

Journal of Materials Chemistry A

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this *Accepted Manuscript* with the edited and formatted *Advance Article* as soon as it is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.

1 **High Energy Density Titanium doped-Vanadium Oxide-Vertically**
2 **Aligned CNT Composite Electrodes for Supercapacitor Applications**

3 **Prashanth H. Jampani^a, Oleg Velikokhatnyi^b, Karan Kadakia^a, Dae Ho Hong^b, Sameer S. Damle^a,**
4 **James A. Poston^c, Ayyakkannu Manivannan^c and Prashant N. Kumta^{a,b,d,e,*}**

5 ^aDepartment of Chemical and Petroleum Engineering, University of Pittsburgh, Pittsburgh, PA
6 15261 (USA)

7 ^bDepartment of Bioengineering, University of Pittsburgh, Pittsburgh, PA 15261 (USA)

8 ^cUS Department of Energy, National Energy Technology Laboratory Morgantown, WV 26507,
9 USA

10 ^dDepartment of Mechanical Engineering and Materials Science, University of Pittsburgh,
11 Pittsburgh, PA 15261 (USA)

12 ^eCenter for Complex Engineered Multifunctional Materials, University of Pittsburgh,
13 Pittsburgh, PA 15261 (USA)

14 Keywords: Vanadium Oxide; Vertically Aligned Carbon Nanotubes (VACNTs);
15 Supercapacitor; Chemical Vapor Deposition; titanium doped vanadium oxide

16 **Corresponding Author**

17 *E-mail: pkumta@pitt.edu

18 **Abstract**

19 We provide herewith the first report of supercapacitance behavior of titanium doped vanadium
20 oxide films grown on vertically aligned carbon nanotubes using a chemical vapor deposition
21 technique (CVD). The capacitance of CVD derived titanium doped vanadium oxide-carbon

22 nanotube composites was measured at different scan rates to evaluate the charge storage
23 behavior. In addition, electrochemical characteristics of the titanium doped vanadium oxide thin
24 films synthesized by the CVD process were compared to substantiate the propitious effect of
25 carbon nanotubes on the doped vanadium oxide capacitance. Attractive capacitance values as
26 high as 310 F/g are reported here-in considering the overall materials loading with good rate
27 capability and excellent charge retention up to 400 cycles. Ab-initio theoretical studies
28 demonstrating the substantial improvement in the electronic conductivity of the vanadium oxide
29 due to titanium doping and oxygen vacancies have also been included corroborating the
30 attractive experimental capacitance response.

31 **1. Introduction**

32 The rapid technological progress of the world in the 19th and 20th century could largely be
33 attributed to the extensive use of electricity in the form of small devices such as light bulbs and
34 extending to very large applications such as those powering large mechanical units in industrial
35 plants. In an increasingly dynamic and electronic age however, a large number of mobile
36 applications still rely heavily on mechanical and thermal energy generation and storage which
37 are both inefficient and subject to generation of pollutants. There is therefore, a pressing need for
38 making electricity mobile and this makes it imperative for the development of charge storage
39 devices capable of delivering energy in a manner suiting for a variety of mobile applications
40 such as automobiles and electronic devices i.e. provide energy rapidly over an extended period of
41 time with minimal loss over time^{1, 2}. The need to tailor energy storage devices to suit such
42 applications is highlighted by the current targets established by the Department of Energy³ for
43 the same i.e. 7.5 Wh/kg at 625 W/kg lasting upto 15 years. An examination of the classic
44 Ragone plot will help identify devices capable of achieving the same⁴.

45 It is distinctly observed that only Lithium-ion batteries (LIBs) and electrochemical
46 capacitors have a reasonable propensity to achieve these targets, at least as far as mobile
47 applications are concerned. Electrochemical capacitors often referred to as ultracapacitors or
48 supercapacitors are electrochemical devices which exploit surface charge storage properties of
49 electrode materials, thus enabling rapid charge-discharge and rendering them very attractive for
50 high rate mobile applications such as automobiles and electronic devices. Supercapacitors are
51 typically categorized into two classes based on the nature of charge storage at the
52 electrochemical interface⁵. It is common knowledge in electrochemistry that an electrochemical
53 double layer is formed at an electrode-electrolyte interface as a result difference in Fermi energy-
54 levels of electrons in the electrode and the electrolyte. This difference in energy acts as a barrier
55 to electron tunneling from electrode to electrolyte or vice-versa until a certain potential is applied
56 to the electrode. The electrochemical interface thus acts as a nanometer sized 'double-layer' of
57 separated charges resembling a conventional dielectric capacitor⁶. Charge storage devices based
58 on this mechanism are aptly named electrochemical double layer capacitors (EDLCs).

59 A second class of capacitors is termed as pseudo-capacitors wherein charge transfer is
60 brought about by virtue of reversible Faradaic electrochemical redox reactions undergone by
61 specific electrochemically active elements within a compound or compounds. The nature of the
62 successive reversible electrochemical reactions of these compounds allows the charge storage
63 response of the device to resemble a capacitor⁷. While double layer capacitors are much more
64 reversible and rapid charge storage devices, pseudocapacitors have much higher energy densities
65 while at the same time maintaining moderately high rate capability on account of the simplistic
66 redox reactions occurring at the surface of the electrode. One could therefore envisage hybrid
67 composite materials capable of both very high rate capability and high energy density.

68 The synthesis of low cost, high energy, high rate hybrid composite materials requires us
69 to identify a high rate double layer material which can be interfaced with a high energy
70 pseudocapacitor material with a synergistic effect occurring as a result of the composite formed
71 by this interaction. Most commonly studied pseudocapacitor electrode materials to date are
72 hydrous ruthenium oxide and manganese oxide, the former limited by the high cost and the latter
73 restricted primarily due to the low energy density, respectively⁸⁻¹¹. There is therefore an urgent
74 need for identification of economically viable pseudo-capacitor materials exhibiting reliably high
75 energy densities. Transition metal oxides are very attractive for this kind of application on
76 account of their ability to accept variable oxidation states. In addition, they exhibit excellent
77 atmospheric stability even in nanoparticulate form making handling and storage on a large scale
78 a non-issue. Of particular interest to pseudocapacitor applications is the ability of the surface
79 oxide to undergo successive reversible electrochemical reactions as is the case in hydrous
80 Ruthenium oxide^{5, 6}. Hydrous Ruthenium oxide is an excellent pseudocapacitor material with
81 the ability to maintain stable capacitances of upto 900 F/g over a large number of cycles¹².
82 However, the prohibitive cost of noble metal oxides compels the search for alternative oxide
83 materials with similar charge storage characteristics. One such oxide resembling the multiple
84 oxidation states of the metal similar to Ru is vanadium oxide. Vanadium is an element belonging
85 to Group 5 of the periodic table with a half empty d-shell allowing the element to exhibit
86 multiple valence states from +2 to +5.

87 Vanadium oxide has previously been studied as a material for lithium ion battery
88 cathodes as well and correspondingly, different morphologies of the oxide including nanotubes
89 have been explored. Amorphous and hydrated vanadium oxide has also been studied widely as a
90 viable supercapacitor electrode in lithium and sodium ion electrolytes¹³⁻²¹. Amorphous

91 vanadium oxide gel/carbon composites have demonstrated capacities ~ 360 mAh/g with very
92 high intercalation rates. Oxide nanotubes on the other hand, have been studied extensively by
93 Nesper et al.^{22, 23} as electrode materials using lithium ion electrolytes and have shown great
94 promise. In addition, a class of so called Lithium-ion based aqueous supercapacitors have been
95 identified by Hu et . al. demonstrating capacitances of upto 740 F/g in aqueous LiCl²⁴⁻²⁷. In all
96 these studies, thin oxide films have shown to result in good electrochemical performance on
97 account of the short diffusion lengths for lithium intercalation. Doping of oxides (Mo, Mn, W,
98 Ti, Cu, and Ag) have been previously studied as an excellent pathway to improve electrical
99 conductivity and electrochemical charge storage properties²⁸⁻³⁵. Silver and copper doped oxides
100 have been shown to increase the electronic conductivities by up to 2 orders of magnitude^{36, 37}. Ti
101 and W doped vanadium oxide has also been shown to lead to improvement in pseudocapacitor
102 characteristics³⁸. Hence, we have chosen Ti doped vanadium oxide as the pseudocapacitor
103 material to be used in the composite hybrid electrode.

104 The identification of ideal double-layer type capacitor materials to couple with doped
105 vanadium oxide in order to form a hybrid electrode capable of high rates is relatively complex to
106 tackle on account of the plethora of carbon based double layer materials identified by the
107 capacitor community at large³⁹⁻⁴¹. The material we choose must be suitable to overcome the low
108 electronic conductivity usually observed in nanoparticulate oxides. We can examine them
109 individually to find suitable support materials for our application needing high rate capability i.e.
110 very high conductivity, capacitor behavior and excellent stability in the voltage of interest. It is
111 also essential to maintain good carbon material loading and good carbon-oxide interface while at
112 the same time avoiding thick oxide film formation on the carbon, making the underlying carbon
113 material substrate effects redundant. Activated carbons are excellent double layer supercapacitor

114 materials on account of their high surface area and large pore volume^{39, 41}. However, they suffer
115 from the problem of wettability, inability to sustain a uniform deposit on the pores of the
116 carbons. In addition, high surface area carbons usually are very porous resulting in poor
117 volumetric capacitance. Another excellent material that has been studied is graphene. Graphene,
118 as the name suggests is a single sheet of graphite which allows for ballistic 1-D transport of
119 electrons leading to electronic conductivities of up to 10^5 S/m in composites made with
120 graphene⁴²⁻⁴⁷. However, a major problem with using graphene is that the individual graphene
121 sheets are too thin to achieve sufficient loading of the oxide to achieve high energy densities in
122 practical devices.

123 On the other hand, an excessively thick deposit of oxide would render the graphene sheet
124 redundant as it would not be able to act as an electron channel to most of the oxide material. A
125 better solution then would be to use carbon nanotubes (CNTs) which have very high conductivity
126 on account of the 1-D electron transport. In addition, it is relatively easy to generate CNTs form
127 vertically aligned arrays which would ensure direct contact of each individual CNT with the
128 current collector as well as creating a 3-D architecture which would act as an ideal substrate for
129 oxide growth⁴⁸. Carbon nanotubes are of two types, SWNTs and MWNTs. SWNTs usually
130 demonstrate slightly better electrical properties than MWNTs⁴⁹. However, the synthesis of pure
131 vertically aligned SWNTs is very expensive and it renders the synthesis of the composite
132 electrodes accordingly prohibitive^{49, 50}. In addition, the electronic properties of single wall
133 carbon nanotubes are very much dependent on the nature of the graphene sheets within the
134 nanotube as reported⁵¹⁻⁵³. Alternatively, multi-wall carbon nanotubes are relatively inexpensive
135 to synthesize and can be synthesized on a variety of substrates.

136 Carbon nanotube based hetero-structures have demonstrated improved rate capabilities
137 for various Li-ion and supercapacitor electrode materials including silicon, ruthenium oxide,
138 manganese dioxide and nickel oxide⁵⁴⁻⁶¹. Electrodeposited amorphous hydrated vanadium oxide
139 coatings on CNTs have shown to exhibit capacitances as high as 910 F/g by Kim et al¹⁴ and
140 Balkus et al.⁶² in lithium electrolyte cells. In a recent report Sathiya et al. reported energy storage
141 behavior of oxide coated carbon nanotubes in a Swagelok type cell with a lithium electrolyte and
142 a lithium counter electrode exhibiting capacities as high as 850 mAh/g with up to 60% of the
143 energy storage coming from a capacitive contribution⁶³. Vanadium oxide/CNT composites have
144 also been explored for capacitance behavior in aqueous electrolytes and capacitances as high as
145 400 F/g have been reported⁶⁴. In all these studies, thin oxide films have shown to result in good
146 electrochemical performance on account of the short diffusion lengths for lithium intercalation.
147 The importance of loading and its influence in achieving attainable capacitance in vanadium
148 nitride based supercapacitor has been highlighted by us in our earlier work⁶⁵. A recent critical
149 review by Gogotsi et. al. has identified the need to understand the importance of thick films for
150 commercial supercapacitor electrodes and the need to report volumetric capacitance along with
151 gravimetric capacitance⁶⁶. Nanoparticulate materials such as oxides/nitrides of transition metals
152 are very advantageous in terms of achieving desirable electrochemical performance on account
153 of the high specific surface area (SSA) and the enhanced activity that is achievable at the nano-
154 scale dimensions. However, the very fact of having a nanoparticulate state which is responsible
155 for enhanced electrochemical activity results in loss of electronic conductivity on account of the
156 large number of grain boundaries. Nanostructured vanadium oxide has a relatively low electronic
157 conductivity^{36, 67, 68} on the order of 10^{-6} - 10^{-2} S/cm. The high electronic conductivity of carbon
158 nanotubes (CNTs) on the other hand, could thus be exploited to create composite 3-D

159 architectures to minimize the oxide–oxide particle contact leading to efficient electron transport
160 from the current collector to the surface of the vanadium oxide wherein the pseudocapacitance
161 behavior is exhibited.

162 In order to achieve this, architecture comprising vertically aligned carbon nanotubes
163 (VACNT) covered with electrochemically active doped-vanadium oxide would be ideal to obtain
164 maximum capacitance while allowing for slight expansion/contraction likely as a result of the
165 ensuing possible phase change. **Figure 1** shows the various architectures commonly used for
166 supercapacitor applications and their accompanying problems. Previous studies have shown the
167 effect of oxide-film thickness on the capacitance of vanadium oxide films prepared using atomic
168 layer deposition⁶⁹. Thicker films of the oxides however do not yield good charge storage
169 behavior due to the expected limitation in electronic conductivity. The aim in this study is
170 therefore to demonstrate a CNT-oxide composite material wherein the use of a moderately thick
171 doped oxide coating on a forest of vertically aligned electronically conducting carbon nanotubes
172 should allow for minimum electron transport limitations while maximizing active material use
173 and reducing inactive phase weight. This should be reflected accordingly as significant
174 improvement in areal capacity, a measure that is more relevant and of higher consequence than
175 gravimetric capacity.

176

177

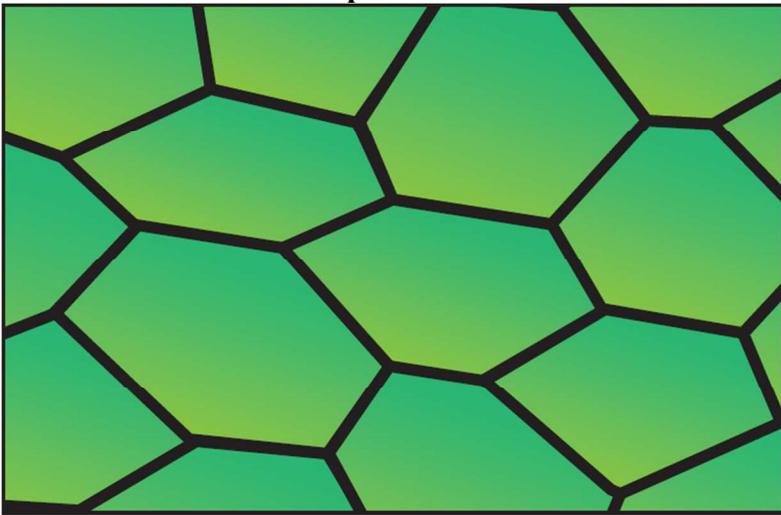
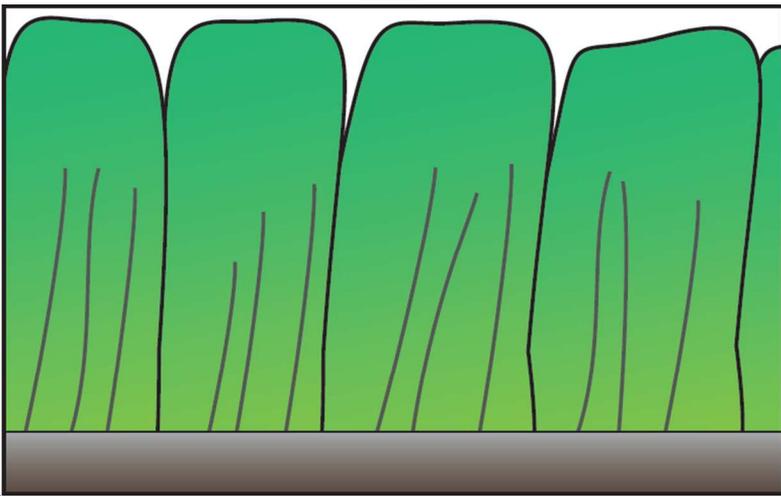
178

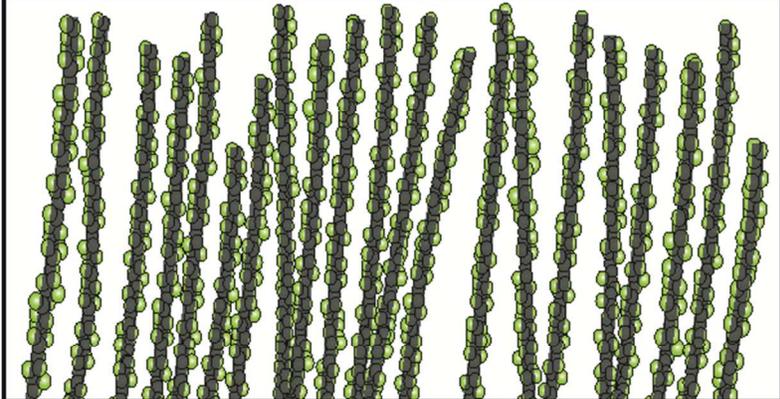
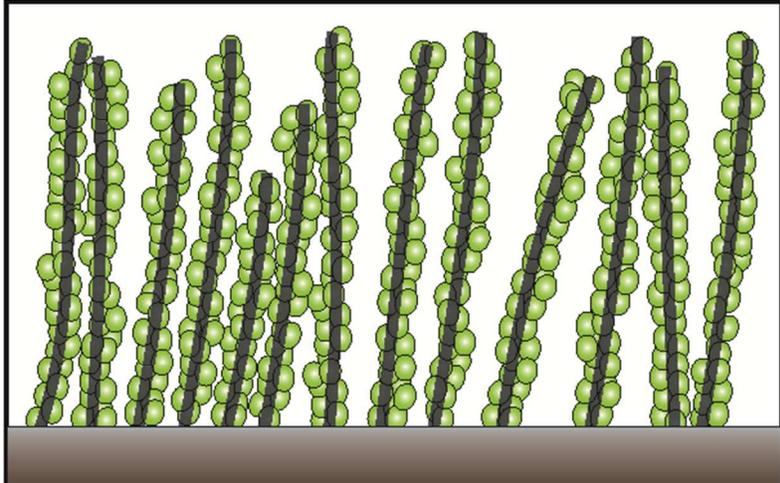
179

180

181

182

Morphology	Characteristics
<p data-bbox="467 531 711 562" style="text-align: center;">Nanoparticle film</p> 	<ul style="list-style-type: none"> <li data-bbox="1062 531 1338 562">(a) High surface area <li data-bbox="1062 569 1305 636">(b) Poor electronic conductivity <li data-bbox="1062 642 1425 709">(c) Large number of grain boundaries limiting flow <li data-bbox="1062 716 1386 852">(d) Inherently low conductivity of oxide materials hampered further
<p data-bbox="201 1081 977 1148" style="text-align: center;">Thick oxide matrix on carbon nanotubes/CNT-oxide mix film.</p> 	<ul style="list-style-type: none"> <li data-bbox="1062 1081 1414 1297">(a) Carbon nanotubes ineffective in reaching all parts of the oxide resulting in large grains and number of grain boundaries <li data-bbox="1062 1304 1414 1482">(b) Acts like thick film of oxide alone with little conduction of electrons to grains far from CNTs/current collector
<p data-bbox="201 1648 977 1715" style="text-align: center;">Thin oxide coating on vertically aligned carbon nanotube matrix.</p>	<ul style="list-style-type: none"> <li data-bbox="1062 1648 1382 1715">(a) Very good electronic contact <li data-bbox="1062 1722 1370 1858">(b) 3-D architecture ensuring ballistic electron transport to oxide film/grains <li data-bbox="1062 1864 1403 1898">(c) High oxide-electrolyte

	<p>contact maximizing pseudocapacitive contribution</p> <p>(d) Low oxide/CNT weight ratio resulting in low overall capacity when CNTs do not possess significant capacitance</p> <p>(e) Very short diffusion distances from oxide-CNT interface to oxide-electrolyte interface</p>
<p>Moderately thick doped oxide coating on vertically aligned carbon array.</p> 	<p>(a) Very good electronic contact between carbon nanotubes and oxide particles</p> <p>(b) 3-D architecture ensuring ballistic electron transport to oxide film/grains</p> <p>(c) High oxide-electrolyte contact maximizing pseudocapacitive contribution</p> <p>(d) Relatively short diffusion distances from oxide-CNT interface to oxide-electrolyte interface</p> <p>(e) Effects of electron mobility reduction due to grains overcome to a large extent by the increased conductivity of the doped oxide</p> <p>(f) Thicker film results in greater oxide/CNT ratio and thus high overall capacitance</p>

183 **Figure 1:** Choice of architecture for oxide-CNT composite materials to maximize areal
 184 capacitance.

