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# Hierarchically structured TiO<sub>2</sub> for Ba-filled skutterudite with enhanced thermoelectric performance

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## Abstract

The influence of hierarchical TiO<sub>2</sub> inclusions on electrical and thermal transport properties has been investigated in the Ba-filled skutterudite compound Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> synthesized by ball milling followed by hot pressing. The hierarchical TiO<sub>2</sub> with morphology of nanocrystallite aggregates were prepared using a carbon sphere-templated method. It was found that TiO<sub>2</sub> inclusions strongly enhance the Seebeck coefficient while they weakly degrade the electric conductivity. In addition, TiO<sub>2</sub> inclusions significantly reduce the lattice thermal conductivity through all-scale length phonon scattering due to these hierarchically-structured TiO<sub>2</sub> particle inclusions. *ZT* values up to 1.2 were obtained for Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> with the optimal amount and dispersion of TiO<sub>2</sub> inclusions. These observations demonstrate an exciting scientific opportunity to raise the figure-of-merit of filled skutterudites via hierarchical dispersion.

**Keywords:** Filled skutterudites; Hierarchical TiO<sub>2</sub>; Phonon scattering; Thermoelectric properties; Nanocomposite.

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## Introduction

The world we are living in has been stressed by the rapidly growing energy consumption and associated with it environmental problems. To solve these problems, or at least mitigating their influence, the new green energy policies and the more efficient ways we use the energy, should be considered. Thermoelectric (TE) materials are a very good candidate to achieve a more efficient usage of energy, especially in the area of waste industrial heat<sup>1-7</sup>. The figure of merit,  $ZT$ , defined as  $ZT=S^2\sigma T/\kappa$ , where  $S$  is the Seebeck coefficient,  $\sigma$  is the electrical conductivity,  $\kappa$  is the thermal conductivity, and  $T$  is the absolute temperature, is the parameter which describes the performance of a TE material. The electronic performance of a thermoelectric material is governed by its power factor  $P=S^2\sigma$ . The challenge to develop efficient TE materials is to achieve a simultaneous enhancement in the power factor while maintaining as low thermal conductivity as possible.

Among promising thermoelectric materials, filled skutterudites based on  $\text{CoSb}_3$  have attracted wide attention not only because of their promising dimensionless figure-of-merit, but also because of their very good mechanical properties<sup>8-18</sup>. Current strategies to improve the figure of merit of  $\text{CoSb}_3$ -based filled skutterudites are focusing on maximizing the power factor through substitutional doping, optimal filling, and reducing the thermal conductivity using mass fluctuation phonon scattering *via* solid-solution alloying. A higher  $ZT$  of 1.7 at 850 K, which is the record figure of merit to date for  $\text{CoSb}_3$ -based skutterudites, was realized with multiple co-fillers Ba, La, and Yb<sup>10</sup>.

Another promising approach is to form bulk nanocomposites where various forms of nanoinclusions are dispersed in the  $\text{CoSb}_3$ -based matrix. Provided the phonon mean-free

path is reduced by enhanced phonon boundary scattering to a greater degree than any reduction in the mean-free path of charge carriers, the net effect is an improvement in the figure of merit. Moreover, there is also a possibility that the presence of dispersed nanoinclusions will lead to an enhancement in the Seebeck coefficient as a result of energy-dependent scattering of electrons at the interface of nanoinclusions and the matrix<sup>19-24</sup>. Such nanostructured composite materials are fabricated using bulk processing routes rather than nanofabrication means. As such, the synthesis is scalable and compatible with industrial processing<sup>17</sup>. Different kinds of nanoinclusions have been embedded in the skutterudite matrix, among them  $\text{TiO}_2$ <sup>20</sup>,  $\text{ZrO}_2$ <sup>21</sup>,  $\text{NiSb}$ <sup>22</sup>,  $\text{Yb}_2\text{O}_3$ <sup>23</sup>,  $\text{Ag}$ <sup>17</sup>,  $\text{C}_{60}$ <sup>24</sup>. They were added to the matrix material *via* mechanical mixing or by *in-situ* reaction methods. However, impurities such as oxygen brought by ball milling, exert a negative impact on the electronic transport properties. While the contamination problem can be substantially avoided by using *in-situ* reaction methods, the control of composition, grain size and distribution of these *in-situ* nanoinclusions is, however, a great challenge. Nevertheless, in view of the fact that the currently achievable thermal conductivity reduction in filled skutterudites is still far from the theoretical limit value of  $\sim 0.3$  W/m-K, further attempts to reduce the thermal conductivity of skutterudites represent worthwhile efforts.

