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# Effect of dopant concentration on the spectroscopic properties in In<sup>3+</sup> doped (0,1,2 and 4 mol%) Yb: Tm: LiNbO<sub>3</sub> crystal

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# Abstract

A series of Yb: Tm: LiNbO<sub>3</sub> crystals with x mol%  $In^{3+}$  ions (x=0, 1,2 and 4 mol %) were grown by conventional Czochralski technique at the first time. The In , Yb and Tm contents in the crystals were measured by an inductively coupled plasma atomic emission spectrometer (ICP-AES). The results revealed that the segregation coefficients of Yb<sup>3+</sup> and Tm<sup>3+</sup> ions decreased with increasing of the In<sup>3+</sup> doping concentration. Also the segregation coefficient of In<sup>3+</sup> ion increases with the increasing In<sup>3+</sup> doping concentration in the melts, and it is close to one. The IR transmission spectra and UV-Vis spectra were measured to analyze the defect structure of the crystals.

Keyword: In: Yb: Tm: LiNbO3 , crystal growth, dopant concentration, spectroscopic properties

# 1. Introduction

Tm<sup>3+</sup>-based lasers emitting in the wavelength region of 2  $\mu$ m have attracted much attention in recent years such as in atmospheric, telecommunications medicinal, and space applications fields [1-3]. The line width of Tm<sup>3+</sup> ions (associated with  ${}^{3}F_{4} \rightarrow {}^{3}H_{6}$  transitions) corresponds to the wavelength range of 2  $\mu$ m, which is similar to the Ho<sup>3+</sup> ions(associated with  ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$  transitions), but the absorption range is wide and the absorption cross section is large. LiNbO<sub>3</sub> crystal is an excellent material exhibiting

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desirable characteristics such as electro-optic, acousto-optic, piezo-electric and nonlinear optical performances [4, 5]. From a materials point of view, several RE-doped LiNbO<sub>3</sub> have been under study in the form of solid state laser devices and planar waveguide lasers. Yb<sup>3+</sup> ion co-doped crystals have received considerable attention due to its high absorption cross-section in the near- infrared, suitable for convenient excitation with commercially available optical sources. Recently, many efforts have been devoted to the upconversion luminescence property in RE-doped LiNbO<sub>3</sub> crystal. The basic optical properties of Tm: LiNbO<sub>3</sub> and Yb: LiNbO<sub>3</sub> crystals have been already investigated [6,7] and are considered to be excellent upconversion media. Unfortunately, to our best knowledge, the properties of Yb:Tm: LiNbO<sub>3</sub> crystal is still to be scrutinized. In recent years, 2 µm lasers have attracted a large amount of interests due to their wide tenability and high efficiency ,which have increasingly important application in remote sensing, lidar and material processing , etc. It is believed that the study of Yb:Tm: LiNbO<sub>3</sub> crystals will be attractive researches.

The light induced refractive index change, called photorefractive effect, limits the performance of Yb: Tm: LiNbO<sub>3</sub> crystals. Until now, the comprehensive studies on the suppression of the optical damage caused by the high laser intensity in LiNbO<sub>3</sub> crystal have been investigated. As an effective solution to this problem, it has been found that the significant resistance to optical damage in LiNbO<sub>3</sub> can be improved by doping with certain divalent , trivalent or tetravalent cations (such as  $Mg^{2+}$ ,  $Zn^{2+}$ ,  $In^{3+}$ ,  $Sc^{3+}$ ,  $Hf^{4+}$  and  $Zr^{4+}$ ) above the threshold concentration, see Ref. [8-13].

In this Letter, In: Yb: Tm: LiNbO<sub>3</sub> crystals with different concentrations of  $In_2O_3$  (0, 1, 2 and 4mol%) were successfully grown by the Czochralski method. The defect structures were analyzed via IR transmission spectrum and the UV-Vis absorption spectra. Finally, the mechanism of the relationship between the spectroscopic properties and the concentration of  $In_2O_3$  was thoroughly discussed.

