







# Intercluster charge ordering in monoclinic and triclinic Ba-Mo-based hollandite phases

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Complete List of Authors:	Elbakry, Eslam; The University of Alabama, Chemistry and Biochemistry Allred, Jared; The University of Alabama, Chemistry and Biochemistry

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# **ARTICLE**

# Intercluster charge ordering in monoclinic and triclinic Ba-Mobased hollandite phases

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Eslam M. Elbakry a and Jared M. Allred +\* a

Reproducible solid-state synthesis methods are presented for the preparation of tetragonal  $Ba_{1.1}Mo_8O_{16}$ , monoclinic  $Na_{0.325(5)}Ba_{1.006(18)}Mo_8O_{16}$ , and triclinic  $Ba_{1.12(3)}Mo_8O_{16}$  hollandite phases, with complete, high-resolution crystal structures of the monoclinic and triclinic phases reported for the first time. The similar synthetic conditions allow direct comparisons between phases; differences between structures are shown to be correlated to subtle changes in the metal-metal bonding of the  $Mo_4$  cluster motif that is unique to the Mo-based hollandites. The trends in the local Mo valence, stoichiometry, and key Mo-Mo bond lengths of these and other reported Mo-based hollandite phases together support an interchain charge ordering model for this family of compounds, which has been previously suggested for the case of  $K_2Mo_8O_{16}$ . An alternate model, where Mott physics dominate the electronic structure near the Fermi level, is not supported by temperature-dependent magnetic susceptibility measurements, which are reported down to 2 K. The incorporation and homogeneity of Na in the monoclinic phase is verified using atom probe tomography.

### Introduction

Metal–metal bonding motifs are of interest as functional structural subunits of metal–organic frameworks, molecular-scale conductors, photosensitizers, catalysts, and so on. The metallic ions of d¹, d², and d³ configurations such as Nb (IV), Mo (IV), and Re (IV), naturally favour the formation of metal-metal bonds. Molybdenum oxides in particular stand out for the delicate balance between different and sometimes coexisting oxidation states. For example, Mo(IV), Mo(V), and Mo(VI) can coexist within the same structural motif, as in La $_3\text{Mo}_4(\text{Al}_{0.67}\text{Mo}_{0.33})\text{O}_{14}.^2$ 

In compounds containing chain-like units of edge-sharing metal-oxygen octahedra, the metal-metal bond formation often becomes energetically favored. The hollandite structure type,  $A_x M_8 O_{16}$ , is one such family of compounds, which has garnered much interest due to the possibility of one-dimensional ionic conductivity through the tunnels located between cages of metal-oxygen octahedra, Figure 1.4 There is substantial current interest in hollandite compounds for their potential as solid-state battery cathodes for trapping hazardous elements such as radioactive wastes.

The strong metal-metal bond motifs that are characteristic of molybdenum-based hollandites provide a unique opportunity to explore the interplay between the oxidation state of Mo, patterns of Mo-clustering, electronic properties of

A-cations, and presence of vacant sites within the hollandite channels. For instance, La $_{1.16}$ Mo $_8$ O $_{16}$ ,  $R_{1.33}$ Mo $_8$ O $_{16}$  (R = Pr, Nd), and Rb $_{1.5}$ Mo $_8$ O $_{16}$  all have tetragonal symmetry, whereas K $_2$ Mo $_8$ O $_{16}$  is monoclinic, $^{9\text{-}11}$  which breaks the usual trend in hollandite crystal symmetry. <sup>12</sup> In Rb $_{1.5}$ Mo $_8$ O $_{16}$  and K $_2$ Mo $_8$ O $_{16}$  compounds, Mo $_4$  rhombohedral clusters form cage-like channels, inside which larger cations are trapped, while Mo $_3$  triangular clusters exist as either isolated units or within infinite chains in La $_{1.16}$ Mo $_8$ O $_{16}$  and  $R_{1.33}$ Mo $_8$ O $_{16}$  (R = Pr, Nd). <sup>9-11</sup> Recent

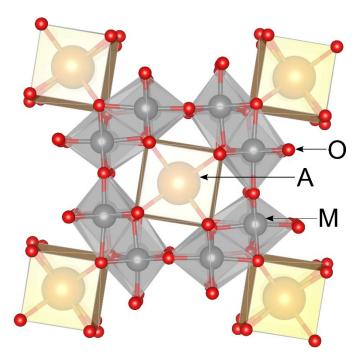


Figure 1: View of the prototypical hollandite structure along the c-axis in the tetragonal setting.

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<sup>&</sup>lt;sup>a.</sup> Department of Chemistry and Biochemistry, The University of Alabama, Tuscaloosa. Alabama 35487. United States

<sup>†</sup> jmallred@ua.edu

electronic structure calculations showed that the monoclinic distortion in  $K_2Mo_8O_{16}$  can be driven electronically: either by charge ordering between  $Mo_4$  clusters that breaks tetragonal symmetry or via strong electronic correlations via a Mott mechanism.<sup>13</sup>

Patterns of Mo-clustering are not the same in Mohollandites that have Ba in their channels. There are two reported ternary Ba-Mo-hollandites, tetragonal  $Ba_{1.14}Mo_8O_{16}$ and triclinic  $Ba_{1.13}Mo_8O_{16}$ , the latter of which being the first and only example of a triclinic distortion of the hollandite structure to our knowledge. The tetragonal form was synthesized electrochemically from a molten mixture of barium and caesium molybdates.<sup>14</sup> It contains Mo<sub>3</sub> triangles and Mo<sub>4</sub> rhomboidal clusters alternating in equal proportions, described by a crystallographic modulation  $\mathbf{q} \approx 4/7 \, \mathbf{c}^*$ . The triclinic phase only contains Mo<sub>4</sub> clusters, though it this information is only derived from a multiphase product and the complete structural details have not been reported. 15 Further complicating the picture is a provisional report of a monoclinic form of the Ba-Mo hollandite stabilized by the addition of sodium and lithium (nominally reported as Na<sub>0.35</sub>BaMo<sub>8</sub>O<sub>16</sub> and Li<sub>0.34</sub>BaMo<sub>8</sub>O<sub>16</sub>), <sup>16</sup> which were respectively indexed and solved in the monoclinic cell as K<sub>2</sub>Mo<sub>8</sub>O<sub>16</sub> and the tetragonal cell of K<sub>2</sub>Cr<sub>8</sub>O<sub>16</sub>, respectively.

Despite the unique diversity of observed crystal symmetries and structural motifs, the origin of the distinct Ba-Mocontaining hollandite polymorphs remains unknown, perhaps due to the difficulties in reproducibly synthesizing them at high phase purities. Alkali salts have previously been used as mineralizers to improve the phase purity of undoped Ba-based hollandites, such as KCl with  $\rm Ba_{1.1}Mn_8O_{16}^{17}$  which makes the role of the Na $^+$  and Cs $^+$  alkali ions in the above examples unclear. The Na atoms in the crystallographic model for monoclinic Na $_{0.35}\rm BaMo_8O_{16}$  are only four coordinate with exceptionally short Na-O bonds (2.20 Å) in a square planar geometry,  $^{16}$  which is unusual enough to require revisiting and validation. The effects of stoichiometry and temperature on phase stability are also essentially unexplored.

The lack of validated phase data also impacts broader materials science activities. The reported Ba-Mo-O phase equilibria diagrams lacked critical information regarding the formation of Ba-Mo-containing hollandite phases.<sup>18</sup>

In this work, we investigate the role of reaction conditions on the polymorph and stoichiometry of Ba-Mo-based hollandites. We develop a consistent and reproducible solidstate synthesis technique for producing high phase-purity Ba-Mo-hollandite compounds separately in all three known crystal systems: triclinic (TRI) phase  $Ba_{1.12(3)}Mo_8O_{16}$ , tetragonal (TET)  $Ba_{1.0+x}Mo_8O_{16}$ phase(s), and a monoclinic  $\text{Na}_{0.324(5)}\text{Ba}_{1.006(18)}\text{Mo}_8\text{O}_{16}$  (or  $\text{Na}_{1/3}\text{Mo}_8\text{O}_{16})$  phase. The findings support treating TRI and MON as essentially line compounds with fixed stoichiometry, while **TET** is found to form as a series of at least five distinct modulated structures, only of which has been previously reported. Extensive optimization in reactant choice, presence and type of mineralizer, and quenching rate reveal what factors control the polymorph type. Atom probe tomography (APT) reveals that the monoclinic symmetry is correlated to the presence of Na<sup>+</sup> ions in the channels.

High-resolution crystal structure solutions of **TRI** and **MON** compounds are also reported. The resulting trends show trends in the metal-metal bonding and local metal atom valence that support an electronically driven charge-ordering model for the structural distortion in **MON**, **TRI**, and  $K_2Mo_8O_{16}$  by extension.

#### Experimental

#### **Materials**

Starting materials include powders of BaCO $_3$  (Strem chemicals, 99.9%), Na $_2$ CO $_3$  (VWR, 99.9%), K $_2$ CO $_3$  (Sigma Aldrich, 99.0%), CaMoO $_4$ , BaMoO $_4$ , Na $_2$ MoO $_4$  (Beantown Chemical, 99.9%), MoO $_2$  powder (Prochem inc., 99.95%), MoO $_3$  (Alfa Aesar, 99.95%), elemental Mo, NaCl (Beantown Chemical, 99.0%), Nal, KI (Alfa Aesar, 99.0%) and MoCl $_3$  (Alfa Aesar, 99.5%). Carbonates were dried in air at 250 °C overnight and MoO $_2$  was further purified by sequential washing with 1 M hydrochloric acid and potassium hydroxide solutions to remove other molybdenum oxide phases present.

