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Polymerization in Soft Nanoconfinements of Lamellar and Reverse Hexagonal Mesophases

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9	Abstract
10	This work describes the kinetics of thermal polymerization in nanoconfined domains of lyotropic
11	liquid crystal (LLC) templates by using chemorheological studies at different temperatures. We
12	investigate lamellar and reverse hexagonal LLC phases with the same concentration of
13	monomeric phase. Results show that the mesophase structures remain intact during thermal
14	polymerization with very slight changes in the domain size. The polymerization rate decreases in
15	the nanoconfined structure compared to the bulk state due to the segregation effect that increases

the nanoconfined structure compared to the bulk state due to the segregation effect that increases the local monomer concentration and enhances the termination rate. Additionally, the polymerization rate is faster in the studied reverse hexagonal systems compared to the lamellar ones due to their lower degree of confinement. A higher degree of confinement also induces a lower monomer conversion. Differential scanning calorimetry confirms the obtained results from chemorheology.

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24 Introduction

Self-assembly of amphiphilic molecules in the presence of two solvents (oil/water) leads
 to the formation of lyotropic liquid crystals (LLCs) with the length scale in the range of 2 50 nm, also called mesophases. Direct templating by preformed LLC phases has widely
 been used for producing organic and inorganic mesoporous materials.^{1–7} For making

organic mesoporous media, after preparation of oil/water/surfactant in a desired phase 29 30 state, one of the phases can be polymerized to obtain a mesoporous polymer. Thus, monomers are included in the oil or water phase prior to the preparation of mesophase. 31 In such cases, we deal with polymerization in nanoconfinement. However, templating 32 method for making organic porous polymers is not a straightforward task. In other words, 33 phase separation during polymerization may change the original nanostructure in terms 34 of shape and size.^{8–13} Studying the kinetics of polymerization in nanoconfinement helps 35 in understanding the effect of different types of confinement on the final polymer 36 properties.¹⁴ 37

38 Several studies have been done on the kinetics of polymerization in hard nanoconfinements.^{15–19} At the early stages of polymerization in the anodic aluminum 39 40 oxide templates, it has been observed that the template walls catalyze the initiation reaction and the polymerization rate increases. However, at the late stages of 41 42 polymerization, the termination reaction increases as the likelihood of radicals being in the close proximity increases, which leads to the lower conversion and reaction rate.¹⁵ 43 Confinement is one of the main factors contributing to the polymerization rate. As the 44 system is more confined, the local monomer concentration increases, which leads to 45 higher termination rate and lower polymerization rate.¹⁵ Figure 1a schematically shows 46 the effect of confinement. 47



50 Figure 1.Schematic illustration of: (a) the effect of confinement on the probability of termination 51 reactions, (b) the increase in domain size upon polymerization due to the tendency of polymer 52 chains to form random coil, and (c) the expected dependence of change in domain size upon 53 polymerization to the template modulus (see eq. 1).

In soft nanoconfinements, the template stiffness also plays a role in the polymerization 54 kinetics. Therefore, the degree of confinement, as the prominent factor affecting the 55 polymerization kinetics, is influenced by the size of the confinement and the template 56 57 stiffness. In other words, as the degree of polymerization increases, polymer chains tend towards random coil conformation to minimize their free energy (Figure 1b). If the 58 template stiffness cannot suppress the tendency of the random coil conformation, the 59 domain size will increase. In an ideal case, the higher the elastic modulus of a soft 60 template, the lower is the change in the domain size upon polymerization, thus, the higher 61 is the effect of confinement at a fixed initial domain size (Figure 1c). Hard templates can 62 be considered to have very large modulus with no change in the domain size, whereas a 63

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template with zero modulus will be destroyed upon polymerization (maximum change in
 domain size). Assuming negligible density change during polymerization, therefore, we
 can propose the following scaling relationship:

$$67 \qquad \frac{d_f - d_0}{d_0} \propto \frac{1}{E^n} \tag{1}$$

where d_f and d_0 are the final and initial domain sizes, *E* is the elastic modulus, and *n* has a value equal or higher than unity. A deviation from this scaling suggests that properties of the template (e.g., surface tension) change during polymerization and extra care should be taken to consider the elastic modulus or change in the domain size as a measure of template stiffness.

73 Kinetics of photopolymerization in the hexagonal and lamellar mesophases formed with small molecule surfactants have been studied via differential scanning calorimetry (DSC). 74 However, the obtained structures from such mesophases after polymerization have 75 usually a domain size bigger than 100 nm,²⁰ and/or have been disrupted due to the 76 polymerization-induced phase separation.^{20,21} Lester et al.² have shown that reactions in 77 the ordered structure of LLCs are highly dependent on the type and degree of order, and 78 are significantly different from the isotropic state. This phenomenon can be attributed to 79 a number of factors including diffusional limitation which reduces termination rates and 80 the segregation of the monomeric species that increases both the apparent propagation 81 and termination rates.^{20,22,23} It has been shown that the rate of polymerization and 82 mesophase structure play important roles in the final properties of resulting polymer.²³ 83 84 Lester et al. observed that for the lyotropic lamellar mesophases made from fluorinated 85 amphiphilic monomers, the termination rate decreases and the polymerization rate increases. Cubic structures, on the other hand, show the slowest kinetics.²³ In another 86 work, they have shown that the polymerization rate of hydrophobic monomers is higher 87 88 in the micellar cubic structures compared to the lamellar and hexagonal phases. They 89 attributed this behavior to the increase in the rate of propagation in the cubic micelles due to higher local monomer concentration as compared to the other systems.² 90

Pluronic block copolymers are amphiphilic molecules that are widely used as
 surfactants.^{24–27} Photopolymerization of hydrophilic and hydrophobic monomers in LLC

structures formed by Pluronic L92 has been studied.²⁸ It has been found that the rate of 93 polymerization of hydrophobic monomers in normal phases increases due to the 94 segregation of monomers during photopolymerization. On the other hand, it has been 95 observed that for hydrophilic monomers, the polymerization rate is higher in the inverse 96 phases, such as inverse hexagonal. These results show that the monomer segregation 97 into confined domains of LLCs affects the polymerization rate due to the increases in the 98 local concentration of monomers and/or radicals. Additionally, Lester et al. observed that 99 the diffusion of propagating sites will be limited in confinement.²⁹ 100

¹⁰¹ Zhao et al.¹⁶ have studied the polymerization of methyl methacrylate in hard nanopores ¹⁰² and found that smaller pore sizes result in shorter autoacceleration times. They have also ¹⁰³ shown that hydrophilic pores have a more significant effect in reducing the polymerization ¹⁰⁴ activation energy compared to the hydrophobic ones. Silies et al.³⁰ compared the ¹⁰⁵ polymerization rate of zwitterionic monomers inside and on the outer surface of a ¹⁰⁶ mesoporous film. They showed that the confinement in mesopores limits the diffusion of ¹⁰⁷ monomers and free radicals into the pores and influences the termination rate.

