



**Efficient polymer light-emitting diodes (PLEDs) based on chiral [Pt(C<sup>^</sup>N)(N<sup>^</sup>O)]-complexes with near-infrared (NIR) luminescence and circularly polarized (CP) light**

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## COMMUNICATION

# Efficient polymer light-emitting diodes (PLEDs) based on chiral [Pt(C<sup>^</sup>N)(N<sup>^</sup>O)]-complexes with near-infrared (NIR) luminescence and circularly polarized (CP) light

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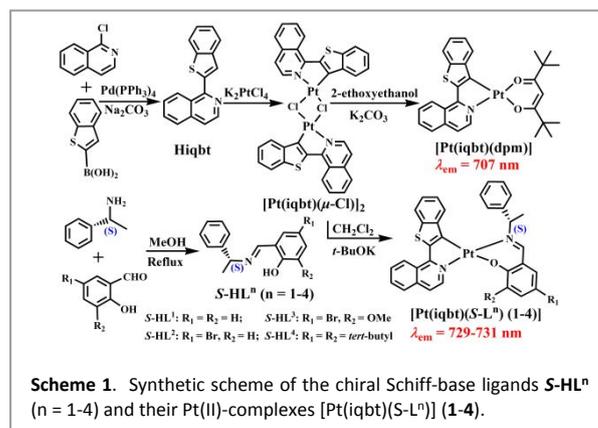
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**Based on molecular design of chiral NIR-emitters [Pt(iqbt)(S-L<sup>n</sup>)] (1-4), the first example CP-NIR-PLEDs with both attractive NIR-emissive properties ( $\lambda_{em} = 732$  nm;  $\eta_{EQE}^{Max} = 0.87-0.93\%$ ) and high dissymmetry  $g_{EL}$  factors up to  $10^{-3}$  are reported.**

Circularly polarized (CP) light-active materials,<sup>1</sup> due to the tracking of specific interactions with chiral environments, have recently attracted particular interest on chiral optoelectronics.<sup>2</sup> Especially independence on the direct generation of CP light, CP organic/polymer light-emitting diodes (OLEDs/PLEDs)<sup>3</sup> advantageous of achieving high contrast 3D images and true backlight, are significantly superior to routine OLEDs/PLEDs. To date, studies on small-molecule<sup>4</sup> or polymer<sup>5</sup> with  $^1\pi-\pi^*$ -transitioned fluorescence have been reported. When thermally-activated delayed fluorescence (TADF)<sup>6</sup> system, transition-metal complex<sup>7</sup> or organo-Ln<sup>3+</sup> complex<sup>8</sup> are employed, their capability of harvesting both  $^1S$  and  $^1T$  excitons for vivid visible ( $\lambda_{em} = 400-700$  nm) CP-OLEDs/PLEDs makes the resulting devices more appealing. In particularly, the CP devices that emit in the NIR (NIR = near-infrared;  $\lambda_{em} > 700$  nm) offer numerous advantages for bioassay, information-security and night-version readable display.<sup>9</sup> However, the study in this direction is much scarce, which can be attributed to several challenges including unavoidable CP-active materials' aggregation and lack of NIR-active chiroptical luminogen.

As a matter of fact, Pt(II)-complex phosphors have achieved great success in the development<sup>10</sup> of efficient NIR-OLEDs/PLEDs. In this context, one elegant strategy toward that restrictive NIR regime for typical vacuum-deposited NIR-OLEDs is focused on the significant bathochromatic-shift of sublimed Pt(II)-complexes caused by unforeseen excimers' formation and strong Pt...Pt-

