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Size Dependent Electronic Band Structures of β- and γ-Graphyne Nanotubes

Baotao Kang a,*, Jong Hun Moon b and Jin Yong Lee b,*

In the present paper, density functional theory calculations have been implemented by using Dmol³ to study the electronic band structures of β-graphyne nanotubes (βGyNTs) and $γ$ graphyne nanotubes (γGyNTs). Our results found different GyNTs show diverse electronic band structures. All βGyNTs have quite small band gaps without any correlation with tube size. Meanwhile, γGyNTs, no matter zigzag or armchair, exhibit semiconductor characteristic with oscillatory band gap ranging from 0.48 eV to 1.20 eV. Furthermore, based on the variation of band gap, both zigzag and armchair γ GyNTs can be divided into two subgroups: 2m and 2m + 1 where n is positive integer, following the order of $2m + 1 > 2m$.

Introduction

During last two decades, carbon-based materials including fullerenes,^{[1,](#page-4-0) [2](#page-4-1)} carbon nanotub[e](#page-4-4)s $(CNTs)$ ^{[3,](#page-4-2) [4](#page-4-3)} and graphene⁵, have received enormous attentions from theorists and experimentalists. Especially, CNTs constituted by only sp^2 -like carbons could be either metallic or semiconducting, and they have been intensively exploited in electronic devices.^{6,} Recently a new carbon allotrope by inserting acetylenic linkage (-C≡C-) into graphene, named graphyne, has been received growing attentions. The coexistence of sp and sp^2 carbons grants graphyne outstanding characteristics.^{[8-12](#page-4-7)} These unique properties could allow graphyne to have some potential applications. ^{[13-15](#page-4-8)} The presence of acetylenic linkages allows graphyne to exist in several different two-dimensional structures, such as α -, β -, and γ - graphyne. ^{[9-11](#page-4-9)} Since CNTs can be considered as seamless cylinders of graphene sheets, we can also build graphyne nanotubes (GyNTs) via the same approach. Recently, graphdiyne nanotubes (GDyNTs) has been successfully synthesized and exhibited high-performance field emission properties^{[16](#page-4-10)}, which gives a possibility for realization of single-walled graphyne nanotubes (GyNTs).

Comparing with CNTs, GyNTs have not received much attention though it deserved.^{[17,](#page-4-11) [18](#page-4-12)} Computational studies have been carried out to investigate the electronic band structure via tight-binding (TB) method. $19, 20$ $19, 20$ However, there are two very import parameters ignored in TB method, σ-π hybridization effect 21 and the coexistence of *sp* and sp^2 hybridized carbon. Such a weakness of TB has been exposed by recent studies on CNTs. Theoretical and experimental approaches have proved that zigzag (n, 0) CNTs are typical semiconductor when $n \neq$ $3m^{22, 23}$ $3m^{22, 23}$ $3m^{22, 23}$ $3m^{22, 23}$ while zigzag (n, 0) CNTs when $n = 3m$ were predicted to be metallic via the TB method. [24,](#page-4-18) [25](#page-4-19) Afterwards, recent research based on DFT calculations disclosed a small energy band gap, 26 , 27 which has been experimentally confirmed. The $(n, 0)$ SWCNTs with $n = 3m$ showed semiconductor property with band gap of 0.080 \pm 0.005 eV for (9, 0), 0.042 \pm 0.004 eV for (12, 0) and 0.029 \pm 0.004 eV for (15, 0) SWCNT. As such different electronic band structure of αGyNTs by TB

method compared with our calculations based on density functional theory (DFT) was revealed. ^{[29](#page-4-23)} To our knowledge, DFT calculations have not been systematically implemented for β-graphyne nanotubes (βGyNTs) and γ-graphyne nanotubes (γGyNTs), which motivates us to undertake the present study. From the calculated electronic band structure, we found that βGyNTs and γGyNTs possess different electronic structures from ordinary CNTs and αGyNTs. βGyNTs have small band gaps without correlation with tube size, while γ GyNTs are typical semiconductor but with different oscillatory band gap behavior compared to the CNTs.

Computational Method

Fig. 1. Band structures of 2D β-graphyne (left) and γ- graphyne (right).

