# Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about *Accepted Manuscripts* in the **Information for Authors**.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.



www.rsc.org/nanoscale

| 1  | Insights from Investigations of Tin Dioxide and Their Composites: Electron-                                   |
|----|---|
| 2  | Beam Irradiation, Fractal Assessment, and Mechanism   |
| 3  |   |
| 4  | Zhiwen Chen, <sup>*a,b</sup> Chan-Hung Shek, <sup>b</sup> C. M. Lawrence Wu <sup>*b</sup>                     |
| 5  | <sup>a</sup> School of Environmental and Chemical Engineering, Shanghai University, Shanghai 200444, People's |
| 6  | Republic of China   |
| 7  | <sup>b</sup> Department of Physics and Materials Science, City University of Hong Kong, Tat Chee Avenue,      |
| 8  | Kowloon Tong, Hong Kong   |
| 9  |   |
| 10 |   |
| 11 |   |
| 12 |   |
| 12 |   |
| 13 |   |
| 14 |   |
|    |   |

\*Corresponding authors. Tel.: +86 21 66137503. Fax: +86 21 66137787. E-mail: <u>zwchen@shu.edu.cn</u>; <u>apcmlwu@cityu.edu.hk</u>.

<sup>&</sup>lt;sup>a</sup> School of Environmental and Chemical Engineering, Shanghai University, Shanghai 200444, People's Republic of China

<sup>&</sup>lt;sup>b</sup> Department of Physics and Materials Science, City University of Hong Kong, Tat Chee Avenue, Kowloon Tong, Hong Kong

Nanoscale Accepted Manuscript

# 15 Abstract

16 Tin dioxide  $(SnO_2)$  is unique strategic functional materials with widespread technological 17 applications, particularly in the fields such as solar battery, optoelectronic devices, and solid-18 state gas sensors owing to advances in optical and electronic properties. In this review, we 19 introduce recent progresses of tin dioxide and their composites, including the synthesis strategies, 20 microstructual evolution, related formation mechanism, and performance evaluation of  $SnO_2$ 21 quantum dots (ODs), thin films, and composites prepared by electron-beam irradiation, pulsed 22 laser ablation, and  $SnO_2$  planted graphene strategies, highlighting contributions from our 23 laboratory. First, we present the electron-beam irradiation strategies for growth behavior of 24 SnO<sub>2</sub> nanocrystals. This method is a potentially powerful technique to achieve the nucleation and 25 growth of SnO<sub>2</sub> ODs. In addition, the fractal assessment strategies and gas sensing behavior of SnO<sub>2</sub> thin films with interesting micro/nanostructures induced by pulsed delivery will be 26 27 discussed experimentally and theoretically. Finally, we emphasize the fabrication process and 28 formation mechanism of SnO<sub>2</sub> QDs planted graphene nanosheets. This review may provide a new 29 insight that the versatile strategies for microstructural evolution and related performance of SnO<sub>2</sub>-based functional materials are of fundamental importance in the development of new 30 31 materials.

33 Keywords: Tin dioxide; Electron-beam irradiation; Fractal assessment; SnO<sub>2</sub> planted graphene;
 34 Microstructure; Mechanism

35

36

# 37 **1. Introduction**

# 38 1.1. Preamble

39 Semiconductor oxides are fundamental to the development of smart and functional materials. devices, and systems.<sup>1-4</sup> These oxide materials have two unique structural features: mixed cation 40 41 valences and an adjustable oxygen deficiency, which are the bases for creating and tuning many novel material properties, from chemical to physical.<sup>5-7</sup> Developing new materials with excellent 42 43 performances depend not only on the multi-components but also on their micro/nanostructures. Multi-44 components and micro/nanostructures are the key issues to realize multifunctional materials. The key 45 scientific issues of tin dioxide (SnO<sub>2</sub>) functional materials and their applications and developments 46 have been driven scientists to explore continuously in depth their preparation, micro/nanostructure and performances. SnO<sub>2</sub>, as a kind of *n*-type wide-band-gap semiconductor ( $E_g = 3.64$  eV at 300 K), has 47 very wide applications in the fields such as solid-state gas sensors,<sup>8</sup> luminescent materials,<sup>9</sup> solar 48 battery,<sup>10</sup> antistatic coating,<sup>11</sup> and optoelectronic devices.<sup>12</sup> Since the properties of materials strongly 49 50 depend on its micro/nanostructures, which all result from the fabrication processes, the influence of 51 micro/nanostructural evolution on material properties is especially remarkable for materials science and 52 engineering. Tin oxide micro/nanostructures provide a versatile, multifaceted platform for a broad range of advanced applications.<sup>13-15</sup> Recent studies have validated their use in many fields of science 53 and technology that integrate the areas of chemistry, physics, materials science and nanotechnology.<sup>16-</sup> 54 <sup>20</sup> For example, they can function as semiconductors for a number of applications in microelectronic 55 56 devices, optoelectronic components, and solar-thermal converter, they can serve as gas-sensors for the 57 detection and forecasting of a variety of combustible gases, environmental pollution gases, industrial 58 emissions and hazardous gases, they can used as electrodes for melting optical glass and electrolytic 59 aluminum industry, they can act as catalysts for the active phase of many heterogeneous catalysis in hydrocarbon oxidation, they can also serve as varistor for the arrester used in power systems, and so on. 60

Nanoscale Accepted Manuscript

However, challenges remain to further improve the material fabrication processes for various advanced
 applications. This optimization requires a clear understanding of the relationship between
 micro/nanostructures and preparation conditions.

64 In this review, we will describe our efforts toward understanding the synthesis strategies, microstructual evolution, related formation mechanism, and performance evaluation of SnO<sub>2</sub> various 65 66 morphologies including SnO<sub>2</sub> quantum dots (QDs), thin films, and composites prepared by electron-67 beam irradiation (EBI), pulsed laser ablation, and SnO<sub>2</sub> planted graphene strategies, highlighting 68 contributions from our laboratory. Fig. 1 shows schematic illustrating in this article including  $SnO_2$ 69 electron-beam irradiation strategies, fractal assessment and SnO<sub>2</sub> planted graphene strategies. In 70 sections 2, we present the EBI strategies for growth behavior of SnO<sub>2</sub> nanocrystals. This method is a 71 potentially powerful technique to achieve the nucleation and growth of SnO<sub>2</sub> QDs. In sections 3, the 72 fractal assessment strategies and gas sensing behavior of SnO<sub>2</sub> thin films with interesting 73 micro/nanostructures induced by pulsed delivery will be discussed experimentally and theoretically. In 74 sections 4, we emphasize the fabrication process and formation mechanism of  $SnO_2$  ODs planted 75 graphene nanosheets. This review may provide a new insight that the versatile strategies for 76 microstructural evolution and related performance of SnO<sub>2</sub>-based functional materials are of 77 fundamental importance in the development of new materials.



78

Fig. 1. Schematic illustrating in this article including SnO<sub>2</sub> electron-beam irradiation strategies, fractal
 assessment and SnO<sub>2</sub> planted graphene strategies.

81

# 82 **1.2. Outline**

This review presents some advanced developments in the synthesis strategies and perspectives of microstructual evolution, related formation mechanism, and performance evaluation for  $SnO_2$  various morphologies including  $SnO_2$  QDs, thin films, and composites. This is an interdisciplinary work that integrates the areas of physics, chemistry and materials science. A brief outline of the contents and organization of each section is presented here to serve as a guide for reading this article.

The major aim in section 2 is to explore the EBI strategies for growth behavior of  $SnO_2$  nanocrystals. In this section, we will describe in detail that the  $SnO_2$  nanocrystals prepared by sol-gel method are radiated by electron accelerator. The microstructural evolution of  $SnO_2$  nanocrystals was evaluated by

XRD, HRTEM, Raman, thermo gravimetric-differential scanning calorimeter, and automated surface area analyzer. The characterization analysis indicated that the crystallinity of the irradiated SnO<sub>2</sub> powders was better than unirradiated one. The clear lattice fringes of SnO<sub>2</sub> QDs were investigated in detail by HRTEM. It was found that the sizes of SnO<sub>2</sub> QDs were localized in the range of 2-4 nm. The Brunauer-Emmett-Teller (BET) surface area analysis indicated that the specific surface area of irradiated sample was much higher, which was almost doubled at 1400 kGy under the 8 mA/s of irradiation rate and 50 seconds of irradiation time.

98 The purpose of the section 3 is to characterize the dependence of gas sensing behavior on fractal 99 dimension in SnO<sub>2</sub> thin films induced by pulsed delivery. SnO<sub>2</sub> thin films with interesting fractal 100 features were prepared by pulsed laser deposition techniques under different substrate temperatures. 101 The fractal assessments indicated that the fractal theory can be applied to the evaluation of this 102 material. The measurements of carbon monoxide gas sensitivity confirmed that the gas sensing 103 behavior was sensitively dependent on fractal dimensions, fractal densities, and average sizes of the 104 fractal clusters. Random Tunneling Junction Network (RTJN) mechanism was suggested to provide a 105 rational explanation for this gas sensing behavior. The formation process of SnO<sub>2</sub> nanocrystals and 106 fractal clusters were reasonably discussed by the use of a model.

107 In section 4, we present the assembling  $SnO_2$  QDs to graphene nanosheets (GNSs) by a facile 108 ultrasonic method. The experimental investigations indicated that the graphene was exfoliated and 109 decorated with SnO<sub>2</sub> QDs, which was dispersed uniformly on both sides of graphene. The size 110 distribution of  $SnO_2$  QDs was estimated to be a ranging from 4 to 6 nm and their average size was 111 calculated to be about  $4.8 \pm 0.2$  nm. This facile ultrasonic route demonstrated that the loading of SnO<sub>2</sub> 112 QDs was an effective way to prevent the GNSs from being restacked during the reduction. During the 113 calcination process, the GNSs distributed between SnO<sub>2</sub> nanoparticles have also prevented the 114 agglomeration of SnO<sub>2</sub> nanoparticles, which were beneficial to the formation of SnO<sub>2</sub> QDs.

# 116 **2. Electron-Beam Irradiation Strategies for Tin Dioxide**

# 117 2.1. Overview on Tin Dioxide Quantum Dots

118 SnO<sub>2</sub>, as a kind of *n*-type semiconductor, has very wide applications in the fields such as solid-state gas sensors,<sup>21</sup> luminescent material,<sup>22</sup> solar battery,<sup>23</sup> antistatic coating,<sup>24</sup> and optoelectronic devices<sup>25</sup> 119 120 owing to their unique electronic, optical, and magnetic properties. During the past decade, a lot of methods such as sol-gel,<sup>26</sup> chemical vapor deposition,<sup>27</sup> spray pyrolysis,<sup>28</sup> hydrothermal method,<sup>29</sup> and 121 microwave technique,<sup>30</sup> have been employed to prepare SnO<sub>2</sub> nanostructures. It is common that the 122 123 grain sizes and morphologies of nanomaterials greatly affect their properties as well as their applications. For example, when the size of the particle decreased, the surface-volume ratio and the 124 band gap will increase correspondingly, such effect will finally lead to the enhancement of the gas 125 126 sensors sensitivity. Therefore, the quantum size effects particularly in the semiconductor quantum dots have enhanced great interest in both basic and applied research.<sup>31</sup> Several studies on quantum dot (OD) 127 systems were carried out in the last decade mainly on GaAs, InP, and Cd chalcogenides due to their 128 129 discrete electronic states, unusual structural transformations and novel physical and chemical properties.<sup>32</sup> It is well-known that the size of QDs is usually among from 1 nm to 10 nm particularly if 130 131 the critical dimensions approach the exciton Bohr radius of nanomaterials which carriers are confined 132 in all three dimensions to a nanometer-sized region. Thus, the QDs exhibit distinctive properties as well as have a wide range applications in optoelectronics, optical sensors and lasers.<sup>33,34</sup> For this reason, the 133 134 properties of oxide semiconductors with wide-band-gap QDs have been growing attracted much 135 attention such as nanocrystalline SnO<sub>2</sub>. In a variety of SnO<sub>2</sub> QD synthesis processes, Castillo and coworkers reported on the preparation of SnO<sub>2</sub> QDs with the sizes in 2-5 nm by the sol-gel method.<sup>35</sup> Zhu 136 and co-workers obtained the SnO<sub>2</sub> QDs with the sizes in 2.3-3.1 nm by hydrothermal process.<sup>34</sup> 137

138 The irradiation technique achieved by electron accelerator is of widespread technological 139 applications in the fields of degraded organic compounds, medical treatment, industrial and agricultural applications.<sup>36</sup> Up to now, there are few reports related tin oxide modified by accelerator radiation. It is 140 141 well-known that the EBI method has a number of high advantageous properties. For instance, (i) this method is simple, rapid and convenient,<sup>37</sup> (ii) it is carried out at room temperature without any kind of 142 143 catalysts, and (iii) the method is useful for the mass-production of nanomaterials. After an extensive 144 search in the published literature, it was found that the previous nanocrystals are assembled either 145 aqueous synthesis, hydrothermal, and pyrolysis or epitaxial approaches through the gas phase. 146 Compared with the previous heat-treatment technique which is a commonly used crystallization 147 method, the production of SnO<sub>2</sub> QDs by the EBI method has also attracted much attention due to above 148 many advantages. In order to investigate the characteristics of these QDs, the microstructure evolution 149 of SnO<sub>2</sub> nanocrystals should be reasonably understood. An integrated device for different QDs is highly desirable for versatile advanced applications. Despite the high sensitivity of SnO<sub>2</sub> to many 150 151 gases, it is often susceptible to electrical drift which requires long stabilization periods, as well as 152 permanent poisoning after extended periods of operation. EBI may be a potentially powerful technique 153 to modify QDs, and we are applying this technique to SnO<sub>2</sub> nanocrystals for the first time in this study.<sup>38</sup> 154

155

# 156 2.2. Electron-Beam Irradiation Strategies

A typical procedure to synthesize  $SnO_2$  nanoparticles via sol-gel method was performed as follows. Under intensive stirring, aqueous ammonia (2 mol/L) was drop wisely added into  $SnCl_4$  solution (0.2 mol/L) until pH reached 7. The reaction temperature was kept at 60 °C. The primal sample was centrifuged and rinsed with alcohol and deionized water several times till the Cl<sup>-</sup> could not be detected

# Page 9 of 58

#### Nanoscale

162 collected after being ground into fine powder. Basic process is shown below:



164 The as-prepared  $SnO_2$  nanoparticles was dispersed as much as possible to thickness with 2-3 mm, 165 which was placed under the Titanium window of GJ-2-II Dynamitron electron accelerator with an 166 accelerating voltage of 2 MeV and a current of 8 mA by different radiation dose.

