2×10⁻³ to 8×10⁻².

104 The aerodynamic chopper³⁰ (rotary disk at 10 Hz) placed at the divergent part of the
105 Laval nozzle was operational under experimental conditions **A** (hitherto pulsed mode for
106 getting 22 K). The pulsing of the gas was not necessary to achieve higher temperatures
107 (conditions **B-D**, continuous mode), since the required mass flows and working pressures
108 were accessible with the pumping capacities of our experimental system. In the generated
109 supersonic gas flow the temperature (*T*), the total gas density (*n*), and the velocity are
110 uniform along the flow.

111

112 **2.2. Aerodynamical characterization of the supersonic flows: Pitot tube measurements**

113 The uniformity of these parameters along the axis of the expansion was confirmed by
114 measuring the impact pressure, *P_i*, in the commonly named *Pitot tube* measurements, as a
115 function of the distance from the exit of the nozzle. The Pitot tube used in the continuous
116 mode consisted in a stainless steel tube (5 cm length) with an orifice of 2 mm, which was
117 connected to a 10 or 100 Torr pressure transducer. It was fixed on two manual translational
118 platforms for aligning the Pitot tube with the center of the gas flow. This system was
119 mounted on a motorized long-travel linear translational platform (Standa, 8MT195-740-
120 2.5).

121 As described in reference 30, the temperature, pressure and gas density of the jet
 122 along the uniform supersonic flow were derived from the Mach number, M , which is
 123 obtained from the average P_1 obtained in the optimal length where the flow is uniform:

$$124 \quad \frac{P_i}{P_{res}} = \left(\frac{(\gamma + 1)M^2}{(\gamma - 1)M^2 + 2} \right)^{\frac{\gamma}{\gamma - 1}} \left(\frac{\gamma + 1}{2\gamma M^2 - \gamma + 1} \right)^{\frac{1}{\gamma - 1}} \quad (2)$$

125 The heat capacity ratio γ for the buffer gas and M values are shown in **Table 2**. In particular
 126 for the gas mixture He/N₂, $\gamma_{mixture}$ was calculated considering the relationship between the
 127 heat capacity of the gas mixture at constant pressure ($C_{p,mixture}$) and at constant volume
 128 ($C_{v,mixture}$),

$$129 \quad C_{p,mixture} = R + C_{v,mixture}, \quad (3)$$

130 where $C_{v,mixture}$ includes the fraction of N₂ (60%) and He (40%) in the mixture. Therefore,
 131 $\gamma_{mixture}$ can be expressed in terms of the individual C_v of the gases present in the mixture
 132 ($C_{v,N_2} = 20.785 \text{ J mol}^{-1} \text{ K}^{-1}$ and $C_{v,He} = 12.471 \text{ J mol}^{-1} \text{ K}^{-1}$):

$$133 \quad \gamma_{mixture} = \frac{R}{0.6 \times C_{v,N_2} + 0.4 \times C_{v,He}} + 1 \quad (4)$$

134 Temperature and gas densities were then obtained as described in reference 30. As an
 135 example, **Fig 2** shows the spatial evolution of these two parameters for conditions **B** (see
 136 **Figs S1** and **S2** of the supporting information for conditions **C** and **D** and reference 30 for
 137 conditions **A**). As can be seen in the figures, the temperature and gas density of the jet are
 138 uniform for several tens of centimeter from the exit of the nozzle, as well as it prolongs
 139 three centimeters inside the divergent part of the nozzle. This length of uniformity is
 140 sufficient to carry out the kinetic experiments in a timescale of several hundreds of

141 microseconds (hydrodynamic time, t_{hydro} , **Table 2**) and the methyl formate concentrations
142 used (see **Table 3**).

143

144 **2.3 PLP-LIF kinetic technique**

145 The conventional PLP-LIF technique has been widely employed in our laboratory for
146 kinetic studies on gas-phase reactions between OH radicals and atmospheric pollutants, as
147 previously described.^{31,32} The OH radical was produced *in situ* by UV photolysis of H₂O₂
148 by a KrF excimer laser at 248 nm. The laser fluence measured by a calorimetric disk at the
149 exit of the Laval nozzle was *ca.* 1 mJ/(cm² pulse) for most of the experiments at 10 Hz. The
150 temporal evolution of OH(X²Π) was monitored by LIF at *ca.* 310 nm after excitation of OH
151 at *ca.* 282 nm using a frequency doubled dye laser (Lambda Physik, model Scanmate)
152 pumped by a XeCl excimer laser (Lambda Physik, model LPX105i). The LIF detection
153 arrangement consists of an optical system coupled with a filtered photomultiplier tube,
154 PMT. The band pass filter is centered at 310 nm with a full width at half maximum of 10
155 nm. The electrical signal from the PMT was transferred to a gated boxcar integration unit
156 and the integrated signal was recorded and processed into a computer by a home-made
157 LabView program.

158 OH radicals were lost mainly by reaction with methyl formate and at a lesser extent
159 by reaction with H₂O₂ (k_{prec}) and other processes such as diffusion of OH radicals out of the
160 detection zone and/or reaction with impurities, if present in the sample. As mentioned in
161 section 2.2, the timescale of the kinetic experiment was restricted by the uniformity of the
162 supersonic flow. For that reason, at all temperatures the maximum reaction time is the

163 hydrodynamic time, mentioned in **Table 2**. The reaction time during the kinetic
164 measurements was varied by changing the delay between the photolysis and the probe
165 lasers from *ca.* 20-40 μs before the trigger of the photolysis laser (to record the background
166 signal) to the hydrodynamic time.