185 Chemical methods can be utilized to synthesize carbon nanotube-doped oxide composite
186 hybrid electrode materials. However, wet chemical methods usually do not afford the freedom
187 for tailoring the morphology of the composite electrode as thickness of oxide film deposited on
188 the carbon nanotubes is usually very high as the oxide usually forms a matrix around the CNTs
189 rather than retaining nanoparticle nature. Synthesis of Ti doped oxide has usually involved the
190 sol-gel method which is non-ideal for obtaining thin films on carbon nanotubes⁷⁰. This leads to
191 increased oxide-oxide particle contact rather than carbon-nanotube-oxide electron transport. In
192 addition, carbon nanotubes usually need to be dispersed in an oxide sol, leading to loss in
193 alignment of the CNTs. The very purpose of using CNTs i.e. one dimensional transport of
194 electrons to nano-sized film of the oxide is lost. Chemical vapor deposition (CVD) is an ideal
195 and versatile method to obtain uniform architectures on carbon structures^{71, 72} while also
196 tailoring the nature of the deposited film depending on the deposition conditions. Various groups
197 have previously reported CVD as a technique for the growth of vanadium oxide thin films^{54, 73-79}
198 using both chloride and metal organic precursors. Oxide films with varying V oxidation states
199 have been reported with ratios of oxides depending on the various deposition conditions
200 including temperature and reagent concentrations⁷⁸.

201 Though some improvement in electronic conductivity and evidence of improvement of
202 charge storage by doping of oxides has previously been demonstrated, a detailed fundamental
203 study has not been undertaken correlating supercapacitor behavior with doping in both
204 nanoparticulate oxide and nitride materials. On the basis of the existing evidence in literature, it
205 is to be expected that a first principles computational screening study to identify suitable dopants
206 could be used to predict and engineer high capacity, high rate, and stable supercapacitor
207 materials. The aim of this work is thus to use doping of oxides as an approach that allows for

208 modifying the electronic conductivity and thus the electrochemical activity, as well as the
209 electrochemical stability of nanoparticulate oxides and nitrides. Using a theoretical simulation
210 involving d-band assessment of the structure, suitable dopants have been identified to improve
211 the electronic conductivity and thus the supercapacitor behavior of nanostructured oxides. Using
212 such approaches, the current study provides a unique methodology for generating thick film
213 architectures of doped vanadium oxides grown on vertically aligned CNTs (VACNTs) resulting
214 in capacitance values that are uniquely distinct from currently reported values in the literature.
215 Accordingly, in this work, an atmospheric pressure chemical vapor deposition (APCVD) method
216 has been explored to deposit titanium doped-vanadium oxide on vertically aligned carbon
217 nanotubes for use as a supercapacitor electrode in aqueous media. Implementation of CVD as a
218 synthesis method also provides the advantage of avoiding the incorporation of binders for
219 generating the electrodes which can consume upto 50% of the total weight of the electrode in
220 high surface area systems contributing largely to inactive dead weight^{55, 72}. As a result of the
221 improved electronic conductivity of doped vanadium oxide and high conductivity of vertically
222 aligned carbon nanotubes, we are able to demonstrate high gravimetric capacitances in thick
223 electrodes resulting in high areal capacities.

224 **2. Experimental Methods**

225 **2.1 Materials Synthesis**

226 **2.1.1 Chemical Vapor deposition**

227 A two-step CVD process was used to obtain doped-vanadium oxide nanosphere coated
228 CNT hetero-structures on nickel substrate disks. First, the nickel disks were properly cleaned
229 using dilute acids and acetone, and then weighed. Vertically aligned carbon nanotubes

230 (VACNTs) were then deposited on the Ni disks (Alfa Aesar Inc., 0.05 mm thick, annealed, 99+%

231 (metals basis)) using a xylene-ferrocene mixture following a hot wall chemical vapor deposition

232 (CVD) procedure similar to that reported by Ajayan et al. using an Ar/H₂ mixture^{80, 81}. m-Xylene

233 (ACROS organics, extra pure, >99%) was used as the carbon source with ferrocene (ACROS

234 Organics, 98%) acting as the catalyst for CNT nucleation and growth (**Figure 2a**). The mixture

235 was vaporized and led into a hot-wall CVD tube maintained at 770° C. Deposition was carried

236 out for 30 minutes resulting in the growth of a uniform forest of vertically aligned carbon

237 nanotubes. The VACNT coated Nickel disks were then weighed and a low temperature hot-wall

238 APCVD setup was used to deposit the doped vanadium oxide directly onto the vertically aligned

239 carbon nanotubes. To generate the doped vanadium oxide, appropriate mixtures of vanadium

240 tetrachloride (VCl₄, Sigma Aldrich) and metal dopant introduced as a chloride precursor, i.e.

241 titanium tetrachloride (TiCl₄, Sigma Aldrich, 99.9% metals basis) were dissolved in an aprotic

242 solvent, chloroform (CHCl₃, ACROS Organics, 99.9% Extra dry) in an inert atmosphere glove

243 box (Vacuum Atmospheres Inc.: model EE-493, below 5ppm O₂ + H₂O) and recovered in a

244 sealed stainless steel evaporator bottle.

245 The doped vanadium oxide nanospheres were deposited on the VACNTs using a hot wall

246 tubular APCVD setup (**Figure 2b**) consisting of two separate nitrogen (UHP N₂, Grade 5, Valley

247 National gas) carrier gas lines running through the evaporator vessels, one containing the mixture

248 of chloride precursors and another containing a distilled water (Millipore QGARD,

249 Resistivity=18.2 MΩ-cm) line driven by a syringe pump (Thermo Scientific Orion M365 Sage

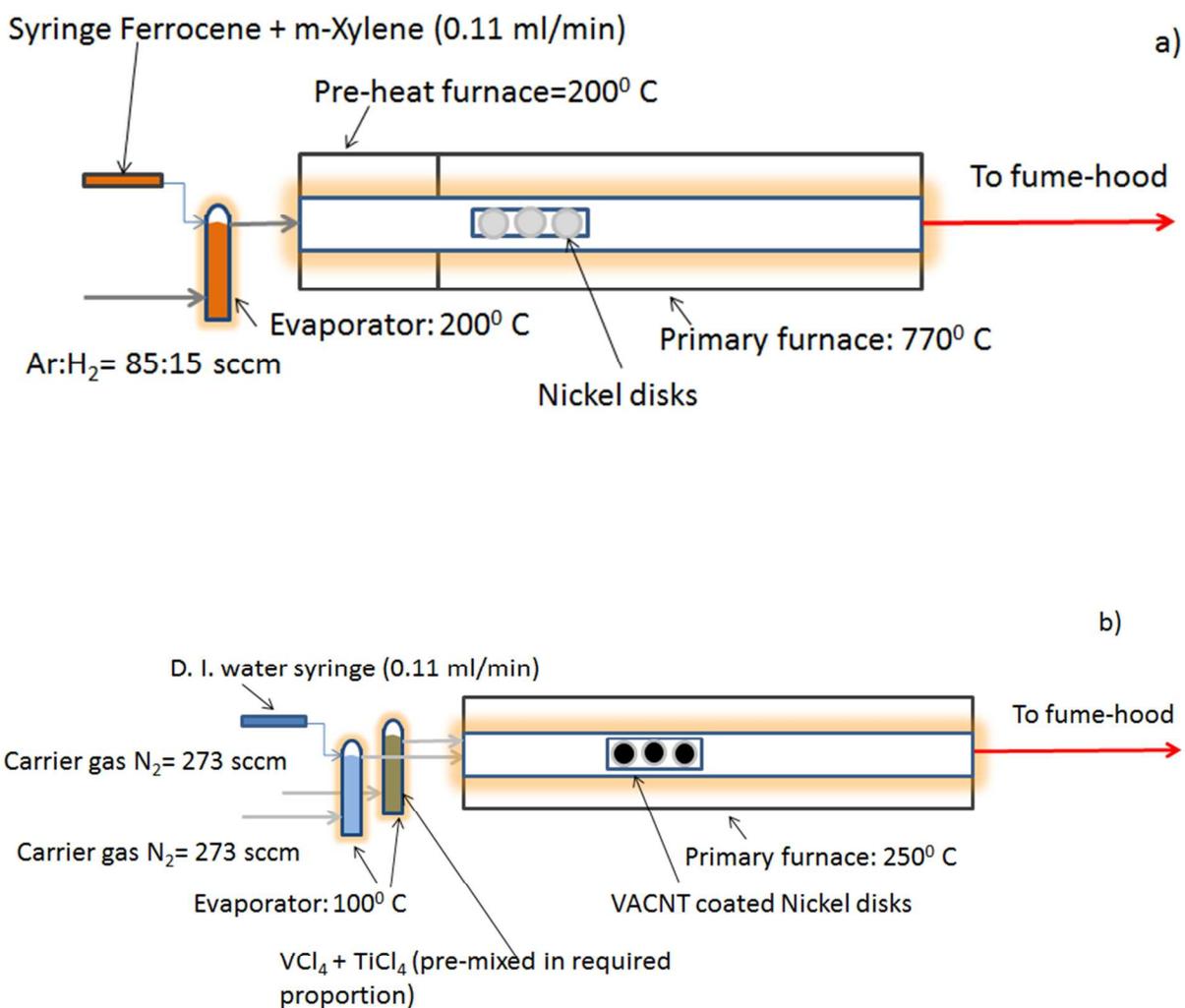
250 syringe pump). Both precursors were evaporated and carried into a hot wall CVD reactor tube

251 wherein the CNT coated nickel (Ni) disks were placed^{17, 38, 78, 82-84}. Deposition was carried out at

252 250°C for 20 minutes. The oxide coated disks were then weighed and characterized. Doped

253 Vanadium oxide thin film was also generated by depositing the doped oxide directly on the Ni
 254 substrate disks by CVD using the same precursors to compare the electrical and electrochemical
 255 properties of the thin films of the doped oxide with that of the composite heterostructure
 256 comprised of doped oxide coated on VACNTs.

257



261 **Figure 2:** Chemical vapor deposition setup for deposition of (a) VACNT forest (b) VO_x:Ti-
 262 VACNT composite.

263 2.2 Materials Characterization

264 The nature of the deposit including morphology of the oxide on the CNTs needs to be
265 analyzed and examined to obtain a better understanding of the observed capacitance behavior. To
266 accordingly investigate and compare the microstructure of the CVD deposited doped vanadium
267 oxide-CNT composites and thin films, scanning electron microscopy (SEM) analysis was
268 conducted using a Philips XL30 machine operating at 20 kV. The atomic composition of the
269 films was also verified using EDAX analysis. To study the nature of the doped oxides coated on
270 the CNTs and to identify presence of any amorphous deposits other than the oxide present on the
271 CNTs, micro-Raman spectroscopy was performed on a Reinshaw inVia Raman microscope.
272 Raman spectra were collected in air at room temperature using a 633 nm laser. Low laser
273 intensities were chosen to avoid beam modification of the samples and accordingly, the spectra
274 were collected by averaging multiple scans over a frequency range of 3500-100 cm^{-1} .

275 To further characterize the doped oxide-CNT interface and morphology of the doped
276 oxides, high-resolution transmission electron microscopy (HRTEM) was performed on a JEOL
277 JEM-210F microscope using beam energy of 200 keV. Samples for HRTEM were prepared from
278 the CVD synthesized composites by scraping off the doped vanadium oxide deposited VACNTs
279 from the Ni substrate, followed by suspending in methanol and dispersing them on to 3.05 mm
280 diameter Nickel grids (JEOL, 400 mesh nickel) containing a holey carbon film by sonication.
281 The grids were then vacuum-dried for 24 h and subjected to conventional and HRTEM analyses.

282 X-ray photoelectron spectroscopy (XPS) was performed on a Perkin-Elmer PHI-5600
283 ESCA system to analyze the surface composition and determine the valence state of V ions in the
284 oxide-VACNT composite. XPS was conducted using a Physical Electronics (PHI) model 32-096

285 x-ray source control and a 22-040 power supply interfaced to a model 04-548 x-ray source with
286 an Omni Focus III spherical capacitance analyzer (SCA). The system is routinely operated within
287 the pressure range of 10^{-8} to 10^{-9} Torr (1.3×10^{-6} to 1.3×10^{-7} Pa). The system was calibrated in
288 accordance with the manufacturer's procedures utilizing the photoemission lines, E_b of $\text{Cu}_{2p}^{3/2}$ -
289 932.7 eV, E_b of $\text{Au}_{4f}^{7/2}$ - 84 eV and E_b of $\text{Ag}_{3d}^{5/2}$ - 368.3 for a magnesium anode. All the reported
290 intensities are in effect experimentally determined peak areas divided by the instrumental
291 sensitivity factors. Charge correction was obtained by referencing the adventitious C_{1s} peak to
292 284.8 eV.

293 The electronic conductivity is a crucial parameter affecting the supercapacitor
294 performance and hence, four point probe conductivity set up was used to measure the
295 conductivity of the vanadium oxide films prepared by CVD. A Veeco FPP 5000 four-probe
296 conductivity instrument (Miller Design & Equipment, Inc.) was used for the conductivity
297 measurements and the sheet resistance values obtained were then converted to electrical
298 conductivity. Four-probe conductivity data is usually obtained by preparing pellets from powders
299 and supporting them on insulators while maintaining contact with four equally spaced probes,
300 two of which supply current while the other two measure voltage across them following standard
301 known procedures. The measurement of resistivity of thin films deposited on conducting metal
302 substrates however is a difficult proposition as the flow of electrons from the two current probes
303 would probably occur through the substrate current collector itself offering the least resistance
304 path. This current flow does not however occur in the CVD derived doped oxide and the
305 VACNT- doped oxide composite as the deposit is sufficiently thick and crack free to avoid
306 interference from the substrate.

307 **2.3 Electrochemical Testing**

308 Cyclic voltammetry (CV) was performed on the synthesized doped vanadium oxide thin
309 film and doped vanadium oxide-VACNT heterostructure materials to test their viability as high
310 energy density supercapacitor materials. A three electrode cell configuration was used to perform
311 cyclic voltammetry measurements and electrochemical impedance spectroscopy analysis (EIS)
312 on the coated Ni disks to both understand the electrochemical mechanisms involved as well as
313 evaluate the specific capacitance of both, the VACNT- doped vanadium oxide nanosphere
314 composite, and the doped vanadium oxide thin films prepared by chemical vapor deposition. The
315 cell was assembled using platinum as the counter electrode, Ag/AgCl (in 3M KCl) as the
316 reference electrode, and the active material coated Ni disks described above as the working
317 electrode. An aqueous 1M Na₂SO₄ solution was used as the electrolyte.

318 Cyclic voltammetry tests were performed on a Princeton Applied Research (PAR)
319 Versastat-3 potentiostat in the -1 V to 0.8 V potential ranges with respect to the reference
320 electrode to evaluate the specific capacitance of the doped vanadium oxide-VACNT composite
321 heterostructures generated by CVD. It was observed that no significant water electrolysis
322 occurred in this extended voltage window. Cycling was therefore performed at different scan
323 rates and area of the cyclic voltammogram was used to calculate the total capacitance.
324 Gravimetric capacitance values were evaluated on the basis of both the oxide weight alone as
325 well as the weight of the total active material.

326 The capacitance was calculated using the formula:

$$327 \quad \text{Capacitance } C \text{ (in F/g)} = \frac{\int I \, dV}{s \cdot 2 \cdot V \cdot m} \quad \text{i);}$$

328 where

329 $\int I dV$ = area generated by cyclic voltammogram

330 s = scan rate in V/s,

331 V = voltage window = 1.8 V,

332 and m = mass of the active material in gram.

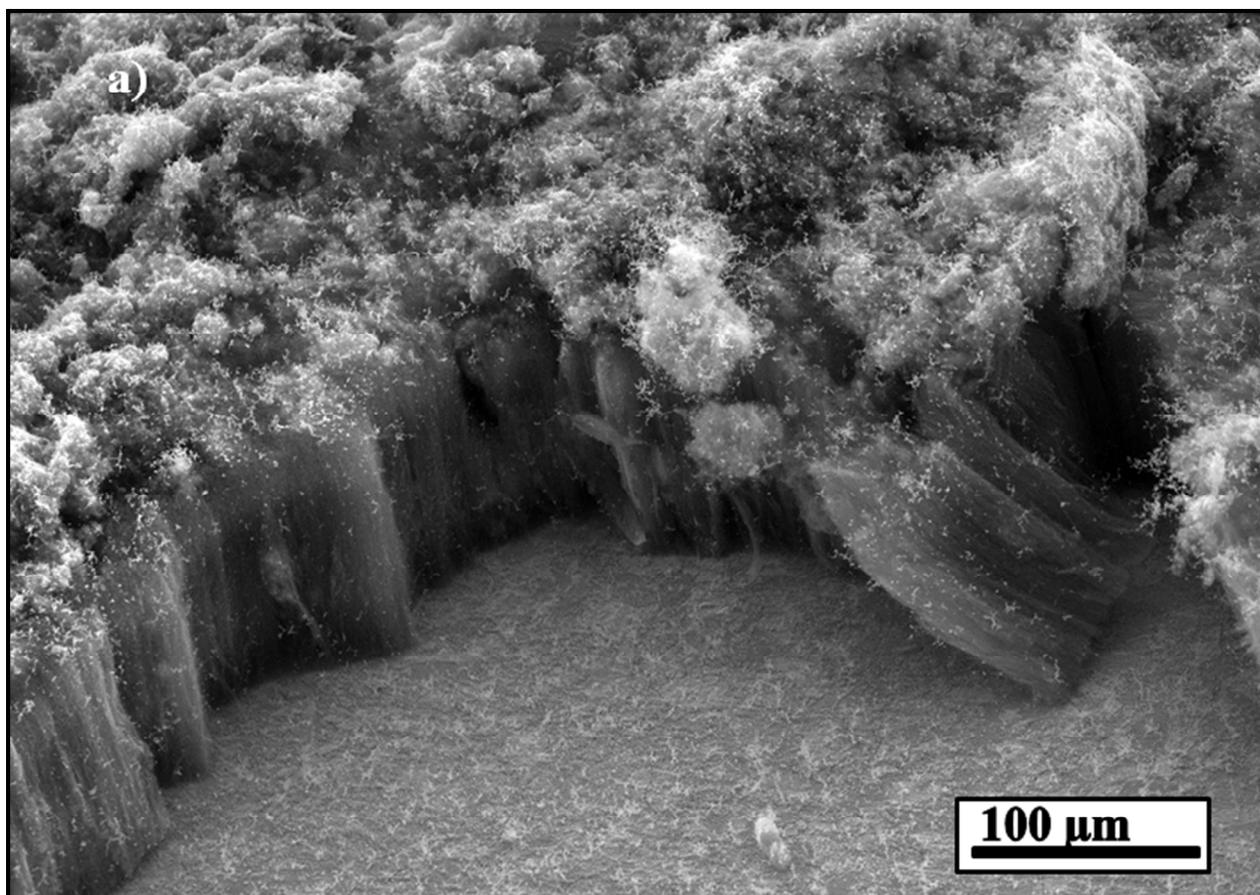
333 In addition to cyclic voltammetry, galvanostatic charge-discharge cycling was performed
334 on an Arbin BT-2000 cycler at different current densities in the voltage window of -0.5 to 1 V
335 w.r.t Ag/AgCl (3M KCl) to observe the nature of the voltage-time (V-t) characteristic and to
336 obtain capacitance values under galvanostatic conditions which resembles more closely
337 commercial applications of supercapacitors. The capacitance accordingly was calculated by
338 integrating the current-time data over the voltage window.

339 To further evaluate the supercapacitor response of the CVD-deposited doped vanadium
340 oxide thin films, doped vanadium oxide-VACNT heterostructures, electrochemical impedance
341 spectroscopy was performed on the Applied Research (PAR) Versastat 3 over a frequency range
342 of 0.1Hz-100 KHz. A 5 mV amplitude was used and the spectra were collected between -1.0 and
343 0.8 V at the specific voltages where Faradaic reactions were observed to be predominant.
344 Impedance spectroscopy is a versatile technique which can be used to quantify charge transfer
345 parameters in electrochemical systems. The impedance spectra were obtained on the titanium
346 doped vanadium oxide thin film and the VACNT-titanium doped vanadium oxide nanosphere
347 composite heterostructure hybrid electrodes to observe any noticeable changes in charge storage
348 mechanism induced by the presence of carbon nanotubes. The ZView-2 (Scribner Associates,
349 Inc.) software was used for theoretical modeling and fitting using the Randall's equivalent circuit
350 models, and the charge transfer parameters corresponding to the CVD synthesized doped

351 vanadium oxide-VACNT nanostructures were obtained and correspondingly compared to the
352 doped oxide thin film to understand the effect of VACNTs on the performance of the Ti doped
353 vanadium oxide.

