

# Nanoscale

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

*Accepted Manuscripts* are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this *Accepted Manuscript* with the edited and formatted *Advance Article* as soon as it is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.

# Closed-Loop ARS Mode for Scanning Ion Conductance Microscopy with Improved Speed and Stability for Live Cell Imaging Applications

Goo-Eun Jung<sup>\*,†,‡</sup>, Hanaul Noh<sup>\*,†</sup>, Yong Kyun Shin<sup>†</sup>, Se-Jong Kahng<sup>‡</sup>, Ku Youn Baik<sup>§</sup>, Hong-Bae Kim<sup>||</sup>, Nam-Joon Cho<sup>#,+</sup>, Sang-Joon Cho<sup>#,†,‡</sup>

<sup>†</sup>Research and Development Center, Park Systems, Suwon 443-270, Korea

<sup>‡</sup>Department of Physics, Korea University, Seoul 136-713, Korea

<sup>§</sup>Department of Electrical and Biological Physics, Kwangwoon University, Seoul 139-701

<sup>||</sup>College of Agriculture and Life Sciences, Seoul National University, Seoul 151-921, Korea

<sup>+</sup>School of Materials Science and Engineering, Nanyang Technological University, 50 Nanyang Avenue 639798, Singapore

<sup>‡</sup>Advanced Institute of Convergence Technology, Seoul National University, Suwon 443-270, Korea

\*Goo-Eun Jung and Hanaul Noh contributed equally to this work

#Corresponding Author

E-mail: njcho@ntu.edu.sg and msjcho@snu.ac.kr

## **Abstract**

Scanning ion conductance microscopy (SICM) is an increasingly useful nanotechnology tool for non-contact, high resolution imaging of live biological specimens such as cellular membranes. In particular, approach-retract-scanning (ARS) mode enables fast probing of delicate biological structures by rapid and repeated approach/retract of a nano-pipette tip. For optimal performance, accurate control of the tip position is a critical issue. Herein, we present a novel closed-loop control strategy for the ARS mode that achieves higher operating speed with increased stability. The algorithm differs from that of most conventional (i.e., constant velocity) approach schemes as it includes a deceleration phase near the sample surface, which is intended to minimize the possibility of contact with the surface. Analysis of the ion current and tip position demonstrates that the new mode is able to operate at approach speeds of up to 250  $\mu\text{m/s}$ . As a result of the improved stability, SICM imaging with the new approach scheme enables significantly improved, high resolution imaging of subtle features of fixed and live cells (e.g., filamentous structures & membrane edges). Taken together, the results suggest that optimization of the tip approach speed can substantially improve SICM imaging performance, further enabling SICM to become widely adopted as a general and versatile research tool for biological studies at the nanoscale level.

**Keywords:** Scanning Probe Microscopy, Scanning Ion Conductance Microscopy, Biological Imaging, Approach-Retract-Scanning (ARS) mode, Live Cell Imaging

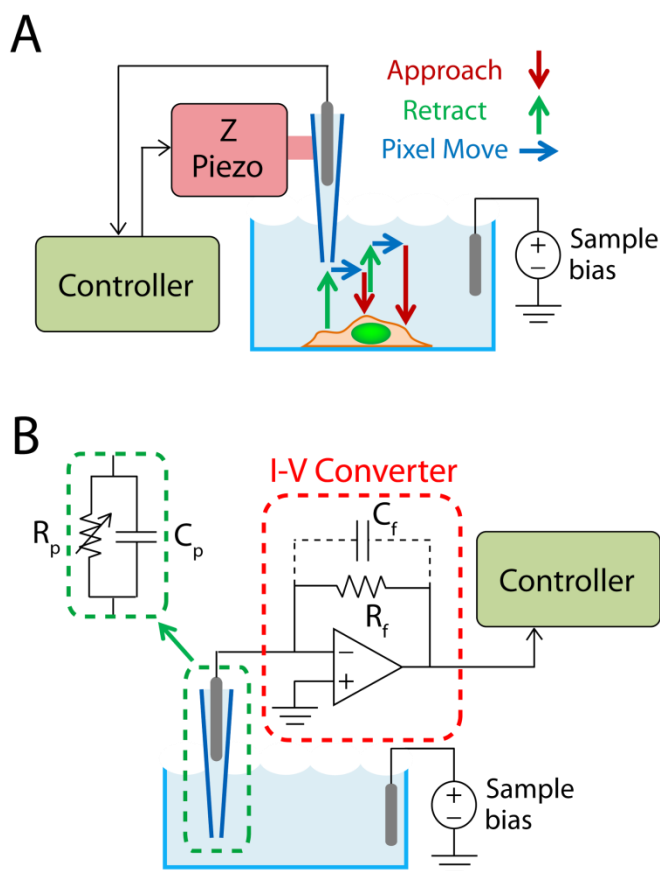
## **Introduction**

With the growth of nanotechnology solutions in medicine, there is increasing demand for understanding of the basis for diseases at the cellular and sub-cellular levels. Improved capabilities to scrutinize in detail the membrane properties of individual cells would be a significant advance towards this goal because many important cellular activities are mediated via cellular membranes. Although much effort in this direction has been made with atomic force microscopy (AFM), successful characterization of membrane topography with AFM remains challenging due to the highly contoured shape and soft nature of most biological specimens<sup>1-3</sup>.

On the other hand, scanning ion conductance microscopy (SICM)<sup>4,5</sup> has become a powerful alternative tool<sup>6</sup> for nanoscale investigations of cell membranes and associated structural features<sup>7</sup>. The SICM device employs a scanning nano-pipette<sup>8,9</sup> to probe the sample surface, ideally without physical contact, by sensing changes in the nano-pipette's ion current near the sample surface (i.e., typically within the tip diameter range<sup>10</sup>). This high resolution, non-contact imaging capability and suitability for liquid measurements makes the SICM device an ideal technique for non-invasive<sup>11</sup> and nanoscale investigation of live cells<sup>12-15</sup> with the vertical approach mode (including versions such as hopping<sup>16</sup>, backstep<sup>17,18</sup>, standing approach<sup>19,20</sup>, and approach-retract- scanning (ARS)<sup>21</sup>). In the vertical approach mode, a nano-pipette repeatedly approaches and retracts from the surface across each imaging pixel as shown in **Figure 1A**. By contrast, traditional direct current (DC)<sup>4</sup> and alternating current (AC)<sup>22,23</sup> modes continuously scan across the sample surface while maintaining a setpoint distance. As such, the vertical approach operating scheme is advantageous for imaging soft and delicate samples because it decouples the lateral and vertical motions of the pipette, minimizing possible tip-sample interactions. A critical drawback of the vertical approach mode, however, is that the addition of

an approach step at each imaging pixel significantly increases the total time for image collection. While increasing the approach speed can improve the performance, there exists a finite current signal delay generated by the pipette capacitance<sup>24</sup>, as well as the parasitic capacitance of the current-voltage converter as presented in **Figure 1B**. Combined with the fairly short sensing range of the SICM device, these constraints limit the performance of existing vertical approach modes<sup>16-21</sup>.

