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# Stronger-than-Pt hydrogen adsorption in a Au<sub>22</sub> nanocluster for hydrogen evolution reaction

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Abstract: Atomically-precise metal nanocluster have recently emerged as a novel class of catalysts for hydrogen evolution reaction. From first-principles density functional theory, we show that the eight coordinatively unsaturated (*cus*) Au atoms in the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster [L = 1,8-bis(diphenylphosphino) octane] can adsorb H stronger than Pt, thereby being a potentially promising catalyst for hydrogen evolution reaction (HER). We find that up to six H atoms can adsorb on the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster and they all have close-to-zero Gibbs free adsorption energies ( $\Delta G_H$ ). From the HOMO-LUMO gaps, frontier orbitals, and Bader charge analysis, we conclude that H behaves as a hydride or electron-withdrawing ligand in the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> clusters, in contrast with the metallic H in thiolate-protected Au nanoclusters. Our study demonstrates that ligand-protected Au clusters with *cus* Au sites will be the most promising candidates for realizing Au-H nanoclusters and as excellent electrocatalysts for HER.

# Introduction

The behavior of H in metals has been of great interest for many decades.<sup>1-3</sup> Metal-hydrogen systems are widely utilized in energy-storage systems, sensor applications, as well as catalysis.<sup>4</sup> In contrast with bulk metal hydrides, the structures of H-metal nanoclusters can provide insight from a molecular perspective of how H interacts with metals and how the magic sizes or superatoms evolve.<sup>5</sup> The hydride ligand has been used in the synthesis of ligand-protected Cu or Ag nanoclusters.<sup>6-13</sup> Especially for the Cu clusters, a series of air- and moisture-stable Cu hydride clusters ranging from Cu<sub>7</sub>H to Cu<sub>32</sub>H<sub>20</sub> have been successfully synthesized and the structures have been determined.<sup>14</sup> More recently, three new hydride-rich Ag clusters, [Ag<sub>18</sub>H<sub>16</sub>(TPP)<sub>10</sub>]<sup>2+</sup>, [Ag<sub>25</sub>H<sub>22</sub>(DPPE)<sub>8</sub>]<sup>3+</sup>, and [Ag<sub>26</sub>H<sub>22</sub>(TFPP)<sub>13</sub>]<sup>2+</sup>, have been synthesized.<sup>15</sup>

Atomically precise Au nanoclusters have attracted intensive research interest over the past two decades due to their potential applications in catalysis, biology, and nanotechnology.<sup>16-</sup> <sup>18</sup> However, to the best of our knowledge, H has not been observed in ligand-protected Au nanoclusters, but some intermediate states of H adsorbed in thiolate-protected Au clusters have been proposed to explain their roles in the hydrogen evolution reaction (HER).<sup>19</sup> It has been shown that both Au<sub>25</sub>(SR)<sub>18</sub> and PtAu<sub>24</sub>(SR)<sub>18</sub> exhibit high HER activity and the activity of  $PtAu_{24}(SR)_{18}$  is much higher than that of  $Au_{25}(SR)_{18}$ . In fact,  $PtAu_{24}(SR)_{18}$  is among the most active molecular catalysts for HER. To understand the underlying reason, the behavior of H in  $[Au_{25}(SR)_{18}]^{q}$ mono-atom-doped bimetallic  $[M_1Au_{24}(SR)_{18}]^q$  clusters and including PtAu<sub>24</sub>(SR)<sub>18</sub> has been investigated.<sup>20</sup> It was found that H behaves as a metal in these clusters and contributes its 1s electron to the superatomic free-electron count. But H adsorption is still uphill for at least 200 meV in free energy on these Au clusters, indicating their thermodynamic

instability. One strategy to drive down the free energy is to create the coordinatively unsaturated (*cus*) or *in situ* uncoordinated sites in the Au clusters.

