

# Nanoscale

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

*Accepted Manuscripts* are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this *Accepted Manuscript* with the edited and formatted *Advance Article* as soon as it is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.

Cite this: DOI: 10.1039/c0xx00000x

www.rsc.org/xxxxxx

ARTICLE TYPE

# Tunable Ferromagnetic Ordering in MoS<sub>2</sub> nanosheets with Fluorine Adsorption

Daqiang Gao, Shoupeng Shi, Kun Tao, Baorui Xia and Desheng Xue\*

Received (in XXX, XXX) Xth XXXXXXXXX 20XX, Accepted Xth XXXXXXXXX 20XX

DOI: 10.1039/b000000x

Two-dimensional ferromagnetic ultrathin nanosheets hold great promise for next generation electronics and spintronics. Here, intrinsic ferromagnetism were achieved through a new effective strategy by fluorine adsorption on MoS<sub>2</sub> nanosheets, where the fluorinated MoS<sub>2</sub> nanosheets exhibit stable ferromagnetic hysteresis at room temperature with saturation magnetization of 0.06 emu/g and the magnetoresistance of 4.1 %. The observed ferromagnetism can be tuned by changing the concentration of the adatom. On the basis of first-principle calculations, it is shown that not only fluorine absorbed MoS<sub>2</sub> monolayer favours spontaneous spin polarization and local moment formation, but also that the spin moments can exhibit long range magnetic ordering. This work paves a new pathway to engineer the magnetic properties of the two-dimensional nano-materials.

Two-dimensional (2D) crystals of transition metal dichalcogenides (TMDs) have received significant interest due to their potential in a wide range of novel applications as well as in basic research.<sup>1-3</sup> As the member of the TMDs family, molybdenum disulfide (MoS<sub>2</sub>) isolated in the single-layer form was first to attract attention, demonstrating high on/off current ratio and low off-state current.<sup>4</sup> Recently, integrated circuits based on single-layer MoS<sub>2</sub>, bilayer MoS<sub>2</sub>, signal amplifiers, and nonvolatile memory cells were also demonstrated.<sup>5-7</sup> What's more, the presence of a direct band gap in MoS<sub>2</sub> due to the dimensional confinement makes it interesting for applications in optoelectronics.<sup>8-10</sup>

In addition, studies on the magnetic response of MoS<sub>2</sub> were also proposed recent, proving its potential application in spintronics and nanoelectronic devices. Many theoretical studies suggested that, although bulk MoS<sub>2</sub> is a diamagnetic material like graphite, the MoS<sub>2</sub> becomes ferromagnetic when MoS<sub>2</sub> nanoribbons are formed with zigzag edges, or defects are induced.<sup>11,12</sup> Experimentally, ferromagnetism was obtained in freestanding nanosheets,<sup>13</sup> the vertically aligned nanosheets,<sup>14</sup> and in the bulk MoS<sub>2</sub> irradiated by proton,<sup>15</sup> where the existence of ferromagnetism was partly attribute to the presence of zigzag edges or defects states at the grain boundaries. However, for its use in nanoelectronics, the magnetic coupling modulated by the edges and defects states

in MoS<sub>2</sub> nanosheets is always somewhat rough as compared to the ideal requirement (stable and precision magnetic spin ordering) and, thus, is inconvenient to control in the actual application, which hampers the development of this material in nanoelectronic devices.<sup>16</sup>