Herein, we describe a control strategy to improve stability and tip approach speed in the ARS mode, which we refer to as the closed-loop (CL) ARS mode. The key differences between the CL-ARS mode and existing hopping/open-loop (OL) ARS modes are the novel use of differential servo gain feedback control for approach speed optimization. In essence, this approach method enables fast approach with rapid deceleration near sample surface, elevating the overall approach speed with reduced tip-sample interaction. In this study, we first measure the ion current signal delay that is inherent in SICM in order to understand its physical limitation. Based on the current delay and the non-linearity of the current-distance relationship, we introduce a differential servo gain profile tailored to SICM in order to achieve high approach speed with rapid deceleration using feedback control. The current overshoot measured in the CL-ARS mode is within an acceptable range for approach speeds of up to 250  $\mu\text{m/s}$ . We verify improved ARS performance by comparing SICM images of human cancer H460 cells and live rat cardiomyocytes and show that soft and subtle biological features can be visualized with much improved structural clarity in the CL-ARS mode, indicating improved stability and non-contact condition.



**Figure 1. Schematic illustration of SICM-ARS mode. (A)** Operational logic of the ARS mode. The pipette repeatedly approaches the sample and then retracts at each imaging pixel position. **(B)** Simplified I-V converter showing the SICM pipette tip modeled as an equivalent circuit consisting of a parallel resistor and capacitor. The feedback resistor has a small parasitic capacitance associated with high resistance.

### Measurement Strategy

**Open-loop (OL) ARS.** Conventional vertical approach/open-loop (OL) ARS mode was implemented and tested as a reference for evaluation of the proposed CL-ARS method. Both ARS methods consisted of three steps that were repeated in succession: approach, retract, and move pixel (see **Figure 1A**). Besides the approach step, identical implementation and parameters were used for experimental consistency. In addition to the constant approach speed and 2% reduction from the normalized saturated current (i.e., current far away from the surface) as the approach setpoint, the approach step of OL-ARS consisted of a 1% mid-setpoint at which point

the nano-pipette was slightly raised and then continued approaching the sample with half of the original speed. Starting with a fairly slow velocity ( $\leq 75 \mu\text{m/s}$ ), the approach speed was incrementally raised until the ion current response resulted in a significant drop near the sample surface (10% or more), which could result in tip-sample contact and damage the pipette. In the retract step, the height and speed for both methods were set at  $2 \mu\text{m}$  and  $480 \mu\text{m/s}$ , respectively.

**Closed-loop (CL) ARS.** The approach speed of the ARS mode is fundamentally limited by a finite delay in ion current signal and piezo response time. Due to these natural limitations in SICM devices, it is very important to optimally utilize the relationship between the tip-sample distance and ion current signal (i.e. current-distance  $I$ - $D$  curve, see **Figure 2**<sup>16</sup>) to achieve higher imaging speed. The main challenges to optimizing the approach velocity are: 1) ion current noise, typically near 0.3% of the normalized saturated current  $i_{\infty}$ , and 2) severe non-linearity in the  $I$ - $D$  curve (up to 0.5% in **Figure 2B**). Since the current drop after the non-linear region implies a very narrow safety distance between the SICM pipette and the sample, smooth deceleration is needed as soon as any subtle change of current is detected in order account for slow current and piezo response. In this regard, existing vertical approach methods that use constant approach speed with set-point higher than 1 or 2% due to large ion current noise fail to sufficiently decelerate the pipette.

The CL-ARS algorithm uses a modified version of classical proportional-integral (PI) feedback control theory with differential servo gain for velocity profile tuning specific to SICM characteristics. Contrary to the vertical approach/OL-ARS methods that move down at a constant speed until set-point current (typically 1% or 2%) is reached, CL-ARS uses the difference between set-point and measured ion current at the present time as servo feedback error. **Figure 2A** shows the servo gain profile with respect to the normalized ion current. The piece-wise shape

of the servo gain is designed in consideration of the  $I$ - $D$  curve non-linearity, as shown in **Figure 2B**. For most of the approach length  $L$  that is far from the sample, which we refer to as coarse approach (shown in **Figure 2C**), ion current does not decrease from the initial saturated current. However, as pipette nears the sample, ion current starts to decrease drastically, beyond approximately 0.5% shown in **Figure 2B**. We refer to this region as the fine approach. The proposed CL ARS uses the maximum servo gain to amplify the servo error and therefore maintain a high approach speed for the coarse approach, followed by a rapid attenuation of the servo gain to the minimum value of 1 in order to significantly slow down the pipette for the fine approach. Between a very short distance (i.e. 10 - 20 nm) from 0.5% and 2 %, the pipette travels with minimum servo gain to ensure smooth approach near the sample.

Numerically, the error signal for feedback control used by CL-ARS is given by

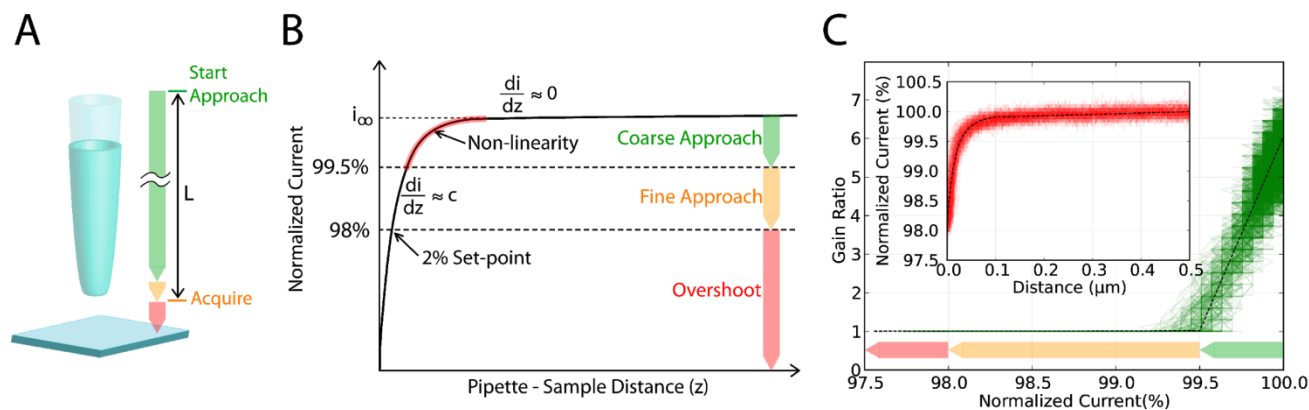
$$e(t) = g(i(t)) \times (i(t) - i_{2.0})$$

