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Abstract

The metallo–formate anions, $M(CO_2)^-$, $M = Ni$, Pd, and Pt, were formed by electron– induced $CO₂$ activation. They were generated by laser vaporization and characterized by a combination of mass spectrometry, anion photoelectron spectroscopy, and theoretical calculations. While neutral transition metal atoms are normally unable to activate $CO₂$, the addition of an excess electron to these systems led to the formation of chemisorbed anionic complexes. These are covalently bound, formate–like anions, in which their CO₂ moieties are significantly reduced. In addition, we also found evidence for an unexpectedly attractive interaction between neutral Pd atoms and $CO₂$.

Introduction

The activation of carbon dioxide underpins its chemistry. Since the carbon atom in $CO₂$ is in its highest oxidation state, the activation of $CO₂$ inevitably involves reducing it, and that implies $CO₂$ accepting some degree of negative charge. Accomplishing this, however, requires at least a partial bending of CO_2 . The CO_2^- anion that results from CO_2 having accepted a full negative charge is metastable; the electron affinity of CO₂ is -0.6 eV. While some studies have dealt with free carbon dioxide anions, most have focused on anionic complexes consisting of $CO₂$ and various other atoms and molecules.^{1–43} Anion photoelectron studies of N–heterocycle–CO₂ heterogeneous anionic dimers by Kim *et al* showed significant covalent character in their intermolecular bond.6,7 Subsequent work by Johnson *et al*⁸ and by ourselves⁹ added additional dimensions to this topic. In all cases, however, the $CO₂$ moieties were found to be partially negatively–charged and bent. Infrared photodissociation studies of transition metal–CO₂ anionic complexes by Weber *et al* further explored this topic.^{10–23} Both electrostatically–bound, metal atom–multiple $CO₂$ anionic complexes (physisorption) and covalently–bound, metal atom–multiple $CO₂$ anionic complexes (chemisorption) were found. In the latter cases, the $CO₂$ moieties were partially bent and had accepted some significant portion of the negative charge, while in the former cases, the $CO₂$ were only very slightly bent, suggesting insignificant $CO₂$ activation. Calculations implied that the metal atoms were far away from the $CO₂$ moieties in the neutral complexes, and that their $CO₂$ moieties were structurally identical to isolated $CO₂$ molecules. Subsequent work in our laboratory measured the anion photoelectron spectra of copper–, silver–, and gold– $CO₂$ anionic dimers, finding only physisorption in the case of silver, i.e, $Ag(CO₂)$, only chemisorption in the case of copper, i.e., $Cu(CO_2)$ ⁻, and both physisorbed and chemisorbed isomers for gold, i.e., $Au(CO_2)$ and $Au(CO_2)$ ⁻ , respectively.²⁴

Here, we present our study of the Group 10 transition metal– $CO₂$ anionic complexes: [Ni(CO₂)]⁻, [Pd(CO₂)]⁻, and [Pt(CO₂)]⁻, using anion photoelectron spectroscopy and theoretical calculations. In contrast to our previous work with the Group 11 coinage $(s¹)$ metals, the Group 10 metals adopt richer outer electron shell configurations ($d^{8}s^{2}$ for Ni, d^{10} for Pd, and $d^{9}s^{1}$ for Pt) with significantly greater prospects for complex chemical bonding. Indeed, strong evidence for chemisorption was found in all three of the Group 10 metal– $CO₂$ anionic complexes studied here, implying that Ni⁻, Pd⁻ and Pt⁻ all activated CO₂ to form Ni(CO₂)⁻, Pd(CO₂)⁻, and Pt(CO₂)⁻, respectively. Thus, all three of these Group 10 metal anions were seen to be able to both reduce and activate $CO₂$.

Methods

Experimental

Anion photoelectron spectroscopy is conducted by crossing a beam of mass–selected negative ions with a fixed–frequency photon beam and energy–analyzing the resultant photodetached electrons. The photodetachment process is governed by the energy–conserving relationship: $hv = EBE + EKE$, where *hv* is the photon energy, EBE is the electron binding (photodetachment transition) energy, and EKE is the electron kinetic energy. Our apparatus consists of a laser vaporization cluster anion source, a time–of–flight mass spectrometer, a Nd:YAG photodetachment laser, and a magnetic bottle electron energy analyzer⁴⁴. The photoelectron spectrometer resolution is \sim 35 meV at EKE = 1 eV. The third harmonic output of a Nd:YAG laser (355 nm) was used to photodetach electrons from mass–selected M– and $[M(CO₂)]$ ⁻ anions, where M = Ni, Pd, Pt. Photoelectron spectra were calibrated against the wellknown atomic transitions of atomic Cu⁻⁴⁵ The $[M(CO₂)]$ ⁻ (M = Ni, Pd, Pt) anion complexes were generated in a laser vaporization ion source. It consisted of rotating, translating nickel, palladium, or platinum rods, which were ablated with second harmonic (532 nm) photon pulses from a Nd:YAG laser, while $10\%/90\%$ He/CO₂ gas mixtures at 60 psi were expanded from a pulsed valve over the rods.