CaMoO $_4$  and BaMoO $_4$  were prepared by solid-state reaction between MoO $_3$  and CaCO $_3$  or BaCO $_3$ , respectively, at 800 °C for 24 hr. Elemental Mo was obtained *via* the reduction of MoO $_2$  using 5% H $_2$ /Ar gas, at 900 °C, for 48 hr. The purity of all reagents was confirmed using powder X-ray diffraction (PXRD) with no impurities detectable above a 2% threshold.

Prior to use, copper tubes (McMaster Carr, Superconductive 101 copper) were heated to 900  $^{\circ}$ C under 5% H<sub>2</sub>/Ar flowing gas to remove the oxide surface.

## Synthetic procedures

Starting materials in amounts of 0.250 mmol (about 300 mg) were finely ground using an agate mortar and pestle. Unless otherwise stated, the mixture was placed in a cylindrical copper tube with 0.08 cm thick walls and having the approximate dimensions of 2 cm length and 0.79 cm inner diameter. The copper tubes were sealed under a partial atmosphere of 50 kPa flowing Ar using an electric arc welder, which was then sealed inside of an evacuated quartz ampoule prior to reaction. Since the melting point of copper is 1085 °C, samples (re)fired at temperatures higher than 1000 °C were wrapped in nickel foil. Niobium tubes resulted in the formation of niobium oxides and Mo metal, so they were not used.

Unless otherwise specified, reaction vessels were heated at a rate 83 °C/hr up to 1000 °C and fired for 4 days. Cooling rates were varied between reactions. Rapid quenching steps were performed by moving the reaction vessel from the furnace at 1000 °C directly to a water bath at room temperature.

# X-ray diffraction

Samples were ground to powder to carry out PXRD. The PXRD patterns were measured on a Bruker D2-Phaser diffractometer equipped with a Cu-K $\alpha$  source and a Lynx-Eye detector. The lattice parameters were determined for each composition by fitting the hollandite crystallographic models to the PXRD patterns in the Rietveld method implemented through GSAS II  $_{19,\,20}$ 

Single crystal diffraction experiments were performed using a Rigaku Synergy-DW, equipped with a HyPix detector and

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operated in the Mo-K $\alpha$  mode. The refinement, absorption correction, and data reduction of the Bragg peak intensity data obtained from single crystal X-ray diffraction were performed using CrysAlisPro software and the crystal structures were refined using SHELX, WinGX, and OLEX2 programs  $^{21-23}.$  Visualizations of crystal structures were carried out with VESTA.  $^{24}$ 

Bond valence sum parameters for Mo were taken from Chen and coworkers<sup>25</sup>. These values provide the most accurate description of extended molybdenum oxide solids such as MoO<sub>2</sub>, which is the closest binary to the studied hollandite compounds in terms of both structure and stoichiometry.

#### Atom probe tomography

A three-dimensional compositional map for the Na-doped sample was constructed by atom probe tomography (APT) technique. This was performed on a LEAP 6000 XR operated in laser mode with a wavelength of 267 nm. Prior to sample preparation, the sample was washed with dilute HCl to remove Na-containing impurities. Atom probe samples were prepared as needle-shaped specimens using a Tescan Lyra FIB-FESEM instrument. Dual-beam FIB protocol was applied <sup>26</sup>. FIB technique was utilized to prepare sample needles, where the needle direction was perpendicular to the z-direction of the sample crystals, at a shank angle ranging from 3° to 5°. 16 million ions were detected, using pulse rate and energy at 67 kHz and 30 PJ, respectively. The direct current (DC) voltage varied freely to maintain a stable detection rate. Data analysis was performed using the IVAS 6.3 software. Spatial distribution maps (SDM) were constructed by combining two radial distribution functions: one represents the distribution of interatomic separations in the z-direction and the other is a 2D map describing the atomic positions in the corresponding x-y lateral plane.<sup>27</sup> The composition of the **MON** specimen was studied down to about 500 nm deep, and the top 50 nm was excluded since the sample surface appears to be oxidized.

# **Magnetization measurements**

The temperature-dependent magnetization measurements were performed using a Quantum Design Dynacool Physical Property Measurement System (PPMS) equipped with a 9 Tesla magnet. A group of randomly oriented single crystals of the

same batch with a total mass of 10–25 mg was wrapped together in a plastic film and confined in a plastic straw.

# Electron microscopy, thermogravimetric analysis, and atomic emission spectroscopy.

Scanning electron microscope (SEM) images was collected using an Apreo FE SEM at an energy of 20.0 kV. Thermogravimetric analysis was performed using a Simultaneous Thermal Analyzer 8000 (PerkinElmer Inc.). Heating was carried out from 30 to 800  $^{\circ}\text{C}$  in the  $N_2$  atmosphere, at a 20  $^{\circ}\text{C/min}$  heating rate. Stoichiometries of the synthesized compounds were determined by inductively coupled plasma atomic emission spectroscopy (ICP-AES) technique, using the Agilent 5800 ICP-OES instrument.

#### Results

In this work, compounds are only referred to by their specific stoichiometries when measured experimentally here or when referencing a reported phase from prior work. This is done to avoid confusion between nominal and actual stoichiometry and further conflation between different crystallographic phases with similar stoichiometry.

Each individual hollandite phase synthesized is referenced by the crystal system that it belongs to, with the terms **TET**, **MON**, and **TRI** each denoting a compound with the previously referenced tetragonal, monoclinic, and triclinic structure types. The crystal system is chosen for convenience because the PXRD patterns are easily distinguished from each phase. This is illustrated in Figure 2, which shows simulated PXRD patterns for each compound derived from the single crystal diffraction results of **TRI** and **MON** reported in section 3.3, and of **TET** as reported by Barrier. Single crystal diffraction experiments were used to confirm the crystal system of specimens as well.

#### **General synthetic observations**

Alternatives to sealed ampoule synthesis methods generally failed in this study. For example, open alumina crucibles under flowing inert gases did not react at 600 °C. At 800 °C, volatile  $MoO_3$  crystals form under inert gas flow while reducing conditions tend to reduce Mo past the range needed here.

Hollandite phases are successfully prepared when the

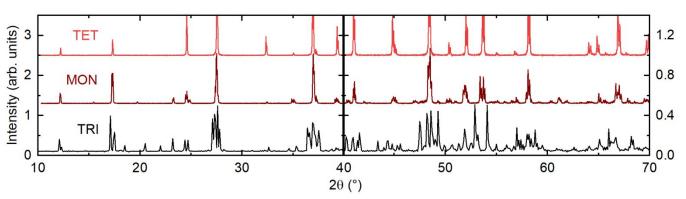


Figure 2: Simulated PXRD patterns of TRI, TET and MON phases using single crystal structure solutions. High angle (right panel) and low angle (left panel) of scanning are plotted on different scales for clarity.

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reaction mixture was sealed in copper tubes and heated at 1000 °C, in keeping with prior literature to prepare molybdenum hollandites. <sup>16</sup> Refiring the product at a higher temperature, up to 1150 °C, increases impurity fractions and does not enhance the crystallinity of the target phase(s).

The purity and form of reaction products were found to be strongly affected by the identity of reactant species.  $MoO_2$  is a common ingredient in all cases, since it gave a higher yield of products, compared to  $MoO_3$  + Mo as precursor compounds. Elemental Mo is required to balance the metal-oxygen stoichiometry, *i.e.*, as the main reducing agent for Mo (IV) and (VI) oxides. Thus, to prepare the Ba-Mo-hollandite phases,  $MoO_2$  and Mo were mixed with  $BaCO_3$  or  $BaMoO_4$ , as a reaction mixture. The reduced molybdenum oxide phases  $Ba_{0.62}Mo_4O_6$ ,  $Ba_3Mo_{18}O_{28}$  and  $BaMo_6O_{10}$  could not be distinguished from each other as minority phases, and so they are collectively referred to as  $BaMo_6O_{10}$  throughout.

The ensuing synthetic results focus on clearly showing what conditions stabilize each of the three known structure types of Ba-Mo-based hollandites. The section is organized to describe the four major factors found to affect the outcome of the reaction: 1) the effects of metal source type (e.g. BaMoO<sub>4</sub> vs. BaCO<sub>3</sub>), 2) input stoichiometry, 3) quenching rate, and 4) whether and how a mineralizer is used. Finally, the synthetic conditions specifically resulting in the **MON** phase formation are described last as it builds on findings from the other steps.

In these sections, the results of each reaction are discussed with regards to the hollandite phase's crystal size and morphology, diffracted crystal symmetry, yield, and type of impurities present. More detailed characterizations of selected products including specimen stoichiometry, crystal structure determination, and atomic-scale microscopy, are reported in their respective sections below.

#### Effects of Ba-source without a mineralizer

The first set of reactions reported are for those with no kind of mineralizer or dopant is added, which nearly always produce small single crystals of **TRI** phase, though **TET** phases are sometimes observed as a secondary phase.