Because the heat in solids is carried by phonons with a rather broad range of wavelengths<sup>25</sup>, it is believed that hierarchically structured thermoelectric materials possessing multi-scale microstructures will have the best chance of greatly reducing the lattice thermal conductivity and thus improving the overall thermoelectric performance. This has been demonstrated with n-type Na-doped PbTe containing  $\text{MgTe}$ <sup>26</sup> and p-type

Na-doped PbSe containing CdS<sup>27</sup>. Here Na doping generates point defects, the second phases form ubiquitous endotaxial nanostructures dispersed in the respective matrices, and the mesoscale microstructure is achieved by spark plasma sintering. Such all-scale hierarchical structure scatters a broad range of phonons leading to a very low (lattice) thermal conductivity.

In this paper we report on a significantly improved performance of Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> achieved by adding hierarchically structured TiO<sub>2</sub> inclusions. This is accomplished because (a) the Seebeck coefficient is enhanced due to the energy filtering effect whereby the low energy electrons are scattered by the interface of inclusions and the matrix; (b) the thermal conductivity is greatly reduced by enhanced scattering of a broad range of phonons on hierarchically structured TiO<sub>2</sub> inclusions. These two factors boost the *ZT* of Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> up to 1.2 at 823K. Our results show clearly that optimal dispersion of the second phase can greatly improve thermoelectric performance.

## Experiment

TiO<sub>2</sub> nanocrystallite aggregates were prepared using a carbon sphere-templated method. In a typical fabrication, a precursor sol of TiO<sub>2</sub> was prepared first by dissolving 3.75 mL of titanium isopropoxide in an acidic aqueous solution that contained 2.25 mL of HNO<sub>3</sub> and 7.5 mL of C<sub>2</sub>H<sub>4</sub>O<sub>2</sub> in 15 mL of DI-water. To the as-prepared TiO<sub>2</sub> sol, 5.25 g of sucrose, 1.5 g of cetyltrimethylammonium bromide (CTAB), and 3.75 mL of ethanol were added to obtain a mixture of TiO<sub>2</sub> sol and sucrose plus surfactant. The mixture was then transferred into an autoclave for hydrothermal growth at 225 °C for 10 h. When the temperature of the autoclave cooled down to room temperature naturally, the precipitate was collected and washed with water and ethanol for several times through a centrifuge

method. The product was dried at 100 °C and the as-received powder was then sintered at 450 °C for 5 h, leading to the formation of TiO<sub>2</sub> nanocrystallite aggregates.

Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> was synthesized *via* a solid state reaction followed by ball milling and hot pressing, similar to the process reported in our previous work<sup>17</sup>. Ba (pieces 99.9%), Co (powder 99.5%), and Sb (shot 99.999%) which were weighed and placed inside carbon coated quartz ampoules sealed under a vacuum. The sealed ampoules were heated to 1373 K at a rate not exceeding 5 K/min and kept at that temperature for 5 hours, and then quenched in a supersaturated salt water solution. The resultant samples were further annealed at 1023 K for 7 days. A Retsch P100 ball milling machine was subsequently used to mill the ingots at 300 rpm for 6-8 hours to form the Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> powder matrix. Two forms of TiO<sub>2</sub> particles were dispersed into the matrix at 300 rpm for 15 min using a ball milling process. One is TiO<sub>2</sub> nanocrystallite aggregates prepared by a carbon sphere-templated method. The other is commercially available TiO<sub>2</sub> nanoparticles (Degussa Aeroxide P25), which consist of ca. 30% rutile and 70% anatase and were used without any purification or treatment. The as-prepared mixture was finally hot pressed at 873 K for 15 minutes. The density of our skutterudite samples is around 7.4 g/cm<sup>3</sup>, larger than 95% of the theoretical density. It should be noted that (1) the low temperature and short soaking time of hot pressing are necessary to maintain the TiO<sub>2</sub> nanocrystallite aggregate's structure; (2) ball milling and hot pressing was carried out inside the glove box to prevent oxidation of the prepared powder.