Thermal polymerization is favored over photopolymerization for large-scale production of mesoporous structures. However, there are only few studies addressing the kinetics of thermal polymerization in different mesophases.^{31,32} DePierro et al.¹⁰ have compared the structures resulting from photo and thermal polymerizations of acrylamide mixtures in different mesophases formed from small molecule surfactants. They have found that thermal polymerization yields less ordered to disordered structures with larger feature sizes due to the slower kinetics of polymerization.

115 While chemorheology has been used for studying the kinetics of polymerization for decades,^{33–36} only a few studies have been done on chemorheology 116 in nanoconfinements. Peng et al.³⁷ have studied the photopolymerization of a polymerizable 117 118 surfactant in binary mixtures of water and Brij 97 in order to maintain the structure upon curing. They have measured the change in the mechanical properties during 119 photopolymerization using in situ photo-rheology and found out that the dynamic moduli 120 121 and viscosity of polymerizing system increase upon exposure to UV light. However, they have not quantified the kinetics of polymerization using rheology. Chow et al.^{38,39} have 122

studied the chemorheology of rod-like poly(p-phenylenebenzobisthiazole) made by polycondensation reactions. They have observed that when the polymer mixture becomes isotropic, the polymerization rate increases as alignment of rods facilitates the condensation reaction.

Recently, we have shown that mesophases of Pluronic/water/monomer with lamellar and hexagonal structures can retain their mesostructure upon thermal polymerization.⁴⁰ In this paper, we study the rheological behavior of mesophases during polymerization. Mesophases with the same amount of monomer phase but different structures are used to cancel the effect of monomer and initiator concentration on the polymerization rate. DSC is used to confirm the kinetics of polymerization derived from chemorheology.

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134 **Experimental**

Pluronic block copolymers, P84 (M_w=4200 g/mol), L121 (M_w=4400 g/mol), and L64 135 $(M_w = 2900 \text{ g/mol})$ are kindly provided by BASF. Butyl acrylate ($\geq 99\%$, Sigma-Aldrich) and 136 137 ethylene glycol dimethcrylate (EGDMA, purified, Electron Microscopy Sciences) are used as monomer and crosslinker, respectively. Azobisisobutyronitrile (AIBN, 98%, Sigma-138 Aldrich) is used as the thermal initiator. Deionized (DI) water (0.055 µs/cm, EMD 139 Millipore Direct-Q3) is used as the aqueous phase. The monomer and crosslinker are 140 141 purified by passing through a silica column. All other chemicals are used as received without further purification. 142

Method of mesophase preparation has been explained before.⁴⁰ Simply, desired amounts 143 of components are mixed using centrifugation at alternative direction until a transparent 144 mesophase is obtained. Table 1 shows the ratio of materials used in each sample. The 145 monomer to crosslinker weight ratio is held constant at 3:1 and the initiator concentration 146 is 3 wt% of the monomer in all samples. Samples are formulated in a way that the 147 monomer plus crosslinker concentration is constant, but different mesostructures are 148 made by varying water/block copolymer ratio. It should be noted that it is almost 149 impossible to have different mesostructures with the same Pluronic at a fixed oil phase 150 (monomer) concentration. Thus, we have used different Pluronic block copolymers while 151

keeping the concentration of the polymerizing phase at 25 wt%. We prepare two lamellar, L_α, mesophases with P84/water/monomers 60/15/25 wt% and L121/water/monomers 57/18/25 wt%, and two inverse hexagonal, H₂, mesophases with P84/water/monomers 40/35/25 wt% and L64/water/monomers 55/20/25 wt%.

Small angle X-ray scattering (SAXS) is used to confirm the mesostructure of each sample.
 Measurements use Cu Kα X-rays radiation source with the wavelength of 1.54184 Å using
 Bruker Nanostar System. All samples are tested before and after polymerization to ensure
 desired mesophases are made and samples retained their structures upon
 polymerization.

161 A stress-controlled rheometer DHR-3 (TA Instruments, New Castle, DE) is used to study the rheological behavior and chemorheology of mesophases. A 40 mm sand blasted 162 parallel plate geometry with 1 mm gap is used in all experiments. All tests are performed 163 in the linear viscoelastic region (0.5% strain, confirmed from amplitude sweep tests). First, 164 dynamic frequency sweep tests are performed at 25 °C in the frequency range of 0.1 to 165 166 600 rad/s. For chemorheology studies, a solvent trap filled with DI water is used. Time tests in small amplitude oscillatory shear mode are done on mesophases at three different 167 temperatures, 60, 65, and 70°C to determine the kinetics of polymerization through 168 evolution of elastic and loss moduli.³³ Data are collected in the linear viscoelastic region 169 (strain amplitude of 0.5%) at constant frequency of 1 Hz. Polymerization of the pure 170 171 monomer phase, consist of butyl acrylate, EGDMA (33 wt% of butyl acrylate), and AIBN (3 wt% of butyl acrylate) is also studied as the control sample to define the kinetics of 172 173 polymerization in the non-confined state.

