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Direct observation of structural and defect evolution in C-rich SiC using in situ helium ion microscopy

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The microstructural effects of SiC swelling, mechanisms of He diffusion and
aggregation in C-rich SiC are studied using an in situ helium ion microscope. The
additive carbon interface provides improved swelling resistance in SiC to ~270 nm,
and defect formation is not observed until very high He implantation doses.

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Continuous silicon carbide (SiC) fiber reinforced SiC matrix composites 1 2 (SiC/SiC) are being considered as candidate structural materials for advanced fission reactors and future fusion reactors. Their primary advantages are their excellent 3 mechanical properties and chemical stability at elevated temperatures, while 4 maintaining low neutron activation and good radiation damage tolerance ^{1,2}. With 14 5 MeV neutrons in a fusion reactor, SiC and other ceramic materials produce higher 6 7 concentrations of transmutation gases (helium (He) and hydrogen (H)) than in ferritic steels and vanadium allovs ^{3,4}. The reported He and H production rates in SiC in the 8 first wall are approximately 130 and 40 atomic parts per million (appm), respectively, 9 for a damage level of one displacement per atom (dpa)³. It is widely recognized that 10 the limited solubility of He can enhance cavity formation in irradiated materials, 11 12 degrade structural properties, and greatly affect the safety in power devices. However, a fundamental understanding of the physical process of He interacting in materials 13 during irradiation is limited. The specific questions of interest include: What are the 14 principal mechanisms to control the diffusion and aggregation of He? How are they 15 influenced by the microstructure? How can the detrimental effects of He be 16 diminished and handled via proper material design? 17

Over the last fifty years, a kind of in-situ facility has been investigated and 18 developed to reveal the mechanisms and kinetics underlying damage production, 19 accumulation and evolution, combined with a real-time transmission electron 20 microscopy (TEM) observation with in situ ion irradiation⁵. Previous studies on He⁺ 21 22 irradiated SiC showed that He bubbles form above a specific temperature-dependent fluence and grow gradually as the implantation proceeds ^{6, 7}. Helium ion microscopes 23 24 (HIM) are recently developed scanning ion microscopes based on a gas field ion 25 source with high resolution (≤ 0.35 nm), which is close to the resolution in TEM. They

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also have low sample preparation requirements. Recent literature has given a detailed
introduction to HIM ⁸. In the present work, HIM was used as a powerful tool to
bombard the sample surface by a focused He-ion beam of 30 keV while recording
data in real-time. Extra phenol-formaldehyde (PF) resin was added to the SiC to
imitate the common carbon (C)-rich surroundings in industrial SiC composites. SiC
grain size and grain boundary effects were studied through a wide size range from
tens of nanometers to several microns.

C-rich SiC with an atomic ratio of C/Si \approx 2.3 was prepared via solid-state 8 sintering. Two β -SiC powders of 1 µm and 45~55 nm from Alfa Aesar were mixed 9 with a mass ratio of 80:20. PF resin was used to homogeneously coat SiC powders 10 using ethanol as a solvent. After low energy ball-milling at 80 revolutions per minute 11 12 (rpm) for 70 h, the mixture powder was consolidated using a sing cycle with cold uniaxial pressing of 300 MPa and sintering at 1500 °C for 6 h in the argon atmosphere. 13 The density of as-received sample was ~ 1.7 g/cm³ by the Archimedes method. Fig. 1a 14 shows two sharp peaks centered at 1357 and 1591 cm⁻¹ using Raman spectroscopy, 15 corresponding to the disorder induced D band and Raman-allowed G band in typical 16 carbonaceous materials, respectively. By analyzing the I(D)/I(G) intensity ratio, the 17 structure of pyrolytic carbon (PyC) from PF resin exhibits a similar disorder to 18 nanocrystalline-graphite (NC-graphite)⁹. The surfaces of specimens were flatted and 19 cleaned to ensure minimal contamination before being loaded into the main chamber 20 of HIM. 21

The in-situ HIM (Orion Plus, Carl Zeiss SMT) was performed at room temperature (RT) in the Environmental Molecular Sciences Laboratory (EMSL) within the Pacific Northwest National Laboratory (PNNL). An acceleration voltage of 30 kV was used to irradiate the C-rich SiC specimens. Fig. 1b shows that the

maximum irradiated depth of 30 keV He⁺ was ~570 nm using the SRIM-2008 1 simulation with threshold displacement energies of 35 eV for Si and 20 eV for C 2 atoms ¹⁰. For a fluence of 1×10^{17} He/cm², the damage values are 0.457~1.248 dpa at 3 4 8-16 nm and 4.86 dpa at 288 nm of the damage peak. In-situ observation was 5 achieved by using the internal patterning software to raster the focused He ion beam over an area of 1×1 or 2×2 µm² with an image size P of 1024×1120 pixels. A beam 6 current about 20 pA was used and the dwell time per pixel t was selected as 10 µs or 7 20 μ s to achieve the required fluence (ion/cm²) based on its definition as the particle 8 number N per area A: 9

10
$$\frac{dN}{dA} = \frac{\frac{I \times t}{ne} \times F}{A}$$

11 (1)

where *I* is the beam current, *n* is the charge number and *e* is the elementary charge of approximately 1.6×10^{-19} C. It suggests that appropriate dwell time and focused area are crucial to capture the slightest shift in surface morphology when other factors are constant.