High temperature electrical conductivity ( $\sigma$ ) and the Seebeck coefficient ( $S$ ) were measured on rectangular-shape samples in the interval 300 K to 823 K using a home built apparatus under the protective atmosphere of argon. Thermal diffusivity ( $\lambda$ ) was obtained

by the laser flash method (Flashline 5000 from Anter), and converted into thermal conductivity ( $\kappa$ ) using  $\kappa = \lambda C_p d$ , where  $C_p$  is the specific heat measured using a Netzsch, 404 Pegasus apparatus. Sample density  $d$  was calculated from sample dimensions and mass. We use pressure contacts and at all times the sample is under the atmosphere of Ar to prevent oxidation at elevated temperatures.

X-ray diffraction patterns for all samples were collected on powders in a Scintag X1 using Cu K $\alpha$  radiation. Electron-probe microanalysis (EPMA) with a Cameca Camebax EPMA was employed to determine the actual composition of the sintered material. The morphology and microstructure were investigated with field emission scanning electron microscopy (SEM, FEI NOVA) and transmission electron microscopy (JEOL 2100F S/TEM microscope operating at 200kV accelerating voltage).

## Results and Discussion

As shown in Figure 1(a) and (b), the hierarchical structure consists of the primary TiO<sub>2</sub> nanocrystallites with the diameter of  $12.2 \pm 0.3$  nm. However, these primary nanostructures form a secondary structure with mesoporous beads typically of  $830 \pm 40$  nm diameter, and a BET surface area of  $108.0 \text{ m}^2 \text{ g}^{-1}$ . As shown in Figure 1(c), TiO<sub>2</sub> nanocrystallite aggregates heat treated at 450°C for 5h possess pure anatase structure as documented by the peaks indexed to the anatase phase (PDF# 002-0387) without any presence of the rutile and brookite phases. As to the formation mechanism of this hierarchical TiO<sub>2</sub> nanocrystallite aggregates, sucrose is used in this method as a carbon source to generate spherical carbon through a hydrolysis process which occurs at  $\sim 180$  °C; the spherical carbon works as a template during the growth of TiO<sub>2</sub> nanoparticles, resulting in the formation of spherical aggregates of TiO<sub>2</sub> nanocrystallites. It has been

found that the size of the nanocrystallite aggregates is a function of the pH value of the precursor solution, while the size of TiO<sub>2</sub> nanocrystallites is dependent on both the concentration of TiO<sub>2</sub> sol and the temperature adopted for the hydrothermal growth. Sintering at 450 °C was purposely designed and carried out to remove carbon from the product and, meanwhile, improve the crystallinity of the anatase TiO<sub>2</sub>.

Powder XRD patterns for these TiO<sub>2</sub>-containing Ba-filled skutterudite bulk samples are shown in Figure 1(d). All major reflections are indexed to the skutterudite phase. No trace of TiO<sub>2</sub> is detected in any sample due to the small volume fraction of TiO<sub>2</sub> inclusions. However, our EPMA data (error bar is around 0.02) clearly reveal the presence of close to the nominal amount of TiO<sub>2</sub>. Regarding the Ba content, the actual amounts of Ba in each TiO<sub>2</sub>-containing sample are very similar but a bit smaller than the nominal composition. We observe no shift of peaks to higher angles, which we attribute to the relatively short ball milling time and rate that is chosen to minimize the contamination from milling media and prevent the lattice strain build-up during ball milling which might otherwise have a detrimental effect on the electronic mobility<sup>28-31</sup>.

