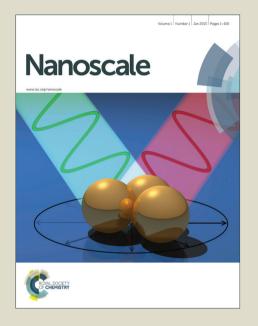
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Tuning and Understanding the Phase Interface of TiO₂ Nanoparticles for More Efficient Lithium Ion Storage†

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We demonstrate that mixed-phase anatase- $TiO_2(B)$ nanoparticles can provide an interesting interphase interface with atomic-level contact for achieving more efficient Li ion storage with high capacity and cycle life. A novel lithium storage mode - "interfacial charge storage in allomorphs" (ICSA)- plays an important role in enhancing Li ion storage.

Transition metal oxide nanoparticles (TMO NPs) have been widely investigated as anode materials for lithium ion batteries (LIBs) because they show the obvious advantages in shortening Li+ ions diffusion distance, and reducing the volume expansion/contraction during continuous charge/discharge process due to their large surface area and good conductivity.¹⁻⁶ Recent studies reveal that the interface of TMO NPs and another assistant materials (e.g. nanostructured metal/Li₂O composite) can provide the additional Li ion storage, in which Li ions are stored on Li₂O side of the interface and electrons are localized on the metal side, leading to a strong charge separation and thus the enhanced Li ion storage.⁷⁻⁹ Therefore, creating the more efficient interface between two materials will be very important for achieving high-performance lithium ion storage. However, the biggest challenge of creating the interface by use of two immiscible materials is that the interface contact is still very limited, which is not suitable for maximizing the lithium ion storage. 10-13 In this regard, the rational design of interphase interface with atomic-level contact

by the use of individual nanomaterials may provide the best choice for maximizing the lithium ion storage.

It is known that mainly TiO₂ exists in four polymorphs: anatase, rutile, brookite and TiO2(B). The structure of anatase can be considered as a stacking of one dimensional zigzag chains consisting of distorted edge-sharing TiO₆ octahedra, which provides empty zigzag channels for Li⁺ insertion to form LixTiO₂. Therefore, anatase phase is generally used as host for solid-state diffusion of Li⁺. TiO₂ (B) can also be employed as anode materials for LIBs because of its structure particularity. The ideal TiO₂(B) is composed of corrugated sheets of edgeand corner-sharing TiO₆ octahedrals. In this special framework, TiO₂(B) possesses one-dimensional parallel infinite channels, providing more open structure than other polymorphs. The diffusion of Li+ in TiO2(B) is a pseudocapacitive faradic process, which is faster than that in anatase. 14,15 Considering these, creating the interesting atomic-level phase interface between anatase and TiO2(B) may provide new opportunity for achieving more efficient lithium ion storage. We recently demonstrated that TiO2 NPs with the mixed TiO2(B) and anatase phases could form the firm connection at surface and interface via crystallographic similarity of two different nanocrystals, instead of a random aggregation of two crystals in allomorphic NPs. 16-18 This indicates that the maximized Li ion storage might be achieved at the phase boundaries of allomorphic TMO NPs. Another unanswered question for allomorphic TMO NPs for lithium ion storage is that the accurate Li storage behaviour at the phase interface (crystal interface or interphase boundaries) in allomorphic TMO anode materials is still hard to be determined because the preferred orientation and interfaces between the crystals of different phases in mixed phase materials are extremely difficult to

Herein, we demonstrate a new concept that TiO_2 NPs with mixed $TiO_2(B)$ and anatase phases, prepared from protonated titanate ($H_2Ti_3O_7$) nanosheets/graphene composite (denoted as T(AB)/G) can have the interesting interphase interface for achieving more efficient Li ion storage with high capacity and cycle life. In our previous work, $^{19-22}$ based on the phase