167

# 168 2.3. XRD and BET Analysis

169 It is known that the SnO<sub>2</sub> has a tetragonal rutile crystalline structure (known in its mineral form as cassiterite) with point group  $D_{4h}^{14}$  and space group  $P4_2 / mnm$ . The lattice parameters are a = b =170 4.7382(4) Å, and c = 3.1871(1) Å.<sup>39</sup> Fig. 2 shows the typical XRD patterns of SnO<sub>2</sub> powders (a) as-171 172 prepared (unirradiated) sample and (b) irradiated sample which was irradiated at 1400 kGy under the irradiation rate 8 mA/s and irradiation time 50 seconds. Through a series of experiments, it can be 173 174 found that the SnO<sub>2</sub> powders have obviously changed under 1400 kGy dose comparing with the as-175 prepared (unirradiated) powders. All diffraction peaks can be attributed to the tetragonal SnO<sub>2</sub> structure with cell parameters a = b = 4.738 Å and c = 3.187 Å, which is consistent with the standard card 176 (JCPDS file No.41-1445). The XRD patterns as shown in Fig. 2b reveal the stronger peak intensity and 177 smaller peak width, which indicated that the irradiated SnO<sub>2</sub> nanoparticles have better crystallizing.<sup>40</sup> 178 179 The presence of broad and weak peaks as shown in Fig. 2a also indicates that the unirradiated SnO<sub>2</sub> 180 nanoparticles smaller than irradiated one because of the width of XRD peaks is related to particle size 181 through Scherrer's equation:  $D = K\lambda / \beta \cos \theta$ , where D is the diameter of the nanoparticles, K = 0.9,  $\lambda$ 

**Nanoscale Accepted Manuscript** 

182 (Cu  $K_{\alpha}$ ) = 1.5406 Å, and  $\beta$  is the full width at half maximum of the diffraction peak.<sup>41</sup> Comparing with 183 Fig. 2a and b, it confirmed that the growth of SnO<sub>2</sub> nanocrystals is influenced significantly by the EBI 184 techniques. All peaks as shown in Fig. 2b became sharp and stronger, and appeared some new peaks 185 proved that the EBI is a potentially powerful technique to improve the growth of SnO<sub>2</sub> nanocrystals.



186

Fig. 2. The typical XRD patterns of SnO<sub>2</sub> powders (a) as-prepared (unirradiated) sample, (b) irradiated
 sample which was irradiated at 1400 kGy under the irradiation rate 8 mA/s and irradiation time 50
 seconds.<sup>38</sup>

190

191 Thermo-analysis of the as-prepared (unirradiated) and irradiated SnO<sub>2</sub> powders has been given in 192 Fig. 3a and b, respectively. As shown in Fig. 3b, the thermo gravimetric (TG) analysis of irradiated 193  $SnO_2$  powders showed that the amount of weight loss (the mass loss ~17%) is smaller than unirradiated 194 one (the mass loss  $\sim 22.3\%$ ), which indicated that the irradiated SnO<sub>2</sub> powders have been further 195 crystallized during EBI processes. For both samples studied in TG-DTG experiments, it was found that 196 the shapes of the TG-DTG curves are similar but exist still obviously difference. For instance, the observed weight loss of two samples between room temperature to 350 °C may be attributed to the loss 197 of ammonia, physically absorbed water and chemically bonded water,<sup>30</sup> which corresponds with the 198 199 endothermic peak at about 200 °C in DSC curves. Moreover, the DTG peaks of the irradiated SnO<sub>2</sub> 200 have a minor right shift comparing with the as-prepared SnO<sub>2</sub>, which indicated that the as-prepared SnO<sub>2</sub> is poor crystalline.<sup>42</sup> According to the formula:  $(dH/dT)/(dT/dt) = dH/dT = C_{p}$ , while dH/dt for the 201 DSC spectrum of the vertical axis, dT/dt is rise/cooling rate,  $C_p$  is specific heat capacity at constant 202 203 pressure, showing the specific heat of the grains in the relative size at this time. Thus, it determined that 204 the heat absorption capacity during crystallization of irradiated sample is much larger than the 205 unirradiated one. Comparing with previous researches such as Zhu and co-workers reported an endothermic peak at a temperature about 90 °C of the DSC curve for SnO<sub>2</sub> nanoparticles<sup>43</sup> and Zeng 206 and co-workers showed two endothermic peaks at 532 and 556 K.<sup>44</sup> There were significant differences 207 as the presented irradiated  $SnO_2$  that revealed two obvious exothermic peaks at temperatures of 78  $^{\circ}C$ 208 209 and 910 °C. However, the exothermic peak of the as-prepared SnO<sub>2</sub> shifted to 71 °C, moreover it was 210 not obviously. In addition, one of endothermic peaks nearly in 180 °C corresponds to the removal of 211 water while the endothermic peak of DSC curves for the as-prepared SnO<sub>2</sub> is flatter, which indicated that less activity center is formed.<sup>45</sup> 212



Fig. 3. The TG-DTG-DSC curves of  $SnO_2$  powders (a) as-prepared (unirradiated) sample, (b) irradiated sample which was irradiated at 1400 kGy under the irradiation rate 8 mA/s and irradiation time 50 seconds.<sup>38</sup>

217

213

218 Fig. 4 shows the specific surface areas of the as-prepared (unirradiated) sample and irradiated 219 samples with different radiation doses under the irradiation rate 8 mA/s. It can be seen that the as-220 prepared sample has the lower surface area than the irradiated samples. When the irradiated dose 221 increased to 700 kGy under the irradiation rate 8 mA/s and irradiation time 25 seconds, the BET surface area increased from 105.39  $m^2/g$  to 170.47  $m^2/g$ . However, the BET value has slightly 222 decreased to 165.42 m<sup>2</sup>/g at 980 kGy radiation dose under the irradiation rate 8 mA/s and irradiation 223 224 time 35 seconds. From view in the overall trend of BET values, it can be reasonable to speculate that 225 the more SnO<sub>2</sub> crystal nucleus could be formed in irradiated samples, which can also be proved by the 226 HRTEM results as shown in Fig. 5. The more smaller SnO<sub>2</sub> crystal nucleus, which may be called as

227 SnO<sub>2</sub> QDs, led to the increasing of the BET surface area of irradiated samples. However, when the 228 irradiated dose increased to 980 kGy, the rotation and coalescence processes of the initial SnO<sub>2</sub> crystal 229 nucleus (e.g. at 700 kGy) culminating in a low-energy configuration may be directly related to the reduction of surface energy, aimed at minimizing the area of high-energy interfaces.<sup>46,47</sup> The possible 230 231 formation of a coherent boundary between grains due to the grain rotation, with the consequence of 232 removing the common grain boundary and culminating in a single larger SnO<sub>2</sub> nanocrystal, may lead to 233 the decreasing of the BET value. With the increasing of the irradiated doses, for example, 1260 and 234 1400 kGy under the irradiation rate 8 mA/s and irradiation times 45 and 50 seconds, respectively, a 235 large number of SnO<sub>2</sub> crystal nucleus could be formed in irradiated samples. Therefore, the BET values 236 increased. The BET analysis indicated that the specific surface area of the irradiated samples tends to 237 increase as the increasing of the irradiated doses. It implied that the EBI method can significantly 238 enhance the BET surface area.



239

Fig. 4. Specific surface area of as-prepared (unirradiated) sample and irradiated samples at 700, 980,
1260 and 1400 kGy under the irradiation rate 8 mA/s and irradiation times 25, 35, 45, and 50 seconds,
respectively.<sup>38</sup>

Nanoscale Accepted Manuscript

243

# 244 **2.4. HRTEM and Raman Examination**

245 Since detailed and refined experiments on SnO<sub>2</sub> nanocrystals permitted us to assess the physics of 246 quantum-confined excitations, the geometry and nature of the contact areas between SnO<sub>2</sub> particles 247 (unirradiated and irradiated samples) are especially important for the steady state optical gain in small 248 semiconductor wide-band-gap QDs. Fig. 5 displays the HRTEM images of SnO<sub>2</sub> powders (a) as-249 prepared (unirradiated) sample and (b) irradiated sample which was irradiated at 1400 kGy under the 250 irradiation rate 8 mA/s and irradiation time 50 seconds. It was found that the SnO<sub>2</sub> QDs can be formed 251 in the as-prepared (unirradiated) and irradiated samples (indicated by the blue circles). The clear lattice 252 fringes demonstrated that the SnO<sub>2</sub> QDs are composed of ultrafine nanoparticles with a diameter below 253 4 nm. Comparison of Fig. 5a and b, the samples modified by EBI revealed clear-cut crystalline features 254 and more clearly crystal planes which could correspond to the lattice parameters of the rutile structure 255 of  $SnO_2$  cassiterite phase. The results indicated that the density and size of the  $SnO_2$  QDs increase after 256 the EBI. HRTEM analysis indicated that the crystal planes of the QDs become more complete. The 257 inset at the lower right-hand corner in Fig. 5b shows two ODs with size of about 4 nm and 3.5 nm. The 258 interplanar spacings are about 0.33 nm and 0.26 nm which respectively correspond to the (110) and 259 (101) planes of tetragonal SnO<sub>2</sub>. More importantly, HRTEM observation in the sample modified by 260 EBI is very interesting. Fig. 5b shows the initial stage of the SnO<sub>2</sub> crystal nucleus and QD growth 261 processes. The white arrow in blue circles indicated the formation of necks between QDs. An analysis 262 suggested that the first step in the QD growth process is the formation of necks. After this step, a 263 rotation may occur to decrease the angle of misorientation, changing the QD orientation. When the 264 QDs approach the same orientation, i.e., a coherent grain-grain boundary, the grain boundary must 265 migrate toward the smaller crystal nucleus, resulting in a single larger QD. Therefore, the EBI method 266 is a potentially powerful technique to achieve SnO<sub>2</sub> nucleation and QD growth.