Unlike many other clusters where the surface Au atoms are fully coordinated by the protecting ligands,<sup>21-24</sup> the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster (L = 1,8-bis(diphenylphosphino) octane) is unique in that it contains eight *cus* or so-called in situ uncoordinated Au atoms.<sup>25</sup> As shown in Fig. 1, the rod-like Au<sub>22</sub> core of the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster consists of two Au<sub>11</sub> units clipped together by four L<sup>8</sup> ligands (ball-and-stick ligands in Fig. 1a), while the additional two ligands each coordinate to a Au<sub>11</sub> unit in a bidentate fashion. Significantly, eight Au atoms (highlighted in red in Fig. 1b) at the interface of the two Au<sub>11</sub> units are not coordinated by any ligands. The presence of the *cus* Au atoms is promising for catalysis. For example, it was found that the intact Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> nanoclusters readily catalyze CO oxidation at room temperature; DFT modeling of adsorption of CO and O<sub>2</sub> suggested that the *cus* Au atoms are the active sites.<sup>26</sup>



Fig. 1 Total structure (a) and core structure (b) of the  $Au_{22}(L^8)_6$  cluster. *Cus* Au, red; other Au, yellow; P, magenta; C, grey; H, light grey.

Due to its many *cus* Au sites, we hypothesize that  $Au_{22}(L^8)_6$  offers a unique opportunity to observe H in Au nanoclusters and can be an excellent electrocatalyst for HER. To test this hypothesis, here we studied the structures and energetics of H adsorption on the  $Au_{22}(L^8)_6$  cluster by first-principles density functional theory (DFT) and evaluated the Gibbs free energies of H adsorption ( $\Delta G_H$ ) to predict the HER activity. In addition, the HOMO-LUMO gaps, frontier orbitals, and partial atomic charges of the H-adsorbed Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> clusters were analyzed to better understand the behavior of H in the cluster.



Fig. 2 Differential adsorption energy  $\Delta E_{\rm H}$  (blue) and adsorption free energy  $\Delta G_{\rm H}$  (red) as a function of number of hydrogen adsorbed on the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster.

## **Results and discussion**

Energetics and structures of H adsorption on the  $Au_{22}(L^8)_6$  cluster. To find the optimal adsorption sites of H, we tested all the possible sites on the  $Au_{22}$  cluster, including top, verticalbridge, horizontal-bridge, and hollow sites on the  $Au_8$  cube at the center of the cluster (highlighted in red in Fig. 1) as well as the top, bridge, and hollow sites on the two  $Au_{11}$  units (highlighted in yellow in Fig. 1). The H atoms first adsorb at the lowest energy sites, and then fill the next lowest and so on.

Fig. 2 shows the calculated differential H adsorption energy ( $\Delta E_{H}$ ; blue) and free energy ( $\Delta G_{H}$ ; red) on the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster relative to  $\frac{1}{2}$  H<sub>2</sub>. One can see that up to six H atoms can be

favorably adsorbed on the cluster, while the second H binds strongest.  $\Delta G_{H}$  has been considered as a good descriptor for HER and a close-to-zero  $\Delta G_{H}$  would suggest a good HER catalyst.<sup>27</sup> One can see that  $\Delta G_{H}$  on the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster is close-to-zero for the first six H atoms, except the second one. This result has two important implications: (i) the two-H state or Au<sub>22</sub>H<sub>2</sub>(L<sup>8</sup>)<sub>6</sub> cluster is highly favorable and probably can be detected experimentally; (ii) the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster should be a highly active HER catalyst since it offers so many sites that have close-to-zero  $\Delta G_{H}$ . For comparison, we have also calculated  $\Delta G_{H}$  on the Pt(111) surface (at the same level of theory and with the same approach) and found it to be -0.14 eV. Hence, the second H adsorption on the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster ( $\Delta G_{H} = -0.41$  eV) is much stronger than H adsorption on the Pt(111) surface, indicating the high stability of the Au<sub>22</sub>H<sub>2</sub>(L<sup>8</sup>)<sub>6</sub> cluster.

H adsorption on the  $Au_{22}(L^8)_6$  cluster shown in Fig. 2 is drastically different from previously studied  $[Au_{25}(SR)_{18}]^q$  and mono-atom-doped bimetallic  $[M_1Au_{24}(SR)_{18}]^q$  clusters.<sup>20</sup> For the thiolate-protected Au clusters, all the surface gold atoms are protected by thiolate ligands already, so there is no *cus* Au site. As a result, the bonding of H in thiolate-protected Au clusters is generally weaker ( $\Delta G_H > 0.200 \text{ eV}$ ) than in the  $Au_{22}(L^8)_6$  cluster ( $\Delta G_H \sim -0.40$  to 0.10 eV).