354 **3. Results & Discussion**

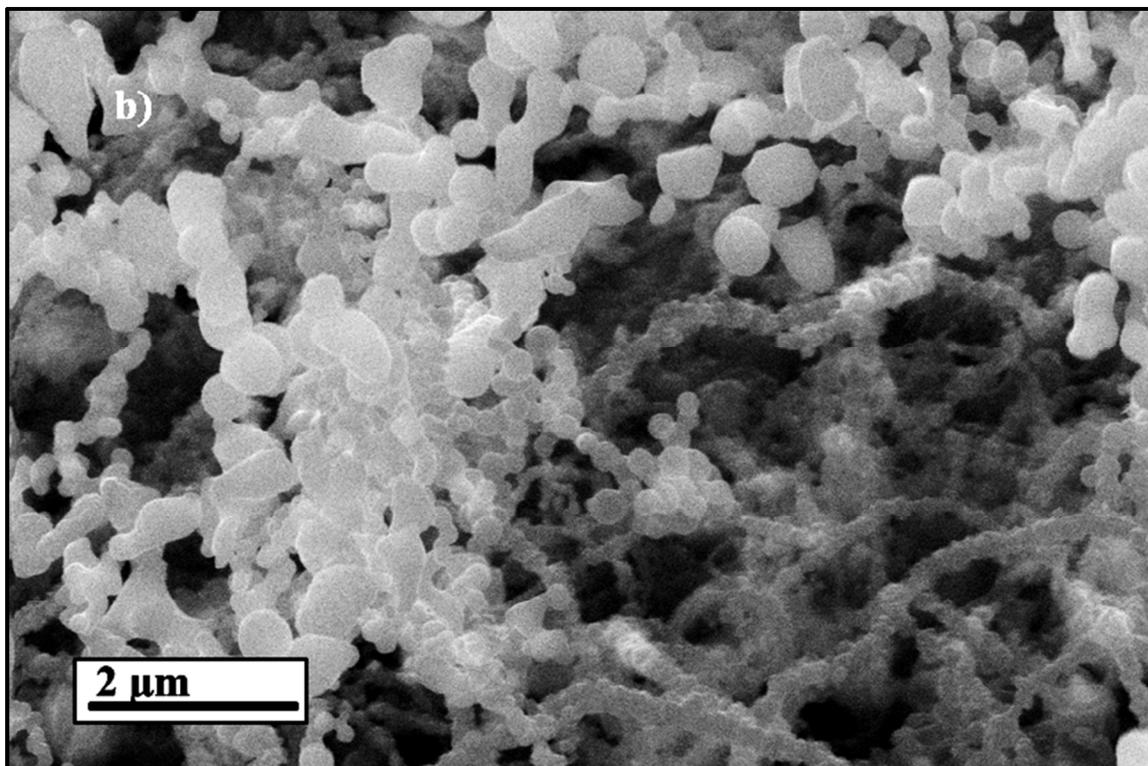
355 **Figure 3a** shows the scanning electron microscope (SEM) image of vertically aligned
356 CNTs (VACNTs) deposited by chemical vapor deposition on Ni substrate using a deposition
357 time of 30 minutes for growth of carbon nanotubes. The carbon nanotubes grow on the nickel
358 substrate supported by iron nanoparticles (from Ferrocene) which nucleate on the surface of the
359 nickel as islands and in turn acting as growth sites for carbon^{85, 86}. The carbon then grows in the
360 form of nanotubes by progressive rise of the nucleated iron islands which continue to act as
361 nucleation sites (floating catalyst mechanism) of carbon. Hence the length of the tubes can be
362 controlled by varying the CVD deposition time. Accordingly, optimized deposition time of 30
363 minutes was used to obtain the vertically aligned tubes 120 micron in length. The hypothesis is
364 that the prolific length of the CNTs would help create a very good 3-D architecture which when
365 covered with a thin layer of the oxide would maximize the volumetric loading while at the same
366 time ensuring that a very thick film does not cover the entire surface of the carbon nanotubes
367 impeding electronic transport within the oxide itself.



368
369

370 **Figure 3b** shows the SEM image of titanium doped vanadium oxide ($\text{VO}_x\text{:Ti}$) coating on the
371 vertically aligned CNTs deposited by chemical vapor deposition using the pre-mixed chloride
372 precursors and water which were aspirated using the CVD system described in **Figure 2b**.
373 Deposition of the doped oxide using the pre-mixed chloride precursors and water aspirated in the
374 CVD system was carried out for 20 minutes ensuring that the oxide is coated on the VACNTs in
375 the form of a contiguous film consisting of a number of nanometer sized globules. **Figure 4**
376 shows the bright field high resolution transmission electron microscopy (HRTEM) image of
377 $\text{VO}_x\text{:Ti}$ -VACNT hybrid composite scrapped off the nickel substrate. The image confirms that the
378 $\text{VO}_x\text{:Ti}$ is deposited over the VACNTs in the form of a semi-contiguous film consisting of ~250

379 nm diameter amorphous globules evident by the lack of visible lattice fringes in the oxide
380 particle seen in **Figure 4**.

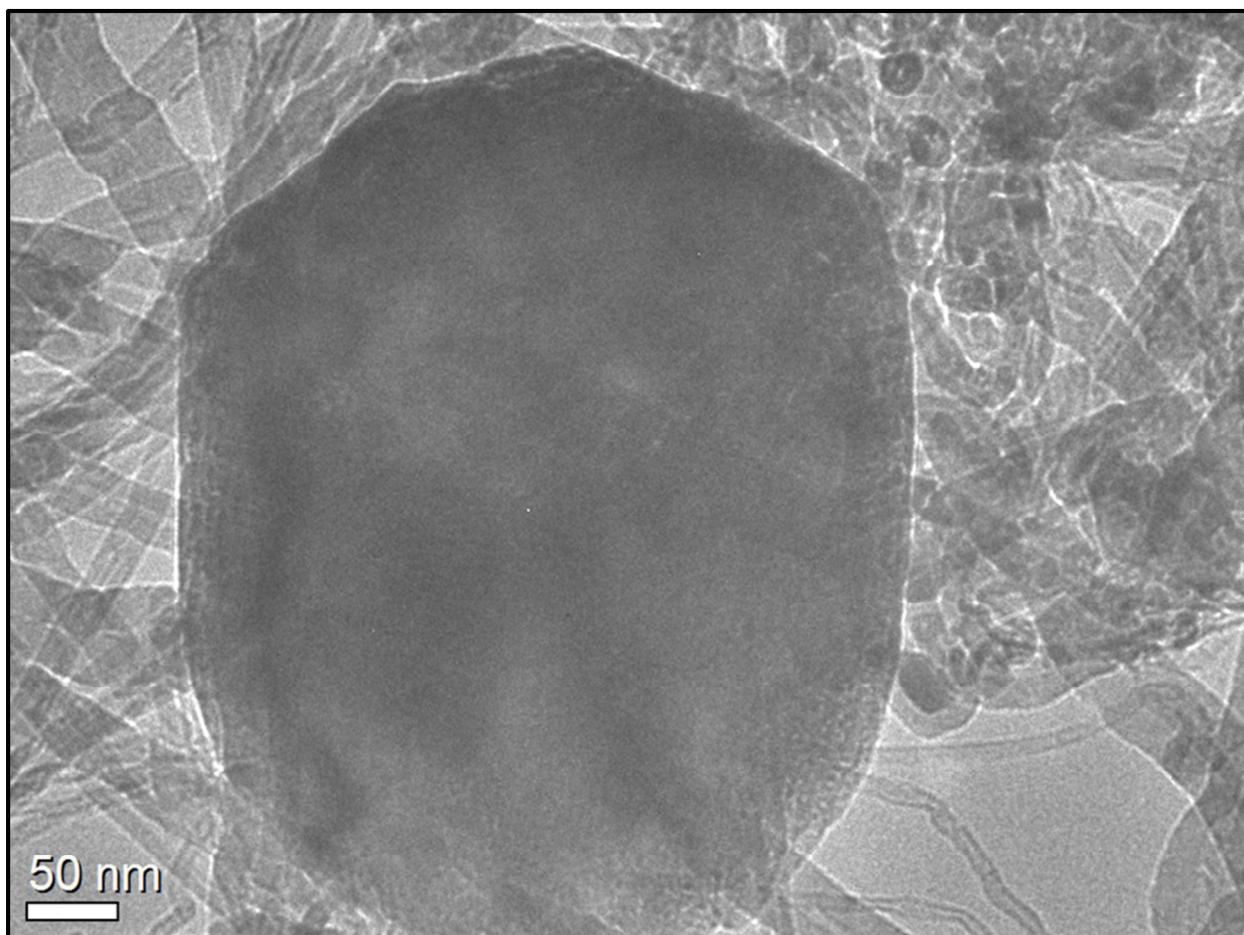


381

382

383 **Figure 3:** Scanning Electron Microscope image of (a) CVD derived Vertically Aligned Carbon
384 nanotubes coated with titanium doped vanadium oxide showing forest of vertically aligned
385 carbon nanotubes, and (b) Nanoparticulate titanium-doped vanadium oxide coated on the carbon
386 nanotubes grown by chemical vapor deposition.

387

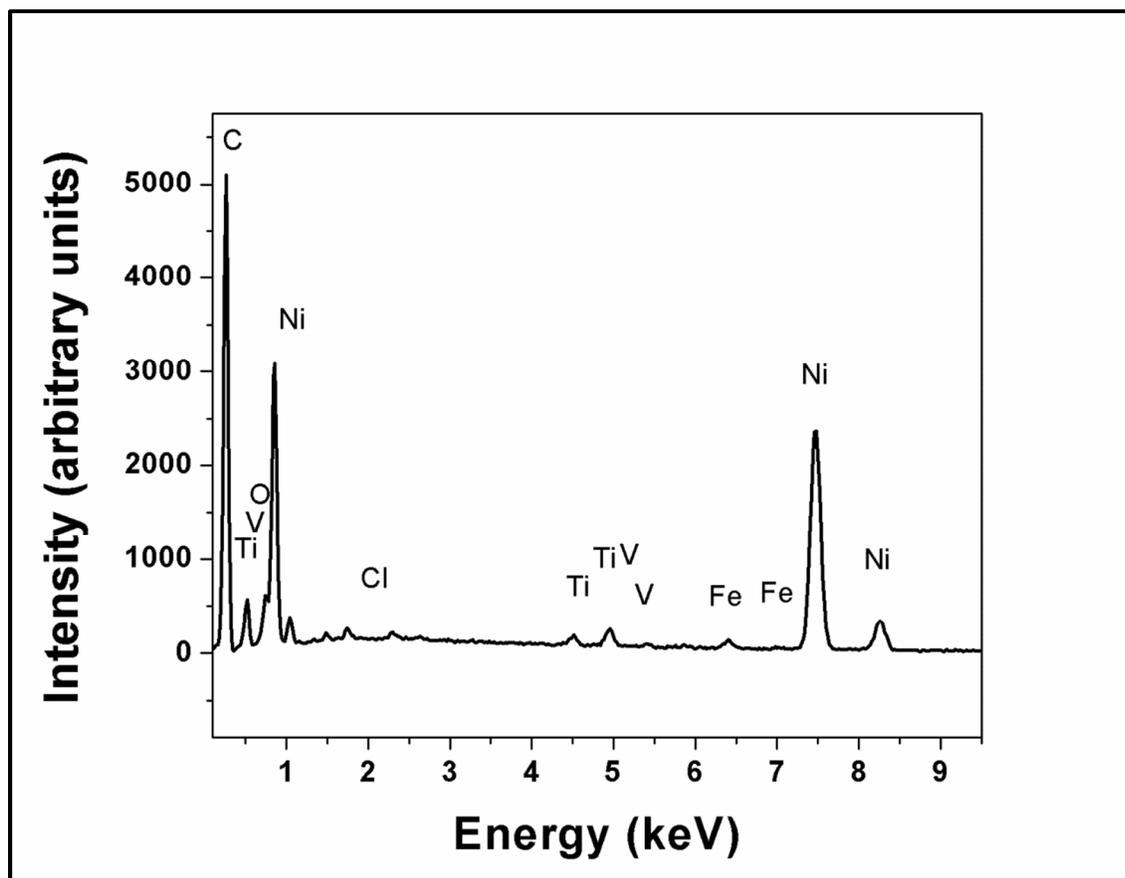


388
389

390 **Figure 4:** HRTEM images of the $\text{VO}_x\text{:Ti}$ coated on the vertically aligned carbon nanotubes
391 (VACNTs) showing nanoparticulate size of the oxide particles grown on the carbon nanotubes.

392 Energy dispersive x-ray analysis (EDX) was performed on the sample as shown in
393 **Figure 5** since the composition of the CVD derived $\text{VO}_x\text{:Ti-VACNT}$ could not be accurately
394 confirmed by x-ray diffraction due to the amorphous nature of the oxide deposited on the
395 VACNTs (**Figure 4**). The presence of elemental V, O and Ti in the $\text{VO}_x\text{:Ti-VACNT}$
396 heterostructures has been confirmed using the EDX configured to the SEM. The V and Ti ratio
397 in the $\text{VO}_x\text{:Ti-VACNT}$ is similar to the nominal composition used and is confirmed by EDX.
398 Nickel is also observed in the EDX which is due to the substrate on which $\text{VO}_x\text{:Ti-VACNT}$ is

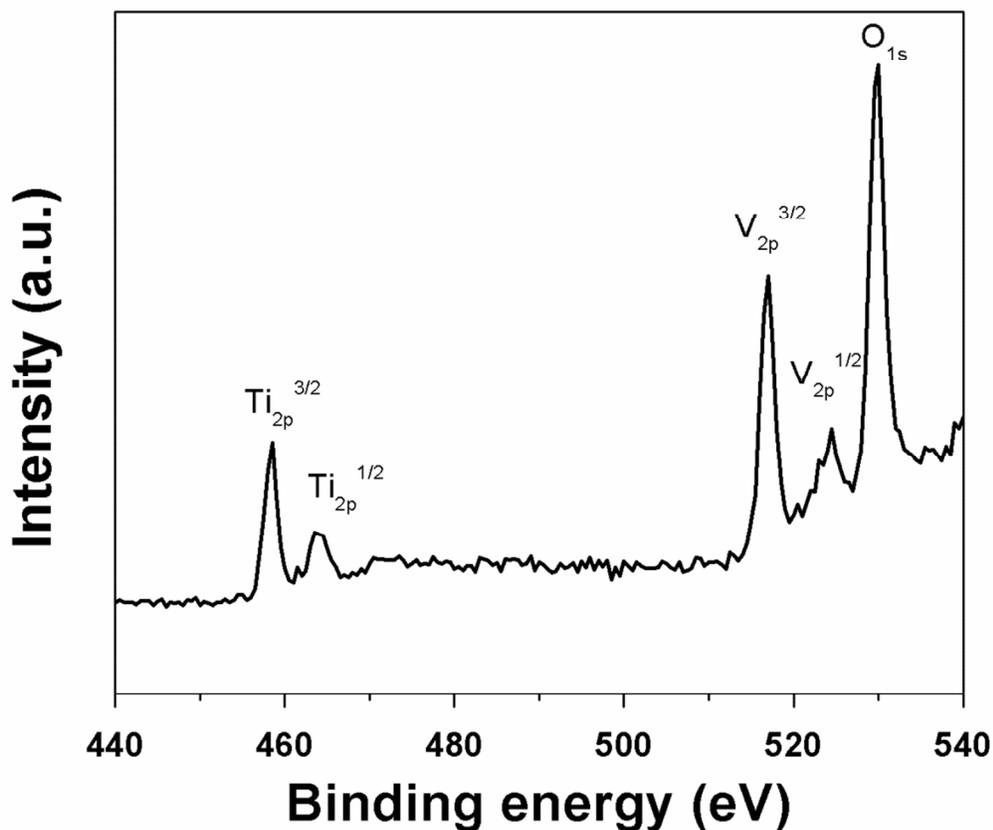
399 synthesized. It was found that the CVD derived doped oxide-VACNT composite heterostructure
 400 had a nominal composition of $V_{1.6}Ti_{0.4}O_{5-x}$ with $x \sim 0.4$. The vacancy rich oxide thus formed
 401 would correspond to a high oxidation state of vanadium which has previously been identified to
 402 be essential for achieving high capacity in vanadium oxide^{24, 87}.



403
 404 **Figure 5:** EDX spectra of the $VO_x:Ti$ -VACNT composite heterostructures synthesized by CVD.

405 In order to determine the chemical oxidation states of V, Ti and O in the VACNT-doped oxide
 406 composite electrode, x-ray photoelectron spectroscopy (XPS) was conducted using a Phillips
 407 PHI 5600 system. The XPS spectra of V_{2p} and O_{1s} peak of vanadium oxide obtained are shown
 408 in **Figure 6**. As shown in the figure, the binding energy of vanadium [$V_{2p}^{3/2}$ (516.95 eV) and
 409 $V_{2p}^{1/2}$ (524.49 eV)] is indicative of V in the oxidation state between +4 and +5^{88, 89}. Further,

410 reported $V_{2p}^{3/2}$ values of vanadium oxide for the +4 oxidation state i.e. VO_2 vary between 515.7
411 and 516.5 eV, and for the +5 oxidation state i.e. V_2O_5 vary between 514.9 and 518.3 eV.
412 Therefore an oxidation state of +4/+5 has been attributed to the vanadium in this study⁹⁰. The
413 binding energy of titanium [$Ti_{2p}^{3/2}$ (458.51 eV) and $Ti_{2p}^{1/2}$ (463.83 eV)] in the CVD derived
414 VACNT-doped oxide corresponds very closely to Ti existing in the stable +4 oxidation state⁷⁶.
415 The binding energy of oxygen O_{1s} (529.96 eV) corresponds to a high oxidation state of the metal,
416 namely in the +4/+5 oxidation state. Both the vanadium and titanium are thus in their highest
417 oxidation states which is conducive for achieving the observed and improved supercapacitor
418 response reported herein. Further analysis including deconvolution of XPS peaks has not been
419 undertaken herein since XPS results have been used only as an indicator of the presence of a
420 mixed oxide of vanadium and titanium in the $VO_x:Ti$ -VACNT composite heterostructure. The
421 Ti doped vanadium oxide coated on the VACNTs synthesized by the CVD approach exhibits V
422 displaying a mixture of high oxidation state of V^{5+}/V^{4+} similar to that of other promising oxide
423 systems used in supercapacitors^{24, 87}. The XPS spectrum and binding energies thus match very
424 closely that of Ti doped vanadium oxide reported earlier⁹¹.



425
 426 **Figure 6:** X-ray photoelectron spectrum of CVD deposited VACNT-titanium doped vanadium
 427 oxide composite heterostructure showing the peaks corresponding to the binding energies of
 428 vanadium, titanium and oxygen.

429 Raman spectra was collected on the VO_x:Ti-VACNT which is shown in **Figure 7a** since the x-
 430 ray diffraction pattern of the titanium doped vanadium oxide could not be acquired on account of
 431 the amorphous nature of the oxide grown on the VACNTs. The D, G and G' bands are observed
 432 at 1250, 1650 and 2650 cm⁻¹. The G/D ($I_g/I_d=1.16$) ratio is greater than one indicating few lattice
 433 defects in the graphene layers resulting in MWCNTs with relatively high conductivity^{92, 93}. This
 434 would also have a positive effect on the capacitance behavior of the Ti doped vanadium oxide.
 435 **Figure 7a** shows several peaks marked that have been identified to correspond and match closest

436 to those observed and reported in thin films of V_2O_5 ⁹⁴ and those of vanadium oxide
437 corresponding to VO_x -carbon nanofibers reported previously⁹⁵. **Table 1** also shows the tabulated
438 peak positions observed in **Figure 7a** and indicates the nature of the vibration resulting in the
439 peaks seen in the Raman pattern of VO_x :Ti-VACNT. The V bonds which are characteristic of
440 double and multilayer structures usually observed in V_2O_5 appeared at 404, 528 and 996 cm^{-1} ⁹⁶,
441 ⁹⁷. The peak at 700 cm^{-1} is indicative of stretching vibration of oxygen ions in bridging position
442 between the three vanadia centers⁹⁷. The peak observed at 477 cm^{-1} is at a similar position to that
443 observed in V_2O_5 where oxygen ions are situated between the two vanadium centers⁹⁷. Titanium
444 oxide conventionally has Raman peaks corresponding to the anatase phase at 142 cm^{-1} (E_g), 395
445 cm^{-1} (B_{1g}), 515 cm^{-1} ($B_{1g}+A_{1g}$) and 636 cm^{-1} (A_{1g}). Brookite (TiO_2) on the other hand, usually has
446 a strong band at 153 cm^{-1} (E_g) and another at 320 cm^{-1} (B_{1g})^{98, 99}. No Raman lines due to
447 titanium oxide can be observed in the VO_x :Ti-VACNT samples, which proves that the titanium
448 is dispersed very well in a homogeneous manner with no observed phase separation and
449 formation of any random individual titanium oxide clusters. It also confirms that titanium may be
450 present in the substitutional positions in the lattice of the parent vanadium oxide. As seen in
451 **Table 1**, every single peak is seen to match with that of the linkages seen in the vanadia lattice.
452 However, the peak occurring at 830 cm^{-1} does not match any of the vibrations of vanadium and
453 oxygen. Nevertheless, it matches closely with that of V-O-Ti as reported by Krykova et al.¹⁰⁰
454 (Ti-O-V linkage at a wavenumber of 860 cm^{-1} has been reported). The wavenumber of 830 cm^{-1}
455 also matches that of rutile TiO_2 seen in TiO_2 supported V_2O_5 ¹⁰¹. However, none of the other
456 vibrations of the rutile phase are present leading us to conclude that it is indeed a Ti-O-V bond
457 rather than a secondary rutile phase. Titanium and vanadium have similar ionic radii which
458 allows for the relative ease of doping and substitution of titanium on V sites in the vanadium

oxide lattice. **Figure 7b** shows the comparison of Raman spectra seen in $\text{VO}_x\text{:Ti-VACNT}$ and $\text{VO}_x\text{-VACNT}$. It can distinctly be observed that the peak at 830 cm^{-1} is absent in the $\text{VO}_x\text{-VACNT}$ further corroborating the hypothesis that the peak position occurs as a result of a Ti-O-V vibration.