267

Fig. 5. HRTEM images of  $SnO_2$  powders (a) as-prepared (unirradiated) sample, (b) irradiated sample which was irradiated at 1400 kGy under the irradiation rate 8 mA/s and irradiation time 50 seconds.<sup>38</sup>



270

Fig. 6. Room-temperature Raman spectra of (a) as-prepared (unirradiated) sample, (b) irradiated
 sample which was irradiated at 1400 kGy under the irradiation rate 8 mA/s and irradiation time 50
 seconds.<sup>38</sup>

274

In order to obtain further information on the microstructure evolution of  $SnO_2$  powders modified by the EBI, Raman techniques could effectively reveal local structural features of amorphous and poorly crystallized materials owing to its fine sensitive to crystal surface area.<sup>41</sup> Fig. 6 shows roomtemperature Raman spectra of (a) as-prepared (unirradiated) sample and (b) irradiated sample which was irradiated at 1400 kGy under the irradiation rate 8 mA/s and irradiation time 50 seconds. It can be seen that the irradiated sample shown in Fig. 6b presents five Raman peaks at 468, 538, 619, 680, 755 and 851 cm<sup>-1</sup> in contrast to the as-prepared (unirradiated) sample shown in Fig. 6a just presents three

Raman peaks at 427, 574 and 879 cm<sup>-1</sup>. In general, the SnO<sub>2</sub> unit cell consists of two metal atoms and four oxygen atoms. Each metal atom is situated amidst six oxygen atoms which approximately form the corners of a regular octahedron. Oxygen atoms are surrounded by three tin atoms which approximate the corners of an equilateral triangle. As a result, there are 18 branches for the vibrational modes in the first Brillouin zone. The normal vibration modes at the  $\Gamma$  point at the center of the Brillouin zone is given by<sup>47</sup>

288 
$$\Gamma = \Gamma_{l}^{+}(A_{lg}) + \Gamma_{2}^{+}(A_{2g}) + \Gamma_{3}^{+}(B_{lg}) + \Gamma_{4}^{+}(B_{2g}) + \Gamma_{5}^{+}(E_{g}) + 2\Gamma_{l}^{+}(A_{2u}) + 2\Gamma_{4}^{+}(B_{lu}) + 4\Gamma_{5}^{+}(E_{u}).$$

According to these modes,  $A_{1g}$ ,  $B_{1g}$ ,  $B_{2g}$  and  $E_g$  are Raman active,  $A_{2u}$ ,  $E_u$  and  $B_{1u}$  are active in infrared 289 (IR), one  $A_{2u}$  and two  $E_u$  modes are acoustic. Raman peaks at 468, 619 and 777 cm<sup>-1</sup> in Fig. 6b can be 290 indexed to  $E_g$ ,  $A_{Ig}$  and  $B_{2g}$  vibration modes, respectively.<sup>48</sup> In addition, the other Raman peaks at 538 291 and 680 cm<sup>-1</sup> also were observed, which may be close to the IR-actived  $A_{2u}$  modes at 512 cm<sup>-1</sup> for 292 transverse-optical (TO), and at 687 cm<sup>-1</sup> for longitudinal-optical (LO).<sup>49</sup> In a perfect crystal, only is the 293 294 center of Brillouin zone (q = 0) phonon can be excited because the correlation length of the particular 295 phonon is infinite. In the nanocrystals, some IR active modes may transform into Raman active with 296 decreasing size and increasing disorder for the nanomaterials. Thus, Raman scattering intensity of lattice vibration is<sup>50</sup> 297

298 
$$I(\omega) \propto \int_0^1 \exp \left| -\frac{q^2 L^2}{4} \right| \frac{dq^3}{\left[ \omega - \omega(q) \right]^2 + \left[ \Gamma_0 / 2 \right]^2},$$

299 and 
$$\omega(q) = A + B\cos(\Pi q)$$
,

where  $\omega$  is phonon frequency, q is phonon wave vector,  $\omega(q)$  is the phonon dispersion curve of the infinite crystal,  $\Gamma_0$  is the natural line width, and L for the coherence length. Abello and co-workers have reported that some vibration modes forbidden in the theory would be activated when the dimension of the grains is small enough, especially, some IR-active modes may be weakened to show Raman-

active.<sup>51</sup> In our Raman experiments, the as-prepared (unirradiated) sample shown in Fig. 6a has three peaks, which there is a few research on the peak of 574, it may be the peak displacemen of amorphous  $SnO_2$ .<sup>52</sup> However, the irradiated sample shown in Fig. 6b has a obvious blue shift compared with Fig. 6a. This may attributed to relaxation of Raman selection rule due to the combined effects in the increasing of QD density and the high concentration of defects in surface site, such as oxygen vacancies and lattice disorder.<sup>51</sup>

310

# 311 **3. Fractal Assessment Strategies for Tin Dioxide**

# 312 **3.1. Overview on Tin Dioxide Tin Films**

Semiconductor oxides are fundamental to the development of smart and functional materials, 313 devices, and systems.<sup>53-55</sup> These oxides have two unique structural features: mixed cation valences and 314 315 an adjustable oxygen deficiency, which are the bases for creating and tuning many novel material properties, from chemical to physical.<sup>56-58</sup> Due to the increasing importance of air pollution and the 316 need to monitor concentration levels of gases such as CO, CO<sub>2</sub>, NO<sub>x</sub>, O<sub>3</sub>, SO<sub>2</sub> etc., the development of 317 318 many kinds of sensors and control systems has been jolted into action in recent years. Tin dioxide 319 (SnO<sub>2</sub>) has been used as a gas sensor material to detect combustible and toxic gases such as CO, NH<sub>3</sub>, NO<sub>2</sub>, H<sub>2</sub>S and CH<sub>4</sub>. Commercial sensors typically use sintered SnO<sub>2</sub> powders, but thin films SnO<sub>2</sub> are 320 gaining increasing popularity.<sup>59,60</sup> With the advent of advanced thin film technology more cost-321 322 effective, reproducible devices can be constructed with a reduction in device size and a concomitant 323 increase in the speed of response by using SnO<sub>2</sub> thin films.

It is known that  $SnO_2$  is used as a gas sensor because the number of electrons in the conduction band is affected by the adsorption of gaseous species on its surface.<sup>61</sup> Reducing gas molecules, such as CO, react with the oxygen species ( $O_2$ ,  $O_2^-$ ,  $O^-$ , and O) on the semiconductor surface.<sup>62</sup> This lowers the height of the Schottky barrier and increases the conductance of the material.<sup>63,64</sup> Gas sensors using

SnO<sub>2</sub> are widely used due to its high sensitivity to humidity and inflammable gases. In this type of 328 329 sensors, gas concentration is related to the material's electrical impedance due to the adsorption of gas 330 molecules on the SnO<sub>2</sub> surface. In a pure air environment, SnO<sub>2</sub> adsorbs oxygen that captures its 331 electrons, thereby raising its resistivity. When a reducing gas is present, it competes for the adsorbed 332 oxygen and hence the SnO<sub>2</sub> resistivity decreases. However, the electrical properties of SnO<sub>2</sub> are 333 strongly dependent on material fabrication parameters. Since gas sensing is based on adsorption 334 mechanisms on the SnO<sub>2</sub> grain surface, for high sensitivity a small grain size is desirable in order to achieve a high specific area, i.e. adsorption area per unit volume.<sup>65,66</sup> 335

336 The surface conductance of semiconducting oxide is affected by the concentration of ambient gases. 337 Resistive gas sensors are based on this principle and the nature of the sensing mechanism is related to 338 the electrical response of gas sensors to reactive gases. The change in the sensor resistance provides an indication of the gas concentration.<sup>67</sup> These sensors can be quite versatile as they may be used to detect 339 340 oxygen, flammable gases and common toxic gases. Their mechanism of operation is complex, 341 involving interactions between gaseous molecules and defects on the surface and grain boundaries. The sign of a change in resistance depends on whether the solid has *n*-type or *p*-type conductivity.<sup>68,69</sup> 342 343 Structural properties such as grain size, grain geometry as well as specific surface area can significantly 344 affect the gas sensing properties of semiconducting SnO2. In order to control these structural 345 characteristics, the microstructure evolution of SnO<sub>2</sub> thin films should be understood. Fractal method is 346 a potentially powerful technique to characterize microstructures, and we are applying this technique to SnO<sub>2</sub> thin films for the first time in this study. Besides showing some examples of geometric structures 347 348 of SnO<sub>2</sub> thin films, we shall discuss in detail the applicability and relevance of fractal theory to 349 studying the microstructure and gas sensing behavior of SnO<sub>2</sub> based environmental functional 350 materials. An integrated device for different gas species is highly desirable for versatile advanced 351 applications. Despite the high sensitivity of SnO<sub>2</sub> to many gases, it is often susceptible to electrical drift

Nanoscale Accepted Manuscript

which requires long stabilization periods, as well as permanent poisoning after extended periods of operation. New fractal assessment strategies for this material formed at different substrate temperatures are of fundamental importance in the development of micro-devices.<sup>70</sup>

355

356 **3.2. Fractal Assessment Strategies** 

357 In order to obtain the sintered SnO<sub>2</sub> target for pulsed laser deposition (PLD), we synthesized a pure nanocrystalline SnO<sub>2</sub> powder by the sol-gel method.<sup>71</sup> The fabrication method is described in the 358 359 following. Meta-stannic acid sol (parent sol) was precipitated by treating a cold ethanol solution of 360 SnCl<sub>4</sub> (27 %) with an aqueous ammonia solution (28 %) until a suitable pH value was reached. Dry powder with average grain size of about 4 nm was obtained by drying the parent sol, which had been 361 washed repeatedly with de-ionized water. The SnO<sub>2</sub> discs, 15 mm in diameter and 4 mm in thickness, 362 363 were prepared by compacting the powder under uniaxial pressure of 0.4 GPa, and sintered at 1150 °C for 2 h. The sintered disc consisted of high-purity cassiterite structure SnO<sub>2</sub> (99.8 %). 364

SnO<sub>2</sub> thin film was prepared by PLD techniques using the above sintered SnO<sub>2</sub> disc.<sup>55</sup> The target was 365 366 cleaned with methanol in an ultrasonic cleaner before installation to minimize contamination. The laser 367 was a KrF excimer laser (Lambda Physik, LEXtra 200, Germany) producing pulse energies of 350 mJ 368 at a wavelength of 248 nm and a frequency of 10 Hz. The duration of every excimer laser pulse was 34 ns. The laser energy was transmitted onto the target in a high-vacuum chamber through an ultraviolet 369 370 (UV)-grade fused silica window using an UV-grade fused silica lens. During the experiment, the target was kept rotating at a rate of 15 rpm to avoid drilling. The fluence was set at 5 J/cm<sup>2</sup> per pulse, 371 corresponding to a total of approximately  $1.5 \times 10^5$  laser pulses. The growth rate was estimated to be 372 373 about 0.3 nm/s (or about 1 µm/h). The ablated substance was collected on a Si (100) substrate mounted 374 on a substrate holder 4 cm away from the target. The high vacuum in the deposition chamber was Page 21 of 58

# Nanoscale

achieved by using a cryopump (Edwards Coolstar 800). The base pressure prior to laser ablation was about  $1 \times 10^{-6}$  mbar, and the oxygen partial pressure during laser ablation was set about  $3 \times 10^{-2}$  Pa. All deposition processes were carried out by in-situ operation on the substrate at temperatures of 300 °C,  $350 \,^{\circ}$ C,  $400 \,^{\circ}$ C, and  $450 \,^{\circ}$ C.

379 Scanning electron microscopy (SEM) images were digitized by using the Fractal Images Process 380 Software (FIPS). These digitized images were divided into boxes of  $360 \times 360$  size and then processed by the fractal theory.<sup>72</sup> Four intact fractal patterns were selected from these digitized images. The 381 382 average value of the fractal dimensions (D), the fractal density and the average size of the fractal clusters for these digitized fractal patterns were obtained by using the box-counting method.<sup>73</sup> The 383 384 carbon monoxide (CO) gas sensing property in the sensor was measured by a simple electrical 385 measuring system. The test CO gas was introduced in the chamber by an injector with variable volume 386 which facilitated control of gas concentration in the range 25-500 ppm. After the sensor was stabilized, 387 the process was repeated by injecting a higher amount of the CO gas.

388

# 389 **3.3. XRD and SEM Analysis**

390 It is known that tin dioxide has a tetragonal rutile crystalline structure (known in its mineral form as cassiterite) with point group  $D_{4h}^{14}$  and space group  $P4_2/mnm$ . The unit cell consists of two metal 391 392 atoms and four oxygen atoms. Each metal atom is situated amidst six oxygen atoms which 393 approximately form the corners of a regular octahedron. Oxygen atoms are surrounded by three tin 394 atoms which approximate the corners of an equilateral triangle. The lattice parameters are a =4.7382(4) Å, and c = 3.1871(1) Å. Fig. 7A, B, C, and D show the XRD patterns of the SnO<sub>2</sub> thin films 395 prepared on Si (100) substrate at 300 °C, 350 °C, 400 °C, and 450 °C respectively. The major 396 397 diffraction peaks of some lattice planes can be indexed to the tetragonal unit cell structure of SnO<sub>2</sub> with lattice constants a = 4.738 Å and c = 3.187 Å, which are consistent with the standard values for bulk 398

399 SnO<sub>2</sub> (International Center for Diffraction Data (ICDD), PDF File No. 77-0447). The (hkl) peaks 400 observed are (110), (101), (200), (211), (220), and (002). No characteristic peaks belonging to other tin 401 oxide crystals or impurities were detected. The high intensity of these peaks suggests that these thin 402 films mainly consist of the crystalline phase. As the substrate temperature increased, the crystallinity of 403 the thin films was enhanced as manifested by the intensity and sharpness of the XRD peaks of the  $SnO_2$ 404 thin films. The substrate temperature dependence can be interpreted mainly by the mobility of the 405 atoms in the thin films. At low substrate temperatures, the vapor species have a low surface mobility 406 and are located at different positions on the surface. The low mobility of the species will prevent full 407 crystallization of the thin films. However, at high substrate temperatures, the species with high enough mobility will arrange themselves at suitable positions in the crystalline cell.<sup>74,75</sup> The SnO<sub>2</sub> average grain 408 sizes were calculated using the Scherrer formula:  $D = K\lambda / \beta \cos \theta$ , where D is the diameter of the 409 nanoparticles, K = 0.9,  $\lambda$  (Cu  $K_{\alpha}$ ) = 1.5406 Å, and  $\beta$  is the full-width-at-half-maximum of the 410 411 diffraction lines. The results show that the average grain sizes of the SnO<sub>2</sub> nanoparticles at different 412 substrate temperatures are in the range of 25.3-27.8 nm. SnO<sub>2</sub> nanoparticle size increases from 25.3 nm 413 at 300 °C to 26.2 nm at 350 °C. It then increases to 27.0 nm at 400 °C and finally to 27.8 nm at 450 °C. 414 In fact, SnO<sub>2</sub> nanostructures can work as sensitive and selective chemical sensors. SnO<sub>2</sub> nanostructure 415 sensor elements can be configured as resistors whose conductance can be modulated by charge transfer 416 across the surface or as a barrier junction device whose properties can be controlled by applying a 417 potential across the junction. Functionalizing the surface further offers a possibility to improve their 418 sensing ability based on a better understanding of the influence of significant microstructural features, 419 for example, the development of gas sensors for the detection of environmentally harmful gases.