167

168 **2.4 Chemicals**

169 He (99.999%, Praxair) and N_2 (99.999%, Praxair) were used as supplied. Liquid sample of
170 anhydrous HC(O)OCH_3 (99%, Sigma-Aldrich) was placed in a flask ($V = 250 \text{ mL}$)
171 degassed by repeated freeze-pump-thaw cycles prior to its use. During the preparation of
172 diluted methyl formate in the buffer gas, the flask containing the liquid was introduced in a
173 beaker with water at room temperature to thermalize the sample, since there is a small gas
174 expansion and, consequently, temperature and vapor pressure of the methyl formate
175 decrease. Aqueous solution of H_2O_2 (Sharlab, initially at 50% w/v) was pre-concentrated as
176 described earlier.^{31,32}

177

178 **3 Results and Discussion**

179 **3.1. Determination of $k_{\text{OH}}(T)$**

180 The kinetic study of reaction (1) was performed under *pseudo*-first order conditions
181 ($[\text{OH}]_0 \ll [\text{HC(O)OCH}_3]$ and $[\text{H}_2\text{O}_2]$). $[\text{HC(O)OCH}_3]$ concentration was calculated from the
182 dilution factors, the mass flow rates, and T and P of the jet. Concentration ranges used in
183 the kinetic analysis are shown in **Table 3**. In order to check that the *pseudo*-first order
184 conditions were accomplished, the gas-phase concentration of H_2O_2 in the jet was optically
185 measured before entering the reservoir (at room temperature) by UV spectroscopy at 254

186 nm using a gas cell with a path length of 108 cm. After correction of temperature and
187 pressure (for conditions **D**), $[\text{H}_2\text{O}_2]$ in the jet was found to be lower than 10^{13} cm^{-3} in this
188 case. Accurate calculation of $[\text{OH}]_0$ would require the knowledge of absorption cross
189 section of H_2O_2 at 248 nm and the OH quantum yield at lower temperatures than 298 K.
190 Nicovich and Wine³³ measured the temperature dependence of σ_λ for $230 \text{ nm} \leq \lambda \leq 295 \text{ nm}$
191 between 285 and 381 K. No appreciable change was observed in σ_λ in this wavelength
192 range. In any case, at the photolysis laser fluence used and considering the absorption cross
193 section of H_2O_2 and quantum yield of OH at room temperature,³⁴ the initial OH
194 concentration between 22 K and 64 K is expected to be several orders of magnitude lower
195 (below 10^{10} cm^{-3}) than that of its precursor. Consequently, the *pseudo*-first order conditions
196 are achieved at all reagent concentrations.

197 In the absence of OH-precursor, Le Calvé *et al.*²⁶ detected a fluorescence emission
198 at 310 nm attributed to CH_3O radicals formed in the photolysis of methyl formate, $(1.94\text{-}$
199 $6.85) \times 10^{15} \text{ cm}^{-3}$, at 248 nm. In the present work, nonetheless, no emission was observed in
200 the photolysis of methyl formate at 248 nm in the absence of H_2O_2 at the concentrations
201 and laser fluences employed. Therefore, the recorded LIF temporal profiles correspond
202 exclusively to OH radicals. **Fig 3** shows an example of the temporal evolution of OH
203 radicals in the presence of methyl formate recorded at 43 K. The temporal profiles of the
204 OH LIF signal (I_{LIF}) were analyzed after 20-40 μs depending on the methyl formate
205 concentration, because some rotational relaxation occurs at shorter times. At all
206 temperatures of the jet, I_{LIF} did not fully decay in the available timescale at the methyl
207 formate concentrations, which had to be low enough to avoid secondary chemistry (as
208 discussed below).

209 Under these conditions and after rotational relaxation of OH, the non-linear least
210 square fit of the LIF data to a single exponential function (solid line in **Fig 3**) yields the
211 pseudo-first order rate coefficient, k' , for a fixed concentration of methyl formate (and
212 H_2O_2):

$$213 \quad k' = k_{\text{OH}}(T)[\text{HC}(\text{O})\text{OCH}_3] + k_0 \quad (5)$$

214 where $k_0 (= k_{\text{prec}}(T)[\text{H}_2\text{O}_2] + k_{\text{other losses}})$ is the *pseudo*-first order coefficient obtained in the
215 absence of $\text{HC}(\text{O})\text{OCH}_3$. During a kinetic experiment at a single temperature, $[\text{H}_2\text{O}_2]$ was
216 kept constant. Averaged k_0 (in s^{-1}) were 6400 at 22 K, 3700 at 43 K, 2100 at 52 K and 5100
217 at 64 K. Individual $k_{\text{OH}}(T)$ were obtained from the slope of the plots of the *pseudo*-first
218 order rate coefficients k' corrected by k_0 against the $\text{HC}(\text{O})\text{OCH}_3$ concentrations.

$$219 \quad k' - k_0 = k_{\text{OH}}(T)[\text{HC}(\text{O})\text{OCH}_3] \quad (6)$$

220 All obtained $k' - k_0$ are plotted as a function of reactant concentration in **Fig 4** for all
221 experiments carried out. As it can be seen in **Fig 4.a**, at 22 K a curvature in the plot of $k' - k_0$
222 versus $[\text{HC}(\text{O})\text{OCH}_3]$ was observed above $3 \times 10^{13} \text{ cm}^{-3}$, indicating that an additional loss of
223 methyl formate is occurring by complexing formation in the ultracold jet ($T=22 \text{ K}$). This
224 effect was also observed in the reaction of OH with 1-butene at 22 K.³⁰ For that reason the
225 kinetic analysis was constrained to the range of $[\text{HC}(\text{O})\text{OCH}_3]$ where *pseudo*-first order
226 plots were linear. As can be seen in **Table 3** and **Fig 4**, higher concentrations of
227 $\text{HC}(\text{O})\text{OCH}_3$ could be accessible at higher temperatures, because the formation of
228 aggregates is kinetically less efficient at the total gas density of these experiments with
229 respect to 22 K.

230

231 **3.2. Temperature Dependence of $k_{\text{OH}}(T)$: Enhancement at ultra-low temperatures**

232 $k_{\text{OH}}(T)$ values obtained are listed in **Table 3** and depicted in **Fig 5**. As it can be seen, the
233 reactivity of HC(O)OCH_3 toward OH radicals is high in the investigated temperature range,
234 decreasing rapidly from 22 to 64 K. In **Table 3**, the enhancement factor relative to
235 $k_{\text{OH}}(64\text{K})$ is also shown for all temperatures. The observed enhancement of $k_{\text{OH}}(T)$ over 40
236 K in the T -range of interest in the ISM is of one order of magnitude. The enhancement
237 factor is *ca.* 4 and 2 at 43 and 52 K, respectively.