Isothermal DSC is carried out utilizing Q2000 (TA Instruments, New Castle, DE). Approximately 10 mg of mesophase is placed in the aluminum pans and the heat of reaction is recorded with time. All measurements are performed under a nitrogen gas atmosphere at elevated temperatures (60, 65, and 70 °C) to determine the rate of polymerization according to the procedure established by Guymon and coworkers.^{1,2,21,22} Having the heat flow, ΔQ , the polymerization rate, R_p , can be calculated as:

180
$$\frac{R_p}{[M]_0} = \Delta Q \left[\left(\frac{M_W}{n \Delta H_p m} \right)_{monomer} + \left(\frac{M_W}{n \Delta H_p m} \right)_{crosslinker} \right]$$
(2)

where M_w , $[M]_0$, ΔH , n, and m are molecular weight, initial concentration, theoretical 181 reaction enthalpy (86,200 J/mol for acrylate and 56,000 J/mol for methacrylate),⁴¹ 182 functionality, and mass, respectively.²⁰ Polymerization of pure oil phase is also studied 183 184 as control sample to define the kinetics of polymerization in the non-confined state. The polymerization rate is normalized to the total reactive species' concentration in the 185 formulation. The reaction temperature is maintained constant (within ± 0.1 °C) during the 186 measurements. The degree of monomer conversion is calculated by integrating the area 187 between the DSC curves and the baseline established by extrapolation from the trace 188 189 produced after complete polymerization. The final conversion is experimentally obtained by washing the mesophases with soxhlet to remove residual monomer, initiator, and 190 crosslinker. There is also a possibility of block copolymer removal during washing. 191 Samples are first washed with water for 12 h, and then with methanol for another 24 h. 192 Thermal gravimetric analysis (TGA) is done on the washing solvents after soxhlet to 193 confirm that Pluronic block copolymer has not been washed out, as shown in Figure S1 194 in the Electronic Supplementary Information (ESI). After washing, the samples are dried 195 in the vacuum oven at 40 °C for 48 h. Dried samples are weighed and the conversion is 196 obtained using the gravimetric analysis. 197

198

199 **Results and discussion**

200 **SAXS**

To confirm the retention of mesostructures upon thermal polymerization, SAXS studies

are done on the mesophases before and after polymerization as shown in Figure 2.



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Figure 2. 1D SAXS graphs of mesophase systems with different Pluronic block copolymer/water/oil (wt%) compositions before and after polymerization: (a) P84-L_{α} (60/15/25), (b) L121-L_{α} (57/18/25), (c) P84-H₂ (40/35/25), and (d) L64-H₂ (55/20/25). Oil phase consists of butyl acrylate, EGDMA, and AIBN.

Lamellar structures have 1:2:3:4:... relative positions of Bragg peaks (q/q*), while hexagonal mesostructures have 1: $\sqrt{3}$:2: $\sqrt{7}$:... relative positions of Bragg peaks, where q* is the principal peak. Polymerized samples at elevated temperatures show the same pattern in the peak positions, while there is a slight shift of peaks to the left. According to Bragg's equation, in the lamellar mesophases, lattice parameter, *d*, can be calculated as:⁴²

$$214 \qquad d = \frac{2\pi}{q^*} \tag{3}$$

Apolar domain volume fraction, ϕ , is defined as the volume fraction of the polymerizing phase and the PPO block, and the polar domain volume fraction, $1 - \phi$, is the volume

fraction of the water and the PEO block. Knowing the lattice parameter and the volume fractions, we can calculate the lamellar apolar domain size, δ , as:

$$\delta = \phi d \tag{4}$$

In case of hexagonal mesophases, lattice parameter, a, and apolar domain size, α , are calculated as follows:

222
$$a = \frac{4\pi}{\sqrt{3}q^*}$$
 (5)

223
$$\alpha = a \left(\frac{\sqrt{3}}{2\pi}\phi\right)^{1/2}$$
(6)

In all the above calculations, we assume that different phases of water, oil, PEO, and 224 PPO are completely segregated and the bulk density of each component is 225 considered.^{40,42} This assumption might not be accurate as both PEO and PPO partially 226 227 segregate in water and oil, respectively and EGDMA may have a tendency towards water phase. However, the assumption helps with simplifying the calculations without losing a 228 great deal of information. ^{40,43} The calculated parameters obtained from SAXS 229 experiments on various samples are schematically shown in Figure 3 and their values are 230 summarized in Table 1. Apolar domain size and its change after polymerization are two 231 important factors in the kinetics of polymerization and will be discussed later. 232

Critical molecular weight entanglement (M_c) for polyethylene oxide (PEO) and 233 polypropylene oxide (PPO) are 10,000 g/mol and 7,000 g/mol, respectively.^{44,45} All the 234 Pluronic block copolymers in our study have the molecular weight well below M_c. 235 236 Therefore, block copolymers are not entangled in the system. Additionally, the radius of gyration of Pluronic P84, Pluronic L121, and Pluronic L64 block copolymer are 237 approximately 17 Å,²⁵ 18 Å,⁴⁶ and 19 Å,⁴⁷ respectively, which are smaller than the domain 238 size of the micelles. Thus, there is no chain crossing the oil phase from one hydrophilic 239 domain to another one. 240

Sample	Pluronic/water/ monomers (wt%)	φ	<i>d</i> or a, unpolymerized (nm)	δ or α, unpolymerized (nm)	Intermicellar distance (nm)
P84- L _α	60/15/25	0.62	6.0	3.7	0.9
L121-L _{α}	57/18/25	0.31	10.0	3.1	1.6
P84- H ₂	40/35/25	0.50	6.6	2.4	0.6
$L64-H_2$	55/20/25	0.60	7.4	3.0	0.8

Table 1. Composition of the samples and their calculated SAXS parameters.

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242



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Figure 3. Schematic illustration of the parameters obtained from SAXS and confinement size in (a) lamellar and (b) reverse hexagonal mesophases.

The polymerized samples show a slight shift in the peaks to smaller scattering vectors, 247 248 which indicates an increase in the lattice parameter, thus, in the apolar domain size. The change in domain size is attributed to the competition between thermodynamics and 249 kinetics. By progression of polymerization, the molecular weight and degree of 250 polymerization, N, increase, thus, the enthalpic penalty, χN , outweighs the entropic 251 contribution to the Gibbs free energy. Therefore, the system is driven towards the 252 increase in domain size (and ultimately phase separation). On the other hand, the density 253 254 of the monomer phase increases upon polymerization ($\Delta \rho \sim 10\%$), which leads to the shrinkage and decrease in the domain size. The presence of Pluronic block copolymers 255 with slow dynamics decreases the rate of phase separation. In addition, crosslinking 256 arrests and preserves the structure during thermal polymerization. 257