where  $i(t)$  is the ion current measured at time  $t$ ,  $i_{2.0}$  is the current corresponding to 2.0% setpoint (or 98% of the saturated current), and  $g$  is the differential servo gain introduced to account for the non-linearity in the  $I$ - $D$  curve. The servo gain is determined based on the measured current during the approach and is defined by the following piece-wise linear function:

$$g(i) = \begin{cases} 1 & : \frac{i}{i_{\infty}} < i_{0.5} \\ 1 + (g_m - 1) \times \frac{i - i_{0.5}}{i_{\infty} - i_{0.5}} & : \frac{i}{i_{\infty}} \geq i_{0.5} \end{cases}$$

where  $i$  is the measured current,  $i_{\infty}$  is the saturated current,  $i_{0.5}$  is the current corresponding to 99.5% of  $i_{\infty}$ , and  $g_m$  is the maximum value of the servo gain to be used during the approach. An experimentally measured  $I$ - $D$  curve and the corresponding servo gain profile are shown in **Figure 2C**. Overall, the proposed approach method based on differential gain feedback control

offers two benefits: 1) Robustness to ion current noise, allowing smaller current reduction to be reliably used as feedback signal, and 2) Rapid deceleration by using smaller servo gain near sample ( $i < i_{0.5}$ ), allowing higher approach speed to be used when far away from sample ( $i \sim i_{\infty}$ ).



**Figure 2. Measurement strategy for the CL-ARS mode.** (A) Schematic of Proposed CL-ARS algorithm. (B) Normalized ion current through pipette,  $i/i_{\infty}$ , as a function of pipette-surface distance,  $z$ . (C) Measured current-distance curve (upper) and gain multiplier profile with respect to ion current (lower).

## Materials and Methods

**Buffer Solution.** Deionized water (resistivity:  $18.2 \text{ M}\Omega\cdot\text{cm}$  at  $25^{\circ}\text{C}$ ) was obtained from a Milli-Q water (Millipore Corp.) and used to prepare all buffer solutions. phosphate buffered saline (PBS) (10 mM phosphate buffer, 137 mM sodium chloride, and 2.7 mM potassium chloride) was used as the standard electrolyte solution for the SICM measurements. **SICM Probe** Nano-pipettes were pulled from borosilicate capillaries of 1.0 mm outside diameter, 0.58 mm inside diameter (Warner Instruments) using a P-2000  $\text{CO}_2$  laser puller (Sutter Instrument). Fabricated nano-pipette tips have a nominal inner diameter of 100 nm (laser puller pulling parameters: Heat 265, Fil 4, Vel 30, Del 225, Pul 150). The nano-pipette tip and petri dish used in experiments were also filled with PBS solution.

**H460 cell culture.** H460 cells were maintained in high glucose Dulbecco's Modified Eagle Medium (DMEM; SH30243.01, Hyclone) supplemented with 10% fetal bovine serum (FBS;



SH30979.03, Hyclone) and 1% penicillin/streptomycin (15140, Gibco) at 37°C in humidified air with 5% CO<sub>2</sub>. Cells were cultured on a petri dish (20035, SPL, diameter: 35 mm) for 24 hours before experiment in order to promote cell adhesion. An applied electric field is known to cause ultra-structural changes and create nano-pores in the plasma membrane of H460 cells<sup>25</sup>. In order to observe such effects, nine pulses of an electrical stimuli with a 1kV/cm electric field and 100 μs pulse width were applied to cells by using an electroporator (ECM830, Havard Apparatus, Inc.) and petri pulse electrodes (450130, Havard Apparatus, Inc.). According to a general electroporation protocol<sup>25</sup>, 1 ml of serum-free DMEM media was used to enhance the electric pulse effects. Within the time scale of several seconds, cells were fixed with Karnovskys fixative (18505 & 18420, TedPella Inc.) for 1 hour and with 1% osmium tetroxide (184450, TedPella Inc.) for the following 30 minutes. Dulbecco's phosphate buffered saline (LB001-04, Welgene) was used to wash cells between steps. For scanning electron microscope (SEM) observation, cells were further dehydrated by a series of ethanol-water mixtures and dried in hexamethyldisilazane and then coated with platinum (Bal-TEC/SCD 005).

**Rat cardiomyocyte cell culture.** Rat cardiomyocytes were isolated from the hearts of two day-old rats, and then the cells were maintained in Dulbecco's Modified Eagles Medium (DMEM; SH30243.01, Hyclone) supplemented with 5% fetal bovine serum (FBS; SH30979.03, Hyclone), 200 μg/ml streptomycin, 200 units/ml penicillin and non-essential amino acids (Gibco). 100 μg/ml genitacin (G418; Gibco) was added in order to inhibit fibroblast growth. Cells were maintained at 37°C in a humidified air with 5% CO<sub>2</sub>. Cells were used within 2 days (discrete cells) and cultured on petri dishes<sup>26</sup>.

**Instrumentation.** Experiments were performed using an NX-Bio SICM instrument (Park Systems Corp.). A schematic of the NX-Bio instrumentation is shown in **Figure 1**. The nano-

