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## **Predictive approach to CVD of crystalline layers of TMDs: The case of MoS<sup>2</sup>**

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## **Abstract**

Layered transition metal dichalcogenides (TMDs), such as  $MoS<sub>2</sub>$ , are candidate materials for next generation 2-D electronic and optoelectronic devices. The ability to grow uniform, crystalline, atomic layers over large areas is the key to developing such technology. We report a chemical vapor deposition (CVD) technique which yields n-layered  $MoS<sub>2</sub>$  on a variety of substrates. A generic approach suitable to all TMDs, involving thermodynamic modeling to identify the appropriate CVD process window, and quantitative control of the vapor phase supersaturation, is demonstrated. All reactant sources in our method are outside the growth chamber, a significant improvement over vapor-based methods for atomic layers reported to date. The as-deposited layers are p-type, due to Mo deficiency, with field effect and Hall hole mobilities of up to 2.4  $\text{cm}^2/\text{V-s}$  and 44  $\text{cm}^2/\text{V-s}$  respectively. These are among the best reported yet for CVD  $MoS<sub>2</sub>$ .

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**Keywords**: MoS<sub>2</sub>, chemical vapor deposition, thermodynamic modeling, supersaturation

## **1. Introduction**

2-D or layered materials came into recent reckoning with the observation of ambipolar electric field effect in graphene flakes.<sup>1,2</sup> Following graphene,<sup>3,4</sup> a single sheet of carbon atoms, TMDs have emerged as a widely studied family of 2-D materials.<sup>5,6</sup> Given the diversity of their properties, ranging from semiconducting to metallic to superconducting,  $7-9$ the dichalcogenides or their combinations with layered conductors like graphene, or insulators like BN,<sup>10</sup> have already been explored for many novel hybrid electronic and optoelectronic devices. 11-13

As with graphene, most initial research on dichalcogenides was confined to micronsized flakes obtained by exfoliation from the bulk.<sup>1,14,15</sup> Such exfoliation is, however, not scalable. For integration of dichalcogenide-based devices into existing microelectronics technology, a CVD technique that allows controlled deposition of uniform layers over large areas, akin to that developed by Li et al. for graphene is required.<sup>16</sup> Currently, layers of MoS<sub>2</sub>, the most widely studied dichalcogenide, have been obtained over large areas by sulfurization of evaporated MoO<sub>3</sub><sup>17-21</sup> and MoCl<sub>5</sub><sup>22</sup> dip-coated (NH<sub>4</sub>)<sub>2</sub>MoS<sub>4</sub><sup>23</sup> sputtered Mo metal<sup>24</sup> and Mo-Au alloys.<sup>25</sup> Most of these efforts<sup>17-24</sup> involve vaporization of the starting materials within a reactor at fairly elevated temperatures of 500-900 °C. This progress is significant as the structure of dichalcogenides is more complex than that of graphene rendering their nucleation and growth in the desired orientation more difficult.<sup>19</sup> However, technology integration requires a CVD method in which all precursors are introduced into the reactor in the gaseous form, from sources that can be vaporized externally at temperatures as low as possible. Although MoS<sub>2</sub> has been synthesized using vapor-transported Mo(CO)<sub>6</sub><sup>26</sup> and  $(NH_4)_2MOS_4$ <sup>27</sup> the end product did not consist of a controlled number of uniform layers over large areas. More importantly, hybrid devices such as those consisting of multiple

chalcogenides, graphene, and BN layers can also be synthesized in the same reactor without interruption or exposure to air, by the appropriate introduction of precursors. This would not be possible in the vapor transport processes, such as the ones previously mentioned, in which one of the precursors is always present in the reactor.  $17-24$ 

We have developed a CVD process using  $Mo(CO)_{6}$  and  $H_{2}S$  as precursors, which enables the uniform deposition of n-layered,  $n=1, 2, 3...>6$ ,  $MoS<sub>2</sub>$  over large areas on a wide variety of substrates, including sapphire, oxidized silicon and fused quartz. The low boiling point of  $Mo(CO)_{6}$  (156°C) and  $H_{2}S$  gas make these ideal precursors for CVD. Thermodynamic modeling was used to first deduce the growth window in which  $MoS<sub>2</sub>$  is the only solid phase formed. Kinetic control was then applied in this growth window by controlling supersaturation to obtain crystalline layers parallel to the substrate surface, as required for 2-D electronics. In the process developed, all precursors are brought into the vapor phase outside of the reactor at temperatures  $\langle 100^{\circ}$ C. This is not only critical to reducing nucleation and growth rates but also to the synthesis of doped device structures. A combination of thermodynamic modeling and kinetic control similar to that used in this study to deposit  $MoS<sub>2</sub>$  layers would be applicable to all TMDs of interest.