Theoretical

Density functional theory calculations were performed with the ORCA computational chemistry software package.⁴⁶ All calculations were carried out with the B3LYP functional⁴⁷ with the D3 dispersion correction⁴⁸ and the RIJCOSX approximation⁴⁹. The Ahlrichs Def2 basis sets were used throughout our calculations⁵⁰. For geometry optimization, Def2–TZVP and auxiliary Def2–TZVP/J basis sets⁵¹ were chosen for carbon, oxygen and nickel atoms; the Stuttgart effective core potential SDD⁵² and ECP basis set Def2–TZVP|Def2–TZVP/J were used for palladium and platinum atoms. The potential energy surfaces of neutral Ni, Pt - $CO₂$ along the M–C coordinate were computed by scanning the M–C bond length with a step width of 0.1 Å, while relaxing the rest of the cluster. Single–point calculations were then improved with Def2–QZVPP|Def2– QZVPP/J basis sets and all–electron relativistic calculations (ZORA). The structure of neutral $PdCO₂$ was also checked using the PBE0 and M06-L functionals.⁵³ The vertical detachment energy (VDE) is the energy difference between the ground state anion and its corresponding neutral at the geometry of the anion, i.e., these are vertical photodetachment transitions. The adiabatic detachment energy (ADE) is the energy difference between the lowest energy, relaxed geometry of the anion and the relaxed geometry of a structurally similar isomer (nearest minimum) of its neutral counterpart. The adiabatic electron affinity (EA) is the energy difference between the lowest energy, relaxed geometry of the anion and the relaxed geometry of the lowest energy isomer (the global minimum) of its neutral counterpart. When the nearest local minimum and the global minimum are one and the same, ADE = EA. In this work, we calculated ADE values. In the systems

studied here, there is only one credible minimum for the neutral species, and it is therefore the global minimum. For that reason, we report ADE values as EA values here. Also, note that since the Franck–Condon principle governs which spectral features are seen within its anion–to–neutral wavefunction overlap window, there is often a correspondence between the lowest EBE transition observed experimentally and the properly calculated ADE value. Franck-Condon simulation was performed for the PdCO₂⁻ spectrum. This simulation, however, was not practical for NiCO₂⁻ and $PtCO₂$ due to the large structural changes between anions and neutrals. Frequency calculations were performed to verify that no imaginary frequencies existed for any of the optimized structures.

Results and Discussion

The photoelectron spectra of $[M(CO₂)]$ ⁻ (M = Ni, Pd, Pt) are presented in Figure 1. For comparison, the atomic anion photoelectron spectra of Ni– , Pd– , and Pt– are also presented above each $[M(CO₂)]$ ⁻ spectrum there. In the anionic complexes, $[M(CO₂)]$ ⁻, the $CO₂$ moiety can be either physisorbed or chemisorbed to M⁻. For physisorbed complexes, the interaction between M⁻ and CO_2 is weak; they can be considered to be M⁻ anions "solvated" by CO_2 molecules, i.e., M⁻ $(CO₂)$. In physisorbed anion–molecule complexes such as these, their M⁻ anion moieties act as chromophores for photodetachment, with the resulting photoelectron spectra closely resembling the photoelectron spectral patterns of M– , just shifted to higher electron binding energies (EBE) with their spectral features slightly broadened.^{54,55} These shifts are typically a few tenths of an eV, corresponding approximately to ion–solvation stabilization energies. This spectral behavior provides a distinctive spectroscopic signature for physisorbed (solvated) anion complexes. The photoelectron spectra of $[Ni(CO₂)]^-$, $[Pd(CO₂)]^-$, and $[Pt(CO₂)]^-$, however, do not exhibit this behavior in relation to their atomic anion photoelectron spectra, i.e., Ni-, Pd-, and Pt-. They are not physisorbed species. Moreover, their spectral shifts are far too large to be due to weak, solvation–like interactions. The spectral shifts between the lowest EBE features of the atomic anions and those of their corresponding anionic complexes are about an electron volt. The anionic complexes studied here, i.e., $[Ni(CO_2)]^-$, $[Pd(CO_2)]^-$, and $[Pt(CO_2)]^-$, are chemisorbed complexes, i.e., NiCO₂⁻, PdCO₂⁻, and PtCO₂⁻. Significant chemical interactions have occurred, resulting in bonds between the metal atoms and their $CO₂$ moieties.