Trials using  $BaMoO_4$  and  $BaCO_3$  as a Ba source distinguish between the different effects introduced by a carbonate reactant and 42 °C/hr cooling rate. Generally,  $BaCO_3$  was found to give a purer phase when used as a Ba source, instead of  $BaMoO_4$ . When  $BaMoO_4$  is used, split peaks characteristic of phase **TRI**, appear as a single broad peak, indicating that other hollandite phases, probably of **TET** form, exist as a major species, figure (S1).

Figure 3 summarizes the effects of changing input Ba:Mo stoichiometry and Mo:O stoichiometry. The right panel shows mixtures of BaCO<sub>3</sub>:MoO<sub>2</sub>, and Mo molar ratio was varied as x:7.15:0.85, where x ranges from 0.9 to 1.3. The referenced oxygen stoichiometry assumes that BaCO<sub>3</sub> is effectively BaO and that the resulting CO<sub>2</sub> is inert. However, tubes are not pressurized after the reaction is complete, which implies that either the carbon dioxide is reduced to carbon or that the CO<sub>2</sub> intercalated into the hollandite channels. Most of the evidence

supports the former scenario, with additional details summarized as supplementary information.

Multiple iterations of this study give the optimal target stoichiometry of  $\mathsf{Ba}_{1.12}\mathsf{Mo}_8\mathsf{O}_{15.44}$  to synthesize compound **TRI**. The fact that the purity of **TRI** is extremely sensitive to conditions and that the lattice parameters do not vary together suggest that **TRI** is a line compound of approximate stoichiometry  $\mathsf{Ba}_{1.12}\mathsf{Mo}_8\mathsf{O}_{16}.$ 

# **Effects of mineralizers**

A systematic study into the effects of different salts added to the reaction mixture is summarized in Table 1. Compounds such as KCl have previously been successfully added as mineralizers in, for example,  $\mathrm{Ba}_{1.2}\mathrm{Mn}_8\mathrm{O}_{16}$ ,  $^{17}$  to suppress intermediate phase formation. In the present system, the study is complicated by the fact that Na and Li have previously been reported to substitute into the channels, yielding the **MON** form. Thus, one primary goal of this work is to differentiate between reactive vs. mineralizing salts.

The preparation of  $A_x \text{Mo}_8 \text{O}_6$  hollandite phases, where A = K, Ca or Sr, and x ranges from 0 to 2 was attempted using the compounds  $\text{K}_2\text{CO}_3$ ,  $\text{CaCO}_3$ ,  $\text{CaMoO}_4$ ,  $\text{SrCO}_3$  and  $\text{SrMoO}_4$  as starting materials and with no  $\text{BaCO}_3$  or  $\text{BaMoO}_4$  present. No hollandite phases were identified across the range of conditions attempted: cooling rates between 21 °C/hr and 167 °C/hr and annealing temperatures from 540 °C to 1150 °C.

Attempts to create mixed cation phases of the form  $A_x Ba_y Mo_8 O_{16}$  with y ranges from 1.0 to 1.2 also showed no signs of changes in the crystal system or lattice parameters compared to the Ba-only experiments. The hollandite phases produced with or without adding these salts to the recipe were indistinguishable. These salts were found to be detrimental to purity and crystallinity of the Ba-Mo-O hollandite phase(s) formed, and thus their utilization as mineralizers was also further excluded.

Halide salts were also investigated as potential mineralizers, including NaCl, KCl, NaI, and KI, with the iodides serving the

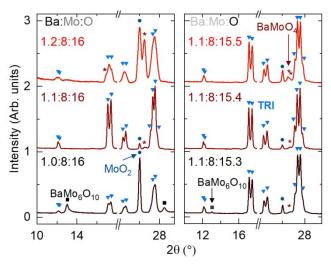


Figure 3: PXRD patterns demonstrating the respective effects of input Ba content (left) and O content (right) on purity of phase TRI when using BaCO3 as a barium source and no additional salts as mineralizers.

double purpose of being mild reducing agents as well. While

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most choices suppressed phase purity and crystallinity, the addition of NaI to the reaction mixture was found to favour the formation of **TET** phases in the form of relatively large, well-faceted single crystals. In this case, using BaMoO<sub>4</sub> yields > 98% pure **TET**, based on Rietveld refinement, whereas using BaCO<sub>3</sub> yields hollandite as a minority phase. The reaction in the presence of NaI was additionally found to be much less sensitive to oxidative effects and that a slight excess of BaMoO<sub>4</sub> further suppresses impurities. The remaining BaMoO<sub>4</sub> impurities were then leached out by washing the product with dilute HCI. The change in cooling rate affected the crystals' morphology and product yield, as rapid quenching yielded larger, often intertwined crystals, accompanied by a higher fraction of impurities, figure (S4).

The stronger dependence of **TET** phase formation on cooling rate compared to **TRI** suggests that **TET** is thermally unstable at lower temperatures.

#### MON phase formation

The previous section illustrates the need to better understand the role of Na in the formation of hollandite phases. Trials were carried out using combinations of either  $BaMoO_4$  or  $BaCO_3$  as Ba sources and either  $Na_2CO_3$  or  $Na_2MoO_4$  as a Na source at 1000 °C and an 83 °C/hr cooling rate. For this compound, MON,  $BaMoO_4$  and  $Na_2MoO_4$  are added instead of  $BaCO_3$  and  $Na_2CO_3$ ,

since adding any of the latter reagents to the recipe gave a higher fraction of  $MoO_2$  and  $BaMoO_4$  impurities.

Figure 4 summarizes the results of this series of experiments. The delicate balance between stoichiometry and the redox environment was again not subjected to simple lever-rule laws. Either  $MoO_2$  or a combination of  $BaMoO_4$  and  $BaMo_6O_{10}$  are produced when there either is deficiency or surplus in Ba content, respectively.  $Na_2MoO_4$  was added in a molar ratio of 1:4 to  $BaMo_8O_{16}$ , to make the stoichiometry of input reactants in the form  $Na_xBaMo_8O_{16}$ , where x ranges between 0.35 to 1 (Figure 4A). Impurities are minimized with the input stoichiometry  $Na_{0.5}BaMo_8O_{16}$ , in accordance with prior work by  $Lii^{16}$ . Input Ba content was then varied from nominally  $Na_{0.5}Ba_{0.97}Mo_8O_{16}$  to  $Na_{0.5}Ba_{1.1}Mo_8O_{16}$ , where adding 1.02 Ba atoms per  $Mo_8O_6$  was found to give the highest purity product, Figure 4B.

Either modifying cooling rates or the input Mo:O molar ratio (from to 8:15.3 to 16: 8) made minor difference to the phase purity or crystallinity of products, and so those results are not shown. No significant effect on product yield was found upon changing the cooling rate, and thus the cooling rate of 33 °C/hr used in initial trials was not further adjusted.

In summary, the **MON** phase wase prepared in optimal form by adding  $Na_2MoO_4$ ,  $BaMoO_4$ ,  $MoO_2$ , and Mo, in a molar ratio of 0.25:1.02:5.46:1.27, respectively. This is equivalent to a Na:Ba:Mo:O molar ratio of 0.5:1.02:8:16.

Table 1: The effect of mineralizer species on product hollandite phase identity.

Salt/Mineralizer <sup>a</sup>	Ba source	Ba:Mo:O (mol)	Cool rate	Hollandite product	Comments; Impurities
None	BaCO <sub>3</sub>	1.1:8:16	50 °C/hr	TRI	Highest <b>TRI</b> yield <sup>b</sup>
None	BaCO <sub>3</sub>	1.2:8:16	Quenched	TET (trace TRI)	Poor cryst; MoO <sub>2,</sub> BaMoO <sub>4</sub>
None	BaMoO <sub>4</sub>	1.1:8:16	Quenched	TET	Poor cryst.; MoO <sub>2,</sub> BaMoO <sub>4</sub>
None	BaMoO <sub>4</sub>	1.0:8:16	42 °C/hr	TET + TRI	Very Poor cryst;
MgO	BaCO <sub>3</sub>	1.0:8:16	10 °C/hr	TRI	Mg <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub> , MoO <sub>2</sub> , BaMoO <sub>4</sub>
KI	BaMoO <sub>4</sub>	1.1:8:16	33 °C/hr	TRI	MoO <sub>2</sub>
K <sub>2</sub> CO <sub>3</sub>	BaCO <sub>3</sub>	1.0:8:16	10 °C/hr	TRI	Poor cryst.; MoO <sub>2</sub>
K <sub>2</sub> CO <sub>3</sub>	BaMoO <sub>4</sub>	1.0:8:16	33 °C/hr	TRI	Poor cryst.; BaMo <sub>6</sub> O <sub>10</sub> , MoO <sub>2</sub>
CaCO <sub>3</sub>	BaCO <sub>3</sub>	1.0:8:16	10 °C/hr	TRI	Poor cryst.; MoO <sub>2</sub> , BaMoO <sub>4</sub>
Na <sub>2</sub> MoO <sub>4</sub>	BaCO <sub>3</sub>	1.0:8:16	33 °C/hr	TRI	MoO <sub>2</sub>
Na <sub>2</sub> MoO <sub>4</sub> (excess)	BaCO <sub>3</sub>	1.0:8:16	33 °C/hr	MON	MoO <sub>2</sub> , BaMoO <sub>4</sub>
$Na_2MoO_4$	BaMoO <sub>4</sub>	1.1:8:16	33 °C/hr	MON	highest <b>MON</b> yield <sup>b</sup>
Na <sub>2</sub> CO <sub>3</sub>	BaMoO <sub>4</sub>	1.0:8:16	33 °C/hr	MON	MoO <sub>2</sub> , BaMoO <sub>4</sub>
Nal	BaMoO <sub>4</sub>	1.2:8:16	83 °C/hr	TET	highest <b>TET</b> yield <sup>b</sup>
Nal	BaMoO <sub>4</sub>	1.2:8:16	Quenched	TET	Large cryst.; MoO <sub>2</sub> , BaMoO <sub>4</sub>
Nal	BaCO <sub>3</sub>	1.0:8:16	83 °C/hr	TET	Poor cryst.; MoO <sub>2</sub> , BaMoO <sub>4</sub>
$CaMoO_4$	BaMoO <sub>4</sub>	1.0:8:16	33 °C/hr	TET (tr. TRI)	MoO <sub>2</sub>
NaCl	$BaMoO_4$	1.0:8:16	33 °C/hr	TET (tr. TRI)	MoO <sub>2</sub> / BaMoO <sub>4</sub>

