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The anomalous compressive strain is observed in thicker  $VO_2/Al_2O_3$  epitaxial films, which is ascribed to the surface growth mode. However, the formation mechanism of the strain in the thinner  $VO_2/Al_2O_3$  epitaxial films can be explained by conventional epitaxial lattice-mismatch. Strain engineering may be served as an effective avenue for manipulating MIT behaviors of the  $VO_2$  epitaxial films.

## Surface-growth-mode-induced Strain Effects on the Metal-Insulator Transition in Epitaxial Vanadium Dioxide Thin Films

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### **ABSTRACT**

A series of high-quality vanadium dioxide (VO<sub>2</sub>) epitaxial thin films on (0001)-oriented sapphire substrates with various thicknesses were fabricated using radio frequency (RF) magnetron sputtering techniques. Structural analysis revealed that an out-of-plane tensile strain ( $\sim$ +0.035%) in the thinner VO<sub>2</sub> epitaxial films was induced by epitaxial lattice mismatch between the monoclinic  $VO_2$  films and  $Al_2O_3$ substrates. However, an anomalous compressive strain (~-0.32%) was accumulated along the out-of-plane direction in the thicker  $VO_2$  films. This result contradicts with the conventional epitaxial lattice-mismatch mechanism for strain formed in epitaxial films. We attribute this anomalous strain to the surface growth mode (island growth) in the thicker VO<sub>2</sub> films, especially those sputtered from the metal target at low pressure. Furthermore, the metal-insulator transition (MIT) temperature shifted to lower temperature with decreasing thickness, which is attributed to modulation of the orbital occupancy through the epitaxial strain and growth-mode-induced strain in the  $VO_2$  epitaxial films. Moreover, the very large resistance change (on the order of magnitude  $\sim 10^3$ ) in the VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> epitaxial heterostructures is promising for electrical switch applications.

**Keywords**: VO<sub>2</sub>; epitaxial thin film; strain effects; metal-insulator transition; surface growth mode

### 1 Introduction

Since Morin discovered vanadium dioxide (VO<sub>2</sub>) exhibiting a first-order phase transition from a monoclinic to a tetragonal crystalline structure at approximately 68 °C in 1959<sup>1</sup>, researchers worldwide have been investigating the mechanism which is responsible for the phase transition. The phase transition well known as metal-insulator transition (MIT) can be triggered not only by temperature<sup>2</sup> but also by an electric field<sup>3</sup>, light<sup>4,5</sup>, and strain<sup>6,7</sup>. The phase transition is accompanied by abrupt several-orders-of-magnitude changes in the optical and electrical properties of VO<sub>2</sub>. These properties make VO<sub>2</sub> a promising candidate for various applications, such as smart windows<sup>8-10</sup> and lithium-ion battery<sup>11,12</sup>.

The thin film of a material is often grown on a certain substrate. Single-crystal TiO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub> substrates are most commonly used for growing high-quality epitaxial VO<sub>2</sub> thin films.<sup>13-22</sup> Therefore, strains due to lattice mismatches between the VO<sub>2</sub> thin films and TiO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub> substrates must be considered, especially for practical applications based on VO<sub>2</sub> thin films. In addition, the strain provides another freedom to manipulate the metal-insulator transition.<sup>17,18</sup> Recently, TiO<sub>2</sub> substrates were chosen for growing epitaxial VO<sub>2</sub> films, where the considerable epitaxial strain can be sustained across the entire thickness of the material due to the relative small lattice mismatch (~-0.65%) between the VO<sub>2</sub> thin films and TiO<sub>2</sub> substrates. For example, Parkin *et al.*<sup>20</sup> changed the MIT temperature of VO<sub>2</sub>/RuO<sub>2</sub>/TiO<sub>2</sub> films from ~12 °C to ~72 °C by varying the thickness of the RuO<sub>2</sub> buffer layer (thus resulting in different epitaxial strains in the VO<sub>2</sub> thin films). L. L. Fan *et al.*<sup>21</sup> studied the strain states in

