Nanoscale

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about *Accepted Manuscripts* in the **Information for Authors**.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.



www.rsc.org/nanoscale

Journal Name

ARTICLE

Received 00th January 20xx, Accepted 00th January 20xx

DOI: 10.1039/x0xx00000x

www.rsc.org/



Wei Sun, ^a Chenxi Qian, ^a Xiao Sherri Cui, ^b Liwei Wang, ^{a,c} Muan Wei, ^{a,d} Gilberto Casillas, ^e Amr S. Helmy, ^b and Geoffrey A. Ozin*^a

While silicon nanocrystals (ncSi) embedded in silicon dioxide thin films have been intensively studied in physics, the potential of batch synthesis of silicon nanocrystals from the solid-state disproportionation of SiO powder has not drawn much attention in chemistry. Herein we describe some remarkable effects observed in the diffraction, microscopy and spectroscopy of SiO powder upon thermal processing in the temperature range 850-1100 °C. Quantum confinement effects and structural changes of the material related to the size of the silicon nanocrystals nucleated and grown in this way were established by Photoluminescence (PL), Raman, FTIR and UV-Visible spectroscopy, PXRD and STEM, pinpointing that the most significant disproportionation transformations happened in the temperature range between 900 and 950 °C. With this know-how a high yield synthesis was developed that produced polydispersions of decyl-capped, hexane-soluble silicon nanocrystals predominantly with near infrared (NIR) PL. Using size-selective precipitation, these polydispersions were separated into monodisperse fractions, which allowed their PL absolute guantum yield (AQY) to be studied as a function of silicon nanocrystal size. This investigation yielded volcano-shaped plots for the AQY confirming the most efficient PL wavelength for ncSi to be located at around 820-830 nm, which corresponded to a size of 3.5-4.0 nm. This work provides opportunities for applications of size-selected near infrared emitting silicon nanocrystals in biomedical photothermal imaging and therapy.

1. Introduction

Among the variety of semiconductor quantum dots targeted for applications in optoelectronics,¹ biological imaging² and chemical sensing,³ silicon nanocrystals (ncSi) are drawing increasing attention because of their earth-abundance, low cost, and purported lack of toxicity.^{4,5} Also improvements in luminescent quantum yield of ncSi through surface-capping control,⁶⁻⁸ makes them interesting for quantum dot applications.

The synthetic strategies for making ncSi have been recently

summarized by Mangolini.⁹ The groups of Ozin,⁸ Korgel¹⁰ and Veinot (the pioneering group)¹¹ have favoured batch syntheses that involve thermally treating hydridosilicate solids with formulae $(HSiO_{1.5})_8$ or $(HSiO_{1.5})_n$ in which silicon nanocrystals nucleate and grow within a silica matrix. Extraction with aqueous HF yields ncSi:H that can subsequently be surface functionalized by hydrosilylation with a terminal olefin to produce colloidally-stable, visibly-emitting ncSi:R. For many applications in biological imaging, near infrared (NIR) emission (from about 800 nm to 2500 nm) is necessary to avoid overlap with biological background autofluorescence.¹² Such ncSi were also demonstrated to exhibit a more effective photothermal effect than the visible-emitting analogues, which enhances the potential application of an effective photothermal therapy in conjunction with photoluminescence (PL) imaging.^{13, 14} Unfortunately, production of near infrared emitting ncSi by the above methods usually requires very high temperatures above 1100 °C. For example, Henderson et al reported the heat treatment of (HSiO_{1.5})n sol-gel glass at 1200 °C for 1h followed by annealing at 1100 °C for another 10 h resulted in ncSi with a PL peak maximum at 955 nm. The maximum percentage mass yield of ncSi by this method was 10% after liberation from the silicon oxide matrix,¹⁵ which is considerable since the mass percent composition of Si in the precursor before heating is only slightly above 50%, and there is additional loss of SiH₄ cogenerated during heating. Evidently, precursors with higher Si



^{a.} Department of Chemistry, University of Toronto, 80 St. George Street, Toronto, Ontario, M5S 3H6. E-mail: gozin@chem.utoronto.ca

