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## Structure and electrical properties of lead-free $\text{Sr}_{1-x}(\text{K,Ce})_{x/2}(\text{Na}_{0.5}\text{Bi}_{0.5})\text{Bi}_4\text{Ti}_5\text{O}_{18}$ piezoelectric ceramics

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**Abstract:** Bismuth layer-structured ferroelectrics,  $\text{Sr}_{1-x}(\text{K,Ce})_{x/2}(\text{Na}_{0.5}\text{Bi}_{0.5})\text{Bi}_4\text{Ti}_5\text{O}_{18}$  (KCSNBT- $x$ ,  $x=0.0, 0.1, 0.2$  and  $0.3$ ), were prepared by a conventional solid-state reaction method. The X-ray diffraction analysis suggested that the substitution formed layered perovskite structure. Plate-like morphologies for the grains were clearly observed for all samples. The activation energy value of the ionic conductivity suggested defects were related to oxygen vacancies. Excellent electrical properties (e.g.,  $d_{33}\sim 21\text{pC/N}$ ,  $2P_r\sim 16.4\mu\text{C/cm}^2$  and  $T_c\sim 567^\circ\text{C}$ ) are simultaneously obtained in the ceramic with  $x=0.3$ . Additionally, thermal annealing studies indicated that piezoelectric constant ( $d_{33}$ ) of KCSNBT-0.1 ceramic remains almost unchanged (22 pC/N, only decrease by 4 %) at temperatures below  $400^\circ\text{C}$ , demonstrating the ceramic is the promising candidates for high-temperature applications.

**Keywords:** Bismuth layer-structured; solid-state reaction; X-ray diffraction; electrical properties

### 1. Introduction

Bismuth layer-structured ferroelectrics (BLSFs), known as Aurivillius family of oxides, have been considered as one kind of promising lead-free piezoelectrics in

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high-temperature applications<sup>1,2</sup>. Compared with traditional perovskite structure, the structure of these ferroelectrics, represented as  $(\text{Bi}_2\text{O}_2)^{2+}(\text{A}_{m-1}\text{B}_m\text{O}_{3m+1})^{2-}$ , consists of  $[\text{Bi}_2\text{O}_2]^{2+}$  layers interleaved with perovskite block of  $[\text{A}_{m-1}\text{B}_m\text{O}_{3m+1}]^{2-}$  units along the *c*-axes, where the *A*-site in the perovskite block can be occupied by 12-coordinated cations such as  $\text{Bi}^{3+}$ ,  $\text{La}^{3+}$ ,  $\text{Ba}^{2+}$ ,  $\text{Sr}^{2+}$ ,  $\text{Ca}^{2+}$ , etc, while the *B*-site can be occupied by 6-coordinated cations such as  $\text{Ti}^{4+}$ ,  $\text{Nb}^{5+}$ ,  $\text{Ta}^{5+}$ ,  $\text{W}^{6+}$ ,  $\text{Mo}^{6+}$ ,  $\text{Co}^{3+}$ , etc. In addition, the *m*, in the range of 1-5, represents the number of perovskite blocks, and the ferroelectric and dielectric properties of BLSFs are strongly dependent on the value of it<sup>3-7</sup>. The BLSFs ceramics possess some outstanding electrical properties, such as high Curie temperature, low dielectric loss, excellent fatigue endurance and low aging rate, making BLSFs a promising candidate for ferroelectric non-volatile random access memory (FRAM) storage devices<sup>8</sup>.

$\text{SrNa}_{0.5}\text{Bi}_{4.5}\text{Ti}_5\text{O}_{18}$  (abbreviated as SBT), as the *m*=5 member of the Aurivillius family where  $\text{Bi}_2\text{O}_2$  layers alternate with  $(\text{SrNa}_{0.5}\text{Bi}_{2.5}\text{Ti}_5\text{O}_{16})$  perovskite blocks built by five  $\text{TiO}_6$  octahedral layers and hosting Sr and Bi at the *A* site, has drawn a special attention because of its high Curie temperature<sup>9</sup>. However, as other BLSF compounds, this material exhibits relatively low piezoelectric coefficients ( $d_{33}$ ) and low remanent polarization ( $2P_r$ ) because of the spontaneous polarization is restricted in *a-b* plane<sup>9,10</sup>. To overcome this shortcoming, many efforts have been paid to improve its dielectric and piezoelectric properties by adding some additives in *A*-site and/or *B*-site<sup>11,12</sup>. In those studies, *A*-site substitution is more effective than that of *B*-site substitution because that the cations in *B*-site have the similar size and hardly make a major contribution to the polarization process for BLSFs. For instance, Ji et al<sup>13</sup> reported that the  $\text{SrNa}_{0.5}\text{Bi}_{4.5}\text{Ti}_5\text{O}_{18}+1\text{wt}\%\text{CeO}_2$  sample shows excellent Curie temperature (586□) and very large piezoelectric coefficient (24pC/N). Chen et al<sup>10</sup>. have