# 420

Fig. 7. XRD patterns of SnO<sub>2</sub> thin films prepared on Si (100) substrate at temperatures of (A) 300 °C;
(B) 350 °C; (C) 400 °C; and (D) 450 °C.<sup>70</sup>

423

424 Fig. 8 presents SEM images of SnO<sub>2</sub> thin films prepared on Si (100) substrate at temperatures of (A) 300 °C, (B) 350 °C, (C) 400 °C, and (D) 450 °C respectively. The SEM observation indicated that all 425 426 thin films produced under different substrate temperatures exhibited self-similar fractal patterns. It can 427 be seen from Fig. 8 that the fractal patterns are open and loose structure with increasing substrate 428 temperature. The average sizes of the fractal patterns (or clusters) are about 0.307 µm (see Fig. 8A), 429 0.906 µm (see Fig. 8B), 1.202 µm (see Fig. 8C) and 1.608 µm (see Fig. 8D). The average sizes of the 430 fractal clusters for four thin films were estimated by measurement on the fractal regions. The 431 measuring procedure is as follows: for each SEM image, we chose ten fractal patterns at random to get 432 an average value. The average sizes of the fractal patterns were obtained by averaging the values of

- 433 SEM images with different orientations. It was found that the average sizes of the fractal clusters
- 434 increase with increasing substrate temperature.



435

Fig. 8. SEM images of SnO<sub>2</sub> thin films prepared on Si (100) substrate at temperatures of (A) 300 °C; (B)
350 °C; (C) 400 °C; and (D) 450 °C.<sup>70</sup>

438

# 439 **3.4. Fractal Dimension Calculation**

Fig. 9 shows that the plots of  $\ln(N)$  versus  $\ln(1/L)$  of the fractal cluster regions in Fig. 8, where L is 440 the box size and N is the number of boxes occupied by the  $SnO_2$  clusters. It can be seen that all plots 441 442 show good linearity, which means that the morphologies of SnO<sub>2</sub> clusters have scale invariance within 443 these ranges. So the SnO<sub>2</sub> clusters can be regarded as fractals. In order to obtain the fractal dimension 444 (D), we fit a linear relationship for the function  $\ln(N)$  versus  $\ln(1/L)$ . The results show that the fractal dimension (D) is 1.896 at 300 °C as shown in Fig. 9A, 1.884 at 350 °C as shown in Fig. 9B, 1.865 at 445 446 400 °C as shown in Fig. 9C, and 1.818 at 450 °C as shown in Fig. 9D. We found that the fractal 447 dimension (D) decreases with increasing substrate temperature. The smaller fractal dimension means

# Page 25 of 58

# Nanoscale

that the SnO<sub>2</sub> thin films are composed of the open and loose fractal structure with finer branches. Fig. 448 449 10A to C shows the distribution of the fractal average size, fractal dimension and fractal density for 450 different substrate temperatures. It can be seen that there is an obvious increase in average fractal size 451 (see Fig. 10A), and the fractal dimension generally decreases (see Fig. 10B) with increasing substrate 452 temperature. In general, the fractal density is determined by the initial nucleation probability of the core crystal. From Fig. 10C, the fractal density was calculated to be 18, 6, 3, and 2 mm<sup>-2</sup> at 300 °C, 350 °C, 453 400 °C, and 450 °C respectively. It was found that the fractal density gradually decreases with 454 455 increasing substrate temperature. In the present work, the initial increase in nucleation probability was 456 due to strain relaxation caused by the low short-range temperature field at 300 °C, so that the fractal 457 density and their occupation area were high. With the increase of substrate temperature, the higher 458 long-range temperature field may promote new nuclei and subsequent growth, which leads to the fractal growth of the fine branches and a lower fractal density. This fractal structure may lead to 459 460 improvement in the design of gas sensors for the monitoring of environmental pollutants.



461

462 **Fig. 9.** Plots of  $\ln(N)$  versus  $\ln(1/L)$  of the fractal cluster regions in Fig. 3, where *L* is the box size and *N* 463 is the number of boxes occupied by the SnO<sub>2</sub> crystalline structure for substrate temperatures at (A) 300 464 °C; (B) 350 °C; (C) 400 °C; and (D) 450 °C.<sup>70</sup>



465

466 Fig. 10. (A) The fractal average size; (B) the fractal dimension; (C) the fractal density versus the
 467 substrate temperature.<sup>70</sup>

468

# 469 **3.5. Fractal Formation Mechanism**

470 On the basis of our experimental observation, the formation process of  $SnO_2$  nanocrystals and fractal 471 clusters could be reasonably described by a novel model, and be separated into eight steps, which 472 illustrated in detail in Fig. 11.

Nanoscale Accepted Manuscript

(i) Operation of the KrF excimer laser at a repetition rate of 10 Hz at an incident angle of  $45^{\circ}$  to the 473 474 polished sintered cassiterite SnO<sub>2</sub> target rotating at a rate of 15 rpm to avoid drilling. 475 (ii) Production of the high-temperature and high-pressure tin dioxide plasma at the solid-liquid interface quickly after the interaction between the pulsed laser and SnO<sub>2</sub> target. 476 (iii) Subsequent expansion of the high-temperature and high-pressure tin dioxide plasma leading to 477 cooling of the tin dioxide plumes.<sup>76-79</sup> In our case, the interval between two successive pulses is much 478 479 longer than the life of the plasma. Therefore, the next laser pulse had no interaction with the former 480 plasma. 481 (iv) Deposition of the tin dioxide plume on the Si (100) substrate after the disappearance of the 482 plasma, inducing the initial nucleation of SnO<sub>2</sub> nanocrystals. 483 (v) Grain rotation culminating in a low-energy configuration. This process is directly related to the reduction of surface energy, aimed at minimizing the area of high-energy interfaces.<sup>45,46</sup> 484 (vi) Possible formation of a coherent boundary between grains due to grain rotation, with the 485 486 consequence of removing the common grain boundary and culminating in a single larger SnO<sub>2</sub> nanocrystal. This is the coalescence process. 487 488 (vii) Growth of SnO<sub>2</sub> nanocrystals along preferred crystallographic directions which could be 489 predicted by an analysis of the surface energy in several crystallographic orientations. 490 (viii) Formation of the fractal structure as SnO<sub>2</sub> crystallizes and nucleates at high energy interfaces 491 such as grains boundaries. According to the fractal theory,<sup>80,81</sup> the heat released by crystallization leads to a local temperature 492 rise in the surrounding area and this temperature field can propagate quickly and stimulate new nuclei 493

494 appearing randomly in nearby regions. The stimulated nuclei of the next generation can also cause a

28

local temperature rise and repeat the above process many times until  $SnO_2$  fractal patterns are formed. Based on the above proposed formation mechanism, we characterize the formation processes of  $SnO_2$ nanocrystals and fractal structure in Fig. 11A-I. We believe that laser ablation technique is an appropriate method to synthesize a series of environmental functional materials with controlled composition, morphology and nanocrystal size, which are of important in the study of the sensitivity of SnO<sub>2</sub> thin films.



501

Fig. 11. The formation process of SnO<sub>2</sub> nanocrystals and fractal clusters. (A) Laser; (B) Target; (C)
Plasma; (D) Plume; (E) Nucleation; (F) Grain Rotation; (G) Coalescence; (H) Growth; and (I)
Fractal.<sup>70</sup>

Nanoscale Accepted Manuscript

505

# 506 **3.6. Gas Sensing Behavior**

507 To verify the gas sensing behavior of these SnO<sub>2</sub> thin films with interesting features of the fractal 508 structure, we investigate the sensitivity dependence on carbon monoxide (CO) concentration, so as to 509 achieve the aim of monitoring environmental pollutants. Fig. 12 shows the CO gas sensing behavior of 510 the SnO<sub>2</sub> thin films prepared on Si (100) substrate at (A) 300 °C, (B) 350 °C, (C) 400 °C, and (D) 450 511 <sup>o</sup>C respectively. The measurement was performed at room temperature with CO concentrations of 25, 512 50, 75, 100, 200, 300, 400, and 500 ppm. We observe that the sensitivity increases with increasing CO 513 concentration and substrate temperature. Similarly Cooper and Cicera found their SnO<sub>2</sub> thin film sensor possessed higher sensitivity to CO by using different procedures.<sup>82,83</sup> Further advancement of this gas 514 515 sensor fabricated by the SnO<sub>2</sub> thin films with fractal structure to detect environmental harmful gases 516 such as CO requires a clear understanding of its gas sensing mechanism. Our experimental results show 517 that the CO gas sensing behavior clearly depends on the fractal dimension, fractal density, and average 518 sizes of the fractal clusters (see Fig. 10 and 12). We propose a Random Tunneling Junction Network 519 (RTJN) mechanism to explain this gas sensing behavior. After the fractal formation, the fractal clusters 520 consist of the SnO<sub>2</sub> grains with the morphology of fine dendrite-like nanocrystals incorporating many 521 tunneling junctions of varying sizes. From the view of electron transport, the whole thin film is made 522 up of a series of tunneling junctions. For the SnO<sub>2</sub> thin films deposited at different substrate 523 temperatures, the sizes of the fractal branches with different fractal dimensions are different, leading to 524 differences in the height of the Schottky barrier of the tunneling junctions, with the consequence that 525 the breakdown voltages are also different. During the measurement of the gas sensitivity, the reducing 526 gas molecules such as CO react with the oxygen species  $(O_2, O_2, O_2, O_2, O_2)$  and O) ionized on the surface of the SnO<sub>2</sub> particles. This lowers the height of the Schottky barrier, and increases the conductance.<sup>62-64</sup> 527 For example, for the SnO<sub>2</sub> thin film deposited at the lower substrate temperature (e.g. at 300 °C) with 528

529 the larger fractal dimension, the junction *i* will have the higher resistance state due to the thicker fractal 530 branches, so the external voltage  $V_i$  cannot lower the Schottky barrier  $S_i$  and the junction *i* cannot be 531 broken. The gas sensitivity is then lowered (see Fig. 12A). Conversely, for the SnO<sub>2</sub> thin film deposited 532 at the higher substrate temperature (e.g. at 450  $^{\circ}$ C) with the smaller fractal dimension, the junction *i* 533 will have the lower resistance state due to the finer fractal branches, so the external voltage  $V_i$  can 534 lower the Schottky barrier  $S_i$  and the junction i will be broken. Therefore, the gas sensitivity would be 535 higher (see Fig. 12D). As mentioned above, there is a relationship between the fractal dimension and 536 the size of the fractal branches in that the number of the fine branches increases with decreasing fractal 537 dimension. Therefore, the smaller the fractal dimension, the larger the number of junctions with the 538 smaller Schottky barrier  $S_i$  and lower resistance state. The present findings reveal new opportunities for 539 future study of fractal structure tin dioxide architectures, with the goal of optimizing environmental 540 functional material properties for specific applications.