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transformation invariant line strain (ILS) theory and highresolution transmission electron microscopy (HRTEM) investigation, we have determined that (001) plane of anatase phase has the same orientation with (100) plane of TiO₂(B) phase, hence the (010) plane of anatase can extend across into the adjacent TiO2(B) phase. The determined interphase interface provides a very good opportunity to observe the novel interfacial Li storage behaviour in mixed phase TiO2 anode materials. Indeed, herein we found such interphase interface allows T(AB)/G exhibit higher performance for LIBs in the term of capacity and cycle life than those of the pure anatase (T(A)/G) and pure TiO₂(B) (T(B)/G) anodes. The firstprinciples simulations based on density-functional theory (DFT) shows a new charge separation mechanism- "interfacial charge storage in allomorphs" (ICSA)-should be responsible for the additional Li storage at the phase interface between the anatase and TiO₂(B), in which the dominant Ti atoms on (001) plane of anatase are beneficial for accepting electrons, followed by the conversion of Ti⁴⁺ into Ti³⁺ and the Li⁺ ions can be accommodated on the (100) plane of TiO2(B), which is occupied by O atoms. Therefore, the separation of Li⁺ ions and electrons finally results in a reduction of recombination between the Li⁺ ions and the electrons.

In order to prove our concept, T(AB)/G was prepared as follows (Fig. 1a). H₂Ti₃O₇ nanosheets were first synthesized from TiOSO₄ by a hydrothermal reaction (see the details in the experimental section), and then modified with (3-aminopropyl) triethoxysilane (APTES) coupling agent to graft -NH2 groups on the H₂Ti₃O₇ nanosheets (proved by the infrared spectroscopy of Fig. S1). Then, 1-pyrenebutyric acid (PA) on H₂Ti₃O₇ nanosheets was used to react with the grafted -NH2 groups to distribute four fused benzene rings onto the surface of H₂Ti₃O₇ nanosheets. Subsequently, the modified H₂Ti₃O₇ nanosheets could automatically attach to the G layers via strong π - π stacking interactions (proved by the infrared spectroscopy of Fig. S2). Finally, the as-prepared hybrids were calcined at 600 °C in a N₂ atmosphere, leading to T(AB)/G. For the comparison, the as-prepared hybrids was also carried out at 300 and 700 °C to obtain pure TiO2(B) (T(B)/G) and anatase (T(A)/G) anodes, respectively.

X-ray diffaction (XRD) technique was used to detect the phase-transition process. Fig. 1b and Fig. S3 show the XRD pattern of as-prepared TiO₂/G NPs, H₂Ti₃O₇/G and H₂Ti₃O₇ nanosheets. The diffraction peaks of anatase (JCPDS 21-1272) and TiO₂(B) (JCPDS 74-1940) were observed, indicating that the H₂Ti₃O₇ nanosheets were fully converted to TiO₂ NPs with mixed anatase and TiO₂(B) phases. The phase transformation was also proved by Raman spectra of the TiO2/G NPs composite (Fig. 1c). The peaks representing Ti-O vibration of both anatase and TiO2(B) crystals were in good agreement with the XRD result. In addition to phonon modes arising from the TiO2, two broad peaks at about 1342 and 1592 cm-1 were observed in the Raman spectra (the inset in Fig. 1c), assigned to the D (defect-related) peak and the G (graphitic, sp²) peak of the graphene.23-25 This indicates that the TiO2 NPs were attached on the G.

The N_2 absorption-desorption isotherms of T(AB)/G, T(A)/G and T(B)/G (Fig. 2a and Table S1) exhibit the type IV isotherms with type H1 hysteresis loop according to IUPAC, ²⁶ and