investigated the properties of  $\text{Sr}_{2-x}(\text{Na}, \text{K})_x\text{Bi}_4\text{Ti}_5\text{O}_{18}$  composite ceramics and found that the piezoelectric constant  $d_{33}$  and remnant polarization  $2P_r$  were improved to be 20 pC/N and  $20\mu\text{C}/\text{cm}^2$ , respectively, whereas the Curie temperature is only  $324^\circ\text{C}$ . Additionally, it was found that (M Ce) (M = Li, Na, K) dopants can efficiently enhance the piezoelectric coefficients for even-layer structured Aurivillius type compounds compared to the pure composition<sup>14, 15</sup>.

Nevertheless, very little is known about the detailed electrical properties on (K Ce) modified  $\text{Sr}(\text{Na}_{0.5}\text{Bi}_{0.5})\text{Bi}_4\text{Ti}_5\text{O}_{18}$  (SNBT) ferroelectric ceramics. Therefore, in the present work, A-site (K Ce)-doped  $\text{Sr}(\text{Na}_{0.5}\text{Bi}_{0.5})\text{Bi}_4\text{Ti}_5\text{O}_{18}$  ceramics were prepared by conventional solid state sintering method, and the effect of (K Ce) modification for A-site on their electrical properties was investigated. Additionally, the underlying physical mechanisms for enhanced piezoelectricity and electrical conductivity have been addressed.

## 2. Experimental details

$\text{Sr}_{1-x}(\text{K}, \text{Ce})_{x/2}(\text{Na}_{0.5}\text{Bi}_{0.5})\text{Bi}_4\text{Ti}_5\text{O}_{18}$  (abbreviated as KCSNBT- $x$ ,  $x=0, 0.1, 0.2$  and  $0.3$ ) ceramics were prepared by conventional solid-state reaction method using reagent-grade metal oxides or carbonate powders of  $\text{K}_2\text{CO}_3$ (99%, Sinopharm Chemical Reagent Co., Ltd., China),  $\text{CeO}_2$ (99.95%, Sinopharm Chemical Reagent Co., Ltd., China),  $\text{Bi}_2\text{O}_3$  (99.99%, Sinopharm Chemical Reagent Co., Ltd., China),  $\text{SrCO}_3$  (99%, Sinopharm Chemical Reagent Co., Ltd., China),  $\text{TiO}_2$  (99.5%, Sinopharm Chemical Reagent Co., Ltd., China) and  $\text{Na}_2\text{CO}_3$  (99.49%, Sinopharm Chemical Reagent Co., Ltd., China) as the starting materials. All raw materials were weighed at stoichiometric proportion and then mixed by planetary ball milling in a polyethylene with stabilized zirconia balls for 15 h, using ethanol as the solvent. After drying, the mixed powders were calcined at  $800^\circ\text{C}$  for 2 h. After calcination, the

mixture was milled again for 12 h, and then dried. The powders were mixed with an appropriate amount of polyvinyl butyral (PVB) binder, and pressed into pallets with a diameter of 12 mm and a thickness of 0.7 mm under the pressure of about 200MPa. After burning off PVB at 850°C, the ceramics were sintered in an alumina crucible at 1180°C for 3 h. For the electric measurements, disk samples with about 0.3 mm in thickness were used.

The density of the sintered ceramics was measured by means of the Archimedes method. The crystal structure of the ceramics was determined by X-ray diffraction (XRD) using a Cu K $\alpha$  radiation ( $\lambda=1.54178 \text{ \AA}$ ) (D8 Advance, Bruker Inc., Germany). The surface morphology of the ceramics was observed by scanning electron microscope (SEM) (JSM-6380, Japan). The ferroelectric hysteresis loops were measured through standardized ferroelectric test system (TF2000, Germany). The temperature dependence of dielectric properties and impedance spectroscopy for the samples was performed using a Broadband Dielectric Spectrometer (Novocontrol Germany). The samples were polarized in silicon oil in the range of 150-180°C for 20 min, and piezoelectric measurements were carried out with a quasi-static  $d_{33}$ -meter YE2730 (SINOCERA, China).