238 The observed temperature dependence of $k_{\text{OH}}(T)$ between 22 and 64 K was first fitted
239 to the following expression commonly used in kinetic databases employed in astrophysical
240 models:³⁵

$$241 \quad k_{\text{OH}}(T) = \alpha (T/300 \text{ K})^\beta \exp(-\gamma/T) \quad (7)$$

242 The fit to this three-parameter expression yielded a $\gamma \sim 0$ K. Therefore, we deliberately
243 fixed this parameter to 0 K without affecting significantly the accuracy of the fit. As $\beta < 0$
244 (see **Fig 5**), extrapolation of Eq.(7) to temperatures close to absolute zero ($T < 22$ K)
245 provides however unrealistic results ($T \rightarrow 0$ K, $k_{\text{OH}}(T) \rightarrow \infty$). Interestingly, a single
246 exponential expression of the form,

$$247 \quad k_{\text{OH}}(T) = a \exp(-b \times T), \quad (8)$$

248 describes perfectly the observed T -dependence of $k_{\text{OH}}(T)$ in the investigated range. The
249 resulting parameters a and b are given in **Fig 5**. Note that even though Eq.(8) has no
250 physical meaning, it predicts a reasonable k_{OH} at 10 K ($2.38 \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$) and at 0 K
251 ($3.84 \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$). Of course, the “real” kinetic behavior may be different as that
252 predicted, hence Eq.(8) should be exclusively used for the temperature range 22-64 K.

253 As shown in **Table 4** and **Fig 6**, the trend in the temperature dependence of $k_{\text{OH}}(T)$
254 is noticeably different at $T < 64$ K (of interest in the ISM, this work), between 233 and 372 K

255 (of interest in terrestrial atmosphere)^{23-27,29} and at higher temperatures (of interest in
256 combustion chemistry).^{27,28} At room temperature, Wallington *et al.*²³, Good *et al.*²⁴ and
257 Szilagyi *et al.*²⁵ reported experimental rate coefficients $k_{\text{OH}}(T)$ in quite good agreement
258 with each other. The typical measured value was about $2 \times 10^{-13} \text{ cm}^3 \text{ s}^{-1}$. Le Calvé *et al.*²⁶
259 extended the kinetic study of reaction (1) to temperatures between 233 and 372 K. These
260 authors observed a small positive temperature dependence with an activation energy of *ca.*
261 1 kcal/mol. At temperatures of interest in combustion chemistry, Lam *et al.*²⁷
262 experimentally reported a much stronger temperature dependence with activation energy of
263 4 kcal/mol between 880 and 1337 K.

264 From a theoretical point of view, a few works can be found in the literature. In 1999
265 Good *et al.*²⁴ used *ab initio* molecular orbital theory to calculate the rate coefficient at 298
266 K obtaining a value of $1.88 \times 10^{-13} \text{ cm}^3 \text{ s}^{-1}$, in excellent agreement with their experimental
267 contribution and other studies. The rate coefficient was calculated by transition state theory
268 (TST) taking into account the Wigner correction accounting for the effect of tunneling.
269 More recently, Tan *et al.*²⁸ performed *ab initio* calculations in the temperature range 300-
270 1500 K, and found that $k_{\text{OH}}(300\text{K}) \sim 8 \times 10^{-15} \text{ cm}^3 \text{ s}^{-1}$ which is lower, by more than one order
271 of magnitude, than the experimental ones. TST in conjunction with the CANTHERM code
272 was used by Tan *et al.*²⁸ to calculate these rate coefficients and the Eckart correction was
273 applied to take into account tunneling. A third study was provided by Elm *et al.*²⁹ The
274 originality of their work was to calculate the rate coefficient at 298 K and 1 atm for several
275 functionals and tunneling corrections using the conventional TST. They concluded that
276 BH&HLYP and MP2 functionals were those leading to the best agreement with
277 experiments, provided that the tunneling Bell correction was used. Interestingly, the rate

278 coefficient was increased by a factor of 36 (BH&HLYP) or 90 (MP2) when the Bell
279 correction was applied to their UCCSD(T)-F12a single point corrected $k_{\text{OH}}(298\text{K})$
280 indicating the very large effect of tunneling in this reaction.

281 To our knowledge, no attempt has been undertaken to determine rate coefficients at
282 sub-ambient temperatures using these theoretical methods. As it can be seen in **Fig 6**, an
283 increase of roughly three orders of magnitude for $k_{\text{OH}}(T)$ is observed between 298 K and 22
284 K, whereas the rate coefficient shows a classical Arrhenius increase at higher
285 temperatures.²⁷ This kind of behavior has been seen for the first time only very recently for
286 other oxygenated compounds³⁶⁻⁴⁰ reacting with OH. This spectacular change in k_{OH} could
287 be explained by a change in the mechanism driving the dynamics of the reaction comparing
288 to the suprathreshold domain.

289

290 **3.3. Reaction mechanisms and products**

291 The reaction pathways of the OH + CH₃C(O)OH reaction have been investigated by
292 Good *et al.*²⁴ and more extensively by Tan *et al.*²⁸. In the latter study, the authors found that
293 the reaction proceeds through H abstraction at either the methyl group or the formyl group
294 of methyl formate, leading to the exothermic formation of CH₂C(O)OH + H₂O or
295 CH₃C(O)O + H₂O, respectively. Four reaction pathways were identified, all starting with
296 the reversible formation of a weakly bound pre-reactive complex followed by an
297 irreversible reaction through a transition state to the above mentioned products. As two pre-
298 reactive complexes have been found to stand in a very shallow well (~1 kcal/mol), only
299 those channels passing through the deeper wells (3.7 and 3.4 kcal/mol, respectively) have
300 been considered by the authors. The barrier heights for the two corresponding transition

301 states were 3.5 and 4.0 kcal/mol respectively. This situation is qualitatively similar to that
302 observed for the reaction of OH with methanol^{37,39}, acetone^{36,38} and dimethyl ether^{36,38} and
303 can qualitatively explain our observations. The pre-reactive complexes are expected to be
304 short-lived at "high" temperatures with the propensity to break back to the reactants,
305 whereas the activation barrier to the transition state leads to a standard Arrhenius behavior.
306 At lower temperatures, the pre-reactive complex lifetime should raise giving a chance to the
307 rate coefficient to increase significantly either by adduct stabilization (three body process)
308 or tunneling through products (binary process). Although the present experimental study
309 cannot discriminate from these two options, the theoretical work carried out by Elm *et al.*²⁹
310 showed that tunneling has a major contribution, at least at 298 K. This should be even
311 stronger at lower temperatures. Note that terms such as "high" and "low" are presently quite
312 vague, because there is no sufficient data available to identify the temperature at which the
313 rate coefficient would reach its minimum value. Further experimental and theoretical works
314 are evidently needed to make these aspects clearer.

