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Molecular Orientation within Thin Films of Isomorphic Molecular Semiconductors

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Four isomorphic organic semiconductors are compared to map out how the precision of chemical structures determines solid state molecular organization. Intramolecular electronic structure and intermolecular packing preference in the solid state were shown exclusively dependent on the relative location of the electron accepting fragments within a given molecule.

It is possible under certain processing conditions for conjugated molecules and polymers to approach the semiconducting properties of their inorganic counterparts.¹⁻³ However, a significant challenge for practical implementation via solution processing arises from the absence of universal guidelines for managing the solid state organization as a function of molecular structure and deposition conditions.⁴⁻⁷ This lack of control prevents one to *a priori* preprogram design features in new molecules so to tailor the electronic coupling between molecules or polymer segments through control of their relative positions.⁸⁻¹¹ The latter problem is particularly exacerbated if one would like optimal electronic coupling to occur along a particular charge transport direction, *i.e.*, lateral or vertical, within a device configuration.

Organizing organic semiconductors from solution is difficult because one needs to manage weak intermolecular forces between the individual semiconducting subunits and solute/solvent interactions.¹² The general aggregation and crystallization tendencies of molecular systems relate to the former and are dependent on structural parameters including the flexibility and size of solubilizing units, electrostatic distributions and molecular topology.^{13,14} Given that synthetic protocols are available for stitching together many already-developed aromatic subunits for creating a larger delocalized

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molecule,¹⁵⁻¹⁷ one often finds that the main challenge is not how to achieve a particular target, but rather how to down select from a vast matrix of possible structural variations.

What we have presented in this paper concerns a new class of molecular frameworks that has been successfully used in the fabrication of thin film transistors and solar cells.¹⁸⁻²⁰ The general molecular structure can be expressed as D-A-D'-A'-D"-A'-D'-A-D, where D corresponds to electron rich modules (donor) and A represents electron deficient units (acceptor). The current study provides a more detailed investigation on how these molecules self-assemble into useful bulk structures and inform the readers that achieving order does not require delicate control of thin film processing. They are also useful study subjects to understand to what extent one can isolate organizational and optoelectronic properties as a function of different molecular fragments.

We provide two new molecules (NO and N2b) that were designed in view of the questions on self-assembly stated above. Specifically, we wish to probe how structural precision can be used to track the preference of molecular ordering in the bulk. Molecules in the study (N0, N2a, N2b, and N4, see Figure 1a for chemical structures) were designed so to contain similar donor units, while altering the acceptor counterparts by interchanging between pyridyl[2,1,3]thiadiazole (PT) and benzo[c][1,2,5]thiadiazole (BT). The synthetic approach to molecules NO and N2b can be found in the Supporting Information. Chemical structures were confirmed by ¹H/¹³C nuclear magnetic resonance spectroscopy and field desorption time-of-flight mass spectrometry. The sequence of structures in Figure 1a allows us to compare differences in optical properties and solid-state arrangements between four nearly isomorphic organic semiconductors. As discussed in more detail below, "N" vs "C-H" exchange within a relatively large molecule framework (molecular weight over 2200 Da) dominates molecular orientation, which emphasizes the importance of structural precision when a new material is designed.

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Figure 1. (a) Chemical structures of molecules **N0**, **N2a**, **N2b**, and **N4**, with respective illustration highlighting their structural relationship. Green filled square represents BT and red filled circle represents PT (donor segments are omitted for clarity). (b) UV/vis absorption spectra of the corresponding four molecules in dilute chloroform solutions and in the solid state. (c) Dominant hole (left) and particle (right) natural transition orbitals for S0 \rightarrow S1 transition of each molecule (from top to bottom: **N0**, **N2a**, **N2b**, and **N4**) in chloroform as calculated from TD-DFT under CAM-B3LYP/6-31G(d,p) basis set.