### 3. Results and discussion

The X-ray diffraction spectra of KCSNBT- $x$  ceramics ( $x=0.0, 0.1, 0.2, 0.3$ ) in the  $2\theta$  range of (a) 20-50° and (b) 28-32° was plotted in **Fig.1**. As shown in **Fig.1(a)**, the Aurivillius structure can be identified by indexing all the diffraction peaks on the basis of an orthorhombic cell (PDF#No.05-0626)<sup>13</sup>, indicating that the K<sup>+</sup> and Ce<sup>3+</sup> ions diffused into the *A*-site lattice and formed solid solutions as expected. Meanwhile, it is obvious that the highest diffraction peak of KCSNBT- $x$  ceramics is (1011) orientation, which is consistent with the fact that the strongest diffraction peak in

BLSFs<sup>16</sup>. In addition to the main Aurivillius phase, secondary phase:  $\text{Bi}_2\text{Ti}_2\text{O}_7$  (PDF#No. 32-0118, labeled by \*) in the KCSNBT- $x$  samples can also be detected regardless the existence of (KCe), as shown in **Fig.1(b)**. It seems that this peak could be related to the volatilization of bismuth species at elevated temperatures<sup>17</sup>. Similar deficiency behavior of bismuth species was also found during the sintering of another bismuth-layered compound-SrBi<sub>4</sub>Ti<sub>4</sub>O<sub>15</sub><sup>18, 19</sup>. In order to further confirm the crystallographic evolution of KCSNBT- $x$  ceramics, the lattice parameters of the samples were calculated and shown in the inset part of **Fig. 1(a)**. It is found that the lattice parameters  $a$ ,  $b$ , and  $c$  varied with increasing  $x$  values, demonstrating that the (KCe) substitution reduced the lattice distortion of SNBT-based ceramics. Such a lattice distortion can be attributed to the replacement of the  $A$ -site  $\text{Sr}^{2+}$  by  $\text{K}^+$  and  $\text{Ce}^{3+}$ , and it is supposed to lead to enhanced electrical properties.

The SEM images of the surface morphologies of KCSNBT- $x$  ceramics ( $x=0.0, 0.1, 0.2, 0.3$ ) were displayed in **Fig.2**. It is found that the grains of all the samples are all well-packed and the pore-free microstructures were presented. In addition, (KCe)-substituted SNBT ceramics have uniform grain size, indicating that the increasing doping with (KCe) does not change the microstructure dramatically. Due to the structurally highly anisotropic grain growth rate in the direction of the  $a$ - $b$  plane and the lower surface energies of the (001) planes induced predominantly during sintering, all the specimens are composed of plate-like grains, which is a typical characteristic of Aurivillius ceramics<sup>13</sup>. Moreover, all ceramics have a high relative density  $\rho_{\text{rd}}$  (>96%), suggesting all samples have been well sintered.

Temperature dependence of dielectric constant ( $\epsilon$ ) and loss ( $\tan\delta$ ) of KCSNBT- $x$  ceramics ( $x=0.0, 0.1, 0.2, 0.3$ ) measured at 10kHz is depicted in **Fig.3**. Two phase transitions were observed at  $\sim 400^\circ\text{C}$  and  $\sim 550^\circ\text{C}$ , corresponding to the phase

transitions of ferroelectric-ferroelectric transition and ferroelectric-paraelectric<sup>20</sup>. Such phenomenon was also reported in Ca-doped  $\text{SrNa}_{0.5}\text{Bi}_{4.5}\text{Ti}_5\text{O}_{18}$  ceramics<sup>13</sup>. In addition, **Fig.3** clearly shows that the value of Curie temperature ( $T_c$ ) corresponding to the ferro-paraelectric phase transition of ferroelectrics is slightly increased from  $562^\circ\text{C}$  to  $570^\circ\text{C}$ , which is higher than the reported results in other Aurivillius material systems (**Table 1**). Moreover, as shown **Fig. 3**, the dielectric loss values were found to be very low (less than 7%) and very stable when the measurement temperature is below  $400^\circ\text{C}$  for all compositions. These results exhibit that the KCSNBT- $x$  ceramics possessed the high stability of dielectric behavior, which is of great importance for high-temperature devices applications. When the temperature is above  $500^\circ\text{C}$ , the  $\tan\delta$  value significantly increased arising from the space charge carriers induced by the increase of electrical conductivity<sup>21</sup>.

