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# Light-Induced Degradation of Mixed-cation, Mixed-halide Perovskite: Observed Rates and Influence of Oxygen

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## Abstract

Formamidinium-rich lead halide perovskite semiconductors comprise the absorber layer in the most efficient single-junction perovskite solar cells (PSCs) but suffer from chemical instability when exposed to high temperatures, moisture, oxygen, and light. Light-induced degradation (LID) is unavoidable in PSCs and can be slowed only by limiting the escape of decomposition products. Here, we study the LID of  $\text{FA}_{0.8}\text{Cs}_{0.2}\text{Pb}(\text{I}_{0.83}\text{Br}_{0.17})_3$  thin films using *in-situ* and *ex-situ* optical spectroscopy, microscopy, and x-ray diffraction. The results reveal that the primary decomposition products under LID conditions are reduced lead-containing species that have broadband optical absorption. XPS reveals the presence of  $\text{Pb}^0$ , but it may coexist with partially reduced lead-containing species. We use *in-situ* sub-bandgap optical absorbance measurements to selectively

1 probe and quantitatively measure the formation rate of reduced lead species. We derive a rate law  
2 for reduced-Pb formation ( $r_{Pb^0}$  estimated at  $\sim 3 \times 10^{-10}$  mol/(m<sup>2</sup>s) at 25 °C in N<sub>2</sub> under 1 sun photon  
3 flux that would result in complete conversion of a 300 nm film in  $\sim 78$  days), determine an  
4 activation energy ( $\sim 0.61$  eV), determine an effective reaction order with respect to the flux of  
5 above bandgap photons ( $r_{Pb^0} \propto I_{in}^{0.72}$ ), and find that the wavelength of above bandgap photons  
6 minimally affects the rate, suggesting that PbI<sub>2</sub> photolysis is not the mechanism for the formation  
7 of decomposition products. These observations represent the first quantitative measurements of  
8 the rate of formation of reduced lead species in perovskites and emphasize a unique interplay  
9 among environmental stressors and degradation pathways for commercially relevant perovskite  
10 materials.

Hybrid organic-inorganic perovskites with the chemical formula  $ABX_3$  are a unique class of semiconductors that exhibit simultaneous ease of fabrication (solution and/or vapor processing<sup>1</sup>) and excellent optoelectronic quality (long diffusion lengths and carrier lifetimes,<sup>2</sup> high absorption coefficients,<sup>3</sup> and tunable bandgap<sup>4</sup>). Perovskite single-junction and perovskite-on-silicon tandem photovoltaic (PV) cells have achieved 26.7% and 34.6% power conversion efficiency<sup>5</sup>, respectively, as of December 2024. Despite their high defect concentration ( $\sim 10^{15}$ - $10^{16}$  cm<sup>-3</sup>) compared to conventional covalent semiconductors such as c-Si or GaAs ( $< 10^{12}$  cm<sup>-3</sup>),<sup>6-8</sup> the ionic nature of halide perovskites causes most of the defect energy levels to lie near the band edges or within the bands, giving these materials their notable defect tolerance.<sup>9</sup> However, halide perovskites are notoriously unstable in the presence of excess heat, above-bandgap photons, oxygen, and moisture. In contrast to current commercial PV materials, halide perovskites readily undergo chemical decomposition reactions when exposed to these stressors that can degrade optoelectronic performance and have, up to this point, limited the lifetime of PSCs well below the multidecade timescale necessary for commercial viability.<sup>10</sup> Stability of the absorber is thus a primary factor hindering PSCs from realizing their commercial potential.

Mixed-cation, mixed-halide perovskites in which the A-site comprises mixtures of formamidinium (FA), methylammonium (MA), and/or cesium, and the X-site comprises mixtures of iodide and bromide, form the absorber layer in some of the most efficient lead-based PSCs and display significant improvements in thermal stability and performance compared with the archetypal perovskite  $MAPbI_3$ .<sup>11-14</sup> These compositions nevertheless exhibit low chemical stability relative to conventional inorganic semiconductors.<sup>15</sup> While strong encapsulation may prevent the ingress of oxygen and moisture for some time during the deployment of a PSC module,<sup>16</sup> illumination is necessary for electrical energy production. Thus, studies of perovskite

1 degradation in inert environments under illumination are vital to probe what may represent the  
2 ultimate performance-limiting case.

3 Previously, we reported the degradation kinetics of  $\text{MAPbI}_3$ <sup>17</sup> and the low bandgap lead-tin  
4 perovskite composition  $\text{FA}_{0.75}\text{Cs}_{0.25}\text{Pb}_{0.5}\text{Sn}_{0.5}\text{I}_3$ <sup>18</sup> in response to oxygen, moisture, and  
5 illumination. While measurable LID was not observed in these previous compositions, initial  
6 investigations into the degradation of  $\text{FA}_{0.8}\text{Cs}_{0.2}\text{Pb}(\text{I}_{0.83}\text{Br}_{0.17})_3$  (referred to as FACsPbIBr)  
7 revealed that unencapsulated films placed in an inert environment undergo decomposition when  
8 exposed to illumination alone. In this study, we investigate this decomposition pathway with the  
9 aims of identifying decomposition products and quantifying the rate of perovskite loss as a  
10 function of environmental stressors, since quantitative kinetic measurements of light-induced  
11 degradation do not currently exist in the literature.

12 The FACsPbIBr films used throughout this study were fabricated by an antisolvent wash method  
13 (described in SI section 1). The pristine films are of high optoelectronic quality. Characterization  
14 by x-ray diffraction (XRD), scanning electron microscopy (SEM), UV-Vis-NIR spectroscopy  
15 (UV-vis), and photoluminescence (PL) spectroscopy are shown in Figure S1. The films have an  
16 optical bandgap of  $\sim 1.66$  eV and display steady-state quasi-Fermi level splitting of 1.29 eV under  
17 1 sun equivalent illumination, which is  $>90\%$  of the radiative limit (1.38 eV). Single-junction solar  
18 cells made from the films with a p-i-n architecture (ITO/Poly-TPD/perovskite/ $\text{C}_{60}$ /BCP/Ag) have  
19 efficiencies of  $\sim 17\%$  (Fig. S2).