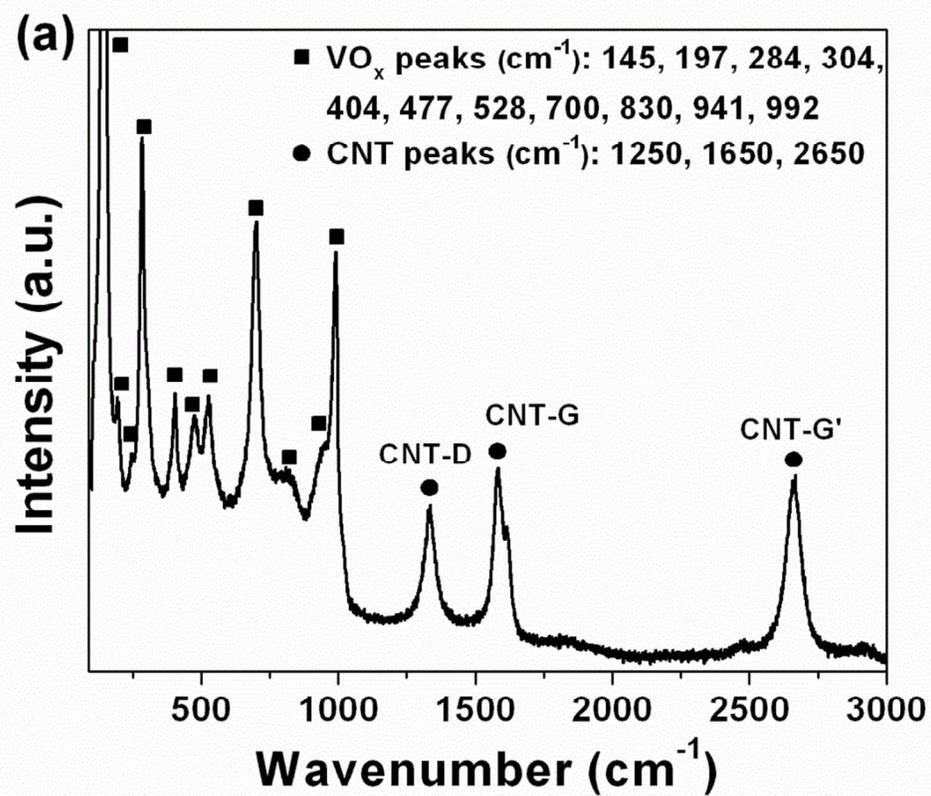
Table 1: Peak positions observed in Raman spectrum of $\text{VO}_x\text{:Ti-VACNT}$ and corresponding vibrations.

Wavenumber (cm^{-1})	Vibration	References
145	$(\text{V}_2\text{O}_2)_n$, bending	85
197	$(\text{V}_2\text{O}_2)_n$, bending	85
284	V=O, bending	85
304	$\text{V}_3\text{-O}$, bending	95, 101
404	V=O, bending	95, 101
477	V-O-V, bending	95
528	$\text{V}_3\text{-O}$, stretching	95
700	$\text{V}_2\text{-O}$, stretching	95
830	V-O-Ti	96, 100
941	Bridging V-O-V vibrational bands	100-102
992	Symmetric stretching of V=O bond	100, 101
1250	CNT-D-band	78, 79
1650	CNT-G-band	78, 79

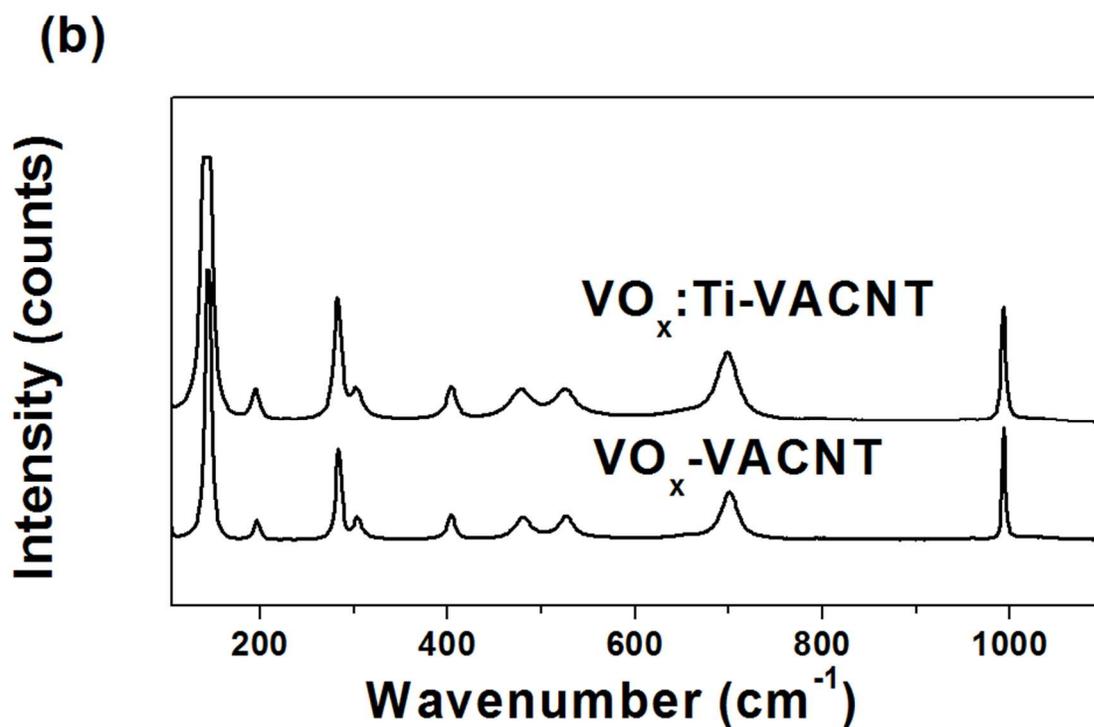
2650

CNT-G'-band

78, 79



465



466

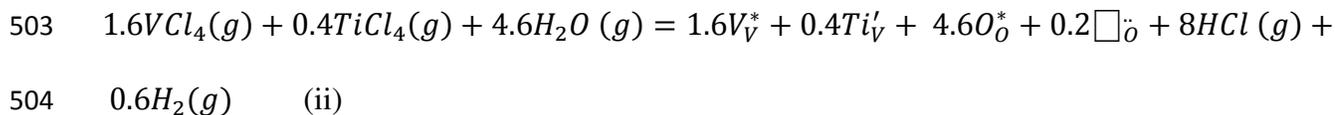
467 **Figure 7:** Raman spectra of vertically aligned CNTs coated with titanium doped vanadium oxide
 468 ($\text{VO}_x\text{:Ti-VACNT}$) over a frequency window of $100\text{-}3000\text{ cm}^{-1}$ showing the presence of oxide
 469 and CNT peaks (a), and comparison of Raman spectra of CVD derived $\text{VO}_x\text{:Ti}$ and VO_x coated
 470 on VACNTs (b).

471 Four-probe conductivity measurements were also performed on both the thin film oxide
 472 and VACNT-supported $\text{VO}_x\text{:Ti}$ nanocomposite heterostructures. Sheet conductivity data was
 473 obtained using a Veeco- FPP 500 instrument wherein the coated Ni disks were suppressed onto
 474 standalone tips and voltage vs. current was measured. Sheet resistance data was multiplied with
 475 the thickness to obtain bulk resistivity which is reported in **Table 2**. The electronic conductivity
 476 of the nickel disk was found to be almost 5 orders of magnitude larger than that of either the
 477 VO_x , $\text{VO}_x\text{:Ti}$, $\text{VO}_x\text{-VACNT}$ or $\text{VO}_x\text{:Ti-VACNT}$ composites indicating that the four-probe

478 conductivity data does not bear any artifact arising from the current collector. The conductivity
479 values fall within the range of resistivities reported (10^1 and $10^4 \text{ } \Omega\text{-cm}^{36, 103}$) in the literature for
480 various oxidation states of vanadium at room temperature. As expected, the electronic
481 conductivity of the VACNT supported VO_x and $\text{VO}_x\text{:Ti}$ is much higher than that of the
482 respective thin film undoped and Ti doped vanadium oxides. The vertically aligned support
483 structure provided by the VACNTs clearly allows for a thin film of oxide to be supported by
484 highly conductive CNTs resulting in more effective transport of electrons through the CNT
485 channels reflected in the improved electronic conductivity. The Ti-doped VO_x supported on
486 VACNT clearly demonstrates the combined benefits of doping as well as the CNT architecture
487 together adding to the improved electronic conductivities even better by more than order of
488 magnitude compared to undoped VO_x supported on VACNTs as seen in Table 2.

489 The deposition of oxide film on CNTs therefore leads to conductivities similar to those of
490 vanadium oxide nanowires reported previously¹⁰⁴. In addition, the electronic conductivity of the
491 $\text{VO}_x\text{:Ti}$ -VACNT can be seen to be almost one order higher than that of the VO_x -VACNT as
492 mentioned above. Moreover, titanium and vanadium have similar ionic radii which allows for the
493 relative ease of doping of titanium into the vanadium oxide lattice. Vanadium doping of titanium
494 oxide has also previously been proposed to result in increased electronic conductivity^{100, 105}. The
495 doping of titanium oxide with vanadium substantially changes its properties; i.e., the life time of
496 charge carriers increases¹⁰⁶ resulting in increased electronic conductivity. It is expected that a
497 similar phenomenon would be reflected to occur in the case of titanium doped vanadium oxide as
498 well. Vanadium oxide is a well-known n-type semiconductor¹⁰⁷. The presence of oxygen
499 vacancies as observed in the EDX data would also contribute to increased electronic
500 conductivity.

501 It is believed the doping of titanium into the vanadium oxide lattice would result in a vacancy
 502 mechanism as proposed in Equation ii.



505 The above proposed mechanism would result in an excess of n-type carriers further increasing
 506 the electronic conductivity of vanadium oxide. Furthermore, the doping of titanium into the
 507 vanadium oxide lattice also results in a shift of the d-band towards a more metallic state resulting
 508 in a further increase in electronic conductivity. This is reflected in Table 2, wherein both the
 509 VO_x : Ti-VACNT and VO_x -Ti have superior electronic conductivity of more than an order of
 510 magnitude in comparison to their undoped counterparts. This increase in electronic conductivity
 511 by doping with titanium would thus lead to enhanced inherent oxide electron transport and hence
 512 enhanced supercapacitor performance.

513

514 **Table 2:** Electronic conductivity data of various titanium doped vanadium oxide and undoped
 515 vanadium oxide and oxide-CNT composite electrodes.

Material	Resistivity (Ω -cm)
VO_x :Ti-VACNT	0.0182
VO_x -VACNT	0.46
VO_x :Ti	0.9
VO_x	7.5

Nickel current collector

$9.3 * 10^{-6}$

516

517 The marked improvement in electronic conductivity by serendipitous introduction of
518 oxygen vacancies and Ti into the V_2O_5 oxide can be clearly demonstrated using first-principles
519 computational approaches implemented for calculations of the electronic structure of the pure
520 and doped vanadium oxide. Changes in the electronic structure of vanadium oxide with the
521 introduction of oxygen vacancies and Ti-atoms in the oxide lattice results in substantial
522 enhancement in the electronic conductivity of the material. For calculations of the total energy,
523 density of electronic states and the cohesive energy (E_{coh} reflecting overall mechanical and
524 chemical stability), the Vienna Ab-initio Simulation Package (VASP) based on the projector-
525 augmented wave (PAW) formalism¹⁰⁸⁻¹¹⁰ within the generalized gradient approximation (GGA),
526 for conducting the exchange-correlation corrections proposed by Perdew *et al.*¹¹¹, has been used
527 in the present study. A plane-wave basis set with a cutoff energy of 520 eV and 184 k -points in
528 the irreducible part of the Brillouin zone have been used to achieve an accuracy in the total
529 energy of $\sim 10^{-2}$ eV per formula unit, which is sufficient for a confident evaluation of the
530 electronic structure, the cohesive energy and comparison of the results with corresponding
531 experimental data.

532 V_2O_5 adopts an orthorhombic crystal structure with a space group $D_{2h}^{13} - Pmmn$ (#59)
533 with two formula units in the elementary unit cell. All structural parameters along with internal
534 atomic Wyckoff positions are summarized in **Table 3** taken from the published report¹¹². It
535 should be noted, that there are three non-equivalent oxygen positions in the unit cell designated
536 as bridge (O_b), vanadyl (O_v), and chain (O_c) oxygen sites.

537 **Table 3:** Crystal structure of V_2O_5 ¹¹²

Lattice constants (Å)		$a = 11.512 \text{ \AA}$	$b = 3.564 \text{ \AA}$	$c = 4.368 \text{ \AA}$
Atoms	Wyckoff positions	x	y	z
O_b	(2a)	0.0	0.0	0.001
O_v	(4f)	0.8543	0.0	0.531
O_c	(4f)	0.6811	0.0	0.003
V	(4f)	0.85118	0.0	0.8917

538

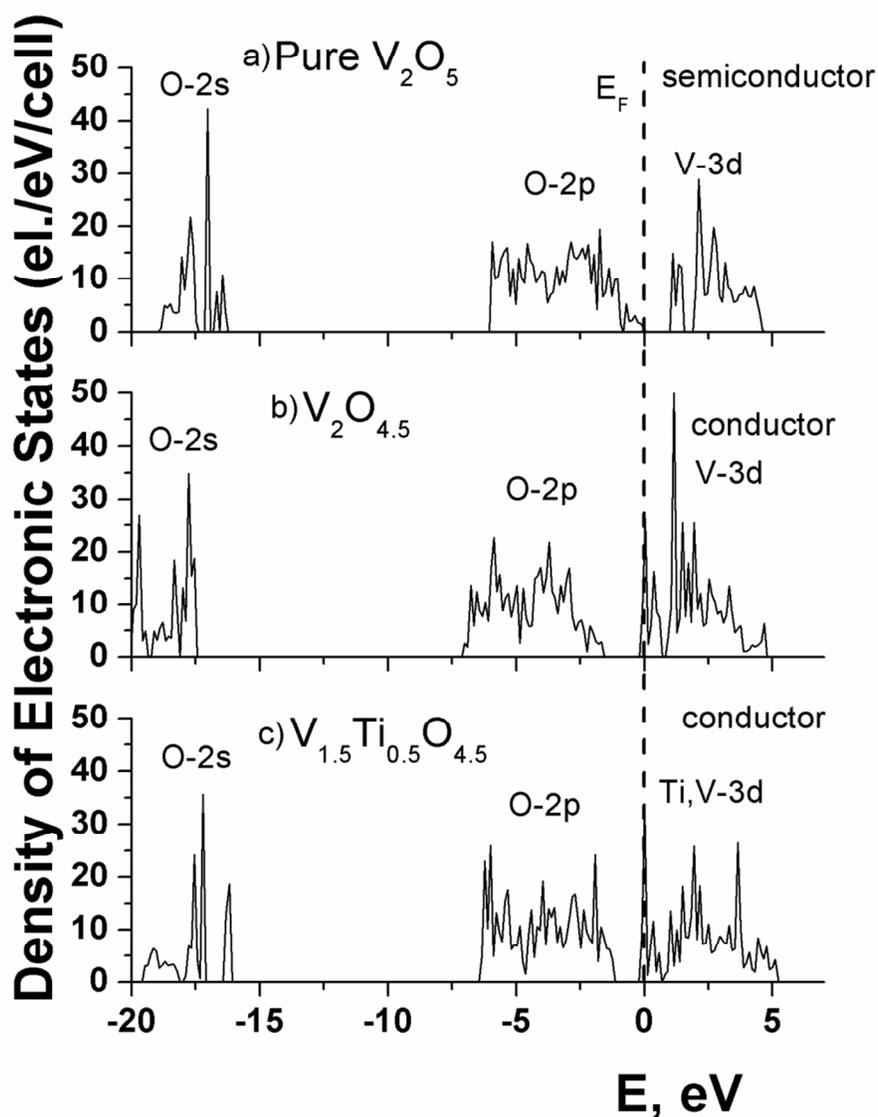
539 To simulate the experimentally studied materials, the following compounds have been
 540 constructed within one elementary unit cell containing 4 V and 10 O atoms: pure V_2O_5 , $V_2O_{4.5}$
 541 with one oxygen atom missing, and $V_{1.5}Ti_{0.5}O_{4.5}$ with one oxygen atom missing, and one V atom
 542 replaced for Ti. Such compositions although not exactly matching the actual experimentally
 543 generated systems, nevertheless could be very close approximations of the materials discussed in
 544 the present study: $V_2O_{4.6}$ and $V_{1.6}Ti_{0.4}O_{4.6}$. For $V_2O_{4.5}$ one vanadyl (O_v) atom has been removed
 545 to generate controlled defects, since previously there was shown in the computational study¹¹³
 546 that vanadyl vacancy at the O_v site is the most favored among all other possible sites in the
 547 vanadium pentoxide. Also, for Ti-doped $V_2O_{4.5}$ a vanadium atom next to the vanadyl vacancy
 548 has been replaced to create the Ti substitution. All the considered structures were fully relaxed
 549 with respect to change in volume, shape, and ion positions within a certain unit cell during the
 550 self-consistent energy minimization.

551 The total density of electronic states calculated for pure V_2O_5 , $V_2O_{4.5}$ and $V_{1.5}Ti_{0.5}O_{4.5}$ are
 552 shown in **Figure 8** with the Fermi level E_F set for zero. One can see that V_2O_5 (**Figure 8a**) is a
 553 semiconductor with a forbidden energy gap between the valence and conduction zones. The
 554 valence band consists primarily of strongly hybridized oxygen 2p- and vanadium 3d- electronic

555 states with chiefly oxygen 2p-states, while the conduction band mainly consists of the vanadium
556 3d-states. This electronic structure correlates very well with results obtained in other theoretical
557 studies previously published¹¹³⁻¹¹⁵. Our studies involving $V_2O_{4.5}$ and $V_{1.5}Ti_{0.5}O_{4.5}$ are unique
558 however and provide new insight into the behavior of such doped structures. An internal analysis
559 of the electronic structure showed that the valence band consisting primarily of O 2p states
560 demonstrates a pronounced bonding character, while the conduction band comprises the anti-
561 bonding states which was also noted in the work by Eyert et al.¹¹⁴.

562 **Figure 8b** shows the total density of electronic states of $V_2O_{4.5}$. Since, one oxygen atom
563 is missing, formally two vanadium 3d-electrons become free (unbounded) and thus fill the
564 bottom of the previously empty conduction zone with these corresponding two electrons indeed
565 rendering the material metallic. This phenomenon correlates well with results on the conductivity
566 obtained in present study for $V_2O_{4.6}$ and discussed earlier.

567 **Figure 8c** demonstrates the total density of states for $V_{1.5}Ti_{0.5}O_{4.5}$. One can see that in
568 general, the electronic structure is very similar to the previous case of $V_2O_{4.5}$ containing the
569 oxygen vacancy with the only difference that now there is one Ti atom substituting in place of a
570 V. Comparing these two cases it can be noted that although one oxygen vacancy releases two
571 vanadium electrons in the unit cell, the introduction of a Ti atom bears one electron less than that
572 of V, therefore contributing to only one electron per unit cell capable of conducting the electric
573 charge through the material. Thus, $V_{1.5}Ti_{0.5}O_{4.5}$ still remains electronically conductive.



574

575 **Figure 8:** Total density of electronic states calculated for pure V_2O_5 , $V_2O_{4.5}$ and $V_{1.5}Ti_{0.5}O_{4.5}$.

576

577 The metallic conductivity is usually proportional to the density of states at the Fermi level
 578 $n(E_F)$ and hence it provides an opportunity to qualitatively evaluate the electronic conductivity of
 579 $V_2O_{4.5}$ and $V_{1.5}Ti_{0.5}O_{4.5}$ oxides. For these purposes **Table 4** collects $n(E_F)$ for these materials
 580 showing that $n(E_F)$ calculated for $V_2O_{4.5}$ is lower than that corresponding to $V_{1.5}Ti_{0.5}O_{4.5}$. The

581 presence of Ti 3d-states increased to some extent the total density of states in the vicinity of the
 582 Fermi level thus rendering the material more electronically conductive than just the defective
 583 $V_2O_{4.5}$ oxide. This result is well confirmed by previously discussed experimental data
 584 summarized in **Table 2**.