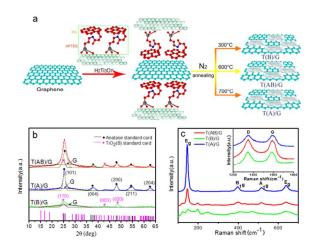


Fig. 1 (a) Schematic illustration for the preparation of TiO₂/graphene. (b) XRD patterns and (c) Raman spectroscopy of T(B)/graphene, T(AB)/graphene and T(A)/graphene nanosheets

display their similar high specific surface area (144-171 m 2 g $^{-1}$), large pore volumes (0.48-0.55 cm 3 g $^{-1}$), and narrow pore size distributions (Table S1). The pore size distributions, obtained from the Barrett–Joyner–Halenda (BJH) method, are shown in Fig. 2b. They contain two types of the pores, in which the small pores (at ~3.8 nm) are attributed to agglomeration of the graphene nanosheets, 29,30 and the mesopore with sizes at ~10.0 nm is assigned to the random aggregation of TiO₂ NPs. 31

The microstructure of the H₂Ti₃O₇/G hybrid before and after thermal treatment was characterized by using transmission electron microscopy (TEM). As shown in Fig. 3a, the H₂Ti₃O₇ nanosheets are overlaid on the top of G. The selected area electron diffraction (SAED) spots of H₂Ti₃O₇/G hybrids in Fig. 3b show two sets of spots, which can be indexed to the (0002) and (1010) planes of G, and the (200), (011), (213) and (223) planes of H₂Ti₃O₇ phase, respectively. Fig. 3c-f shows the energy dispersive spectroscopy (EDS) mapping of of Ti (green), O (blue), and C (red), proving H₂Ti₃O₇ are formed on the surface of G. Furthermore, the high-resolution TEM (HRTEM) image (Fig. 3g) of H₂Ti₃O₇/G hybrids shows a set of lattice fringes with a spacing of 0.23 nm, corresponding to the (1010) plane of G and another set of lattice fringes with a spacing of 0.19 and 0.20 nm, attributed to both (020) and (203) planes of H₂Ti₃O₇.

After the thermal treatment at 600 °C, highly crystalline TiO_2 NPs were observed on the graphene surface (Fig. 4a). The SAED spots of TiO_2 /G hybrids (Fig. 4b) show three sets of spots, which can be indexed to the $(\bar{2}110)$ and $(10\bar{1}0)$ planes of G, the (101), (200) and (120) planes of anatase phase, and the (110), (203), (600), (113) and (424) planes of TiO_2 (B) phase.

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This further indicates that the parent $H_2Ti_3O_7$ nanosheets were converted to the mixed phase anatase/ $TiO_2(B)$ NPs after calcination at 600 °C. The HRTEM image in Fig. 4c provides more subtle structure of the mixed phase anatase/ $TiO_2(B)$ NPs

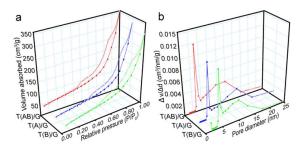


Fig. 2 (a) N_2 adsorption and desorption isotherm and (b) the pore size distribution of T(AB)/G, T(A)/G and T(B)/G.

along the [111] direction of anatase phase. A set of lattice fringes with a spacing of 0.57 nm, corresponding to the (200) plane of TiO₂(B) and another set of lattice fringes with a spacing of 0.35 nm, attributed to both (011) and (101) planes of anatase phase were observed. The Inversed Fast Fourier Transform (IFFT) image of the region enclosed by a rectangle in Fig. 4c reveals the exist of phase interface between the adjacent TiO₂(B) and anatase nanocrystals (Fig. 4d). Apparently, the (010) planes of anatase NPs have a basal spacing of 0.378 nm, close to that of the (010) planes (0.374 nm) of the adjacent TiO₂(B) NPs. With (001) plane of anatase parallel to (100) plane of TiO₂(B), the (001) plane of anatase and the (100) plane of TiO₂(B) have the same orientation so that the(010) plane of anatase extends across into the adjacent TiO2(B) phase. This observation is in good agreement with our previous finding (different TiO2 nanostructures), 18 where anatase and TiO2(B) phases in a single TiO₂ nanofiber should be combined coherently with a close crystallographic registry to reduce the Coulomb forces and minimize the dangling bonds. Furthermore, Fig. 4e-h shows the EDS mapping of TiO2/G hybrids. It is obvious that TiO2 NPs are formed on the surface of G. The microstructures of T(B)/G and T(A)/G samples were also recorded in Fig. S4 and S5, where only pure TiO2(B) or anatase TiO2 NPs was observed on graphene lavers.