To understand the energetic pattern of H adsorption on the  $Au_{22}(L^8)_6$  cluster, we now examine the optimized structures. One can see from Fig. 3 that the first H atom adsorbs at the vertical-bridge site formed by two *cus* Au atoms, while the second H atom also adsorbs at the vertical-bridge site on the opposite side of the Au<sub>8</sub> cube. Then the third and the fourth H atoms occupy the other two vertical-bridge sites subsequently. Interestingly, the fifth and the six H atoms adsorb at the bridge sites at the opposite ends of the cluster where the Au atoms are not *cus* sites but protected by phosphine ligands. If the first H atom directly adsorbs at the non-*cus*-Au bridge site,  $\Delta G_H$  is 0.87 eV. In other words, H adsorption at the *cus* Au sites can activate the non-*cus* Au sites; the interaction between H atoms and non-*cus* Au atoms becomes stronger after four H adsorb at the *cus* Au sites. What's more intriguing is that each cap with their non-*cus* Au atoms can accommodate only one H atom, because adsorption of the seventh H atom at the non*cus* Au bridge site is highly unfavorable (Fig. 2). These observations strongly indicate some magic-number nature with compositions of  $Au_{22}H_2(L^8)_6$  and  $Au_{22}H_6(L^8)_6$  which should correlate with their electronic structure.



Fig. 3 Optimized structures of the  $Au_{22}H_n(L^8)_6$  clusters, n=1~6. Protecting ligands are shown in line mode. *Cus* Au, red; other Au, yellow; H at the *cus* Au sites, blue; H at the non-*cus* Au sites, green; other H, light grey; P, magenta; C, grey.



Fig. 4 Minimum-energy path of  $H_2$  dissociation on the  $Au_{22}(L^8)_6$  cluster. Structures for initial, transition, and final states are also shown. The protecting ligands are eliminated for clarity. *Cus* Au, red; other Au, yellow; H, blue.

Another interesting implication from the highly stable  $Au_{22}H_2(L^8)_6$  cluster is that  $Au_{22}(L^8)_6$  can be a good catalyst for H<sub>2</sub> activation. To test this idea, we located the minimum-energy path for H<sub>2</sub> dissociation on the  $Au_{22}(L^8)_6$  cluster (Fig. 4) and the barrier was found to be 0.76 eV with a reaction energy of -0.43 eV. One can see that the H<sub>2</sub> molecule dissociates at the short-edge site of the *cus* Au<sub>8</sub> center. After dissociation, H atoms adsorb at two adjacent long-bridge sites. We further calculated the barrier of H diffusion from a long-bridge site to the next to form the configuration as in Fig. 3 for  $Au_{22}H_2(L^8)_6$ . This barrier was found to be 0.54 eV ( $\Delta E = -0.19$  eV). These low barriers indicate that the  $Au_{22}(L^8)_6$  cluster should be able to activate H<sub>2</sub> at room temperature.

**Partial atomic charges of H in the Au\_{22}(L^8)\_6 clusters.** To understand the electronic structure of these clusters, we need to determine their free-electron count. Here a key question is whether H behaves as a hydride or a metal. To solve this issue, we performed Bader charge analysis. We

found that charges on H in the phosphine-protected  $Au_{22}(L^8)_6$  cluster are all negative (ranging from -0.17 to -0.10 |e|), like a hydride. In contract, H in thiolate-protected  $Au_{25}$  behaves as a metal with close-to-zero charges.<sup>20</sup> So the behavior of H in Au nanoclusters is ligand-dependent. Due to the electron-donating phosphines, charges on Au are negative (ranging from -0.23 to -0.11 |e|) for the  $Au_{22}(L^8)_6$  cluster. Fig. 5 shows the isosurface plot of charge density difference for the  $Au_{22}(L^8)_6$  cluster with two adsorbed H. As one can see, electrons accumulate on the adsorbed H, while deplete on the Au atoms bonded to H. Our results of the charge analysis clearly show that H in the phosphine-protected Au nanoclusters behaves as a hydride, different from the metallic H in the thiolate-protected Au nanoclusters.