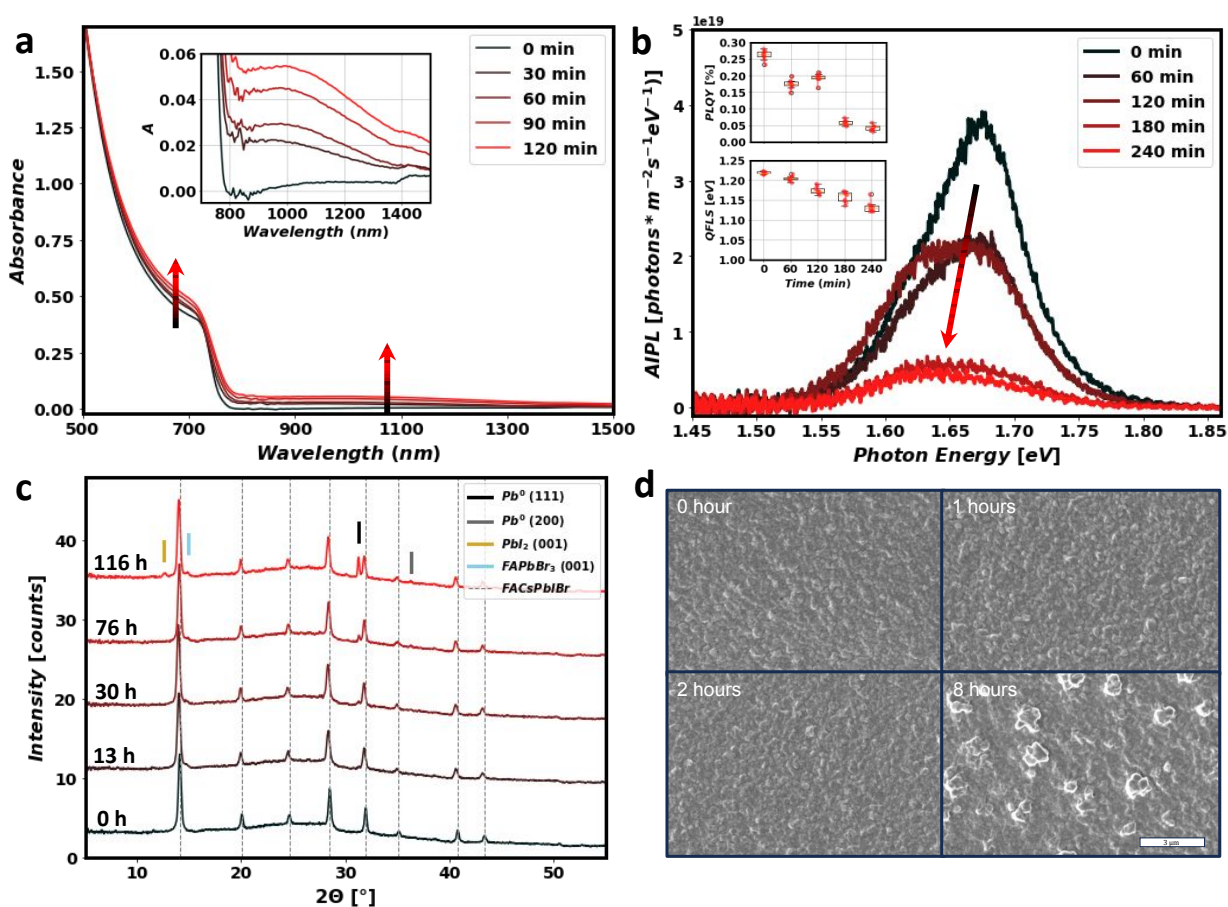
20 *Effects of LID on Mixed-cation, Mixed-halide Thin Films.* To identify the hallmarks of this  
21 decomposition pathway, we performed *ex-situ* characterization of the FACsPbIBr films over a  
22 period of LID. SI section 2 (Fig. S3b) includes the spectra of all illumination sources with a  
23 description of calibration procedures. Sources are calibrated such that the incident above bandgap

absorbed photon flux is identical to that of the AM1.5 G solar spectrum ( $1.5 \times 10^{21}$  photons/m<sup>2</sup>/s for a 1.66 eV absorber).

Unencapsulated FACsPbIBr films deposited on glass substrates were placed in a chemically inert environment (dry N<sub>2</sub> glovebox) on a hotplate (shown schematically in Fig. S3a) set to 65°C under 1 sun equivalent photon flux from a broad-spectrum ‘white’ LED (emission from 400-700 nm) and were periodically removed for characterization. Figure 1 shows the changes of the internal absorbance spectrum (see SI section 4 for further analysis), absolute intensity photoluminescence spectrum (AIPL), XRD pattern, and morphology over the course of degradation and demonstrates that optical absorption is a leading indicator of this decomposition mode. The mean PL peak (Fig. 1b) redshifts modestly (~20 meV). The intensity drops rapidly, and a low energy shoulder develops as degradation proceeds. Since the Br content is relatively low, we do not expect this low energy shoulder to be a result of halide segregation but rather of shallow defect states.<sup>19</sup> The XRD patterns (Fig. 1c) show that, over long periods of degradation, crystalline domains of metallic lead (Pb<sup>0</sup>) begin to form, which first become discernable at the 13-hour timepoint. Aside from small peaks associated with PbI<sub>2</sub> (001-plane reflection at ~12.6°) and  $\alpha$ -FAPbBr<sub>3</sub> (001-plane reflection at ~14.8°) at very late stages of degradation, no other solid degradation products are observed in the XRD patterns (zoomed-in plots and reference patterns are presented in SI section 3 and Fig. S4). We also confirmed the presence of Pb<sup>0</sup> in degraded films with x-ray photoelectron spectroscopy (XPS) (Fig. S5). The morphology, while minimally affected during early degradation times, becomes riddled with micron sized features during later times (Fig. 1d), and the changes become most obvious on similar timescales as those for which Pb<sup>0</sup> begins to appear in the XRD pattern.