**Table 1**

Comparison of polarization ( $2P_r$ ), Curie temperature ( $T_c$ ) and piezoelectric constant ( $d_{33}$ ) of KCSNBT- $x$  ceramics with other Aurivillius material systems.

Materials	$2P_r$ ( $\mu\text{C}\cdot\text{cm}^{-2}$ )	$T_c$ ( $^\circ\text{C}$ )	$d_{33}$ ( $\text{pC}\cdot\text{N}^{-1}$ )	Reference
KCSNBT- $x$	16.4	567	21	Current work
SBTi-(Na,K) $_x$	20	324	20	10
$\text{C}_x\text{SBN}$	15	552	18	16
SBBT	1.4	535	12	22
SNBT	-	535	17	23
SBN	-	440	20	24

It is well known the nature of the phase transition is determined by calculating the degree of diffusion ( $\gamma$ ) in the measured dielectric constant by modifying Curie-Weiss law<sup>5</sup>:

$$\frac{1}{\varepsilon} - \frac{1}{\varepsilon_m} = \frac{(T - T_m)^\gamma}{C} \quad (1 \leq \gamma \leq 2, T > T_m) \quad (1)$$

Where,  $\varepsilon$  is the dielectric constant,  $\varepsilon_m$  is the dielectric constant maximum,  $T$  is the temperature ( $>T_c$ ),  $T_m$  is the temperature at the dielectric peak,  $\gamma$  is the relaxation strength and  $C$  is the Curie constant, respectively. The value of  $\gamma$  is from 1 (normal ferroelectrics) to 2 (complete diffuse phase transition). The inset of **Fig.3** shows the plots of  $\ln(1/\varepsilon - 1/\varepsilon_m)$  versus  $\ln(T - T_m)$  at 10 kHz for SNBT, KCSNBT-0.1, KCSNBT-0.2, and KCSNBT-0.3 ceramics. It is observed that the values of degree of diffusion ( $\gamma$ ) for all ceramics can be described by two different temperature regions with inflexion at  $\ln(T - T_m) \sim 3.0$ , under which the degree of diffusion ( $\gamma$ ) is obtained ( $\gamma = 1.82 - 1.87$ ), while above which the values of degree of diffusion are much lower ( $\gamma = 1.23 - 1.41$ ). This phenomenon could be related to the structural distortion originated from the increasing oxygen vacancies for SNBT ceramic at high temperature<sup>25</sup>. Additionally, the values of  $\gamma$  for all samples are found to be increased with increasing (KCe) concentration and all the values are close to the degree of diffusion ( $\gamma$ ) reported for BLSFs ceramics, indicating that the addition of (KCe) in the SNBT ceramic intensifies the diffusive-type phase transition<sup>22</sup>. This can be attributed to the grain size dependence of relaxor behaviour in micro compositional fluctuations<sup>5</sup>. Such a transition from relaxor-like behavior to relaxor-ferroelectric behavior has been observed in other BLSFs<sup>26</sup>.

## Table 2

Room temperature electrical properties of KCSNBT- $x$  ceramics.



Samples ( <i>x</i> )	$\epsilon$	$\tan\delta$ (%)	$T_c$ (°C)	$E_a$ (eV)	$\rho_{rd}$ (%)	$E_c$ (kV/cm)	$2P_r$ ( $\mu\text{C}\cdot\text{cm}^{-2}$ )	$d_{33}$ ( $\text{pC}\cdot\text{N}^{-1}$ )
0.0	267	3.4	562	0.80	96.7	34	8	20
0.1	291	1.2	563	0.78	96.9	36	11	23
0.2	293	1.4	564	0.66	97.3	50	12.8	21
0.3	291	3.0	567	0.54	97.4	59	16.4	21

**Fig.4** shows the DC resistivity of KCSNBT-*x* ceramics as a function of reciprocal temperature. The resistivity of all compositions is higher than  $10^5 \Omega\cdot\text{cm}$  at  $525^\circ\text{C}$ . This is important for high temperature piezoelectric sensor applications. Furthermore, the behavior of temperature dependent resistivity follows the Arrhenius relationship:

$$\rho = A \exp(-E_a/k_B T) \quad (2)$$

where  $A$  is a pre-exponential factor constant,  $E_a$  is the activation energy of the mobile charge carriers,  $k_B$  is the Boltzmann constant and  $T$  is the absolute temperature<sup>16</sup>. The activation energy is shown in the inset of **Fig. 4** and **Table 2**. It is found that the values of activation energy  $E_a$  of KCSNBT-*x* ceramics, calculated by linear fitting of the data points, were calculated to be 0.8 eV, 0.78 eV, 0.66 eV and 0.54 eV for SNBT, KCSNBT-0.1, KCSNBT-0.2 and KCSNBT-0.3 ceramics, respectively. The results were close to the activation energy values of the ionic conductivity by oxygen vacancies in perovskite type ferroelectric oxides<sup>27-29</sup>.