VO<sub>2</sub>/TiO<sub>2</sub> epitaxial films by changing the VO<sub>2</sub> film thickness and claimed that the electronic orbital occupancy was strongly affected by the interfacial strain. Although there have been many studies focused on the strain states in the typical VO<sub>2</sub>/TiO<sub>2</sub> hetero-epitaxial system, there are few systematic studies of the strain states in VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> system. Recently, Tsung-Han Yang *et al.* <sup>16</sup> deposited VO<sub>2</sub> thin films grown on c- and r-sapphire substrates using RF-magnetron sputtering and attributed the modulation of the MIT to the strain between the film and substrate. However, they didn't study the thickness-dependent strain effect on the MIT behaviors in the VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> epitaxial films. Because of the large lattice mismatch (~13%) and angle mismatch between the VO<sub>2</sub> (monoclinic) and Al<sub>2</sub>O<sub>3</sub> (rhombohedral) substrate, the strain rapidly relaxes above the critical thickness (~2.1 nm) in the viewpoint of the hetero-epitaxial system.<sup>23</sup> Therefore, the strain states of epitaxial VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> thin films have not been considered in detail. The role of strain freedom in the MIT of this system is not yet fully understood.

On the other hand, the underlying physics of the MIT of VO<sub>2</sub> thin films is still debated among the prominent theories involving a thermal-driven structural phase transition<sup>24</sup>, an electron-correlation-driven Mott<sup>25</sup> and coactions of these two mechanisms<sup>26</sup>. Recently, S. Kittiwatanakul *et al.*<sup>27</sup> reported that a large epitaxial bi-axial strain induced a Mott-like phase transition without the concomitant Peierls transition in the hetero-epitaxial system of VO<sub>2</sub>/TiO<sub>2</sub> thin films, which is similar to our findings (which are currently unpublished). This finding indicates that the strain may suppress the structural phase transition and even alter the electronic structure of

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the VO<sub>2</sub>, which makes us re-examine and further understand the mechanism of the MIT of strained VO<sub>2</sub> thin films.

Motivated by the above studies, in this work, we prepared a series of high-quality VO<sub>2</sub> epitaxial films of various thicknesses grown on (0001)-Al<sub>2</sub>O<sub>3</sub> substrates (*c* oriented) using RF-magnetron sputtering technique. Another reason for the choice of Al<sub>2</sub>O<sub>3</sub> substrate is of wide availability for fabricating high quality epitaxial VO<sub>2</sub> films on it. The strain evolutions were investigated in the different thickness VO<sub>2</sub> films. We reported anomalous strain states in this VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> system, which have not been observed previously and ascribed the reason to the specific surface growth mode.

### 2 Experimental Procedures

### 2.1 Film growth and thickness characterization

A series of VO<sub>2</sub> thin films of different thicknesses were deposited on commercial single (0001)-oriented Al<sub>2</sub>O<sub>3</sub> substrates via RF-magnetron sputtering. The base pressure was pumped to  $2 \times 10^{-4}$  Pa. A vanadium metal target (99.99% purity) was sputtered at a 60 W RF power and sputtering pressure of 0.43 Pa with an Ar:O<sub>2</sub> (99.99%) flow ratio of 60:0.5.<sup>19</sup> The growth temperature was set to 350 °C, and the growth times were 15, 20, 30 and 60 min for the VO<sub>2</sub> thin films with various thicknesses, respectively.

The thickness measurement was conducted via small angle X-ray reflectivity (XRR) measurement, which is an effective non-destructive tool to measure the

thickness. With XRR, we have determined the thickness of all samples.

### 2.2 Structure characterization

The structure and quality of the VO<sub>2</sub> films were characterized by high-resolution XRD with Cu K $\alpha_1$  ( $\lambda$ =1.5406 Å) radiation (Rigaku SmartLab Film Version), and the partial structure information was obtained from the 14B beamline of the Shanghai Synchrotron Radiation Facilities with the reciprocal space mapping (RSM) technology. The structural phase transition across the MIT of the VO<sub>2</sub> thin films was investigated by *in situ* temperature-dependent XRD. XRD scans were performed when the *in situ* temperature was stable.

### 2.3 Surface morphology characterization

The surface morphology of the  $VO_2$  thin films were examined by the Atomic Force Microscopy (AFM, Bruker Dimension Icon) in the ScanAsyst mode. Then the surface roughness values (root mean square) were obtained by averaging the total surface area in the AFM images.