259 Rheology

Amplitude sweep test is done on samples at different temperatures (25, 35, and 45°C) 260 and frequencies (1, 10, and 100 rad/s) to verify the linear viscoelastic region. Data are 261 shown in Figures S2-S9 in the ESI, where we conclude that all samples are in the linear 262 263 regime at 0.5% strain. The real-time monitoring of stress and strain signals by time also 264 confirms the linear behavior during the polymerization (i.e., the waves remain sinusoidal). Frequency sweep results for lamellar and reverse hexagonal mesophases in Figure 4 265 show that all mesophases have solid-like behavior, where the elastic modulus is higher 266 than the loss modulus in the studied range of frequency.⁴⁸ Such behavior has been 267 observed for suspensions, block copolymer solutions, and highly concentrated 268 emulsions.^{49–51} In the systems under study, the solid-like behavior is due to the high 269 concentration of block copolymer and the compact LLC structures. It has been shown 270 that lamellar structures have one order of magnitude lower elastic modulus (G') when 271 compared with hexagonal mesophases.^{52,53} As it is evident from Figure 4, the elastic 272 modulus in (c) and (d) plots (H₂ samples) is higher than that of (a) and (b) ones (L_{α} 273 samples). Comparing the lamellar samples, P84-L_{α} (60/15/25) shows higher elastic 274 275 modulus compared to L121-L_a (57/18/25), which is due to the smaller domain size and 276 intermicellar distance in P84-L_{α} sample. According to a model we have developed 277 recently on the basis of van der Waals interactions (which is the main source of intermicellar interactions),⁴³ the elastic modulus is highly dependent on the reciprocal of 278 intermicellar distance. There are two intermicellar distances D_1 and D_2 across the oil 279 phase and water phase in the lamellar mesophases, respectively (see Figure 3). 280 Considering that the van der Waals forces between planes scale with D^{-3} ,⁵⁴ an average 281 intermicellar distance in the lamellar systems, D_L , can be defined as follows: 282

283
$$\frac{1}{D_L^3} = \frac{1}{D_1^3} + \frac{1}{D_2^3}$$
 (7a)
284 $D_1 = \varphi_{oil} d$ (7b)

$$285 D_2 = \varphi_{water} d (7c)$$

where φ_{oil} and φ_{water} are the volume fractions of oil and water, respectively. Average intermicellar distance in P84-L_{α} sample is smaller than that of L121-L_{α} sample and consequently the elastic modulus is higher in P84-L_{α} sample.

Additionally, in the reverse hexagonal mesophases, P84-H₂ (40/35/25) has higher elastic modulus than L64-H₂ (55/20/25). Intermicellar distance in reverse hexagonal systems, D_H , can be defined as follows:⁴³

$$292 D_H = a - 2M_H (8a)$$

293
$$M_H = a_{\sqrt{\frac{\sqrt{3}}{2\pi}}}(\varphi_{Pluronic} + \varphi_{water})$$
(8b)

where $\varphi_{Pluronic}$ is the volume fraction of Pluronic block copolymer. The intermicellar distance in P84-H₂ sample is smaller than that of L64-H₂ sample. Thus, the elastic modulus is higher in P84-H₂ sample.⁴³

A shallow minimum is observed in the loss modulus (G") curves of all samples that is the 297 characteristic of polymeric gels and has been observed for emulsions and soft-glassy 298 materials as well.^{55–59} The minimum in the loss modulus shows the presence of two 299 300 relaxation behaviors in the system and the transition from α -relaxation observed in low 301 frequencies (long time, related to large domain sizes, i.e., grain size scale) to β -relaxation observed in high frequencies, (>10³ rad/s), (short time, related to small domain sizes, i.e., 302 micelle size scale).⁶⁰ We have recently shown that the higher is the probability of finding 303 defects in the LLC systems, the higher is the angular frequency where the minimum in G" 304 takes place.60 305

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307



309Figure 4. Frequency sweep curves of mesophase systems with different Pluronic block310copolymer/water/oil (wt%) compositions: P84-L_{α} (60/15/25), L121-L_{α} (57/18/25), P84-H₂311(40/35/25), and L64-H₂ (55/20/25).

Variations of the dynamic moduli with time at three different temperatures for the lamellar and reverse hexagonal samples are shown in Figure 5. Bulk data are provided in Figure 6a. Three stages of polymerization can be found: (i) induction, where G' is relatively constant in the beginning of the experiment for a period indicated as t_{in} ; (ii) polymerization, where G' sharply increases by time; and (iii) final curing and end of reaction, where G' reaches a plateau.

Gel point terminology in the conventional chemorheology literature cannot be applied here because of the solid-like behavior of the samples (see Figure 3).⁶¹ Additionally, the G'=G'' or tan δ =1 criterion⁶² cannot be used as no crossover between dynamic moduli is observed. Considering the gel point as the time when G'' is maximum does not work either since G'' does not show a maximum in the samples under study.^{33,62,63} Therefore, evolution of dynamic moduli and complex viscosity are fitted using Arrhenius type models to determine the kinetic constants.⁶⁴



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Figure 5. Variation of dynamic moduli with time at three different temperatures (60, 65, and 70 °C) for (a) P84-L_{α}, (b) L121-L_{α}, (c) P84-H₂, and (d) L64-H₂ samples.

329



Figure 6. Variation of (a) dynamic moduli and (b) complex viscosity with time at three different temperatures (60, 65, and 70 °C) for bulk sample.

Activation energy in the induction step can be modeled using the Arrhenius model as follows for the induction time, t_{in} :³³

$$335 t_{in} = A e^{\binom{E_{in}}{RT}} (9)$$

where E_{in} , R, T, and A are the induction activation energy, gas constant, absolute temperature, and the pre-exponential factor, respectively. Having t_{in} at different temperatures, we can calculate the activation energy from the slope of $\ln t_{in}$ versus $\frac{1}{RT}$ graph (Figure 6).



Figure 7. The effect of temperature on the induction time of the polymerization in bulk sample and mesophases of P84-L_{α}, L121-L_{α}, P84-H₂, and L64-H₂. The slopes of the lines are used to estimate the induction activation energy.