interacting-within intermolecular interactions. However, besides the contention<sup>11</sup> of aggregation-induced quenching or enhancement of emission (AIQ or AIE), the realization of their  $\eta_{EQE}$  up to the unit<sup>12</sup> or even the tenth<sup>13</sup> actually trades off the inevitable high-cost of device fabrication and aggregation-induced detrimental efficiency-roll-off. Convincingly, circumventing such problems for the compromise of cost-effectiveness and desirable device performance (high-efficiency while relatively weak efficiency-roll-off), solution-processable NIR-OLEDs/PLEDs<sup>14</sup> based on homogeneous doping of Pt(II)-complex just with the strong MLCT effect into an appropriate host with the deep HOMO-LUMO bandgap should be more considerable as an alternative. Whereas the development of highly emissive NIR-OLEDs/PLEDs without Pt(II)-chromophores' interactions remains a great challenge due to the "energy gap law"<sup>15</sup> limit. Nuzzo *et al.*<sup>16</sup> used a chirally substituted polymer doped with achiral Pt(II)-complex as the emitting material and fabricated the device with emission around 550 nm. Fuchter *et al.*<sup>17</sup> incorporated a chiral Pt(II)-complex into a conventional light-emitting polymer and the CP-electroluminescence was observed around 600 nm. Herein, in light of the deserved NIR ( $\lambda_{em} = 707$  nm) emission<sup>18</sup> of [Pt(iqbt)(dpm)] (Scheme 1), we report the first example of CP-NIR-PLEDs by ingeniously introducing simplex chirality into the N<sup>^</sup>O-Schiff-base ancillary ligands **S-HL<sup>n</sup>** ( $n = 1-4$ ) for their chiral [Pt(iqbt)(S-L<sup>n</sup>)]-heteroleptic complexes (**1-4**; also Scheme 1). Both CP-light activity and NIR electroluminescence (732 nm) were observed. This result demonstrated the potential application of chiral NIR-emitting Pt(II)-complexes in solution-processable CP-NIR-PLEDs.



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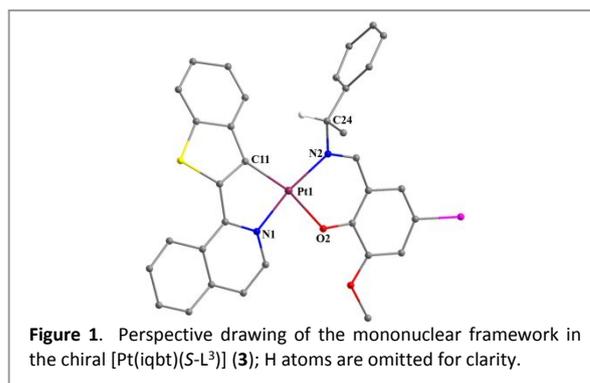
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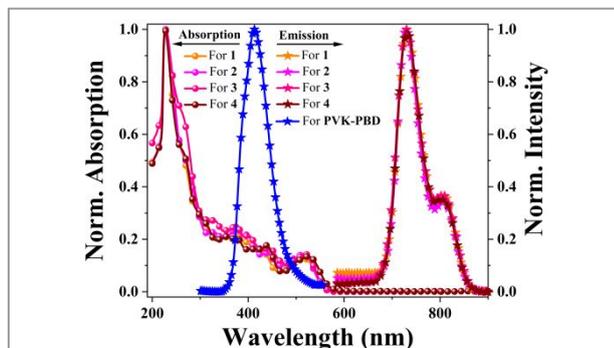
Based on reaction (also Scheme 1) of  $K_2PtCl_4$  with the  $C^{\wedge}N$ -cyclometalated ligand **Hiqbt** synthesized from the Suzuki coupling of cost-effective 2-Cl-isoquinoline<sup>19</sup> with benzo[*b*]thien-2- $\gamma$  boronic acid, the  $\mu$ -chloro-bridged dimer intermediate  $[Pt(iqbt)(\mu-Cl)]_2$  was obtained in 89%. As to the chiral Schiff-base-typed ligands **S-HL<sup>n</sup>** ( $n = 1-4$ ), they were prepared from the rational condensation<sup>20</sup> of *S*-( $-$ )-1-phenylethylamine with one of the salicylaldehyde derivatives, respectively. Further through the treatment of  $[Pt(iqbt)(\mu-Cl)]_2$  with each of the chiral ligands **S-HL<sup>n</sup>** ( $n = 1-4$ ) in the presence of *t*-BuOK, the targeted chiral Pt(II)-complexes  $[Pt(iqbt)(S-L^n)]$  (**1-4**) soluble in common organic solvents, were afforded, respectively.