### **2.4 Electrical transport measurements**

With the help of the multifunction of Physical Properties Measurement System (PPMS, Quantum Design, Inc.), the hysteresis loops of the resistivity as a function of temperature were measured by four-probe method in the temperature range of  $25\sim120$  °C with the accuracy of 0.05 °C.

### **3** Results and Discussion

### **3.1 Epitaxiality characterization**

To evaluate the quality of the VO<sub>2</sub> thin films, XRD patterns were acquired for a typical VO<sub>2</sub> thin film with thickness of  $\sim$ 60 nm and are presented in Fig. 1. As can be observed from Fig. 1(a), only (020) peak of the  $VO_2$  film and three peaks of the  $Al_2O_3$ substrate appeared, thereby indicating the high orientation and good purity of the VO<sub>2</sub> thin films. The inset of Fig. 1(a) is a typical X-ray reflectivity (XRR) curve showing the distinct interference fringes. The corresponding thickness can be calculated to be  $\sim 60$  nm via Bragg formula.<sup>21</sup> According to the rocking curve of the VO<sub>2</sub> (020) peak (see Fig. 1(b)), the full width at half maximum (FWHM) is only about  $0.54^{\circ}$ , which indicates the good quality of the VO<sub>2</sub> thin films. <sup>19,28</sup> From the fine scanning curve shown in Fig. 1(c), it can be found that the  $VO_2$  (020) peak and sapphire (0006) peaks locate at 39.955° and 41.685°, respectively, which is consistent with the results reported in ref. 29. The reciprocal space map (RSM) of the VO<sub>2</sub> (220) peak and Al<sub>2</sub>O<sub>3</sub> (116) obtained after performing  $\varphi$  scans<sup>19</sup> is presented in Fig. 1(d); it demonstrates the perfect epitaxial relationship between the VO<sub>2</sub> film and Al<sub>2</sub>O<sub>3</sub> substrate. Consequently, the epitaxial relationship is derived to be  $[101]_{VO_2} / [11\overline{2}0]_{Al,O_3}$  along the in-plane direction and  $[010]_{VO_2}$  //[0001]<sub> $Al_2O_3$ </sub> along the out-of-plane direction.<sup>28,29</sup> Interestingly, the  $VO_2$  (220) peak splits into three parts representing three different oriented domains, and we will discuss this result in another paper. We deposited four  $VO_2$ films of different thicknesses with different growth times of 15, 20, 30 and 60 min, respectively. The thicknesses were determined to be  $\sim 30, \sim 40, \sim 60, \text{ and } \sim 120 \text{ nm by}$ non-destructive small angle X-ray reflectivity measurement (typically shown in the inset of Fig. 1(a)).



Fig. 1 (color online) (a) Typical  $\theta$ -2 $\theta$  scan of the VO<sub>2</sub> film grown on c-Al<sub>2</sub>O<sub>3</sub> substrate by RF magnetron sputtering, inset: the typical XRR curve of a 60 nm-thickVO<sub>2</sub> thin film. (b) Rocking curve of the VO<sub>2</sub> (020) peak. (c) Fine scan curve of the VO<sub>2</sub> (020) peak from 38.6° to 42.6°. (d) RSM maps of VO<sub>2</sub> (220) and Al<sub>2</sub>O<sub>3</sub> (116).

### 3.2 Thickness-dependent MIT of the VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> epitaxial films

In order to analyze the MIT characteristics of the  $VO_2$  thin films, the resistance-temperature (R-T) curves of the four samples were measured by PPMS. As shown in Fig. 2, during the temperature cycling from 25 to 120 °C and reversing back, the normalized resistances of 40, 60 and 120 nm samples increased by approximately three orders of magnitudes, showing an obvious and sharp transitions with the typical hysteresis appearing between the heating and cooling processes. The thinnest  $VO_2$ 

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film (~30 nm) also shows an obvious MIT; however, the change of resistance across the transition was most likely suppressed by strain in the film.<sup>21</sup>

The differentiations of the R-T curves obtained using the formula  $d \log(R)/dT$  are presented in the insets of Fig. 2. As deduced from the minimum of each differentiation curve, the transition temperatures of the four samples were 60, 63, 64 and 67 °C during the cooling process and 71, 74, 72 and 75 °C during the heating process. It can be inferred from the multiple measurements and averaging the transition temperatures that the MIT temperatures has an increasing tendency with increasing thickness of the VO<sub>2</sub> film in a large content, which is shown in the red curves in Fig. 3(b).