542 Fig. 12. The CO gas sensing behavior of SnO<sub>2</sub> thin films prepared on Si (100) substrate at temperatures
543 of (A) 300 °C; (B) 350 °C; (C) 400 °C; and (D) 450 °C.<sup>70</sup>

544

# 545 4. Tin Dioxide Planted Graphene Nanosheets

# 546 **4.1. Overview on Tin Dioxide Composite**

547 Nanocomposites have attracted widespread attention because of their potential to combine desirable 548 properties of different nanoscale building blocks to achieve advantageous electronic, optical, magnetic, and mechanic properties.<sup>84-86</sup> It has been demonstrated that the high-performance lightweight 549 550 composites could be developed by dispersing strong and highly stiff fibres in a polymer matrix and tailored to individual applications.<sup>87</sup> Graphene as the basic plane of graphite is a single atomic layer of 551  $sp^2$  hybridized carbon atoms arranged in a honeycomb lattice. Graphene materials have attracted 552 special interest due to its excellent optical,<sup>88-90</sup> mechanical,<sup>91-93</sup> and electrochemical properties.<sup>94-98</sup> 553 554 After an extensive search in the published literature, we stand at a similar threshold in the realm of 555 incorporation of well-dispersed graphene-based sheets, produced by mechanical exfoliation, solution 556 exfoliation and reduction of graphite oxide with metal oxides, to result in excellent applications. Many 557 types of metal oxide/graphene nanocomposites have been effectively prepared and extensively CoO/graphene.<sup>99,100</sup> Co<sub>3</sub>O<sub>4</sub>/graphene,<sup>101-104</sup> CuO/graphene,<sup>105,106</sup> 558 discussed, including Mn<sub>3</sub>O<sub>4</sub>/graphene,<sup>107-109</sup> Fe<sub>3</sub>O<sub>4</sub>/graphene,<sup>110,111</sup> NiO/graphene,<sup>112</sup> TiO<sub>2</sub>/graphene.<sup>113,114</sup> 559 and ZnO/graphene,<sup>115</sup> etc. These nanocomposites are achieved, not only by using the inherent properties of 560 561 the nanofiller, but also more importantly by optimizing the dispersion, interface chemistry and 562 nanoscale morphologies to take advantage of the enormous surface area per unit volume of graphene.

563 On the other hand, the size and morphology of metal oxides on graphene greatly affect their 564 properties as well as their applications. For instance, when the size of the particle decreased, the 565 surface-volume ratio and the band gap will increase correspondingly. This effect will finally lead to 566 enhancement of the sensitivity of gas sensors.<sup>70</sup> Therefore, the quantum effects, particularly in size of 567 quantum dots, have enhanced great interest in both basic and applied research.<sup>38,116</sup> However, most

metal oxide/graphene composites prepared so far have relatively larger sizes, over 10 nm up to hundreds of nanometers, with unsatisfactory dispersions. Furthermore, the main routes commonly used for the preparation of the metal oxide/graphene composites were either carried out with complicated processes or often suffered from poor manipulation on metal oxide/graphene composites. Therefore, it is highly desirable to develop a facile and general approach for the synthesis of metal oxide/graphene composites with favored microstructures for high-performance electrochemical properties.<sup>117</sup>

574

# 575 4.2. Tin Dioxide Planted Graphene Strategies

576 The reaction procedure can be described as a two-step method: (i) in situ loading of  $SnO_2$  QDs 577 (SODs) onto graphene oxide (GO) and (ii) GO conversion to graphene. The detailed preparation procedure is as follows: GO solution was synthesized from natural graphite powders by a modified 578 Hummer's method.<sup>99</sup> The GO was washed repeatedly with de-ionized water to completely remove 579 580 metal and acid particles. SnO<sub>2</sub> quantum dots and graphene nanosheets (SODs/GNS) composites were 581 prepared by co-precipitation method of the GO and SnCl<sub>2</sub>·2H<sub>2</sub>O solution. In a typical synthesis, 100 ml 582 GO solution was dispersed in 50 ml distilled water by ultrasonic treatment to form a colloidal 583 suspension. Subsequently, the suspension was mixed with 10 ml aqueous solution of SnCl<sub>2</sub>·2H<sub>2</sub>O (0.45 584 g). The resulting mixture with the Sn/graphene ratio = 1:1 in weight was continually stirred for 30 min. 585 The mixture was then placed in an ultrasonic bath for 60 min. The product was left for 24 h, washed 586 completely with distilled water, and dried in vacuum at 60 °C. Finally, the product was annealed at 500 <sup>o</sup>C for 3 h in N<sub>2</sub> atmosphere. For comparison, the pure graphene nanosheets were prepared by the 587 588 reduction of dried GO in N<sub>2</sub> atmosphere at 500 °C for 3 h.

589 The XRD patterns were obtained from a Japanese Regaku D/max-2500 machine using Cu  $K_{\alpha}$ 590 radiation in reflection geometry. An operating voltage of 40 kV and a current of 40 mA were used.

Nanoscale Accepted Manuscript

XRD patterns were recorded at a scanning rate of 0.08 °s<sup>-1</sup> in the 20 range 10 to 90°. Field emission-591 592 secondary electron microscopy (FE-SEM) imaging was carried out with a field emission scanning 593 electron microanalyzer (JEOL-6700F, 15 kV). HRTEM observations were performed on a JEOL JEM-2010F electron microscope operating at 200 kV. Surface area determinations were performed by 594 595 Brunauer-Emmett-Teller (BET) method using an ASAP-2000 Surface Area Analyzer (Micromeritics 596 Instrument Corporation). Raman scattering measurements were obtained by backscattering geometry 597 with a SPEX-1403 laser Raman spectrometer. The excitation source was an argon-ion laser operated at 598 a wavelength of 514.5 nm in the backscattering configuration and a low incident power to avoid 599 thermal effects. Thermogravimetric analysis (TGA) was conducted on a STA 449C instrument (Netzsch Instrument Co., German) from 30 to 1000 °C at a speed of 10 °C/min under N<sub>2</sub> atmosphere. 600

601

# 602 **4.3. Microstructure Analysis**

603 Fig. 13 shows the typical XRD patterns of (a) bare pure GNS, (b) SQDs/GNS composites, and (c) 604  $SnO_2$  nanoparticles. The pure GNS only shows a strong (002) diffraction peak in the XRD pattern as 605 shown in Fig. 13a. This peak shows that the GNS is multi-layer graphene stacked with an interlayer spacing of about 0.35 nm, indicating that the graphene is similar to natural graphene sheets ( $d_{002} = 0.34$ 606 nm).<sup>118</sup> Comparing Fig. 13a, b and c, it can be seen that the major diffraction peaks of some lattice 607 608 planes of the SQDs/GNS composites as shown in Fig. 13b are similar to the tetragonal unit cell structure of SnO<sub>2</sub> with lattice constants a = 4.738 Å and c = 3.187 Å, consistent with the standard 609 610 values for bulk SnO<sub>2</sub> (International Center for Diffraction Data (ICDD), PDF File No. 77-0447). The 611 observed (*hkl*) peaks are (110), (101), (200), (211), (220), (002), (112) and (202), indicating that the 612 SnO<sub>2</sub> particles in this composite are well-crystallized. The minor peak (particularly at ca. 15 degrees) 613 belonging to other tin oxide crystals or impurities were detected. The high intensity of these peaks

# Page 35 of 58

# Nanoscale

614 suggests that this SQDs/GNS composite mainly consists of the crystalline phase, while the (002)
615 diffraction peak of layered GNS has almost disappeared. Thus, we can speculate that the regular
616 lamellar structure of the GNS sheets has been completely broken in the as-synthesized composites,
617 forming exfoliated GNS nanosheets.



618

Fig. 13. The typical XRD patterns of (a) bare pure GNS, (b) SQDs/GNS composites, and (c) SnO<sub>2</sub>
 nanoparticles.<sup>117</sup>







623

Fig. 14 shows the FE-SEM images of (a) GNS and (b) SQDs/GNS composites on Si (100) 624 625 substrates. Fig. 14a shows a representative SEM image of the GNS from the top view. This image 626 shows layered platelets composed of curled nanosheets, in accordance with the (002) diffraction peak 627 in the XRD pattern of the GNS. Fig. 14b shows the FE-SEM image of the SODs/GNS composites 628 prepared by the facile synthesis approach. It is found that the FE-SEM images of the GNS and 629 SQDs/GNS composites have similar morphology, implying that a fine structural manipulation of the 630 GNS was successfully achieved even after the reassembling process with SnO<sub>2</sub> nanoparticles. 631 Interestingly, it can be seen clearly that the SnO<sub>2</sub> nanoparticles were uniformly distributed on the 632 graphene nanosheets, by which graphene nanosheet restacking and SnO<sub>2</sub> nanoparticles aggregation 633 were prevented. The component of prepared composites was confirmed by energy-dispersive X-ray 634 analysis (EDS) as shown in the inset of Fig. 14b. It can be found that only Sn, C, and O elements were 635 detected and the atomic ratio of Sn and C atoms was 1:10, which was consistent with their weight ratio 636 (1:1). This further proves that the SnO<sub>2</sub> nanoparticles have been successfully assembled to the graphene 637 nanosheets.

638 The surface morphology, particle size and microstructure of these SQDs/GNS composites were 639 further investigated by HRTEM. Fig. 15a shows a typical exfoliated nanostructure. The TEM image 640 clearly illustrates that the nanoparticles on the surface of the GNS were composed of homogeneous 641 ultrafine SnO<sub>2</sub> nanoparticles. It is obvious that they can be called SnO<sub>2</sub> QDs which were uniformly 642 distributed on the surface of the GNS in as-prepared SQDs/GNS composites. During the calcination 643 process, the GNS distributed between SnO<sub>2</sub> QDs prevented the agglomeration of these SnO<sub>2</sub> nanoparticles,<sup>119-121</sup> which was of great benefit to the formation of SnO<sub>2</sub> QDs. The SnO<sub>2</sub> nanoparticles 644 deposited on the GNS have also prevented the GNS from stacking into multi-lavers<sup>121-123</sup> in accordance 645 646 with the result that no obvious diffraction peak attributed to the GNS in the XRD pattern of the SQDs/GNS composites was observed. Fig. 15b displays a HRTEM image of the SQDs/GNS 647

composites. This shows clearly that the SnO<sub>2</sub> QDs can form in the SQDs/GNS composites. The size 648 649 distribution of SnO<sub>2</sub> QDs is estimated to be a ranging from 4 to 6 nm and their average size is calculated to be about  $4.8 \pm 0.2$  nm according to the XRD and HRTEM results. The clear lattice fringes 650 651 demonstrate that the SnO<sub>2</sub> QDs were composed of ultrafine nanoparticles. In addition, the HRTEM 652 image of a typical SnO<sub>2</sub> quantum dot shown in the inset of Fig. 15b proves that the clear crystal lattice 653 with a spacing of 0.336 nm corresponds to (110) face of the SnO<sub>2</sub> rutile phase. Therefore, it is certain 654 that the SnO<sub>2</sub> QDs distributed uniformly on the surface of the GNS were well-crystallized and was pure 655  $SnO_2$ . It is known that the embedding  $SnO_2$  QDs in a graphene nanosheet matrix would be an ideal 656 strategy to overcome the electrode material's poor cycling performance in lithium-ion batteries, which 657 is attributed to the large volume changes and serious aggregation of particles during repeated lithium insertion and extraction reactions.<sup>124</sup> However, the synthesis of a single-layer graphene nanosheet 658 matrix is very challenging due to its high surface energy.<sup>125</sup> Here, the SQDs/GNS composites with very 659 few layers of graphene nanosheets have been successfully synthesized via a facile ultrasonic method. In 660 661 addition, the approach is environmentally friendly because there is no toxic gas release during the 662 reaction. Results show that this method can effectively prevent graphene nanosheets from being 663 restacked, probably due to in situ loading of nanoparticles, which can decrease the surface energy of the graphene nanosheets.<sup>124</sup> 664



Fig. 15. TEM images of SQDs/GNS composites. (a) Low-magnification TEM image, (b) HRTEM
 image.<sup>117</sup>

668

669 In order to obtain further information on the microstructure evolution of chemical processing from 670 GO to SQDs/GNS, Raman spectroscopy was employed owing to its fine sensitivity to the crystal surface. As seen in Fig. 16, the Raman spectrum of the GNS as shown in Fig. 16a contains both D band 671  $(\sim 1335 \text{ cm}^{-1}, \text{k-point phonons of } A_{1g} \text{ symmetry})^{126}$  and G band ( $\sim 1593 \text{ cm}^{-1}, E_{2g} \text{ phonons of } sp^2 \text{ atoms})$ . 672 673 Raman spectra with characteristic D and G bands are sensitive to detects, disorder and carbon grain 674 size, and have been extensively used to characterize carbon materials. The D band in the Raman spectra 675 is an indication of disorder in the GNS originating from defects associated with vacancies, grain 676 boundaries, and amorphous carbon species. The Raman spectrum of SQDs/GNS composites as shown in Fig. 16b contains the G band at 1600.97  $\text{cm}^{-1}$  owing to the presence of isolated double bonds that 677 resonates at higher frequencies than the G band (1593.88 cm<sup>-1</sup>) of the GNS. The  $I_D/I_G$  intensity ratio is 678 a measure of the degree of disorder and average size of the  $sp^2$  domains.<sup>127-130</sup> In Fig. 16, the  $I_D/I_G$  of 679 680 the GNS is 1.09, but the value of  $I_D/I_G$  is 1.20 for the SQDs/GNS composites. Thus, there is an increased  $I_D/I_G$  intensity ratio in SQDs/GNS. When pyrolyzing GO to form GNS in N<sub>2</sub>, the mass of GO 681 sharply declined, accompanied by a large number of gas emissions, such as CO, CO<sub>2</sub> and H<sub>2</sub>O. The 682 683 graphene nanosheets are stripped by breaking the restraint of Van-Der-Waals force in this process.<sup>131</sup> 684 The emission of CO and CO<sub>2</sub> results in a large loss of carbon atoms from the graphene layers. This 685 generates a lot of carbon cavities, leading to disorder in pyrolyzed GNS. It is noticeable that an increased  $I_D/I_G$  ratio in SQDs/GNS (1.20) was observed in comparison with that of GNS (1.02) at the 686 same annealing condition. This change suggests a decrease in the average size of the  $sp^2$  domains,<sup>132</sup> 687 688 and an increase in vacancies, grain boundaries, and amorphous carbon species. Because of the reduction effect of bivalent tin, the  $Sn^{2+}$  is integrated with the oxygen atom from the carboxy group of 689

Nanoscale Accepted Manuscript

- 690 GO to form SnO<sub>2</sub> nanoparticle during stirring and ultrasonic vibration. Therefore, the Raman results are
- 691 consistent with the results of TEM and EDS, indicating the formation of SQDs/GNS composites.