^{b.} The Edward S. Rogers Sr. Department of Electrical and Computer Engineering, University of Toronto. 10 King's College Road, Toronto, Ontario, M5S 3G4, Canada

^{c.} School of Marine Sciences, Guangxi University, Nanning, 530004, China ^{d.} Department of Materials Science and Engineering, University of Toronto, 184

College Street, Toronto, Ontario, M5S 3E4, Canada

^{e.} Electron Microscopy Centre, University of Wollongong, New South Wales 2500, Australia

 $^{^{\}dagger}$ Electronic Supplementary Information (ESI) available: [Experimentals, Technicians' affiliations and instrumental descriptions for STEM imaging, XRD and PL comparison for different heating time at 900 $^\circ C$, a photo of the SiO powders, Raman spectra fitting method, STEM image for stacking fault]. See DOI: 10.1039/x0xx00000x

ARTICLE

content may be advantageous for obtaining higher yields of ncSi. Indeed, silicon nanoclusters formed in silicon monoxide (SiO) films¹⁶⁻¹⁹ and other silicon sub-oxide (SiO_x, x<2) films,²⁰⁻²³ usually prepared via thermal, plasma, and laser deposition, have been of interest in optical physics for applications as protective layers^{24, 25} and in silicon based optoelectronics.²⁶ Notably, SiO is also commercially available as powders in bulk quantities having a structure and properties that appears to be different from thin films. It has however received very little attention for the batch synthesis of freestanding ncSi.¹⁸ In this context, Liu et al. treated SiO powder at 1000 °C in Ar and obtained freestanding ncSi with visible PL.²⁷ Dorofeev et al. expanded upon this work and reported a maximum PL QY of 12% for ncSi.²⁸ Recently we successfully synthesized ncSi from SiO at a specific heating temperature of 900 °C, the size distribution of which enabled a study on the size dependence of the absolute quantum yield (AQY) on passing from bulk to quantum confined silicon.²⁹ However, much detailed information about the chemistry during such synthesis was left to be investigated.

Herein we describe an in-depth diffraction, microscopy and spectroscopy study that probes the effect of heat treatment of bulk quantities of SiO powder in the temperature range 850-1100 °C. This insight enabled optimization of a high yield scalable synthesis for decyl-capped, hexane-soluble silicon nanocrystal (ncSi:C₁₀) polydispersions that predominantly display NIR PL. Size-selective precipitation of these polydisperse ncSi:C₁₀ samples yielded monodisperse fractions, the study of which provided a unique insight into the optimum PL efficiency versus emission wavelength, as well as guidance on the tuning of size distributions through the choice of the synthetic conditions to target a specific range of PL wavelengths.

2. Results and discussions

2.1 Dependence of ncSi formation on thermal treatment temperature

The commercial SiO powder precursor used by Liu²⁷ and Dorofeev²⁸ were claimed to be amorphous,³⁰ however, their powder X-ray diffraction (PXRD) results clearly show a crystalline silicon phase prior to any thermal treatment or at very low temperatures (200 °C). It is conceivable in this earlier work that the size and yield of the resultant ncSi after thermal annealing would sensitively depend on the presence of the initial ncSi, which could act as nucleation sites as well as possibly influence the thermal disproportionation. In contrast, the SiO powder we used in this study were amorphous from the beginning, where only two broad humps centred at ~23 and ~51 degrees could be seen in the PXRD pattern of the commercial SiO precursor.²⁹ This ensures that the production of ncSi mainly originates from the thermally induced disproportionation reaction shown in stoichiometric form below:





Before heating, the SiO might comprise a matrix of small clusters of Si and SiO₂³¹ with insufficient crystalline order or size to be detected by PXRD. When thermally processed at 875 $^{\circ}\mathrm{C}$ in an atmosphere of 5% H₂ in Ar (the optimal protective gas that can also enhance crystal quality^{32, 33}) and kept at this temperature for 1h, only