## **2. Experimental**

**2.1. Thermodynamic Modeling:** A commercially available PC-based software program was used for the thermodynamic calculations [Autokumpu HSC Chemistry 6.1, Finland] carried out to identify the growth window.

**2.2. Methods:**  $Mo(CO)_{6}$  (Sigma Aldrich (99.99%)) powder and  $H_2S$  (Bhoruka, 99.5%) gas were used for CVD growth of MoS<sub>2</sub> films. Growth was carried out in a custom-built

horizontal, hot-wall CVD reactor, shown in Fig. 1. The substrates,  $SiO<sub>2</sub>/Si$ , c-plane sapphire, and fused quartz, were cleaned in acetone, isopropyl alcohol (IPA), and de-ionised (DI) water in an ultrasonic bath for 5 min each, and dried with UHP  $N_2$  stream. Mo(CO)<sub>6</sub> powder was placed in a vaporizer and maintained at  $25{\text -}130^{\circ}\text{C}$  to obtain the required precursor concentration. To ensure complete transfer of the precursor in to the reactor, the vapor transport lines were maintained at the vaporizer temperature. Ar and  $H_2$  were used as carrier gases, and  $H_2S$  as the reacting gas. To control the number of  $MoS_2$  layers, the reactor pressure  $(P_{\text{tot}})$  was maintained between 20 to 600 Torr. The growth temperature (T) was varied from 350°C to 850°C. To control the thickness of the layers formed, deposition duration (t) was varied from 10 s to 10 min. Prior to deposition, the reactor was pumped down to the desired pressure ( $P_{tot}$ ), and the temperature was ramped at 25°C/min under a carrier gas flow of Ar and/or  $H_2$  to the desired T and held for 15 min. The Mo(CO)<sub>6</sub> powder in the vaporizer is heated to the required temperature, maintaining a carrier gas flow of 15 sccm. The total pressure in the vaporizer is maintained at 20 psi, to achieve a flow of 1.8 sccm of  $Mo(CO)<sub>6</sub>$  into the reactor (detailed calculation is in supporting information). The ratios of gas flow rates, Mo:H<sub>2</sub>S:(Ar or H<sub>2</sub>), were maintained at 1:10:100 or 1:10:1000. After deposition, the films formed were sulfurized at the same T and  $P_{tot}$  for 10 min, with only  $H_2S$ flowing through the reactor.

**2.3. Characterization and device fabrication:** The samples were examined by optical microscopy (Leica) and field emission scanning electron microscopy (FESEM, Zeiss, Ultra 55) to check the uniformity of  $MoS<sub>2</sub>$  films deposited. Raman (Horiba, LabRam) spectra were used to determine the number and uniformity of layers. Atomic force microscopy (Bruker) was used to confirm the number of layers. The chemical composition was ascertained by xray photoelectron spectroscopy (Shimadzu, Axis Ultra DLD) (see supporting information for more details). The binding energies have been measured with respect to graphitic carbon at