**Fig.5 (a)** shows  $P$ - $E$  hysteresis of KCSNBT-*x* ceramics measured at 10Hz frequency under a maximum electric field 140kV/cm and at temperature  $180^\circ\text{C}$ . The measured remnant polarization ( $2P_r$ ) and coercive field ( $E_c$ ) values with varying (KCe) content (*x*) in SNBT ceramics are listed in **Table 2**. It is evident that the remanent

polarization ( $2P_r$ ) and coercive field ( $E_c$ ) both increases gradually with the increase of (KCe) doping content. When the doping content is 0.3, the  $2P_r$  and  $E_c$  for KCSNBT- $x$  simultaneously exhibit maximum value of  $16.4\mu\text{C}/\text{cm}^2$  and  $59\text{kV}/\text{cm}$ , respectively. The  $2P_r$  is much higher than that of pure SNBT ceramic, indicating that the ferroelectric property of SNBT ceramic has been enhanced by (KCe)-doping. The remarkable enhancement in ferroelectric polarization could be mainly attributed to the crystal lattice distortion originated from the replacement of the  $A$ -site  $\text{Sr}^{2+}$  by  $\text{K}^+$  and  $\text{Ce}^{3+}$ , which is also substantiated by XRD results that lattice parameters  $a$ ,  $b$ , and  $c$  varied with increasing  $x$  values. Moreover, the ferroelectricity in BLSFs is related to the tilting of oxygen octahedral  $\text{BO}_6$  from the  $c$  axis and the rotation in the  $a$ - $b$  plane<sup>30</sup>.<sup>31</sup> In view of the difference of mean ionic radius caused by the introduction of (KCe) into the SNBT ceramics, the tilting and rotation of oxygen octahedral  $\text{TiO}_6$  have been enlarged, which should be responsible for the enhanced ferroelectric property<sup>32</sup>. Meanwhile, the increased  $E_c$  is consistent with the assumption (**Fig.4**) that oxygen vacancies are formed. Under high electric fields, the mobile oxygen vacancies can assemble at the low energy domain walls, and thereby hinder domain switching due to domain pinning<sup>33</sup>. In order to further realize the polarization state of KCSNBT- $x$  ceramics, the polarization current curves of the ceramics were collected, as shown in **Fig. 5(b)**. One sharp polarization current peak could be observed when the applied electric field reached  $E_c$ , indicating the ferroelectric domain could be easier to switch when the driving electric field increase to  $E_c$ <sup>34</sup>.

**Fig.6** gives the thermal annealing behavior of piezoelectric constant ( $d_{33}$ ) of the KCSNBT- $x$  ceramics depolarized at different temperatures for holding 10 min. The excellent piezoelectric coefficient is founded to be  $23\text{ pC}/\text{N}$  when  $x$  is 0.1 as shown in **Table 2**, which is higher than the reported results in other Aurivillius material systems

(Table 1)<sup>22-24</sup>. In addition, the values of  $d_{33}$  were not zero when the depolarization temperature was higher than the first dielectric anomalies ( $\sim 400^\circ\text{C}$ ). However, when the depolarization temperature was higher than the second dielectric anomalies, the  $d_{33}$  was zero. Therefore, in the  $x$  range of 0~0.3, the KCSNBT- $x$  piezoelectric materials underwent a ferroelectric–ferroelectric transition at the first dielectric anomalies and a ferroelectric-paraelectric transition at the second dielectric anomalies temperatures, as shown in Fig.3. Moreover, the inset of Fig.6 clearly shows that the piezoelectric constant ( $d_{33}$ ) of KCSNBT-0.1 ceramic remains almost unchanged (22pC/N, only decrease by 4 %) at temperatures below  $400^\circ\text{C}$ , indicating that the ceramic has excellent temperature stability, so that it is very tolerant to thermal annealing and might be an appropriate candidate for high temperature applications.