To demonstrate the effects of mixed phase structure on the Li-ion insertion/extration properties, T(AB)/G, T(A)/G and T(B)/G electrodes were fabricated as follows. The working electrodes were prepared by casting a mixture of the active materials, acetylene black, and binder(polytetrafluoroethylene, PTFE) in a weight ratio of 75: 15: 10 onto pure Cu foil. Fig. 5a shows the discharge/charge profiles of three types of TiO₂ electrodes from the voltage 3.0 to 1.0 V vs Li/Li⁺. Apparently, during the initial discharge and charge processes, the capacity of the T(AB)/G electrode can achieve at 164 and 143 mA h g⁻¹, corresponding to a Coulombic efficiency (CE) of 88%, which are better than those of T(B)/G (157 and 130 mA h g⁻¹.

corresponding to a CE of 83%) and T(A)/G (146 and 124 mA h g⁻¹, corresponding to a CE of 85%), respectively. This value is better than those of pure phase T(B)/G (15%) and T(A)/G electrodes (17%). The irreversible capacity of 21 mAh g⁻¹(12%) was mainly attributed to electrolyte reductive decomposition, forming solid electrolyte interphase layer on the electrode surface in the initial discharging step. 31

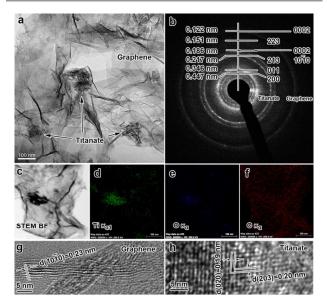


Fig. 3 (a) TEM image, (b) SAED pattern, (c) STEM DF image, (d-f) elemental mapping and (g-h) HRTEM images of H-titanate/G nanohybrids

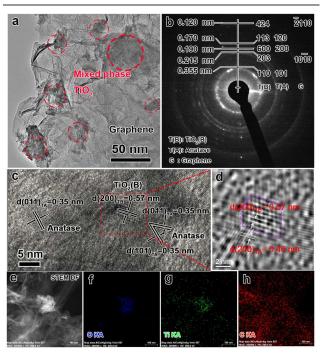


Fig. 4 (a) TEM image, (b) SAED pattern, (c) HRTEM image, (d) IFFT image of the selected area in plane c, (e) STEM dark-field (DF) image, (f-h) elemental mapping of oxygen, titanium and carbon of T(AB)/G nanohybrids.

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Clearly, the mixed phase T(AB)/G electrode exhibited the better specific capacity and irreversibility during the electrochemical lithium insertion/extraction.