Fig. 1 Representative Powder X-ray diffraction patterns of SiO powder processed at 900 °C and 920 °C.

one new broad diffraction peak at ~28 degrees emerged, corresponding to the Si (111) diffraction, indicating the genesis of embryonic ncSi forming inside the matrix. Above 900 °C the (111), (220) and (311) diffraction peaks grow larger and sharper, which continues as the temperature was elevated, and are sensitive to even a small change at this threshold temperature. As shown in Fig. 1, a 20 °C increase of processing temperature even with a shorter processing time led to an observable increase of the (111) peak relative to the amorphous hump adjacent to it. The detailed growth of ncSi was quantified by full profile fitting the (111) peaks to obtain the temperature dependence of the unit cell size of the ncSi and by applying the Scherrer equation to the peak broadening to obtain an estimate of the ncSi size.²⁹ While there could be other sources of peak broadening for example, associated with dislocations, stacking faults and twinning,³⁴ there is no doubt higher temperatures favour the formation of larger ncSi with increasing unit cell dimensions that converge towards the bulk Si value of 5.4307Å.

During the thermal treatment, the broad PXRD peak around 51 degrees gradually diminished and is most likely associated with amorphous Si clusters or a SiO_x interface between domains of ncSi and SiO₂.¹⁸ Note that with increasing reaction time further coarsening³⁵ or "pseudo-Ostwald ripening"²³ of ncSi is not rare. We observe that longer heating times at the same temperature result in larger size ncSi, revealed by sharper peaks in PXRD and a red-shifted peak in PL spectra (see ESI⁺).

As the broad diffraction peak at ~21 degrees seemed slightly shifted to lower angles as the heating temperature was elevated,²⁹ approaching the position for amorphous SiO₂, the redistribution of the structure towards ncSi in SiO₂ matrix is conceivable. So we investigated the characteristic peak

Journal Name



Jou



assigned to Si-O-Si in the IR spectra which reflects the bond angle information.

As can be seen in Fig. 2 as the temperature increases, the Si-O-Si stretching mode appeared sharper and the peak centre shifted to higher wave numbers, from 1028 cm⁻¹ for the unheated sample to 1098 cm⁻¹ for a sample thermally processed at 1100 °C. The position of the peak has been correlated with the angle 2ϑ between the two Si atoms separated by a bridging O atom in Fitch's work: ³⁶

Fig. 2 ATR-FTIR spectra of SiO thermally processed at different temperatures.

 $v = v_0 \sin \vartheta$

where v_0 was 1116.5 cm⁻¹ when ϑ was 75° for a SiO₂ grown at 1150 °C in this work. The blue shifting and sharpening of the Si-O-Si stretching frequency observed during the disproportionation SiO \rightarrow Si + SiO₂ reaction indicates the formation of a more well-defined SiO₂ phase^{20, 37} from a poorly ordered SiO phase with a range of Si-O-Si bond angles.³⁸ With the change of the structure, we also observed an accompanying colour change of the matrix from reddishbrown to dark brown (see photo in ESI⁺), which is manifest as changes in the reflectance spectra especially at longer wavelengths that begins above 850 °C, as shown in Fig. 3.

Another informative technique for probing the formation of ncSi during the thermal induced disproportionation of SiO is Raman spectroscopy.³⁹ All the thermally treated SiO powder samples have multiple broad peaks, shown in Fig. 4. The detailed assignment^{33, 40} of Raman peaks is shown in Fig. 5a: the longitudinal acoustic (LA) mode centered at ~300 cm⁻¹, the longitudinal optical (LO) mode at ~425 cm⁻¹, and the strong transverse optical (TO a-Si) mode at ~475 cm⁻¹ are all from the amorphous Si contribution, while a single asymmetric transverse optical (c-TO) mode, which is related to the ncSi, is slightly downshifted from the crystalline Si wafer mode at 521 cm⁻¹, to around 516-517 cm⁻¹, indicating the effect of phonon confinement.^{41, 42} Multiple peak fitting was performed for each Raman spectrum except for that of the sample thermally processed at 850 °C, as shown in Fig. 5a, following a reported Fano model.⁴³ (see detailed fitting method in ESI⁺).