# 2. Experiments

In this experiment, congruent LiNbO<sub>3</sub> crystals codoped with 0.5 mol. %Yb<sub>2</sub>O<sub>3</sub>, 0.5 mol. % Tm<sub>2</sub>O<sub>3</sub>, and various concentration of In<sub>2</sub>O<sub>3</sub> were grown. The composition of melt was chosen as Li/Nb = 48.6/51.4. The In<sub>2</sub>O<sub>3</sub> concentrations were 0, 1, 2 and 4mol%, respectively, which were labeled as 1#, 2#, 3# and 4#, as are shown in table 1. All raw materials precisely weighed and thoroughly mixed for 24h, heated up to 750°C for 2h to remove CO<sub>2</sub>, and then further sintered at 1050°C for 2h. The crystals were grown in open atmosphere by traditional Czochralski method along ferroelectric z-axis with the rotation rate of 15~20 r/min and pulling speed of 1~1.5mm/h. After growth, the crystals were cooled down to room temperature at a rate of 60 °C/h. The crystals were polarized at 1200°C with a current density of 5 mA/cm<sup>2</sup>. The wafers with diameter of 10×20×2mm (x× z× y) were cut from the middle part along y-face and polished to optical grade both of sides for optical tests. Fig. 1 showed the wafers of four In:Yb:Tm:LiNbO3 samples in experiment.

#### Table 1

Composition of In: Yb: Tm: LiNbO3 crystal and the size of crystal.

Sample	1#	2#	3#	4#
In <sup>3+</sup> in the crystal (mol%)	-	0.762	1.654	3.956
$K_{eff}$ of $\ln^{3+}$	-	0.762	0.827	0.989
Yb <sup>3+</sup> in the crystal (mol%)	0.848	0.793	0.754	0.712
$K_{e\!f\!f}$ of Yb <sup>3+</sup>	0.848	0.793	0.754	0.712
Tm <sup>3+</sup> in the crystal (mol%)	0.877	0.849	0.831	0.816
$K_{eff}$ of $\mathrm{Tm}^{3+}$	0.877	0.849	0.831	0.816
Crystal size(mm <sup>2</sup> )	Φ25×28	Ф25×38	Φ25×35	Ф25×25
Axial temperature gradient(°C/cm)	25-30	30-45	30-35	30-35
Radial temperature gradient (°C/cm)	4-6	4-6	3-5	3-5

We analyzed the concentration of  $In^{3+}$ ,  $Yb^{3+}$  and  $Tm^{3+}$  in the crystals with an inductively coupled plasma atomic emission spectrometer (ICP-AES, Optima 5300DV). The X-ray powder diffraction patterns were recorded using a SHMADZU XRD-6000 diffractometer with monochromatic Cu K $\alpha$ 

radiation. The  $OH^-$  absorption spectra of In: Yb: Tm: LiNbO<sub>3</sub> crystals were recorded by an Avator-360 FT-IR spectrometer at room temperature. The scanning was carried out between 400 and 4000 cm<sup>-1</sup> with the scanning step of 1 cm<sup>-1</sup>. In our later discussion, 3400 to 3600 cm<sup>-1</sup> was chosen as the effective range. The ultraviolet-visible (UV-Vis) absorption spectra of the samples were measured with a CARY UV-Visible spectrometer, with the incident light transmitting along the y axis at room temperature. The measurement range was from 300 to 1100 nm.



(a) 1# sample (b) 2# sample (c) 3# sample (d) 4# sample

Fig.1. Wafers of four In:Yb:Tm:LiNbO3 crystals in experiment

# 3. Results and discussion

Preparation of standards solution: 99.99% purity  $In_2O_3$ ,  $Yb_2O_3$  and  $Tm_2O_3$  were put in Polytetrafluoroethylene (PTFE) beaker, plus nitric acid and hydrofluoric acid solution drops, and then dissolved by heating to make-up 40g / L stock solution and saved in PTFE bottle. From each top, middle and lower parts of the crystal, cut out a small crystal specimens and placed in an agate mortar grind into fine powder, with the balance millionth 0.05g powder sample is weighed out. Powder was placed in a PTFE beaker, added hydrofluoric acid and nitric acid solution each 5mL. After preheated for 25 minutes on a heated plate, using fully automated microwave digestion oven to microwave sample digestion. Then, the sample was heated and titrated with perchloric acid to eliminate hydrofluoric acid and protect the testing equipment from corrosion. At the same time, the digestive process was finished. After cooling to room temperature, the sample solution was mixed with 5% of

nitric acid and played in the 50 ml flask for testing, the volume of the solution was set up by addition of 5% of nitric acid. Fig. 1 showed the photos of the In:Yb:Tm:LiNbO<sub>3</sub> crystals samples.