The chiral Schiff-base-typed ligands **S-HL<sup>n</sup>** ( $n = 1-4$ ) and their Pt(II)-complexes **1-4** were well-characterized by EA, FT-IR, <sup>1</sup>H NMR (Figures S1-2) and ESI-MS. All the Pt(II)-complexes exhibited relatively high thermal stabilities. TGA results (Figure S3) of the chiral Pt(II)-complexes **1-4** show that their high decomposition temperatures ( $T_d$ , corresponding to 5% weight loss; 230 °C for **1**, 267 °C for **2**, 239 °C for **3** or 214 °C for **4**) were observed. As to the more favourable thermal stability with the electron-drawing Br atom for **2** while an opposite trend with the electron-donating *tert*-butyl groups for **4**, it can be reasonably explained by the relatively shorter Pt-C (1.99421 Å), Pt-N (2.05599-2.06234 Å) and Pt-O (2.11856 Å) bond lengths (Table S1) for **2** than those (Pt-C (2.00438 Å), Pt-N (2.05585-2.06706 Å) and Pt-O (2.12124 Å)) for **4** from the following TD-DFT calculations. Molecular structure of  $[Pt(iqbt)(S-L^3)]$  (**3**) was confirmed by X-ray single-crystal diffraction analysis. The complex crystallizes in the triclinic space group of *P1*, whereas the non-centrosymmetric nature originates from the (*S*- $L^3$ )-induced chirality. As shown in Figure 1, one  $C11^{\wedge}N1$ -chelate (**iqbt**) ligand and one ancillary ligand (**S-L<sup>3</sup>**) with the  $N2^{\wedge}O2$ -chelate mode, coordinate to the central  $Pt^{2+}$  ion (Pt1) in a square planar geometry, leading to the formation of a typical asymmetric mononuclear framework.<sup>18</sup> Noticeably, different from the *co*-planar conformation of the cyclometalated (**iqbt**) ligand, the terminal phenyl ring of the chiral (**S-L<sup>3</sup>**) ligand is almost perpendicular to its  $N2^{\wedge}O2$ -chelate plane with a dihedral angle of 84.1(2)°, from which, the *S*-chiroptical character centered at the chiral -C\*24-H atom for  $[Pt(iqbt)(S-L^3)]$  (**3**) is well-retained. Interestingly, complementary intermolecular interactions (3.566(2) Å of  $\pi \cdots \pi$ , 3.791(2) Å of  $\pi \cdots Pt$  and 3.710(30) Å of C18-H18 $\cdots\pi$ ; Figure S4) while no valid Pt $\cdots$ Pt interactions, are observed between every adjacent two mononuclear units, giving rise to the 1D polymeric chain. The crystallographic data of this compound can be found in Tables S2-3.

Photophysical properties of the chiral Pt(II)-complexes **1-4** in solution or solid-state were explored at RT or 77 K (Table S4 and Figures 2-3 and S5-9). As shown in Figure 2, chiral Pt(II)-complexes **1-**