315

316 3.4. Branching ratios

317 Another aspect to consider is the potential organic radical formation in reaction (1).
318 Quantitative information concerning the products of reaction (1) cannot be derived from the
319 present experiments, nevertheless some comments are worthwhile. As mentioned earlier, in
320 the titled reaction, hydrogen atom abstraction by OH radicals can occur at either the formyl
321 group (channel 1a) or the methyl group (channel 1b):



324 At 298 K, Le Calvé *et al.*²⁴ predicted, based on structure-activity (SAR) calculations that
325 both channels accounted for *ca.* 50% of the overall rate coefficient k_{OH} . From *ab initio*
326 calculations, the predicted branching ratio k_{1a}/k_{1b} at room temperature greatly differs. Good
327 *et al.*²⁵ reported that that formation of C(O)OCH₃ radicals (channel 1a) was the major
328 reaction pathway, accounting for 86%, while Tan *et al.*²⁸ predicted that formation of
329 HC(O)OCH₂ radicals (channel 1b) was the major reaction pathway, accounting for 66% of
330 k_{OH} . At higher temperatures (*e.g.* 1168 K), Lam *et al.*²⁷ estimated that channel (1b) is the
331 main reaction pathway accounting for 68% of overall k_{OH} , which is lower than that derived
332 from *ab initio* calculations by Tan *et al.*²⁸ (*ca.* 54% for channel 1b). In light of these
333 discrepant results and given that the formation of C(O)OCH₃ and HC(O)OCH₂ radicals is
334 not quantitatively known, we cannot assert which channel could become more important at
335 lower temperatures than 298 K. However, by comparison with the reaction of OH with
336 CH₃OH, it could be expected a change in the yield of the reaction products at low
337 temperatures with respect to that found at room temperature.³⁷ According to master
338 equation calculations carried out by Shannon *et al.*, the H-atom abstraction by OH radicals
339 seems to exclusively occur at low temperatures from the –OH group of methanol.³⁷ These
340 calculations show that the branching ratio for the CH₃O + H₂O channel reaches unity at
341 temperatures lower than 70 K. In contrast, at room temperature the main reaction product is
342 CH₂OH radical formed through the H-abstraction from methyl group of methanol.⁴¹ Further
343 mechanistic studies on reaction (1), especially at low temperatures are indispensable to
344 check if the C(O)OCH₃ radical can be the main reaction product at temperatures of the
345 ISM.

346

347 4. Conclusions and astrophysical impact

348 The present work constitutes the first determination of the rate coefficients for the gas-
349 phase reaction of methyl formate and OH radicals at ultra-low temperatures. Our results
350 show that, at extremely low temperatures, the investigated reaction (1) is a very efficient
351 loss route for HC(O)OCH₃ in the gas phase, up to three orders of magnitude faster than at
352 room temperature. Acharyya *et al.*⁴² have recently evidenced the importance of determining
353 gas-phase rate coefficients at temperatures found in ISM molecular clouds (close to 10 K)
354 for the reaction of OH with methanol. The acceleration of k_{OH} for this reaction has a
355 marked influence on the modeled abundances of CH₃O radicals. For that reason, the
356 inclusion of our results on the OH+methyl formate in astrochemical models may have a
357 great impact and would help to improve/shed light on the interpretation of the detected
358 abundances of methyl formate in the various environments of the ISM. These
359 measurements can also stimulate attempts to detect the radicals issued from the studied
360 reaction, which have not been looked for until now.

361 The temperature dependence of reaction (1) has been found to be similar to that
362 observed for other oxygenated species reacting with OH³⁵⁻³⁹. With the exception of the
363 reaction of OH with methanol, Shannon *et al.*³⁷ and Caravan *et al.*³⁹ reported that a pressure
364 dependence of $k_{\text{OH}}(T)$ was observed at gas densities between 3.2×10^{16} and 2.6×10^{17} cm⁻³ in
365 the temperature range *ca.* 88 and 140 K^{37,39}. Despite the present work cannot provide any
366 evidence of a pressure dependence of the rate coefficient k_{OH} , this will be explored in the
367 near future with specific Laval nozzles, operating at different gas densities for a same
368 temperature.

369 Further CRESU experiments are also needed at different temperatures to have a
370 complete picture of the kinetic behavior of reaction (1) at sub-ambient temperatures. For
371 this purpose, new Laval nozzles will be constructed in our laboratory to probe temperatures
372 below 20 K and above 64 K.

373

374

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383

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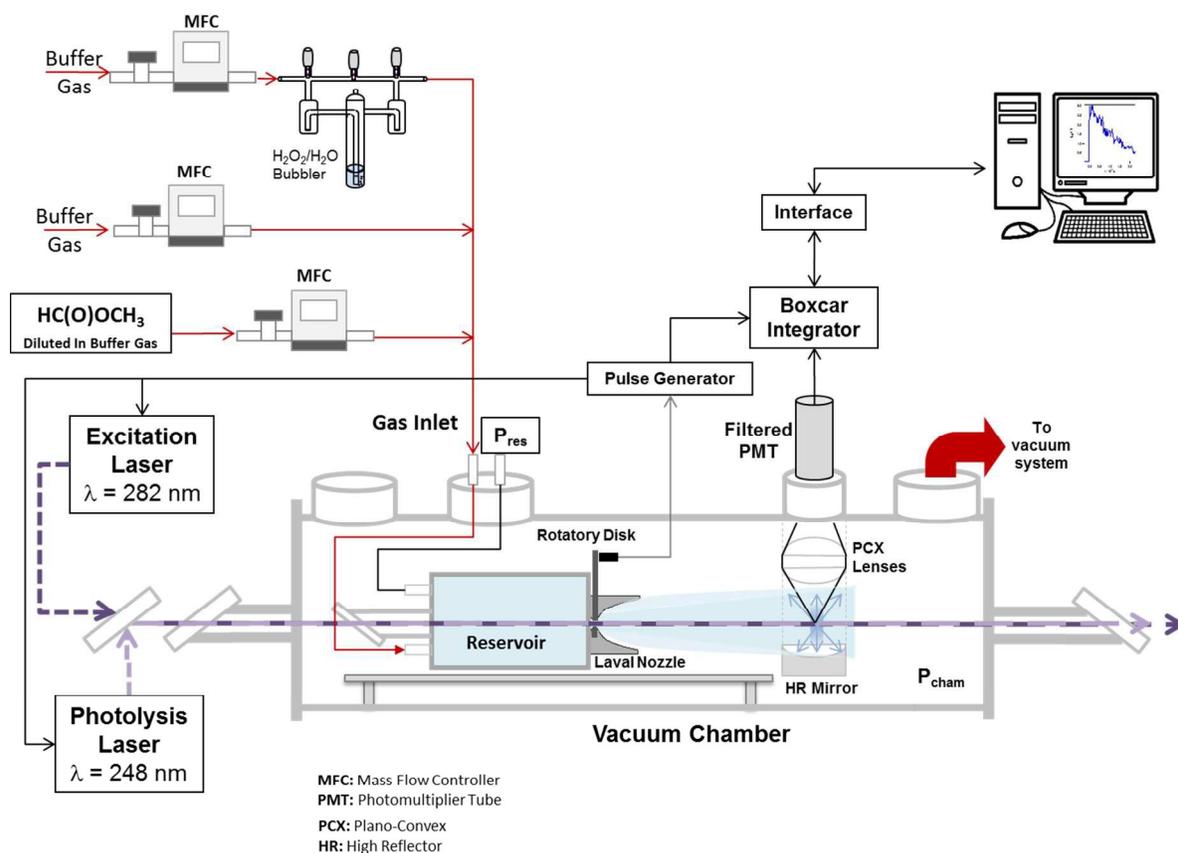
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508