Figure 1. Anion photoelectron spectra of Ni– and NiCO₂⁻, Pd⁻ and PdCO₂⁻, and Pt and PtCO₂⁻. The vertical lines in the $PdCO₂$ spectrum represent Franck-Condon simulated vibrational progression.

Energetic parameters can also be determined from the spectra. When there is sufficient Franck–Condon overlap between the ground state of the anion and the ground state of the neutral, and when vibrational hot bands are absent, the threshold EBE (E_T) is the value of the electron affinity (EA). The E_T in PdCO₂⁻ spectrum can be determined definitively. For NiCO₂⁻ and PdCO₂⁻, however, the relatively weak, shelf–like features on the low EBE side of the major peaks are likely hot bands. Thus, the E_T values for those two systems were determined by extrapolating the low

EBE sides of their major peaks to baseline. In all three anionic complexes, the EBE values of the intensity maxima in their major peaks are their vertical detachment energy (VDE) values, corresponding to the transitions that have maximum Franck–Condon overlap between the ground electronic states of the anionic complexes and their neutral counterparts. The onset of the higher EBE feature in $PdCO_2$ spectrum is due to the vertical photodetachment transition from the ground state anion to its neutral counterpart in its first excited electronic state. The EA and VDE values of NiCO₂^{0/-}, PdCO₂^{0/-}, and PtCO₂^{0/-} are listed in Table 1.

Table 1. Experimental and theoretical EA and VDE values for MCO_2 and MCO_2^- , $M = Ni$, Pd, Pt, respectively. All values are in eV.

$NiCO20/-$ 2.33/2.31 1.9/1.86	
$PdCO20/-$ 1.3/1.10 1.60/1.57	
$PtCO2$ ^{0/-} 3.0/2.81 3.43/3.37	

A synergy between theoretical calculations and anion photoelectron spectroscopy can provide insights into the structures, energetics and the nature of chemical bond of the investigated clusters.55-61 The calculated structures of the anionic complexes and their corresponding neutrals are presented in Figure 2. The first row shows the geomtries of the anionic complexes along with their respective HOMOs, while the second row provides the geometries of their neutral counterparts. The M–C bond length (\AA) , the C–O bond length (\AA) , and the O–C–O bond angle (in degrees) are shown for each case. For the chemisorption species, $NiCO_2^-$, $PdCO_2^-$, and $PtCO_2^-$, the M–C bond lengths are 2.01 Å, 2.03 Å, and 2.02 Å, respectively, suggesting the formation of single bonds between M and C in all cases. The $CO₂$ moiety is significantly bent in all three anionic. complexes, with an O–C–O bond angle of 138.26° for NiCO₂⁻, 140.62° for PdCO₂⁻ and 136.34^o

Figure 2. Calculated structures of $NiCO_2^-$, $PdCO_2^-$, $PtCO_2^-$ (first row) and $NiCO_2$, $PdCO_2$, $PtCO_2$ (second row). The HOMOs of the anionic complexes are also presented.

for PtCO₂⁻, respectively. A natural population analysis shows that the CO₂ moieties in NiCO₂⁻, PdCO₂⁻, and PtCO₂⁻ have negative charges of $-0.75 e$, $-0.59 e$, and $-0.67 e$, respectively. Thus, the CO2 moiety has been significantly reduced in all of these anionic complexes. The O–C–O bond angles in all three systems are similar one another, as are the negative charges on their $CO₂$ moieties. The fact that they are not completely synchronized is likely due to the natural population analysis (NPA) being less reliable than the structural calculations. For the three anionic complexes, the C–O bond lengths are all between 1.22–1.23 Å, which is longer than the 1.16 Å C–O bond

length in isolated CO_2 (1.16 Å). This implies that when negative charge is transferred to the CO_2 moiety, the C–O bond is elongated and weakened. Thus, $NiCO_2^-$, $PdCO_2^-$, and $PtCO_2^-$ all share a metallo–formate geometry. Furthermore, they are structurally quite similar in terms of M–C bond length, C–O bond length and O–C–O bond angle. Note that based on the calculated structures and charge distributions, Pd seems to have the weakest interaction with $CO₂$ of all three metal anions, although Pd has the lowest electron affinity which is expected to facilitate the charge transfer to the CO_2 moiety. One possible reason for Pd⁻ being the outlier is its electron configuration, $d_{10}s$, which is different to that of Ni^- and Pt^- , d_9s_2 . The calculated EA and VDE values are listed in Table 1 along with their corresponding experimental values. Excellent agreement between experimental and theoretical values, is seen for all three anionic complexes, validating the geometry optimizations shown in Figure 2.