SEM images of the fractured sintered bulk sample containing two different kinds of TiO<sub>2</sub> particles, namely 0.5 wt% P25 plus 0.5 wt% TiO<sub>2</sub> nanocrystallite aggregates are shown in Figure 2(a) and (b). The bottom left inset of Figure 2(b) displays the SEM image for Ba-filled skutterudite with 1 wt% TiO<sub>2</sub> nanocrystallite aggregates. These samples possess sharp and clean grain boundaries, which is beneficial for charge transport. In addition, some TiO<sub>2</sub> nanocrystallite aggregates and nanoparticles are observed at the grain boundaries and the surface of the skutterudite sample. Such inclusions are reported to enhance scattering of phonons with mid- to long-wavelengths, as schematically shown in

Figure 2(f), thereby providing a mechanism to significantly reduce the lattice thermal conductivity<sup>32-38</sup>. It is straightforward to identify the inclusions as the dispersed TiO<sub>2</sub> aggregates since their size and morphology are very similar to the as-prepared TiO<sub>2</sub> nanocrystallite aggregates. The preservation of the continuous hierarchical structure is clearly evident in a close-up TEM image of TiO<sub>2</sub> aggregates in Figure 2(c), which shows a single aggregate in the sintered bulk sample. SAED (Selected Area Electron Diffraction) patterns of the TiO<sub>2</sub> aggregates, Figure 2(d), demonstrate that they are composed of TiO<sub>2</sub> nanocrystallites with the anatase structure. HRTEM images show individual nanocrystallites on the scale of 10-20 nm. Unfortunately, until now, we have not been able to provide a single TEM image depicting the presence of commercially-obtained P25 together with TiO<sub>2</sub> nanocrystallite aggregates, likely due to a very small amount of TiO<sub>2</sub> inclusions we employ in our study. Judging from SEM and TEM images, the size and shape of these TiO<sub>2</sub> nanocrystallite aggregates is unchanged, instead of being enlarged after hot pressing. It should be mentioned that the preservation of the nanocrystallite aggregate structure is critical to ensure the enhanced phonon scattering even after long-term exposure to temperature where the material is expected to operate in thermoelectric modules. As such, the hot pressing parameters are not a trivial matter. The processing conditions with respect to the temperature and the soaking time are of prime concern and determine the quality of the final product.

Temperature dependence of the electrical conductivity ( $\sigma$ ), the Seebeck coefficient ( $S$ ), and the power factor ( $PF=S^2\sigma$ ) for Ba-filled skutterudites containing different dispersed forms of TiO<sub>2</sub> inclusions are plotted in Figure 3. Sample designated as Matrix is a hot-pressed ball-milled powder of Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> without TiO<sub>2</sub> inclusions. As shown in Figure

3 (a), all samples display a metallic conduction behavior. The electrical conductivity decreases with increasing temperature over the entire temperature range. In addition, TiO<sub>2</sub>-containing samples exhibit lower electrical conductivity than that of Matrix. Specifically, the electrical conductivity decreases from ~3533 S/cm to ~3199 S/cm and from 2000 S/cm to 1600 S/cm for Matrix, respectively the sample containing 1wt% TiO<sub>2</sub> nanocrystallite aggregates at room temperature and 823 K. This trend is completely opposite to the case of Ag-containing Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub><sup>17</sup>, where the electrical conductivity was significantly enhanced by the presence of Ag nanoinclusions. Such a different behavior is due to Ag and TiO<sub>2</sub> being two entirely different kinds of inclusions, namely, Ag is a metal with the highest electrical conductivity while TiO<sub>2</sub> is a wide band gap semiconductor. Ag nanoparticles, dispersed on surfaces of the grains and at grain boundaries, promote electrical connectivity at grain boundaries and thus improve electronic transport. In contrast, TiO<sub>2</sub> is electrically inert at temperatures of the experiment and merely scatters charge carriers, leading to a slight degradation of the electrical conductivity. It is interesting to note that the skutterudite sample containing the commercial form of TiO<sub>2</sub> (P25) has a lower electrical conductivity than the sample containing an equal amount of prepared TiO<sub>2</sub> aggregates. A possible reason for this difference could be the fact that TiO<sub>2</sub> aggregates represent the pure anatase phase while the commercially obtained TiO<sub>2</sub> contains 30% of the rutile phase. It is considered that anatase phase is more desirable for charge transport, allowing for the improvement of electrical conductivity in the TiO<sub>2</sub> aggregates-containing skutterudites<sup>39</sup>.