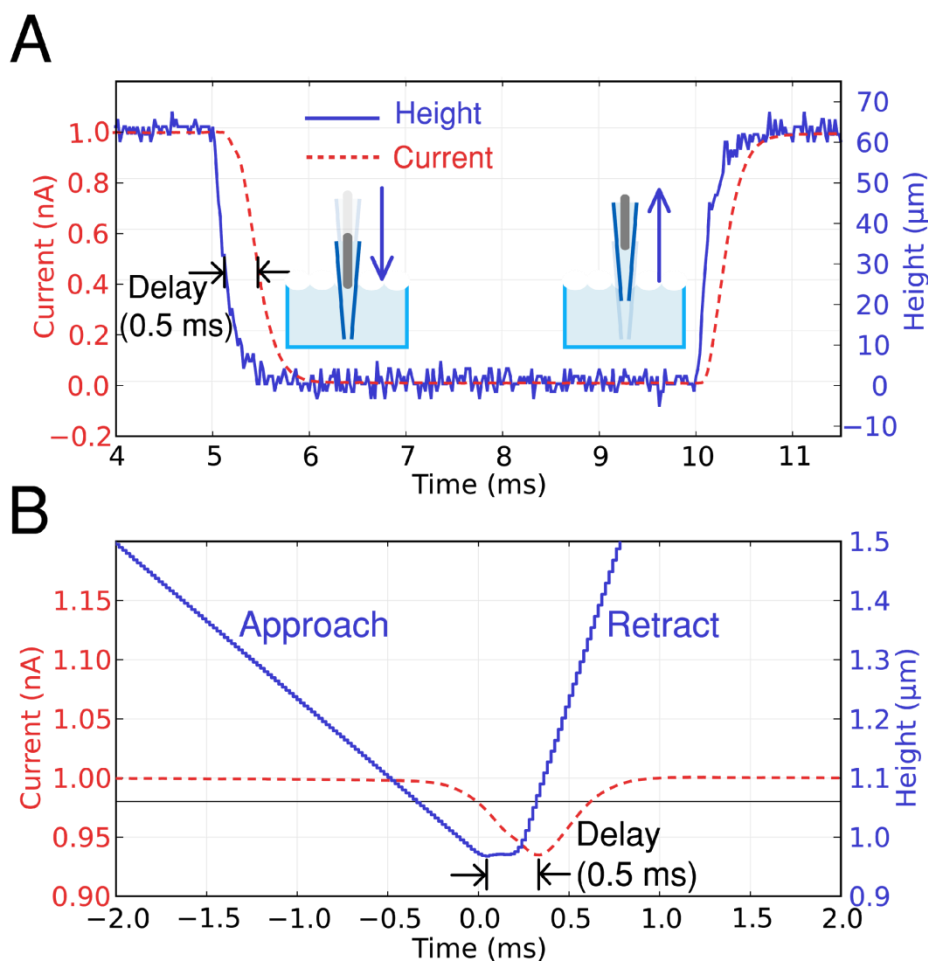
pipette tip was mounted on a piezo scanner (moving range: 25  $\mu\text{m}$ ) for movement in the Z direction. Samples were placed in a small petri dish filled with PBS solution. The petri dish was rested on a XY scanner (100 $\times$ 100  $\mu\text{m}$ ) for SICM imaging. The entire SICM microscope was mounted above an inverted optical microscope (Ti-U; Nikon) stage for positioning of the pipette relative to the sample. For conduction of the ion current, an Ag/AgCl electrode immersed in the bath solution was grounded and served as the reference electrode for all applied potentials. A separate Ag/AgCl electrode was placed inside the nano-pipette tip and biased in order to generate ion current through the tip. For the feedback signal, the ion current was amplified with an analog current-to-voltage converter, as shown in **Figure 1B**. Prior to measurement and imaging, the bias potential was adjusted to yield 1.0 nA of current and kept constant throughout, which roughly corresponded to 100 mV. In addition, cells were imaged by using a field-emission scanning electron microscope (SEM; SUPRA55VP, Carl Zeiss).

## **Results and Discussion**

**Pipette Response.** Unlike the conventional DC and AC modes, the scan speed of the ARS mode is critically determined by the approach speed, which in turn is largely dictated by the response speed of the ion current and the piezo-actuator. Understanding the sources and magnitudes of the time delay present in SICM instrument can thus give an indication of the maximum approach velocity achievable with ARS mode. In order to measure the current delay produced by the pipette capacitance and I-V converter, step signal was used to modulate the tip in the vertical direction on a glass surface and the change in ion current was recorded (**Figure 3A**). The ion current response was additionally measured during OL-ARS operation for further investigation (**Figure 3B**). Previous results from Novak *et al.*<sup>7</sup> have pointed to a slow Z-axis piezo actuator as the bottleneck of the delayed current response (up to 2 ms). With the use of fast secondary piezo-

actuator, they demonstrated that the speed bottleneck due to slow piezo response can be compensated.

In our experiments, however, the ion current delay was found to contribute equally significantly to the overall approach speed limit, generating at least 0.5 ms of time delay with respect to the probe height (see **Figure 3A**). The time delay due to piezo-actuator in our experiment was also measured at approximately 0.5 ms. This means that the use of secondary fast piezo-actuator is not sufficient to compensate for the slow ion current response of SICM. Same result was consistently observed during OL-ARS operation, where the tip more or less stopped at a 2% setpoint height but the current kept diminishing for at least 0.5 ms, indicating that the tip surpassed the setpoint for nearly 100 nm.

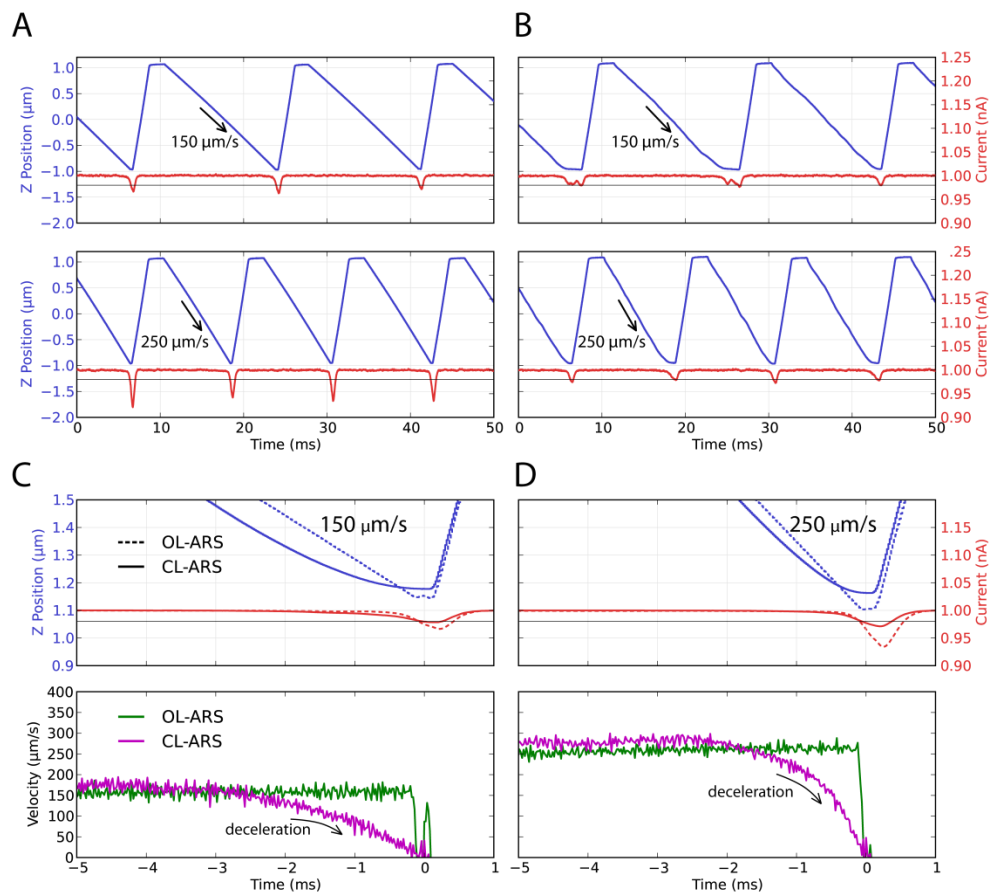


**Figure 3. Response speed of the SICM current signal.** (A) Step signal applied to the Z-axis piezo with 10 ms time interval. (B) Current and height signals measured during operation of the ARS mode with constant approach speed.