Table 1. Physical Properties of Molecules N0, N2a, N2b, and

Molecule	Solution λ_{\max} (nm) $^{[b]}$	Film $\lambda_{ ext{onset}}$ (nm) $^{[c]}$	$E_{\rm g} ({\rm eV})^{\rm [c]}$	$E_{\rm HOMO} (eV)^d$	E _{LUMO} (eV) ^e	7 _m (°C)	<i>Т</i> _с (°С)
N0	642	850	1.46	-4.95	-3.49	258, 269	238
N2a ^[a]	672	880	1.41	-4.96	-3.55	269	254, 244
N2b	712	915	1.36	-5.05	-3.69	259	221
N4 ^[a]	715	915	1.36	-5.19	-3.83	265	246, 230

^[a]Values obtained from reference 18. ^[b]Obtained from dilute chloroform solutions with concentrations of 0.02 mg/mL. ^[c]Thin films were prepared by spin-casting chloroform solutions (10 mg/mL) of each molecule atop pre-cleaned glass slides with film thickness of 60°70 nm as determined by X-ray reflectivity measurement. ^[d]Estimated from UPS measurement with thin films (10°20 nm) atop ITO substrates. ^[e]Calculated by $E_{LUMO} = E_{HOMO} - E_g$.

UV/vis absorption spectra (Figure 1b) of NO, N2a, N2b, and N4 in dilute chloroform solution show molar absorption coefficients (ϵ) of the lowest energy peaks in the order of N2b<N0<N4<N2a, in which N2b shows the lowest ε (1 × 10⁵ $M^{-1}cm^{-1}$) and the highest value (1.5 × 10⁵ $M^{-1}cm^{-1}$) is observed for N2a. All molecules possess distinct charge-transfer absorption bands between 500 and 800 nm principally due to efficient D-A electronic communication.²¹ Given that donor elements are identical across these molecules, it is reasonable that the absorption edges should be dominated primarily by the electron affinity of the acceptor segments. Indeed, one observes that the bandgap (E_a) of N2a is narrower (1.41 eV) than for NO (1.46 eV) because it contains a higher content of the more electron deficient PT fragments. However, molecules N2b and N4 have nearly identical absorption profiles, despite their structural differences: 2 PT + 2 BT (N2b) and 4 PT (N4).

We employed time-dependent density functional theory (TD-DFT) to gain insight into the nature of the excited states.²²⁻ ²⁴ Figure 1c illustrates the calculated dominant hole and particle natural transition orbitals (NTOs) for the S0 \rightarrow S1 transition of each molecule based on the CAM-B3LYP/6-

31G(d,p) basis set. As shown in Figure 1c, the hole is well delocalized across the entire backbone in each molecular system. In contrast, comparison of the particle NTOs shows a more delocalized nature for N2a than for N2b, which qualitatively can be understood by the presence of the more electrophilic PT units at the termini of the N2a conjugated framework. This "stretching" of the particle orbital by the PT fragment is also observed when comparing N4 with N2b.

Table 1 summarizes the physical properties for molecules **N0**, **N2a**, **N2b**, and **N4**. The highest occupied molecular orbital (HOMO) energy levels of these molecules were estimated (in eV) as -4.95 (**N0**), -4.96 (**N2a**), -5.05 (**N2b**), and -5.19 (**N4**) from ultraviolet photoelectron spectroscopy (UPS) on layers obtained by solution-deposition of each molecule atop indium tin oxide. In combination with optical band gaps (1.46 eV, **N0**; 1.41 eV, **N2a**; and 1.36 eV, **N2b** and **N4**) obtained from the solid-state absorption edges (Figure 1b), the lowest unoccupied molecular orbital (LUMO) energy levels were also estimated. This set of data indicates that HOMO and LUMO energy levels are determined not only by the relative numbers

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Figure 2. (a) 2D GIWAXS profiles for as-cast films of N0, N2a, N2b, and N4 atop silicon (100) substrates. Illustrations at the top right of each 2D image highlight the structural features of each molecule and provide possible visual correlations. (b) Out-of-plane and in-plane line-cuts obtained from the corresponding 2D GIWAXS profile. (c) Plots of scattering intensity as a function of azimuthal angle (χ) for both (100) and (010) peaks of each molecule.

the molecular framework. Thermal phase transitions of these materials were studied by differential scanning calorimetry (DSC). Melting (T_m) and crystallization (T_c) temperatures are detailed in Table 1. From the thermal transitions, one observes that all compounds have melting temperatures above 250 °C. The interplay between BT and PT does not seem to have obvious effect on the differences in T_m . However, at least within this series of compounds, the endmost PT units lead to a higher value of T_c (N2a/N4 vs. N0/N2b). Although a clear picture is not currently available to interpret thermal transitions relative to the chemical structures, as discussed in more detail below, a trend emerges that the combination of PT/BT functional groups plays a more significant role in directing the preferred molecular orientations in the solid state.