We found that the SiO powder precursor contains no detectable crystallized Si initially but rather only amorphous Si, even if thermally processed at 880 °C. When the powder was thermally processed above 900 °C, the Raman peak for crystalline Si became observable and could be well-fitted, indicating formation of significant amounts of ncSi, consistent with our PXRD result. As expected, higher temperature led to



Si(0) within the oxide matrix with a concomitantly higher crystal fraction calculated using the equation:

 $F_c = I_c / (I_c + \alpha I_a)$

Fig. 3 Diffuse reflectance spectra of SiO thermally processed at different temperatures.

where α =0.8 is the Raman cross section ratio for phonon excitation of crystalline Si (I_c) to the amorphous (I_a)

components.⁴⁴ Again we consistently found that drastic change occurred between 900 and 950 °C, and plateaued above 950 °C. Though 1h of heating did not lead to 100% yield of the crystalline ncSi phase, it nevertheless did produce significant amounts of ncSi as a precursor for the subsequent production of size-separated, colloidally-stable, organic-capped and brightly-luminescent ncSi:R with the NIR PL range of interest in this investigation as discussed below.

2.2 Size-separated colloidally stable ncSi with NIR PL

With the information derived from the various characterization techniques discussed above, it is clear that the size and yield of ncSi are sensitive to the heating temperature and time: 1h at ~900 °C is a threshold temperature for significant crystallization whereas for the same time below ~850 °C, crystallization barely occurred, and we observed no PL from the ncSi produced from SiO thermally processed at 950 °C within our detector's working wavelength range. Therefore the processing temperatures of 900, 920, and 925 °C were chosen as the standard protocols for producing ncSi. After the liberation of ncSi from the SiO₂ matrix by aqueous HF etching,⁸ turbid dispersions of ncSi:H could be obtained in common organic solvents such as toluene and aliphatic hydrocarbons. Note that this synthesis approach has several advantages over the (HSiO_{1.5})_n precursor, which was synthesized from HSiCl₃ and processed with several steps before HF etching.⁴⁵⁻⁴⁷ Table 1 lists several comparisons between the two methods. The origin of the relatively low yield of ncSi:H product for both methods we believe is during the aqueous HF extraction and transfer of the dispersion to the organic solvent, however the yield of ncSi:H from the SiO precursor is consistently higher than (HSiO_{1.5})_n due to the higher atom ratio of Si, and the experimental procedure is much easier to manage.

The ncSi:H devoid of organic ligands is not colloidally stable whereas ncSi:R with organic capping groups becomes colloidally stable. This was achieved by hydrosilylation using 1-decene as the non-polar solvent for both the extraction and as the reactant at 170 °C.⁴⁸ Following this procedure the dispersion

of



Fig. 4 Raman spectra of SiO powder thermally processed at different temperatures.

Table 1 ncSi from	(HSiO ₁₅), versus SiC	precursors
		p. c.c.a

	⊤ (°C)	Precursor mass (g)	ncSi yield (g)	% mass yield	PL (nm) ^a	Pre- processing	
(HSiO _{1.5}) _n	1050	0.60	0.015	2.5	733	Yes	
SiO	900	0.30	0.020	6.7	869	No	
^a After etching in HF solution for 1h 30 min.							

became clear.

NcSi:C₁₀ remained colloidally stable after re-dispersion in hexane. The size of a statistically meaningful population of the crystalline ncSi:C₁₀ was measured by bright-field (BF) scanning transmission electron microscopy (STEM). The mean size of ncSi:C₁₀ obtained from SiO thermally processed at 900 °C determined in this way is 3.5 nm as we reported elsewhere,²⁹ which is larger than the value calculated by Scherrer analysis of PXRD line widths (~2 nm). This discrepancy might arise from the occurrence of stacking faults and twinning observed in the larger ncSi which can contribute additional line broadening in PXRD³⁴ (see ESI⁺ for STEM image of large ncSi). The PL wavelength and the AQY of the ensemble was determined to be 834 nm and 12-14% respectively, which is consistent with a ncSi:C₁₀ size of 3.5-4 nm.⁴⁹ It is interesting that in the ensemble prepared from SiO some very small ncSi:C₁₀ still displayed crystalline order. In our previous report, using the HSiO₁₅ precursor, we were hardly able to identify crystalline order for