509 **Figures**

510

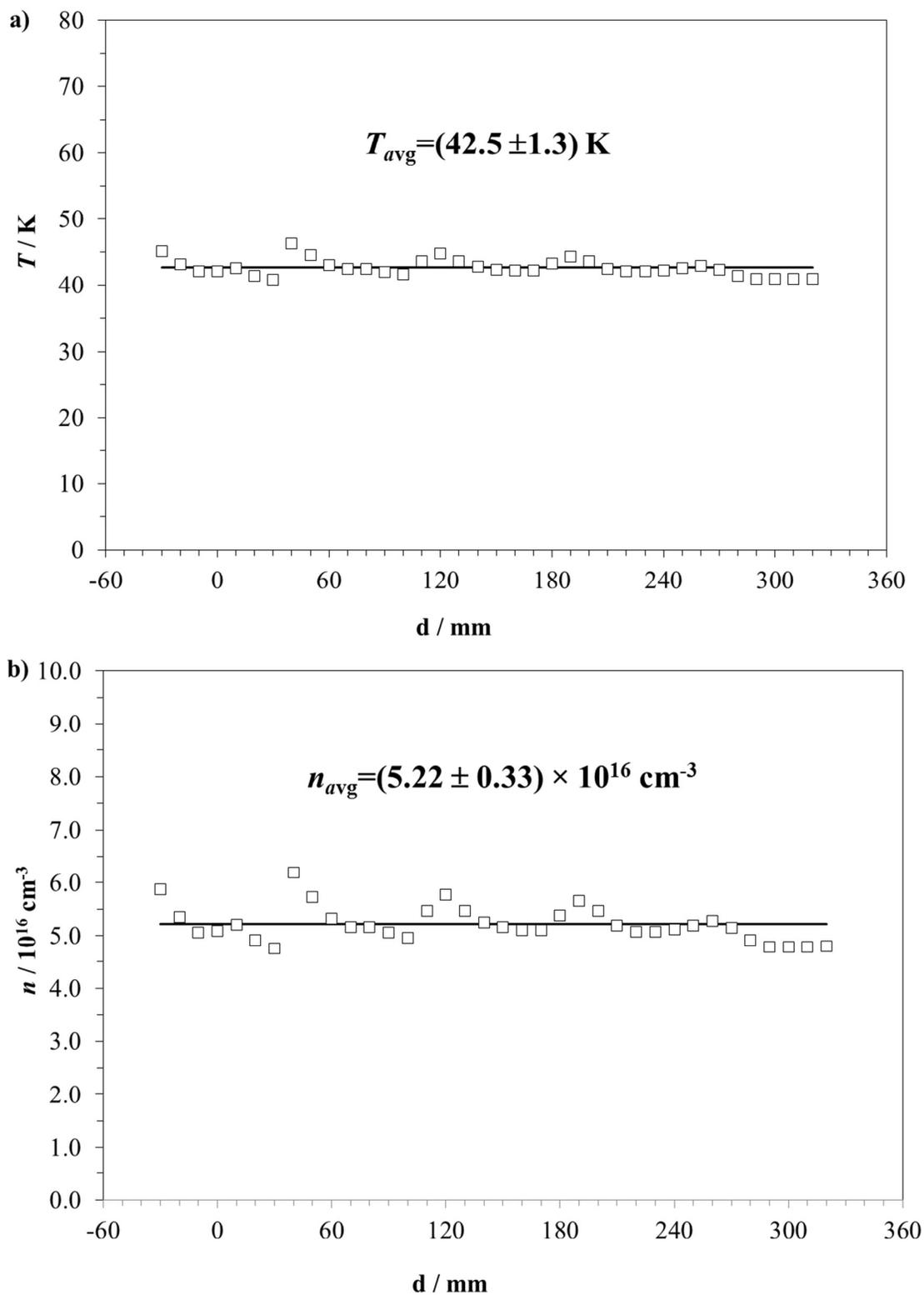


511

512 **Fig 1.** Schematics of the experimental system employed: CRESU apparatus coupled to

513 PLP-LIF.

514



515

516 **Fig 2.** Spatial profiles of the temperature and total gas density of the jet for experimental
517 conditions **B**. Uncertainty is $\pm\sigma$.