How molecular semiconductors orient themselves within thin films is relevant for understanding optoelectronic properties within a device configuration.²⁵ Here, we present a detailed structural characterization on molecules **N0**, **N2a**, **N2b**, and **N4** by using grazing incidence wide-angle X-ray scattering (GIWAXS) on spin coated thin films without subsequent thermal or solvent annealing treatments.²⁶⁻²⁹ Thin films were prepared by spin casting chloroform solution (10 mg/mL) of each material atop silicon (100) substrates with resulting film thickness of 60~70 nm as determined from X-ray reflectivity measurement.¹⁹

Figure 2a provides the 2D GIWAXS plots for thin films of **N0**, **N2a**, **N2b**, and **N4**. A summary of the relevant molecular features is depicted at the top-right corner of each image. Line cuts along both the out-of-plane (q_z) and in-plane (q_{xy}) directions are shown in Figure 2b. All molecules exhibit distinct lamellar packing (denoted as (100)) and π - π stacking (010) diffractions. Interestingly, molecules **N0** and **N2a**, both with two central BT units, show (100) peaks ($q = 4.1 \text{ nm}^{-1}$) preferentially oriented along the silicon substrate normal, with clear higher order peaks observed, *i.e.*, (200) and (300). In contrast, molecules **N2b** and **N4** exhibit (100) peaks ($q = 3.6 \text{ nm}^{-1}$) that are concentrated perpendicular to the silicon substrate normal. The lamellar *d*-spacings (d_{100}) correspond to 1.5 and 1.7 nm for **N0/N2a** and **N2b/N4**, respectively, with similar crystallite coherence lengths (L_c) of 4~7 nm observed for all molecules. However, correlating these molecular organization features with their chemical structures points to how a delicate control of chemical structure ("C-H" vs "N") can lead to completely opposite crystallographic orientations across the majority of the film composition.

In accordance to the preference of (100) peaks from this series of molecules, (010) peaks were observed also exclusively in an opposite orientation, going from N0/N2a to N2b/N4. The same $d_{010} = 0.36$ nm (q = 17.6 nm⁻¹) was obtained regardless of the difference of crystallographic orientations. To illustrate more clearly the molecular orientations, pole figures (Figure 2c) were extracted from each 2D GIWAXS plot (Figure 2a) along both (100) and (010) diffractions.^{26,30,31} The polar angle, χ , is defined as the angle between the scattering vector and the silicon substrate normal. Since electronic coupling occurs more readily along the intermolecular π - π stacking direction rather than within lamellar alkyl chain interaction, the examination of pole figures for the (010) peaks are of particular relevance, where scattering peaks near $\chi = 0^{\circ}$ and χ = $\pm 90^{\circ}$ are associated with molecular face-on and edge-on orientations, respectively. Defining R as the ratio between the integrated area where $0^{\circ} < \chi < 45^{\circ}$ and the area where $45^{\circ} < \chi < 90^{\circ}$,³² one obtains *R* as ~1:9 (10% vs. 90%), ~1:9 (12%

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vs. 88%), ~9:1 (89% vs. 11%), and ~9:1 (92% vs. 8%) for N0, N2a, N2b, and N4, respectively. This set of data illustrates the degree to which the molecular orientation of N0/N2a is predominantly edge-on, while N2b/N4 exhibits primarily face-on with respect to the substrate surface. By looking at the molecular structures one thus determines that the preferential organization is dominated by the nature of the two interior acceptor groups.

In conclusion, we showcase four isomorphic organic semiconductors with the aim of mapping out how functional groups within a given molecule influence molecular aggregation in the solid state. It is generally acknowledged that the transition from solution to the solid state remains poorly understood both theoretically and experimentally. $^{\rm 33-35}$ Simple guidance of how molecular semiconductors organize is not readily available and may ultimately prove to be systemdependent. However, within the class of donor-acceptor narrow band-gap organic semiconductors with intermediate dimensions discussed here, we are able to extract a relationship between the chemical structure and the preferred bulk molecular organization obtained via solution deposition. A "functional group"-dependent molecular organization emerges that the choice of the central electron-accepting groups correlates with molecular face-on/edge-on orientation. With increasing interest on this class of molecular semiconductors, these findings provide valuable insight to manage bridging single-molecule conformation with rather convoluted material bulk structure; a necessary step toward reliable electronic devices using molecular semiconductors.

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Subtle structural variation of isomorphic molecular semiconductors leads to sharp contrast in electronic structures and molecular orientation in the bulk.