Furthermore, the discharge capacities of the TiO₂/G electrodes decrease almost linearly in the first 5 cycles and then are stabilized over the remaining cycles (Fig. 5b). After 40 galvanostatic cycles at a high current density of 1000 mA g⁻¹, the discharge capacities of T(AB)/G, T(A)/G and T(B)/G electrodes were stabilized to be around 134, 123, and 109 mA h g⁻¹, respectively. The capacity losses of T(AB)/G, T(A)/G and T(B)/G electrodes were 15.8, 18.7, and 30.5%, respectively. Obviously, the mixed phase T(AB)/G exhibits better cycling stability in comparison with the pure phase samples. Fig. 5c shows the rate capability of the three electrodes. T(AB)/G shows much better rate ability with the specific discharge capacity being around 164 (1 C), 154 (2 C), 141 (4 C), 135 (6 C) and 130 mA h g⁻¹ (8 C), respectively than the T(A)/G electrode (143, 127, 116, and 106 mA h g⁻¹ at the same current densities) and T(B)/G electrode (122, 105, 93, and 82 mA h g⁻¹ at the same current densities). That means the the mixed phase T(AB)/G electrode exhibits the best capacity retention at high discharge/charge rate, indicating that the mixed phase structure of T(AB)/G facilitates the fast Li-ion diffusion high electron transport. As a comparison, the T(AB)/G exhibits much better rate capability than the pure T(AB) nanosheets (see Fig. S6). This is because the graphene can improve the electron conductivity of TiO2 and thus improve the rate capability. Given T(AB)/G, T(A)/G, and T(B)/G show the similar crystallite size and specific surface area, we think that the size and surface area should be not responsible obvious difference between the mixed phase T(AB)/G and for the difference on lithium storage ability for these three samples. The most the pure phase T(A)/G (or T(B)/G) is the phase interface between anatase and TiO₂(B) in T(AB)/G. Specifically, the (001) plane of anatase and the (100) plane of TiO2(B) have the same orientation to connect the anatase and TiO₂(B) together tightly. Herein, a novel interfacial charge storage in allomorphs (ICSA) was proposed to explain the origin of higher performance of mixed phase TiO2(B)/anatase anode in lithium storage. The (100) plane of TiO₂(B) is occupied by O atoms only (Fig. 6a), and the dominated atoms at the (001) plane of anatase are Ti atoms (Fig. 6b). The delicate atomic arrangement at the phase interface between TiO2(B) and anatase can facilitate an additional Li storage via the charge separation. As depicted in Fig. 6c, the Ti atoms at the (001) plane of anatase act as acceptor to capture electrons, followed by the conversion of Ti⁴⁺ into Ti³⁺. ^{33,34} Meanwhile, besides the insert in the parallel tunnels perpendicular to the (010) plane of the TiO₂(B) phase, the Li⁺ ions can be accommodated on the (100) plane of TiO₂ (B), which is occupied by O atoms. Therefore, the separation of Li⁺ ions and electrons results in a reduction of recombination between the Li⁺ ions and the electrons, leading to the enhanced specific capacity and rate performance of the mixed phase electrode.

To verify the ICSA mechanism, we carried out the first-principles calculations based on density-functional theory (DFT) using a model slab with $\text{TiO}_2(B)$ layer and anatase layer. The

two-dimensional slab contains several layers of Ti-terminated anatase and O-terminated $TiO_2(B)$, respectively. The first-principles simulations of the Li storage in the $TiO_2(B)$ /anatase interface are performed within the local density approximation and ultrasoft pseudopotential method in the CASTEP code. The energy cutoff for plane wave basis is 450 eV and the k-point for the slab is $5 \times 3 \times 1$. The leftmost layer atoms of T(A) and the rightmost layer atoms of T(B) were fixed to their positions in bulk T(A) or T(B) when the slab was optimized. We located a Li atom on various sites of T(B), and found that the tunnel T1 site is more stable for Li ion storage than the tunnel T2 sites (green zone in Fig. 6d&e). According to the present model, the Li⁺ ions would be stored on the T(B) side of the interface, while there is an electron transfer from the Li-doped T(B) to T(A), resulting in a charge separation.

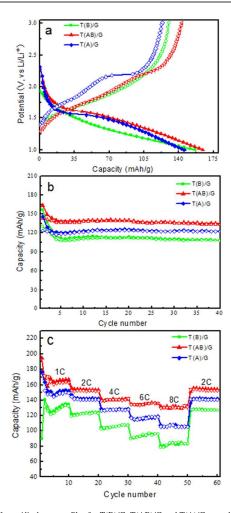


Fig. 5 (a) Charge/discharge profiles for T(B)/G, T(AB)/G and T(A)/G nanosheets between 1.0 to 3.0 V at a current density of 1000 mA g–1 at room temperature, (b) Cycling performance of T(B)/G, T(AB)/G and T(A)/G nanosheets to 40 cycles at 1000 mA g–1 charge/discharge rate, (c) Specific capacity of nanosheets at different C rates: 1 C, 2 C, 4 C, 6 C, and 8 C (T(C)) Table 1.0 (and T(C)) C charge, solid: discharge).