Fig. 2 (color online) R-T curves of the typical batch of the (a) 30-nm, (b) 40-nm, (c) 60-nm and (d) 120-nm samples measured by PPMS by single cycling temperature. Insets: differential curves of

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each R-T curve for all four samples.

### 3.3 Origin of the thickness-dependent transition temperature

According to previous studies<sup>20,21,30</sup>, modulation of the transition temperature across the MIT may be induced by interfacial strains between the VO<sub>2</sub> film and the substrate. Therefore, we attributed the change of transition temperature to the interfacial strain. We performed XRD  $\theta$ -2 $\theta$  scans for the four VO<sub>2</sub> thin films with different thicknesses. The XRD patterns shown in Fig. 3(a) contain sharp peaks at 39.9° that are associated with the Al<sub>2</sub>O<sub>3</sub> substrate (0006) peaks, which do not shift among the four samples. With increasing thickness, the VO<sub>2</sub> diffraction peak exhibits a clear shift to higher angles (2 $\theta$  are 39.785°, 39.825°, 39.875° and 39.955° respectively, which is indicated by the guided line). According to Bragg's law, the lattice constant  $d_f$  were calculated to be 2.26389 Å, 2.2628 Å, 2.26008 Å and 2.25574 Å, respectively. Thus, the strain between the VO<sub>2</sub> thin film and the substrate is defined as:

$$strain = \frac{d_f - d_b}{d_b} \times 100\%, \qquad (1)$$

where  $d_f$  and  $d_b$  ( $d_b$ =2.2631 Å) represent the lattice constants of the VO<sub>2</sub> thin film and the counterpart bulk one, respectively. Following this definition, the strain vs. thickness curve was obtained and is presented in Fig. 3(b). It can be found that the strain changes from tensile (conventionally positive) to compressive (and hence, negative) strain.



Fig. 3 (color online) (a) Fine  $\theta$ -2 $\theta$  scans of the VO<sub>2</sub> (020) peaks of the four different thickness samples. The red guided line represents the peak shift. (b) The out-of-plane strain and transition temperature lines are plotted as a function of the sample thickness. The error bars base on the statistical average of the MIT temperature by multiple measurements. The schematic illustrations of the (c) tensile and (d) compressive strain along the out-of-plane direction in the thin and thick VO<sub>2</sub> thin films, respectively.

Comparing the three curves shown in Fig. 3(b), the transition temperature exhibits a tendency opposite to the strain both at heating and cooling process. It is believed that this behavior is modulated by the interfacial strain, which is consistent with the conclusions of ref. 20 and 21. The modulation mechanism is that the strain in the  $VO_2$  film modifies the orbital occupancy near the Fermi energy level  $E_F$ . The

compressive strain along the out-of-plane direction causes the  $d_{\parallel}$  and  $\pi^*$  orbitals to shift upward and the  $d_{\parallel}^*$  orbital downward. Therefore, in the insulator state, the band gap increases by the out-of-plane compressive strain along the  $b_M$  direction in the (020)-VO<sub>2</sub> epitaxial film, the energy barrier increases and thus the MIT is triggered at higher temperature. Consequently, the electronic orbital occupancy changes due to the epitaxial strain, and the surface-growth-mode-induced strain modulates the phase transition behavior and triggers the onset of the MIT at higher temperature in the thicker and compressively strained VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> thin films.<sup>20,21</sup>

It is noteworthy that the compressive strain state (Fig. 3(d)) along the out-of-plane direction is observed in the thicker films; however, a tensile strain exists in the thinnest film (Fig. 3 (c)). Moreover, the compressive strain becomes even stronger in the thicker films. From hetero-epitaxial system<sup>23</sup>, only tensile but not compressive strain can be maintained in the VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> epitaxial films along the out-of-plane direction, which should be relaxed with increasing film thickness. As for the thinnest sample (~30 nm), the strain along the out-of-plane direction is tensile, which is consistent with the strain mechanism of a hetero-epitaxial system, for which the lattice mismatch (~+13%) between the bulk VO<sub>2</sub> and sapphire substrate has a positive sign and should induce a tensile (positive) strain in the VO<sub>2</sub> thin film (Fig. 3(c)). However, for the thicker samples, a compressive strain developed; this result distinctly contradicts the strain relaxation in the conventional framework of the hetero-epitaxial system. Therefore, there must be another cause of the strain in our sputtered three thick VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> films. It is mentionable that the typical batch of the

samples were selected in this work, although the other batches of the samples showing the similar rules were not shown here.