284.5 eV. The crystallinity of the layers was examined by transmission electron microscopy (FEI, T20). For TEM and electrical characterization,  $MoS<sub>2</sub>$  layers were transferred onto a copper grid or a 500-nm-SiO<sub>2</sub>/Si substrate, accordingly. For transferring onto a grid, a drop of IPA was placed on the substrate along with the grid. After the IPA is allowed to dry, the substrate was immersed in a HF solution (1:10), resulting in transfer of the  $MoS<sub>2</sub>$  layer onto the grid. For transferring onto a  $SiO<sub>2</sub>/Si$  substrate, the sample was spin-coated with PMMA at 6000 rpm for 60 s. It was then baked at  $150^{\circ}$ C for 2 min. The sample was then immersed in HF solution (1:5) until the PMMA layer with  $MoS<sub>2</sub>$  detached from the underlying substrate. FETs were fabricated with 20  $\mu$ m channel length (L) and 200  $\mu$ m width (W) using photolithography. Source and drain pads of Cr/Au, 20/80 nm, were deposited by e-beam evaporation. 30 nm  $Al_2O_3$  deposited by e-beam evaporation was used as the top gate oxide, with Cr/Au of 20/80 nm as the gate electrode. Hall measurements were carried out to determine mobility and carrier concentration.

## **3. Results and discussion**

## **3.1. Thermodynamic modeling to determine growth window**

By assuming the validity of equilibrium thermodynamics for the CVD process, "phase stability diagrams" can be determined, which can predict "process windows" for the deposition of the desired solid. The stability diagrams for the solid products of the reactant system of  $Mo(CO)<sub>6</sub>-H<sub>2</sub>S-Ar$ , for different temperatures (T) and reactant mixtures, obtained through thermodynamic calculations, are plotted in Fig. 1. The details of the thermodynamic calculations are provided in the supporting information, including the list of all the possible solid and gaseous products from the Mo-S-C-H-O system considered in the analysis (Table S1). All depositions for verifying the analytical results were carried out in the reactor shown

in Fig. 1a. Our thermodynamic modeling predicts that only  $MoS<sub>2</sub>$ , carbon, and  $MoS<sub>3</sub>$  would be the solid phases formed. For a molar ratio of 10:1 between the precursors  $Mo(CO)_{6}$  and  $H_2S$ , used for most of the growth conditions reported here, formation of  $MoS_3$  is predicted at T<200 $^{\circ}$ C. However, no deposition of MoS<sub>3</sub> was observed at these temperatures, as shown in Fig. S1. That the thermodynamically predicted formation of  $MoS<sub>3</sub>$  or  $MoS<sub>2</sub>$  at low temperatures is not experimentally observed is attributed to kinetic limitations. Raman peaks corresponding to  $MoS_2$  were observed only at T $>350^{\circ}$ C. Subsequent discussion is therefore confined to depositions at  $T > 350^{\circ}C$ .

The ability of these phase diagrams in predicting growth outcomes is corroborated out by the Raman data from samples deposited over a range of temperatures for a particular pressure and gas composition. The phase diagram of Fig. 1b predicts that, with argon as carrier gas, the carbon content of the solid deposit decreases as the growth temperature (T) is raised. The Raman data, Fig. 1c, confirm this prediction, highlighting the importance of predictive modeling. On the other hand, the presence of minute residual amounts of carbon at 700°C beyond the 575°C limit predicted by phase diagram - shows the need for experimental verification, as the accuracy of modeling is limited by the precision of the thermodynamic data available. At 850°C, there is "no trace" of carbon, and the presence of only  $MoS<sub>2</sub>$  is evidenced by the Raman spectrum.

Total reactor pressure,  $P_{tot}$ , is one of the most important CVD parameters. An optimal choice of  $P_{\text{tot}}$ , as will be seen later, was critical in controlling  $MoS<sub>2</sub>$  layer deposition in this study. Thermodynamic modeling, Fig. 1d, predicts that, for a given  $T$ , an increase in  $P_{tot}$  leads to an increase in the carbon content of the solid phase. Raman spectra (Fig. S2-a) of deposited samples confirm this trend. The corresponding phase stability diagram at  $P_{tot} = 600$  Torr, the

upper bound used in our studies, is included in the supporting information, Fig. S3. Modeling predicts (Fig. 2a), however, that, at a given T, the deposition of carbon at high  $P_{tot}$  can be precluded by using a mixture of argon and hydrogen as the carrier gas (as is common in CVD). Carbon content can be reduced by increasing the partial pressure of hydrogen. Indeed, in pure  $H_2$  ambient, carbon is predicted to be absent at all T when  $P_{tot} > 10$  Torr (Fig. 2b). This prediction is confirmed by the Raman data in Fig. 2c and Fig. S2-b. The ability to predict the completely different outcomes on using Ar versus Ar+H<sup>2</sup> mixtures as carrier gas highlights the importance of such thermodynamic modeling.