The values of the induction activation energy for different samples are shown in Table 2. 344 The activation energies for confined samples are higher than the bulk sample that is due 345 to the higher viscosity of confined mesophases that influences the diffusion-controlled 346 initiation in the systems. The lamellar samples show slightly lower induction activation 347 energy compared to the reverse hexagonal ones because of their lower viscosity as 348 confirmed from rheology. It should be noted that having a relatively large uncertainty in 349 the induction times of the lamellar and reverse hexagonal mesophases (Figure 7), caution 350 should be taken when comparing the activation energies of these two systems. 351

To calculate the polymerization rate, the evolution of LLC complex viscosity, $|\eta^*|$, during polymerization (Figure 8) is used. The first order double Arrhenius model is used to describe the chemorheological data:^{33,64}

355
$$\ln |\eta^*| = \ln |\eta^*_v| + \frac{E_v}{RT} + tK$$
 (10)

where $|\eta_v|$ is the initial complex viscosity before polymerization, and E_v is the viscous activation energy. *K* is the polymerization rate constant, which is related to the curing temperature as:

359
$$K = k_0 \exp\left(\frac{-E_p}{RT}\right)$$
(11)

where E_p is the activation energy of polymerization.

361

Table 2. Model parameters for the induction and propagation steps of polymerization in the bulk,
lamellar and reverse hexagonal samples.

Sample	T (°C)	E _{in} (kJ/mol)	<i>K</i> (10⁻ ³/s)	E _p (kJ/mol)	Confinement size (nm)	Change in apolar domain size (%)	Complex viscosity (Pa.s) at ω=1 Hz
Bulk	60 65 70	38±11	21.6±1.1 24.1±1.0 32.4±1.4	38.4±0.7	NA	NA	NA
Ρ84- L _α	60 65 70	104±12	4.7±0.2 6.5±0.3 9.3±0.2	64.8±2.1	0.8	0.12	9.55
L121-L _α	60 65 70	96±27	4.6±0.1 7.2±0.4 9.4±0.2	67±0.1	1.2	1.85	1.86
P84- H ₂	60 65 70	141±13	5.4±0.3 8.5±0.1 11.8±0.3	74.4±3	1.3	11.02	384.86
L64-H ₂	60 65 70	138±6	5.5±0.2 8.0±0.1 11.2±0.2	69.4±1.8	1.5	0.09	49.93

364



Figure 8. Evolution of complex viscosity at different curing temperatures for (a) P84-L_{α}, (b) L121-L_{α}, (c) P84-H₂, and (d) L64-H₂ samples. The dashed lines show the fitted slope in the propagation step.

From Figure 7, the initial complex viscosity of the mesophase samples, shown by an 370 arrow in the graphs, does not change with temperature, which can be attributed to the 371 confined structure of LLCs that hinders the free movement. Therefore, thermal 372 dependency of their viscosity is negligible ($E_v = 0$). The polymerization stage of $|\eta^*|$ 373 versus time is fitted with eq. (10) and the model parameters are summarized in Table 2. 374 Refer to Figure 6b for the bulk data. The polymerization rate in the reverse hexagonal 375 confinement is higher than that of the lamellar ones. We believe the degree of 376 confinement ^{2,20} contributes to the higher polymerization rate in the reverse hexagonal 377 mesophases. 378

As discussed in the introduction, the degree of confinement can be related to the elastic 379 modulus ^{65,66} and the domain size ^{67,68} of confinement. From the rheological data, the 380 381 reverse hexagonal mesophases have higher elastic modulus compared to the lamellar ones. Therefore, according to eq. (1) we ideally expect to see a smaller change in the 382 domain size of H₂ samples upon polymerization than the L_{α} ones. However, analysis of 383 SAXS results in Table 2 shows that the P84-H₂ mesophase has a higher change in 384 domain size upon polymerization compared to the P84-L_{α} sample. Deviation from eq. (1) 385 can be attributed to dynamic changes in the interfacial properties of the mesophases and 386 shows that factors other than elastic modulus are responsible for the change in domain 387 size. The change in domain size results show that the stiffness contribution to degree of 388 confinement is lower for reverse hexagonal mesophases than the lamellar ones. 389

Additionally, as shown in Figure 3, the size of confinement in the reverse hexagonal and lamellar systems can be calculated from the radius of the biggest circle ($R_{h, max}$ and $R_{l, max}$) that can be contained in the apolar domain (the bigger the sphere, the bigger is the confinement size). The radius of such circle is equal to the half of the monomeric domain size in the lamellar samples, as:

395
$$R_{l,max} = \frac{D_1}{2}$$
 (12)

whereas for reverse hexagonal mesophases, $R_{h, max}$ is obtained as follows:

$$R_{h,max} = \sqrt{\frac{A_h}{\pi}}$$
(13a)

398
$$A_h = \frac{a^2 \sqrt{3}}{4} - \frac{\pi (M_H)^2}{2}$$
 (13b)

³⁹⁹ Parameters are schematically shown in Figure 3. Confinement sizes are calculated for all ⁴⁰⁰ samples based on the SAXS data and reported in Table 2 (detailed calculations are ⁴⁰¹ provided in the ESI). We observe that for the same polymerizing chain, the lamellar ⁴⁰² mesophases provide a smaller confinement size, compared to the hexagonal ones. ⁴⁰³ Therefore, by considering both size and stiffness contributions to the confinement, we ⁴⁰⁴ conclude that the P84- H₂ has the lowest degree of confinement and the P84- L_{α} has the ⁴⁰⁵ highest one.

Our results in Table 2 show that the polymerization rate decreases with increasing the 406 407 degree of confinement. In the confined structures, the probability of two macroradicals to 408 react with each other increases that leads to a higher termination rate. Therefore, the rate of polymerization as well as the conversion decrease.¹⁵ In fact, we can consider the 409 mesophase polymerization similar to the polymerizations of a dispersed phase in 410 411 miniemulsions and emulsions. In such systems, the concentration of radicals in the dispersed phase varies depending on the volume of the monomer droplets/polymer 412 particles dispersed in the continuous phase.^{69,70} In this work, rather than being confined 413 inside droplets of a dispersed phase, the reaction is confined within the structures of the 414 lamellar and reverse hexagonal mesophases. Salsamendi and coworkers have 415 considered the confinement effects experienced by propagating radicals to explain the 416 significantly reduced rate of polymerization.⁷¹ They developed a mathematical model to 417 take into account the segregation effects on the polymerization rate. Their segregation 418 419 model predicts that bulk free radical polymerization would proceed at a much faster rate compared to nanoconfined structures. Our data show that the polymerization proceeds 420 421 significantly faster in the bulk state compared to the reverse hexagonal and lamellar samples. The lamellar systems have the slowest polymerization rates. Therefore, our 422 423 results are in agreement with the segregation model, which predicts that as the degree of confinement increases, the probability of the termination increases due to the higher local 424 425 radical concentration resulting in a slower rate of polymerization.⁷¹ Comparing the polymerization rates of two hexagonal samples together (and likewise the two lamellar 426 427 ones together) suggests that the major contribution to the degree of confinement comes from the size rather than modulus/change in domain size. 428