<sup>&</sup>lt;sup>a</sup>Unless otherwise stated, 5 to 10 mg of mineralizer is added to 0.25 mmol reactant mixture. Impurities and crystallinity are mentioned when excessively present, or being different relative to best results obtained, respectively. <sup>b</sup> Optimal yield and crystallinity conditions for this polymorph.

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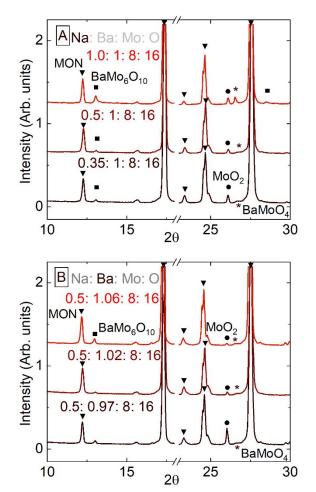


Figure 4: A and B: PXRD patterns demonstrating the respective effects of input Ba content and O content on phase purity of MON. Note that excess BaMoO4 and Na2MoO4 have been removed by washing with dil. HCl.

#### Spectroscopic stoichiometry determination

Output stoichiometry was determined by ICP-AES analysis for selected hollandite samples. The Ba:Mo ratio for phases **TRI** and **MON** was found to be 1.036(3) and 1.083(3) Ba per 8 Mo atoms, respectively. The average Ba: Mo ratio ranges between 1.15-1.18 Ba per 8 Mo atoms, in **TET** phase. These stoichiometries are slightly different from those obtained by XRD and APT techniques, because of the presence of  $MoO_2$  and  $BaMo_6O_{10}$  impurities. Additionally, Mo ions have been reported to be somewhat unstable in aqueous solutions over time, unlike Ba ions. Rietveld analysis of the PXRD patterns of the respective optimal synthetic conditions indicates that **TRI** and **MON** are  $\geq$  98% of the crystalline phases present by mass, with  $MoO_2$  being the main impurity in both products, followed by  $BaMo_6O_{10}$ . No impurity phases were observable in the **TET** PXRD patterns. An illustration of the experimental PXRD patterns with main peak

indices is shown in Figure (S6) of the supplemental documents. These observations are reproducible across multiple reactions under the same conditions.

The three-dimensional APT maps were used to determine the stoichiometry and sub-nanometre homogeneity of a specimen derived from a selected MON crystal, Figure (S7). The stoichiometry of atomic interplanar spacings in the APT measurements is consistent with the X-ray diffraction measurements for the monoclinic unit cell. The chemical formula  $Na_{0.35(2)}Ba_{1.08(3)}Mo_8O_{16},$  was determined using the ratio of Na, Ba, and Mo atoms in the 15 million ions detected. No significant spatial systematic variance was observed in chemical composition. Figure (S8) shows a 3D reconstruction of the specimen.

#### Single crystal structure solutions

Figure 5 shows SEM images of selected **TRI**, **TET**, and **MON** crystals recovered from their sealed copper tubes after heating. Their crystal habits are consistent with their assigned crystal classes, which implies that the phases observed at room temperature remain in the same structure type as when they formed at elevated temperature.

For MON and TRI crystals, there were no significant differences in the measured crystallographic parameters between crystals of the same type made using different reaction conditions. The results reported below are from the batches that had undergone optimal preparation conditions as mentioned earlier. The absence of satellite peaks indicates that MON and TRI compounds do not exhibit long-range structural modulations. MON specimens typically exhibit twinning that mimics the primitive tetragonal cell of Rb<sub>1.5</sub>Mo<sub>8</sub>O<sub>16</sub>, but close inspection of the diffraction peaks finds that a monoclinic distortion is present that is consistent with the original preliminary report of nominal Na<sub>0.35</sub>BaMo<sub>8</sub>O<sub>16</sub> (comparison shown in Table S(1)). TRI crystals also exhibit twinning, but they mimic the higher symmetry body-centred tetragonal cell of K<sub>2</sub>Cr<sub>8</sub>O<sub>16</sub>. In most batches containing TRI, this twinning precluded reliable structure solutions. We were only able to identify suitable crystals in the nearly optimized synthesis batches. Nevertheless, the unit cells and provisional structural models on the lower quality crystals are consistent with the optimal ones reported below. As above, the cell determined here is consistent with the nominal Ba<sub>1.13</sub>Mo<sub>8</sub>O<sub>16</sub> phase reported in the literature. 15 A comparison is shown in Table (S1)

Isolated **TET** crystals are unambiguously of tetragonal metric symmetry, and, unlike **MON** and **TRI** crystals, all **TET** specimens show satellite peaks that are indicative of incommensurate modulated structures. The observations hold true for synthesis

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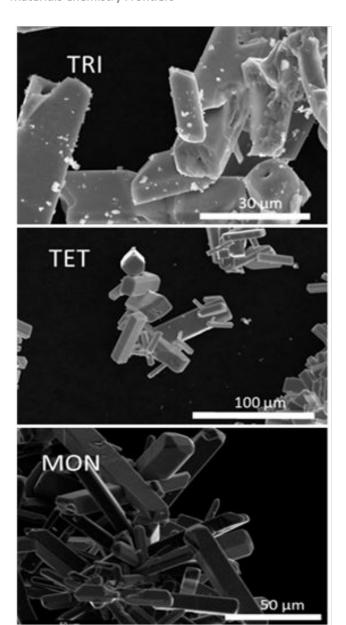


Figure 5: SEM images of the TRI, TET, and MON phases.

methods containing NaI and those without, implying that these exist as ternary Ba-Mo-O phases.

Crystals exhibiting four different **q** modulation vectors, ranging from 0.533 to 0.554 **c\***, can be isolated from the same reaction product. These are different lengths than the 0.571 (4/7) **c\*** one that was previously reported in  $Ba_{8/7}Mo_8O_{16}$ . Assuming a similar modulation model as the latter compound implies a precise composition of  $Ba_xMo_8O_{16}$ , where x=2q with q in reciprocal lattice units. The observed **TET** phases can thus be assigned nominal compositions ranging in discrete steps from  $Ba_{1.067}Mo_8O_{16}$  to  $Ba_{1.108}Mo_8O_{16}$ . Due to the complexity of solving and describing modulated structures, however, the full structure solutions of all four new **TET** phases are the subject of a separate article that is in preparation.<sup>32</sup>

The experimental conditions and model parameters for the complete structural models of **TRI** and **MON** are given in in Table 2, which have the refined stoichiometries  $Ba_{1.12(3)}Mo_8O_{16}$ 

and  $Na_{0.325(5)}Ba_{1.006(18)}Mo_8O_{16}$ , respectively. In both cases, Mo and O positions and anisotropic thermal parameters were freely refined, but the treatment of A-site cations in the channels differs. The atomic positions in the TRI and MON crystallographic models are given in Table 3 and Table 4, respectively. The process for choosing and refining crystallographic models for each crystal structures is described as follows.

In TRI structure, there are two unique channel sites, which can be seen in the right panel of Figure 6A. The cation centred at the unit cell origin requires two inequivalent Ba sites, one (Ba1A) of about 0.62 occupancy on the inversion centre and a second split site (Ba1B) about 0.4 Å away of about 0.125 occupancy each. This gives the cell origin channel site a nearly full combined occupancy of about 0.87. The other cation channel cage only has one inequivalent site (Ba2) that is essentially the same relative position and occupancy as the Ba1B site. This means that the second cage in the model has a much lower average occupancy of 0.25. The overall stoichiometry refines to Ba<sub>1.12(3)</sub>Mo<sub>8</sub>O<sub>16</sub>, which the same as the ideal input stoichiometry within error. The stoichiometry implies that an extra 1 Ba site per 8 vacancies, on average, which also seems to correlate with the approximate occupancies of the three A-sites. An alternate model that uses only one site in the Ba1 cage, which is only stable on the inversion centre, leads to significantly worse overall agreement factor of  $wR_2 = 0.0710$ , rather than 0.0576 for two Ba sites. In the preferred model, the Ba1B site was constrained to have the same thermal parameters as the Ba1A site, as the freely refined model did not improve agreement with the data despite having more refinement parameters.