**ARS mode comparison.** We next investigated the height and current signals for OL-ARS and CL-ARS modes with similar approach speeds in order to compare the relative stability of the two methods (Figure 4) on a flat polydimethylsiloxane (PDMS) surface. For a range of approach speeds much lower than the response delay of the ion current previously discussed, maintaining a 2% setpoint current is rather straightforward for both methods. However, as the speed increases to approximately 150  $\mu\text{m/s}$ , OL-ARS measurements consistently generates an overshoot current of nearly 1%, resulting in a 3% overall current reduction close to the sample surface (**Figure 4A, upper panel**). If the speed is increased to 250  $\mu\text{m/s}$  (**Figure 4A, lower panel**), then the current overshoot increases to approximately 7%. Since the ion current contains a time delay of greater than 0.5 ms, the measured current overshoot of 7% indicates a significantly higher chance of tip-sample contact and potentially tip and/or sample damage. For CL-ARS measurements at 150  $\mu\text{m/s}$  (**Figure 4B, upper panel**), fluctuation of the current near the sample surface can be observed. This indicates that the tip-sample distance has reached the sensing range of the nanopipette tip, where small fluctuations from the Z-axis piezo actuator can be captured in the ion current signal. Nevertheless, the current overshoot remains very close to the 2% setpoint for 150  $\mu\text{m/s}$  speed and within 3% for the higher approach speed of 250  $\mu\text{m/s}$ , which is much smaller than OL-ARS measurements for similar speeds (**Figure 4B, lower panel**). While the single approach current and height signals provide some indication of the stability of a specific ARS mode of interest, the absolute current reduction varies due to the noise contained in the ion current signal (**Figure 2C, inset**)<sup>27</sup>. In our experiment, the current noise generated was as large as 0.3%. Therefore, we repeated the ARS approach experiments for a sufficient period of time

(2.5 s) and analyzed the average height, current and approach velocity signals during this period in order to mitigate the effect of current noise (**Figure 4C,D**). Examining the pipette interaction near the sample surface (between -5 ms and 0ms when acquisition occurs), the differences between OL-ARS and CL-ARS measurements can be clearly discerned. The approach velocity for OL-ARS measurements is constant throughout the entire approach phase, regardless of the chosen approach speed. As a result, the time interval between the initial current reduction and setpoint (2%) current is smaller than 1ms for a relatively slow (150  $\mu\text{m/s}$ ) approach speed. Although 1 ms is sufficient time to detect current changes for commonly used setpoint ranges between 2-3%<sup>28</sup>, this result has important implications for improving ARS performance without using feedback control. For instance, safely increasing the constant approach velocity would require use of a setpoint value lower than 2% (e.g., 1% or below). However, due to the large ion current noise (as much as 0.3% (**Figure 2C, inset**)), the resulting sampling height inevitably becomes particularly sensitive to noise and therefore significant measurement error can be introduced. On the other hand, if a certain threshold value (e.g., between 1% and 2%) prior to the 2% setpoint is used to lower the approach speed, then the current delay becomes greater than the fine approach interval and the tip would retract without entering the fine approach. This was experimentally observed when a threshold of 1% was used in order to retract the tip slightly and re-approach with a lower velocity (**Figure 4C**), but the heights at 1% and 2% current are almost equal (**Figure 4C, upper panel**). For 250  $\mu\text{m/s}$  approach speed (**Figure 4D, upper panel**), a backstep almost does not even take place. Therefore, deceleration must occur much quicker than the location at 1% setpoint current in order to take effect. By contrast, the velocity profile of the CL-ARS measurement approach begins deceleration approximately 4 ms (**Figure 4C, lower panel**) and 2 ms (**Figure 4D**) prior to the setpoint current for 150  $\mu\text{m/s}$  and 250  $\mu\text{m/s}$  approach

speeds, respectively, which indicates smooth transition between the coarse and fine approaches. The velocity profile also shows continuous deceleration as intended, gradually slowing down toward the setpoint current, at which point the velocity is virtually  $0\mu\text{m/s}$ . Since deceleration occurs much earlier than in the OL-ARS mode, the measured current may reflect the actual surface distance and therefore the current overshoot is no more than 3%, even at  $250\mu\text{m/s}$ . In addition, comparing the distance travelled by integrating the velocity profile, it shows that the nano-pipette operated in the OL-ARS mode travels  $150\text{ nm}$  farther than in the CL-ARS mode at  $150\mu\text{m/s}$ , a critical difference that could determine the non-contact imaging condition. For  $250\mu\text{m/s}$  approach speed, the difference is less drastic, but the distance travelled by the nano-pipette in the CL-ARS mode is found to be nearly equal to the sensing range of the pipette.

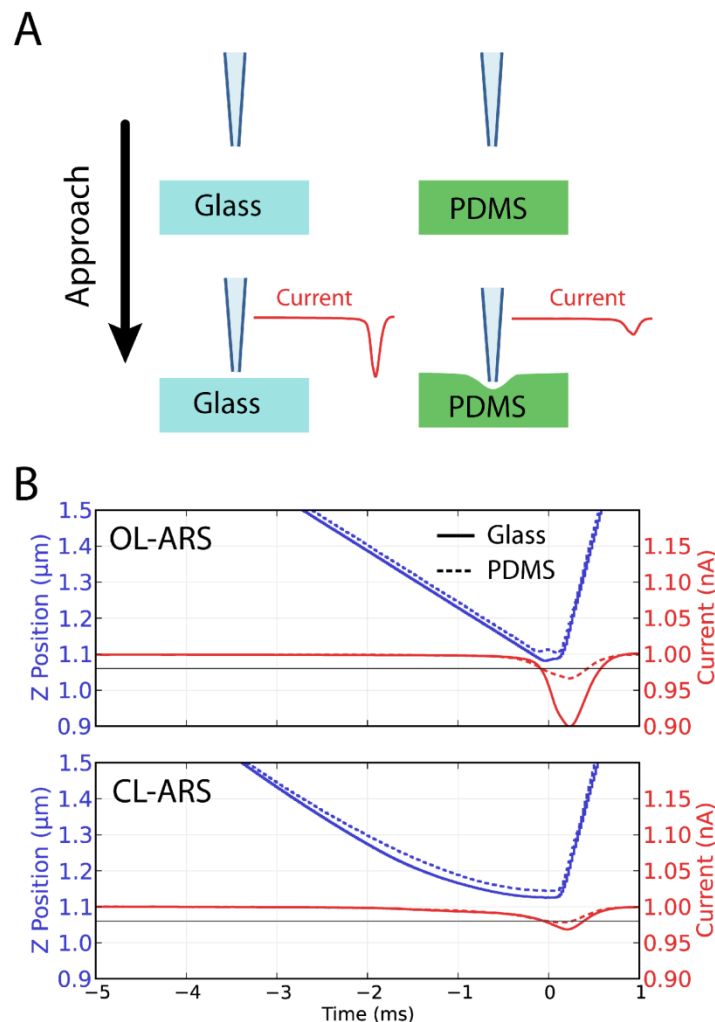


**Figure 4. Comparison of measurement stability in the OL-ARS and CL-ARS modes.** (A) OL-ARS height and current for approach speed of 150  $\mu\text{m/s}$  (upper) and 250  $\mu\text{m/s}$  (lower). (B) CL-ARS height and current for similar approach speeds with (A). (C) Average current and height (upper) and velocity profile (lower) superimposed for OL-ARS and CL-ARS with approach speed of 150  $\mu\text{m/s}$ . (D) Same comparison as (C) with higher approach speed of 250  $\mu\text{m/s}$ .