The concentrations of  $\ln^{3+}$ ,  $Yb^{3+}$  and  $Tm^{3+}$  ions in the In: Yb: Tm: LiNbO<sub>3</sub> crystals and effective distribution coefficient ( $K_{eff}$ ) of ions are listed in Table 1. The effective distribution coefficient of  $\ln^{3+}$ ,  $Yb^{3+}$  and  $Tm^{3+}$  ions are calculated by comparing the measured  $\ln^{3+}$ ,  $Yb^{3+}$  and  $Tm^{3+}$ ions concentration in the crystal and melt. The dependence of the segregation coefficients of  $\ln^{3+}$ ,  $Yb^{3+}$ and  $Tm^{3+}$  ions on the  $\ln_2O_3$  concentration in the melts is shown in Fig. 2. It is clearly seen that the effective distribution coefficient of  $\ln^{3+}$  ion increases with the increasing  $\ln_2O_3$  concentration in the melts, with the maximum close to one, indicating that it distributes uniformly in crystal. It can be further seen from Fig. 2 that the effective distribution coefficient of  $Yb^{3+}$  and  $Tm^{3+}$  ions decreases as the doping concentration of  $\ln^{3+}$  increases in the melts. When  $\ln^{3+}$  ion entered in Yb: Tm: LiNbO<sub>3</sub> crystal, it could readily occupy intrinsic defect ( $Nb_{Li}^{4+}$ ) sites and the  $Nb_{Li}^{4+}$  concentration was sharply reduced, which suppressed the entrance of  $Yb^{3+}$  and  $Tm^{3+}$  ions into crystal. Thus, the effective distribution coefficient of ions  $Yb^{3+}$  and  $Tm^{3+}$  was continuously reduced.



Fig. 2. The dependence of the segregation coefficient of In, Yb, and Tm ions on the In<sub>2</sub>O<sub>3</sub> concentration in the melts.

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Fig. 3. X-ray powder diffraction patterns of In: Yb: Tm: LiNbO<sub>3</sub> crystals.

The XRD experimental results are shown in Fig 3. It can be seen that the dope of  $In^{3+}$ ,  $Yb^{3+}$  and  $Tm^{3+}$  ions did not create new diffraction peak. This verified that after adding the  $In^{3+}$ ,  $Yb^{3+}$  and  $Tm^{3+}$  ions to crystal lattice of LiNbO<sub>3</sub> , the structure did not change dramatically [14]. Based on the ion radius approximation principle in the crystal field theory, the doped ion entered crystal lattice by replacing the  $Li^+$  and  $Nb^{5+}$  ion instead of occupied the slot among crystal lattice. Hence to obtain the formulated In: Yb: Tm: LiNbO<sub>3</sub> crystals solid solution as a single phase.

In course of growth,  $H_2O$  vapors can be introduced into the crystals through formation of O-H bonds. The OH<sup>-</sup> absorption spectra measurement results of In: Yb: Tm: LiNbO<sub>3</sub> crystals are represented in Fig. 4. Fig. 4 summarizes the OH<sup>-</sup> absorption peak of samples 1#, 2#, 3# and 4# are located at 3482, 3482, 3507 and 3508 cm<sup>-1</sup>, respectively. It can be found that the spectra belonging to In<sup>3+</sup> of different concentration form different groups in In: Yb: Tm: LiNbO<sub>3</sub> crystals. According to Li-site vacancy model, in pure congruent LiNbO<sub>3</sub> crystal, the ratio of [Li]/[Nb] is 0.946. There always exist intrinsic