585 **Table 4:** Comparison of density of states at Fermi energy level and cohesive energy between
 586 pure V_2O_5 , vacancy-rich V_2O_5 and titanium doped vanadium oxide

Materials	$n(E_F)$, (el./ $(eV \text{ f.un.})$)	E_{coh} (eV/f.un.)
V_2O_5	0 (band gap)	-42.17
$V_2O_{4.5}$	11.5	-39.05
$V_{1.5}Ti_{0.5}O_{4.5}$	15.5	-39.54

587

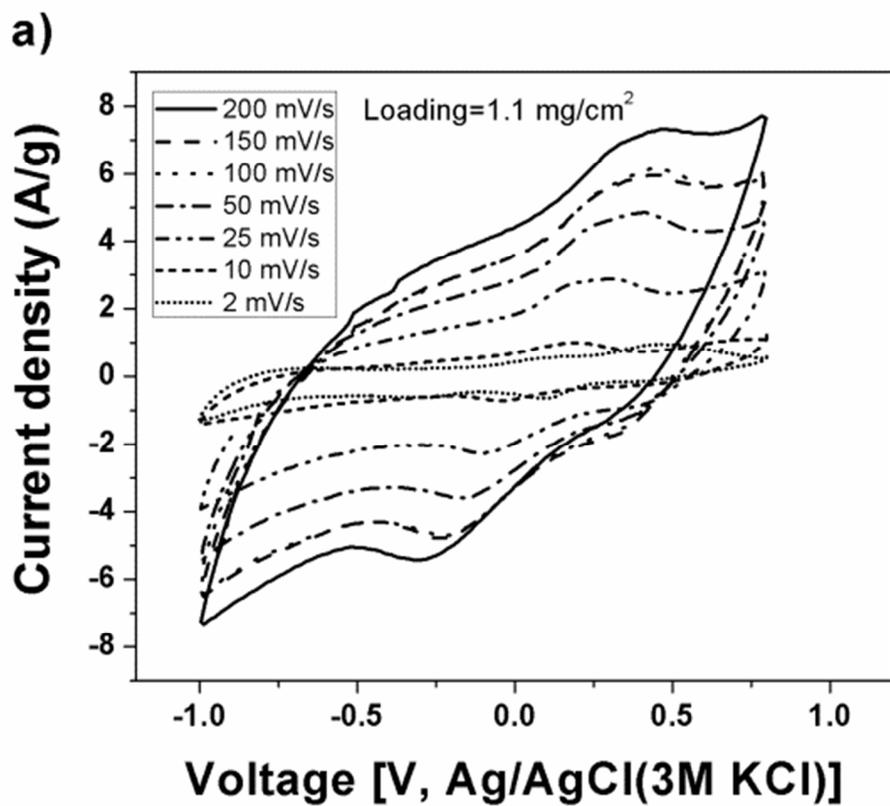
588 It should be noted that in this study the theoretical first principles analysis was used to
 589 explain the underlying phenomena resulting in superior electrical conductivity of the doped
 590 oxide materials. Although the first principles studies were performed on crystalline materials and
 591 not on the experimentally observed amorphous state though similar in composition, the results
 592 can nevertheless be considered as a qualitative representation of the behavior of amorphous
 593 materials corresponding to similar compositions. Amorphous materials as is known can be
 594 considered to be domains of short range periodicity not extending itself into the meso-scale state.
 595 Thus it can be hypothesized that the behavior of the larger body consisting of an infinite number
 596 of such domains would to some extent parallel the behavior of the domain itself. Though such a
 597 consideration is not perfect in the conventional and ideal sense, the experimental evidence
 598 nevertheless (**Table 2**) herein parallels the trends in electronic conductivity (**Table 4**), hence
 599 justifying the assumption to be valid at least for the case applicable to the present system.

600 Further, as mentioned above, the material should demonstrate high stability in the voltage of
601 interest. For these purposes the cohesive energy E_{coh} can be considered a measure of an overall
602 structural and chemical stability of the material. A higher E_{coh} would therefore imply that more
603 energy is required to break the primary chemical bonds of the compound during chemical
604 reaction. Thus, the cohesive energies of all three compounds considered in the present theoretical
605 study have been calculated and summarized in **Table 4**.

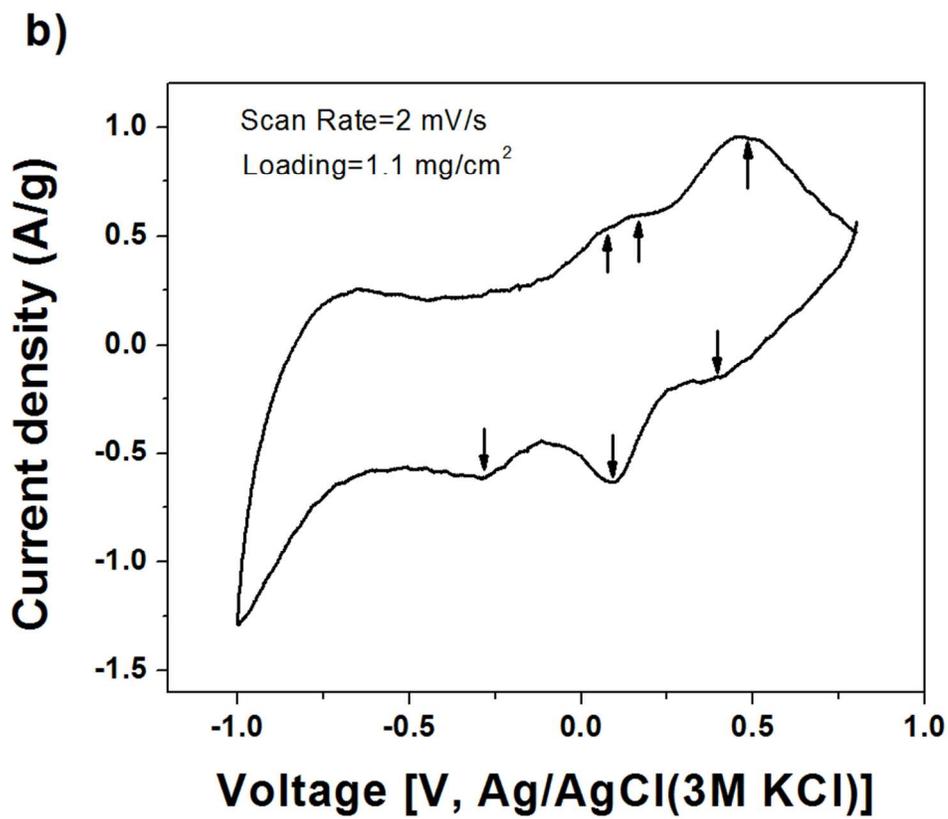
606 From **Table 4** it can be construed that pure V_2O_5 is the most stable material in
607 comparison with both the doped and defective oxides. It is therefore expected that the defective
608 oxide with oxygen vacancies $\text{V}_2\text{O}_{4.5}$ would be less stable due to filling up of the anti-bonding
609 electronic states in the conduction band above the band gap and also because of lowering the
610 total number of V-O bonds in comparison with non-defective V_2O_5 . Thus, presence of O-
611 vacancies decreases the overall stability of the material to some extent. The situation albeit
612 improves with the introduction of the Ti-doping. Because Ti-O bonds are stronger than V-O (668
613 kJ/mol vs. 637 kJ/mol bond strength in Ti-O and V-O diatomic molecules, respectively¹¹⁶) the
614 overall stability of the Ti-doped vanadium oxide is higher than the undoped oxide.

615 Thus, in terms of the electronic conductivity and the overall stability the present
616 theoretical study showcases $\text{V}_{1.5}\text{Ti}_{0.5}\text{O}_{4.5}$ as a promising super-capacitor material with excellent
617 electrochemical properties. To the best of our knowledge, no theoretical studies exist correlating
618 the effect of vacancies on the electronic structure of vanadium oxide and titanium oxide for
619 supercapacitor applications. Based on the insight provided by this novel ab-initio study, we
620 expect $\text{V}_{1.5}\text{Ti}_{0.5}\text{O}_{4.5}$ to be an excellent candidate for supercapacitor applications.

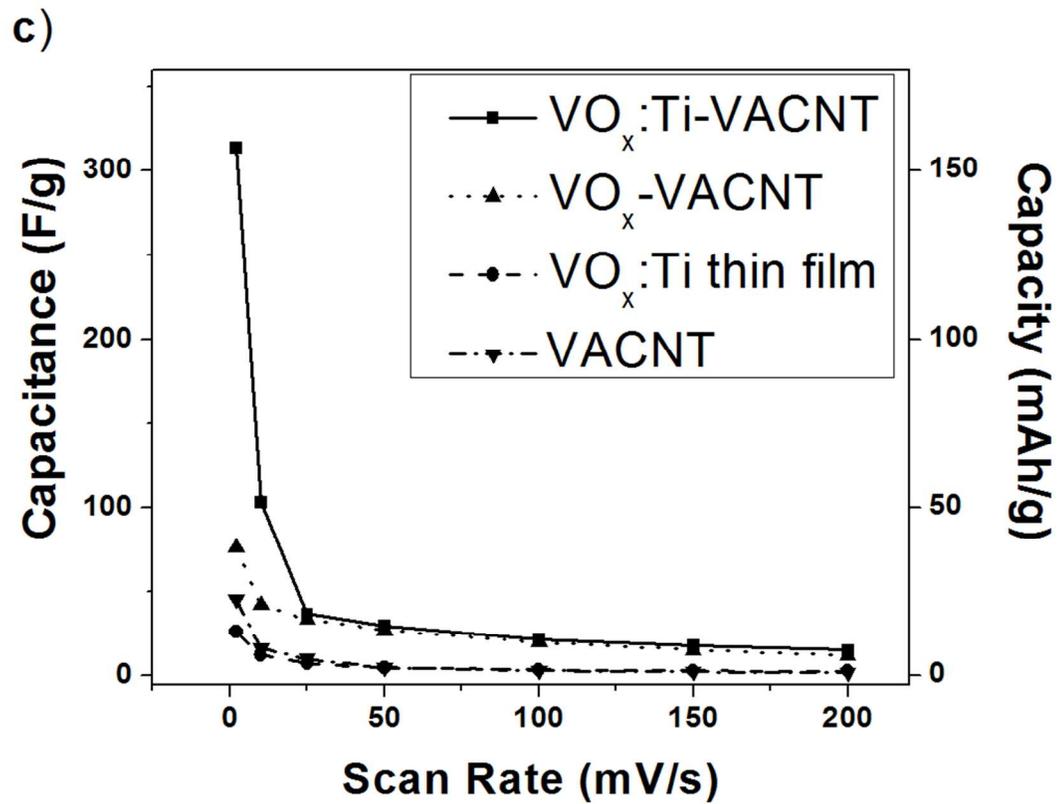
621



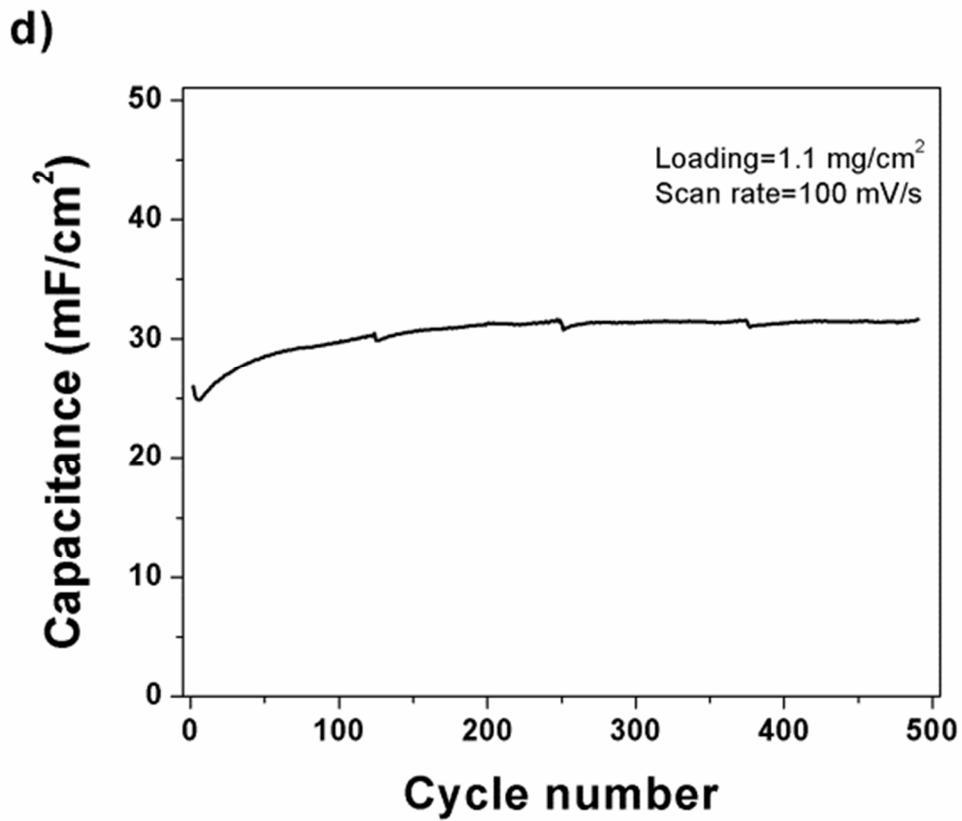
622



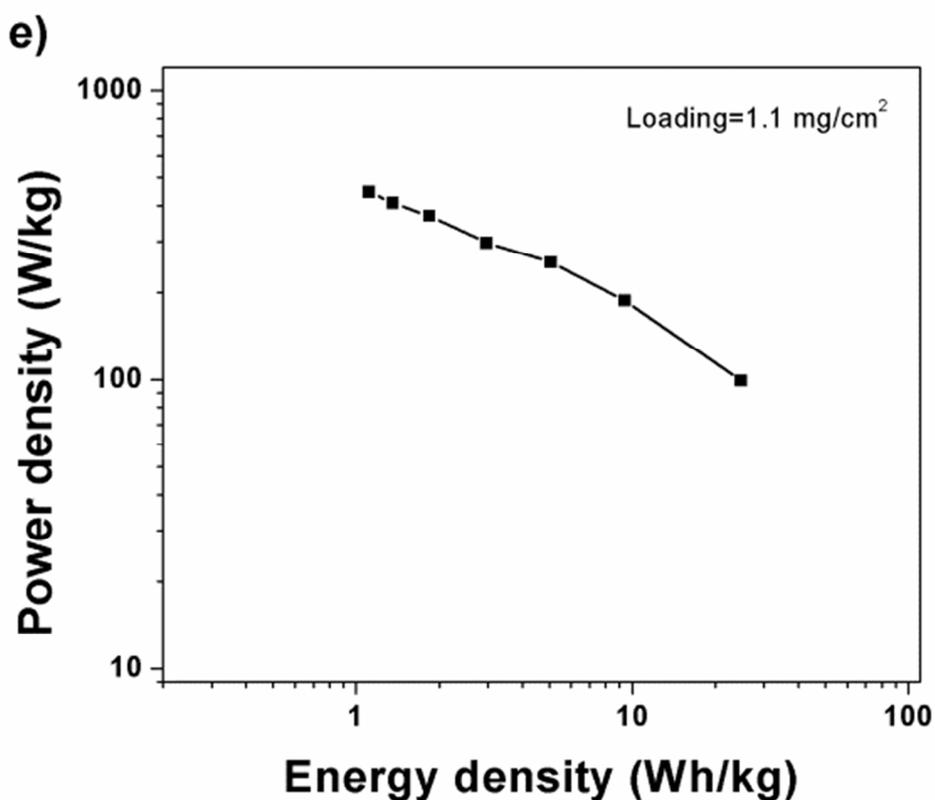
623



624



625



626

627 **Figure 9:** (a) Cyclic voltammograms of the $\text{VO}_x\text{:Ti-VACNT}$ heterostructures at different scan
628 rates, (b) Cyclic voltammograms of the $\text{VO}_x\text{:Ti-VACNT}$ at 2mV/sec showing anodic and
629 cathodic peaks for the reversible pseudocapacitive reactions, (c) Variation of capacitance with
630 respect to scan rate for $\text{VO}_x\text{:Ti-VACNT}$, $\text{VO}_x\text{-VACNT}$ and thin film $\text{VO}_x\text{-Ti}$ (oxide
631 loading=0.55 mg/cm^2 , Oxide-VACNT active material loading=1.1 mg/cm^2) (d) Cycling data
632 showing capacitance as a function of cycle number conducted at a scan rate of 100mV/sec, and
633 (e) Ragone plot (power density vs. energy density) for $\text{VO}_x\text{:Ti-VACNT}$.

634

635 **Figures 9a** shows the cyclic voltammograms of the $\text{VO}_x\text{:Ti-VACNT}$ at different scan rates
636 ranging from 200mV/sec to 2mV/sec. The change in slopes of the I-V curves clearly shows the
637 expected effects of scan rate. As the scan rate increases, the current response on scanning
638 potential undergoes a transition from a flat capacitor type behavior to a resistor-like dependence
639 as is common for most pseudocapacitor materials. The fade in capacitance with increase in scan
640 rate is a phenomenon observed in all typical supercapacitor materials as described by Conway et.

641 al⁶. This is caused by the change in characteristic of the interface from that of a double layer
642 capacitor to a typical resistor due to the inherent limitations and complexities of charge transfer
643 and electron mobility dynamics.

644 It is well known that capacitance changes with scan rate, and is indeed a factor inherently
645 dependent on the specific materials' characteristics. A similar characteristic in Faradaic reactions
646 is the kinetic rate constant (k_0) which is a material property in a given electrode/electrolyte
647 assembly. Pseudocapacitance reactions involving charge transfer by tunneling across the
648 electrochemical interface have an inherent rate constant depending on the various species
649 involved in the charge transfer which could be limiting in the case of the doped oxide coated on
650 the CNTs. It is observed that even though the carbon nanotubes have a vertical alignment, the
651 nature of the tethering of the oxide film on the carbon nanotubes needs to be improved to achieve
652 better rate capability. Optimization of the system to improve the tethering of the doped
653 oxide/VACNT and the CNT inherent conductivity would ideally increase both the initial
654 capacitance as well as the rate capability, two of the most desired signatures characterizing the
655 efficacy of a capacitor system. Such optimization studies are currently ongoing and will be
656 reported in a forthcoming publication.

657 **Figure 9b** shows the slow scan rate cyclic voltammogram of the VO_x:Ti-VACNT
658 composite electrode. Since voltammetry was performed at a slow rate, the Faradaic
659 pseudocapacitance processes are hence not rate limited and distinct peaks occur corresponding to
660 each Faradaic process is clearly visible. Reversible Faradaic peaks are observed at ~0.48V,
661 0.17V and 0.07V (*vs.* Ag/AgCl) on the positive scan and ~0.40V, 0.09V and -0.28V (*vs.*
662 Ag/AgCl) on the negative scan in the VO_x:Ti-VACNT composite (see **Figure 9b and Table 5**).
663 **Table 5** also shows tabulated the possible potential dependent pseudocapacitance type reactions

664 present at these potentials corresponding to a reversible transition of oxidation state of vanadium
 665 from +5 to +3, the most commonly observed in aqueous systems at the slightly alkaline pH
 666 (~8.5) occurring in a 1M Na₂SO₄ solution^{117, 118}. In addition, the onset of water electrolysis is
 667 seen to occur at approximately, -0.75V (vs. Ag/AgCl) in **Figure 9b** at a low scan rate of 2 mV/s.
 668 This however is not observed at higher scan rates indicating that the kinetics of electrolysis are
 669 very sluggish on the VO_x:Ti-VACNT composite heterostructure electrodes, despite the improved
 670 electronic conductivity indicated in Table 2. We can therefore use this material over an extended
 671 window of ~1.8V with minimal or no electrolyte loss caused due to water electrolysis.

672 It should also be noted that in titanium oxide, the electronic state transition of Ti from +4
 673 to +3 occurs generally around at ~-0.6V (vs. Ag/AgCl) at pH ~ 8.5^{77, 78}. However, we can see
 674 that these transitions are not observed (see **Figure 9a** and **9b**) indicating that the titanium is
 675 completely doped into the VO_x lattice forming a solid solution with and no phase separation or
 676 secondary oxide formation. This further complements the Raman results confirming the Ti
 677 doping in the VO_x lattice.