# 692

693

Fig. 16. Room-temperature Raman spectra of (a) GNS and (b) SQDs/GNS composites.<sup>117</sup>

694

# 695 **4.4. BET Investigation**

Thermoanalysis results of the pure GNS and SQDs/GNS composites are given in Fig. 17a and b respectively. As can be seen in Fig. 17, thermogravimetric (TG) analysis of the pure GNS shows that the amount of weight loss (mass loss  $\approx$  16.65%) is smaller than that of the SQDs/GNS composites (mass loss  $\approx$  26.27%). The redundant portion may be attributed to loss of ammonia, physically absorbed water, and chemically bonded water in the SnO<sub>2</sub> QDs.



701

Fig. 17. TG curves of (a) GNS and (b) SQDs/GNS composites.<sup>117</sup>

703

702

 $N_2$  adsorption/desorption isotherms were employed to investigate the pore structures of the GNS and SQDs/GNS composites as shown in Fig. 18a and b, respectively. The Brunauer-Emmett-Teller (BET) surface areas of the GNS and SQDs/GNS composites are about 205 m<sup>2</sup>/g and 250 m<sup>2</sup>/g, respectively. Obviously, the BET specific surface area of the SQDs/GNS composites is bigger than that of the GNS.

708 Because of the confinement effect of the GNS, the SnO<sub>2</sub> QDs can be separated one by one. Thus, the 709 SnO<sub>2</sub> nanoparticles on the both sides of the GNS are expected to become smaller and the specific 710 surface area of the SQDs/GNS composites became larger. A hysteresis loop in the nitrogen 711 adsorption/desorption isotherms of the pure GNS was observed as shown in Fig. 18a, indicating that 712 the graphene sheets are porous. This hysteresis loop resembles type-H3 IUPAC (International Union of 713 Pure and Applied Chemistry) classification, and may result from slit-shaped pores between parallel layers.<sup>133</sup> However, a nitrogen sorption hysteresis loop in the SQDs/GNS composites was also clearly 714 715 observed in Fig. 18b. The differences are that the adsorption branch rises slowly to the higher relative 716 pressure point and then reaches to a constant value, whereas, the desorption branch moves slowly to 717 middle  $P/P_0$  signal point and decreases sharply, showing a clear type-H2 of the SQDs/GNS composites.<sup>134</sup> Furthermore, it should be pointed out that the nitrogen desorption of the SQDs/GNS 718 719 composites occurs at much lower relative pressure than that of the GNS, indicating that desorption of 720 N<sub>2</sub> molecules from the larger cavity is retarded by the smaller necks. For the above evidences, the ink-721 bottle-like pores structure can be concluded for the SQDs/GNS composites. For this structure, the desorption of N<sub>2</sub> occurs from the narrow neck and this is replenished from the larger parts of the 722 pore.<sup>135</sup> Fig. 18c shows that the main pore size distribution of the GNS is about 3-5 nm, while Fig. 18d 723 724 shows that the SQDs/GNS pore size distribution is between 2 and 4 nm. Interestingly, the GNS 725 possesses higher porosity than that of the SQDs/GNS composites. The total pore volume of the GNS  $(0.47 \text{ cm}^3\text{g}^{-1})$  is larger than that of the SQDs/GNS composites  $(0.22 \text{ cm}^3\text{g}^{-1})$ . It can be reasonably 726 727 speculated that the decreased pore volume in the SQDs/GNS composites could arise primarily from the 728  $SnO_2$  QDs, which filled the space of the close stacking of graphene nanosheets, consistent with the 729 average size of  $4.8 \pm 0.2$  nm for SnO<sub>2</sub> QDs in HRTEM observations as shown in Fig. 15b.



Fig. 18. Nitrogen adsorption/desorption isotherms of (a) GNS and (b) SQDs/GNS composites. Pore
diameter distribution patterns of (c) GNS and (d) SQDs/GNS composites.<sup>117</sup>

733

730

# 734 **4.5. Formation Mechanism**

735 Further development of the SQDs/GNS composites requires a clear understanding of their formation 736 mechanism. On the basis of our experimental results, Fig. 19 shows the formation processes of the 737 SQDs/GNS composites, which could be reasonably described by a novel model, and be separated into 738 three steps. A GO sheets solution is fabricated from pristine graphite. Many previous studies have 739 confirmed that the GO sheets are heavily oxygenated with hydroxyl and epoxide functions groups decorated on their basal planes, in addition to carbonyl and carboxyl groups located at the sheets 740 edges.<sup>133,136</sup> Firstly, the GO was well dispersed in distilled water to form a uniform GO nanosheets 741 suspension by the strong hydrophilic of these functional groups.<sup>121,137</sup> Then, the GO solution was 742

mixed with a certain concentration  $SnCl_2$  solution.  $Sn^{2+}$  ion is selectively bonded with carboxyl and 743 carbonyl through electrostatic attraction.<sup>138</sup> As a result of the strong stirring conditions, the mixture 744 solution is homogeneously dispersed and  $\text{Sn}^{2+}$  ions are also uniformly dispersed on the surface of GO 745 sheets. Under ultrasonic vibration conditions, the interlayer spacing gradually increases, and  $\text{Sn}^{2+}$  ions 746 could easily migrate into the enlarged layer. Then Sn<sup>2+</sup> ions, which are adsorbed firmly on the both 747 sides of the GO, prohibit the stacking of the GO by Van Der Waals forces. Sn<sup>2+</sup> ions are not very stable 748 due to electrostatic attraction. Finally, the precursor is heated at 500 °C in N<sub>2</sub> for 3 h. The basic Sn<sup>2+</sup> 749 750 ions should be converted uniformly to small SnO<sub>2</sub> QDs and the reduction of the GO. The GNS should 751 occur simultaneously. We think that the rational selection of the SnCl<sub>2</sub> precursor and the ultrasonic 752 method may be critical for the successful preparation of uniformly dispersed SnO<sub>2</sub> QDs on graphene substrates. First, the basic  $Sn^{2+}$  ions are firmly anchored on both sides of the GO through electrostatic 753 754 attraction, and thus enhance the conjunction stability of the hybrids. This is an essential way to get the highly dispersed  $\text{Sn}^{2+}$  ions on the GNS. Subsequently, the ultrasonic method provides a necessary 755 756 driving force for complete uniform disposition.

757 To sum up, we have reported a facile one-step ultrasonic way to synthesize  $SnO_2 QDs$  on graphene 758 nanosheets by using SnCl<sub>2</sub>·2H<sub>2</sub>O as a precursor at ambient temperature. We believe that the 759 incorporation of well-dispersed graphene-based nanosheets with the metal oxides was an effective way 760 to achieve the SnO<sub>2</sub> nucleation and the SnO<sub>2</sub> QDs growth. In this article, we provided some preliminary 761 investigations in assembling SnO<sub>2</sub> QDs to graphene nanosheets. The further detailed investigations on 762 the influence some specific synthetic conditions (such as Sn:GO ratio, calcination temperature, 763 sonication time, and pH value of the solution) on the experimental results are in progress. The 764 ultrasonic process indicated that the loading of SnO<sub>2</sub> QDs was an effective way to prevent graphene 765 nanosheets from being restacked during the reduction. The calcination process revealed that the graphene nanosheets distributed between SnO<sub>2</sub> nanoparticles also prevented the agglomeration of SnO<sub>2</sub> 766

nanoparticles, which were beneficial to the formation of SnO<sub>2</sub> QDs. The present results indicated that

Nanoscale Accepted Manuscript



769

767

768

Fig. 19. Schematic formation mechanism of SQDs/GNS composites. (a) GO, (b) electrostatic interaction between oxide functional groups of GO and  $\text{Sn}^{2+}$ , (c) graphene decorated with SnO<sub>2</sub> nanoparticles (filled circles) after the calcine treatment, and (d) TEM image of SQDs/GNS composites.<sup>117</sup>

774

# 775 **5. Conclusion and Outlook**

# **5.1. Conclusion**

In this article, we has reviewed the efforts toward understanding the synthesis strategies,
microstructual evolution, related formation mechanism, and performance evaluation of SnO<sub>2</sub> various

morphologies including SnO<sub>2</sub> QDs, thin films, and composites prepared by electron-beam irradiation,
pulsed laser ablation, and SnO<sub>2</sub> planted graphene strategies. The following conclusions can be drawn
from the present work:

782 In the section 2, the electron-beam irradiation strategies for growth behavior of  $SnO_2$  nanocrystals 783 were assessed in detail. It can be emphasized that the crystallinity of irradiated SnO<sub>2</sub> powders is better 784 than unirradiated one. The BET specific surface area of the irradiated sample is much higher, which is 785 almost doubled at 1400 kGy. The appearance of new Raman peaks is closely related to the relaxation of 786 Raman selection rule due to the EBI effects in the increasing of OD density and high concentration of 787 defects in surface site. Under 1400 kGy irradiation dose, the high-quality of SnO<sub>2</sub> ODs with diameter 788 ranging from 2 to 4 nm can be obtained. The clear lattice fringes of the QDs indicated that the crystal 789 planes of the QDs are structurally perfect and uniform. Slight misorientations changing in the QD 790 orientation demonstrated that the grain boundary must migrate toward the smaller crystal nucleus, 791 resulting in a single larger QD. The results revealed that the electron-beam irradiation method is a 792 potentially powerful technique to achieve SnO<sub>2</sub> nucleation and QD growth.

793 In the section 3, SnO<sub>2</sub> thin films prepared at different substrate temperature with interesting fractal 794 features were presented. The experimental evidence indicated that the fractal clusters with various 795 sizes, densities, and fractal dimensions were affected by different substrate temperature. The formation 796 of significant fractal features was described by (i) operation of the KrF excimer laser; (ii) production of 797 the tin dioxide plasma; (iii) cooling of the tin dioxide plumes; (iv) deposition of the tin dioxide plume 798 on the Si substrate; (v) grain rotation; (vi) formation of coherent boundary between grains followed by 799 coalescence; (vii) growth of SnO<sub>2</sub> nanocrystals along preferred crystallographic directions; and (viii) 800 formation of the fractal structure. The CO gas sensing behavior was sensitively dependent on fractal 801 dimensions, fractal densities, and average sizes of the fractal clusters. The sensitivity increases with

Nanoscale Accepted Manuscript

increasing CO concentration and decreasing fractal dimension. The Random Tunneling Junction
 Network (RTJN) mechanism can be applied to explain this gas sensing behavior.

804 In the section 4, a ultrasonic approach for the facile synthesis of SQDs/GNS composites using GO as 805 supporting materials and SnCl<sub>2</sub> as a precursor was presented in detail. The microstructural analysis 806 indicated that the GNS was exfoliated and decorated with SnO<sub>2</sub> QDs, which was dispersed uniformly 807 on both sides of the graphene. The size distribution of  $SnO_2$  QDs was estimated to be a ranging from 4 808 to 6 nm and their average size was calculated to be about  $4.8 \pm 0.2$  nm. The clear lattice fringes of 809 SQDs indicated that the crystal planes of SQDs were structurally perfect and uniform. Investigations of 810 Raman spectra showed that the SQDs/GNS composites contained both D band and G band. The 811 increased  $I_D/I_G$  intensity ratio in SQDs/GNS composites was closely related to the SQDs on the GNS. The BET specific surface area (250  $m^2/g$ ) of the SQDs/GNS was higher than that of the GNS (205 812  $m^2/g$ ). The results revealed that the incorporation of well-dispersed graphene-based sheets with metal 813 814 oxide is an effective way to achieve the SnO<sub>2</sub> nucleation and SnO<sub>2</sub> QD growth.