## **3. 2. Growth of few-layered to mono-layered MoS<sup>2</sup> by kinetic control**

To obtain monolayered or few-layered  $MoS<sub>2</sub>$  films predictably and with precision, kinetic control, within the thermodynamic window, is called for and may be exercised through variation of supersaturation. Under conditions wherein thermodynamic modeling predicts the formation of pure  $MoS<sub>2</sub>$ , the gas phase supersaturation, or  $\Delta G$  for the chemical reaction, Equation 1, was calculated using Equation 2.

$$
Mo(CO)6 + 2H2S \longrightarrow MoS2 + 6CO + 2H2
$$
 (1)

$$
\Delta G = RT \ln \left( \frac{\eta_{Mo(CO)_6} \eta_{H_2S}^2 \eta_{tot}^5}{\eta_{CO}^6 \eta_{H_2}^2 P_{tot}^5 K_{Eq}} \right)
$$
 (2)

where  $P_{\text{tot}}$  is the total reactor pressure in atmospheres, T the growth temperature and  $\eta$ , the flow rates of various gases as indicated by the subscripts.  $K_{Eq}$  calculated from standard data is listed in the supporting information. A decrease in  $\Delta G$  was affected by reducing precursor fluxes suitably and increasing the total pressure (Equation 2). The effect of this reduction is tabulated in the supporting information, Table S2. The  $\Delta G$  values would be a maximum at the

zero-point in the reactor (Fig. 1a), and would decrease along the length of the reactor. According to Equation 2, lower temperatures would result in reduced supersaturation. However, the Raman data in Fig. 1c and 2c show that  $MoS<sub>2</sub>$  films deposited at 500°C have a larger full-width at half maximum (FWHM) than films deposited at  $850^{\circ}$ C. A similar trend is also observed at higher pressure as shown in Fig. S2b. As larger FWHM is indicative of a more defective film, the depositions were carried out at  $850^{\circ}$ C. P<sub>tot</sub> being the only independent variable parameter in the logarithmic term in Equation 2, it is a critical parameter for reducing  $\Delta G$ . At P<sub>tot</sub> = 20 Torr, the largest  $\Delta G$  values used in our study, bulk MoS<sub>2</sub> was obtained in 10 minutes everywhere in the reactor. The effect of reduction in  $\Delta G$  by increasing  $P_{tot}$  from 20 Torr to 850 Torr, in the thermodynamically permissible window at 850 $^{\circ}$ C, spanning the duration of growth (t) is shown in Fig. 3. The reduction in supersaturation achieved by this increase in  $P_{tot}$  was 175 kJ/mol as calculated from Equation 2. The details of the calculations are provided in the supporting information. As summarized in Fig. 3, by the SEM, AFM, optical and spectroscopic information, a combined reduction in  $\Delta G$  and t resulted in uniform, controlled deposition of  $MoS<sub>2</sub>$  from bulk-like to 6 layers to a monolayer. The number of layers obtained for the identical growth conditions was the same for sapphire,  $SiO_2/Si$  and quartz. For a combination of the smallest  $\Delta G$  and t, growth was slowed down well enough to yield discrete hexagonal islands. To put the supersaturation concerned in perspective, we estimate that the large islands demonstrated by recent PVT reports using bulk  $MoS<sub>2</sub>$  as the source were obtained with a  $\Delta G$  of ~131 kJ/mol(see supporting information ). Typical durations required to obtain a full monolayer by PVT is of the order of 20 minutes,<sup>28</sup> versus 1 minute in the present study. This indicates that the actual supersaturation values in the CVD here are much higher. While equipment limitation prevents us from achieving even lower supersaturations currently, the present work shows that, with a further reduction in  $\Delta G$  and hence in nucleation rates, CVD would eventually