The absorbance measurements appear to be the most sensitive for identifying this decomposition mode (Fig. 1a), as indicated by the increase in sub-bandgap absorbance within the first 30 minutes

(inset of Fig. 1a). The only plausible species that could absorb sub-bandgap photons is a reduced lead species. XRD and XPS provides insight that  $\text{Pb}^0$  is present at advanced stages of degradation, but we cannot preclude the presence of other reduced lead species, especially at earlier degradation times.<sup>20</sup> After 120 minutes of exposure, the excess absorbance spectrum (spanning regions both above and below the perovskite bandgap) closely resembles that of a ~5 nm film of elemental  $\text{Pb}^0$  (see Fig. S6 for the absorption spectrum of  $\text{Pb}^0$  thin films). In summary, the optoelectronic properties of the perovskite film degrade significantly as defects associated with reduced lead species are generated during LID. Notably, substantial losses in PL intensity occur well before any significant changes in morphology or crystal structure are detectable, highlighting the importance of optical measurements as essential early indicators of this degradation pathway. The QFLS calculated from the PL quantum yield decays almost linearly, and while a solar cell using the films under study as absorbers would limit degradation by trapping decomposition products and extracting reactive charge carriers, reduced lead species formed during LID would deteriorate photovoltaic performance via increased trap assisted recombination (resulting in  $V_{\text{oc}}$  decline) and parasitic absorption (resulting in  $J_{\text{sc}}$  decline).<sup>21</sup>



**Figure 1. *Ex-situ* characterization during degradation under 1 sun photon flux from a broad-spectrum LED at 65°C in a dry N<sub>2</sub> atmosphere.** The changes in (a) absorbance, (b) AIPL spectrum ( $\lambda_{excitation} = 532\text{ nm}$ ; 1 sun equivalent photon flux), (c) XRD pattern, and (d) morphology (SEM images). Inset in panel (a) is a magnified view of the sub-bandgap region. Insets in panel (b) are the *PLQY* and *QFLS* determined from the Ross relation (refer to SI section 1 for these calculations). Special care is taken to ensure that samples are outside of a controlled N<sub>2</sub> environment for as little time as possible (typical ambient lab conditions are ~40% relative humidity, 20°C; exposure during measurements is <30 minutes for XRD, <10 minutes for UV-Vis-NIR, <10 minutes for SEM, and 0 minutes for AIPL). SEM measurements are taken on separate samples to avoid conflating morphological effects of prolonged periods of vacuum exposure with LID.

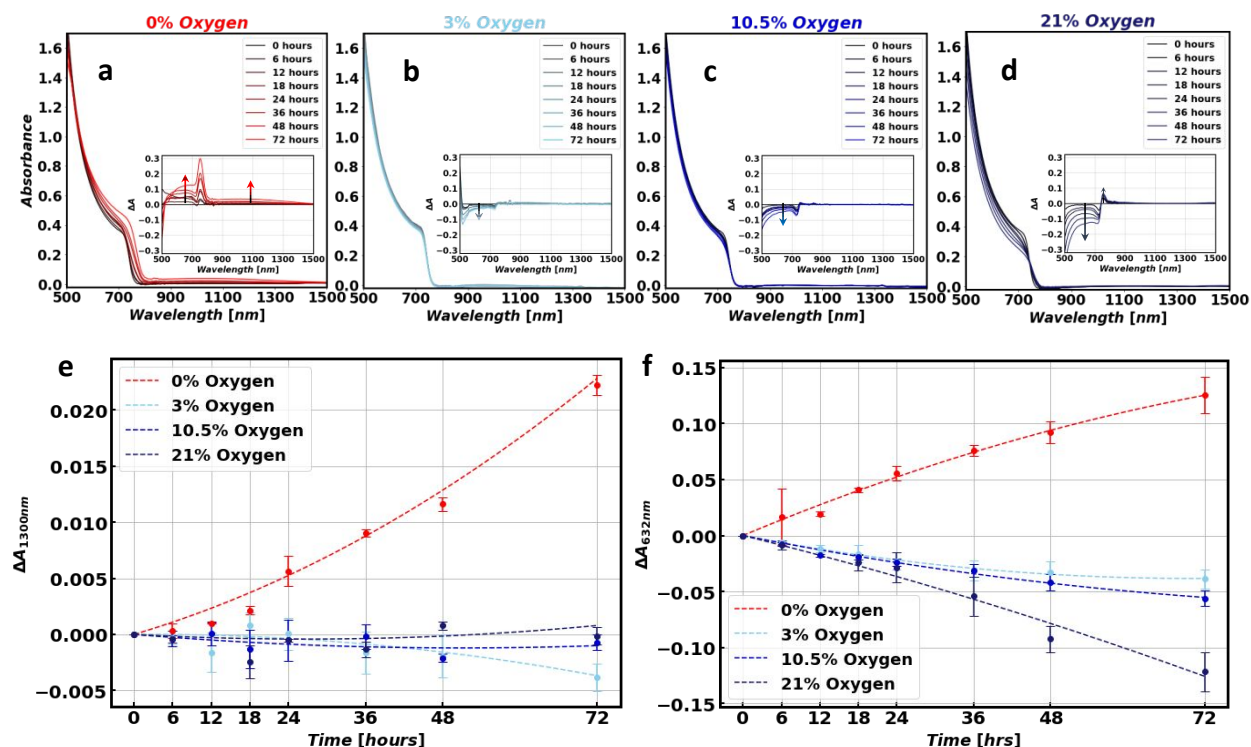


*Mechanism of Light-Induced Degradation.* The first reports of LID were of MAPbI<sub>3</sub> under vacuum.<sup>22,23</sup> However, FA-rich compositions are also susceptible to LID<sup>24</sup> with several reports observing hole trapping at grain boundaries which results in halide gas loss (X<sub>2</sub>) and Pb<sup>2+</sup> reduction to Pb<sup>0</sup>.<sup>25–27</sup> Material quality plays an important role in both initial optoelectronic device performance and stability.<sup>28,29</sup> Defect-dense and/or heterogenous films have reduced barriers for atomic, ionic, and molecular transport which increases the rate of material decomposition and device performance degradation.<sup>30–32</sup> While many of these prior reports suggest that Pb<sup>2+</sup> in the perovskite lattice is directly reduced to Pb<sup>0</sup> during LID, other reports argue that the photolysis of intermediate PbI<sub>2</sub> is the pathway to form Pb<sup>0</sup> (some noting that it is possible<sup>33</sup>).<sup>20,21,34</sup>

Kerner et al.<sup>35</sup> hypothesized that halide oxidation ( $h^+ + X^- \rightleftharpoons X^{\cdot} \rightleftharpoons \frac{1}{2} X_2$ ) by valence band holes (photoexcited or injected) can substantially increase halide species mobility (neutral X<sup>·</sup> and X<sub>2</sub> are assumed to be much more mobile than perovskite bound X<sup>−</sup>) which can yield two related but distinct phenomena. (1) If the halides settle in different domains, phase segregation results. The perovskite lattice remains intact, but different domains of varying halide concentration exist. Evidence of this behavior is typically observed with photoluminescence and electroluminescence spectroscopy<sup>36–38</sup>, and it is usually reversible after the incident light is removed.<sup>19</sup> (2) If halides escape the perovskite lattice in the form of diatomic halide gases (X<sub>2</sub>), decomposition occurs as the perovskite is irreversibly converted to other chemical species.<sup>39,40</sup> It is worth noting that the thermodynamics of halide vacancies<sup>41</sup> and the kinetics of species diffusion<sup>42,43</sup> are temperature dependent with the latter following Arrhenius behavior. These prior reports provide insights into the chemical processes that may be used to propose specific reaction pathways that in turn can be used to derive rate expressions for this decomposition mode. Given that many other studies report

the formation of  $\text{Pb}^0$  during LID, we operate under the assumption that the majority of the light-induced defects are  $\text{Pb}^0$ , accounting for most of the near-infrared light absorption.