**Sample Hardness Comparison.** While the mechanical property of soft PDMS is closer to that of a living cell compared to a hard glass surface, we wanted to investigate the non-contact assumptions of the OL-ARS and CL-ARS modes in more detail. Therefore, we performed similar experiments for both soft PDMS and hard glass surfaces and compared the current stability. It has been previously reported<sup>29,30</sup> that a constant applied pressure on a nano-pipette interacting with a sample could be used for estimating mechanical properties. The SICM current-distance curve has also been reported<sup>31</sup> to depend on sample hardness. The evidence shows that a moving nano-pipette, to a certain degree, may be mechanically coupled with the sample surface exerting a physical influence. From the perspective of ARS performance, this means that the same approach speed would be expected to generate different current overshoot values depending on the sample surface, as shown in **Figure 5A** for the PDMS and glass substrates. The soft PDMS is deformed by the approaching tip, which generates a small current reduction, whereas the hard glass surface does not deform and thus there is no current reduction. Therefore, understanding the differences in current overshoot due to material properties can be used as an indication of the contact state of a particular ARS mode.

**Figure 5B** presents the current and height signals of the OL-ARS and CL-ARS modes near the glass and PDMS sample surfaces at an approach speed of 150  $\mu\text{m/s}$ . As can be observed, the OL-ARS mode shows a significantly greater current overshoot (greater than 10%) on the glass surface, as compared to the PDMS surface (3.5%). This result illustrates that the tip

interaction with the sample surface occurs for soft samples and therefore accurate control of the tip position is critical for a non-contact measurement. For the CL-ARS mode with a similar speed, the difference in the current overshoot between the glass and PDMS surfaces is less than 1%, indicating a very small interaction with the sample surface.



**Figure 5. Effect of surface mechanical properties on OL-ARS and CL-ARS mode operation.** (A) Illustration of current reduction with pipette approach for a glass surface (left) and soft PDMS surface (right). (B) Height and current signal comparison for glass and PDMS surfaces in OL-ARS (upper) and CL-ARS (lower) modes, both at 150  $\mu\text{m/s}$  approach speed.

**Cell Imaging.** We next performed SICM experiments in the OL-ARS and CL-ARS modes in order to characterize the morphology of electroporated H460 cells. When H460 cells are exposed

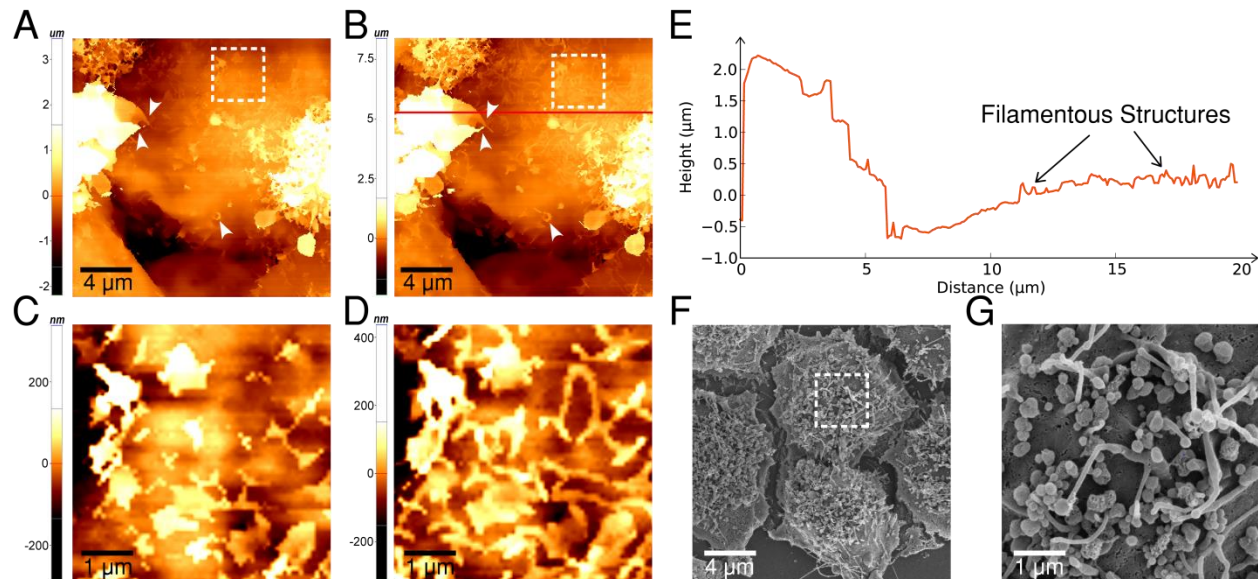


to electrical pulses, phospholipids and actin fibers associated with the plasma membrane are reorganized in order to make pores and filamentous structures<sup>25</sup>. **Figure 6** presents SICM and SEM images of these ultrastructural changes in membrane morphology. The cellular structures are independently visualized in the OL-ARS and CL-ARS modes, and the results are compared with the corresponding SEM images.

As shown in **Figure 6A,B**, measurements in both the OL-ARS and CL-ARS modes capture the general shape of the cell membrane that closely agree with each other and the SEM image (**Figure 6F**). However, upon examining the cell surface and edge features more closely, the difference in clarity of cellular structures is readily observed between the two modes (see white arrows in **Figure 6A, B**). The surface topography obtained with the OL-ARS mode (**Figure 6C**) shows the presence of some additional features, but it is difficult to draw any meaningful conclusion about the identity of these surface features or their characteristics due to the lack of detail. However, the same location imaged with the CL-ARS mode (**Figure 6D**) distinctly reveals that these features are actually connected components that closely resemble filamentous structures, as also visualized in the SEM experiments (**Figure 6G**). With the CL-ARS mode, filamentous structures of approximately 0.2  $\mu\text{m}$  height (**Figure 6E**) could be reliably measured. We stress that the differences in image quality stems from the relative stability of the two methods. While the OL-ARS mode often fails to stop approaches above the surface topography, which in turn generates disconnected or thinner features, the CL-ARS mode overcomes this issue and captures a more authentic topography with minimum artifacts, even for sample boundaries and edges.