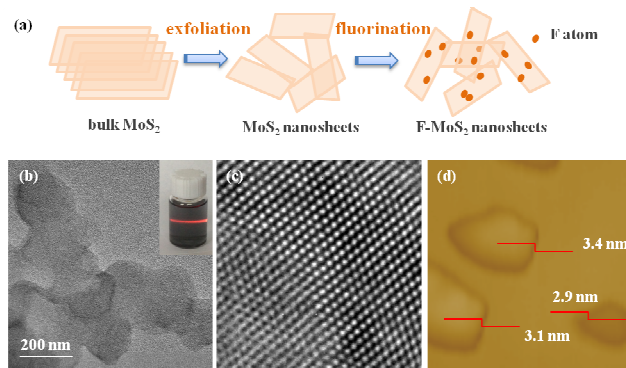
Previous studies have found that the absorption of nonmetal element on the surface of low-dimensional systems can induce local magnetic moment.<sup>17,18</sup> Compared with the magnetic moment from *d* or *f* electrons of transition-metal or rare-earth atoms, the magnetism based on the *sp* states of nonmetal element has some obvious advantage, such as stronger long-range exchange coupling interactions and no clustering of magnetic ions. Fluorination was confirmed to be the most effective route to introduce localized spins in graphene and it has been reported that the fluorination can induce tunable magnetism in the BN nanotubes and the nanosheets.<sup>19,20</sup> Recently, the published studies on the possible magnetic properties of fluorinated TMDs are also performed in theoretical nature. For example, He *et al.* investigated the adsorption of nonmetal atoms on pristine MoS<sub>2</sub> and found that H, B, C, N and F atoms can induce local magnetic moments.<sup>21</sup> Yue *et al.* study the magnetic and electronic properties of MoS<sub>2</sub> sheets with substitution of H, B, C, O, N and F atoms with sulfur defects (*V<sub>S</sub>*). Calculated binding energies show that all the substitutional atoms are strongly bound to *V<sub>S</sub>* defects, resulting in different magnetic behaviors.<sup>22</sup> These anticipate results opening a new way to modify the magnetic properties of the monolayer MoS<sub>2</sub>. However, examining the magnetic properties of non-metal atoms absorbed MoS<sub>2</sub> has not been realized experimentally. Therefore, it is of crucial importance to develop a feasible method for synthesizing them and explore the corresponding magnetic properties.

In this paper, we reported the preparation and corresponding magnetic properties of the fluorinated MoS<sub>2</sub> (F-MoS<sub>2</sub>) nanosheets. First-principles calculations results demonstrate that F atom can chemically absorb on MoS<sub>2</sub> monolayer favouring the top of S atom, where total magnetic moment of 1.0 μ<sub>B</sub> was found in F-absorbed MoS<sub>2</sub> monolayer. Magnetic measurements results indicate that the enhancement of the room temperature ferromagnetism was observed for F-MoS<sub>2</sub> nanosheets compared to the primitive nanosheets, which

material is quite promising for applications in spintronic devices.

## Results and discussions

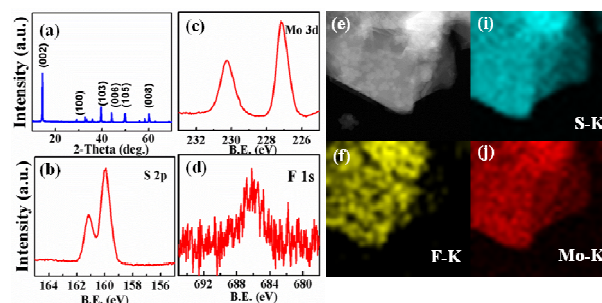
F-MoS<sub>2</sub> nanosheets were prepared by a two-step method as shown in Fig. 1a. In brief, ultrathin MoS<sub>2</sub> nanosheets were first gotten by exfoliating the bulk form in N,N-Dimethylformamide (DMF, 100 mL) solution as our previous report.<sup>13</sup> Then, 1.0 g of the exfoliated MoS<sub>2</sub> nanosheets and 2.0 g ammonium fluoride (NH<sub>4</sub>F) were further ground into a homogeneous form with the deionized water and transferred in a Teflon-lined stainless steel autoclave, which system was heated at 180 °C for 24 h and then cooled down to room temperature. The autoclave container was taken out and the product was centrifuged and washed with ethanol and deionized water repeatedly. Finally, the as-obtained sample was dried in a vacuum oven for 24 h. Subsequently, systematical characterizations were performed to verify that the finally products were in ultrathin nanosheet structure. Transmission electron microscope (TEM, TecnaiTM G2 F30, FEI, USA) image of F-MoS<sub>2</sub> nanosheets shown in Fig. 1b revealed that the lateral diameters of the nanosheets are ~200 nm. Results indicate that the colloidal dispersions of the F-MoS<sub>2</sub> nanosheets are highly stable over a period of several days and display a typical Tyndall effect, reasonably indicating the formation of freestanding and homogeneous 2D ultrathin sheets (inset in Fig. 1b). The further high resolution TEM results for the selected regions of one nanosheet is shown in Fig. 1c, indicating that the nanosheets is well crystallographic. The thickness of the products was measured by the Atomic force microscopy (AFM Dimension 3100 with Nanoscope IIIa controller, Veeco, CA, USA). As can be seen from Fig. 1d, the thicknesses of ultrathin nanosheets are ranging from 2.9 nm to 3.4 nm, which indicates the ultrathin nanosheets are composed of 3~5 atomic monolayers.