*Influence of Oxygen on LID.* Before proceeding with further investigations into the LID of FACsPbIBr films, we seek to address how oxygen presence influences the degradation pathway. While  $\text{O}_2$  and  $\text{H}_2\text{O}$  vapor may be kept out of a PSC during fabrication and for some time with encapsulation, even rigorous glass-glass encapsulation strategies with polyisobutylene (PIB) edge seals cannot fully prevent diffusion of gaseous small molecules.<sup>44,45</sup> It is thus important to understand how degradation will proceed if an environmental reactant diffuses through packaging to meet the perovskite layer. To do so, we analyzed the change in optical absorbance when films were subjected to various concentrations of ambient  $\text{O}_2$  (0%, 3%, 10.5%, and 21%). In these experiments, we placed FACsPbIBr films in environmental control chambers and illuminated them with a broad-spectrum white LED (Fig. S3a). We measured the surface temperature with a contact thermocouple and found it to be stable at 35°C. The flow rate of inlet gases is 3 L/min with the desired oxygen concentration being achieved by mixing dry air and dry  $\text{N}_2$ .



**Figure 2.** *Ex-situ* UV-Vis-NIR measurements during degradation at 35°C under 1 sun equivalent illumination. Film absorbance over the course of degradation in dry atmospheres of (a) 0%, (b) 3%, (c) 10.5%, and (d) 21% oxygen. Insets in a-d are the change in absorbance from the initial time point with narrower vertical axes to emphasize the spectral changes over time. The timeseries absorbance values at (e) 1300 nm and (f) 632 nm. Parabolic fits to the timeseries data are included as dashed lines in (e) and (f) to guide the eye. Uncertainty is quantified by the standard deviation ( $\sigma$ ) of film absorbance change for 3 samples.

The results of these experiments are shown in Figure 2, demonstrating that 3%  $O_2$  in the ambient atmosphere is sufficient to prevent the accumulation of defects with near-infrared absorptivity ( $Pb^0$ ). Although none of the samples degraded in oxygen-containing environments accumulated detectable levels of  $Pb^0$ , we infer that photooxidation processes did take place as evidenced by decreases in above bandgap absorbance, with higher oxygen concentrations yielding higher rates of decomposition. We speculate on the reason for  $O_2$  preventing  $Pb^0$  formation in SI section 8 and

hypothesize that it scavenges photogenerated electrons that would otherwise reduce  $\text{Pb}^{2+}$ . Interestingly, examining Figure 2a and 2b reveals that the net effect of adding small amounts of oxygen to an otherwise inert atmosphere is a slower rate of film decomposition. The inset in Figure 2a reveals an outsized increase in absorption as LID proceeds in the 700-800 nm region. This behavior is observed to a lesser degree for degradation in 21% oxygen (inset of Fig. 2d). We attribute this to an increase in disorder of the semiconductor (either static or dynamic) resulting in increases in sub bandgap absorption in the region just below the bandgap (within  $\sim 0.1$  eV).<sup>46,47</sup> As such, we do not observe measurable absorption increases in this region for the films degraded in 3% or 10.5% oxygen atmospheres because they degrade less significantly (and thus induce less additional disorder) than the films in 0% or 21% oxygen atmospheres. Importantly, the excess absorption beyond  $\sim 800$  nm shown in the inset of Figure 2a agrees with the  $\text{Pb}^0$  absorption spectrum shown in Figure S6, indicating that absorption measurements in the region  $> 800$  nm are accurate for quantitative measurements of  $\text{Pb}^0$  formation.

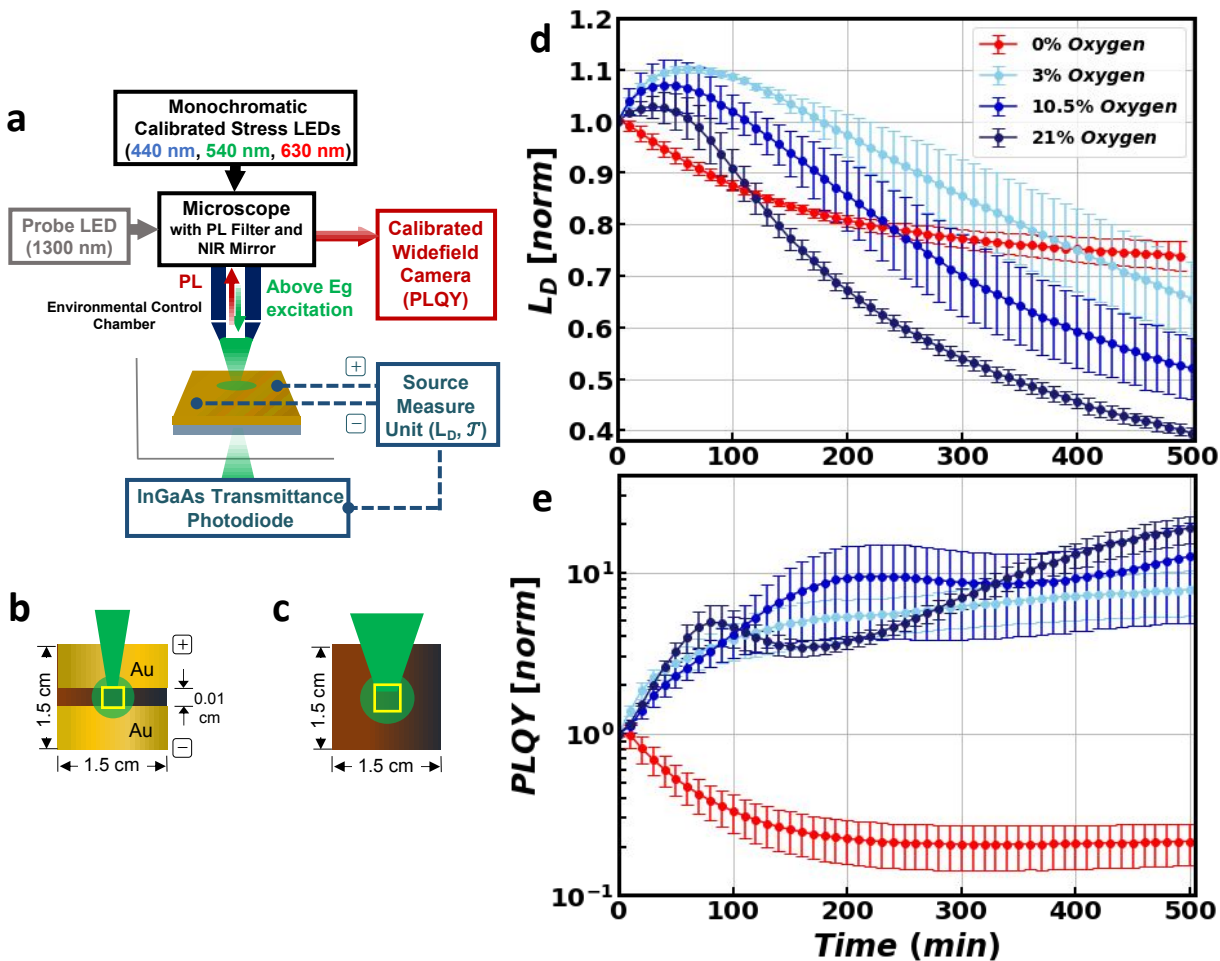
While film absorbance measurements provide valuable information about absorbing species present in the film, the exact effect of any given degradation pathway on device relevant optoelectronic properties is unknown from absorption measurements alone. To address this missing link, we conducted *in-situ* measurements of film ambipolar diffusion length ( $L_D$ ) and photoluminescence quantum yield ( $PLQY$ ) over the course of degradation in the same conditions as the *ex-situ* experiments shown in Figure 2. To perform these measurements, we use a calibrated widefield microscope for  $PLQY$  measurements and deposit gold contact pads on the FACsPbIBr films with a thin channel (0.01 cm) for simultaneous two-point photoconductivity measurements (Fig. 3a,b). From photoconductivity measurements, we can obtain an estimate for the  $L_D$  based on sample geometry and photogeneration rate, which we have detailed in prior reports.<sup>48,49</sup> The results