Fig. 5 Charge density difference for the  $Au_{22}(L^8)_6$  cluster with two adsorbed H at the contour level of 0.003 eV/Å<sup>3</sup>. The charge accumulation region is in light blue and the charge depletion region is in magenta. The protecting ligands are omitted for clarity. *Cus* Au, red; other Au, yellow; H, blue.

**HOMO-LUMO gaps and frontier orbitals of the H-adsorbed**  $Au_{22}(L^8)_6$  clusters. Since H behaves as a hydride, it can localize an electron from the core. By the electron-counting rule from the superatom complex model,<sup>5</sup> the free-electrons of the  $Au_{22}H_n(L^8)_6$  clusters (n=0, 2, 4, 6) are shown in Table 1 together with the calculated HOMO-LUMO gaps. Compared with  $Au_{22}(L^8)_6$  and  $Au_{22}H_2(L^8)_6$ ,  $Au_{22}H_4(L^8)_6$  has a smaller gap. This explained why the second H rather than the fourth H gave the strongest H adsorption (Fig. 2). The  $Au_{22}(L^8)_6$  cluster has 22 electrons, which does not correspond to a spherical shell closing. Instead, one can consider that

each of two Au<sub>11</sub> units contributes three electrons for the bonding between them. The HOMO of Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> shows significant distribution at the interface between the two Au<sub>11</sub> units (Fig. 6). Au<sub>22</sub>H<sub>2</sub>(L<sup>8</sup>)<sub>6</sub> has a more delocalized HOMO, but still with large densities at the interface. However, in the case of Au<sub>22</sub>H<sub>4</sub>(L<sup>8</sup>)<sub>6</sub>, the Au<sub>11</sub> units do not contribute electrons to the interface. It can be viewed as a "8e +10e" system instead of "9e+9e", because the triplet state is 0.17 eV higher than the singlet state. Accordingly, the HOMO of Au<sub>22</sub>H<sub>4</sub>(L<sup>8</sup>)<sub>6</sub> is mainly localized on the two Au<sub>11</sub> units, in the absence of any interactions between them. In addition, the HOMO diagrams also indicate the electron-rich sites that would attract (additional) H atoms.

Table 1 Free-electron counts (N) and HOMO-LUMO gaps of the  $Au_{22}H_n(L^8)_6$  clusters, n=0, 2, 4, 6.

Clusters	$\mathrm{Au}_{22}(\mathrm{L}^8)_6$	$\mathrm{Au}_{22}\mathrm{H}_2(\mathrm{L}^8)_6$	$\mathrm{Au}_{22}\mathrm{H}_4(\mathrm{L}^8)_6$	$\mathrm{Au}_{22}\mathrm{H}_6(\mathrm{L}^8)_6$
N	22	20	18	16
Gap /eV	0.716	0.667	0.520	1.185



Fig. 6 The HOMO diagrams of  $Au_{22}(L^8)_6$ ,  $Au_{22}H_2(L^8)_6$ , and  $Au_{22}H_4(L^8)_6$  clusters.

The Au<sub>22</sub>H<sub>6</sub>(L<sup>8</sup>)<sub>6</sub> cluster has a much larger gap (1.185 eV). Moreover, the nondegenerate, doubly occupied HOMO of Au<sub>22</sub>H<sub>6</sub>(L<sup>8</sup>)<sub>6</sub> resembles two spatially separated P-type orbitals, while the LUMO resembles two D-type orbitals (Fig. 7). This suggests that Au<sub>22</sub>H<sub>6</sub>(L<sup>8</sup>)<sub>6</sub> can be viewed as a "8e+8e" superatom for its 16 free electrons. We also checked the frontier orbitals of another 16-electron system,  $[Au_{22}H_4(L^8)_6]^{2+}$ . We found that its frontier orbitals have very similar features as Au<sub>22</sub>H<sub>6</sub>(L<sup>8</sup>)<sub>6</sub> and it also has a very large gap of 1.399 eV. This "8e+8e" superatomic shell-closing is the main reason that the Au<sub>22</sub>(L<sup>8</sup>)<sub>6</sub> cluster can absorb up to six H atoms to form the Au<sub>22</sub>H<sub>6</sub>(L<sup>8</sup>)<sub>6</sub> cluster, even though the last two H atoms have to go to the phosphine-protected gold atoms where the HOMO has more distribution (Fig. 6).