As shown in Figure 3(b), all samples have a negative Seebeck coefficient and show an increasing trend with increasing temperature, indicating that the majority carriers are

electrons. In contrast to the Seebeck coefficient of  $-167 \mu\text{V/K}$  for the Matrix sample, the Seebeck coefficients of  $-179$  and  $-176 \mu\text{V/K}$  were obtained for samples containing  $0.5\%$  P25 plus  $0.5\%$   $\text{TiO}_2$  aggregates and  $0.5\%$  commercial available P25, respectively. The Seebeck coefficient, according to Mott, can be expressed as<sup>38</sup>

$$S = \frac{\pi^2 \kappa_b^2 T}{3q} \left( \frac{d \ln N(E)}{dE} + \frac{d \ln \tau(E) \nu(E)^2}{dE} \right)_{E=E_F}$$

Here  $N(E)$  and  $\tau$  are the electronic density of states and the relaxation time, respectively. It is difficult to conclude, based on the electronic transport properties of  $\text{TiO}_2$ -containing  $\text{Ba}_{0.3}\text{Co}_4\text{Sb}_{12}$  whether one also benefits from the first term in the above equation (any sharpening of the density of states at the Fermi level). However, band bending at the interface between embedded nanoinclusions and the semiconducting matrix might produce a scattering potential that could preferentially scatter low energy charge carriers, leading to an enhancement of the Seebeck coefficient, just as was the case with Pb nanoparticles in  $\text{PbTe}$ <sup>40</sup>.

The power factor of these samples is plotted in Figure 3(c). The power factor of the Matrix sample is very close to that of the sample containing  $0.5 \text{ wt}\%$   $\text{TiO}_2$  aggregates over the entire temperature range. Matrix possesses the largest electrical conductivity of all samples as shown in Figure 3(a). However, its Seebeck coefficient is inferior to  $\text{TiO}_2$  containing samples, and thus the highest power factor is obtained for samples with  $0.5\%$  P25 plus  $0.5\%$   $\text{TiO}_2$  aggregates, reaching its maximum of  $6 \text{ mW K}^{-2}\text{m}^{-1}$  at around  $675 \text{ K}$ . Figure 4(a) displays the temperature dependence of the thermal conductivity ( $\kappa$ ) for samples containing different amounts of  $\text{TiO}_2$  inclusions calculated from equation  $\kappa = \lambda C_p d$ . The Matrix sample has the largest thermal conductivity, reaching  $4.52 \text{ W/m-K}$  at  $813 \text{ K}$  which is  $16\%$ ,  $5\%$ ,  $14\%$ , and  $14\%$  higher than the samples of  $\text{Ba}_{0.3}\text{Co}_4\text{Sb}_{12}$  with

additions of 0.5% P25, 0.5% TiO<sub>2</sub> aggregates, 0.5% P25 plus 0.5% TiO<sub>2</sub> aggregates, and 1% TiO<sub>2</sub> aggregates, respectively. As such, embedding TiO<sub>2</sub> inclusions is, indeed, effective in reducing the thermal conductivity of Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub>. Thermal conductivity can be divided into two major components,  $\kappa = \kappa_e + \kappa_l$ . The electronic component of the thermal conductivity ( $\kappa_e$ ) is due to the heat transported by electrons. The lattice component of the thermal conductivity ( $\kappa_l$ ) is due to the heat transported by phonons and is low for amorphous or disordered materials. The electronic thermal conductivity was estimated using the Wiedemann-Franz law  $\kappa_e = L\sigma T$ , where the constant,  $L$ , known as the Lorenz number, depends on the degree of elasticity in carrier scattering. We take  $L$  at its fully degenerate value of  $L = 2.44 \times 10^{-8} \text{ V}^2/\text{K}^2$  believing to be reasonable for our degenerate semiconductors<sup>17</sup>.