yield islands as large as those obtained by PVT. For instance, Equation 2 shows that this could be done by going to even higher total pressures. Fig. 3a provides SEM images of  $MoS<sub>2</sub>$ films ranging from few layers to islands  $25 \mu m$  in size. MoS<sub>2</sub> grains and complete monolayer grown at 850 Torr (850°C) for various growth times are shown in Fig. S4. The reduction in FWHM of the Raman spectra shows the improvement in quality. Fig. 3b shows the topography and step height obtained by AFM of crystalline  $MoS<sub>2</sub>$  from few-layers to 1L. The step height of 1L is around 0.75 to 0.77 nm, consistent with previously reported values of 0.72 to 1 nm.<sup>22,29</sup> The step height of 2L and 3L is ~1.55 nm and ~2.4 nm, respectively. The AFM images of  $MoS<sub>2</sub>$  on other substrates are shown in Fig. S5. The Raman spectra of  $MoS<sub>2</sub>$ layers are shown in Fig. 3c, along with that of exfoliated bulk  $MoS<sub>2</sub> (MoS<sub>2</sub>>6$  layers which has been exfoliated using scotch tape) as reference. As the layer thickness decreases, the interlayer coupling decreases, reducing the gap between  $E_{2g}^1$  and  $A_{1g}$ . The difference is 18.9 cm<sup>-1</sup>, 21.8 cm<sup>-1</sup>, 22.3 cm<sup>-1</sup> and 27 cm<sup>-1</sup> for 1L, 2L, 3L and bulk exfoliated MoS<sub>2</sub>, respectively.<sup>22,29-31</sup> Fig. 3d shows variation in the Raman peak separation and position with the number of layers. It is evident that, as thickness increases, the in-plane mode  $E_{2g}^{1}$  redshifts and  $A_{1g}$  mode blue-shifts. The peak separation also increases with thickness, but equalizes for films thicker than 6L. Details about the variation in the FWHM of the peaks with growth conditions are given in Fig. S6. PL spectra, which provide another indication of the number of layers, are shown in Fig. 3e for the  $MoS<sub>2</sub>$  layers grown on sapphire. The two prominent peaks around 627 nm and 667 nm correspond to A1 and B1 direct excitonic transitions. The strength of the PL signal is in the order  $1L > 2L > 3L$ , with bulk MoS<sub>2</sub> yielding a negligible PL signal. This reaffirms that the bandgap  $(E_g)$  in MoS<sub>2</sub> changes from direct to indirect, as layer thickness increases. The peaks undergo a red-shift with increasing thickness, highlighting a reduction in  $E_g$ , as expected.<sup>32-33</sup>

Uniform deposition over large areas, for different layers, is illustrated in Fig. 4a, which shows photographs of  $MoS<sub>2</sub>$  on 1 cm x 1 cm sapphire substrates. The variation in the Raman spectra and the FWHM of monolayer  $MoS<sub>2</sub>$  on different substrates is shown in Fig. 4b. The FWHM of the films deposited on the substrates ( $E_{2g}^1$ : 3.1-3.6 cm<sup>-1</sup> and  $A_{1g}$ : 5.1-5.3 cm<sup>-1</sup>) is higher than that of exfoliated single layer MoS<sub>2</sub> ( $E_{2g}^1$ : 2.8 cm<sup>-1</sup> and  $A_{1g}$ : 4.7 cm<sup>-1</sup>). The variation in the FWHM for the substrates  $SiO<sub>2</sub>/Si$ , sapphire and quartz is however not very significant, indicating that layers are of comparable quality, irrespective of the substrate. These FWHM are comparable with PVT grown monolayer values  $E_{2g}^1$ : 3.5 cm<sup>-1</sup> and  $A_{1g}$ : 6.6 cm<sup>-1</sup>.<sup>18</sup> Fig. 4d-e shows the uniformity in coverage on sapphire substrate through Raman mapping of monolayer MoS<sub>2</sub> over a 100  $\mu$ m x 100  $\mu$ m scan area. The peak separation is in the range of 18.5 to 20 cm<sup>-1</sup>. These values concur with a previous report on monolayer  $M_0S_2$ <sup>22</sup>. The Raman spectra and mapping data for triangular islands of  $MoS<sub>2</sub>$  on  $SiO<sub>2</sub>/Si$  and monolayer films on diverse substrates are given in Fig.  $S7-10$ . Uniformity over entire 1 cm<sup>2</sup> sapphire is shown in Fig. S8.