Fig. 7 Frontier orbitals of the  $Au_{22}H_6(L^8)_6$  cluster. Only one HOMO and one LUMO are shown.

It is interesting to investigate the difference in the orbitals below the HOMO, especially in the region of the H atoms, for H-adsorbed  $Au_{25}(SR)_{18}$  and  $Au_{22}(L^8)_6$  clusters (see Fig. S1 in the electronic supplementary information). The eigenvalue spectra of  $Au_{25}H(SR)_{18}$  and  $Au_{22}H_2(L^8)_6$ are compared in Fig. S1a,b. One can see that the orbital levels in  $Au_{22}H_2(L^8)_6$  have a much lower degeneracy than those in  $Au_{25}H(SR)_{18}$ , due to the former's lower symmetry. By analyzing the local density of states (Fig. S1c,d), we find that the 1s orbital of H is mainly involved in the HOMO-1 of  $Au_{25}H(SR)_{18}$  (Fig. S1e) and HOMO-2 of  $Au_{22}H_2(L^8)_6$  (Fig. S1f). One can see that the HOMO-1 of  $Au_{25}H(SR)_{18}$  is very delocalized, representing a superatomic P-type orbital, while the HOMO-2 of  $Au_{22}H_2(L^8)_6$  is more localized to the *cus* Au atoms in the middle and the two adsorbed H atoms. This orbital picture confirms the difference of hydrogen interaction with  $Au_{25}(SR)_{18}$  and  $Au_{22}(L^8)_6$ : the interaction with  $Au_{25}(SR)_{18}$  is dictated by the superatomic electronic structure due to the lack of *cus* Au sites; the interaction with  $Au_{22}(L^8)_6$  for the first four H atoms is rather local to and dictated by the eight *cus* Au atoms.

Implications of H adsorption on the  $Au_{22}(L^8)_6$  cluster. Our study of H adsorption on the  $Au_{22}(L^8)_6$  cluster has important implications. First, our results show that the  $Au_{22}H_2(L^8)_6$  cluster is thermodynamically very favorable, suggesting that it may be realized experimentally. This can in principle be realized by just bubbling H<sub>2</sub> gas through a solution of the  $Au_{22}H_2(L^8)_6$  cluster. Spectroscopic tools including H-NMR,<sup>28</sup> FT-IR,<sup>29</sup> and inelastic neutron scattering (INS)<sup>30</sup> have been demonstrated the ability in detecting hydride species in both homogeneous gold complexes and heterogeneous gold nanoparticles and can thus be utilized to experimentally verify the ability of  $Au_{22}(L^8)_6$  cluster in forming stable hydride species.

Second, the four *cus*-Au and two non-*cus*-Au bridge sites are potentially active sites for HER. The present work focuses on the thermodynamics of H adsorption on the  $Au_{22}(L^8)_6$  cluster and its implication for HER. The viability of the cluster for HER will of course depend on the barriers of the Volmer, Heyrovsky, and Tafel steps, which warrant further studies.

### Conclusions

In summary, we have investigated how H interacts with the  $Au_{22}(L^8)_6$  clusters by DFT calculations. We find that six H can adsorb on the  $Au_{22}(L^8)_6$  cluster: four H at the *cus* Au sites, while two H at the non-*cus* sites. More interestingly, most of the adsorbed H possess close-to-zero  $\Delta G_H$ , indicating that  $Au_{22}(L^8)_6$  can be a very good electrocatalyst for HER. In addition, we find that H behaves as a hydride or electron-withdrawing ligand in the  $Au_{22}(L^8)_6$  cluster. This is different from the metallic H in thiolate-protected Au nanoclusters, so we conclude that the behavior of H in ligand-protected Au nanoclusters is ligand-dependent. Our study demonstrates the unique behavior of H in atomically precise Au nanoclusters and invite experimental detection of H in these clusters and test of the  $Au_{22}(L^8)_6$  cluster for water splitting.