25

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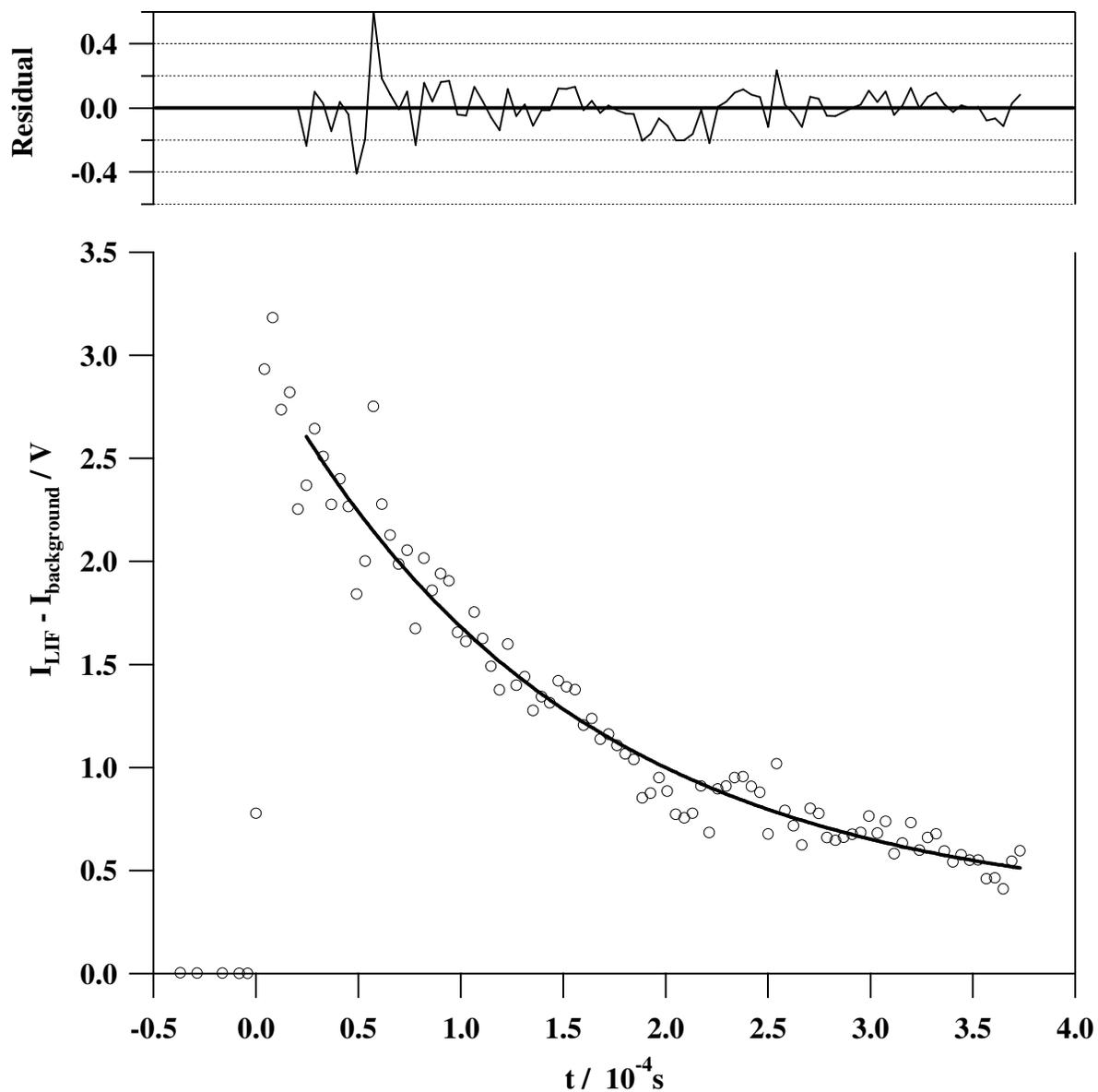


Fig 3. Example of a temporal profile of the LIF signal from OH radicals recorded at $T = 43$ K in the presence of methyl formate ($1.86 \times 10^{13} \text{ cm}^{-3}$). The fit started at $20 \mu\text{s}$ to ensure fully rotational relaxation of OH.

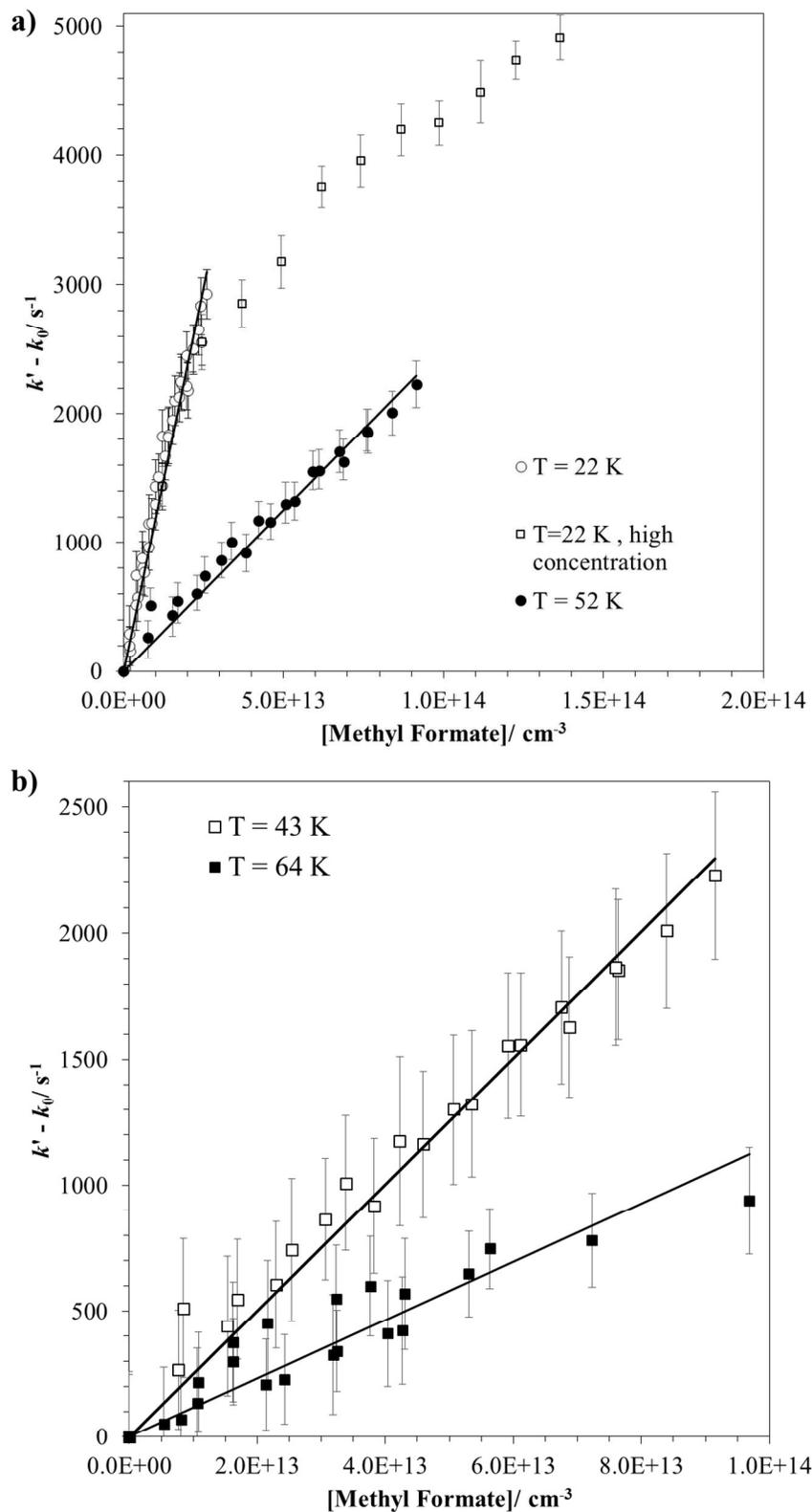


Fig 4. Plots of corrected *pseudo*-first order rate coefficient *versus* methyl formate concentration.