In neutral NiCO₂ and PtCO₂, the metal atom is far away from CO_2 . Also, their CO_2 moieties are structurally identical to an isolated $CO₂$ molecule. The potential energy surfaces of neutral $NiCO₂$ and PtCO₂ seem to show a shallow energy well at a M-C bond length of around 2.0 Å, but the energies of these local minima are higher than when $CO₂$ is far away (Figure 3). The repulsive part of the neutral surfaces occurs at a M-C bond length less than 1.9 Å. Since NiCO₂⁻ and PtCO₂⁻ have a M-C bond length of 2.01 and 2.02 Å, respectively, the repulsive part of each neutral surface is not accessed during the vertical photodetachment process. The structural parameters show that there is little interaction between the neutral atoms of Ni and Pt and $CO₂$, which is as expected based on our previous research.²⁴ Surprisingly, however, the optimized neutral PdCO₂ structure shows incipient chemisorption character, that while much weaker than in its $PdCO_2^-$ anionic counterpart, is significantly stronger than the physisorption interactions seen in neutral $NiCO₂$ and PtCO₂. In neutral PdCO₂, the Pd–C bond length is 2.31 Å, which is characteristic of a metal–

carbon bond. The CO_2 moiety is noticeably bent, with the O-C-O bond angle being 163.42°. The NPA analysis shows the CO_2 moiety as possessing a negative charge of -0.16 e, indicating a degree of charge transfer between neutral Pd and $CO₂$. The neutral PdCO₂ structure was also verified using the PBE0 and M06-L functionals. Both functionals yielded the same $PdCO₂$ structure as the B3LYP one. In addition, the Franck-Condon simulated spectrum of $PdCO_2$ ⁻ reproduces the experimental one (Figure 1), offering a further validation of the neutral $PdCO₂$ structure. Taken together, these features suggest that $CO₂$ can be activated by a single neutral Pd atom.

Figure 3. The potential energy surfaces of neutral NiCO₂ and PtCO₂ with respect to the M-C bond length. The red crosses represent the M-C bond length of optimized anionic $NiCO_2^-$ and $PtCO_2^-$.

In order to provide further insight into the nature of the bonding in neutral $PdCO₂$, we analyzed its molecular orbitals. In most cases, the molecular orbitals of neutral $M(CO_2)$ species are composed of atomic orbitals of M and molecular orbitals of $CO₂$, which are essentially independent of one another. The molecular orbitals of neutral $PtCO₂$ (HOMO to HOMO–8) are

presented on the left side of Figure 4, and they provide typical examples of the molecular orbitals in neutral MCO_2 species. The absence of wavefunction overlap between the M and CO_2 moieties is consistent with a lack of interaction between neutral M atoms and $CO₂$ moieties. In the case of

Figure 4. Selected molecular orbitals of neutral (a) $PtCO₂$ and (b) $PdCO₂$. The overlap between the metal atomic orbitals and the $CO₂$ molecular orbitals is barely present in neutral PtCO₂, but is significant in neutral $PdCO₂$.

neutral PdCO₂, most of its molecular orbitals can also be viewed as independent atomic orbitals of Pd and molecular orbitals of $CO₂$. However, this is not the case for all orbital interactions between Pd and CO₂. The HOMO–4 and HOMO–6 orbitals, seen on the right side of Figure 3, clearly result from the combination of Pd atomic orbitals and $CO₂$ molecular orbitals. There, one observes significant interaction, i.e., overlap, between some of the orbitals of Pd and $CO₂$. For

HOMO–4, the interaction is mainly between the Pd d_z^2 orbital and the B_l orbital of bent CO₂. For HOMO–6, the interaction is through the overlap between Pd d_{xy} orbital and the B_2 orbital of bent $CO₂$. The wavefunction overlap between these Pd atomic and $CO₂$ molecular orbitals is likely the reason for the weak binding interaction and the partial charge transfer between neutral Pd atom and CO2. Although we are unaware of corroborating experimental evidence for their weak binding, our calculations imply a neutral $PdCO₂$ bond dissociation energy of 0.85 eV. For comparison, the bond dissociation energy of the $PdCO_2^-$ anionic complex is predicted to be 1.78 eV.

Conclusion

This combined anion photoelectron spectroscopic and computational study characterized the metallo–formate anions, $M(CO_2)$ ⁻, where $M = Ni$, Pd, and Pt, and demonstrated that the addition of an excess electron led to significant $CO₂$ reduction in these systems. While neutral transition metals are normally incapable of reducing $CO₂$, we found evidence of an unusual attractive interaction between neutral Pd and $CO₂$.

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Notes

The authors declare no competing financial interest.

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