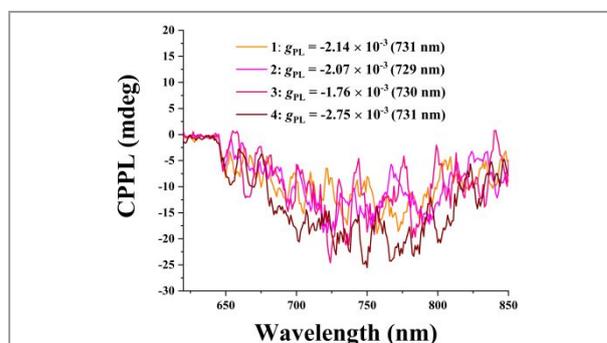


spectra than those ( $\lambda_{ab} < 400$  nm; Figure S5) of the **Hiqbt** and **S-HL<sup>n</sup>** ligands: the high energy-state intense absorption in the 200-400 nm range assigned to the ligands-based  $^1\pi-\pi^*$  transitions, and the relatively weak absorptions ( $\lambda_{ab} > 430$  nm) probably attributed to a mixed  $^1,^3LLCT/^1,^3MLCT$  transitions. The substituents on the (**S-L<sup>n</sup>**) ligands have little effect on the absorption and CD spectra (Figure S6) of their Pt(II)-complexes. Importantly, contrasted to the ligand-relative visible emissions (Figure S7) for both the **Hiqbt** and **S-HL<sup>n</sup>** ligands, all the chiral four Pt(II)-complexes **1-4** display the unique NIR-emissive emission spectra in solution at RT also as shown in Figure 2, featuring with a strong peak located at  $\lambda_{em} = 730$  nm and a shoulder at 806 nm. Among the split two peaks (0-0 transitioned  $\lambda_{em} = 730$  nm and 0-1 transitioned  $\lambda_{em} = 806$  nm) relative to a different degree of vibrational coupling, the lower-energy emission peak is probably constituted by the high-frequency normal modes with larger Huang-Rhys factors.<sup>21</sup> Their phosphorescent properties show a structural profile: a short lifetime of 0.88  $\mu s$  for **1**, and a significantly longer 1.02  $\mu s$  **2** or 0.96  $\mu s$  for **3** should result from the Br-facilitated intersystem-crossing. The electron-donating effect of two *tert*-butyl groups should contribute to the shortest one (0.85  $\mu s$ ) for the (**S-L<sup>4</sup>**)-involved **4**. It is worth noting that all the NIR-emissive efficiencies ( $\Phi_{pl}$ ) of the four chiral Pt(II)-complexes are also regulated from the (**S-L<sup>n</sup>**)-involved electronic effects, and their attractive  $\Phi_{pl}$  centered at 731 nm among reported NIR-emissive Pt(II)-complexes<sup>10</sup> just with the MLCT impaction, should be attributed to their large radiative rate constant ( $k_r = 10^5 s^{-1}$ ). The Br-incorporated chiral  $Pt^{2+}$ -complex **2** gives rise to the best NIR-emissive  $\Phi_{pl}$  of 21% due to the largest  $k_r$  ( $2.10 \times 10^5 s^{-1}$ ) while the smallest  $k_{nr}$  ( $7.75 \times 10^5 s^{-1}$ ). The effective radiative transition with the weak vibrational coupling for the four chiral  $Pt^{2+}$ -complexes **1-4** can be further confirmed by their rigidochromic shifts (Figure S8) at 77 K, whereas the (**S-L<sup>n</sup>**)-incorporation profits to effectively sidestep the intramolecular rotation/distortion for more chance to the  $T_1$ -



**Figure 2.** The normalized UV-visible absorption and emission spectra of the chiral Pt(II)-complexes **1-4** in degassed solution or PVK-PBD in solid-state at RT.

**4** in solution show the significantly broadened UV-visible absorption



**Figure 3.** The CP light spectra of the chiral Pt(II)-complexes **1-4** in degassed solution at RT.

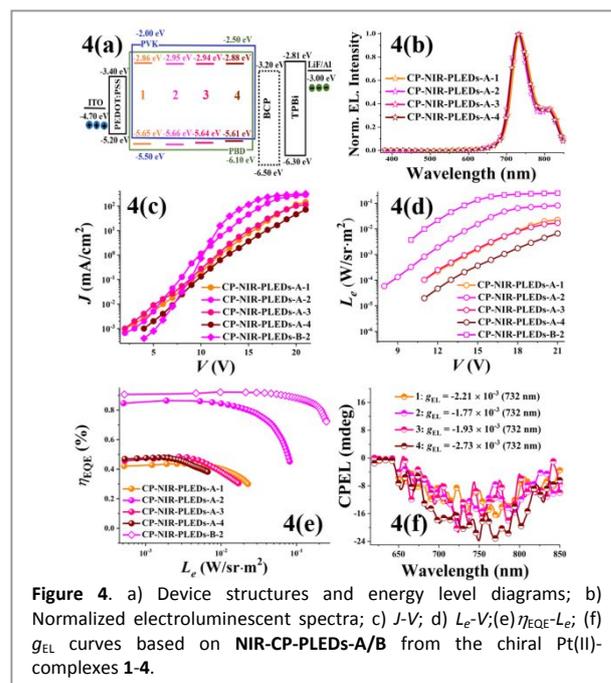
preferred radiative decaying with longer phosphorescent lifetimes.