 ncSi:C_{10} smaller than ~2.5 nm, 50 however in this work the use

Fig. 5 (a) A representative fitting of the Raman spectrum obtained from SiO powder thermally processed at 1050 °C, where a-Si represents amorphous Si, dots are obtained data points, and the curve overlapping with the dots is the cumulative fitting curve. The inset shows the ncSi TO mode position. (b) Calculated crystalline fraction versus processing temperature during the synthesis.

a probe-corrected microscope allowed us to confirm that these small particles on a graphene TEM grid⁵¹ are also crystalline, as shown in the HAADF-STEM image in Fig. 6.

We performed size separation to several ensemble dispersions obtained from SiO processed at different temperatures using the size selective precipitation method.⁴⁵ In the top plot of Fig. 7 it can be seen from the right to the left that the PL wavelength maximum of ncSi:C₁₀ blue shifted monotonically from ~900 nm to ~740 nm, consistent with the quantum size effect (QSE) expected on decreasing the size of ncSi:C₁₀ from the largest to the smallest. Since the measurements were conducted at room temperature and the ncSi are freestanding, temperature and interface effects are excluded. Of special interest is the amonotonic trend observed for the AQY, which switches on from essentially zero for ncSi:C₁₀ at the wavelength of ~900 nm, corresponding to the size around 5 nm,²⁹ increases monotonically to reach a maximum around 820 nm and then decreases monotonically to the smallest studied approaching 725 nm, with a value again close to zero.

Journal Name

700

0.16

900 °C

750

800

850

900

Journal Name

We previously attribute this volcano like behaviour of the size dependent PL AQY of $ncSi:C_{10}$ to how the counterbalancing radiative and non-radiative contributions to



photo-excited

0.12 0.08 0.04 0.16 920 °C 0.12 AQY 0.08 0.04 0.00 925 °C 0.16 0.12 0.08 700 750 800 850 900 Peak Center Wavelength (nm)

Fig. 6 HAADF-STEM image showing ncSi:C_{10} smaller than 2.5 nm displaying clear lattice fringes

electron-hole recombination emission was affected by the size of the ncSi:C₁₀. These three regions of the volcano-shaped plot are believed to respectively define the (i) transition from bulk non-emissive silicon to quantum confined emissive silicon, (ii) increasing spatial confinement of electron-hole pairs enhancing radiative recombination, and (iii) increasing contributions of organic-capping group vibrations and surface defects favouring non-radiative recombination.^{29, 45} The volcano tops at 820-830 nm in Fig. 7 effectively pinpoints where the opposing factors of radiative recombination and non-radiative recombination balance each other. Although the ncSi were produced under different conditions, the position of the maximum in the volcano-shaped plots are very close, thus confirming the credibility of such behaviour, excluding the possibility of artefacts originated from a specific synthetic condition. Therefore, at this wavelength, the PL of ncSi is most efficient. However, the most complete profile of the volcanoshaped plot was only observed for the ncSi from the SiO processed at 900 °C. The temperature of 920 °C and 925 °C only yielded part of the volcano-shape, without showing the point where the AQY switched on. Therefore, the optimum processing temperature to produce colloidally-stable ncSi with PL spanning the whole efficient luminescence wavelength range should be exactly 900 °C. Nevertheless, if one aims at obtaining PL within the NIR range, 920 °C could be optimum, as shown in Fig. 8, in which the most concentrated fractions exhibit longer PL emission wavelength over 800 nm, corresponding with the results in Fig. 7.