The model for **MON** is shown in Figure 6B. In this case, Ba and Na atoms are both present in the channels and the model requires two inequivalent channel cation sites per cage, about 0.30 Å apart, giving a total of four inequivalent channel cation sites in the model. Unlike the **TRI** structure, there is no special site along the channel to generate extra disorder nor are there any atomic positions on the cage centres. Like the **TRI** structure, attempts to combine neighbouring partially occupied sites gave many problems. The numerical label refers to the inequivalent cages 1 and 2, with A and B denoting comparable sites for each cage. In this case, the Ba1A and Ba2A sites are the ones farther from their respective cage centres, and they are of slightly lower occupancy. This is shown in the right panel of Figure 6B.

All four Ba sites are treated as a disordered mixture of Na and Ba at a ratio fixed to the APT experimental results (0.35:1.08 or 0.24:0.76), with the total occupancy of each site being freely refined as a single parameter per site. The 1A and 2A site pair and 1B and 2B pair were respectively constrained to have equivalent thermal parameters. More details about these choices are available in the electronic supplementary information.

A final complication in the **MON** structural model is the observation of excess electron density in the unoccupied part of the hollandite channel (between the 1B and 2B sites in Figure 6B), which is most likely a manifestation of additional local

Table 2: Crystallographic data and parameters of synthesized monoclinic and triclinic phases.

	MON	TRI
Refined formula	Na <sub>0.305(8)</sub> Ba <sub>0.94(2)</sub> Mo <sub>8</sub> O <sub>16</sub>	Ba <sub>1.12(3)</sub> Mo <sub>8</sub> O <sub>16</sub>
Space group type	P2/n	P-1
Z	2	1
а	10.22775(19) Å	5.73179(12) Å
b	5.73224(12) Å	7.31811(17) Å
с	10.2849(2) Å	7.45804(14) Å
α	90°	89.2492(17)°
в	90.2821(18)°	78.4863(17)°
Υ	90°	80.4477(18)°
V	602.98(2) Å <sup>3</sup>	302.224(11) Å <sup>3</sup>
Temp.	299.4(6) K	301.15(10) K
Size	63 × 59 × 32 μm	79 × 24 × 20 μm
λ	0.71073 Å	0.71073 Å
μ	11.169 mm <sup>-1</sup>	11.689 mm <sup>-1</sup>
Abs corr	Multi-scan	Multi-scan
$T_{min}$ , $T_{max}$	0.589, 0.750	0.508, 0.840
Scan mode	ω-scans	ω-scans
$artheta_{\sf max}$ , $d_{\sf min}$	42.980°, 0.521 Å	44.876°, 0.504 Å
(h k I) <sub>max</sub>	(19 10 19)	(11 14 14)
Unique ref.	4294	4985
Completness	97.4 %	99.9 %
$R_{ m int}$ , $R_{ m \sigma}$	0.0247, 0.0260	0.0264, 0.0285
F <sub>000</sub>	1041	527
R <sub>1</sub> (all)	0.0286	0.0416
$R_1(F_o > 4\sigma F_o)$	0.0245	0.0339
$wR_2$ (all)	0.0634	0.0576
Max. Diff. peaks	8.02, -1.98 (e <sup>-</sup> /ų)	1.52, -2.03 (e <sup>-</sup> /ų
CSD code	2371927	2371928

ordering correlations that are not accounted in the model. Attempts to include the residual density as an additional site show that it accounts for about 5% of the total A-site electron density. These attempts required placing the A-site exactly between two cages with an extremely anisotropic thermal parameter. This happens to be the location of the Na-site in the prior report of the same compound. However, the resulting A-O bond lengths are around 2.1 Å, which is much too short even for Na and the observations contraindicate fractional O-site distortions nearby.

It is our view that there are significant local structure correlations that preclude choosing a single crystallographic model that best explains all features. Instead, specific parameters should be drawn from the model that they are tailored for. We conclude that the approximate A-site coordination is best described by the four-site model described above and that the fifth A-site between cages does not correspond to an actual crystallographic site. On the other hand, the overall occupation of the hollandite channels is better described by modelling the electron density, which does require using a fifth site. This model yields the final stoichiometry  $Na_{0.325(5)}Ba_{1.006(18)}Mo_8O_{16}$ , which is a closer match to the mass spectrometry results and within error of the fractional stoichiometry  $Na_{1/3}BaMo_8O_{16}$ . Though the latter point may be a mere coincidence, we have provided a simple schematic showing the two simplest local vacancy filling models at the determined ratio in Figure 6C.

Table 3: Atomic parameters for the TRI structure.

Atom	x	у	Z	$U_{11} \cdot 10^2$	$U_{22} \cdot 10^2$	$U_{33} \cdot 10^2$	$U_{23} \cdot 10^2$	$U_{13} \cdot 10^2$	$U_{12} \cdot 10^2$	Осс.
Ba1A	1	0	0	1.69(17)	1.08(3)	0.068(2)	0.168(18)	0.21(5)	0.51(7)	0.74(2)
Ba1B	1.071(3)	0.0187(9)	0.0139(8)	0.2(2)	0.85(10)	0.090(10)	-0.17(8)	-0.20(10)	0.02(10)	0.066(10)
Ba2	0.5923(2)	-0.02740(17)	-0.01957(16)	1.29(5)	0.90(5)	0.067(4)	0.16(3)	-0.37(3)	-0.45(4)	0.1240(10)
Mo1	0.35067(3)	0.15006(3)	0.48775(3)	0.280(6)	0.488(7)	0.0380(6)	0.021(5)	-0.052(5)	0.021(5)	1
Mo2	0.79853(3)	0.17005(3)	0.47590(3)	0.293(6)	0.512(7)	0.0444(6)	0.014(5)	-0.065(5)	-0.006(5)	1
Mo3	0.44095(3)	0.49163(3)	0.84392(3)	0.341(7)	0.509(7)	0.0393(6)	0.014(5)	-0.108(5)	0.017(5)	1
Mo4	0.89350(3)	0.48111(3)	0.81555(3)	0.342(7)	0.551(7)	0.0442(6)	0.058(5)	-0.109(5)	-0.020(5)	1
01	0.6461(3)	0.0304(3)	0.2992(2)	0.55(6)	0.91(6)	0.039(5)	0.00(5)	-0.06(5)	0.06(5)	1
02	0.8610(3)	-0.0352(2)	0.6563(2)	0.46(6)	0.68(6)	0.062(6)	0.04(5)	-0.05(5)	-0.13(5)	1
03	0.6680(3)	0.6436(3)	0.7007(2)	0.41(6)	0.76(6)	0.066(6)	0.30(5)	-0.08(5)	-0.01(5)	1
04	0.8351(3)	0.3668(3)	0.2848(2)	0.50(6)	0.76(6)	0.067(6)	0.29(5)	-0.23(5)	-0.08(5)	1
<b>O</b> 5	0.6745(3)	0.3045(2)	0.9576(2)	0.49(6)	0.65(6)	0.061(6)	0.12(5)	-0.06(5)	-0.01(5)	1
06	0.8376(3)	0.6453(3)	1.0448(2)	0.39(6)	0.81(6)	0.067(6)	0.13(5)	-0.13(5)	-0.08(5)	1
07	1.0123(3)	0.2978(3)	0.6272(2)	0.34(6)	0.89(7)	0.072(6)	-0.19(5)	-0.08(5)	0.10(5)	1
08	0.5144(3)	0.3083(3)	0.6303(2)	0.43(6)	0.77(6)	0.057(6)	-0.23(5)	-0.15(5)	0.16(5)	1