678 **Table 5:** Faradaic charge transfer processes observed in the titanium doped vanadium oxide-
 679 VACNT composite electrode

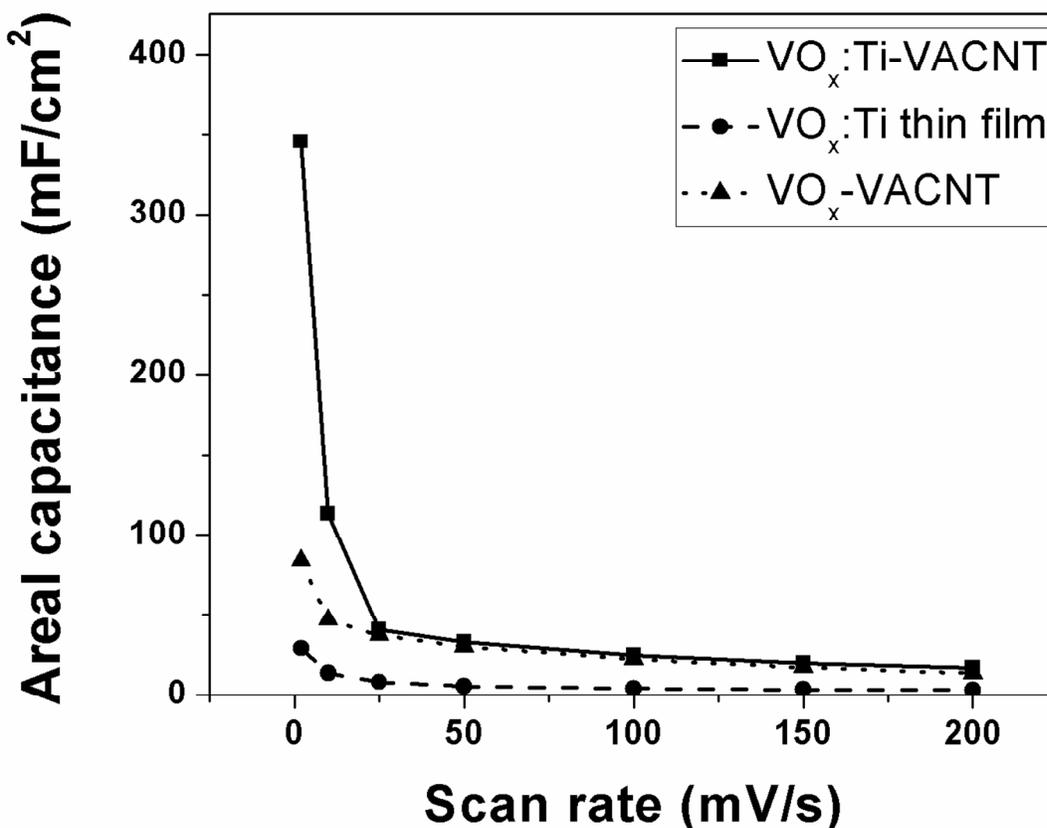
Anodic scan peak potential (vs. Ag/AgCl)	Cathodic scan peak potential (vs. Ag/AgCl)	Possible reaction mechanism
~-0.4839V	~-0.3997V	$3\text{V}_2\text{O}_5(\text{s}) + 2\text{H}_2\text{O} + 4\text{e}^- \leftrightarrow \text{V}_6\text{O}_{13}(\text{s}) + 4\text{OH}^-$ $E = 0.8425 - 0.0591 \times \text{pH} \text{ (Volts vs. SHE)}$ $E = 0.1302 \text{ (Volts vs. Ag/AgCl (3M KCl))}$
~-0.1669V	~-0.0922V	$\text{V}_6\text{O}_{13}(\text{s}) + \text{H}_2\text{O} + 2\text{e}^- \leftrightarrow 3\text{V}_2\text{O}_4(\text{s}) + 2\text{OH}^-$ $E = 0.4352 - 0.0591 \times \text{pH} \text{ (Volts vs. SHE)}$ $E = -0.2768 \text{ (Volts vs. Ag/AgCl (3M KCl))}$

~0.074V	~-0.276V	$3\text{V}_2\text{O}_4(\text{s}) + 2\text{H}_2\text{O} + 4\text{e}^- \leftrightarrow 2\text{V}_3\text{O}_5(\text{s}) + 4\text{OH}^-$ $E = 0.3901 - 0.0591 \times \text{pH} \quad (\text{Volts vs. SHE})$ $E = -0.3222 \quad (\text{Volts vs. Ag/AgCl (3M KCl)})$

680

681 The variation of the capacitance of VO_x:Ti-VACNT, VO_x-VACNT and VO_x:Ti on Ni
682 substrate with respect to scan rate is plotted in **Figure 9c** (oxide loading~0.55 mg/cm², oxide-
683 VACNT loading~1.1 mg/cm²) in addition to that of plain VACNT electrodes (used as a control).
684 It should be noted that the capacitances of ~315 F/g (~157 mAh/g) shown in **Figure 9c** on the
685 basis of the total active material loading (~1.1 mg/cm²) (not just vanadium oxide loading) has
686 been observed at a scan rate of 2 mV/s in the VO_x:Ti-VACNT nanocomposite material with
687 excellent capacitance retention at higher rates. It can be seen that the VACNT electrode used as a
688 control has a very low capacitance of ~45 F/g which drops with scan rate indicating that a very
689 small part of the overall charge storage arises on account of the carbon nanotubes itself which is
690 to be expected given that very small part of the CNTs are actually exposed to the electrolyte on
691 account of being covered with the VO_x:Ti. The drop in capacitance with scan rate is due to the
692 effect of the resistance/capacitance behavior previously described. Also, VO_x:Ti-VACNT has a
693 higher capacitance not only at low scan rates but also at higher scan rates compared to VO_x: Ti
694 thin film. In addition to capacitance, we have reported capacity in mAh/g for the various
695 electrodes in **Figure 9c**. Vanadium oxide materials have previously been explored as charge
696 storage electrodes in aqueous solutions in the form of aqueous batteries^{110, 114}. However, it
697 should be noted that the capacity obtained herein is almost twice that of previous studies wherein
698 capacities of ~60-70 mAh/g have been reported with the electrode capacity rapidly fading to ~22

699 mAh/g over just a few cycles¹¹⁹. The areal capacitance for the VO_x:Ti-VACNT is between 350
700 mFcm⁻² and 30 mFcm⁻² (see **Figure 10**) which is superior to that reported for other
701 supercapacitor materials at similar scan rates^{120, 121}. The performance of supercapacitor materials
702 at high rates is most-often limited by the electronic conductivity of the material. The presence of
703 VACNTs as a support structure for the oxide leads to unidirectional electron transfer through the
704 VACNTs¹²² which then travel through the semi-contiguous film and interact at the oxide-
705 electrolyte interface where the charge is stored. This allows for maximum CNT-oxide electron
706 transport while minimizing electron transport through the VO_x grain boundaries which act as a
707 barrier to electron transport. The present morphology also leads to increased surface exposure of
708 the oxide to the electrolyte *i.e.* maximizing the electrochemical surface area and affording
709 maximum exploitation of the capacitance behavior.



710

711 **Figure 10:** Variation of areal capacitance with respect to scan rate for VO_x:Ti-VACNT , VO_x-
 712 VACNT and thin film VO_x-Ti (oxide loading=0.55 mg/cm², Oxide-VACNT active material
 713 loading=1.1 mg/cm²)

714 The cycling stability of the VO_x:Ti-VACNT supercapacitor material to retain charge over
 715 400 cycles at a scan rate of 100mV/sec is shown in **Figure 9d**. Previous studies of vanadium
 716 oxide electrodes in aqueous electrolytes for charge storage have seldom yielded such superior
 717 electrochemical stability^{119, 123}. It can be observed that the areal capacitance increases initially
 718 stabilizing at a value of ~31 mF/cm². This excellent cyclability can be attributed to the reversible
 719 nature of the surface pseudocapacitance reactions of the oxide (see **Table 5**). The slight initial
 720 increase in capacity is possibly a result of the gradual improvement in wetting of the VO_x:Ti-

721 VACNT composite with the electrolyte. No significant difference is seen in the cycling behavior
722 of VO_x-VACNT (previously reported by us¹⁷) and VO_x:Ti-VACNT reported herein indicating
723 that both structures are stable in the electrolyte over the chosen voltage window. The oxidation
724 state of the CVD deposited vanadium oxide is ~4.6 consisting predominantly of a +4/+5 oxide.
725 According to the potential-pH diagram, at the pH of ~8.5, vanadium oxide has stable insoluble
726 phases between the potentials of ~0.5 V and -1 V wrt Ag/AgCl¹¹⁷. Above 0.5 V, the higher
727 oxidation state oxide is considered soluble to form H₂VO₄⁻ at a neutral pH. This should be
728 reflected as a loss in capacity with cycling. However, the onset potential of the electrochemical
729 reactions is kinetically limited and is known to vary with particle size and the ensuing
730 microstructure¹²⁴. The Ti doped oxide is amorphous in the present work and most amorphous
731 materials will be expected to exhibit lower stability compared to the crystalline counterpart due
732 to the higher reactivity and defective state. However, the stronger Ti-O bonds and the overall
733 higher cohesive energies for the crystalline counterpart translated to the amorphous state could
734 likely contribute to the better stability of the system. Hence it is believed that the amorphous
735 nature of the oxide nanoparticles supported on the vertically aligned carbon nanotubes could
736 offset the onset of such dissolution resulting in the observed very stable performance of the oxide
737 in aqueous Na₂SO₄ shown in **Figure 9d**. It is also possible that the kinetics of the dissolution
738 reaction might be sluggish and such a fade in capacity would probably be observed if the
739 material were cycled over several thousands of cycles which was not conducted in the present
740 study since the objective was to demonstrate the performance of the doped system supported on
741 VACNTs rather than demonstrate optimization of the system which will be part of the continued
742 study that will be reported in subsequent publications. No post-cycling solution analysis was
743 performed since significant fade was not observed over 500 cycles as seen in **Figure 9d**.

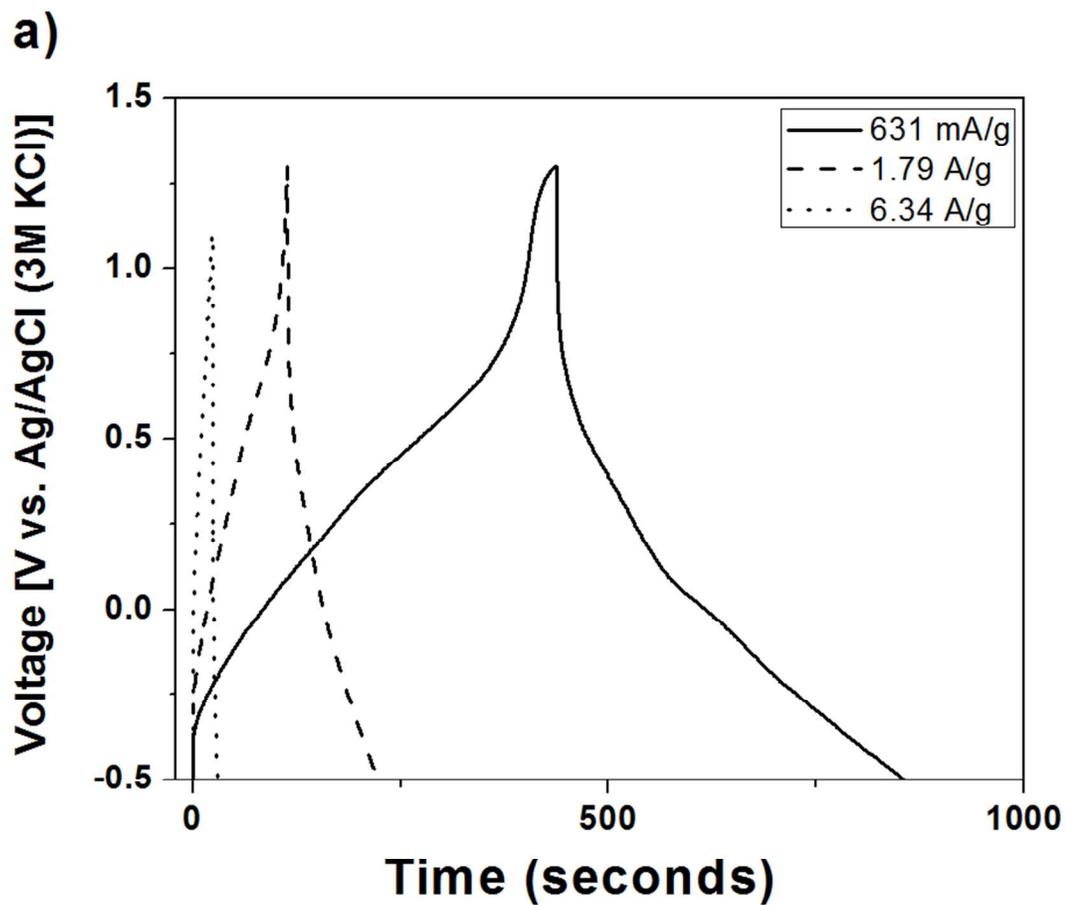
744 The synergistic effect of the carbon nanotube-vanadium oxide junction has previously
745 been studied for gas-sensing purposes^{95, 125, 126}. It has been suggested that a p-n junction type
746 interaction between the VO_x and the CNT leads to increased catalytic response of the oxide
747 while the multi walled CNTs, which acts as a p-type semiconductor facilitates the fast electron
748 channels. A similar p-n type interaction can be expected to be observed in VO_x:Ti-VACNT
749 composite as it is known that VO_x:Ti tends to exhibit predominantly n-type carrier behavior. All
750 these factors are reflected in the improved electronic conductivity (see **Table 2**) contributing to
751 achieving the improved rate capability, excellent supercapacitor response and cycling stability as
752 seen in **Figure 9c and 9d**. Improved charge storage behavior may also be attributed to
753 electronic sensitization of the doped oxide by this effect as commonly observed in doped oxides
754 used for gas sensing applications¹²⁷.

755 It is common knowledge that thick films of nanoparticulate oxides especially amorphous
756 oxides would be non-ideal configurations to obtain high capacitances on account of the low
757 inherent electronic conductivity, grain boundary diffusion and exposed surface area. On the other
758 hand, the use of a relatively thin film on a supported structure affords the benefits of high surface
759 area, small particle size and direct contact of grains with the highly conductive carbon nanotube
760 surface. This accessible capacitance of vanadium oxide is very much limited when the electrode
761 is cast as a thick film. However, thin film titanium vanadium oxide has a high capacitance when
762 supported on VACNTs as demonstrated in **Figure 9c**. The VACNTs act as a 3-dimensional
763 electronically conductive support resulting in enhanced electrochemical surface area reflected
764 both in the improved performance and rate capability of the VO_x:Ti-VACNT heterostructure
765 composite. The overall capacitance reported in this manuscript is however lower than that
766 reported in LiCl based systems by Hu et al.²⁶, as there is no intercalation mechanism involved in

767 tandem with the pseudocapacitance processes combined with double layer charge storage seen in
768 the present VO_x:Ti-VACNT composite. Evaluation of VO_x:Ti-VACNT and other doped
769 vanadium oxide systems is ongoing in intercalation based electrolytes and it is expected that
770 higher energy density can be obtained using this system. The unique synthesis method,
771 composition, morphology, electrode/electrolyte combination and an in-depth electrochemical
772 characterization provided in this detailed study are original contributions enhancing the
773 understanding of supercapacitor behavior in early transition metal oxide systems to the best of
774 the understanding known from the published literature to date.

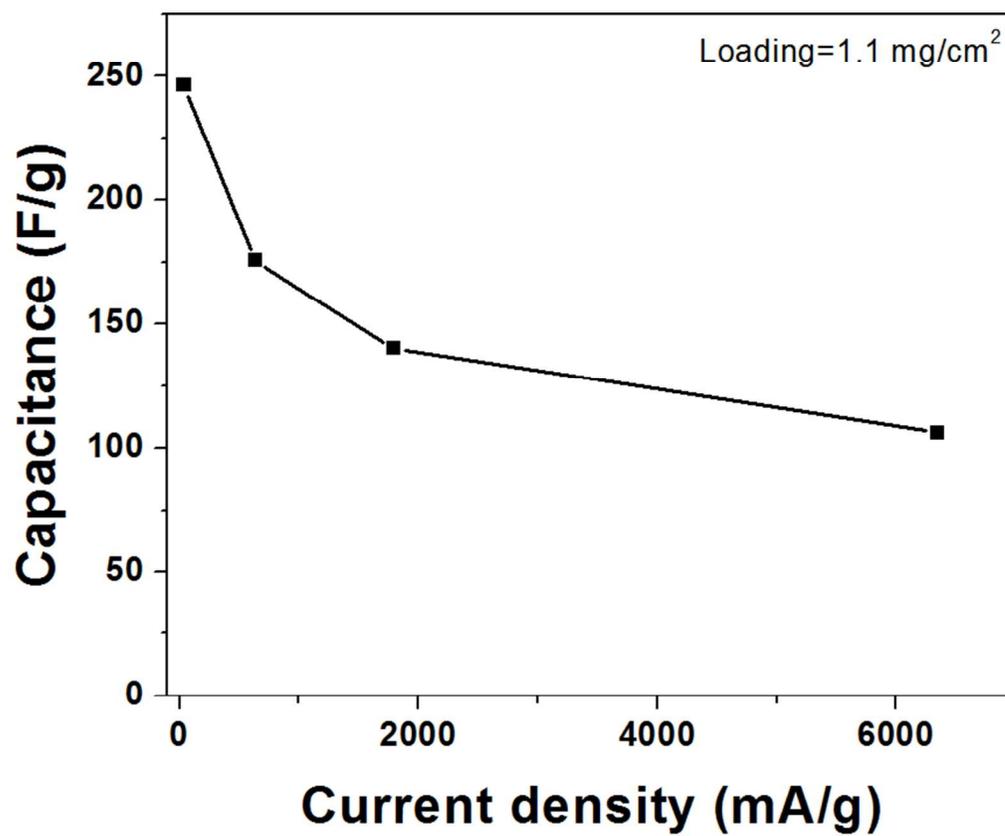
775 The above-stated improvements in the VO_x:Ti-VACNT are reflected as the excellent
776 supercapacitor performance seen in the Ragone plot (**Figure 9e**). The Ragone plot is
777 conventionally used to represent the variation of energy density as a function of power density.
778 An ideal device would maintain its energy density upon increasing power density. However, this
779 is not very commonly observed as higher power densities indicate a greater diversion from
780 equilibrium resulting in an L-shaped curve commonly seen showing higher power densities with
781 low energy density and corresponding decrease in power densities with increasing energy
782 density. In the case of the VO_x:Ti-VACNT high energy densities are sustained by the system
783 with increasing power density although there is also observed a fade in energy density with rise
784 in power density. The highest energy density of ~25 Wh/kg was obtained at a power density of
785 ~100 W/kg. This is comparable to noble metal oxide (~25 Wh/kg) and very thin film based oxide
786 (~30 Wh/kg) supercapacitor electrodes¹²⁸⁻¹³⁰ and superior to other oxide materials(10-20 Wh/kg)
787 representing the novelty of the work conducted in this study¹³¹.