Conclusions

Fig. 7 Absolute quantum yield of photoluminescence as a function of peak centre wavelength. The SiO powder was processed at 900, 920, 925 °C for 1 h, 35 min, and 30 min, and etched for 2 h 50 min, 2 h 50 min, and 2 h 30 min, respectively. The dashed column indicates the samples where the absolute quantum yield was maximum.

The thermally-induced solid-state disproportionation of commercially available SiO, according to the simplified reaction scheme 2SiO \rightarrow Si + SiO₂ provides a convenient gram scale batch synthesis of ncSi encapsulated by a SiO₂ matrix. The 900 °C processing temperature serves as the threshold temperature for ncSi nucleation and growth, with the most drastic changes of structure occurring between 900 °C and 950 °C. Above 950 °C, large ncSi with no detectable PL dominate, with properties resembling those of bulk Si. Subsequent aqueous HF etching of the SiO₂ matrix liberates ncSi:H, which can be hydrosilated with a terminal olefin such as 1-decene to form polydispersions of colloidally-stable, organic-capped brightly-photoluminescent ncSi:C₁₀. These polydispersions can be sorted by size selective precipitation methods using a solvent anti-solvent strategy into monodispersions of ncSi:C₁₀ that display characteristic QSE in their PL spectra. By controlling the temperature and time of the thermal treatment of SiO as well as the aqueous HF etching time it has proven possible to obtain monodispersions of ncSi:C₁₀ that display PL predominantly throughout the NIR wavelength range. Here it has been discovered that the size-dependence of the PL AQY displays a volcano-shape, which maximises at a wavelength of 820-830 nm. This all bodes well for the development of a scalable synthesis of size-selected near

Journal Name

ARTICLE

infrared emitting silicon nanocrystals with potential applications in biomedical imaging and photothermal therapy.

Concentrated



> 800 nm

Fig. 8 Photos of size separated $ncSi:C_{10}$ fractions under room light (upper) and under UV light (lower).

Acknowledgements

GAO is Government of Canada Research Chair in Materials Chemistry and Nanochemistry. The strong and sustained support of the Natural Sciences and Engineering Research Council of Canada is deeply appreciated. Chenxi Qian and Wei Sun are grateful to the Connaught Foundation for graduate fellowships. This research used equipment funded by the Australian Research Council (ARC) – Linkage, Infrastructure, Equipment and Facilities (LIEF) grant LE120100104 located at the UOW Electron Microscopy Centre. STEM imaging conducted by Dr. Larios-Rodriguez, E is appreciated.

References

- D. V. Talapin, J.-S. Lee, M. V. Kovalenko and E. V. Shevchenko, *Chem. Rev.*, 2010, **110**, 389-458.
- X. Michalet, F. F. Pinaud, L. A. Bentolila, J. M. Tsay, S. Doose, J. J. Li, G. Sundaresan, A. M. Wu, S. S. Gambhir and S. Weiss, *Science*, 2005, **307**, 538-544.
- 3. W. C. W. Chan and S. Nie, Science, 1998, 281, 2016-2018.
- A. N. Shubenkov, S. B. Korovin, E. R. Andreeva, L. B. Buravkova and V. I. Pustovoy, *BIOPHYSICS*, 2014, **59**, 105-109.
- 5. A. M. Derfus, W. C. W. Chan and S. N. Bhatia, *Nano Letters*, 2003, **4**, 11-18.