The crystalline structure of the grown layers wa investigated by TEM. Fig. 5a shows a lowmagnification image of a monolayer grown on  $SiO_2/Si$  at 500 Torr at 850°C. The highresolution image (Fig. 5b) shows the hexagonal ordering present at the atomic scale. The inset in Fig. 5b highlights the single-crystalline diffraction pattern obtained. The chemical composition of the deposited layers was ascertained by XPS analysis. Fig. 5c shows the spectra obtained from 3L films as-deposited. The Mo spectra show evidence of the presence of  $Mo^{4+}$ , corresponding to  $MoS_2$ , and  $Mo^{6+}$ , corresponding to  $MoO_3$ . The formation of  $MoO_3$ is attributable to the oxidation outside the CVD chamber - of metallic Mo arising either from incomplete sulfurization or from the partial reduction of the deposited  $MoS<sub>2</sub>$  while being cooled in  $H_2$  ambient. When the MoS<sub>2</sub> films are annealed post growth in a  $H_2S$  ambient for

10 min and cooled in the same ambient, only the peak corresponding to  $Mo^{4+}$  is seen in the XPS spectrum (Fig. 5d) The sulfurization step is thus crucial to achieving the "right" chemical composition. After sulfurization, the Mo:S ratio is 0.8:2, indicating Mo deficiency. This ratio calculated for  $MoS_2 > 3L$ (detailed growth conditions in Table S2) also gave similar results. The Mo error bar is 0.82±0.08 for 3L( 5 samples), 0.85±0.09 for 4L(5 samples) and 0.79±0.06 for 5L(2 samples). This Mo deficiency, as discussed in the next section, results in p-type behavior of the as-grown films. For reference purposes, bulk exfoliated  $MoS<sub>2</sub>$  sample was also characterised by XPS (see Fig. S11). The calculated atomic concentration of Mo:S was 1:1.912 thereby implying it should be n-type behavior. Device characterstics, see Fig. S11, indicate that this is indeed so. This exercise thus serves as an internal calibration of our XPS data and shows that the correlation with device behaviour is consistent. Most data in the literature, except in one report,  $24$  concern n-type MoS<sub>2</sub>.<sup>34-38</sup> The ability to vary stoichiometry by varying the vapor phase composition is one of the strengths of CVD and is very important for obtaining p- and n-type layers for device applications.

Preliminary electrical characterization of the CVD  $MoS<sub>2</sub>$  - field effect transistors (FET) fabricated both in top-gate and bottom-gate configurations also confirm the Mo deficiency shown by XPS. Fig. 6a is a schematic representation of a such a  $MoS<sub>2</sub> FET.$  Fig. 6b shows the output characteristics of the device and Fig. 6c shows transfer characteristics of a 20  $\mu$ m channel-length device. The top-gate characteristics are in the inset. The linear output characteristics for a low drain voltage indicate good ohmic contact at the source and drain pads. The transfer characteristics of the  $MoS<sub>2</sub>$  device show the aforementioned p-type behaviour. The field effect mobility of the charge carriers is extracted from the slope of the linear regime, using Equation 3

$$
\mu = \left(\frac{\partial I_{ds}}{\partial V_g}\right) \frac{L}{WC_{ox}V_{ds}}
$$
\n(3)

where L is the channel length, W the channel width, and  $C_{OX}$  the gate capacitance.