788

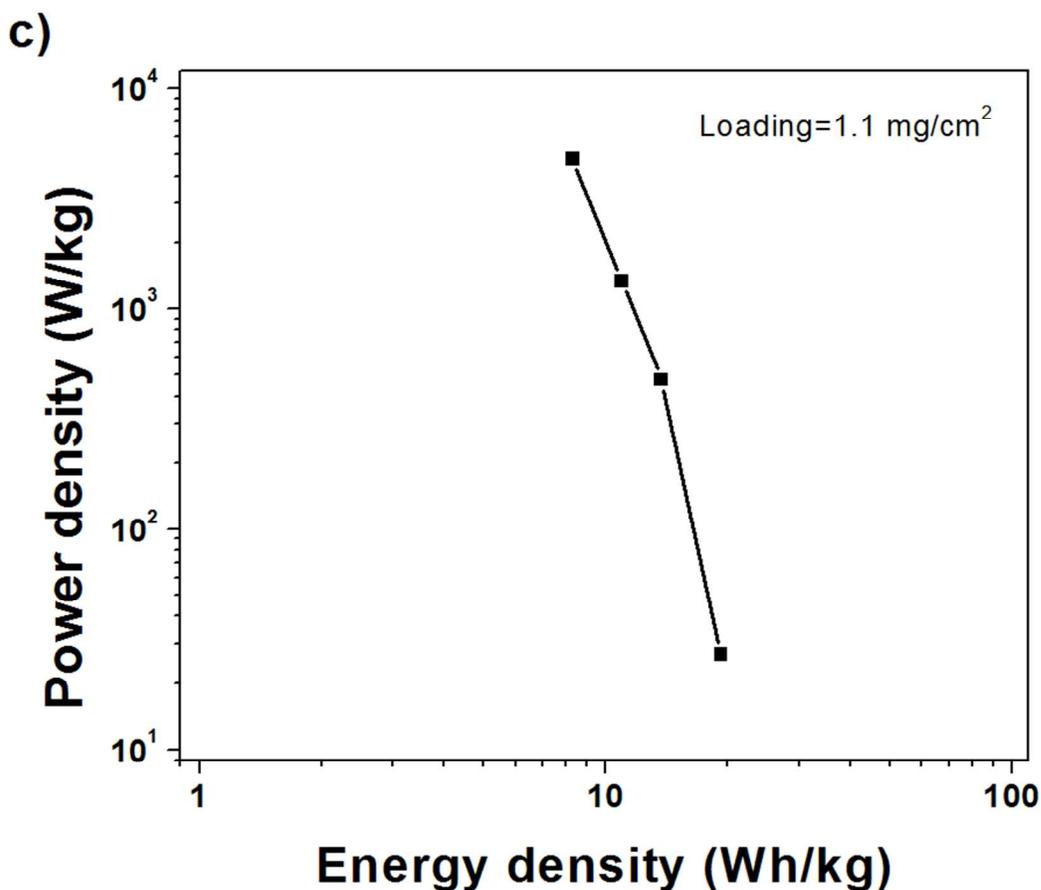


789

b)



790



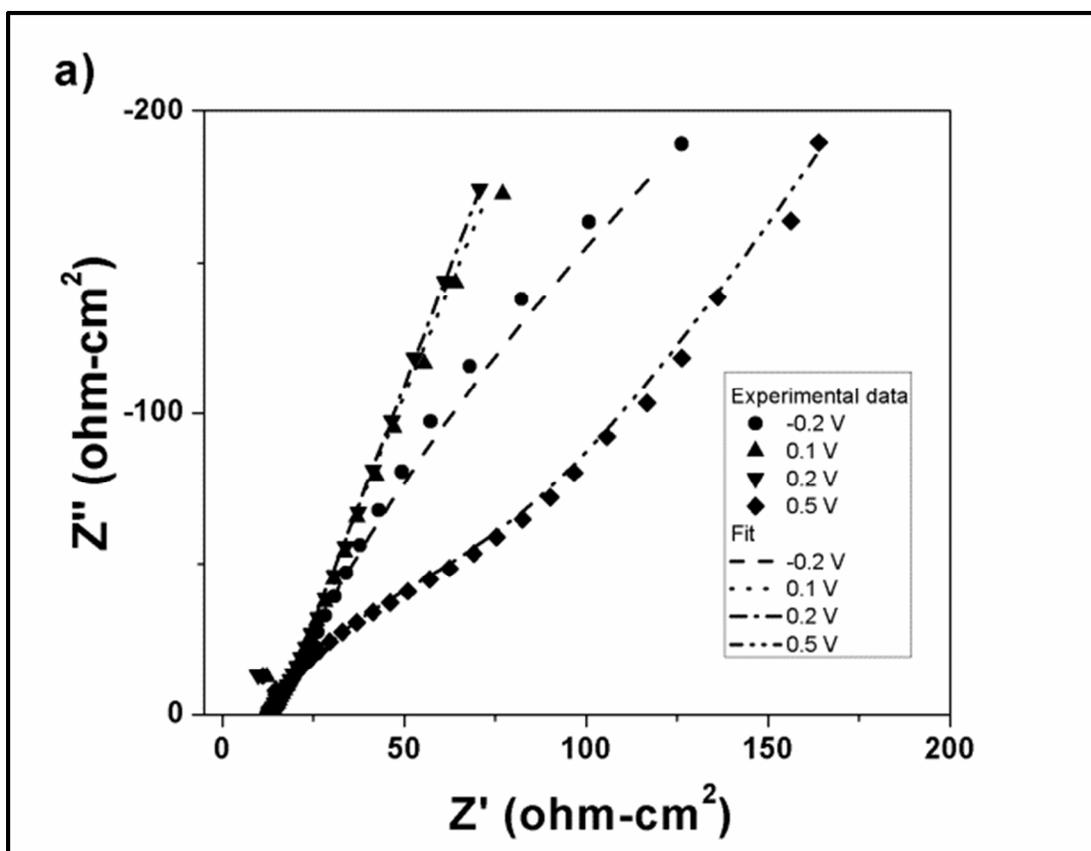
791
 792 **Figure 11:** (a) Galvanostatic charge-discharge profiles of the VO_x:Ti-VACNT material (oxide
 793 loading=0.55 mg/cm², Oxide-VACNT active material loading=1.1 mg/cm²) at different current
 794 densities (b) Capacitance dependence of the VO_x:Ti-VACNT material on current density
 795 calculated from the charge-discharge profiles in (a) (c) Ragone plot (power density vs. energy
 796 density) for VO_x:Ti-VACNT calculated from the galvanostatic charge-discharge profiles.

797
 798 Galvanostatic charge-discharge was performed on the most optimal material i.e. VO_x:Ti-
 799 VACNT in order to simulate the more practical commercial conditions and to examine the
 800 performance under high load. **Figure 11a** depicts the characteristic charge-discharge curves at
 801 different current densities. The material exhibits pseudocapacitor type behavior with very slight

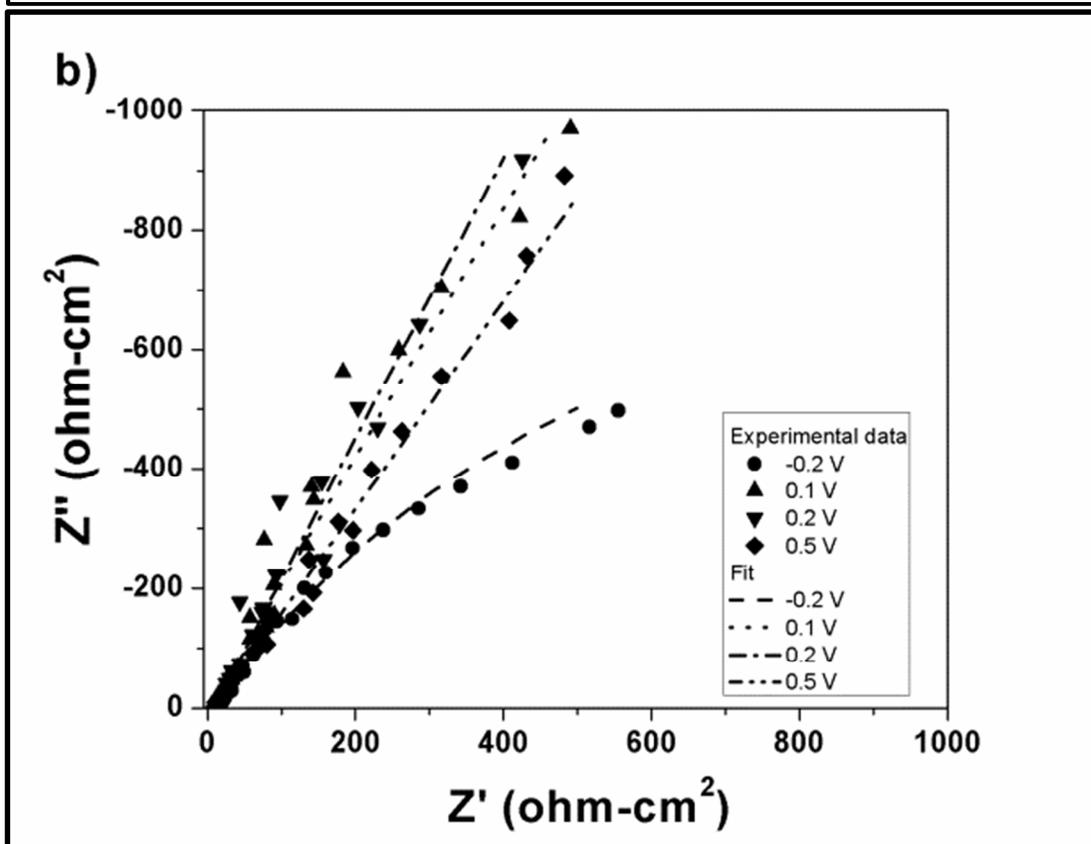
802 shoulders at the potentials corresponding to the Faradaic reactions described in **Table 5**. It
803 should be noted herein that the charge-discharge experiments were performed over a voltage
804 window of -0.5 V to 1 V with respect to Ag/AgCl reference electrode. This was because onset of
805 hydrogen evolution was found to occur at various potentials beyond ~ -0.6 V depending on the
806 current density. This was to be expected given that the window for water electrolysis is normally
807 at 1.2 V. However, a stable reproducible charge-discharge behavior was observed over 1.5 V
808 with no onset of electrolysis indicating that this material can be used over such an extended
809 voltage window on account of the sluggish kinetics of electrolysis occurring on this material
810 described earlier (**Figure 9b**). **Figure 11b** depicts the capacitance calculated from the charge-
811 discharge profiles shown in **Figure 11a** as a function of current density. It can be seen that the
812 $\text{VO}_x\text{:Ti-VACNT}$ material exhibits high capacitance values of ~ 250 F/g at low scan rates which is
813 retained very well even at higher scan rates. Capacitances of ~ 125 F/g are observed at very high
814 current densities of ~ 6 A/g demonstrating the superior charge storage characteristics of the
815 material as a result of the improved electronic conductivity and improved Faradaic reaction
816 kinetics yielded by doping of the oxide as well as coating on the vertically aligned carbon
817 nanotubes. The Ragone plot obtained from the galvanostatic charge-discharge data has been
818 plotted in **Figure 11c**. Excellent energy densities of ~ 11 Wh/kg are retained at very high power
819 densities of ~ 4.5 kW/kg, a demonstration of the excellent charge storage characteristics of the
820 doped $\text{VO}_x\text{:Ti-VACNT}$ material.

821 In order to further establish the advantages yielded by the CVD deposition of $\text{VO}_x\text{:Ti}$ on
822 VACNTs, electrochemical impedance spectroscopy (EIS) analysis was conducted on the
823 nanospheres as well as the doped oxide film and charge transfer characteristics were accordingly
824 compared. **Figures 12a** and **12b** show the Nyquist plots of the $\text{VO}_x\text{:Ti-VACNT}$ hybrid and the

825 VO_x:Ti thin film oxide, respectively at different voltages over the window of interest. It can be
826 directly observed that the impedance behavior of both materials is different especially at voltages
827 close to the regions of the Faradaic peaks observed at ~0.48 V, 0.17 V and 0.07 V on the anodic
828 scan and ~-0.40 V, 0.09 V and -0.28 V in the cyclic voltammograms (see **Figure 9b**). At these
829 potentials corresponding to the Faradaic charge transfer reactions, it can be seen upon comparing
830 the response of VO_x:Ti and VO_x:Ti-VACNT composite that the thin film oxide has a much
831 larger charge transfer resistance. This is inferred from the large semi-circular arcs seen in **Figure**
832 **12b** as compared to **Figure 12a** where one can see the Warburg tail being more prominent in
833 **Figure 9a** as compared to the relatively much smaller semi-circular arcs (oxide loading~0.55
834 mg/cm², oxide-VACNT loading~1.1 mg/cm²). The large semi-circular arcs in a Nyquist plot are
835 indicative of high charge transfer resistance which would result in the reduced Faradaic response
836 in VO_x:Ti thin film oxide at higher scan rates as seen in **Figure 9c**. Pseudocapacitance
837 performance which is usually prevalent in the 10³ to 10⁴ Hz region on account of the electrode
838 kinetic rate limitations⁵ is much more prominent in the VO_x:Ti-VACNT as compared to the
839 VO_x:Ti.



840



841

842 **Figure 12:** (a) Impedance behavior of VO_x:Ti-VACNT (oxide loading~0.55 mg/cm², oxide-
843 VACNT loading~1.1 mg/cm²) and (b) VO_x:Ti.

844 Equivalent Circuit modeling was performed using the Z-view software to obtain charge-
845 transfer parameters of the titanium doped vanadium oxide thin film and titanium doped
846 vanadium oxide-VACNT composite electrodes. An under-potential deposition mechanism
847 described by Conway et. al.^{5, 6} was used to model the pseudocapacitance reactions at both sets of
848 electrodes. The mechanism considers both double layer capacitance (usually modeled as a
849 capacitor (C_{dl}) in parallel with a resistor (R_F)) and pseudocapacitance arising from surface
850 reactions (C_p) on the electrode of interest. In addition, an electrochemical series resistance (R_s)
851 common to all electrochemical systems is also considered. However, in porous electrode
852 systems, the capacitor elements are usually modeled as constant phase elements^{7, 132}. Constant
853 phase elements (CPE) are usually used to fit data having depressed semi-circular arcs on account
854 of electrode porosity and roughness which leads to fractal character and inhomogeneity in
855 response on account of presence of different crystallographic reaction sites. On account of such
856 behavior noticeable in **Figures 12a and 12b**, we have used constant phase elements, CPE-dl and
857 CPE-p to fit our data. The values of various charge transfer parameters obtained by fitting have
858 been summarized in **Table 6**.

859 **Table 6:** Charge-transfer properties of the CVD derived titanium doped vanadium oxide
860 materials.

Voltage	Series Resistance (R _s)-Ωcm ²		Double layer Constant Phase Element (CPE-dl)		Pseudocapacitance Constant phase element (CPE-F)		Faradaic resistance (R _F)-Ωcm ²	
	Ti-V oxid e-	Ti-V oxide thin	Ti-V oxide-VACNT composite	Ti-V oxide thin film	Ti-V oxide-VACNT composite	Ti-V oxide thin film	Ti-V oxide-VACNT	Ti-V oxide thin film

	VAC NT com posit e	film	film				composit e					
			T*10 ⁴	φ								
-0.2	2.46	1.67	6.14	0.86	4.22	0.60	23.37	0.71	0.24	0.98	0.86	9.68
0.1	2.51	1.7	3	0.98	0.04	0.96	19	0.77	1.17	0.72	0.94	1.06
0.2	2.51	1.69	6	0.90	0.05	0.94	16	0.77	1.01	0.73	1.44	1.09
0.5	2.46	2.01	23	0.74	0.57	0.84	56.95	0.71	1.11	0.75	24.06	37.83

861

862 It can be seen that the series resistance (R_s) is fairly constant in both systems and varies
863 between 2 and 3 Ωcm^2 which is common in carbon based systems¹³³. The most notable
864 difference in behavior between $\text{VO}_x:\text{Ti}$ and the $\text{VO}_x:\text{Ti-VACNT}$ composite is the charge transfer
865 resistance R_F . R_F values are reduced by up to an order on account of the presence of VACNTs.
866 This can also be seen in **Figure 12a** where the oxide nanosphere-VACNT heterostructure based
867 material has an almost 90 degree slope at low frequencies as compared to large semi-circles
868 indicating high R_F values seen in **Figure 12b**¹³⁴. R_F is the charge transfer resistance for electron
869 tunneling from the electrode to the electrolyte i.e. the Faradaic pseudocapacitance process. R_F
870 can usually be correlated with exchanged current density through the expression⁵:

$$871 \quad R_F \text{ (in ohms)} = \frac{RT}{i_0 F} \quad \text{iii)}$$

872 where

873 R = Gas constant (J/mol-K)

874 T = Temperature (in Kelvin)

875 i_0 = Exchange current density (in Amperes)

876 F = Faraday constant (Coulomb/mol)

877 R_F as defined in **Equation iii** is the Faradaic charge-transfer resistance at reversible
 878 potential. A decrease in charge-transfer resistance thus indicates an increase in exchange current
 879 density resulting in increased rates of the Faradaic reactions. The value of R_F as defined in
 880 Equation iii is limited to the reversible potential of the particular reaction. Therefore, in addition
 881 to the above conclusion, one can understand that the Faradaic leakage current density and thus
 882 Faradaic charge-transfer resistance are overvoltage dependent⁵. The dependence of R_F on voltage
 883 can be seen in **Table 7**. **Table 7** compares the charge transfer resistance of the VO_x :Ti-VACNT
 884 with that of VO_x :Ti at different voltages and also to that of undoped VO_x thin film material
 885 prepared using a similar approach reported earlier¹⁷. The charge transfer resistance of VO_x :Ti –
 886 VACNT is less than or equal to that of VO_x :Ti at all the potentials indicating that electron
 887 tunneling across the electrochemical interface is more easily activated in the case of VO_x :Ti-
 888 VACNT than in the case of VO_x :Ti. The improvement in pseudocapacitance behavior can also
 889 be inferred by referring to the quantitative values of the constant phase elements (CPE-F) seen in
 890 **Table 6**. It can also be seen that there is an almost one order improvement in Faradaic charge-
 891 transfer resistance upon doping VO_x thin film material with titanium. The doping of titanium into
 892 the vanadium oxide lattice thus enhances not only the electronic conductivity but also the
 893 electrochemical performance to a certain extent resulting in the improved response displayed in
 894 **Figure 9c**. The presence of carbon nanotubes further enhances both the electronic conductivity
 895 and ease of electron transfer confirming that the VO_x :Ti –VACNT is indeed an excellent
 896 supercapacitor material.

897 **Table 7:** Comparison between the Faradaic charge-transfer resistance (R_F) (Ωcm^2) of the CVD
 898 derived VO_x :Ti and VO_x ¹⁷

Voltage	VO_x :Ti-VACNT	VO_x :Ti	Voltage	VO_x
---------	------------------	------------	---------	--------

-0.2	0.86	9.68	-0.8	968.93
0.1	0.94	1.06	-0.3	58.802
0.2	1.44	1.09	0.2	58.197
0.5	24.06	37.83	0.3	67.12

899

900 **Conclusions**

901 Hot-wall chemical vapor deposition has for the first time been used to create a unique
902 morphology consisting of titanium doped vanadium oxide nanospheres on vertically aligned
903 carbon nanotubes (VACNTs). Such a CVD derived $\text{VO}_x\text{:Ti-VACNT}$ heterostructured composite
904 has been shown to exhibit excellent charge storage characteristics with capacitances of up to 313
905 F/g achieved at a scan rate of 2 mV/s particularly in a thick oxide nanosphere film coated on the
906 VACNTs. The doping of titanium into the vanadium oxide lattice has been shown to have a
907 salutary effect on the electrochemical charge storage capability of vanadium oxide due to the
908 improved electronic conductivity, an essential parameter for achieving fast electron transport and
909 charge transfer response. Another highlight of this study is the ease of electrode fabrication.
910 Chemical vapor deposition (CVD) is a method amenable for continuous production of electrodes
911 in an assembly-chain configuration in contrast with chemical methods used to produce oxides
912 which is a batch process, and in-turn also involve a second heat treatment step. Moreover, the
913 need for slurry preparation is eliminated and binder-free electrodes are thus generated. The
914 formation of a thick film with high loadings (upto 1.1 mg/cm²) by a chemical vapor deposition
915 method allows for tailoring the morphology of the oxide, while also maintaining sufficient active
916 materials resulting in high areal capacitance of upto ~350 mF/cm². The CVD method could

917 further be optimized to lead to electrodes with superior rate capability by tailoring the nature of
918 the CNT-oxide interface by additional functionalization to attain better interface interaction and
919 consequently, more improved electron transport. In addition, a complete survey of the titanium
920 doped vanadium oxide capacitor behavior spectrum is in order and efforts are on to achieve the
921 same, the results of which will be reported in subsequent manuscripts.

922 **Acknowledgement**

923 Research supported by the National Science Foundation, under Award CBET-0933141. PNK
924 acknowledges the Edward R. Weidlein Chair Professorship funds and the Center for Complex Engineered
925 Multifunctional Materials (CCEMM) for procuring the electrochemical equipment used in this research
926 work. The authors would like to acknowledge the support of Dr. Susheng Tan, Nanoscale Fabrication and
927 Characterization Facility at the University of Pittsburgh for help with Transmission Electron Microscopy
928 data collection and analyses.

929 **References**

- 930 1. H. Chen, T. N. Cong, W. Yang, C. Tan, Y. Li and Y. Ding, *Progress in Natural Science*,
931 2009, 19, 291-312.
- 932 2. P. Jampani, A. Manivannan and P. N. Kumta, *Electrochemical Society Interface*, 2010,
933 19, 57-62.
- 934 3. U. S. A. B. Consortium, Department of Energy, 2006.
- 935 4. P. Simon and Y. Gogotsi, *Nature Materials*, 2008, 7, 845-854.
- 936 5. B. E. Conway, *Electrochemical Supercapacitors: Scientific Fundamentals and*
937 *Technological Applications*, Kluwer Academic/Plenum Publishers, New York, 1999.
- 938 6. B. E. Conway, *J. Electrochem. Soc.*, 1991, 138, 1539-1548.
- 939 7. P. S. Germain, W. G. Pell and B. E. Conway, *Electrochim. Acta*, 2004, 49, 1775-1788.
- 940 8. C. C. Hu and T. W. Tsou, *Electrochemistry Communications*, 2002, 4, 105-109.
- 941 9. R. Kötz and M. Carlen, *Electrochimica Acta*, 2000, 45, 2483-2498.
- 942 10. J. M. Miller, B. Dunn, T. D. Tran and R. W. Pekala, *J. Electrochem. Soc.*, 1997, 144,
943 L309-L311.
- 944 11. M. Toupin, T. Brousse and D. Belanger, *Chemistry of Materials*, 2002, 14, 3946-3952.
- 945 12. J. W. Long, K. E. Swider, C. I. Merzbacher and D. R. Rolison, *Langmuir*, 1999, 15, 780-
946 785.
- 947 13. T. Kudo, Y. Ikeda, T. Watanabe, M. Hibino, M. Miyayama, H. Abe and K. Kajita, *Solid*
948 *State Ionics*, 2002, 152, 833-841.
- 949 14. I. H. Kim, J. H. Kim, B. W. Cho, Y. H. Lee and K. B. Kim, *Journal of the*
950 *Electrochemical Society*, 2006, 153, A989-A996.