- R. J. Anthony, D. J. Rowe, M. Stein, J. Yang and U. Kortshagen, *Adv. Funct. Mater.*, 2011, **21**, 4042-4046.
- M. L. Mastronardi, E. J. Henderson, D. P. Puzzo, Y. Chang, Z. B. Wang, M. G. Helander, J. Jeong, N. P. Kherani, Z. Lu and G. A. Ozin, *Small*, 2012, 8, 3647-3654.
- M. L. Mastronardi, F. Maier-Flaig, D. Faulkner, E. J. Henderson, C. Kübel, U. Lemmer and G. A. Ozin, *Nano Letters*, 2011, **12**, 337-342.
- L. Mangolini, Journal of Vacuum Science & Technology B, 2013, 31, 020801.
- C. M. Hessel, D. Reid, M. G. Panthani, M. R. Rasch, B. W. Goodfellow, J. Wei, H. Fujii, V. Akhavan and B. A. Korgel, *Chemistry of Materials*, 2011, 24, 393-401.
- E. J. Henderson, J. A. Kelly and J. G. C. Veinot, *Chemistry of Materials*, 2009, **21**, 5426-5434.
- 12. A. M. Smith, M. C. Mancini and S. Nie, *Nat Nano*, 2009, **4**, 710-711.
- 13. S. Regli, J. A. Kelly, A. M. Shukaliak and J. G. C. Veinot, *The Journal of Physical Chemistry Letters*, 2012, **3**, 1793-1797.
- 14. L. Han, M. Zeman and A. H. M. Smets, *Nanoscale*, 2015, **7**, 8389-8397.
- C. M. Hessel, E. J. Henderson and J. G. C. Veinot, *Chemistry of Materials*, 2006, **18**, 6139-6146.
- 16. E. A. Mikhantiev, S. V. Usenkov and N. L. Shwartz, 2012.
- 17. K. Schulmeister and W. Mader, *Journal of Non-Crystalline Solids*, 2003, **320**, 143-150.
- A. Hohl, T. Wieder, P. A. van Aken, T. E. Weirich, G. Denninger, M. Vidal, S. Oswald, C. Deneke, J. Mayer and H. Fuess, *Journal of Non-Crystalline Solids*, 2003, **320**, 255-280.
- 19. J. Wang, X. F. Wang, Q. Li, A. Hryciw and A. Meldrum, *Philosophical Magazine*, 2007, **87**, 11-27.
- 20. T. R. Roschuk, J. Wojcik, E. A. Irving, M. Flynn and P. Mascher, 2004.
- 21. N. Saxena, A. Agarwal and D. Kanjilal, *Physica B: Condensed Matter*, 2011, **406**, 2148-2151.
- 22. N. Daldosso, G. Das, S. Larcheri, G. Mariotto, G. Dalba, L. Pavesi, A. Irrera, F. Priolo, F. Iacona and F. Rocca, J. Appl. Phys., 2007, 101, 113510.

23. D. Yu, S. Lee and G. S. Hwang, J. Appl. Phys., 2007, **102**, 084309.

- 24. C. A. Wang, R. K. Huang, D. A. Shiau, M. K. Connors, P. G. Murphy, P. W. O'Brien, A. C. Anderson, D. M. DePoy, G. Nichols and M. N. Palmisiano, *Applied Physics Letters*, 2003, 83, 1286-1288.
- 25. H. Shibayama, S. Hasuo and T. Yamaoka, *Applied Physics Letters*, 1985, **47**, 429-430.
- L. Pavesi, L. Dal Negro, C. Mazzoleni, G. Franzo and F. Priolo, Nature, 2000, 408, 440-444.
- 27. S.-M. Liu, Yang, S. Sato and K. Kimura, *Chem. Mater.*, 2006, **18**, 637-642.
- 28. S. G. Dorofeev, A. A. Ischenko, N. N. Kononov and G. V. Fetisov, *Current Applied Physics*, 2012, **12**, 718-725.
- 29. W. Sun, C. Qian, L. Wang, M. Wei, M. L. Mastronardi, G. Casillas, J. Breu and G. A. Ozin, *Adv. Mater.*, 2015, **27**, 746-749.
- 30. M. Nagamori, J. A. Boivin and A. Claveau, *Journal of Non-Crystalline Solids*, 1995, **189**, 270-276.