16 SiC grains (see the convex regions) with sizes ranging from 35nm to 2µm were 17 embedded in the PyC matrix by comparing the secondary electron (SE) and Rutherford backscattered ion (RBI) modes under HIM in Supplementary Fig. 1. 18 Based on this, real-time observation of the SiC grain growth with different size at RT 19 under in-situ 30 keV He⁺ ion irradiation is shown in Fig. 2a to 2f. The focused area 20 was $2 \times 2 \text{ }\mu\text{m}^2$ and the dwell time was 10µs. A movie revealing the dynamic growth of 21 SiC grains is shown in Supplemetary Video 1 to a total fluence of 5.605×10^{18} He⁺/cm². 22 Three SiC grains, A(>450 nm, in red contours), B(~270 nm, in blue contours) and 23 $C(\sim 65 \text{ nm}, \text{ in green contours})$ are selected as our subjects in this study. From 24

 1.48×10^{17} to 2.22×10^{18} He⁺/cm², the large grain A presents a typical swelling due to 1 2 the fact that interstitial defects quickly move to the surface, while small grains B and C undergo no swelling or slowly shrinking under RT irradiation because the 3 increasing proportion of grain boundary captures defects and re-emits them to 4 annihilate with vacancies, showing an improved radiation resistance. F. Gao¹¹ and W. 5 Jiang ¹² have both found an "interface-driven-shrinking" in nanocrystalline SiC via 6 7 theoretical and experimental approaches. However, in the previous study, the effects of structure and chemistry of the interface are overlooked. Moreover, there is a 8 significant improvement in the critical size for interface-driven-shrinking by adding C 9 into SiC, from 12 nm (simulated value ¹¹) or 3.8 nm (measured values ¹²) to 270 nm. 10 To examine the growth rate in large SiC grains and the influence of diffusion and 11

12 aggregation of He in more detail, we adjusted the dwell time to 20 μ s and the focus area to $1 \times 1 \mu m^2$. Fig. 3a-3h shows a series of in-situ HIM images containing two SiC 13 microparticles and a C interface in between. The corresponding dynamic process at 14 RT under 30 keV He⁺ ion irradiation with a total fluence of 8.73×10^{18} He⁺/cm² is 15 given in Supplemetary Video 2. Grain growth occurs after necks are formed, thus the 16 grain growth rate in SiC can be obtained from the measurement of the neck length (L) 17 between the two SiC particles. The neck length (illustrated with the green dotted line) 18 presents continuous growth in the first five frames of Video 2, from 0 to 1.13×10^{18} 19 He^{+}/cm^{2} (the neck will be out of the field of view from the sixth frame). And the 20 corresponding swelling in SiC microparticles appears to be linear in Fig. 4a, with a 21 growth rate of 11.37 %. At present, we cannot characterize the amorphous 22 23 transformation using the in-situ HIM, but the above growth rate of SiC under He^+ ion irradiation is consistent with previous reports of 10.8 % in neutron-amorphized SiC 24 under 343 K irradiation ¹³. The implanted He tends to move to the sample surface for 25

extremely low solubility in C-rich SiC, and a small cavity of ~6.5 nm first appears at 1 the C interface at a fluence of 3.10×10^{18} He⁺/ cm². As the irradiation proceeds, the 2 small cavity becomes spherical in shape, such as defects A and B (circled in blue), and 3 gradually grows to the maximum diameter as He gas releases from the cavity when it 4 5 reaches saturation. The shrinkage leads to a sudden increase in the cavity mobility at the C interface. According to an earlier report 14 , we attribute this type of defect, 6 7 which moves freely on the surface of the C-rich SiC, to He interstitial bonds (He atoms on interstitial sites, type I) that diffuse quickly even below room temperature. 8 The cavity size distribution was analyzed by Nano Measurer software. For type I 9 10 cavity which gas-filled (a "bubble"), the maximum dimension was less than 40 nm (Fig.4b). Then He cavities move to the middle and coalesce into very large cavities 11 12 (defects 1 and 2, circled in red) in the C interface with increasing radius from 18 to 80 nm at fluences from 6.47×10^{18} to 8.73×10^{18} He⁺/cm². Since this type of defect is 13 immobile and without a maximum dimension (type II), we believe that 14 "bubble-to-void (cavities without gas)" transitions happen in the process. The 15 subsequent implanted He interstitials are trapped by the voids to make the cavities 16 even larger. According to the latest study by I.J. Beyerlein, et al.¹⁵, this 17 expansion-shrinkage in He cavities is a new morphological change when there exists 18 19 grain boundaries or interfaces, and is driven by a competition between three kinds of pressures acting on the cavity. In the equilibrium condition, these pressures are: 20