- 951 15. S. D. Perera, B. Patel, J. Bonso, M. Grunewald, J. P. Ferraris and K. J. Balkus, *Acs*
952 *Applied Materials & Interfaces*, 2011, 3, 4512-4517.
- 953 16. S. Tepavcevic, H. Xiong, V. R. Stamenkovic, X. Zuo, M. Balasubramanian, V. B.
954 Prakash, C. S. Johnson and T. Rajh, *ACS Nano*, 2011, 6, 530-538.
- 955 17. P. H. Jampani, K. Kadakia, D. H. Hong, R. Epur, J. A. Poston, A. Manivannan and P. N.
956 Kumta, *Journal of The Electrochemical Society*, 2013, 160, A1118-A1127.
- 957 18. W. Wang and P. N. Kumta, *Journal of Power Sources*, 2007, 172, 650-658.
- 958 19. I. Shakir, Z. Ali, J. Bae, J. Park and D. J. Kang, *Nanoscale*, 2014, 6, 4125-4130.
- 959 20. M. Lee, B. H. Wee and J. D. Hong, *Advanced Energy Materials*, 2014.
- 960 21. H. Wang, H. Yi, X. Chen and X. Wang, *Journal of Materials Chemistry A*, 2014, 2,
961 1165-1173.
- 962 22. M. E. Spahr, P. Bitterli, R. Nesper, M. Muller, F. Krumeich and H. U. Nissen,
963 *Angewandte Chemie-International Edition*, 1998, 37, 1263-1265.
- 964 23. M. E. Spahr, P. Stoschitzki-Bitterli, R. Nesper, O. Haas and P. Novak, *Journal of the*
965 *Electrochemical Society*, 1999, 146, 2780-2783.
- 966 24. C. C. Hu, C. M. Huang and K. H. Chang, *Journal of Power Sources*, 2008, 185, 1594-
967 1597.
- 968 25. C. M. Huang, C. C. Hu, K. H. Chang, J. M. Li and Y. F. Li, *J. Electrochem. Soc.*, 2009,
969 156, A667-A671.
- 970 26. J. M. Li, K. H. Chang and C. C. Hu, *Electrochemistry Communications*, 2010, 12, 1800-
971 1803.
- 972 27. J. M. Li, K. H. Chang, T. H. Wu and C. C. Hu, *Journal of Power Sources*, 2013, 224, 59-
973 65.
- 974 28. M. E. Franke, T. J. Koplín and U. Simon, *Small*, 2006, 2, 36-50.
- 975 29. G. Korotcenkov, *Sensors and Actuators B: Chemical*, 2005, 107, 209-232.
- 976 30. Y. Wei, C.-W. Ryu and K.-B. Kim, *Journal of Power Sources*, 2007, 165, 386-392.
- 977 31. D. Hara, H. Ikuta, Y. Uchimoto and M. Wakihara, *Journal of Materials Chemistry*, 2002,
978 12, 2507-2512.
- 979 32. H. K. Park, *Solid State Ionics*, 2005, 176, 307-312.
- 980 33. L. Q. Mai, W. Chen, Q. Xu, J. F. Peng and Q. Y. Zhu, *Chem. Phys. Lett.*, 2003, 382, 307-
981 312.
- 982 34. K. Kadakia, M. K. Datta, P. H. Jampani, S. K. Park and P. N. Kumta, *Journal of Power*
983 *Sources*, 2013, 222, 313-317.
- 984 35. K. Kadakia, M. K. Datta, O. I. Velikokhatnyi, P. Jampani, S. K. Park, P. Saha, J. A.
985 Poston, A. Manivannan and P. N. Kumta, *International Journal of Hydrogen Energy*,
986 2012, 37, 3001-3013.
- 987 36. F. Coustier, J. Hill, B. B. Owens, S. Passerini and W. H. Smyrl, *J. Electrochem. Soc.*,
988 1999, 146, 1355-1360.
- 989 37. F. Coustier, S. Passerini and W. H. Smyrl, *Solid State Ionics*, 1997, 100, 247-258.
- 990 38. Y. Takasu, S. Mizutani, M. Kumagai, S. Sawaguchi and Y. Murakami, *Electrochemical*
991 *and Solid State Letters*, 1999, 2, 1-2.
- 992 39. E. Frackowiak and F. Béguin, *Carbon*, 2001, 39, 937-950.
- 993 40. E. Frackowiak, K. Metenier, V. Bertagna and F. Béguin, *Appl. Phys. Lett.*, 2000, 77,
994 2421-2423.

- 995 41. P. Kurzweil, in *Encyclopedia of Electrochemical Power Sources*, ed. G. Editor-in-
996 Chief: Jürgen, Elsevier, Amsterdam, 2009, DOI: 10.1016/b978-044452745-5.00352-x,
997 pp. 634-648.
- 998 42. S. Stankovich, D. A. Dikin, G. H. B. Dommett, K. M. Kohlhaas, E. J. Zimney, E. A.
999 Stach, R. D. Piner, S. T. Nguyen and R. S. Ruoff, *Nature*, 2006, 442, 282-286.
- 1000 43. S. Stankovich, D. A. Dikin, R. D. Piner, K. A. Kohlhaas, A. Kleinhammes, Y. Jia, Y.
1001 Wu, S. T. Nguyen and R. S. Ruoff, *Carbon*, 2007, 45, 1558-1565.
- 1002 44. A. K. Geim and K. S. Novoselov, *Nature Materials*, 2007, 6, 183-191.
- 1003 45. K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V.
1004 Grigorieva and A. A. Firsov, *Science*, 2004, 306, 666-669.
- 1005 46. S. Murali, J. R. Potts, S. Stoller, J. Park, M. D. Stoller, L. L. Zhang, Y. Zhu and R. S.
1006 Ruoff, *Carbon*, 2012, 50, 3482-3485.
- 1007 47. Y. Zhu, S. Murali, M. D. Stoller, K. J. Ganesh, W. Cai, P. J. Ferreira, A. Pirkle, R. M.
1008 Wallace, K. A. Cychosz, M. Thommes, D. Su, E. A. Stach and R. S. Ruoff, *Science*,
1009 2011, 332, 1537-1541.
- 1010 48. Z. Chen, V. Augustyn, X. L. Jia, Q. F. Xiao, B. Dunn and Y. F. Lu, *ACS Nano*, 2012, 6,
1011 4319-4327.
- 1012 49. Q. W. Li, Y. Li, X. F. Zhang, S. B. Chikkannanavar, Y. H. Zhao, A. M. Dangelewicz, L.
1013 X. Zheng, S. K. Doorn, Q. X. Jia, D. E. Peterson, P. N. Arendt and Y. T. Zhu, *Advanced*
1014 *Materials*, 2007, 19, 3358-3363.
- 1015 50. P. L. McEuen, M. S. Fuhrer and H. K. Park, *IEEE Trans. Nanotechnol.*, 2002, 1, 78-85.
- 1016 51. T. W. Odom, J.-L. Huang, P. Kim and C. M. Lieber, *Nature*, 1998, 391, 62-64.
- 1017 52. P. Ajayan, *Chemical reviews*, 1999, 99, 1787-1800.
- 1018 53. J. W. G. Wilder, L. C. Venema, A. G. Rinzler, R. E. Smalley and C. Dekker, *Nature*,
1019 1998, 391, 59-62.
- 1020 54. T. Szorenyi, K. Bali and I. Hevesi, *Journal of Non-Crystalline Solids*, 1980, 35-6, 1245-
1021 1248.
- 1022 55. S. D. Perera, B. Patel, N. Nijem, K. Roodenko, O. Seitz, J. P. Ferraris, Y. J. Chabal and
1023 K. J. Balkus, *Advanced Energy Materials*, 2011, 1, 936-945.
- 1024 56. C. Y. Lee, H. M. Tsai, H. J. Chuang, S. Y. Li, P. Lin and T. Y. Tseng, *Journal of the*
1025 *Electrochemical Society*, 2005, 152, A716-A720.
- 1026 57. A. L. M. Reddy, M. M. Shaijumon, S. R. Gowda and P. M. Ajayan, *Nano Letters*, 2009,
1027 9, 1002-1006.
- 1028 58. S. R. Sivakkumar, J. M. Ko, D. Y. Kim, B. C. Kim and G. G. Wallace, *Electrochimica*
1029 *Acta*, 2007, 52, 7377-7385.
- 1030 59. Z. Fan, J. Chen, K. Cui, F. Sun, Y. Xu and Y. Kuang, *Electrochimica Acta*, 2007, 52,
1031 2959-2965.
- 1032 60. I. H. Kim, J. H. Kim and K. B. Kim, *Electrochemical and Solid State Letters*, 2005, 8,
1033 A369-A372.
- 1034 61. I. H. Kim, J. H. Kim, Y. H. Lee and K. B. Kim, *Journal of the Electrochemical Society*,
1035 2005, 152, A2170-A2178.
- 1036 62. S. D. Perera, B. Patel, N. Nijem, K. Roodenko, O. Seitz, J. P. Ferraris, Y. J. Chabal and
1037 K. J. Balkus, *Advanced Energy Materials*, 2011, 1, 936-945.
- 1038 63. M. Sathiyaa, A. S. Prakash, K. Ramesha, J. M. Tarascon and A. K. Shukla, *J. Am. Chem.*
1039 *Soc.*, 2011, 133, 16291-16299.

- 1040 64. Z. Chen, Y. C. Qin, D. Weng, Q. F. Xiao, Y. T. Peng, X. L. Wang, H. X. Li, F. Wei and
1041 Y. F. Lu, *Adv. Funct. Mater.*, 2009, 19, 3420-3426.
- 1042 65. D. Choi, G. E. Blomgren and P. N. Kumta, *Advanced Materials*, 2006, 18, 1178-+.
- 1043 66. Y. Gogotsi and P. Simon, *Science*, 2011, 334, 917-918.
- 1044 67. J. C. Badot and N. Baffier, *Journal of Materials Chemistry*, 1992, 2, 1167-1174.
- 1045 68. Y. Wang and G. Cao, *Chemistry of Materials*, 2006, 18, 2787-2804.
- 1046 69. S. Boukhalifa, K. Evanoff and G. Yushin, *Energy & Environmental Science*, 2012, 5,
1047 6872-6879.
- 1048 70. N. Ozer, S. Sabuncu and J. Cronin, *Thin Solid Films*, 1999, 338, 201-206.
- 1049 71. W. Wang, R. Epur and P. N. Kumta, *Electrochemistry Communications*, 2011, 13, 429-
1050 432.
- 1051 72. W. Wang and P. N. Kumta, *ACS Nano*, 2010, 4, 2233-2241.
- 1052 73. D. Vernardou, P. Paterakis, H. Drosos, E. Spanakis, I. M. Povey, M. E. Pemble, E.
1053 Koudoumas and N. Katsarakis, *Solar Energy Materials and Solar Cells*, 2011, 95, 2842-
1054 2847.
- 1055 74. L. Michailovits, K. Bali, T. Szorenyi and I. Hevesi, *Acta Physica Academiae Scientiarum
1056 Hungaricae*, 1980, 49, 217-221.
- 1057 75. R. Binions, G. Hyett, C. Piccirillo and I. P. Parkin, *Journal of Materials Chemistry*, 2007,
1058 17, 4652-4660.
- 1059 76. T. D. Manning and I. P. Parkin, *Polyhedron*, 2004, 23, 3087-3095.
- 1060 77. M. N. Field and I. P. Parkin, *Journal of Materials Chemistry*, 2000, 10, 1863-1866.
- 1061 78. T. D. Manning, I. P. Parkin, R. J. H. Clark, D. Sheel, M. E. Pemble and D. Vernadou,
1062 *Journal of Materials Chemistry*, 2002, 12, 2936-2939.
- 1063 79. I. P. Parkin and G. S. Elwin, *Journal of Materials Chemistry*, 2001, 11, 3120-3124.
- 1064 80. S. K. Pal, S. Talapatra, S. Kar, L. Ci, R. Vajtai, T. Borca-Tasciuc, L. S. Schadler and P.
1065 M. Ajayan, *Nanotechnology*, 2008, 19.
- 1066 81. S. Talapatra, S. Kar, S. K. Pal, R. Vajtai, L. Ci, P. Victor, M. M. Shaijumon, S. Kaur, O.
1067 Nalamasu and P. M. Ajayan, *Nature Nanotechnology*, 2006, 1, 112-116.
- 1068 82. T. D. Manning, I. P. Parkin, M. E. Pemble, D. Sheel and D. Vernardou, *Chem. Mat.*,
1069 2004, 16, 744-749.
- 1070 83. C. S. Blackman, C. Piccirillo, R. Binions and I. P. Parkin, *Thin Solid Films*, 2009, 517,
1071 4565-4570.
- 1072 84. D. Vernardou, M. E. Pemble and D. W. Sheel, *Chemical Vapor Deposition*, 2007, 13,
1073 158-162.
- 1074 85. C. Emmenegger, J. M. Bonard, P. Mauron, P. Sudan, A. Lepora, B. Grobety, A. Zuttel
1075 and L. Schlapbach, *Carbon*, 2003, 41, 539-547.
- 1076 86. G. Atthipalli, R. Epur, P. N. Kumta, M. Yang, J. K. Lee and J. L. Gray, *J. Phys. Chem. C*,
1077 2011, 115, 3534-3538.
- 1078 87. M. Sathiya, A. S. Prakash, K. Ramesha, J. M. Tarascon and A. K. Shukla, *Journal of the
1079 American Chemical Society*, 2011, 133, 16291-16299.
- 1080 88. E. Hryha, E. Rutqvist and L. Nyborg, *Surface and Interface Analysis*, 2012, 44, 1022-
1081 1025.
- 1082 89. A. V. Naumkin, A. Kraut-Vass, S. W. Gaarenstroom and C. J. Powell, in *NIST Standard
1083 Reference database 20*, National Institute of Standards & Technology, 2012, vol. Version
1084 4.1.

- 1085 90. N. I. S. Testing,
1086 [http://srdata.nist.gov/xps/elm_Spectra_query.aspx?Elm1=V&LD1=2p1%2f2&Elm2=V&](http://srdata.nist.gov/xps/elm_Spectra_query.aspx?Elm1=V&LD1=2p1%2f2&Elm2=V&LD2=2p3%2f2&Elm3=&LD3=&Elm4=&LD4=&sType=PE)
1087 [LD2=2p3%2f2&Elm3=&LD3=&Elm4=&LD4=&sType=PE](http://srdata.nist.gov/xps/elm_Spectra_query.aspx?Elm1=V&LD1=2p1%2f2&Elm2=V&LD2=2p3%2f2&Elm3=&LD3=&Elm4=&LD4=&sType=PE), Accessed 11/24/2014,
1088 2014.
- 1089 91. V. L. Volkov, G. S. Zakharova, E. G. Volkova and M. V. Kuznetsov, *Russ. J. Inorg.*
1090 *Chem.*, 2006, 51, 847-851.
- 1091 92. A. C. Ferrari and J. Robertson, *Phys. Rev. B*, 2000, 61, 14095-14107.
- 1092 93. A. C. Ferrari and J. Robertson, *Phys. Rev. B*, 2001, 64.
- 1093 94. E. Cazzanelli, G. Mariotto, S. Passerini, W. H. Smyrl and A. Gorenstein, *Solar Energy*
1094 *Materials and Solar Cells*, 1999, 56, 249-258.
- 1095 95. S. Santangelo, G. Messina, G. Faggio, M. G. Willinger, N. Pinna, A. Donato, A. Arena,
1096 N. Donato and G. Neri, *Diamond and Related Materials*, 2010, 19, 590-594.
- 1097 96. A. Löfberg, T. Giornelli, S. Paul and E. Bordes-Richard, *Applied Catalysis A: General*,
1098 2011, 391, 43-51.
- 1099 97. M. G. Nobbenhuis, A. Baiker, P. Barnickel and A. Wokaun, *Applied Catalysis A:*
1100 *General*, 1992, 85, 157-172.
- 1101 98. H. Liu, Y. Wu and J. Zhang, *Acs Applied Materials & Interfaces*, 2011, 3, 1757-1764.
- 1102 99. V. Stengl, D. Popelkova and P. Vlacil, *J. Phys. Chem. C*, 2011, 115, 25209-25218.
- 1103 100. G. N. Kryukova, G. A. Zenkovets, G. Mestl and R. Schlogl, *Reaction Kinetics and*
1104 *Catalysis Letters*, 2003, 80, 161-169.
- 1105 101. G. T. Went, S. T. Oyama and A. T. Bell, *J. Phys. Chem.*, 1990, 94, 4240-4246.
- 1106 102. N. Y. Topsoe, M. Anstrom and J. A. Dumesic, *Catal. Lett.*, 2001, 76, 11-20.
- 1107 103. J. B. MacChesney and H. J. Guggenheim, *Journal of Physics and Chemistry of Solids*,
1108 1969, 30, 225-234.
- 1109 104. J. Muster, G. T. Kim, V. Krstic, J. G. Park, Y. W. Park, S. Roth and M. Burghard,
1110 *Advanced Materials*, 2000, 12, 420-+.
- 1111 105. L. A. Balagurov, I. V. Kulemanov, A. F. Orlov and E. A. Petrova, *Russ Microelectron*,
1112 2012, 41, 503-507.
- 1113 106. N. Serpone, D. Lawless, J. Disdier and J.-M. Herrmann, *Langmuir*, 1994, 10, 643-652.
- 1114 107. J. D. Butler, *Transactions of the Faraday Society*, 1960, 56, 1842-1850.
- 1115 108. G. Kresse and J. Furthmüller, *Physical Review B*, 1996, 54, 11169-11186.
- 1116 109. G. Kresse and J. Furthmüller, *Computational Materials Science*, 1996, 6, 15-50.
- 1117 110. G. Kresse and D. Joubert, *Physical Review B*, 1999, 59, 1758-1775.
- 1118 111. J. P. Perdew and W. Yue, *Physical Review B*, 1986, 33, 8800-8802.
- 1119 112. R. Enjalbert and J. Galy, *Acta Crystallographica Section C*, 1986, 42, 1467-1469.
- 1120 113. S. Laubach, P. C. Schmidt, Thi, F. J. Fernandez-Madrigal, Q.-H. Wu, W. Jaegermann, M.
1121 Klemm and S. Horn, *Physical Chemistry Chemical Physics*, 2007, 9, 2564-2576.
- 1122 114. V. Eyert and K. H. Höck, *Physical Review B*, 1998, 57, 12727-12737.
- 1123 115. A. Chakrabarti, K. Hermann, R. Druzinic, M. Witko, F. Wagner and M. Petersen,
1124 *Physical Review B*, 1999, 59, 10583-10590.
- 1125 116. D. R. Lide, *CRC Handbook of chemistry and physics, 86th edn, 2005-2006*, CRC Press,
1126 Boca Raton, 2005.
- 1127 117. M. Pourbaix, *Atlas of electrochemical equilibria in aqueous solutions*, Pergamon Press,
1128 Oxford, 1966.
- 1129 118. N. Takeno, *Atlas of Eh-pH diagrams*, National Institute of Advanced Industrial Science
1130 & Technology, 2005.

- 1131 119. S. Zhang, Y. Li, C. Wu, F. Zheng and Y. Xie, *The Journal of Physical Chemistry C*,
1132 2009, 113, 15058-15067.
- 1133 120. H. Kim and B. N. Popov, *Journal of Power Sources*, 2002, 104, 52-61.
- 1134 121. G. Zheng, L. Hu, H. Wu, X. Xie and Y. Cui, *Energy & Environmental Science*, 2011, 4,
1135 3368-3373.
- 1136 122. M. Zhang, M. Yudasaka, S. Bandow and S. Iijima, *Chem. Phys. Lett.*, 2003, 369, 680-
1137 683.
- 1138 123. H. Q. Yang, D. P. Li, S. Han, N. Li and B. X. Lin, *Journal of Power Sources*, 1996, 58,
1139 221-224.
- 1140 124. D. Choi, Doctor of Philosophy, Carnegie Mellon University, 2005.
- 1141 125. M.-G. Willinger, G. Neri, A. Bonavita, G. Micali, E. Rauwel, T. Hertrich and N. Pinna,
1142 *Physical Chemistry Chemical Physics*, 2009, 11, 3615-3622.
- 1143 126. L. Niinisto, M. Ritala and M. Leskela, *Mater. Sci. Eng. B-Solid State Mater. Adv.*
1144 *Technol.*, 1996, 41, 23-29.
- 1145 127. A. Ponzoni, E. Comini, I. Concina, M. Ferroni, M. Falasconi, E. Gobbi, V. Sberveglieri
1146 and G. Sberveglieri, *Sensors*, 2012, 12, 17023-17045.
- 1147 128. J. M. Sieben, E. Morallón and D. Cazorla-Amorós, *Energy*, 2013, 58, 519-526.
- 1148 129. X. Liu and P. G. Pickup, *Journal of Power Sources*, 2008, 176, 410-416.
- 1149 130. N. Nagarajan, M. Cheong and I. Zhitomirsky, *Materials Chemistry and Physics*, 2007,
1150 103, 47-53.
- 1151 131. Y. J. Kang, H. Chung and W. Kim, *Synthetic Metals*, 2013, 166, 40-44.
- 1152 132. G. J. Brug, A. L. G. Vandeneeden, M. Sluytersrehabach and J. H. Sluyters, *Journal of*
1153 *Electroanalytical Chemistry*, 1984, 176, 275-295.
- 1154 133. C. Portet, G. Yushin and Y. Gogotsi, *Carbon*, 2007, 45, 2511-2518.
- 1155 134. V. Subramanian, H. W. Zhu and B. Q. Wei, *Electrochemistry Communications*, 2006, 8,
1156 827-832.

1157

1158