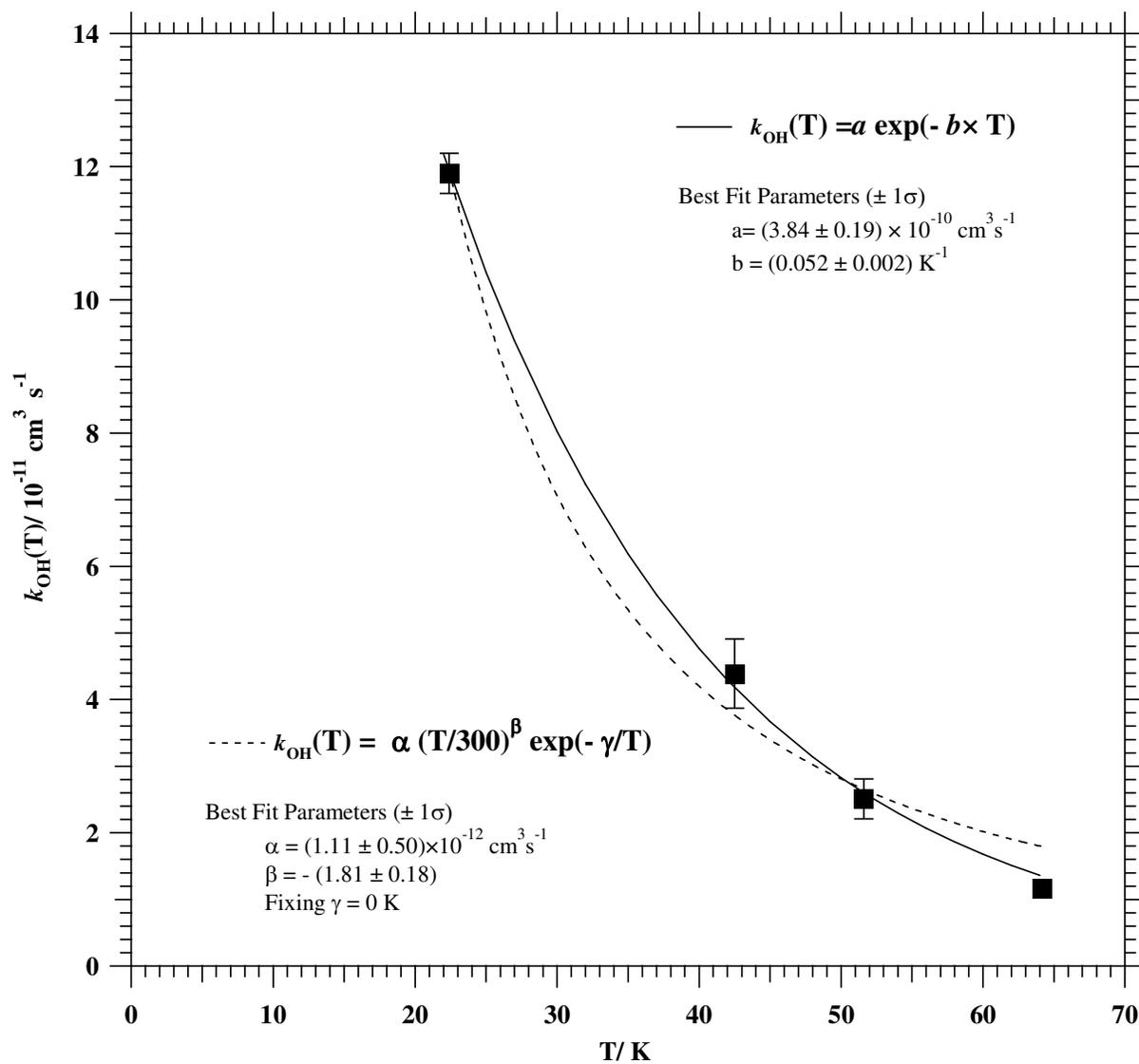


Fig 5. Temperature dependence of $k_{\text{OH}}(T)$ for the reaction of OH with methyl formate between 22 and 64 K.