Chemorheology results show that final values of dynamic moduli increase with curing 429 430 temperature (Figure 5). Mechanical properties of the crosslinked polymer control the elastic modulus of each sample that in turn is proportional to the reaction conversion. In 431 order to confirm this hypothesis, the conversion of the samples cured at different 432 temperatures is measured using the gravimetric technique and plotted against curing 433 temperature (Figure 9). After the polymerization is ended, we measured the rheological 434 properties of each cured sample at different temperatures to investigate the temperature 435 dependency of the final elastic modulus (data shown in Figure S10 of the ESI). We found 436

that the elastic modulus of polymerized samples decreases with increasing temperature 437 from 25 to 70 °C, because the polymer network becomes softer by increasing the 438 439 temperature. In addition, according to poroelasticity theory, the decrease in elastic modulus of polymerized mesophases could be due to the decrease in the water viscosity 440 by temperature.^{72,73} Final elastic modulus of the samples, all measured at 70 °C (to 441 remove the effect of temperature), are plotted against the conversion in Figure 9b. The 442 results suggest that G'_{final} is linearly correlated with the reaction conversion within the 443 studied range of temperature. The enhanced conversion with raising curing temperature 444 is attributed to the increase in polymerization rate and monomer diffusion coefficients.³³ 445





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451 **DSC Measurements**

To confirm the kinetic parameters obtained from chemorheological studies, DSC is performed on two typical samples with lamellar and hexagonal structures (P84-L_{α} and P84-H₂). The polymerization rate versus time (Figure 10) and conversion versus time plots (Figure 11) confirm that the radical polymerization in mesophases is influenced by the diffusion-controlled phenomena.



In the first stage of polymerization (low conversions), an almost linear dependence of conversion, and an approximately constant R_p appears, indicating purely chemicalcontrolled nature of the polymerization.⁷⁴ In the region of 10-20 % conversion, a sharp increase in the reaction rate (autoacceleration) starts followed by an increase in the conversion values.^{75,76} The autoacceleration (gel effect) is attributed to the effect of diffusion-controlled phenomena on the termination reaction.

By considering the steady-state hypothesis for the free radical concentration, the polymerization rate, R_p , is given as a function of conversion, X:⁷⁶

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$$\frac{R_p}{[M]} = \frac{dX}{dt} = k_p \left(\frac{fk_d}{k_t}\right)^{1/2} [I]^{1/2} (1-X) \cong K' (1-X)$$
(14a)

473
$$K' = k_p \left(\frac{fk_d}{k_t}\right)^{1/2} [I]^{1/2}$$
 (14b)

where k_d represents the kinetic rate constant of initiator decomposition and f is the initiator efficiency. The propagation and termination rate constants are k_p and k_d , respectively.

To investigate the effect of temperature on the reaction kinetics, equation (14a) can be integrated by considering that all the kinetic rate coefficients, initiator concentration, and efficiency are constant. Therefore, an expression which directly correlates the monomer conversion with an observed overall kinetic rate coefficient, K', will be obtained:

$$481 -ln(1-X) = K't (15)$$

It should be noted that mentioned assumptions are valid only for the low degrees of monomer conversion.⁷⁶ The slope of the initial linear part (between 2 to 10 % where autoacceleration is negligible) of the plot of -ln(1 - X) versus *t*, gives the overall kinetic rate constant.^{16,76} Accordingly, the overall kinetic rate values are measured at different temperatures for the bulk state and two typical samples with lamellar and hexagonal structures. The overall activation energy of polymerization is obtained by considering an Arrhenius type dependency of reaction rate to temperature.

Table 3 shows the kinetic parameters derived from DSC. As seen, the polymerization rate 489 490 in nanoconfined structures is significantly (one order of magnitude) lower than that of the bulk polymerization which is in agreement with chemorheology results. As mentioned 491 earlier, the DSC results (Figure 10 and 11) show that the radical polymerization in 492 mesophases is controlled by the diffusion-controlled phenomena. The activation energies 493 494 obtained from chemorheology (Table 2) and DSC (Table 3) experiments show that the reverse hexagonal mesophases have higher activation energy than lamellar ones (bulk 495 has the lowest activation energy). This is attributed to the higher viscosity of the reverse 496 hexagonal systems compared to the lamellar samples at the early stages of 497 polymerization. On the other hand, it is observed that by increasing the degree of 498 confinement in the system the gel effect decreases during the polymerization. It is 499 confirmed that the lamellar mesophase with the highest degree of confinement has the 500

Iowest conversion (Figure 9) and polymerization rate, which are induced by the increasein the termination rate.

503 These results show that the confinement effect is competing with the gel effect when the 504 polymerization is proceeding within the mesophase system. The high degree of 505 confinement increases the probability of two radicals to react with each other and 506 consequently increases the termination rate.

- 507
- 508 509

Table 3. Activation energy and kinetic rate constants of polymerization for bulk and confinedstructures at different temperatures derived from DSC measurements.