815

# 816 **5.2. Outlook**

In fact, the rise of low dimensional  $SnO_2$  nanostructured materials has opened up new routes for their use in various materials engineering applications. Significant impacts have also been made along the way. Of course, investigation in the  $SnO_2$  based applications still poses many critical problems that need to be addressed and understood. In particular, these crucial problems include:

(i) Development of improved methods of structurally controlled growth of semiconductor SnO<sub>2</sub>
 nanomaterials based on adjustment of structural parameters. For example, the electron-beam irradiation
 techniques should be further applied to the preparation and modification of materials.

(ii) Challenging issue for future research in the field of the fractal semiconductors is the fabrication of the related advanced controllable micro/nanostructures with novel and improved properties. Most surface, interface, and quantum structures can be good functional parts of advanced applications. For example, the formation processes of  $SnO_2$  nanocrystals and fractal clusters can be reasonably described by the fractal theory.

(iii) Achievement of better understanding of the interaction between low-dimensional SnO<sub>2</sub>-based nanomaterials and other elemental species arising from the process of doping, modification, and formation of composites, and associated micro/nanostructure changes caused by composition modulation. For example, the incorporation of well-dispersed graphene-based nanosheets with metal oxide is an effective way to achieve the nucleation and growth of the metal oxide nanocrystals and quantum dots.

There is still a lot of room for development in semiconductor  $SnO_2$  nanomaterials to realize its full potential. Future work in this field should focus on the above mentioned areas to enable commercial  $SnO_2$  micro/nanodevices to be available for practical applications.

838

# 839 Acknowledgements

The work described in this article was financially supported by the National Natural Science Foundation of China (Project Numbers: 11375111, 11428410 and 11074161), the Research Fund for the Doctoral Program of Higher Education of China (Project Number: 20133108110021), the Key Innovation Fund of Shanghai Municipal Education Commission (Project Numbers: 14ZZ098 and 10ZZ64), the Science and Technology Commission of Shanghai Municipality (Project Numbers: 14JC1402000 and 10JC1405400), the Shanghai Pujiang Program (Project Number: 10PJ1404100), and the Program for Innovative Research Team in University (Project Number: IRT13078). This work was also supported by a General Research Fund from the Research Grants Council, Hong Kong (Project
Number: CityU 119212).

849

# 850 Notes and references

- 851 (1) Z. W. Chen, J. K. L. Lai and C. H. Shek, Chem. Phys. Lett., 2006, 422, 1.
- (2) L. P. Qin, J. Q. Xu, X. W. Dong, Q. Y. Pan, Z. X. Cheng, Q. Xiang and F. Li, *Nanotechnology*,
  2008, 19, 185705.
- 854 (3) Z. W. Chen, J. K. L. Lai and C. H. Shek, J. Non-Cryst. Solids, 2005, 351, 3619.
- 855 (4) Z. W. Chen, J. K. L. Lai and C. H. Shek, *Phys. Lett. A*, 2005, **345**, 391.
- (5) Z. W. Chen, Z. Jiao, D. Y. Pan, Z. Li, M. H. Wu, C. H. Shek, C. M. L. Wu and J. K. L. Lai, *Chem. Rev.*, 2012, **112**, 3833.
- 858 (6) J. Du, J. Wang, Z. Jiao, M. H. Wu, C. H. Shek, C. M. L. Wu, J. K. L. Lai and Z. W. Chen, J.
  859 Nanosci. Nanotechnol., 2011, 11, 9709.
- 860 (7) Z. W. Chen, Q. B. Li, D. Y. Pan, Z. Li, Z. Jiao, M. H. Wu, C. H. Shek, C. M. L. Wu and J. K. L.
  861 Lai, *J. Phys. Chem. C*, 2011, **115**, 9871.
- 862 (8) Z. Ying, Q. Wan, Z. T. Song and S. L. Feng, *Nanotechnology*, 2004, **15**, 1682.
- 863 (9) P. C. Pandey, B. C. Upadhyay, C. M. D. Pandey and H. C. Pathak, *Sens. Actuators B*, 1999, 56,
  864 112.
- 865 (10) W. A. Badawy, Sol. Energy Mater. Sol. Cells, 2002, 71, 281.
- 866 (11) Y. S. Cho, G. R. Yi, J. J. Hong, S. H. Jang and S. M. Yang, *Thin Solid Films*, 2006, **515**, 1864.

- 867 (12) F. Gu, S. F. Wang, H. M. Cao and C. Z. Li, *Nanotechnology*, 2008, **19**, 095708.
- 868 (13) Z. R. Dai, Z. W. Pan and Z. L. Wang, Adv. Funct. Mater., 2003, 13, 9.
- 869 (14) X. Wang, J. Zhuang, Q. Peng and Y. D. Li, *Nature*, 2005, 437, 121.
- 870 (15) G. R. Patzke, F. Krumeich and R. Nesper, Angew. Chem. Int. Ed., 2002, 41, 2446.
- 871 (16) Y. Xia, Y. J. Xiong, B. Lim and S. E. Skrabalak, Angew. Chem. Int. Ed., 2009, 48, 60.
- 872 (17) Y. Yin and A. P. Alivisatos, *Nature*, 2005, **437**, 664.
- 873 (18) Y. Shimizu, A. Jono, T. Hyodo and M. Egashira, Sens. Actuators B, 2005, 108, 56.
- 874 (19) X. W. Lou, L. A. Archer and Z. Yang, *Adv. Mater.*, 2008, **20**, 3987.
- (20) Z. W. Chen, C. M. L. Wu, C. H. Shek, J. K. L. Lai, Z. Jiao and M. H. Wu, *Crit. Rev. Solid State Mater. Sci.*, 2008, **33**, 197.
- 877 (21) Z. Ying, Q. Wan, Z. T. Song and S. L. Feng, Nanotechnology, 2004, 15, 1682.
- 878 (22) P. C. Pandey, B. C. Upadhyay, C. M. D. Pandey and H. C. Pathak, *Sens. Actuators B*, 1999, 56,
  879 112.
- 880 (23) W. A. Badawy, J. Alloys Compd., 2008, 464, 347.
- 881 (24) Y. S. Cho, G. R. Yi, J. J. Hong, S. H. Jang and S. M. Yang, *Thin Solid Films*, 2006, **515**, 1864.
- 882 (25) F. Gu, S. F. Wang, H. M. Cao and C. Z. Li, *Nanotechnology*, 2008, 19, 095708.
- (26) L. Kőrösi, S. Papp, V. Meynen, P. Cool, E. F. Vansant and I. Dékány, *Colloids Surf. A*, 2005,
  268, 147.
- 885 (27) Y. Liu, E. Koep and M. Liu, Chem. Mater., 2005, 17, 3997.
- 886 (28) F. Paraguay-Delgado, W. Antúnez-Flores, M. Miki-Yoshida, A. Aguilar-Elguezabal, P.
  887 Santiago, R. Diaz and J. A. Ascencio, *Nanotechnology*, 2005, 16, 688.

- 888 (29) B. Cheng, J. M. Russell, W. S. Shi, L. Zhang and E. T. Samulski, *J. Am. Chem. Soc.*, 2004, **126**,
  889 5972.
- (30) T. Krishnakumar, R. Jayaprakash, M. Parthibavarman, A. R. Phani, V. N. Singh and B. R.
  Mehta, *Mater. Lett.*, 2009, 63, 896.
- 892 (31) G. C. Xi and J. H. Ye, *Inorg. Chem.*, 2010, 49, 2302.
- 893 (32) E. J. H. Lee, C. Ribeiro, T. R. Giraldi, E. Longo, E. R. Leite and J. A. Varela, *Appl. Phys. Lett.*,
  894 2004, 84, 1745.
- 895 (33) Z. W. Chen, J. K. L. Lai and C. H. Shek, Appl. Phys. Lett., 2006, 88, 033115.
- 896 (34) H. L. Zhu, D. R. Yang, G. X. Yu, H. Zhang and K. H. Yao, *Nanotechnology*, 2006, 17, 2386.
- 897 (35) J. D. Castillo, V. D. Rodríguez, A. C. Yanes, J. Méndez-Ramos and M. E. Torres,
  898 *Nanotechnology*, 2005, 16, S300.
- (36) H. M. Dahlan, M. D. Khairul Zaman and A. Ibrahim, Radiat. Phys. Chem., 2002, 64, 429.
- 900 (37) Z. Li, Z. Jiao, M. H. Wu, Q. Liu, H. J. Zhong and X. Geng, Colloids Surf. A, 2008, 313, 40.
- (38) J. Wang, J. Du, C. Chen, Z. Li, Z. Jiao, M. H. Wu, C. H. Shek, C. M. L. Wu, J. K. L. Lai and Z.
  W. Chen, J. Phys. Chem. C, 2011, 115, 20523.
- 903 (39) Z. M. Jarzebski and J. P. Marton, J. Electrochem. Soc., 1976, 123, 199C.
- 904 (40) Z. W. Chen, J. K. L. Lai, C. H. Shek and H. D. Chen, *Appl. Phys. A*, 2005, **81**, 959.
- 905 (41) B. D. Cullity, *Elements of X-ray Diffractions*, Addition-Wesley: MA, 1978.
- 906 (42) Y. J. Lin and C. J. Wu, Surf. Coat. Technol., 1996, 88, 239.

- 907 (43) J. J. Zhu, J. M. Zhu, X. H. Liao, J. L. Fang, M. G. Zhou and H. Y. Chen, *Mater. Lett.*, 2002, 53,
  908 12.
- 909 (44) Y. F. Zeng, Z. L. Liu and Z. Z. Qin, J. Hazard. Mater., 2009, 162, 682.
- 910 (45) R. L. Penn and J. F. Banfield, *Geochim. Cosmochim. Acta*, 1999, **63**, 1549.
- 911 (46) R. L. Penn and J. F. Banfield, Science, 1998, 281, 969.
- 912 (47) A. Diéguez, A. Romano-Rodríguez, A. Vilà and J. R. Morante, J. Appl. Phys., 2001, 90, 1550.
- 913 (48) W. Z. Wang, C. K. Xu, G. H. Wang, Y. K. Liu and C. L. Zheng, J. Appl. Phys., 2002, 92, 2740.
- 914 (49) S. H. Sun, G. W. Meng, G. X. Zhang, T. Gao, B. Y. Geng, L. D. Zhang and J. Zuo, *Chem. Phys.*915 *Lett.*, 2003, **376**, 103.
- 916 (50) Y. J. Jiang, L. J. Liao, G. Chen and P. X. Zhang, *Phys. Stat. Sol. B*, 1989, **156**, 145.
- 917 (51) L. Abello, B. Bochu, A. Gaskov, S. Koudryavtseva, G. Lucazeau and M. Roumvantesva, *J. Solid*918 *State Chem.*, 1998, 135, 78.
- 919 (52) K. N. Yu, Y. H. Xiong, C. S. Xiong, Y. T. Qian and Y. Hu, NanoStruct. Mater., 1995, 5, 819.
- (53) Z. L. Wang and Z. C. Kang, *Functional and Smart Materials-Structural Evolution and Structure Analysis*, Plenum Press, New York, 1998.
- 922 (54) Z. W. Pan, Z. R. Dai and Z. L. Wang, Science, 2001, 291, 1947.
- 923 (55) Z. W. Chen, J. K. L. Lai and C. H. Shek, *Phys. Rev. B*, 2004, 70, 165314.
- 924 (56) Z. R. Dai, J. L. Gole, J. D. Stout and Z. L. Wang, J. Phys. Chem. B, 2002, 106, 1274.
- 925 (57) Z. L. Wang, Adv. Mater., 2003, 15, 432.

- 926 (58) H. T. Ng, J. Li, M. K. Smith, P. Nguyen, A. Cassell, J. Han and M. Meyyappan, *Science*, 2003,
  927 **300**, 1249.
- 928 (59) U. Mikko, L. Hanna, V. Heli, N. Lauri, R. Resch and F. Gernot, *Mikrochim. Acta*, 2000, 133,
  929 119.
- 930 (60) P. Hidalgo Falla, H. E. M. Peres, D. Gouvêa and F. J. Ramirez-Fernandez, *Mater. Sci. Forum*,
  931 2005, **498-499**, 636.
- 932 (61) T. Oyabu, J. Appl. Phys., 1982, 53, 2785.
- 933 (62) Z. W. Chen, J. K. L. Lai and C. H. Shek, *Appl. Phys. Lett.*, 2006, **89**, 231902.
- 934 (63) V. Demarne and A. Grisel, *Sens. Actuators*, 1988, **13**, 301.
- 935 (64) O. Dos Santos, M. L. Weiller, D. Q. Junior and A. N. Medina, Sens. Actuators B, 2001, 75, 83.
- 936 (65) S. Shukla, S. Seal, L. Ludwing and C. Parish, Sens. Actuators B, 2004, 97, 256.
- 937 (66) S. Harbeck, A. Szatvanyi, N. Barsan, U. Weimar and V. Hoffmann, *Thin Solid Films*, 2003, 436,
  938 76.
- 939 (67) H. Windischmann and P. Mark, J. Electrochem. Soc., 1979, 126, 627.
- 940 (68) V. Lantto and P. Romppainen, Surf. Sci., 1987, 192, 243.
- 941 (69) P. Romppainen and V. Lantto, J. Appl. Phys., 1988, 63, 5159.
- 942 (70) Z. W. Chen, D. Y. Pan, B. Zhao, G. J. Ding, Z. Jiao, M. H. Wu, C. H. Shek, L. C. M. Wu and J.
  943 K. L. Lai, *ACS Nano*, 2010, 4, 1202.
- 944 (71) C. H. Shek, J. K. L. Lai and G. M. Lin, J. Phys. Chem. Solids, 1999, 60, 189.
- 945 (72) Z. W. Chen, S. Y. Zhang, S. Tan, J. G. Hou and Y. H. Zhang, *Thin Solid Films*, 1998, **322**, 194.