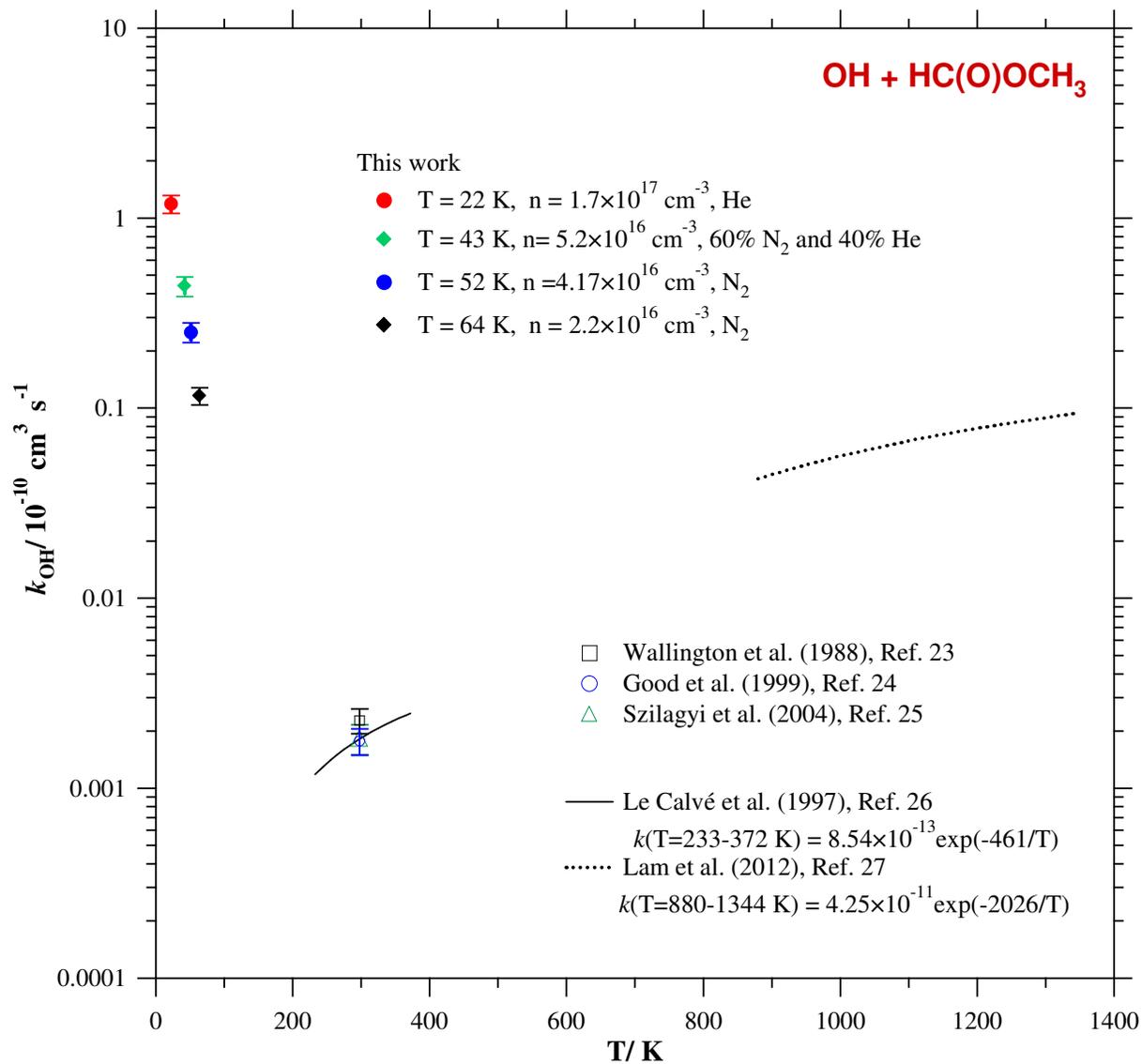


Fig 6. Temperature dependence of $k_{\text{OH}}(T)$ for the investigated reaction (1) between 22 and 1344 K.

Tables

Table 1. Summary of the experimental conditions employed in this work ($T_{\text{res}} = 297 \pm 1$ K).

Conditions	Gas Flow	Buffer Gas (%)	$F_{\text{He}}/$ slpm	$F_{\text{N}_2}/$ slpm	$F_{\text{buffer}}/$ slpm	$F_{\text{H}_2\text{O}_2^*}/$ sccm	$F_{\text{R}}/$ sccm	$f \times 10^{-3}$	$P_{\text{res}}/$ mbar	$P_{\text{cham}}/$ mbar
A	Pulsed	He (100%)	8.4	0	8.4	95-190	48-592	1.9-12	337.0	0.621
B	Continuous	N ₂ (60%) He (40%)	6.0	9.0	15.0	190	45-550	37-77	127.1	0.296
C	Continuous	N ₂ (100%)	0	12.6	12.6	190	47-562	50-55	136.2	0.279
D	Continuous	N ₂ (100%)	0	4.4	4.4	47	48-573	19-38	41.67	0.183

* Flow of buffer gas through the bubbler containing the aqueous solution of OH precursor; F_{R} is the flow rate of the diluted methyl formate through the reservoir; f is the dilution factor of the methyl formate in the storage bulb.

Table 2. Characteristics of the gas flows: average impact pressure (P_i), heat capacity ratio (γ), Mach number (M), temperature of the jet (T), hydrodynamic time (t_{hydro}) and the length of the flow uniformity (d).

Conditions	P_i /mbar	γ	M	T /K	$t_{\text{hydro}}/\mu\text{s}$	d /cm
A	29.1±2.6	1.6667	6.1±0.2	22.4 ±1.4	244	40
B	10.5±0.6	1.4762	5.0±0.1	42.5±1.3	414	32
C	9.3±0.7	1.4000	4.9±0.1	51.6±1.7	777	52
D	4.7±0.3	1.4000	4.2±0.1	64.2±1.7	202	11

Table 3. Rate coefficients for the gas-phase reaction of OH with methyl formate between 22 and 64 K and enhancement factors with respect to k_{OH} at 64 K.

Conditions	T/ K	$n/$ 10^{16}cm^{-3}	$[\text{HC(O)OCH}_3]/$ 10^{13}cm^{-3}	$k_{\text{OH}}(T)^{\text{a}}/ 10^{-11} \text{cm}^3 \text{s}^{-1}$	$k_{\text{OH}}(T)/ k_{\text{OH}}(64\text{K})$
A	22.4 ±1.4	17.2±1.6	0.20-3.00	11.9±0.36	10.3
B	42.5±1.3	5.22±0.33	0.62-9.20	4.39±0.52	3.8
C	51.6±1.7	4.17±0.35	0.77-9.20	2.51±0.30	2.2
D	64.2±1.7	2.24±0.15	0.54-9.70	1.16±0.12	1

^a Uncertainties are $\pm 2\sigma$, statistical and 10% due to systematic errors

Table 4. Literature values of the Arrhenius parameters, $k_{\text{OH}}(T) = A \exp(-E_a/RT)$, and experimental room temperature rate coefficient for the titled reaction.

T/ K	$k_{\text{OH}}(298\text{K})/$ $10^{-13} \text{ cm}^3 \text{ s}^{-1}$	$A/ 10^{-12} \text{ cm}^3 \text{ s}^{-1}$	$(E_a/ R)/ \text{K}$	Reference
296	2.27±0.34			Wallington <i>et al.</i> ²³
298	1.77±0.28			Good <i>et al.</i> ²⁴
298	1.83±0.33			Szilagyi <i>et al.</i> ²⁵
233 - 372	1.84	0.854±0.198	461±70	Le Calvé <i>et al.</i> ²⁶
880-1344	-	42.5	2026	Lam <i>et al.</i> ²⁷