Sample	E (kJ/mol)	<i>K</i> ′ (10 ⁻³ /s) (60	K′ (10⁻³/s) (65	<i>K</i> ′ (10 ⁻³ /s) (70
		°C)	°C)	°C)
P84- H ₂	84	3.6	8.5	9.1
Ρ84- L _α	75	3.4	7.6	8.2
Bulk polymerization	52	21.4	25.3	37.3

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511

512 Conclusion

In this paper, we described a detailed analysis of the thermal polymerization kinetics of 513 monomers in ternary lyotropic liquid crystal systems. The degree of confinement in these 514 systems is controlled by the confinement size and the stiffens of the template. The 515 516 stiffness can be measured by the elastic modulus or change in domain size upon 517 polymerization, which are inversely proportional except if the properties of the template change during polymerization. We discussed that the reverse hexagonal structures offer 518 less degree of confinement than the lamellar samples. Three stages of polymerization 519 (i.e., induction, polymerization and final curing) were observed in the chemorheology. We 520 521 found that the viscosity of the mesophases contributes to the initial and total activation energy of the polymerization. The results from the chemorheology and DSC also showed 522 523 that the rate of polymerization decreases significantly in confined structures compared to the bulk polymerization. Additionally, the polymerization rate was higher in reverse 524 525 hexagonal mesophases than that of the lamellar structures. This phenomenon was

attributed to the segregation effects in which as the degree of confinement increases, the probability of the termination increases due to the higher local radical concentration. At the final stages of the polymerization, a lower limiting conversion was observed in the lamellar systems compared to the reverse hexagonal ones, because the higher degree of confinement in the lamellar mesophases leads to a higher rate of termination in the late

531 stages of polymerization.

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543 **References**

- 544 1 B. S. Forney and C. A. Guymon, *Macromolecules*, 2010, **43**, 8502–8510.
- 545 2 C. L. Lester, C. D. Colson and C. A. Guymon, *Macromolecules*, 2001, **34**, 4430–4438.
- W. Srisiri, T. M. Sisson, D. F. O'Brien, K. M. McGrath, Y. Han and S. M. Gruner, *J. Am. Chem. Soc.*, 1997, **119**, 4866–4873.
- 548 4 C. T. Kresge, M. E. Leonowicz, W. J. Roth, J. C. Vartuli and J. S. Beck, *Nature*, 1992,
 549 359, 710–712.
- J. S. Beck, J. C. Vartuli, W. J. Roth, M. E. Leonowicz, C. T. Kresge, K. D. Schmitt, C. T.
 W. Chu, D. H. Olson and E. W. Sheppard, *J. Am. Chem. Soc.*, 1992, **114**, 10834–10843.
- 552 6 P. Feng, X. Bu and D. J. Pine, *Langmuir*, 2000, **16**, 5304–5310.
- 553 7 S. Grubjesic, B. Lee, S. Seifert and M. A. Firestone, *Soft Matter*, 2011, **7**, 9695.
- K. S. Worthington, C. Baguenard, B. S. Forney and C. A. Guymon, *J. Polym. Sci. Part B Polym. Phys.*, 2017, **55**, 471–489.

L. Sievens-Figueroa and A. Guymon, *Macromolecules*, 2009, **42**, 9243–9250.

M. A. DePierro, K. G. Carpenter and C. A. Guymon, Chem. Mater., 2006, 18, 5609–5617. J. Zhang, Z. Xie, M. Hoang, A. J. Hill, W. Cong, F. H. She, W. Gao and L. X. Kong, Soft Matter, 2014, 10, 5192-5200. J. Zhang, Z. Xie, A. J. Hill, F. H. She, A. W. Thornton, M. Hoang and L. X. Kong, Soft Matter, 2012, 8, 2087–2094. C. C. Co, Soft Matter, 2008, 4, 658. N. Massad-Ivanir, T. Friedman, A. Nahor, S. Eichler, L. M. Bonanno, A. Sa'ar and E. Segal, Soft Matter, 2012, 8, 9166. B. Sanz, N. Ballard, J. M. Asua and C. Mijangos, *Macromolecules*, 2017, **50**, 811–821. H. Zhao and S. L. Simon. Polvmer (Guildf)., 2011, 52, 4093–4098. S. Alexandris, P. Papadopoulos, G. Sakellariou, M. Steinhart, H.-J. Butt and G. Floudas, Macromolecules, 2016, 49, 7400-7414. C. Mijangos, R. Hernández and J. Martín, Prog. Polym. Sci., 2016, 54–55, 148–182. Y. Suzuki, H. Duran, M. Steinhart, H.-J. Butt and G. Floudas, Soft Matter, 2013, 9, 2621. M. A. DePierro and C. A. Guymon, *Macromolecules*, 2014, **47**, 5728–5738. J. D. Clapper, L. Sievens-Figueroa and C. A. Guymon, *Chem. Mater.*, 2008, **20**, 768–781. M. A. DePierro and C. A. Guymon, Macromolecules, 2006, 39, 617-626. L. C. Lester and C. A. Guymon, *Macromolecules*, 2000, 33, 5448–5454. P. Alexandridis, U. Olsson and B. Lindman, J. Phys. Chem., 1996, 100, 280–288. P. Alexandridis, U. Olsson and B. Lindman, Langmuir, 1998, 14, 2627–2638. P. Holmgvist, P. Alexandridis and B. Lindman, *Langmuir*, 1997, **13**, 2471–2479. P. Alexandridis and B. Lindman, Amphiphilic Block Copolymers: Self-assembly and Application, Elsevier, 2000. D. T. McCormick, K. D. Stovall and C. A. Guymon, Macromolecules, 2003, 36, 6549-6558. C. L. Lester, S. M. Smith and A. Guymon, *Macromolecules*, 2001, **34**, 8587–8589. L. Silies, H. Didzoleit, C. Hess, B. Stü and A. Andrieu-Brunsen, Chem. Mater., 2015, 27, 1971-1981. M. Antonietti, C. Göltner and H.-P. Hentze, *Langmuir*, 1998, **14**, 2670–2676. H.-P. Hentze and E. W. Kaler, Curr. Opin. Colloid Interface Sci., 2003, 8, 164–178. R. Foudazi, P. Gokun, D. L. Feke, S. J. Rowan and I. Manas-Zloczower, Macromolecules, 2013, 46, 5393-5396. B. Strachota, L. Matějka, A. Sikora, J. Spěváček, R. Konefał, A. Zhigunov, M. Šlouf and L. Matějka, Soft Matter, 2017, 13, 1244–1256.