of these experiments are shown in Figure 3. Diffusion length is a good proxy for how we expect the film degradation to affect device performance as it unites both carrier transport and carrier lifetime. The *PLQY* serves as a proxy for carrier lifetime alone, which helps disentangle changes in diffusion length due to lifetime from those caused by changes in carrier mobility.

Degradation in pure  $N_2$  results in monotonic decreases in both  $L_D$  and *PLQY*. The evolution of the  $L_D$  is strongly correlated with that of the *PLQY*, suggesting that it may be almost entirely explained by  $Pb^0$  trap state formation. In contrast, the films placed in oxygen containing atmospheres see an initial period of passivation (increasing  $L_D$  due to increasing minority carrier lifetime) followed by degradation (decreasing  $L_D$  attributable to reduction in effective mobility that likely arise from the buildup of insulating phases at grain boundaries). The kinetics of this behavior accelerate as the environmental oxygen content increases, consistent with a photooxidation mechanism that generates insulating byproducts that impede carrier transport.<sup>48,50</sup> Note that the *PLQY* generally rises for each of the oxygen containing runs, indicating that the local minority carrier lifetime is constant or increasing and suggesting that defect formation is not a critical issue in these cases. While these photooxidation phenomena are not unique to these perovskite compositions and have been reported previously for  $MAPbI_3$ ,<sup>42</sup> the stark differences in optoelectronic decay from relatively small differences in ambient atmosphere contents for this perovskite composition is important because it indicates that relatively small amounts of oxygen are sufficient to shift the dominant mode of decomposition. Again, this may signify a role of oxygen as a particularly efficient scavenger of photogenerated electrons.<sup>51</sup> From the data shown in Figures 2 and 3, we infer that there is likely an oxygen concentration between 0% and 3% at which both mechanisms occur simultaneously but that concentrations  $\geq 3\%$  appear to prevent  $Pb^0$  formation and are within the regime of purely photooxidation decomposition (all  $Pb^{2+}$  remains

oxidized). However, it is important to note that even for low-oxygen environments in which LID is suppressed, apparent mobility losses can still quickly reduce the perovskite’s optoelectronic quality, highlighting the importance of retarding photooxidation as well as LID.

To assess and compare the relevant time scales of diffusion length decay among the different degradation conditions, we define metrics of  $TL_D$ -80 and  $TL_D$ -90 which are defined as the time it takes for the diffusion length to decrease to 80% and 90% of its original value, respectively. The average  $TL_D$ -80 and  $TL_D$ -90 times ( $\pm \sigma$ ) are displayed in table 1. Importantly, the films placed in 0%  $O_2$  reach  $L_D$ -90 faster on average than the films placed in any of the oxygen containing environments, highlighting the rapid optoelectronic performance decline that occurs during LID.



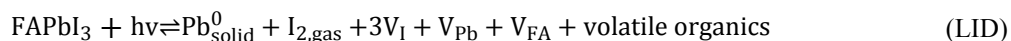
**Figure 3. *In-situ* optoelectronic performance trends for degradation in atmospheres of varying oxygen concentration.** (a) Schematic of the *in-situ* microscope degradation experimental apparatus which features a calibrated LED light source fed through an upright microscope and shone through an objective lens onto the sample. Top-down view of perovskite samples with (b) gold contacts (used for simultaneous collection of  $PLQY$  and  $L_D$ ) and (c) bare (used for  $T_{1300\text{ nm}}$  measurements). The normalized (d) ambipolar diffusion length and (e) photoluminescence quantum yield under 10 sun equivalent photon flux at 65°C. Uncertainty is quantified by the standard error of the mean ( $\sigma/\sqrt{n}$ ) for 3 samples.

**Table 1.  $TL_{D-80}$  and  $TL_{D-90}$  Times for Degradation at 65°C, 1 sun under Different Ambient Oxygen Concentrations**

[O <sub>2</sub> ]	$TL_{D-90}$ [min]	$TL_{D-80}$ [min]
0%	78 ± 20	293 ± 179
3%	285 ± 95	363 ± 115
10.5%	186 ± 74	242 ± 85
21%	97 ± 22	132 ± 20

*Measuring the Kinetics of Light-Induced Defect Formation.* Having established that light-induced defect formation dominates the optoelectronic evolution of FACsPbIBr films in an inert N<sub>2</sub> atmosphere under illumination, we proceed with the aim of measuring and predicting the formation rate of these defects (with the assumption that most are Pb<sup>0</sup>). We propose an overall net reaction for this process that uses FAPbI<sub>3</sub> as an approximation of the mixed-cation, mixed-halide composition. We do not expect the Cs or Br minorities to participate in important ways. To test this assumption, we also test the photostability of FA<sub>0.8</sub>Cs<sub>0.2</sub>PbI<sub>3</sub> and FAPbI<sub>3</sub> thin films (Fig. S8).