#### **Computational methods**

We employed the Vienna *ab initio* simulation package (VASP) to perform DFT calculations with periodic boundary conditions (PBC) and plane-wave bases.<sup>31</sup> We put the cluster in a 30×30×30 Å<sup>3</sup> cubic box, and the initial structure of the cluster is form the crystallographic information file provided by the experimental researchers.<sup>25</sup> Phenyl groups were replaced by hydrogens for the purpose of simplifying the calculations. The Perdew–Burke–Ernzerhof (PBE) form of the

generalized gradient approximation was used for electron exchange and correlation.<sup>32</sup> The PBE-GGA functional has been demonstrated to achieve a good balance between accuracy and cost for the ligand-protected Au nanoclusters.<sup>33</sup> The ion-electron interaction was described by projector augmented wave (PAW),<sup>34</sup> and the wave function was expanded by plane waves with a cutoff energy of 400 eV. Spin polarization was used for system of odd-number electrons. Geometry relaxations were carried out using the conjugate-gradient algorithm with a criterion that all the residual force components on each atom are less than 0.03 eV/ Å. Γ point only was used to sample the k space. Partial atomic charges were obtained using Bader charge analysis as implemented by Henkelman and co-workers.<sup>35</sup> The climbing-image nudged elastic band (CI-NEB) method implemented in VASP was used to determine the energy barriers.<sup>36</sup> The transition states were obtained by relaxing the force below 0.05 eV/Å. Periodic DFT calculations at the same level of theory were performed for the Pt(111) surface.

The Pt(111) surface was modeled with four layers of slab in  $(4 \times 4)$  lateral cells, and with 15 Å of vacuum along the z-direction. During the structure optimizations, the top two layers of the slab were allowed to relax together with the adsorbed H.

The frontier orbitals were generated from the Turbomole program (Version 6.5)<sup>37</sup> without the periodic boundary conditions at the same DFT-PBE level with the def2-SVP basis sets,<sup>38</sup> and a scalar-relativistic effective-core potential with 60 core electrons was used for Au.<sup>39</sup>

The differential hydrogen adsorption energy  $\Delta E_H$  was calculated by

$$\Delta E_H = E(catalyst + nH) - E[catalyst + (n-1)H] - \frac{1}{2}E(H_2)$$
(1)

where E(catalyst + nH) and E[catalyst + (n - 1)H] represent the total energies of the catalyst with n and n-1 adsorbed H, respectively.  $E(H_2)$  represents the total energy of a gas

phase H<sub>2</sub> molecule. A negative value of  $\Delta E_H$  suggests favorable absorption. The differential Gibbs free energy of hydrogen adsorption  $\Delta G_H$  was obtained by

$$\Delta G_H = \Delta E_H + \Delta E_{ZPE} - T \Delta S_H \tag{2}$$

where  $\Delta E_{ZPE}$  is the difference in zero-point energy between the adsorbed H and H in the gas phase H<sub>2</sub> molecule, and  $\Delta S_H$  is the entropy difference between the adsorbed H and  $\frac{1}{2}$ H<sub>2</sub> in the gas phase at the standard condition. The zero-point energy was calculated by summing vibrational frequencies  $\omega_v$  over all normal modes v:  $E_{ZPE} = \frac{1}{2} \sum \hbar \omega_v$ . The entropy of the free H<sub>2</sub> molecule at 298.15 K and 1 atm was taken from the NIST database,<sup>40</sup> while for the adsorbed H, the vibrational entropy was considered and calculated by the methods described by Cramer.<sup>41</sup>

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**Electronic supplementary information (ESI) available**: Comparison of  $Au_{25}H(SR)_{18}$  and  $Au_{22}H_2(L^8)_6$  clusters in terms of eigenvalue spectra; the local density of states; H-involving orbitals.

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Ligand-protected Au clusters with coordinatively-unsaturated (*cus*) Au atoms are predicted to be promising for realizing Au-H nanoclusters and water splitting.