- 946 (73) J. Feder, *Fractal*, Plenum Press, New York, 1988, p15.
- 947 (74) H. R. Fallah, M. Ghasemi and A. Hassanzadeh, *Phys. E*, 2007, **39**, 69.
- 948 (75) Y. J. Kim, Y. T. Kim, H. K. Yang, J. C. Park, J. I. Han, Y. E. Lee and H. J. Kim, *J. Vac. Sci.*
- 949 *Tech. A*, 1997, **15**, 1103.
- 950 (76) S. Zhu, Y. F. Lu, M. H. Hong and X. Y. Chen, J. Appl. Phys., 2001, 89, 2400.
- 951 (77) D. Kim and H. Lee, J. Appl. Phys., 2001, 89, 5703.
- 952 (78) S. Zhu, Y. F. Lu and M. H. Hong, Appl. Phys. Lett., 2001, 79, 1396.
- 953 (79) L. Berthe, R. Fabbro, P. Peyre, L. Tollier and E. Bartnicki, J. Appl. Phys., 1997, 82, 2826.
- 954 (80) J. G. Hou and Z. Q. Wu, *Phys. Rev. B*, 1989, 40, 1008.
- (81) Z. W. Chen, S. Y. Zhang, S. Tan, J. G. Hou and Y. H. Zhang, *J. Vac. Sci. Technol. A*, 1998, 16, 2292.
- 957 (82) R. B. Cooper, G. N. Advani and A. G. Jordan, J. Electron. Mater., 1981, 10, 455.
- 958 (83) A. Cicera, A. Dieguez, R. Diaz, A. Cornet and J. R. Morante, Sens. Actuators B, 1999, 58, 360.
- 959 (84) D. H. Wang, R. Kou, D. Choi, Z. G. Yang, Z. Nie, J. Li, L. V. Saraf, D. H. Hu, J. G. Zhang, G.
- 960 L. Graff, J. Liu, M. A. Pope and I. A. Aksay, *ACS Nano*, 2010, **4**, 1587.
- 961 (85) A. C. Balazs, T. Emrick and T. P. Russell, *Science*, 2006, **314**, 1107.
- 962 (86) H. Zeng, J. Li, J. P. Liu, Z. L. Wang and S. H. Sun, *Nature*, 2002, **420**, 395.
- 963 (87) T. Ramanathan, A. A. Abdala, S. Srankovich, D. A. Dikin and M. Herrera-Allonso, *Nature*964 *Nanotechnology*, 2008, **3**, 327.
- 965 (88) D. S. L. Abergel and V. I. Fal'ko, *Phys. Rev. B*, 2007, **75**, 155430.
- 966 (89) F. Wang, Y. B. Zhang, C. S. Tian, C. Girit, A. Zettl, M. Crommie and Y. R. Shen, Science, 2008,

- **967 320**, 206.
- 968 (90) R. R. Nair, P. Blake, A. N. Grigorenko, K. S. Novoselov, T. J. Booth, T. Stauber, N. M. R. Peres
  969 and A. K. Geim, *Science*, 2008, **320**, 1308.
- 970 (91) M. S. Dresselhaus and G. Dresselhaus, Adv. Phys., 2002, 51, 1.
- 971 (92) M. Hirata, T. Gotou, S. Horiuchi, M. Fujiwara and M. Ohba, *Carbon*, 2004, 42, 2929.
- 972 (93) M. F. Yu, O. Lourie, K. Moloni, T. F. Kelly and R. S. Ruoff, *Science*, 2000, 287, 637.
- 973 (94) C. Berger, Z. M. Song, T. B. Li, X. B. Li, A. Y. Ogbazghi, R. Feng, Z. T. Dai, A. N.
  974 Marchenkov, E. H. Conrad, P. N. First and W. A. D. Heer, *J. Phys. Chem. B.*, 2004, **108**, 19912.
- 975 (95) Y. Zhang, J. P. Small, M. E. S. Amori and P. Kim, *Phys. Rev. Lett.*, 2005, **94**, 176803.
- 976 (96) Y. Zhang, Y. W. Tan, H. L. Stormer and P. Kim, *Nature*, 2005, **438**, 201.
- (97) K. S. Novoselov, A. K. Geim, S. K. Morozov, D. Jiang, M. I. Katsnelson, I. V. Grigorieva, S. V.
  Dubonos and A. A. Firsov, *Nature*, 2005, 438, 197.
- 979 (98) S. Stankovich, D. A. Dikin, G. H. B. Dommett, K. M. Kohlhaas, E. J. Zimney, E. A. Stach, R. D.
  980 Piner, S. T. Nguyen and R. S. Ruoff, *Nature*, 2006, 442, 282.
- (99) C. P. Peng, B. D. Chen, Y. Qin, S. H. Yang, C. Z. Li, Y. H. Zuo, S. Y. Liu and J. H. Yang, ACS
  Nano, 2012, 6, 1074.
- 983 (100) J. X. Zhu, Y. K. Sharma, Z. Y. Zeng, X. J. Zhang, M. Srinivasan, S. Mhaisalkar, H. Zhang, H. H.
- 984 Hng and Q. Y. Yan, J. Phys. Chem. C, 2011, 115, 8400.
- 985 (101) H. Kim, D. H. Seo, S. W. Kim, J. Kim and K. Kang, *Carbon*, 2011, **49**, 326.
- 986 (102) B. J. Li, H. Q. Cao, J. Shao, G. Q. Li, M. Z. Qu and G. Yin, *Inorg. Chem.*, 2011, 50, 1628.
- 987 (103) S. Q. Chen and Y. Wang, J. Mater. Chem., 2010, 20, 9735.
- (104) Z. S. Wu, W. C. Ren, L. Wen, L. B. Gao, J. P. Zhao, Z. P. Chen, G. M. Zhou, F. Li and L. M.
  Cheng, *ACS Nano*, 2010, 4, 3187.
- 990 (105) B. Wang, X. L. Wu, C. Y. Shu, Y. G. Guo and C. R. Wang, J. Mater. Chem., 2010, 20, 10661.
- 991 (106) C. Xu, X. Wang, L. C. Yang and Y. P. Wu, J. Solid State Chem., 2009, 182, 2486.

Page 57 of 58

- 992 (107) H. L. Wang, L. F. Cui, Y. Yang, H. S. Casalongue, J. T. Robinson, Y. Y. Liang, Y. Cui and H. J.
  993 Dai, J. Am. Chem. Soc., 2010, 132, 13978.
- 994 (108) L. Li, Z. P. Guo, A. J. Du and H. K. Liu, J. Mater. Chem., 2012, 22, 3600.
- 995 (109) J. W. Lee, A. S. Hall, J. D. Kim and T. E. Mallouk, Chem. Mater., 2012, 24, 1158.
- 996 (110) X. J. Zhu, Y. J. Zhu, S. Murali, M. D. Stoller and R. S. Ruoff, ACA Nano, 2011, 5, 3333.
- 997 (111) B. J. Li, H. Q. Cao, J. Shao, M. Z. Qu and H. H. Warner, J. Mater. Chem., 2011, 21, 5069.
- 998 (112) Y. Q. Zou and Y. Wang, *Nanoscale*, 2011, **3**, 2615.
- 999 (113) Y. C. Qiu, K. Y. Yan, S. H. Yang, L. M. Jin, H. Deng and W. S. Li, ACS Nano, 2010, 4, 6515.
- 1000 (114) S. J. Ding, J. S. Chen, D. Y. Luan, F. Y. C. Boey, S. Madhavi and X. W. Lou, *Chem. Commun.*,
  1001 2011, 47, 5780.
- 1002 (115) S. L. Zhang, Y. Z. Zhang, S. P. Huang, H. Liu, P. Wang and H. P. Tian, *J. Phys. Chem. C*, 2011,
  1003 114, 19284.
- 1004 (116) Z. W. Chen, Z. Jiao, M. H. Wu, C. H. Shek, C. M. L. Wu and J. K. L. Lai, *Prog. Mater. Sci.*,
  1005 2011, 56, 901.
- 1006 (117) C. Chen, L. J. Wang, Y. Y. Liu, Z. W. Chen, D. Y. Pan, Z. Li, Z. Jiao, P. F. Hu, C. H. Shek, C.
- 1007 M. L. Wu, J. K. L. Lai and M. H. Wu, *Langmuir*, 2013, **29**, 4111.
- (118) K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V.
  Grigorieva and A. A. Firsov, *Science*, 2004, **306**, 666.
- 1010 (119) H. P. Liu, D. H. Long, X. J. Liu, W. M. Qiao, L. Zhan and L. C. Ling, *Electrochim. Acta*, 2009,
  1011 542, 5782.
- 1012 (120) X. Y. Wang, X. F. Zhou, K. Yao, J. G. Zhang and Z. L. Liu, *Carbon*, 2011, 49, 133.
- 1013 (121) P. C. Lian, X. F. Zhu, H. F. Xiang, Z. Li, W. S. Yang and H. H. Wang, *Electrochim. Acta*, 2010,
  1014 56, 834.
- 1015 (122) L. S. Zhang, L. Y. Jiang, H. J. Yan, W. D. Wang, W. Wang, W. G. Song, Y. G. Guo and L. J.
  1016 Wan, *J. Mater. Chem.*, 2010, **20**, 5462.

Nanoscale Accepted Manuscript

- 1017 (123) Z. F. Du, X. M. Yin, M. Zhang, Q. Y. Hao, Y. G. Wang and T. H. Wang, *Mater. Lett.*, 2010, 64,
  1018 2076.
- 1019 (124) A. Zhong, J. Z. Wang, Z. X. Chen and H. K. Liu, J. Phys. Chem. C, 2011, 115, 25115.
- 1020 (125) Y. X. Xu, H. Bai, G. W. Lu, C. Li and G. Q. Shi, J. Am. Chem. Soc., 2008, 130, 5856.
- 1021 (126) J. F. Liang, W. Wei, D. Zhong, Q. L. Yang, L. D. Li and L. Guo, ACS Appl. Mater. Interfaces,
  1022 2012, 4, 454.
- 1023 (127) H. He, J. Klinowski, M. Forster and A. Lerf, Chem. Phys. Lett., 1998, 287, 53.
- 1024 (128) A. Lerf, H. He, M. Forster and J. Klinowski, J. Phys. Chem. B, 1998, 102, 4477.
- 1025 (129) M. Hirata, T. Gotou and M. Ohba, *Carbon*, 2005, **43**, 503.
- 1026 (130) T. Szabo, A. Szeri and I. Dekany, *Carbon*, 2005, **43**, 87.
- 1027 (131) L. S. Zhong, J. S. Hu, Z. M. Cui, L. J. Wan and W. G. Song, Chem. Mater., 2007, 19, 4557.
- 1028 (132) F. Tuinstra and J. L. Koenig, J. Chem. Phys., 1970, 53, 1126.
- 1029 (133) S. M. Paek, E. J. Yoo and I. Honma, *Nano Lett.*, 2009, 9, 72.
- 1030 (134) C. E. Salmas and G. P. Androutsopoulos, *Langmuir*, 2005, **21**, 11146.
- 1031 (135) D. Y. Zhao, J. L. Feng, Q. S. Huo, N. Melosh, G. H. Fredrickson, B. F. Chmelka and G. D.
  1032 Stucky, *Science*, 1998, **279**, 548.
- 1033 (136) L. Noerochim, J. Z. Wang, S. L. Chou, H. J. Li and H. K. Liu, *Electrochim. Acta*, 2010, 56, 314.
- 1034 (137) Y. Wang, H. C. Zeng and J. Y. Lee, Adv. Mater., 2006, 18, 645.
- 1035 (138) P. C. Lian, X. F. Zhu, S. Z. Liang, Z. Li, W. S. Yang and H. H. Wang, *Electrochim. Acta*, 2011,
- 1036 **56**, 4532.