#### 4. Conclusion

KCSNBT- $x$  ceramics were prepared by conventional solid state sintering method, and their structure and electrical properties were studied. The KCSNBT- $x$  ceramics presented a typical layered perovskite structure. The morphologies of Aurivillius ceramics show the grains of all the samples are all well-defined and the plate-like morphologies of samples can be obtained. The  $T_c$  increases slightly from  $562^\circ\text{C}$  to  $567^\circ\text{C}$ . The activation energy ( $E_a$ ) values of the ionic conductivity suggested defects were related to oxygen vacancies. The  $2P_r$  increase up to be  $16.4\mu\text{C}\cdot\text{cm}^{-2}$  with (KCe)-modifications. Meanwhile, the ceramics show excellent thermal stability when the annealing temperature below  $400^\circ\text{C}$ . As a result, a high  $T_c$ ,  $d_{33}$  and good thermal stability have been attained in KCSNBT- $x$  ceramics. Therefore, such a material system is a potential candidate for high-temperature piezoelectric applications.

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## Figures captions

**Fig.1.** XRD patterns of KCSNBT- $x$  ceramics sintered at 1180 °C; Inset: Detailed information on the response of the variation of lattice parameters as a function of  $x$

**Fig.2.** SEM micrographs of KCSNBT- $x$  ceramics: (a)  $x=0$ , (b)  $x=0.1$ , (c)  $x=0.2$ , (d)  $x=0.3$ ,

**Fig.3.** Temperature dependence of dielectric constant ( $\epsilon$ ) and loss ( $\tan\delta$ ) of KCSNBT- $x$  ceramics measured at 10 kHz. Inset:  $\ln(1/\epsilon-1/\epsilon_m)$  as a function of  $\ln(T-T_m)$  at frequency of 10 kHz of KCSNBT- $x$  ceramics.

**Fig.4.** Temperature dependence of resistivity of KCSNBT- $x$  ceramics

**Fig.5.** (a)  $P$ - $E$  hysteresis loops of KCSNBT- $x$  ceramics with different  $x$  at 180 °C; (b)  $I$ - $E$  loops of KCSNBT- $x$  ceramics at 180 °C.

**Fig.6.** Annealing temperature dependence of piezoelectric coefficient ( $d_{33}$ ) for KCSNBT- $x$  specimens, the inset shows the relative  $d_{33}$  (%) values of KCSNBT-0.1 sample.