As to the slight blue-shifts ( $\lambda_{em} = 725\text{--}727\text{ nm}$ ) for **1-2** while red-shifts ( $\lambda_{em} = 741\text{--}751\text{ nm}$ ) of **3-4** in solid-state (Figure S9) relative to that ( $\lambda_{em} = 732\text{ nm}$ ) in solution at RT, it should probably be caused by strong intermolecular interactions (also Figure S4) as in the case of routine  $\text{Pt}^{2+}$ -complexes.

In their CD spectra (also Figure S6), besides the Cotton effects with two positive signals at 259–268 and 328–331 nm and an negative one at 286–313 nm like the free **S-L<sup>n</sup>** ligands, additional broad Cotton-effect bands came from the (**S-L<sup>n</sup>**)-regulated exciton-coupling  $^1\text{LLCT}/^1\text{MLCT}$  transitions with three negative signals and two positive signals within 355–600 nm. As to the emissive dissymmetry factor ( $g_{\text{PL}}$ ; Figure 3) of the (**S-L<sup>n</sup>**)-involved four chiral Pt(II)-complexes, it is at the degree of  $10^{-3}$ , endowing a (**S-L<sup>n</sup>**)-regulated size ( $-2.14 \times 10^{-3}$  (731 nm) for **1**;  $-2.07 \times 10^{-3}$  (729 nm) for **2**;  $-1.76 \times 10^{-3}$  (730 nm) for **3** or  $-2.75 \times 10^{-3}$  (731 nm) for **4**) at the corresponding NIR-emission peak probably due to the differential electron cloud shares<sup>22</sup> especially within the chiral (**S-L<sup>4</sup>**) portion (also Table S4 and Figure S10).

The theoretical studies provided further insights of the photophysical properties of these chiral complexes [Pt(iqbt)(**S-L<sup>n</sup>**)] (**1-4**). The TD-DFT calculations (Table S5 and Figure S10) indicated that each of chiral Pt(II)-complexes **1-4** has the similar electron density distribution patterns to the LUMO+1 or LUMO, corresponding to the major contribution from the (**S-L<sup>n</sup>**) ligand (90.20–93.04%) or the (**iqbt**) ligand (90.40–90.68%), respectively. The electron density of the HOMO is dominated by the ligands with the increased contribution from the (**iqbt**) ligand (43.90%) than that from the (**S-L<sup>1</sup>**) ligand (38.36%) for **1**. This is different to the situation in **2-4** (47.81% for **2**, 59.87% for **3** or 60.48% for **4**) from the (**S-L<sup>n</sup>**) ( $n = 2\text{--}4$ ) ligand. It was found that the introduction of 3-position MeO/*tert*-butyl and/or 5-position Br/*tert*-butyl groups for **2-4** gives a preferred contribution (53.47% for **2**, 62.32% for **3** or 63.88% for **4**) of the (**iqbt**) ligand to their HOMO-1s. The Br substantially causes the destabilizations of both the HOMOs and LUMOs levels while an opposite stabilization from the MeO- or *tert*-butyl- group is observed. Hence, the relatively narrow HOMO-LUMO bandgaps (2.86 eV for **2**; 2.83 eV for **3** or 2.80 eV for **4**) relative to **1** (2.91 eV) were observed in **2-4**. Even though the calculated HOMO  $\rightarrow$  LUMO transition (514 nm for **1**; 524 nm for **2**; 526 nm for **3** or 532 nm for **4**) came mainly (over 97%) from the  $S_0 \rightarrow S_1$  transition, the experimental lower-energy ( $\lambda_{\text{ab}} > 430\text{ nm}$ ; 518–521 nm for **1-4**) absorptions should also include the LLCT/MLCT transitions. The NTO calculations (Table S6 and Figure S11) for  $S_0 \rightarrow T_1$  excitations indicated that about 80% hole orbital contribution from (**S-L<sup>n</sup>**) and 92% particle orbital contribution from the (**iqbt**) ligand for each of the chiral Pt(II)-complexes **1-4**, constituted almost the entire (*ca.* 99%) Hole  $\rightarrow$  Particle transition, suggesting that the  $^3\text{LLCT}$ -dominated and the less  $^3\text{MLCT}$  transitions, should be responsible for their NIR-emissive phosphorescence. This bandgap is quite similar to those determined electrochemically (Figure S12 and also Table S4).