lattice defects, such as anti-site Nb  $(Nb_{Li}^{4+})$  and Li vacancy  $(V_{Li}^{--})$  defects. The electropositive H<sup>+</sup> ions easily attracted by the electronegative  $V_{Li}^{--}$  and formed  $V_{Li}^{--}$  OFF complexes, whose vibrational absorption peak is located at 3482 cm<sup>-1</sup> [15]. In Yb: Tm: LiNbO<sub>3</sub> crystals (i.e. for 1# sample), Yb<sup>3+</sup> and Tm<sup>3+</sup> replace  $Nb_{Li}^{4+}$ , forming Yb<sup>2+</sup><sub>Li</sub> and Tm<sup>2+</sup><sub>Li</sub>. When the concentration of In<sup>3+</sup> ions is Imol% in the melt, below its threshold value (i.e. for 2# sample) , Yb<sup>3+</sup> and Tm<sup>3+</sup> ions substitute for  $Nb_{Li}^{4+}$  defects and exist in the form of Yb<sup>2+</sup><sub>Li</sub> and Tm<sup>2+</sup><sub>Li</sub> in LiNbO<sub>3</sub>, while these defects repel the H<sup>+</sup> and affect the  $V_{Li}^{--}$ -OFF complexes only as ions around them and the peak position is at about 3482 cm<sup>-1</sup>. When the concentration of In<sup>3+</sup> ions exceeded the threshold (i.e. for 3# and 4# samples), all  $Nb_{Li}^{4+}$  defects disappear and In<sup>2-</sup><sub>Sb</sub> defect structures appears as the In<sup>3+</sup> ions enter the normal Nb sites. Meanwhile, the amount of Yb<sup>3+</sup> and Tm<sup>3+</sup> located at Nb sites increases, namely, the amount of Yb<sup>2+</sup><sub>Nb</sub> and Tm<sup>2+</sup><sub>Nb</sub> increases. The In<sup>2-</sup><sub>Nb</sub> in the In<sub>Nb</sub>-H-O defect structure has intensive attraction to H<sup>+</sup>, more energy is required for O-H vibration, and the corresponding absorption is dominant at about 3507cm<sup>-1</sup>. According to the analysis above, we can conclude that the In doping threshold concentration is about 1.5-2.0mol% as previously reported in [16-18].



Fig.4. The IR absorption band related to the OH vibration in In:Yb:Tm:LiNbO3 crystals

To confirm our significant conclusion, we further studied the UV-Visible absorption spectra of four

samples cut from the top, center and bottom of the grown crystals. Fig. 5 shows the UV-Visible absorption spectra of the In:Yb:Tm:LiNbO<sub>3</sub> crystals with different concentration of  $In^{3+}$ . The absorption edge of LiNbO<sub>3</sub> crystals is used to determine the composition [19]. The absorption edge at the level of a=15 cm<sup>-1</sup> or a=20 cm<sup>-1</sup> is usually used to quantify the defects in a LiNbO<sub>3</sub> crystal [20]. In our experiment, the referred UV-Vis optical absorption coefficient for determination of the position defined as 20 cm<sup>-1</sup>. It can be seen from Fig.5 that the UV-Vis absorption edges of the 1#, 2#, 3# and 4# samples are located at 353, 346, 335 and 379 nm, respectively. The absorption peak of In: Yb: Tm:  $LiNbO_3$  crystals is first violet-shifted with the increasing In<sub>2</sub>O<sub>3</sub> doping concentration in the melt, then red-shifted when the doping concentration of  $In^{3+}$  is 4 mol%. The shift of absorption edge (4# sample) is up to 26 nm with respect to position of the Yb:Tm: LiNbO3crystal. It is a remarkable fact that when  $In^{3+}$  doping concentration reaches 4.0 mol%, namely 4# sample, the absorption edge of shifts to the red substantially relative to 3#. There is some absorption band located at 300-1050nm, which is attributed to absorption of the Tm<sup>3+</sup>and Yb<sup>3+</sup>ions. The three absorption bands of Tm<sup>3+</sup> ions centered at 465nm, 692 and 798nm are assigned to the  ${}^{3}H_{6} \rightarrow {}^{1}G_{4}$ ,  ${}^{3}H_{6} \rightarrow {}^{3}F_{23}$  and  ${}^{3}H_{6} \rightarrow {}^{3}H_{4}$  transitions [21]. And 980nm absorption peak correspond to transition from the  ${}^{2}F_{7/2}$  ground state of Yb<sup>3+</sup> ion to the excited state  ${}^{2}F_{5/2}$  .