Table 4: Atomic parameters for the  $\bf MON$  structure.  $^a$ 

Atom	Х	у	Z	$U_{11} \cdot 10^{2}$	$U_{22} \cdot 10^{2}$	$U_{33} \cdot 10^{2}$	$U_{13} \cdot 10^{2}$	Ва осс.	Na occ.
Ba/Na1A	1/4	-0.0430(9)	3/4	0.63(3)	0.61(10)	0.59(3)	0.05(2)	0.146(12)	0.047(4)
Ba/Na1B	1/4	-0.1008(11)	3/4	0.60(2)	3.25(16)	0.64(2)	0.028(15)	0.381(13)	0.124(4)
Ba/Na2A	1/4	0.2721(8)	3/4	0.63(3)	0.61(10)	0.59(3)	0.05(2)	0.185(11)	0.060(4)
Ba/Na2B	1/4	0.3239(18)	3/4	0.60(2)	3.25(16)	0.64(2)	0.028(15)	0.231(11)	0.075(3)
Atom	X	у	Z	$U_{11}\cdot 10^2$	$U_{22}\cdot 10^2$	$U_{33}\cdot 10^2$	$U_{23}\cdot 10^2$	$U_{13}\cdot 10^2$	$U_{12} \cdot 10^{2}$
Mo1	0.43134(2)	0.61618(3)	0.42042(2)	0.267(6)	0.218(6)	0.345(6)	0.048(4)	-0.006(5)	0.031(4)
Mo2	0.42496(2)	0.16683(3)	0.40104(2)	0.314(6)	0.293(7)	0.388(6)	0.002(4)	0.004(5)	0.001(4)
Mo3	0.41723(2)	0.40093(3)	0.07320(2)	0.371(6)	0.277(7)	0.304(6)	0.044(4)	-0.043(5)	-0.047(4)
Mo4	0.39595(2)	0.84549(3)	0.08001(2)	0.392(6)	0.314(6)	0.342(6)	0.021(4)	-0.021(5)	-0.038(5)
01	0.41767(16)	0.6064(3)	0.61926(17)	0.52(5)	0.46(6)	0.54(5)	-0.05(4)	0.06(4)	0.10(4)
02	0.44024(16)	0.1098(3)	0.59622(16)	0.53(5)	0.33(6)	0.49(5)	0.05(4)	0.02(4)	-0.01(4)
03	0.41906(17)	0.6272(3)	0.22081(16)	0.88(6)	0.40(6)	0.33(5)	0.09(4)	-0.01(4)	0.04(4)
04	0.41867(17)	0.1244(3)	0.20812(16)	0.81(6)	0.47(6)	0.35(5)	0.01(4)	-0.01(4)	-0.07(4)
05	0.61745(17)	0.3767(3)	0.07986(16)	0.66(6)	0.54(6)	0.42(5)	0.03(4)	0.00(4)	0.16(4)
06	0.59344(16)	0.8777(3)	0.05329(16)	0.52(5)	0.35(6)	0.58(6)	0.02(4)	0.07(4)	0.02(4)
07	0.28124(16)	0.3839(3)	0.41589(17)	0.24(5)	0.48(6)	0.80(6)	-0.07(4)	0.01(4)	-0.02(4)
08	0.21132(16)	0.8831(3)	0.08279(17)	0.28(5)	0.47(6)	0.80(6)	-0.01(4)	0.01(4)	0.02(4)

 $^{a}U_{ij}$  constrained to zero by symmetry are not included

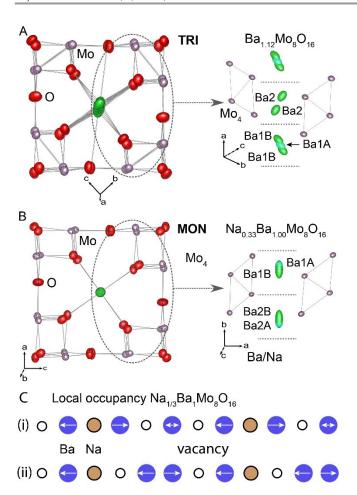


Figure 6: Crystal structure of (A) **TRI** and (B) **MON** phases. The left side shows the view projected down the Ba channels. The right portion shows a view perpendicular to the channels, which emphasizes the different stackings of the Mo<sub>4</sub> clusters relative to the Ba sites. (C) Shows two local structures of Na and Ba occupancies that were identified as consistent with the approximate rational stoichiometry of **MON**.

Table 5: The Mo-Mo bond lengths of K<sub>2</sub>Mo<sub>8</sub>O<sub>16</sub>, **TRI** and **MON** phases.

Bond (Å)	TRI	MON	$K_2Mo_8O_{16}$
$A_1$ (Mo <sub>1</sub> -Mo <sub>1</sub> )	2.5797(5) Å	2.5308(4)	2.551(3)
$B_1$ (Mo <sub>1</sub> -Mo <sub>2</sub> )	2.6170(4) Å	2.6562)3)	2.6966(17)
$C_1$ (Mo <sub>1</sub> -Mo <sub>2</sub> )	2.5794(4)	2.5842(3)	2.5965(17)
$D_1$ (Mo <sub>1</sub> -Mo <sub>2</sub> )	3.1626(4)	3.1634(3)	3.1765(18)
$E_1$ (Mo <sub>2</sub> -Mo <sub>2</sub> )	3.1716(5)	3.1827(4)	3.199(3)
$A_2$ (Mo <sub>3</sub> -Mo <sub>3</sub> )	2.5605(5)	2.5392(5)	2.527(3)
$B_2$ (Mo <sub>3</sub> -Mo <sub>4</sub> )	2.8473(3)	2.8558(3)	2.8371(17)
$C_2$ (Mo <sub>3</sub> -Mo <sub>4</sub> )	2.5483(4)	2.5585(3)	2.5653(17)
$D_2$ (Mo <sub>3</sub> -Mo <sub>4</sub> )	3.2005(4)	3.1922(3)	3.2085(18)
$E_2$ (Mo <sub>4</sub> -Mo <sub>4</sub> )	3.2570(5)	3.2264(4)	3.195(3)

# **Magnetic Properties**

The magnetization of **TRI** and **MON** phases is overall quite low, as shown in Figure 7, with each having similar temperature-independent components,  $\chi_o$  = 4.50 x 10<sup>-5</sup> and 7.20 x 10<sup>-5</sup> emu/mol<sub>Mo</sub>. The fitted Curie constants are 1.238(3) x 10<sup>-2</sup> and 2.235(7) x 10<sup>-3</sup> emu · K/mol<sub>Mo</sub>, for **TRI** and **MON** species, respectively, which correspond to effective magnetic moments of 0.315 and 0.137  $\mu_B$  in the Curie-Weiss model. Along with a larger spin density, the **TRI** sample also exhibits a larger magnitude Curie-Weiss constant of -12.75(6) K. It also shows a small deviation from Curie Weiss behaviour below 10 K, which is consistent with weak antiferromagnetic correlations. The fitted  $\theta_W$  in **MON**, 1.1(2) K, is effectively zero and is consistent with a very dilute magnetic impurity.

### Discussion

A goal of this work is to understand the nature of the observed Mo-Mo bonding motifs, with a particular focus on the relationships between electron count, A-site steric interactions, and thermodynamic stability. Since the A-site filling is equivalent to adding/removing electrons from the chains of  $[\text{Mo}_4\text{O}_8]^{n_-}$  clusters, the synthetic portion of the study provides the foundational details.

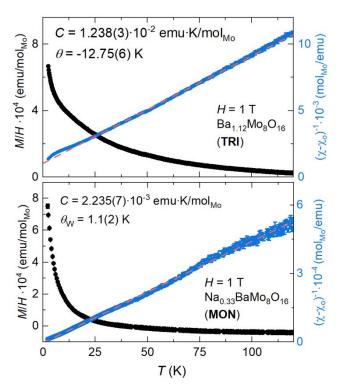


Figure 7: Temperature dependence of *M/H* of **TRI** (top) and **MON** (bottom) phases measured using an applied field of 1 T.

We found that both **TRI** and **MON** phase purity are extremely sensitive to input stoichiometry, requiring greater precision in composition than the sensitivity of these experiments, meaning they can be treated as compounds with fixed stoichiometries at about  $Ba_{1.12(3)}Mo_8O_{16}$  and  $Na_{0.325(5)}Ba_{1.006(18)}Mo_8O_{16}$ .

The **TRI** phase Ba-site occupancy is also crystallographically disordered, which implies that overall having 56% Ba and 44% vacancies strikes the optimal balance between  $[Mo_4O_8]^{n-1}$  reduction and Ba-Ba inter-site repulsion. The present results also favour treating the (as of now) five known **TET** phases as fixed stoichiometries that are derived from modulated *A*-site occupancies ranging from 1.067 to 1.143 Ba per 8 Mo. This implies that conditions favouring long-range ordering provide additional stabilization concerns. We speculate that the **TRI** phase could be the kinetic product that is metastable with respect to the ordered **TET** phases, but more work is required.

Since there is only one stable composition of Na-Ba-Mo-O hollandite, it is apparently best to treat **MON** as a unique quaternary phase and not "Na-substituted  $Ba_xMo_8O_{16}$ ". The determined stoichiometry's proximity to small rational numbers might be coincidental, as we found no evidence for additional crystallographic order in the dozens of single-crystal diffraction experiments we performed using a high flux rotating anode source and low-background detector.

While this rules out long-range crystallographic ordering as stabilizing the specific stoichiometry, short-range ordering is still expected. A simple interpretation is that a lattice of alternating Ba<sup>2+</sup> ions and vacancies can accommodate one Na<sup>+</sup> ion per six vacancies without disrupting the Ba occupancies. Figure 6C shows two local structures that are consistent with

this scenario. Both local structures appear susceptible to disordered atomic displacements, especially when mixed. We note that this filling of A-sites donates 2.33 e. per formula unit (f.u.), which is higher than any of the ternary Ba-Mo-O phases. It thus appears that thermodynamic stability of **MON** is driven by the reduction of the  $[Mo_4O_8]^{n}$ - sublattice, which is enabled by the lower electrostatic repulsion of the Na<sup>+</sup> ion compared to the Ba<sup>2+</sup> ion.

# **Analysis of crystal symmetries**

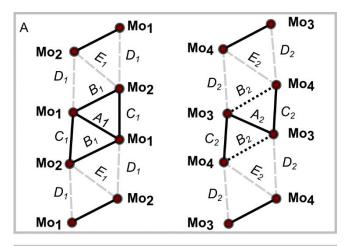
The above analysis helps explain the compositions to a certain degree, but they do not explain the origin of the different structure types. To do this, one must consider trends in the crystal structures of these and related compounds. In the general case, the symmetry of hollandite compounds is known to correlate to the ionic size of A-cations. If the cation is small enough, the octahedral walls undergo twisting distortions to accommodate minimum energy cation positions closer to the coordinating O atoms and monoclinic symmetry will be favoured. Otherwise, a tetragonal symmetry is maintained when larger A-cations occupy tunnels within the hollandite structure.