In this study, DFT-D calculations were carried out by DMol³ module^{[30](#page-4-24)} in Material Studio 5.5. In earlier paper, different DFT method including PWC functional ^{[31](#page-4-25)} belonging to LDA, PW91, ³¹ BLYP 32 and PBE fuctionals $33, 34$ $33, 34$ belonging to GGA, were implemented to investigate the band structures of αGyNTs and gave nearly the same results. Thus, in the present paper PBE method was utilized with the consideration of dispersive interaction correction (PBE-D). Such PBE-D method has been intensively applied to study carbon related materials.[35,](#page-4-29) [36](#page-4-30) Moreover, there is no spin restriction imposed during calculations. For all calculations, 20 Å vacuum space was imposed to avoid interlayer interactions. All electron treatment was performed and a double numerical plus polarization (DNP) basis set was used. The convergence tolerance of energy was 10-5 hartree, and the maximum allowed force and displacement were 0.002 hartree/Å and 0.005 Å, respectively. The unit cell of β- and γgraphyne was fully optimized with 2D hexagonal symmetry with kpoint meshes of $7 \times 7 \times 1$, and the lattice constant (a) was computed to be 9.527 and 6.923 Å, respectively. As noted in Figure 1, the band structure of β-graphyne has a zero band gap, while that of γgraphyne has a band gap (E_g) of 0.47 eV at M high symmetric point. Dispersive interactions are important in electronic structure calculations of carbon materials such as graphene. 37 Herein we compared the *Eg* obtained with and without dispersion correction. The E_g without dispersion correction was reported to be 0.42 eV, 20 20 20 which was further verified by our result of 0.43 eV by merely PBE method. Since pure DFT often underestimates band gap, dispersion correction could relieve such a problem.

Based on the optimal unit cell of $β$ - and γ-graphyne, the corresponding zigzag and armchair nanotubes were built by rolling up film along different orientation. The structure of each nanotube was first fully optimized with Brillouin zone k-point meshes of $1 \times 1 \times 4$. Then the k-point meshes was increased to $2\times2\times50$, which should be large enough to obtain reliable electronic band structure (50 is the maximum value allowed by Dmol³). For each tube, the cohesive energy (*Ecoh*) was calculated as follows:

 $E_{coh} = E_{total}/N_C - E_C$

where *Etotal*, *N^C* and *E^C* stand for the total energy, number of carbon atoms and energy of isolated carbon atom, respectively. More negative *Ecoh* implies energetically more stable structure.

Results and Discussion

Fig. 2. Optimized structures of (4, 0)- and (4, 4)-βGyNT and γGyNT.

Fig. 3. Cohesive energy vs. tube size of zigzag and armchair graphyne nanotubes.

As we did in earlier study 29 , the chirality of tubes was characterized by applying nomenclature (n, m). As shown in Figure 2, the chirality of β GyNTs and γ GyNTs is slightly different from that of αGyNTs possessing hexagonal carbon ring. To maintain consistency with αGyNTs, zigzag and armchair were still used to name the chirality of βGyNTs and γGyNTs. Following such notation, (n, 0) and (n, n) stand for zigzag and armchair nanotubes, respectively. These tubes were denoted by N_z -Z- β (γ) GyNTs and N_a -A- β (γ) GyNTs, where N_z and *N^a* represent the tube size of zigzag and armchair nanotubes, respectively. Figure 2 shows the optimized structures of βGyNTs and γGyNTs with *N^z* and *N^a* being 4 for instance. The relationship between cohesive energy (*Ecoh*) and tube size (*N^z* or N_a) was pictured in Figure 3. The cohesive energy for β GyNTs and γGyNTs monotonically decreased as the tube size increases implying gradual reduction in surface strain. For both βGyNTs and γGyNTs, the cohesive energy of armchair was always smaller than that of zigzag when $N_z = N_a$ because A-GyNTs has larger diameter. The cohesive energy of βGyNTs and γGyNTs gradually converged to -7.11 and -7.32 eV/atom since surface strain tends to decrease. Those values are larger than that of αGyNTs with -7.02 eV/atom. Such values are comparable to that of fullerene with -7.29 eV/atom 38 that makes the single walled GyNT to be realized in near future.

Fig. 4. The band structures of β GyNT and γ GyNT for $N_z/N_a = 3$ and 4.

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Then the corresponding electronic band structure of βGyNT with N_z = 3 and 4 were plotted in Figure 4. For 3-Z- β GyNT, one can see that the valence band maximum (VBM) reaches to the Fermi level (*EF*) and the conduction band minimum (CBM) locates above E_F , inducing a direct band gap (E_g) of 0.03 eV between Γ and Z point. Meanwhile, the location of such direct band gap moves to the top of Γ point for 4-*Z-*βGyNT whose *E^g* was 0.05 eV. Further analysis confirmed the dependence of band gap's location upon the parity of N_z . To elucidate the size dependence of band gap, the variation of E_g as a function of N_z was shown in Figure 5 and no correlation between E_g and tube size was found for Z-βGyNTs.

The band structure of Z-βGyNTs is distinctly [diverse](http://www.iciba.com/diverse) from zigzag CNTs and α GyNTs. It was determined that zigzag $(n, 0)$ CNTs are semiconductor and show an oscillatory band gap dependence on tube size (n). The band gap follows the rank order: $3m - 1 > 3m + 1 > 3m$. ^{[22,](#page-4-16) [23,](#page-4-17) [28](#page-4-22)} Recently, we carried out reliable DFT calculations to investigate the electronic structures of αGyNTs and found Z-αGyNTs possess qualitatively similar band gap behavior.^{[29](#page-4-23)} All Z- α GyNTs are semiconductors and exhibit strong dependence on tube size. N_z -Z- α GyNTs can be classified into three families according to band gap variation: *N^z* $= 3m - 1$, 3m, and 3m + 1 resulting in the following rank order of band gap magnitudes: $3m - 1 > 3m + 1 > 3m$.