Fig. 4 (color online) Schematic diagrams of the  $VO_2$  film surface growth modes: (a) two-dimensional growth mode, (b) and (c) island growth mode. The schematic diagrams in Fig. 4(a), (b) and (c) is not to scale.

According to previous literature<sup>31-37</sup>, the anomalous strain behavior is closely related to the growth mode of the VO<sub>2</sub> films. At the very initial stage, under the critical thickness, the VO<sub>2</sub> films grow two-dimensionally on the sapphire substrate (see Fig. 4(a)). The evolution of the strain follows the mechanism of hetero-epitaxial system. Therefore, the strain state in the thinnest VO<sub>2</sub> films is tensile along the out-of-plane direction from the viewpoint of the lattice-mismatch-induced epitaxial strain, as illustrated in Fig. 3(c). With increasing growth time, the strain starts to relax in the thicker films, as illustrated by the dotted lines in Fig. 4(a). Simultaneously, the VO<sub>2</sub> nucleates discretely and forms isolated three-dimensional islands above the critical thickness (see Fig. 4(b)). At this moment, the VO<sub>2</sub> islands are isolated each

other and there is no strain between them. Then the islands coalesce with further increasing thickness, and for balancing the surface energy of the islands and the free energy of the grain boundaries<sup>33,36-38</sup>, the lattice is extend along the in-plane direction and thus the in-plane tensile strain appears. According to the volume conservation in the solid thin films, the out-of-plane compressive strain is thus induced as shown in Fig. 4(c). This effect becomes obvious especially in the sputtering process from the metal targets at the low pressure.<sup>33,38</sup> Therefore, an out-of-plane compressive strain, which balanced the in-plane tensile strain, was observed in the thicker three films and when the films become thicker, the compressive strain becomes stronger. This anomalous strain is the so-called surface-growth-mode-induced strain as reported in refs. 31 and 34.



Fig. 5 (color online) Surface morphology observed using AFM and surface roughness values

(RMS: root mean square) of (a) 30-nm, (b) 40-nm, (c) 60-nm, and (d) 120-nm samples.

To verify the mechanism described above, we performed AFM studies on the VO<sub>2</sub> film surfaces of various thicknesses. Fig. 5 shows the surface morphologies of the four samples and clearly supports the above explanation. In Fig. 5(a) and (b), the roughness of 30-nm and 40-nm samples is less than that of 60-nm and 120-nm samples. Moreover, the islands are formed individually and appear not to overlap with each other in the thinner 30-nm and 40-nm films. The islands start to coalesce in 60-nm films, as evident from Fig. 5(c), and the three-dimensional islands become much clearer in 120-nm films, as can be observed from Fig. 5(d). This explanation has also been demonstrated experimentally in ultra-thin FeSe<sub>0.5</sub>Te<sub>0.5</sub>/LaAlO<sub>3</sub> and MgB<sub>2</sub>/SiC epitaxial films, although an in-plane compressive strain could be expected in these films under the conventional framework of the lattice mismatch mechanism.<sup>33,34</sup>

### **4** Conclusions

We deposited a series of VO<sub>2</sub> epitaxial films of different thicknesses via RF magnetron sputtering and studied their structural and electrical properties. We observed an out-of-plane tensile strain (~+0.035%) in the thinnest VO<sub>2</sub> epitaxial films that was induced by an epitaxial lattice mismatch and a considerable out-of-plane compressive strain (~-0.32%) that was induced by surface growth mode. Furthermore, the MIT temperature shifted to lower temperature with decreasing thickness, which is

attributed to modification of the orbital occupancy caused by the hetero-epitaxial strain and growth-mode-induced strains. Our results indicate that strain engineering may be served as an effective avenue for manipulating MIT behaviors of the VO<sub>2</sub> epitaxial films.

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