We observe that the absorbance of both compositions increases in the sub bandgap region after a period of photo-stress, indicating the universal susceptibility to LID of FA-rich compositions. Notably, the rate of sub-bandgap absorption increase is slower for these compositions compared to mixed-halide FACsPbIBr which agrees with the finding that halide alloying with bromine yields less photostable materials.<sup>52</sup> Thus, the net total LID reaction can be written as the following:



Section 5 of the SI discusses a plausible complete reaction mechanism with elementary steps for the LID reaction. We arrive at a functional form for a mathematical expression to describe the initial rate of  $\text{Pb}^0$  formation ( $r_{\text{Pb}^0}$ ) as a function of temperature and photon flux. Note that this equation also happens to be identical to an empirical expression that could be written for the total LID reaction.

$$r_{\text{Pb}^0} = k_{0,\text{LID}} \exp\left(-\frac{E_{\text{A,LID}}^{\text{eff}}}{k_{\text{B}}T}\right) I_{\text{in}}^n \quad (1)$$

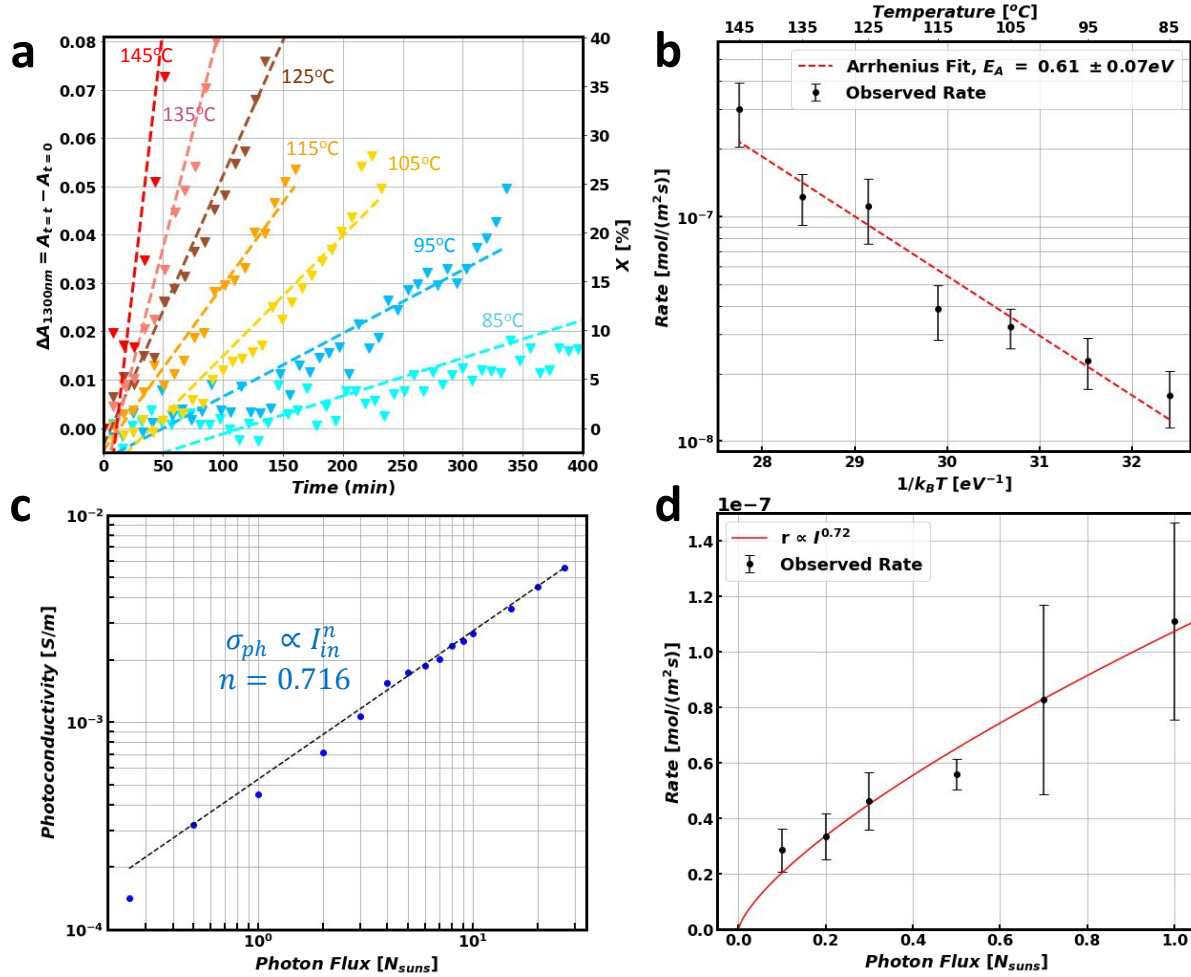
Here,  $k_{0,\text{LID}}$  is a temperature independent prefactor,  $E_{\text{A,LID}}^{\text{eff}}$  is the effective activation energy,  $I_{\text{in}}$  is the incident above bandgap photon flux, and  $n$  is the effective order of illumination.

We next proceed to measure the kinetics of the LID reaction to determine the various unknown parameters contained within equation 1 for FACsPbIBr. We use a similar technique as in our previous kinetic studies<sup>17,18</sup> on initial perovskite decomposition rate to determine the rate of  $\text{Pb}^0$  formation ( $r_{\text{Pb}^0}$ ). Information regarding the initial perovskite reaction rate has been shown to be highly relevant for understanding and predicting device stability.<sup>48,53</sup> In short, we measure the sample transmittance ( $\mathcal{T}$ ) at 1300 nm (where reflectance changes minimally over the early times of degradation, Fig. S7) in an *in-situ* degradation chamber (Fig. 3a,c) to obtain an estimate for initial rate of absorbance change  $\left(\frac{dA}{dt}\right)$ . This allows us to quantify  $r_{\text{Pb}^0}$  by the following:



$$r_{Pb^0} = \frac{\rho_{Pb^0} \cdot \log_e(10)}{M_{Pb^0} \cdot \alpha_{1300nm}} \cdot \frac{dA}{dt}$$

where  $\rho_{Pb^0}$  is the density,  $M_{Pb^0}$  is the molecular weight, and  $\alpha_{1300nm}$  is the absorption coefficient at 1300 nm for  $Pb^0$ . Refer to SI section 4 for a detailed discussion of this derivation.