6 | J. Name., 2012, 00, 1-3

This journal is © The Royal Society of Chemistry 20xx

Journal Name

- A. Hohl, T. Wieder, P. A. van Aken, T. E. Weirich, G. Denninger, M. Vidal, S. Oswald, C. Deneke, J. Mayer and H. Fuess, *Journal of Non-Crystalline Solids*, 2003, **320**, 255-280.
- 32. Y. Li, P. Liang, Z. Hu, S. Guo, Q. You, J. Sun, N. Xu and J. Wu, Applied Surface Science, 2014, 300, 178-183.
- L. J. Borrero-González, L. A. O. Nunes, M. R. B. Andreeta, J. Wojcik, P. Mascher, Y. A. Pusep, D. Comedi and F. E. G. Guimarães, *J. Appl. Phys.*, 2010, **108**, 013105.
- 34. A. K. Singh and D. M. R. Laboratory, Advanced X-ray Techniques in Research and Industry, Capital Publishing Company, 2006.
- 35. S. Mukherjee and B. R. Cooper, *Materials Science and Engineering: A*, 1998, **248**, 101-114.
- 36. J. T. Fitch, G. Lucovsky, E. Kobeda and E. A. Irene, *Journal of Vacuum Science & Cherry Technology B*, 1989, **7**, 153-162.
- 37. A. Sassella, A. Borghesi, F. Corni, A. Monelli, G. Ottaviani, R. Tonini, B. Pivac, M. Bacchetta and L. Zanotti, *Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films*, 1997, **15**, 377-389.
- 38. K. AlKaabi, D. L. V. K. Prasad, P. Kroll, N. W. Ashcroft and R. Hoffmann, *Journal of the American Chemical Society*, 2014, 136, 3410-3423.
- 39. R. Anthony and U. Kortshagen, *Physical Review B*, 2009, **80**, 115407.
- 40. Y. Duan, J. F. Kong and W. Z. Shen, Journal of Raman Spectroscopy, 2012, 43, 756-760.
- 41. K. W. Adu, H. R. Gutiérrez, U. J. Kim, G. U. Sumanasekera and P. C. Eklund, *Nano Letters*, 2005, **5**, 409-414.
- 42. L. Han, M. Zeman and A. H. M. Smets, *Appl. Phys. Lett.*, 2015, **106**, 213106.
- D. M. Sagar, J. M. Atkin, P. K. B. Palomaki, N. R. Neale, J. L. Blackburn, J. C. Johnson, A. J. Nozik, M. B. Raschke and M. C. Beard, *Nano Lett.*, 2015, **15**, 1511-1516.
- L. J. Borrero-González, L. A. O. Nunes, M. R. B. Andreeta, J. Wojcik, P. Mascher, Y. A. Pusep, D. Comedi and F. E. G. Guimarães, *Journal of Applied Physics*, 2010, **108**, -.
- 45. M. L. Mastronardi, F. Maier-Flaig, D. Faulkner, E. J. Henderson, C. Kubel, U. Lemmer and G. A. Ozin, *Nano Lett.*, 2012, **12**, 337-342.
- 46. W. Sun, C. X. Qian, M. L. Mastronardi, M. Wei and G. A. Ozin, *Chem. Commun.*, 2013, **49**, 11361-11363.
- C. Qian, W. Sun, L. Wang, C. Chen, K. Liao, W. Wang, J. Jia, B. D. Hatton, G. Casillas, M. Kurylowicz, C. M. Yip, M. L. Mastronardi and G. A. Ozin, *J. Am. Chem. Soc.*, 2014, **136**, 15849-15852.
- W. Sun, C. Qian, M. L. Mastronardi, M. Wei and G. A. Ozin, Chemical Communications, 2013, 49, 11361-11363.
- W. D. A. M. de Boer, D. Timmerman, K. Dohnalova, I. N. Yassievich, H. Zhang, W. J. Buma and T. Gregorkiewicz, *Nat. Nanotechnol.*, 2010, 5, 878-884.
- M. L. Mastronardi, F. Hennrich, E. J. Henderson, F. Maier-Flaig, C. Blum, J. Reichenbach, U. Lemmer, C. Kübel, D. Wang, M. M. Kappes and G. A. Ozin, *Journal of the American Chemical Society*, 2011, **133**, 11928-11931.
- 51. M. G. Panthani, C. M. Hessel, D. Reid, G. Casillas, M. José-Yacamán and B. A. Korgel, *J. Phys. Chem. C*, 2012, **116**, 22463-22468.