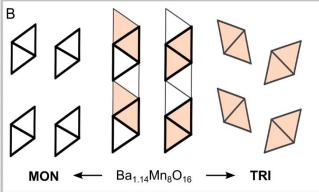
While the above works well for many hollandites, the opposite trend is seen in  $Rb_{1.5}Mo_8O_{16}$  and  $K_2Mo_8O_{16}$ , which are tetragonal and monoclinic at room temperature, though  $Rb^+$  is the larger ion. It thus seems likely that there is an electronic component in the case of Mo-based hollandites, with the electron configuration of the Mo clusters playing a major role.<sup>34</sup>

It is necessary to look more closely at the Mo<sub>4</sub> clusters to uncover the primary driver of the different structural orderings. Despite the different symmetries and stoichiometries, the main Mo<sub>4</sub> cluster motif in MON and TRI is very similar. In both determined crystal structures, the Mo atoms form strong Mo-Mo bonds that are reminiscent of molybdenum oxide clusters embedded in the hollandite chains. The intra-cluster Mo-Mo bonds range from 2.50 A to 2.90 Å, while the distance between two adjacent Mo atoms in different clusters is 3.18-3.20 Å. There are also two distinct types of chains in each compound, shown in Figure 8A as cluster 1 on left and cluster 2 on right. One chain is composed of regular rhombohedra Mo<sub>4</sub> clusters, where the five Mo-Mo bonds are nearly equal. The other chain has clusters that have three short Mo-Mo bonds (bonds A and C in Figure 8A) and two long bonds on the outer edges of the rhombus (bond B in Figure 8A). The intercluster distances are also shown in Figure 8A as pale dashes and labelled as D and E.

Error! Reference source not found. summarizes the bond lengths of both cluster types for the compounds  $Ba_{1.12}Mo_8O_{16}$  (TRI),  $Na_{1/3}BaMo_8O_{16}$  (MON), and  $K_2Mo_8O_{16}$ . There are evidently only subtle differences between compounds when comparing equivalent bonds, suggesting that such differences are not relevant to the relative stability of triclinic and monoclinic structural forms. The trends in bond lengths are discussed in detail in the electronic structure section below.

Figure 8B compares the pattern of Mo<sub>4</sub> clusters in **MON** and **TRI** phases. The triclinic distortion in **TRI** is a result of a unique set of displacement distortions affecting both Ba<sup>2+</sup> ions and Mo-





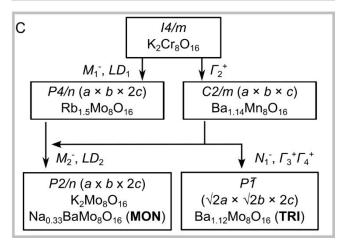


Figure 8: A) Schematic of the two distinct types of Mo4 cluster arrays in **TRI**, **MON**, and  $K_2Mo_8O_{16}$  phases showing bond labels. B) Mo clustering patterns of **TRI** and **MON**. C) Distortion pathways from higher symmetry hollandite types to **TRI** and **MON** phases. Arrows are labelled with the irreducible representations contributing to the distortion.

O octahedra that results in a staggered tiling pattern. This differs from the other Mo-based hollandites,  $Na_{1/3}BaMo_8O_{16}$  (MON) and  $K_2Mo_8O_{16}$ , where both chains of clusters are lined up along the channel axis. The different  $Mo_4$  cluster packing in TRI apparently induces shifts in the oxygen sublattice that can be associated with the distinctive Ba-site displacements away from the central channel axis. This might also allow the unusually high Ba occupancy at the cage centre for the remaining site that we observe in the triclinic form.

Figure 8C illustrates the triclinic distortion pathway determined using the Isodistort software suite. Starting from hollandites of higher symmetry, the highest symmetrical hollandite structure is represented by the space group type I4/m and the compound  $K_2Cr_8O_{16}$ . The distortion complexity is the same in **TRI** and **MON** compounds (index 8), and they are both subgroups of the monoclinic Ba<sub>1.14</sub>Mn<sub>8</sub>O<sub>16</sub> structural type, of space group type C2/m. The latter is the type of monoclinic distortion found in ordinary hollandites that was mentioned above, and it is active on the  $\Gamma_2^+$  irrep.

Considering the active mode populations, we find that the orthorhombic strain magnitude, (a-b)/(a+b), in both **MON** and  $K_2Mo_8O_{16}$  is less than a quarter of the magnitude observed in  $Ba_{1.14}Mn_8O_{16}$ , suggesting that the connection may only be circumstantial. The P2/n distortion is also a subgroup of the tetragonal P4/n phase unique to  $Rb_{1.5}Mo_8O_{16}$ , which is a more descriptive pathway. First, the  $M_1^-$  and  $LD_1$  irreps are together responsible for the cluster formation in  $Rb_{1.5}Mo_8O_{16}$ , yielding crystallographically equivalent chains. The chain degeneracy is then broken by introducing the  $M_2^-$  and  $LD_2$  irreps, resulting in the P2/n structures. In other words, the  $\Gamma_2^+$  irrep is not needed to understand the primary cluster distortion, though it does contribute as a secondary irrep. Symmetry allows the distortion from P4/n to P2/n to be continuous.

The **TRI** phase is not a subgroup of Rb<sub>1.5</sub>Mo<sub>8</sub>O<sub>16</sub> due to the different tiling of clusters. The  $N_1^-$  irrep simultaneously yields the triclinic-type cluster tiling and breaks the chain degeneracy. In other words, only inversion symmetry is compatible with the triclinic-type cluster tiling. We also note that the  $\Gamma_2^+$  type orthorhombic distortion magnitude is nearly equivalent in C2/m type Ba<sub>1.14</sub>Mn<sub>8</sub>O<sub>16</sub> and Ba<sub>1.12</sub>Mo<sub>8</sub>O<sub>16</sub> (**TRI**). We suggest the possibility that the different cluster tiling between **MON** and **TRI** is linked to the same steric considerations that drive the I4/m to C2/m distortion in ordinary hollandite phases.

#### Trends in Mo-Mo bonding

Having laid out the symmetry considerations above, this section considers viable electronic models for the origin of the structural distortions in the various Mo-based hollandites. The prior computational study by Toriyama *et al.* of  $K_2Mo_8O_{16}$  is a useful framework for the discussion.<sup>13</sup>

In that work, density functional theory (DFT) calculations found a number of flat bands around the Fermi level that are mostly localized to the Mo₄ clusters' d orbitals. They interpreted the bands as corresponding to quasi-molecular Mo-Mo orbitals: four bonding orbitals per cluster, 4 non-bonding orbitals per cluster, and 4 anti-bonding orbitals per cluster. The clusters in  $K_2Mo_8O_{16}$  are formally  $[Mo_4O_8]^-$  if treated as equivalent, meaning there are 9 d electrons per cluster and the Fermi level lies in the middle of the lowest energy non-bonding band ( $d_{vz}$  in their nomenclature). In other words, a chain would have a halffilled highest occupied molecular orbital (HOMO), and the configuration might be written d<sub>yz</sub>1. If both chains are equivalent, there would be a degenerate pair of said orbitals. Two models were then provided under this approach for interpreting the electronic environment of the Mo4 clusters and their effect on the structure and properties of K<sub>2</sub>Mo<sub>8</sub>O<sub>16</sub>.

The first model interpreted  $K_2Mo_8O_{16}$  as a band insulator, with symmetry-breaking modifications relieving chain degeneracy and opening a band gap of about 80 meV (when using GGA).<sup>13</sup> The broken degeneracy yields a fully occupied HOMO that is primarily on chain 1 (in our nomenclature) and lowest unoccupied molecular orbital (LUMO) that is primarily on chain 2. In other words, the chains 1 and 2 might be thought of respectively as  $d_{yz}^2$  and  $d_{yz}^0$ . This is the same as saying there is a net charge transfer that yields one chain with closer to 10 electrons per cluster and the other chain with 8 electrons per cluster.

The second model provided by Toriyama et~al applied Mott physics to the system and found that a Hubbard  $U_{eff} \geq 0.3~{\rm eV}$  is sufficient to break the  $d_{yz}^1$  orbital degeneracy , driving the formation of an antiferromagnetic insulating (AFI) phase. <sup>13</sup> In this model, a single spin-polarized state is filled on each cluster, giving  $S = \frac{1}{2}$  per cluster in  $K_2Mo_8O_{16}$ . Notably, the antiferromagnetic exchange interactions in this model were found to lead to geometric frustrations of the spins, which was used to explain the lack of observed long-range magnetic ordering <sup>13</sup>.

These two models have very different predictions about the relationship between electron count per  $Mo_4$  cluster and electronic properties. We show below that the band insulator model provides a good description of the observed compounds with only a minor change in interpretation, while we are unable to understand our observations within the Mott insulator model.