**Figure 4. Determining model parameters.** (a) Change in absorbance over time (with linear fits to the data used to determine  $\frac{dA}{dt}$ ) during degradation with 1 sun photon flux. Percent film conversion to  $Pb^0$  ( $X$ ) is displayed on the right vertical axis. (b) Arrhenius plot for degradation experiments under 1 sun illumination at varying temperatures. (c) Photoconductivity versus photon flux with sublinear fit. (d) Initial  $Pb^0$  formation rate as a function of incident photon flux at 125°C with fit (exponent fixed to 0.72).

**Table 2. Best-Fit Parameters for the Light-Induced Pb<sup>0</sup> Formation Kinetic Model**

Parameter	Value	Unit
$E_{A,LID}^{eff}$	$0.61 \pm 0.07$	eV
$k_{0,LID}$	$(3.55 \pm 0.58) \cdot 10^{-15}$	$\text{mol} \cdot \text{m}^{-2} \cdot \text{s}^{-1} \cdot (\text{photons} \cdot \text{m}^{-2} \cdot \text{s}^{-1})^{-0.72}$
$n$	0.72	unitless

To determine parameters that define equation 1 ( $E_{A,LID}^{eff}$ ,  $n$ ,  $k_{0,LID}$ ), we used monochromatic 540 nm illumination over a range of temperature and illumination conditions. We choose to measure the rate of degradation at elevated temperatures (85 -145°C) due to the relatively slow rate of this reaction. We observe an Arrhenius relationship with respect to the effect of temperature on  $r_{Pb^0}$  which suggests that this reaction is thermally activated. Thus, we expect the experimentally determined parameters at elevated temperatures to remain accurate across a wide range of temperature and illumination conditions. We demonstrate that decomposition due to elevated temperature exposure, resulting in apparent loss of formamidineum, is negligible in comparison to LID in SI section 9 (particularly with regards to its effects on film optical properties).

Figure 4a displays typical data for the change in 1300 nm absorbance over time with linear fits to that data. Degradation experiments from 85-145°C under 1 sun equivalent photon flux are used to determine the activation energy as ~0.61 eV (Fig. 4b). To determine the effective order of above bandgap photons ( $n$ ), we performed photoconductivity measurements as a function of incident photon flux (Fig. 4c). We have shown previously that the activity of a photoexcited species is directly proportional to the incident above bandgap photon flux raised to a constant value.<sup>17</sup> Additional information on the relationship between photoconductivity and photoexcited species activity is included in SI section 6. Here, we determine this value to be ~0.72 when measuring across a range of photon fluxes (0.25 to ~29 suns). To confirm that this exponent accurately

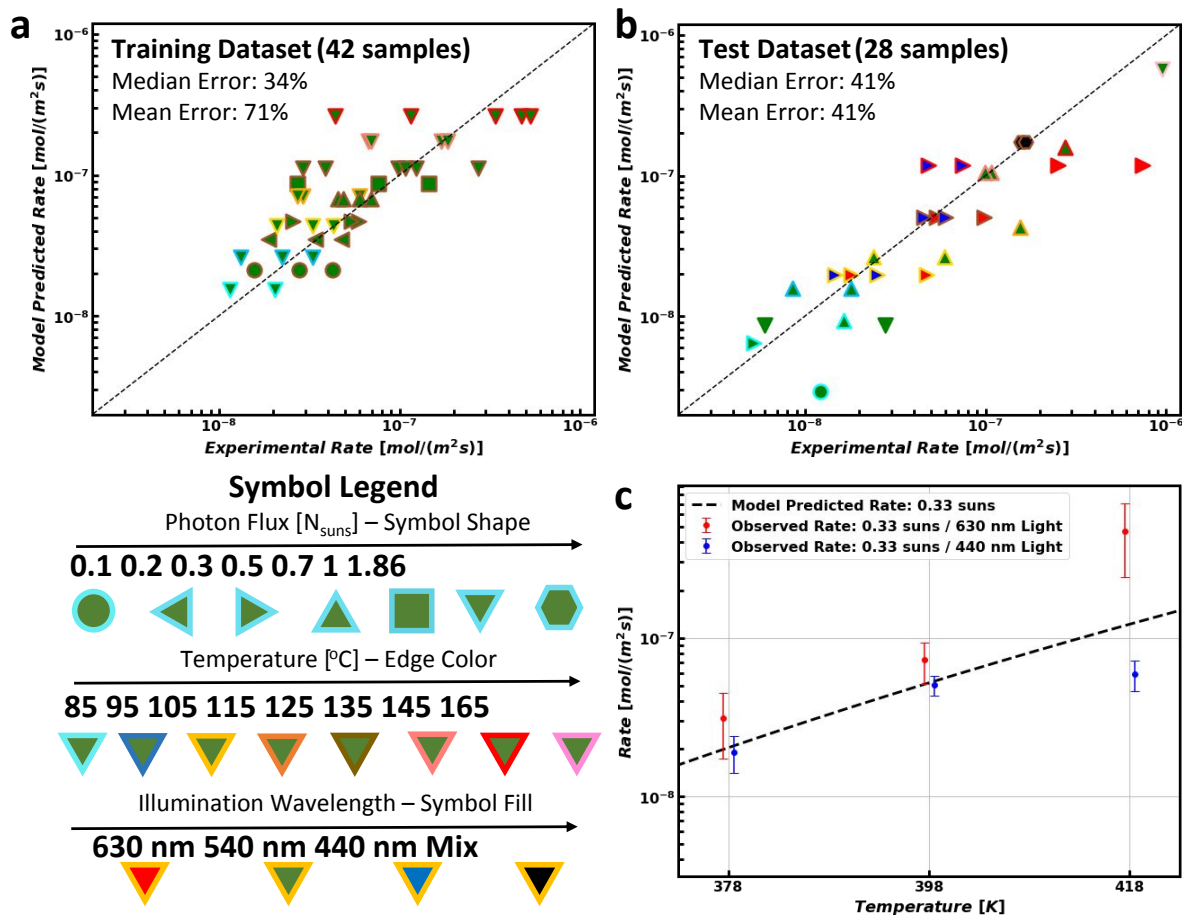
captures the change in hole activity (posited to be the relevant photoexcited species) during light-induced degradation, we measure  $r_{Pb^0}$  as a function of photon flux from 0.1-1 suns at 125°C (Fig. 4d). A fit of the data is excellent (refer to SI section 7 for further fitting details), indicating that we have accurately captured the incremental activity of holes for increases in photon flux with our power law model.