It is observed that on-off ratios for the trilayer  $MoS_2$  device is around  $10^5$ , with an estimated mobility of 2.4  $\text{cm}^2\text{V}^{-1}\text{s}^{-1}$  in the bottom-gate configuration, and  $\text{I}7.1 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$  in the top-gate configuration. Fig. 6d summarizes the mobility obtained from FETs with  $MoS<sub>2</sub>$  of differing number of layers, including electron mobility of exfoliated bulk  $MoS<sub>2</sub>$ . The mobility is seen to decrease with the increase in layer number. This is contrary to reported literature trends.34-38 The layers used for the electrical characterization were grown under different supersaturations(details in Table S2). For instance, 6 layered  $MoS<sub>2</sub>$  has been grown under higher supersaturation than 3 layered  $MoS<sub>2</sub>$ . As a result, the grain size in these films are also different. FWHM of these layers, given in Figure S6, shows that with increasing layer numbers, crystallinity degrades. Consequently, we believe that the precise mechanism of mobility degradation originates from defect scattering as indicated by the reduction in crystallinity. As suggested by the results in this paper, even better layers can be achieved by tuning the physico-chemical parameters during growth. Hall measurements, which are independent of contact resistance were also performed(included in the supporting information), confirms p-type behavior. A carrier concentration of 3 x 10<sup>-16</sup> cm<sup>-3</sup> and a mobility of 44  $\text{cm}^2/\text{V}$ -sec was measured on 3L samples. These Hall mobilities are among the highest reported for p-type  $MoS<sub>2</sub>$ . (in Table S4 of supporting information). The low hole concentrations point to a deep acceptor level introduced by Mo vacancies. Our calculations (see supporting information) indicate that this level should lie at 0.62 eV above the valence band edge in reasonable agreement with theoretical estimates which point to their being at 1  $eV<sup>39</sup>$ 

## **4. Conclusions**

In summary, a CVD technique compatible with traditional device integration has been developed to deposit  $MoS_2$  atomic layers over large areas using  $Mo(CO)_6$  and  $H_2S$ . Thermodynamic modeling has been used to deduce the growth window most conducive to the deposition of pure  $MoS<sub>2</sub>$ . Control of supersaturation within this growth window was critical to obtaining uniform layers of the desired thickness, from discrete crystalline islands to monolayers to films >6 layers. The layers deposited have been investigated by Raman spectroscopy, XPS, AFM and TEM to ascertain the number of layers, their uniformity and crystallinity. Monolayers of  $MoS<sub>2</sub>$  grown at lower supersaturation have single-crystalline domains as large as 25 um. The layers are p-type, due to Mo deficiency. Field effect and Hall hole mobilities of up to 2.4 cm<sup>2</sup>/V-s and 44 cm<sup>2</sup>/V-sec respectively were measured at a carrier concentration of  $3x10^{16}/cm^3$ .

## **ACKNOWLEDGEMENTS**

The authors would also like to acknowledge the Department of Science and Technology, Government of India for the financial support under the Grant No SR/S2/CMP-02/2007, SR/S3/ME/045/2007, SR/WOS-A/CS-69/2012 and the TUE project for nanoscience and technology. The authors would like to acknowledge the National Nanofabrication Center (NNfC), Micro and Nano Characterization Facility (MNCF) and the Advanced Facility for Microscopy and Microanalysis (AFMM) for the fabrication and characterization work done. I would like to thank Prof. Navakanta Bhat, CeNSE IISc for his suggestions.

## **SUPPORTING INFORMATION**

Methodology for thermodynamic modeling, supersaturation calculation, Raman, optical microscopy and AFM analysis of  $MoS<sub>2</sub>$  single layers are available in the supporting information.

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## **Figure captions**

**Fig. 1** CVD phase stability in Ar ambient (a) Schematic representation of the CVD system used for the synthesis of  $MoS<sub>2</sub>$  (b) Calculated stability windows for various condensed phases at P<sub>tot</sub> = 20 Torr. (c) Raman spectra (with exfoliated bulk  $MoS<sub>2</sub>$  as reference) of multilayer MoS<sub>2</sub> films deposited in argon ambient on sapphire over a wide range of T, with  $P_{tot} = 20$ Torr and  $Mo(CO)<sub>6</sub>:H<sub>2</sub>S = 1:10$  with indices in bracket showing FWHM of corresponding  $E_{2g}^{1}$  and A<sub>1g</sub> peaks. (d) Contour plot showing the variation in the mole percentage of MoS<sub>2</sub> in the  $MoS<sub>2</sub>-C$  film at different temperatures and pressures for the composition represented by the dotted line in (b). The colored dots in (b) and (d) correspond to Raman spectra in (c) of samples 1, 2, 3 and 4 deposited at different temperatures.