For the band structure of A-βGyNTs, as displayed in Figure 4, the VBM reaches to *EF*, and the CBM locates above *EF*, inducing a direct band gap between Γ and Γ point. The E_g of AβGyNTs was calculated to be about 0.01 eV except for the cases of $N_z = 4$ and 9. The E_g of 4- and 9-A- β GyNT were 0.06 eV and 0.11 eV, respectively. Considering DFT usually underestimates band gap, A-βGyNTs could be concluded as semiconductor. Similar to Z-βGyNTs, A-βGyNTs do not show a size dependence of band gap as noted in Figure 5 and have different electronic band structure compared to armchair CNTs and A-αGyNTs. Armchair CNTs had been experimentally determined to be metallic.^{[28,](#page-4-22) [39](#page-4-33)} Furthermore, A- α GyNTs were concluded to be semiconductors when tube size is small and showed downward trend on tube size, then becomes metallic as the tube size increases.

The band structure of γGyNTs showed that the VBM reaches to *E^F* and the CBM locates above *EF*, thus resulting in direct band gaps. Similar to Z-βGyNTs, the band gap location of γGyNTs was found to be related with the parity of tube size. When N_z or N_a was odd, the direct band gap located at Γ point, while that of even N_z or N_a located at Z point. Furthermore, all γGyNTs were semiconductors with moderate E_g ranging from 0.48 eV to 1.20 eV as pictured in Figure 5. Moreover, it was further revealed that the band gap of Z-γGyNTs exhibited oscillatory and size dependent behavior, which is similar to zigzag CNTs and αGyNTs. However, differently from zigzag CNTs and αGyNTs, Z-γGyNTs could be classified into two families with $N_z = 2m$ and $2m + 1$ (n is a positive integer) following the order of band gap magnitudes: $2m + 1 > 2m$ according to band gap variation. Additionally, armchair CNTs, αGyNTs and βGyNTs are either semiconductor with very small band gap or metal, while A-γGyNTs are typical semiconductor

with moderate band gaps. Similar to Z-γGyNTs, A-γGyNTs can also be divided into two subgroups with $N_a = 2m$ and $2m + 1$ following the order of band gap magnitudes: $2m + 1 > 2m$. The oscillatory behavior of the band gap of γGyNTs was also observed from DFTB study.[17](#page-4-11) But one point should be emphasized that the band gap of A-γGyNTs in present paper is always larger than that of Z- γ GyNTs when $N_a = N_z$, while it is overturned by DFTB method.[17](#page-4-11) Such disharmony may originate from the inadequacy of tight-binding-link treatment in DFTB method. It would be elusive until more expensive calculations like GW method ^{[40](#page-4-34)} or experiments report the band gaps of γGyNTs in the future. In addition, the band gaps of both ZγGyNTs and A-γGyNTs gradually decreased as tube size increased, and were predicted to converge to that value of γ graphyne film.

Fig. 5. The variation of β and γGyNTs' band gaps as a function of tube size.

Conclusions

Our DFT calculation showed different GyNTs have diverse electronic band structures. αGyNTs possess qualitatively similar band gap behaviour to ordinary CNTs. Z-αGyNTs were disclosed as semiconductors with band gap which is strongly dependent on tube size. N_z -Z- α GyNTs can be classified into three families: $N_z = 3m - 1$, 3m, and 3m + 1 resulting in the following rank order of band gap magnitudes: $3m - 1 > 3m +$ $1 > 3$ m. Meanwhile, A- α GyNTs are semiconductors with small tube size then becomes metallic as the tube size increases. [29](#page-4-23) However, βGyNTs appeared to be atypical among CNTs and GyNTs. All βGyNTs have quite small band gaps but without any correlation with tube size. For γGyNTs, no matter zigzag or armchair, semiconductor character was clearly observed with moderate band gap of $0.48 \sim 1.20$ eV. Both Z- and A-γGyNTs can be divided into two subgroups: $2m$ and $2m + 1$ where n is positive integer, following the order of band gap magnitudes: $2n + 1 > 2n$. Their band gaps decreased as tube size increases, and were predicted to converge to the value of 2D γ -graphyne. Moreover, the band gap of A-γGyNTs is always larger than that of Z-γGyNTs. According to the electronic properties, γGyNTs would be a potential candidate as semiconductors with controllable band gap by tuning the tube size.

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Notes and references

^a School of Chemistry and Chemical Engineering, University of Jinan, Jinan, Shandong 250022, People's Republic of China, *E-mail: kangbaotao@hotmail.com

^b Department of Chemistry, Sungkyunkwan University, Suwon, 440-746, Republic of Korea, *E-mail: jinylee@skku.edu

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