To determine  $k_{0,LID}$ , we take all degradation runs used to determine  $E_{A,LID}^{eff}$  (1 sun, 85-145°C, 540 nm) and all runs used to confirm  $n$  (125°C, 0.1-1 sun, 540 nm) and perform a minimization on the sum of the squared errors, which gives a result of  $3.55 \times 10^{-15} \frac{mol}{m^2 s (photons \cdot m^2 s)^{0.72}}$ . The dataset used for this fitting is referred to as the ‘training dataset’ (42 samples). The best-fit parameters for the kinetic rate model for  $Pb^0$  formation are displayed in Table 2. Note that the uncertainty estimates for  $E_{A,LID}^{eff}$  and  $k_{0,LID}$  are quantified by the standard error of the linear fits to the Arrhenius plot and to the complete training dataset, respectively.

As mentioned previously, the debate around  $PbI_2$  photolysis remains unanswered in the perovskite community. First, we note that significant  $Pb^0$  is formed during irradiation with 540 nm light ( $E_{photon}=2.30$  eV) which itself cannot excite  $PbI_2$  by a band-band transition since it has an energy below the bandgap of  $PbI_2$  ( $E_g \sim 2.35$  eV). Thus, while  $PbI_2$  photolysis may occur in perovskite samples with excess  $PbI_2$  impurities, we note that this is not necessary to induce  $Pb^0$  formation in our films and is not the mechanism that we observe since  $Pb^0$  is formed even when the incident beam has photon energies less than the bandgap of  $PbI_2$ . To confirm this finding, we collected a new dataset (referred to as the ‘test dataset’ (28 samples)) that incorporates alternate temperatures and illumination conditions for 540 nm illumination as well as 440 nm (well above  $PbI_2$  bandgap) and 630 nm illumination (well below  $PbI_2$  bandgap). SI section 7 provides further

1 explanation of the conditions for the testing and training datasets and an expanded discussion of  
2 the fitting procedure for the parameters of equation 1.

3 Figure 5a shows the results for the training dataset with median and mean error of 34% and 71%,  
4 respectively. Figure 5b shows the results for the test dataset with median and mean error of 41%  
5 and 41%, respectively. Despite some sample-to-sample variation, we see that the model  
6 determined from the training dataset, which only considers a relatively small sliver of the  
7 parameter space, can accurately predict  $r_{Pb^0}$  for varying temperatures, photon fluxes, and photon  
8 energies. We emphasize that illumination wavelength within the visible regime minimally  
9 influences the degradation rate. Figure 5c shows the degradation data for the 630 nm and 440 nm  
10 illumination. This plot further emphasizes the goodness of the model fit and suggests that, if  
11 anything, the lower energy 630 nm light produces higher degradation rates than 440 nm. These  
12 results highlight that  $PbI_2$  photolysis is not the mode of degradation that we observe en route to  
13  $Pb^0$  formation and suggest that degradation mitigation strategies outside that of the  $PbI_2$  photolysis  
14 suppression framework must be considered to enhance perovskite longevity.



**Figure 5. Light-induced  $\text{Pb}^0$  formation rate model performance.** (a) Parity plot for the training dataset and (b) for the test dataset. (c) Comparison of 630 nm monochromatic illumination (red) and 440 nm monochromatic illumination (blue) on  $\text{Pb}^0$  rate. 630 nm and 440 nm photon fluxes are both set equal to 0.33 suns ( $5 \times 10^{20}$  photons/ $\text{m}^2/\text{s}$ ). Datapoints are slightly offset along the temperature axis for clarity.

Mitigation strategies for suppressing halide segregation and LID center around reducing halide mobility and escaping tendency from the perovskite. Device stacks with strong encapsulation will certainly slow volatile halide species loss and thus promote reverse decomposition reactions (e.g.  $\text{Pb}^0 + \text{I}_2 \rightleftharpoons \text{Pb}^{2+} + 2\text{I}^-$ ), slowing LID processes. We demonstrate that encapsulation slows  $\text{Pb}^0$  formation in SI section 10. However, the intrinsic material instability related to mobile and reactive

ions will remain and prevent the realization of multidecade module service lifetimes if a promising solution cannot be found. Promising paths towards reducing halide movement are via passivation (which will necessarily slow diffusion due to fewer vacancies present in the lattice<sup>54,55</sup>) and by tin substitution for lead (which shifts the valence band maximum above the halide defect energies and slows halide oxidation<sup>56,57</sup>). Ultimately, mobile halides that facilitate decomposition present a significant challenge for lead-based perovskite photovoltaic commercialization. We also would like to emphasize that the trapping of volatile I<sub>2</sub> vapor can promote other decomposition pathways (not just the reverse LID reaction) in a complete device stack.<sup>58,59</sup> As such, a complete picture of PSC device degradation under inert conditions would include LID, its reverse reaction, and I<sub>2</sub> related oxidation reactions with the perovskite itself and other layers in the device stack.

In summary, we report, for the first-time, quantitative measurements of light-induced degradation in a halide perovskite semiconductor. The degradation products of LID are reduced lead species with broadband optical absorption and are very likely fully reduced Pb<sup>0</sup>. We have studied the effects of oxygen on this decomposition pathway and have derived a rate law for the initial rate of Pb<sup>0</sup> formation and fit it to our quantitative measurements. We also show that the wavelength distribution of the incident above bandgap photons negligibly influences the formation rate. This work raises important questions regarding how to address oxygen ingress through advanced encapsulation techniques, as well as how to manage the iodine species generated during LID, both of which remain significant challenges for enhancing the stability of perovskite-based devices.

## Supporting Information

Detailed experimental methods, characterization of perovskite films and devices, data on light sources used for degradation experiments, degradation product identification with XRD and XPS, further information on the estimation of initial  $\text{Pb}^0$  formation rate calculated from sub bandgap transmittance and absorption coefficient of  $\text{Pb}^0$ , derivations of a rate expression for light-induced degradation and for variation of photoactive species activity with photon flux, details of rate model fitting and testing, discussion of the role of oxygen in preventing metallic lead formation, comparison of degradation at high temperatures under illumination and in the dark, effects of encapsulation on the rate of  $\text{Pb}^0$  formation.

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#### **Conflicts of Interest**

There are no conflicts of interest to declare.

#### **Data Availability**

Data for this article are available at [https://osf.io/n7u2t/?view\\_only=4a251b2ae5f94aa0adb408efc7b0d4ca](https://osf.io/n7u2t/?view_only=4a251b2ae5f94aa0adb408efc7b0d4ca).

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