$$P_{He} + P_V = P_c \tag{2}$$

where P_{He} is the mechanical pressure of the trapped He gas, P_{V} is the osmotic pressure due to the flux of radiation-induced vacancies within the crystal to the cavity, and P_{c} is the capillary pressure arising from the surface energy of the cavity. P_{He} and P_{V} tend to expend the cavity while P_{c} tends to shrink it. Under RT irradiation, the flux

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1 of radiation-induced vacancies is very low because vacancies are practically immobile, 2 which we hypothesize that Pv remains the same and could be ignored during the entire test. In the first stage, small cavities trap He gas to cause $P_{\text{He}} > P_{\text{c}}$, and expansion 3 occurs. Carbon reconstructs with He ion beams due to the continuous loss of atoms¹⁶, 4 leading to an increase of surface energy. Thus, in the second stage, $P_{\rm He} < P_{\rm c}$, and 5 6 shrinkage occurs. In the third stage, P_c stops increasing when C possesses the highest 7 disorder since there is no more reconstruction. Meanwhile, the $P_{\rm He}$ in large cavities gradually increase as more and more He interstitials are trapped, $P_{\text{He}} > P_{\text{c}}$, and 8 expansion occurs again. Fig. 4b shows the evolutions of defect diameters and defect 9 10 number with fluences. The migration and coalescence processes, as well as the 11 aforementioned expansion-shrinkage in He cavities, are illustrated in Fig. 5. The small cavity ~8 nm within SiC grain is first observed at a fluence of $7.88{\times}10^{18}~\text{He}^{+}/~\text{cm}^2,$ 12 which is much higher than that of the single crystal SiC (usually $\sim 1 \times 10^{17}$ He⁺/ cm²)⁶. 13 ¹⁷. Combined with the aforementioned result in C interface $(3.10 \times 10^{18} \text{ He}^+/\text{ cm}^2)$, the 14 threshold fluence for defect formation has greatly increased in C-rich SiC, which 15 indicates that the radiation resistance can be improved by an additive C interface. 16

17 Conclusion

The in situ HIM observation of irradiation induced structural and defect 18 evolutions in C-rich SiC has been performed for the first time. The 19 high resolution of HIM offers the intriguing possibility to detect small cavities as 20 small as 6.5 nm, and accurately measure the swelling rate in SiC grains. Grain 21 boundaries are effective sinks for defects. The NC-graphite-like C interface in C-rich 22 23 SiC effectively increases the swelling resistance in large SiC grains (<270nm). It also 24 delays the emergence of He cavities and controls the number of cavities even at high 25 He doses. Therefore, tailoring the interface with a C phase may offer a promising

- 1 approach in SiC composite design for radiation resistance, and meet the demand for
- 2 next-generation nuclear reactor applications.
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Figures



Fig. 1 (a) Raman spectrum of the prepared C-rich SiC using 100 mW, 488 nm excitation with a spot size of 1 μ m, and (b) SRIM predicted collision(left) and damage profile(right) at XY longitudinal for the sample irradiated by 30 keV He⁺ to a dose of 1×10¹⁷ /cm².



Fig. 2 Selected HIM images showing SiC grain shape evolutions to different doses: (a)
0, (b) 1.48×10¹⁷ He⁺/cm², (c) 4.44×10¹⁷ He⁺/cm², (d) 8.88×10¹⁷ He⁺/cm², (e)
1.48×10¹⁸ He⁺/cm² and (f) 2.22×10¹⁸ He⁺/cm². Three grains A (~450 nm), B
(~270 nm) and C (~65 nm) use solid lines for original samples and dashed lines for irradiated samples.



Fig. 3 He post-irradiation on Au-RT-irradiated sample in HIM focus on carbon interphase between two SiC micron-grains to different doses: (a) 0, (b) 8.45×10^{17} He⁺/cm², (c) 3.10×10^{18} He⁺/cm², (d) 4.22×10^{18} He⁺/cm², (e) 5.07×10^{18} He⁺/cm², (f) 6.19×10^{18} He⁺/cm², (g) 7.32×10^{18} He⁺/cm² and (h) 8.73×10^{18} He⁺/cm².



Fig. 4 (a) Expansion of micron-SiC under in-situ He irradiation, (b) Defect diameters and number with increasing fluence.



Fig. 5 Schematic illustrations of the helium cavities evolutions in C interface between

two SiC grains under in-situ HIM.