Fig. 2 CVD phase stability in Ar/H<sub>2</sub> and H<sub>2</sub> ambient (a) Contour plot showing the formation of MoS<sub>2</sub> and C, at P<sub>tot</sub> = 20 torr, when the Ar:H<sub>2</sub> ratio is changed from 75:25 to 0:100 with  $Mo(CO)<sub>6</sub>:H<sub>2</sub>S = 1:10$ . (b) The contour plot for the formation of  $MoS<sub>2</sub>$  and C in H<sub>2</sub> ambient, as a function of T and P<sub>tot</sub>. (c) Raman spectra (with exfoliated bulk  $MoS<sub>2</sub>$  as reference) of multilayer  $MoS<sub>2</sub>$  films deposited on sapphire in  $H<sub>2</sub>$  ambient over a wide range of T, keeping  $P_{tot}$  = 20 Torr and Mo(CO)<sub>6</sub>:H<sub>2</sub>S = 1:10. The indices in bracket show FWHM of corresponding  $E_{2g}^1$  and  $A_{1g}$  peaks. The colored dots in (b) correspond to Raman spectra in (c) of samples 1, 2 and 3 deposited at different temperatures. As predicted carbon(G) is absent.

**Fig. 3** Few-layers to monolayer through supersaturation control (a) SEM images of  $MoS<sub>2</sub>$ , from few to monolayer, deposited in hydrogen ambient at  $850^{\circ}$ C, by varying supersaturation. The  $P_{tot}$  which controls the supersaturation for given flux, and growth duration (t) are indicated. (b) AFM images, with superimposed line scan, showing  $MoS<sub>2</sub>$  films, from three-

layered (3L) to a monolayer (1L). (Scale bar corresponds to 2  $\mu$ m) (c) Raman spectra of  $MoS<sub>2</sub>$  films, from 1L to 6L, together with that of exfoliated bulk  $MoS<sub>2</sub>$  (d) The variation in the peak separation and position of the Raman spectra, as a function of the number of layers (e) The PL spectra of  $MoS<sub>2</sub>$  films as the number of layers L increases, showing the change from direct to indirect band gap for large L.

Fig. 4 Uniformity of the deposited layers. (a) Photograph of MoS<sub>2</sub> films with different L, grown on sapphire substrate. The color is an indication of the number of layers grown. (b) Raman spectra of  $MoS<sub>2</sub>$  monolayers on various substrates along with their FWHM. The Raman spectra of exfoliated monolayer and bulk  $MoS<sub>2</sub>$  is also included (c) Optical image of single layer of  $MoS<sub>2</sub>$  grown on sapphire, with a 100  $\mu$ m x 100  $\mu$ m area selected for mapping of the (d)  $E_{2g}^{1}$  mode (e)  $A_{1g}$  mode and (f) peak difference. The mapping highlights the uniformity in layer thickness.

**Fig. 5** Structural and chemical analysis of MoS<sub>2</sub> layers grown at 850 $^{\circ}$ C by TEM and XPS. (a) Low- magnification image of a monolayer. (b) High-resolution image of a monolayer. The inset shows the characteristic single crystalline pattern, with hexagonal symmetry. XPS spectra of MoS<sub>2</sub> films 3L grown at 20 Torr, before (c) and after (d) sulphurization is presented. The Mo 3d spectrum obtained before sulphurization shows the presence of  $Mo^{6+}$ , corresponding to  $MoO<sub>3</sub>$ , which is absent after sulphurization.

Fig. 6 Electrical characterization of MoS<sub>2</sub> MOSFETs. (a) Schematic representation of a dualgate  $MoS<sub>2</sub> FET transistor.$  (b) Output characteristics of the device. (c) Transfer characteristics of a device with 20  $\mu$ m channel length, with the bottom-gate configuration. The top-gate characteristics are shown as an inset. (d) Mobility as a function of the number of layers.











**Figure 3.**







## **Figure 6.**



## **Table of Contents entry image:**

A generic approach suitable to all transition metal dicalchogenides has been demonstrated. Thermodynamic modeling has been used to identify the appropriate CVD process window. The growth has been controlled by tuning the vapor phase supersaturation. As as result  $MoS<sub>2</sub>$ growth from bulk to monolayers has been acheived.