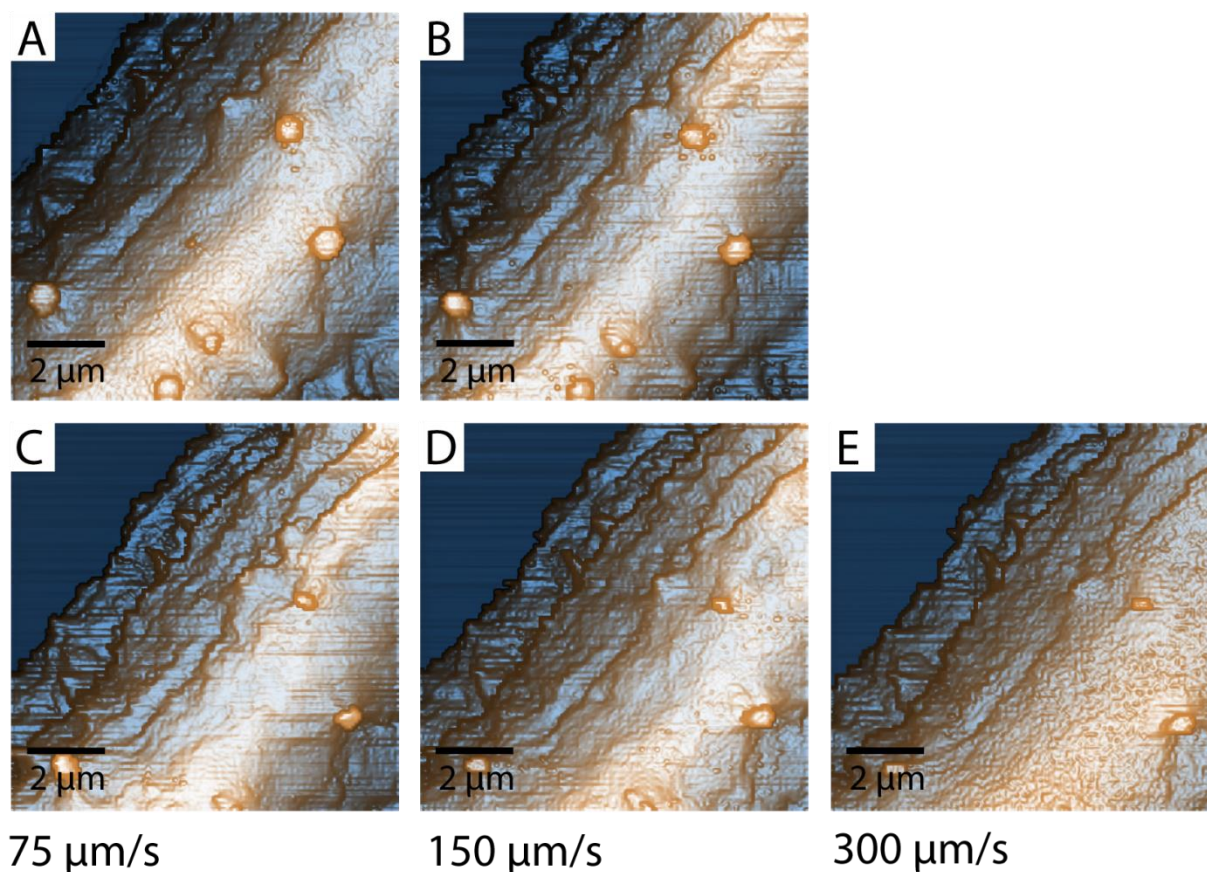
In order to evaluate the performance and stability of the OL-ARS and CL-ARS modes for imaging live cells, we also performed SICM experiments on live rat cardiomyocyte membrane

edges deposited on glass substrates (**Figure 7**). Images are displayed in an enhanced color scheme by computing the differential component of adjacent pixels to accentuate the topography changes of both glass and cell surfaces. While maintaining image quality, we investigated the maximum approach speed achievable in the OL-ARS and CL-ARS modes. Images were taken in the sequence of increasing speed in the OL-ARS mode, and then increasing speed in the CL-ARS mode. Neglecting minor deterioration of cell features due to the live imaging condition, we observed that the maximum approach speed which can be used with the OL-ARS mode is no more than  $150 \mu\text{m/s}$ . The pipette was broken during imaging at higher approach speeds. For the CL-ARS mode, features remained visible without pipette damage for approach speeds of up to  $300 \mu\text{m/s}$ , indicating at least a 50% increase in the maximum approach speed. At this imaging speed, however, image artifacts such as dense speckle noise began to appear on the surface of the cardiomyocytes as a result of the reduced stability (**Figure 7E**). This effect was not observed at approach speeds up to  $250 \mu\text{m/s}$  (**Figure 4B, lower**).



**Figure 6. SICM and SEM imaging on fixed H460 cells immediately post-electroporation. (A)** Image obtained with SICM OL-ARS mode, scale bar  $4 \mu\text{m}$ . **(B)** Image obtained with SICM CL-ARS mode, scale bar  $4 \mu\text{m}$ . **(C)** Zoomed-in image of Panel (A), scale bar  $1 \mu\text{m}$ . **(D)** Zoomed-in image of Panel (B), scale bar  $1 \mu\text{m}$ . **(E)** Line profile of the red line in the SICM image from

Panel (B). (F) Image obtained with SEM, scale bar  $4\ \mu\text{m}$ . (G) Zoomed-in image of Panel (F), scale bar  $1\ \mu\text{m}$ . The approach speed and total imaging time in Panels A-D were  $150\ \mu\text{m/s}$  and 52 minutes, respectively. The pixel resolution was  $256 \times 256$  and the approach/retract height was  $4\ \mu\text{m}$  in order to ensure safe and high-quality imaging.



**Figure 7. SICM imaging on live rat cardiomyocytes with different approach speeds.** OL-ARS mode with (A)  $75\ \mu\text{m/s}$  and (B)  $150\ \mu\text{m/s}$  approach speed. CL-ARS mode with (C)  $75\ \mu\text{m/s}$ , (D)  $150\ \mu\text{m/s}$ , and (E)  $300\ \mu\text{m/s}$  approach speed. The imaging time was 36 minutes for Panels A and C, 20 minutes for Panels B and D, and 15 minutes for Panel E. The pixel resolution was  $256 \times 256$  and the approach/retract height was  $2\ \mu\text{m}$ .

## **Conclusion**

In this work, we have proposed a closed-loop (CL) ARS method that is based on feedback control of the ion current for controlling the approach velocity of the pipette tip. We observed improved ion current stability (i.e., less overshoot) with the CL-ARS mode, as compared to the conventional, constant speed approach method (OL-ARS). Taking advantage of very small current changes, the new algorithm reduced the approach velocity prior to the setpoint current in order to minimize the chance for tip-sample interaction. This refinement of the tip approach resulted in reduced current overshoot and more consistent measurements of nanoscale sample topography. Since the physical interaction with soft biological samples is a key factor to minimize for optimal ARS performance, we performed SICM experiments with human lung cancer cells and rat cardiomyocytes. Importantly, these experiments demonstrated improved stability and speed for fixed and live cell imaging in the new CL-ARS mode, thereby demonstrating strong potential for wider application of the SICM technique in biological research and related fields which demand nanoscale characterization of soft materials.