## Figure captions

Figure 1

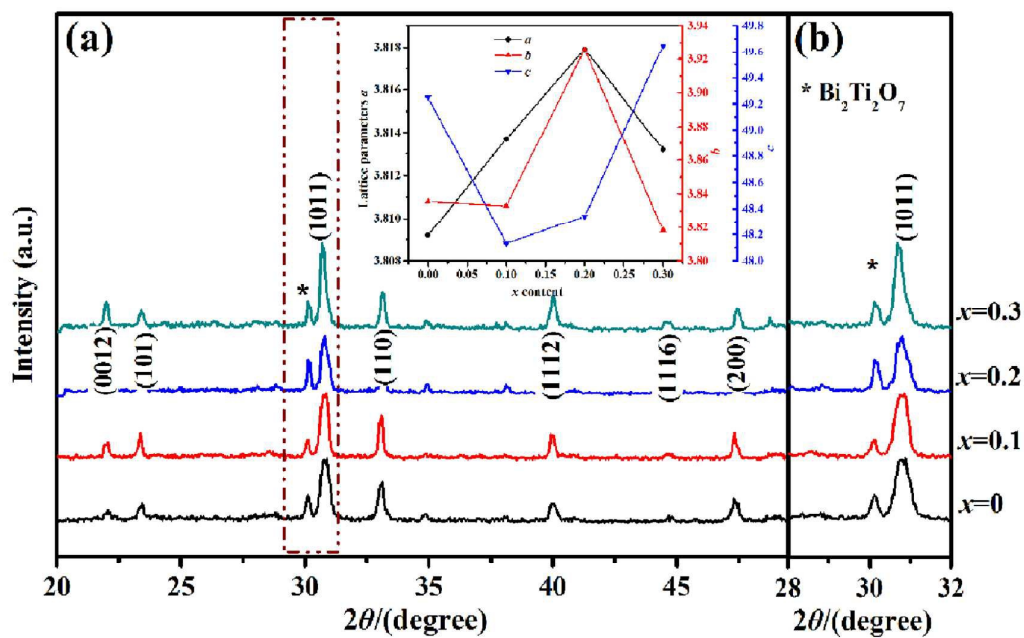


Figure 2

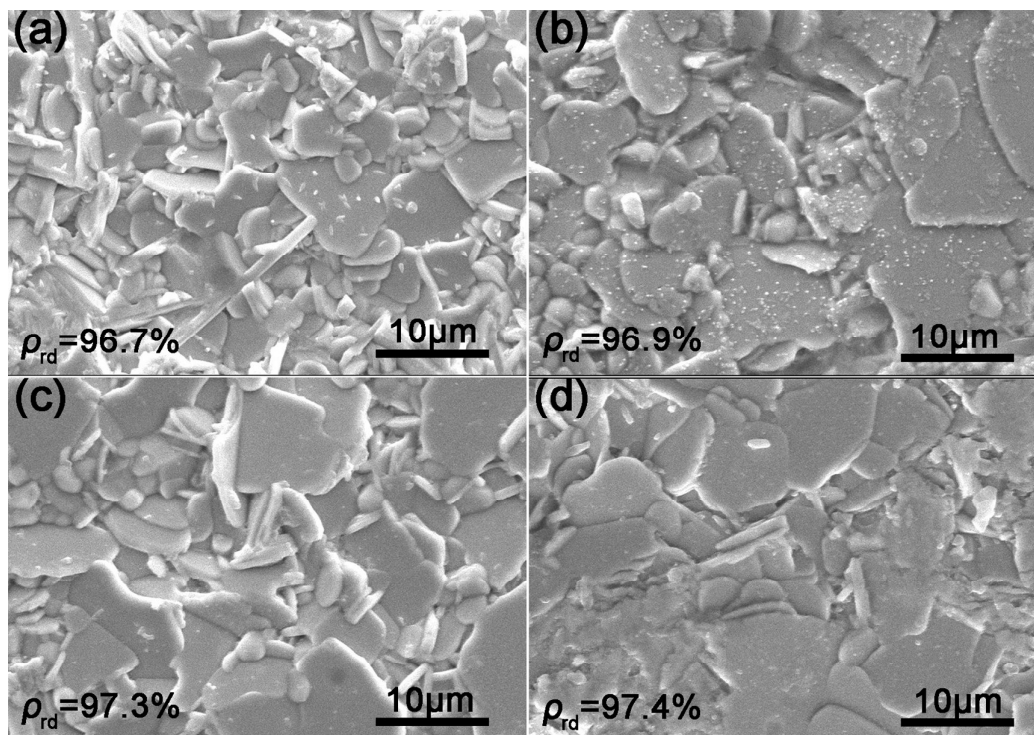




Figure 3