Fig.5. The UV-Vis absorption spectra of the In:Yb:Tm:LiNbO<sub>3</sub> crystals.

It's common knowledge that the shift of absorption edge can be explained by the polarization ability of the doped ions [22, 23]. The valence electron transition energy, which is from the 2p orbits of  $O^{2-}$  to the 4d orbits of Nb<sup>5+</sup>, determines the fundamental absorption edge of the LiNbO<sub>3</sub> crystal [23]. It can be concluded from the previous researches that the violet shift of the absorption edge is attributed to the increasing in the ability of adjacent dopant ions to polarize an  $O^{2-}$  ion shortens the band gap of the valence electron transition. Otherwise, the absorption edge of sample shifts toward the infrared band. In fact, the polarization ability of Nb<sup>5+</sup> is smaller than that of In<sup>3+</sup>, Yb<sup>3+</sup>, Tm<sup>3+</sup> and Li<sup>+</sup>. In Yb:Tm:LiNbO<sub>3</sub> crystal, Yb<sup>3+</sup> ions and Tm<sup>3+</sup> ions replace  $Nb_{Li}^{4+}$  and form Yb<sup>2+</sup><sub>Li</sub> and Tm<sup>2+</sup><sub>Li</sub>. When the In<sup>3+</sup> enter Yb: Tm: LiNbO<sub>3</sub> crystals, because of the polarization ability of In<sup>3+</sup> ions is much smaller than that of  $Nb_{Li}^{4+}$ , the absorption edge of In:Yb:Tm:LiNbO<sub>3</sub> crystal (2# sample ) shifts to the ultraviolet band compared with that of Yb:Tm:LiNbO<sub>3</sub> crystal (1# sample ). When the concentration of In<sup>3+</sup> ions in the melt exceeded the threshold value (3# sample ),  $Nb_{Li}^{4+}$  defects almost disappeared, the In<sup>3+</sup> ions began to enter the normal Li and a little bit of Nb sites. As the polarization ability of In<sup>3+</sup> was lower than that of Li<sup>+</sup>, which caused the deformation of O<sup>2-</sup> ions electron clouds, the width of the forbidden band

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increases and absorption edge continuous shifts to the violet. When the  $In_2O_3$  doping concentration in the melts is up to 4.0mol% (4# sample ), part of  $In^{3+}$ , Yb<sup>3+</sup> and Tm<sup>3+</sup> ions begin occupies Nb sites form  $In_{Nb}^{2-}$ , Yb<sup>2-</sup><sub>Nb</sub> and Tm<sup>2-</sup><sub>Nb</sub>. As the polarization ability of  $In^{3+}$  ion is smaller than that of Nb<sup>5+</sup>, red shift of the absorption edge in 4# relative to 1#. This discussion of dopant occupancy is consistent with the analysis result of experimental shown in OH<sup>-</sup> absorption spectra.

# 4 .Conclusions

In conclusion, In: Yb: Tm: LiNbO<sub>3</sub> crystals with different concentrations of  $In_2O_3$  (0, 1, 2 and 4 mol%) were successfully grown by the Czochralski method. The inductively coupled plasma optical emission spectrometry, IR transmission spectrum and the UV-Vis absorption spectra were measured and discussed in terms of the spectroscopic properties in  $In^{3+}$  doped (0, 1, 2 and 4 mol%) Yb: Tm: LiNbO<sub>3</sub> crystal. The effective distribution coefficient of In ion increases with the increasing  $In_2O_3$  concentration in the melts, and it is close to one, indicating that  $In^{3+}$  distributes uniformly in crystal. The UV-Vis absorption spectra showed that the absorption edges of In:Yb:Tm: LiNbO<sub>3</sub> crystals is first violet-shifted with the increasing  $In_2O_3$  doping concentration in the melt , then red-shifted when the doping concentration of  $In^{3+}$  is 4 mol%. It was suggested that for low concentration of  $In^{3+}$  ion, the all dopant ions first occupy the Nb anti-sites or Li vacancies and when the concentration of  $In^{3+}$  was reaches 4.0 mol% they may replace normal Nb sites.

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