It is of particular interest on the cost-effective and large-area flexible CP-NIR-PLEDs based on these chiral Pt(II)-complexes **1-4**. PVK-PBD was used as the appropriate host due to its good hole-electron transport, and the significant spectra overlap (also Figure 2) between the emission of PVK-PBD and the LLCT/MLCT absorption of the chiral Pt(II)-complexes **1-4** for effective Förster energy transfer<sup>23</sup> from PVK-PBD to the chiral Pt(II)-complex species. Since the LUMO ( $-2.95 \sim -2.86\text{ eV}$ ) and HOMO ( $-5.66 \sim -5.61\text{ eV}$ ) levels of the chiral Pt(II)-complexes **1-4** perfectly within those ( $-2.50 \sim -2.00\text{ eV}$  of LUMOs and  $-6.10 \sim -5.50\text{ eV}$  of HOMOs) of PVK and PBD, respectively, the injected holes from PEDOT:PSS and the injected



**Figure 4.** a) Device structures and energy level diagrams; b) Normalized electroluminescent spectra; c)  $J$ - $V$ ; d)  $L_e$ - $V$ ; e)  $\eta_{\text{EQE}}$ - $L_e$ ; f)  $g_{\text{EL}}$  curves based on NIR-CP-PLEDs-A/B from the chiral Pt(II)-complexes **1-4**.

electrons from TPBi pass through the PVK-PBD host to be trapped by the chiral Pt(II)-complex species, and subsequently, direct charge carrier trapping and recombination should take place. Through the solution-processable procedure with the same configuration shown in Figure 4(a), series of CP-NIR-PLEDs-A/B using the chiral Pt(II)-complexes **1-4** as the dopants with a stipulated doping of 8 wt% are fabricated. The difference lies in the fact that a hole blocking BCP-layer was added in the B series.