## **Acknowledgements**

This work was supported in part by the Basic Science Research Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Science, ICT & Future Planning (grant number: 2013R1A1A2011526) and Advanced Technology Center (ATC) Program funded by the Ministry of Trade, Industry & Energy (grant number: 10045812) to S.-J.C. This work was also supported by the National Research Foundation (NRF-NRFF2011-01), and the National Medical Research Council (NMRC/CBRG/0005/2012) to N.J.C.

## References

- (1) Rheinlaender, J.; Geisse, N. A.; Proksch, R.; Schaffer, T. E. *Langmuir : the ACS journal of surfaces and colloids* **2011**, *27*, 697-704.
- (2) Jiao, Y.; Schäffer, T. E. *Langmuir : the ACS journal of surfaces and colloids* **2004**, *20*, 10038-10045.
- (3) Hennesthal, C.; Drexler, J.; Steinem, C. *Chemphyschem : a European journal of chemical physics and physical chemistry* **2002**, *3*, 885-889.
- (4) Hansma, P. K.; Drake, B.; Marti, O.; Gould, S. A.; Prater, C. B. *Science* **1989**, *243*, 641-643.
- (5) Korchev, Y. E.; Milovanovic, M.; Bashford, C. L.; Bennett, D. C.; Sviderskaya, E. V.; Vodyanoy, I.; Lab, M. J. *J Microsc* **1997**, *188*, 17-23.
- (6) Zhang, S.; Cho, S. J.; Busuttill, K.; Wang, C.; Besenbacher, F.; Dong, M. *Nanoscale* **2012**, *4*, 3105-3110.
- (7) Novak, P.; Shevchuk, A.; Ruenraroengsak, P.; Miragoli, M.; Thorley, A. J.; Klenerman, D.; Lab, M. J.; Tetley, T. D.; Gorelik, J.; Korchev, Y. E. *Nano letters* **2014**, *14*, 1202-1207.
- (8) Caldwell, M.; Del Linz, S. J.; Smart, T. G.; Moss, G. W. *Analytical chemistry* **2012**, *84*, 8980-8984.
- (9) Morris, C. A.; Friedman, A. K.; Baker, L. A. *The Analyst* **2010**, *135*, 2190-2202.
- (10) Edwards, M. A.; Williams, C. G.; Whitworth, A. L.; Unwin, P. R. *Analytical chemistry* **2009**, *81*, 4482-4492.
- (11) Del Linz, S.; Willman, E.; Caldwell, M.; Klenerman, D.; Fernandez, A.; Moss, G. *Analytical chemistry* **2014**, *86*, 2353-2360.
- (12) Shevchuk, A. I.; Frolenkov, G. I.; Sánchez, D.; James, P. S.; Freedman, N.; Lab, M. J.; Jones, R.; Klenerman, D.; Korchev, Y. E. *Angewandte Chemie* **2006**, *118*, 2270-2274.
- (13) Anariba, F.; Anh, J. H.; Jung, G.-E.; Cho, N.-J.; Cho, S.-J. *Modern Physics Letters B* **2012**, *26*, 1130003.
- (14) Lee, Y.; Jung, G. E.; Cho, S. J.; Geckeler, K. E.; Fuchs, H. *Nanoscale* **2013**, *5*, 8577-8585.
- (15) Clarke, R. W.; Zhukov, A.; Richards, O.; Johnson, N.; Ostanin, V.; Klenerman, D. *Journal of the American Chemical Society* **2013**, *135*, 322-329.
- (16) Novak, P.; Li, C.; Shevchuk, A. I.; Stepanyan, R.; Caldwell, M.; Hughes, S.; Smart, T. G.; Gorelik, J.; Ostanin, V. P.; Lab, M. J.; Moss, G. W. J.; Frolenkov, G. I.; Klenerman, D.; Korchev, Y. E. *Nature methods* **2009**, *6*, 279-281.
- (17) Happel, P.; Hoffmann, G.; Mann, S. A.; Dietzel, I. D. *J Microsc* **2003**, *212*, 144-151.
- (18) Happel, P.; Dietzel, I. D. *Journal of nanobiotechnology* **2009**, *7*, 7.
- (19) Takahashi, Y.; Murakami, Y.; Nagamine, K.; Shiku, H.; Aoyagi, S.; Yasukawa, T.; Kanzaki, M.; Matsue, T. *Physical chemistry chemical physics : PCCP* **2010**, *12*, 10012-10017.
- (20) Yamada, H. *Electrochimica Acta* **2014**, *136*, 233-239.
- (21) Ushiki, T.; Nakajima, M.; Choi, M.; Cho, S. J.; Iwata, F. *Micron* **2012**, *43*, 1390-1398.
- (22) Proksch, R.; Lal, R.; Hansma, P. K.; Morse, D.; Stucky, G. *Biophysical journal* **1996**, *71*, 2155-2157.
- (23) Pastre, D.; Iwamoto, H.; Liu, J.; Szabo, G.; Shao, Z. *Ultramicroscopy* **2001**, *90*, 13-19.
- (24) McKelvey, K.; Perry, D.; Byers, J. C.; Colburn, A. W.; Unwin, P. R. *Analytical chemistry* **2014**, *86*, 3639-3646.
- (25) Gerisch, G.; Ecke, M.; Neujahr, R.; Prassler, J.; Stengl, A.; Hoffmann, M.; Schwarz, U. S.; Neumann, E. *Journal of cell science* **2013**, *126*, 2069-2078.

- (26) Williamson, C.; Gorelik, J.; Eaton, B. M.; de Swiet, M.; Korchev, Y. *Clinical science (London, England: 1979)* **2001**, *100*, 363-369.
- (27) Li, C.; Johnson, N.; Ostanin, V.; Shevchuk, A.; Ying, L.; Korchev, Y.; Klenerman, D. *Progress in Natural Science* **2008**, *18*, 671-677.
- (28) Liu, B. C.; Lu, X. Y.; Song, X.; Lei, K. Y.; Alli, A. A.; Bao, H. F.; Eaton, D. C.; Ma, H. P. *Frontiers in physiology* **2012**, *3*, 483.
- (29) Schaffer, T. E. *Analytical chemistry* **2013**, *85*, 6988-6994.
- (30) Rheinlaender, J.; Schäffer, T. E. *Soft Matter* **2013**, *9*, 3230.
- (31) Mizutani, Y.; Choi, M.-H.; Cho, S.-J.; Okajima, T. *Applied physics letters* **2013**, *102*, 173703.