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Using the optimized structure of the T(A)/T(B) interface, we calculated the Mulliken charge distributions in both the slabs with and without the Li-doping interface. The charges of the T(A)/T(B) interface atoms were obtained, as listed in Table 1. The extra Li atom considered as an electron donor loses 1.33e. On the T(A) (001) double-atom layer, there is an average increase of electrons of $\sim 0.16e$ on the Ti atom and $\sim 0.03e$ on the O atom. Therefore, the T(A) interface layer can serve as an electron acceptor. Meanwhile, on the T(B) (100) interface layer, there is also an average increase of electrons of $\sim 0.2e$ on the Ti atom and $\sim 0.09e$ on the O atom. However, Li⁺ is also located at the side of T(B). Therefore, the T(B) and Li lose electrons as a whole, and the T(A) gains electrons, resulting in a charge separation at the T(A)/T(B) interface. This result is in

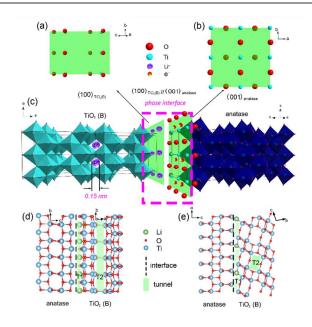


Fig. 6 A proposed interfacial storage model in mixed phase TiO₂(B)/anatase system: (a) 2-D atomic profile of (100) plane of TiO₂(B). (b) 2-D atomic profile of (001) plane of anatase. (c) Charge separation at the phase interface between TiO₂(B) and anatase crystals. Optimized atomic arrangement of the model slab with Li atom in the interface between TiO₂(B) and anatase phase: (d) View along [100] direction of anatase and [001] direction of TiO₂(B), (e) View along [010] direction of anatase and TiO₂(B).

 $\label{thm:charge} \textbf{Table 1} \ Mulliken \ charge \ distribution \ for \ TiO_2(B) \ / \ anatase \ interface \ with \ and \ without \ Li \ atom \ insertion.$

Slab	Interfacial Li	O-av ^a	Ti-av ^a	Doped Li ^b
TiO ₂ (B)	undoped	-0.64	1.06	/
$TiO_2(B)$	doped	-0.73	0.86	1.33
Anatase	undoped	-0.66	0.91	/
Anatase	doped	-0.69	0.75	/

 a O-av and Ti-av are average effective charge per atom on O and Ti atom in the interface layer of TiO₂(B) and TiO₂(A). b Doped Li is effective charge on lithium per atom.

Good agreement with the proposed ICSA mechanism. Furthermore, a typical Nyquist plot of T(AB)/G, T(A)/G, T(B)/G and T(AB) electrodes was given in Fig. S7. Obviously, T(AB)/G displays the lowest resistance, further supporting our ICSA mechanism.

To summarize, we demonstrate a multi-step process for the controlled synthesis of T(AB)/G through annealing the H₂Ti₃O₇/G composite nanosheets in N₂ at high temperature, which provide an interesting interphase interface with atomiclevel contact for achieving more efficient Li ion storage with high capacity and cycle life. We found that the mixed phase T(AB) exhibits higher specific capacity and better rate performance than those of the pure T(A) and T(B). The highperformance Li ion storage ability of the mixed phase electrode should be ascribed to the additional Li storage at the phase interface between the T(B) and T(A) (novel ICSA mechanism). This has been well verified by the first-principles calculations based on DFT, in which the interfacial (001) plane of anatase acts as acceptor to capture electrons, and the (100) plane of TiO₂ (B), occupied by O atoms can accommodate theLi⁺ ions. The proposed ICSA mechanism open a new avenue for tuning and understanding new multi-phase materials for future more efficient LIBs.

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