As shown in Table S7 and Figures 4(b)/S13, all the normalized electroluminescent NIR-emissive ( $\lambda_{em} = 732\text{ nm}$ ) spectra for the series of CP-NIR-PLEDs-A are well resembled those of chiral Pt(II)-complexes **1-4** in solution at RT, which should arise from the low-concentration dispersion of the dopants **1-4** into the corresponding PVK-PBD host in the absence of those strong intermolecular interactions (also Figure S4). It also shows structure-featureless and an independence on the applied bias voltage. Within Figures 4(c-e) for the  $J$ - $V$ ,  $L_e$ - $V$  and  $\eta_{\text{EQE}}$ - $L_e$  correlations for CP-NIR-PLEDs-A with different chiral Pt(II)-complexes **1-4**, differently from the monotonous increasing of both the  $J$  and  $L_e$  values with the increase of the applied bias voltage, all the  $\eta_{\text{EQE}}$  increase first and insistently decrease throughout the corresponding illuminating regime. The highest device efficiency is observed for the CP-NIR-PLED-A-2, giving the  $\eta_{\text{EQE}}^{\text{Max}}$  of 0.87% with  $L_e = 0.002\text{ W/sr}\cdot\text{m}^2$  at  $J = 6.02\text{ mA/cm}^2$  after 10.5 V of the  $V_{\text{on}}$  (the voltage of  $L_e = 5 \times 10^{-4}\text{ W/sr}\cdot\text{m}^2$ ). In comparison, both the  $\eta_{\text{EQE}}^{\text{Max}}$  and  $L_e^{\text{Max}}$  (0.44–0.49% and 0.007–0.023  $\text{W/sr}\cdot\text{m}^2$ ) of all the other CP-NIR-PLEDs-A based on chiral Pt(II)-complexes **1** and **3-4** are relatively lower than those (0.87% and 0.081  $\text{W/sr}\cdot\text{m}^2$ ) of the CP-NIR-PLED-A-2, which is consistent with the measured order of their  $\Phi_{\text{PL}}$  results afforded by (**S-L<sup>n</sup>**)-involved electronic effect for their chiral Pt(II)-complexes (**1-4**) in solution (also Table S4). Nonetheless, the longest  $T_1$ -decayed phosphorescent lifetime for the (**S-L<sup>2</sup>**)-induced **2** with the 5-position electron-drawing -Br effect, renders its CP-NIR-PLED-A-2 a significant efficiency-roll-off (48%) as compared to those (20–37%) of the other three devices. For optimization, the CP-NIR-PLED-B-2 (also Figure 4(a)) with an additional hole-blocking BCP layer was fabricated. Delightfully, contributing from more excitons confined within the broadening recombination zone,<sup>24</sup> the distinctively

improved electroluminescent NIR-emissive property (also Figures 4(b-e)) is approached. First, besides the illuminating voltage down to 6.0 V, the  $L_e^{\text{Max}}$  is almost three times higher than that of the **CP-NIR-PLED-A-2**. Moreover, the reformed carrier balance within the **CP-NIR-PLED-B-2** also endows an attractive  $\eta_{\text{EQE}}^{\text{Max}}$  up to 0.93% among previous Pt(II)-complex-based NIR-PLEDs.<sup>10,14</sup> Interestingly enough, its efficiency-roll-off (23%) is also relieved. Worthy of notice, as shown in Figure 4(f), the dissymmetry factors ( $g_{\text{EL}}$ , also Table S7) of all the **CP-NIR-PLEDs** are also ruled by the (**S-L**<sup>n</sup>)-relativity also within the 10<sup>-3</sup> order as the reported Pt(II)-complex-based visible CP-OLEDs.<sup>7</sup>

In conclusion, through the involvement of (**S-L**<sup>n</sup>)-N<sup>^O</sup> ancillary ligands with different electronic effects, series of chiral [Pt(iqbt)(**S-L**<sup>n</sup>)] complexes (**1-4**) with NIR-emissive ( $\lambda_{\text{em}} = 729\text{-}731\text{ nm}$ ) and CPL-active properties were obtained. Moreover, based on their doping into the PVK-PBD host, the **CP-NIR-PLEDs** with the much attractive device performance ( $\lambda_{\text{em}} = 732\text{ nm}$ ;  $\eta_{\text{EQE}}^{\text{Max}}$  up to 0.87-0.93%) and the high  $g_{\text{EL}}$  up to 10<sup>-3</sup>, were first reported. This result suggests that the (**S-L**<sup>n</sup>)-involved chiral [Pt(iqbt)(**S-L**<sup>n</sup>)]-complexes are promising candidates for future **CP-NIR-PLEDs**.

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## Table of content only

### Efficient polymer light-emitting diodes (PLEDs) based on chiral [Pt(C<sup>^</sup>N)(N<sup>^</sup>O)]-complexes with near-infrared (NIR) luminescence and circularly polarized (CP) light

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Using chiral NIR-emitters [Pt(iqbt)(S-L<sup>1</sup>)] (**1-4**) as the dopants, their CP-NIR-PLEDs with  $\lambda_{em} = 732$  nm,  $\eta_{EQE}^{Max} = 0.87-0.93\%$  and  $g_{EL}$  up to  $10^{-3}$  are realized.