The lattice thermal conductivity is calculated by subtracting the electronic part from the total thermal conductivity as shown in Figure 4(b). Lattice thermal conductivities follow approximately the  $1/T$  relation and samples containing TiO<sub>2</sub> inclusions display much reduced values, except for Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> with 1% aggregates. A contributing factor in this reduction is the enhanced phonon scattering associated with the presence of TiO<sub>2</sub> inclusions. The heat flow is carried by a spectrum of phonons with a rather broad range of wavelengths. Commercial P25 with the size ranging from ~20-40 nm is an efficient scatterer of shorter wavelength phonons, leading to a much lower lattice thermal conductivity of Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub>-containing 0.5% P25. However, in the case of samples with 0.5% hierarchically structured TiO<sub>2</sub>, due to the inclusions being spheres with the diameter of ~600-800 nm, they preferentially scatter phonons with longer wavelength. While the thermal conductivity is reduced, it is decreased to a lesser degree because such phonons

carry a smaller fraction of the overall heat. Noticeably, the lowest lattice thermal conductivity is obtained in samples with 0.5% P25 plus 0.5% hierarchically structured  $\text{TiO}_2$ , reaching 0.50 W/m-K at 813 K and reflecting a wider frequency range of phonons that is affected by scattering on inclusions in  $\text{Ba}_{0.3}\text{Co}_4\text{Sb}_{12}$ . However, when excessive amount of  $\text{TiO}_2$  aggregates of up to 1% are added to  $\text{Ba}_{0.3}\text{Co}_4\text{Sb}_{12}$ , the lattice thermal conductivity becomes larger than that of the Matrix sample. Therefore, only optimized combinations of inclusions in  $\text{Ba}_{0.3}\text{Co}_4\text{Sb}_{12}$ , representing equal amounts of P25 and hierarchically structured  $\text{TiO}_2$ , can reduce the lattice thermal conductivity to the greatest extent. This observation is in a good agreement with our previous study<sup>17</sup>.

The dimensionless thermoelectric figure of merit  $ZT$  of our samples is calculated based on the above transport data and is plotted in Figure 5. It is found that the  $ZT$  values increase with the increasing temperature in the measured temperature range. Among samples with different kinds of  $\text{TiO}_2$  inclusions, the sample with 0.5% P25 plus 0.5% hierarchically structured  $\text{TiO}_2$  aggregates reaches the highest  $ZT$ , around 1.20 at ~813 K. This is approximately 12%, 9%, 20% and 40% higher than the  $ZT$  value of the Matrix, and  $\text{Ba}_{0.3}\text{Co}_4\text{Sb}_{12}$  containing 0.5% P25, 0.5% and 1%  $\text{TiO}_2$  aggregates, respectively. The improvement in the figure of merit is attributed to: (a) embedded  $\text{TiO}_2$  inclusions which provide extra phonon scattering centers and lower the lattice thermal conductivity; (b) the energy filtering effect caused by the interface between nanoinclusions and the matrix which enhances the Seebeck coefficient. This simultaneous enhancement in the Seebeck coefficient and the diminished ability of phonons to conduct heat results in an improved thermoelectric figure of merit of  $\text{Ba}_{0.3}\text{Co}_4\text{Sb}_{12}$  containing  $\text{TiO}_2$  inclusions, in spite of the reduced electrical conductivity in these samples.