591 592	35	C. A., S. K., SM. J. F. A., W. Q., P. J. A. and MM. J. D., <i>RSC Adv.</i> , 2016, 6 , 81694– 81702.
593 594	36	M. Di Biase, P. de Leonardis, V. Castelletto, I. W. Hamley, B. Derby and N. Tirelli, <i>Soft Matter</i> , 2011, 7 , 4928.
595 596 597	37	S. Peng, P. G. Hartley, T. C. Hughes, Q. Guo, M. Y. Guo, S. B. Wei, S. L. Qiu, R. H. Templer, T. H. Nguyen, T. M. Hinton, L. J. Waddington, N. Kirby, D. K. Wright, H. X. Wang, G. E. Egan and B. A. Moffat, <i>Soft Matter</i> , 2015, 11 , 6318–6326.
598	38	A. W. Chow, J. F. Sandell and J. F. Wolfe, Polymer (Guildf)., 1988, 29, 1307–1312.
599	39	A. W. Chow, R. D. Hamlin, J. F. Sandell and J. F. Wolfe, MRS Proc., 1988, 134, 95.
600 601	40	S. Qavi, A. P. Lindsay, M. A. Firestone and R. Foudazi, <i>J. Memb. Sci.</i> , 2019, 580 , 125– 133.
602	41	G. Odian, Principles of Polymerization, John Wiley & Sons, 2004.
603	42	P. Alexandridis, U. Olsson and B. Lindman, <i>Macromolecules</i> , 1995, 28, 7700–7710.
604	43	S. Qavi, M. A. Firestone and R. Foudazi, Soft Matter, 2019, 15, 5626–5637.
605 606	44	B. A. Smith, E. T. Samulski, LP. Yu and M. A. Winnik, <i>Phys. Rev. Lett.</i> , 1984, 52 , 45–48.
607	45	A. K. Fritzsche and F. P. Price, <i>Polym. Eng. Sci.</i> , 1974, 14 , 401–412.
608	46	N. Ileri Ercan, P. Stroeve, J. W. Tringe and R. Faller, Langmuir, 2016, 32, 10026–10033.
609 610	47	K. Bryskhe, K. Schillen, JE. Lofroth and U. Olsson, <i>Phys. Chem. Chem. Phys.</i> , 2001, 3 , 1303–1309.
611 612	48	R. Foudazi, S. Qavi, I. Masalova and A. Y. A. Y. Malkin, <i>Adv. Colloid Interface Sci.</i> , 2015, 220 , 78–91.
613 614	49	K. Hyun, J. G. Nam, M. Wilhellm, K. H. Ahn and S. J. Lee, <i>Rheol. Acta</i> , 2006, 45 , 239– 249.
615	50	S. R. Raghavan and S. A. Khan, J. Colloid Interface Sci., 1997, 185, 57–67.
616 617	51	C. Bower, C. Gallegos, M. R. Mackley and J. M. Madiedo, <i>Rheol. Acta</i> , 1999, 38 , 145– 159.
618 619	52	J. Zipfel, J. Berghausen, G. Schmidt, P. Lindner, P. Alexandridis, M. Tsianou and W. Richtering, <i>Phys. Chem. Chem. Phys.</i> , 1999, 1 , 3905–3910.
620	53	JP. Habas, E. Pavie, A. Lapp and J. Peyrelasse, <i>Rheol. Acta</i> , 2008, 47 , 765–776.
621	54	J. N. Israelachvili, Intermolecular and Surface Forces, Academic Press, third edit., 1992.
622	55	R. Foudazi, I. Masalova and A. Y. Malkin, <i>J. Rheol. (N. Y. N. Y</i>)., 2012, 56 , 1299–1314.
623	56	T. Mason, J. Bibette and D. Weitz, <i>Phys. Rev. Lett.</i> , 1995, 75 , 2051–2054.
624	57	R. J. Ketz, R. K. Prud'homme and W. W. Graessley, <i>Rheol. Acta</i> , 1988, 27 , 531–539.
625 626	58	M. Laurati, G. Petekidis, N. Koumakis, F. Cardinaux, A. B. Schofield, J. M. Brader, M. Fuchs and S. U. Egelhaaf, <i>J. Chem. Phys.</i> , 2009, 130 , 134907(1-14).

- 627 59 P. Sollich, *Phys. Rev. E*, 1998, **58**, 738–759.
- 628 60 S. Qavi and R. Foudazi, *Rheol. Acta*, 2019, **58**, 483–498.
- 629 61 P. J. Flory, *Principles of Polymer Chemistry*, Cornell University Press, 1953.
- 630 62 H. H. Winter and F. Chambon, *J. Rheol. (N. Y. N. Y).*, 1986, **30**, 367–382.
- 63 A. Y. Malkin and S. G. Kulichikhin, *Polymer Compositions Stabilizers/Curing SE-5*,
 63 Springer-Verlag, Berlin/Heidelberg, 1991, vol. 101.
- 633 64 P. J. Halley and M. E. Mackay, *Polym. Eng. Sci.*, 1996, **36**, 593–609.
- 634 65 I. Koltover, T. Salditt, J. O. Rädler and C. R. Safinya, *Sci.*, 1998, **281**, 78–81.
- 635 66 V. Allain, C. Bourgaux and P. Couvreur, *Nucleic Acids Res.*, 2012, **40**, 1891–1903.
- 636 67 Y. Li, O. S. Sarıyer, A. Ramachandran, S. Panyukov, M. Rubinstein and E. Kumacheva,
 637 Sci. Rep., 2015, 5, 17017.
- 68 B. Yu, P. Sun, T. Chen, Q. Jin, D. Ding, B. Li and A.-C. Shi, *Phys. Rev. Lett.*, 2006, 96, 138306.
- 640 69 C. Autran, J. C. de la Cal and J. M. Asua, *Macromolecules*, 2007, **40**, 6233–6238.
- T. G. T. Jansen, J. Meuldijk, P. A. Lovell and A. M. van Herk, *J. Polym. Sci. Part A Polym. Chem.*, 2016, **54**, 2731–2745.
- 643 71 M. Salsamendi, N. Ballard, B. Sanz, J. M. Asua and C. Mijangos, *RSC Adv.*, 2015, 5,
 644 19220–19228.
- 645 72 J. Dvorkin, R. Nolen-Hoeksema and A. Nur, *Geophysics*, 1994, **59**, 428–438.
- 646 73 E. Detournay and A. H. Cheng, 1993, II, 113–171.
- 647 74 D. S. Achilias, *Macromol. theory simulations*, 2007, **16**, 319–347.
- 648 75 P. Siafaka and D. S. Achilias, *Macromol. Symp.*, 2013, **331–332**, 166–172.
- 649 76 D. S. Achilias, J. Therm. Anal. Calorim., 2014, **116**, 1379–1386.