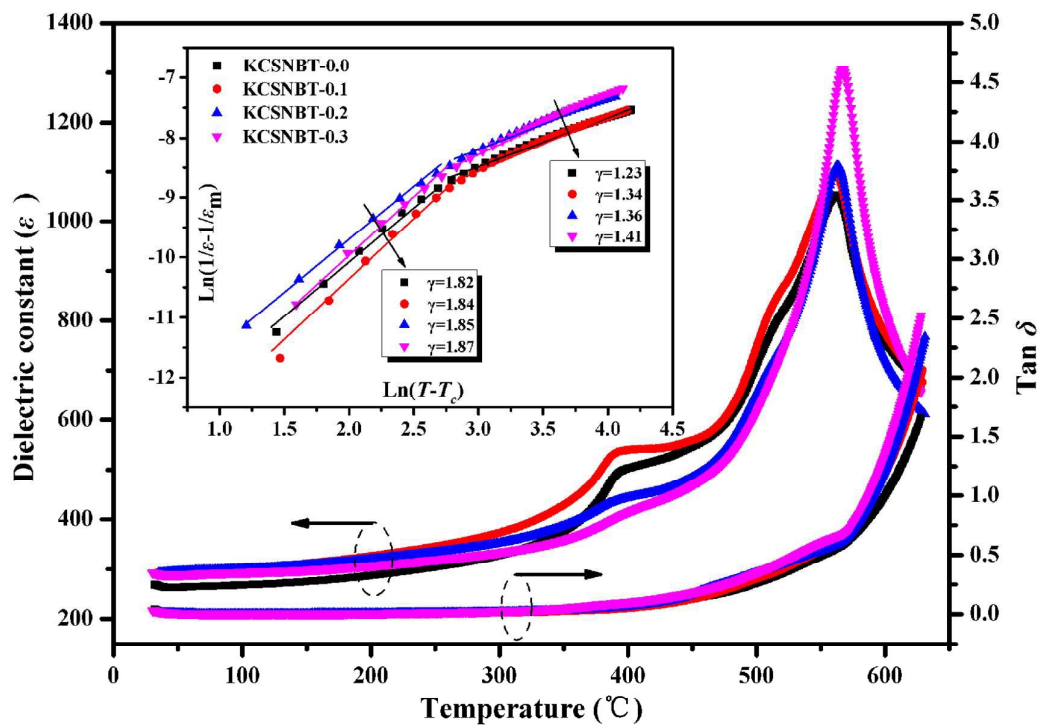


Figure 4

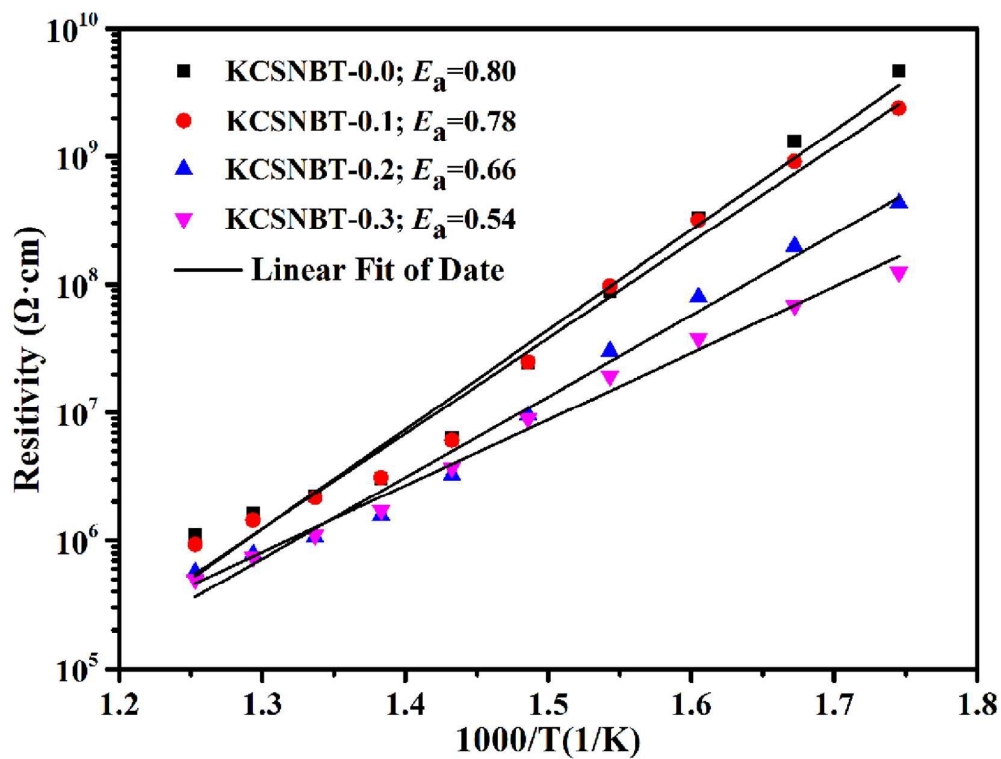


Figure 5

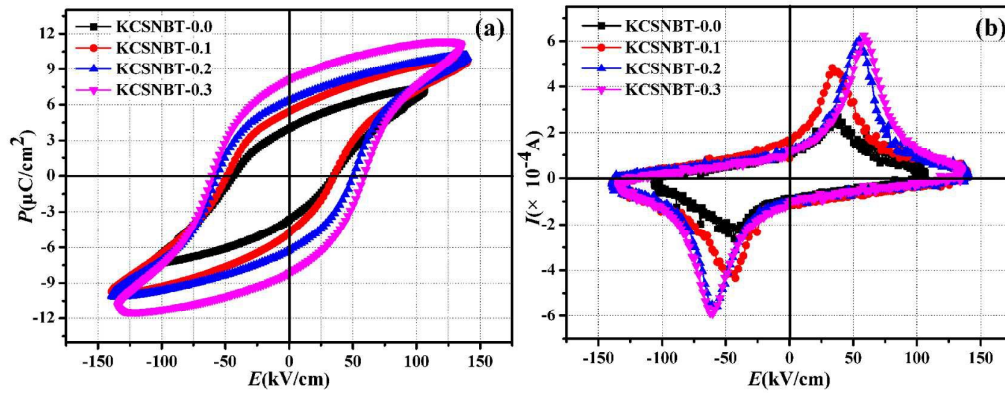


Figure 6