#### 4. Conclusion

Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> samples containing hierarchically structured TiO<sub>2</sub> were synthesized by ball milling and subsequent hot press sintering. They show a very good thermoelectric performance with the highest *ZT* value reaching 1.2 at 813K when 0.5% P25 plus 0.5% TiO<sub>2</sub> aggregates are added to the Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> matrix. The TiO<sub>2</sub> nanocrystallite aggregates were prepared using a carbon sphere-templated method. It is found that TiO<sub>2</sub> inclusions are dispersed between the grain boundaries and on the surfaces of grains. The observed enhancement in the Seebeck coefficient together with a significant reduction in the thermal conductivity leads to an improved *ZT* in these TiO<sub>2</sub>/Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> bulk composite structures. This method is a good example of optimizing inclusions to achieve desirable thermoelectric properties in n-type Ba-filled skutterudites.

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## Table caption

Table 1: EPMA data for the sintered TiO<sub>2</sub>-included skutterudites materials.

## Captions for Figures

Figure 1 (a), (b) and (c) SEM images and XRD pattern of TiO<sub>2</sub> nanocrystallite aggregates, (d) XRD patterns of TiO<sub>2</sub>-containing Ba-filled skutterudite bulk samples.

Figure 2 SEM and TEM micrographs showing TiO<sub>2</sub> aggregates dispersed in the matrix: (a), (b) SEM images of Ba-filled skutterudite with aggregates of 0.5% P25+0.5% TiO<sub>2</sub>. The bottom left inset of Fig 2(b) is SEM image of Ba-filled skutterudite with 1% TiO<sub>2</sub> aggregates; (c) a close-up image of TiO<sub>2</sub> sphere, (d), (e) SAED (Selected Area Electron Diffraction) patterns and HRTEM imaging of the TiO<sub>2</sub> aggregates, (f) a schematic diagram illustrating various phonon scattering mechanisms taking place on embedded TiO<sub>2</sub> inclusions.

Figure 3 (a) Electrical conductivity, (b) Seebeck coefficient, and (c) power factor for Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> (Matrix) and Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> with 0.5% P25, 0.5% TiO<sub>2</sub> aggregates, 0.5% P25+0.5% TiO<sub>2</sub> aggregates, and 1% TiO<sub>2</sub> aggregates, respectively.

Figure 4 (a) Thermal conductivity and (b) lattice thermal conductivity for the Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> Matrix, and skutterudite samples containing 0.5% P25, 0.5% TiO<sub>2</sub> aggregates, 0.5% P25+0.5% TiO<sub>2</sub> aggregates, and 1% TiO<sub>2</sub> aggregates. The inset of (b) shows the corresponding electronic thermal conductivity.

Figure 5 *ZT* for the Ba<sub>0.3</sub>Co<sub>4</sub>Sb<sub>12</sub> Matrix and skutterudite samples containing 0.5% P25, 0.5% TiO<sub>2</sub> aggregates, 0.5% P25+0.5% TiO<sub>2</sub> aggregates, and 1% TiO<sub>2</sub> aggregates, respectively.

Table 1: EPMA data for the sintered skutterudite material containing TiO<sub>2</sub>. (The error bar is ~0.02)

Sample	Composition
Matrix	Ba <sub>0.22</sub> Co <sub>3.96</sub> Sb <sub>12</sub>
0.5% P25	Ba <sub>0.21</sub> Co <sub>3.95</sub> Sb <sub>12</sub> Ti <sub>0.09</sub> O <sub>0.19</sub>
0.5% TiO <sub>2</sub> Agg.	Ba <sub>0.21</sub> Co <sub>3.96</sub> Sb <sub>12</sub> Ti <sub>0.10</sub> O <sub>0.19</sub>
0.5% (TiO <sub>2</sub> Agg.+P25)	Ba <sub>0.22</sub> Co <sub>3.95</sub> Sb <sub>12</sub> Ti <sub>0.18</sub> O <sub>0.37</sub>
1 % TiO <sub>2</sub> Agg.	Ba <sub>0.22</sub> Co <sub>3.96</sub> Sb <sub>12</sub> Ti <sub>0.19</sub> O <sub>0.38</sub>