### Mott-type model

We consider the utility of the Mott-type model first. The model assumes each cluster has a net  $S=\frac{1}{2}$  for the special case of 2 A-site electrons per formula unit (as in  $K_2Mo_8O_{16}=[K^+]_2[Mo_4O_8^-]_2$ ). It is not easy to reconcile this idea with the experimental susceptibilities normally observed in Mocontaining hollandite compounds, which is generally small and more reminiscent of dilute impurity spins.  $K_2Mo_8O_{16}$ , for example, shows one spin per 12 Mo atom,  $^{34}$  meaning 2 out of every 3 clusters are missing the spins predicted by the Mott-like model.

The formula of **MON** indicates 2.33  $e^-$  donated per formula unit, which the Mott model would predict leads to a modest change in the number of spins:  $[Na^+]_{1/3}[Ba^{2+}][Mo_4O_8^-]_{1.67}[Mo_4O_8^{2-}]_{0.33}$ , with the  $[Mo_4O_8^{2-}]$  clusters being diamagnetic S=0. Yet, we found that an effective magnetic moment of 0.137  $\mu_B$  per Mo in the **MON** phase. Treating this as an impurity spin of  $J=\frac{1}{2}$  would imply 1 defect per 160 Mo atoms. In contrast with expectation, there are almost no free spins available.

Applying the same arguments to **TRI** gives a nearly equivalent electron donation count (here, 2.24  $e^{-}$  / f.u.) as **MON**. The experimental effective magnetic moment is 0.315  $\mu_{B}$  per f.u in **TRI**, implying one free spin per 30 Mo atoms, or about one per eight Mo<sub>4</sub> clusters.

In summary, the magnetic properties appear to be dominated by defects. We see no obvious way to reconcile the

observations with the Mott-like model and so it is not discussed further.

#### **Band insulator model**

While Toriyama *et al.* identified the resulting HOMO and LUMO bands as being essentially non-bonding  $^{13}$ , we argue that the broken-degeneracy HOMO shows clear evidence of bonding character that is mostly localized on the outer Mo-Mo bonds ( $B_1$  in Figure 8A). This implies a net increase in bond order along the same bond whenever the state is filled, which indeed matches observation: the  $B_1$  orbital is much shorter than  $B_2$  in all compounds with broken chain degeneracy, as shown in **Error! Reference source not found.**. There are also subtle differences between compounds in the relative  $B_1$  and  $B_2$  bond lengths, which can be understood as arising from different populations of the relevant states in each cluster type. That is, the band character of the states is non-negligible and so there is non-integer electrons per cluster.

This interpretation predicts that the bond-lengths should be correlated to the net valence of the Mo atoms in the cluster, since the states at the Fermi level are primarily Mo character. This is tested by applying the bond valence sum (BVS) analysis to the experimental crystal structures of **TRI**, **MON**,  $K_2Mo_8O_{16}$ , and room-temperature  $Rb_{1.5}Mo_8O_{16}^{37}$ , the results of which are shown in Table 6. The overall BVS of a compound decreases as the A-site electron count increases, as expected, and each of the monoclinic and triclinic phases show a significant difference in overall bond valence between clusters. The *P4/n* symmetry of  $Rb_{1.5}Mo_8O_{16}$  causes the two clusters to be symmetry equivalent, so both chains are approximately equivalent to chain 1 in the other compounds. This is consistent with electron transfer between chains being a significant part of the symmetry-breaking interactions.

The final test of the model is shown in Figure 9A, which shows the difference in total valence per  $Mo_4$  cluster plotted against the difference in the B Mo-Mo bond lengths from the same clusters. A direct relationship is observed, meaning a net increase in electrons in cluster 1 does in fact correlate to a net decrease in bond length (increase in bond order). This model is depicted schematically in Figure 9B, which also considers the different number of electrons donated by the A-site.

In summary, the band insulator model adapted from Toriyama *et al.* provides a helpful framework for understanding the subtle difference in molybdenum cluster formation between hollandites. The fact that the per cluster electron populations can be continuously varied (i.e. non-integer) delocalized, at least within the same chain. Similar effects have been observed in the vanadium-based Magnéli phases  $V_8O_{15}$  and  $V_9O_{17}^{36}$ , which each exhibit a metal-to-insulator transition

Table 6: Results of BVS analysis for Mo ions for the TRI and MON phases.

_								
	Bond valence							
	TRI MON K <sub>2</sub> Mo <sub>8</sub> O <sub>16</sub> Rb <sub>1.5</sub> Mo <sub>8</sub> O <sub>1</sub> .							
	$Mo_1$	3.44	3.47	3.49	3.48			
	$Mo_2$	3.59	3.61	3.67	3.61			
	Mo <sub>3</sub>	3.57	3.55	3.49	-			
	$Mo_4$	3.91	3.90	3.89	-			
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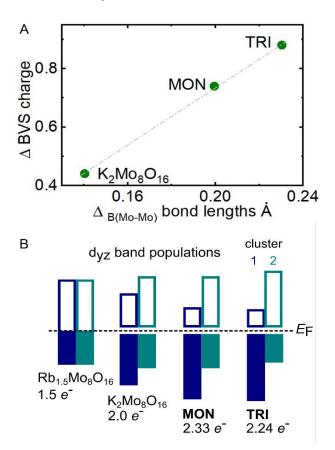


Figure 9: A) Change in BVS charge per cluster plotted against difference in bond length B for **TRI**, **MON**, and  $K_2Mo_8O_{16}$ . B) Schematic showing the interpretation of panel A using the band insulator model. The electron number refers to the total band population across both clusters, which is calculated from the stoichiometry. The cluster filling ratios are taken from Table 6.

(MIT) and also show evidence of electron delocalization between local metal-metal bonded motifs.

Further, our group isotropy analysis is also consistent with charge ordering driving the observed electrostructural transition in  $Rb_{1.5}Mo_8O_{16}$  at 208 K. Though the low-temperature structure has not yet been proven to be the same as  $K_2Mo_8O_{16}$ , the transition is apparently of second order, which is allowed by symmetry. We also speculate that the  $K_2Mo_8O_{16}$  and  $Na_{1/3}BaMo_8O_{16}$  ( $\pmb{MON}$ ) may be induced to transform into the tetragonal P4/n structure either above room temperature or under uniaxial strain.

# **Conclusions**

Distinct solid-state synthesis pathways were determined for the closely related triclinic  $\mathsf{Ba}_{1.12}\mathsf{Mo}_8\mathsf{O}_{16}$  (TRI), tetragonal  $\mathsf{Ba}_{1+x}\mathsf{Mo}_8\mathsf{O}_{16}$  (TET), and monoclinic  $\mathsf{Na}_{1/3}\mathsf{Ba}\mathsf{Mo}_8\mathsf{O}_{16}$  (MON) hollandites. TRI is optimally synthesized using  $\mathsf{BaCO}_3$  as the Ba source in the Ba:Mo:O molar ratio 1.12:8:15.44 and an optimal cooling rate around 50 °C/hr. The reducing conditions seem to be necessary due to the potential for carbothermal oxidation of Mo from the evolved  $\mathsf{CO}_2$ . No evidence was found to suggest there is compositional variance in the Ba occupancy in TRI, and the sensitivity of phase stability to reaction conditions suggests that TRI may be metastable with respect to modulated

tetragonal phase formation. The addition of Nal to the reaction mixture enhances crystal quality while simultaneously stabilizing various modulated **TET** phases not previously reported. Further work is required to fully characterize the complex structure of the modulated structures.

The addition of Na<sub>2</sub>MoO<sub>4</sub> instead of NaI results in the formation of the sodium barium molybdate **MON**, which is also apparently a line compound with the refined formula Na<sub>0.325</sub>Ba<sub>1.005</sub>Mo<sub>8</sub>O<sub>16</sub>. BaMoO<sub>4</sub> is the optimal Ba source and Na<sub>2</sub>MoO<sub>4</sub> should be added in excess, using the Na:Ba:Mo:O molar ratio 0.5:1.02:8:16, which is in accordance with prior work. Atom probe tomography confirms the presence and homogeneous distribution of Na atoms within the hollandite channels of the monoclinic compound.

Detailed analysis of Mo-O and Mo-Mo bond lengths reveal clear evidence of intercluster charge transfer. This observation supports the charge ordering model proposed by Toriyama et al for the related  $K_2Mo_8O_{16}$ , indicating that broken symmetry is electronically driven. The effects of electron doping imply that the Mo states near the Fermi level remain at least partially delocalized in the low symmetry form, which highlights the complexity of metal-metal bonding in early transition metal oxides. Additional characterization of the Mo-based hollandites is required to determine the detailed relationship between metal bonding, charge ordering, and the electronic properties.

## **Author Contributions**

Eslam Elbakry: conceptualization, data curation, formal analysis, investigation, methodology, software, validation, visualization, writing - original draft, writing - review and editing.

Jared M. Allred: conceptualization, data curation, formal analysis, funding acquisition, methodology, project administration, resources, software, supervision, validation, visualization, writing – review and editing.

## Conflicts of interest

There are no conflicts to declare.

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# **Data Availability Statement**

CIF files containing crystallographic data and observed structure factors are included as part of the supplementary information. They have also been deposited with the Inorganic Crystal Structure Database (ICSD) and Cambridge Crystal Structure Database (CCSD) under deposition numbers 2371927 and 2371928.

Other data presented in the manuscript will be made available upon request.