**Fig. 1.** (a) Schematic illustration for the synthesizing of fluorination-MoS<sub>2</sub> nanosheets. (b) TEM and the corresponding colloidal DMF dispersion displaying the Tyndall effect. (c) HRTEM (d) AFM image and the corresponding height profile for the fluorination-MoS<sub>2</sub> nanosheets.

Typical X-ray diffraction (XRD, X' Pert PRO PHILIPS with Cu K $\alpha$  radiation) spectrum of the F-MoS<sub>2</sub> nanosheets is shown in Fig. 2a, it is possible to assign the reflection peaks to the family lattices planes of hexagonal MoS<sub>2</sub> (JCPDS Card No.77-1716). Result indicates there is no other new phase introducing for the fluorination progress. What's more, the

bonding characteristics and the composition of the exfoliated MoS<sub>2</sub> samples were captured by X-ray photoelectron spectroscopy (XPS, Kratos Axis Ultra). Besides element of C, the wide XPS spectra of the F-MoS<sub>2</sub> sample show only signals arising from elements Mo, S and F. The peaks, corresponding to the S 2p<sub>1/2</sub> and S 2p<sub>3/2</sub> orbital of divalent sulfide ions (S<sup>2-</sup>), are observed at 163.3 and 162.0 eV shown in Fig. 2b. The Mo 3d XPS spectrum of F-MoS<sub>2</sub> nanosheets, reported in Fig. 2c shows two strong peaks at 229.3 and 232.5 eV, respectively, attributed to the doublet Mo 3d<sub>5/2</sub> and Mo 3d<sub>3/2</sub>. Besides, the peak at the binding energy of 686.9 eV attributes to the F 1s spectre. All these results are consistent with the reported values.<sup>23</sup> Elemental mapping results shown in Fig. 2 f~i also give the information that the samples contain element Mo, S and F, where the elemental distribution are uniformly.

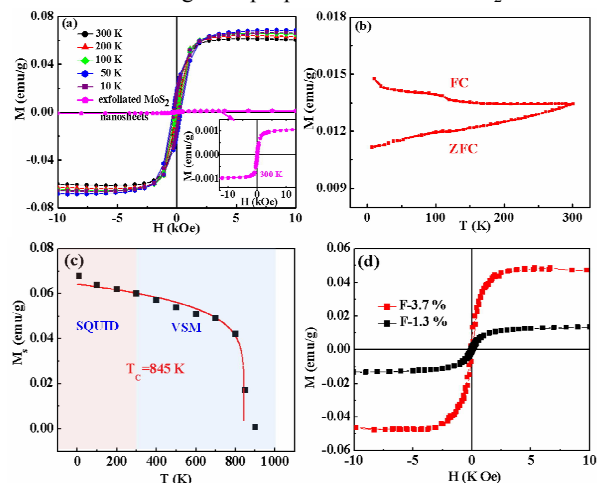


**Fig. 2.** (a) XRD (b)~(d) XPS results, and (e)~(j) TEM image and the elemental mapping for the fluorination-MoS<sub>2</sub> nanosheets.

The magnetic properties of the F-MoS<sub>2</sub> nanosheets were studied by using a superconducting quantum interferometer device (SQUID) magnetometer from Quantum Design. Fig. 3a shows the magnetization curves ( $M$ - $H$ ) of the F-MoS<sub>2</sub> nanosheets measured at different temperatures (10~300 K) and the magnetization curve for the exfoliated MoS<sub>2</sub> nanosheets was also given to compared, where the diamagnetic background have been subtracted (the primitive curves are shown in Fig.S1†). It can be seen that the exfoliated MoS<sub>2</sub> nanosheets show the saturation magnetization ( $M_s$ ) of 0.001 emu/g at 300 K, which is corresponding to our previous result, where the ferromagnetism is considered to be related to the edge states.<sup>13</sup> However, it is interesting that the sample's  $M_s$  shows the great enhancement after fluorination. As shown in Fig. 3a, the F-MoS<sub>2</sub> nanosheets have the  $M_s$  of 0.07 and 0.06 emu/g at 10 and 300 K, respectively. The magnetic hysteresis loops measured in the low field range are plotted in Fig. S2†. Obviously, ferromagnetic behaviour of the F-MoS<sub>2</sub> is robust at this temperature range with the remanent magnetic moment ( $M_r$ ) and coercivity ( $H_c$ ) decreases with the increase of temperature. In addition, temperature dependent magnetization ( $M$ - $T$ ) curves of the F-MoS<sub>2</sub> nanosheets presented in Fig. 3b provide another evidence to prove the sample is intrinsically room temperature ferromagnetism. Obviously, zero field cooling (ZFC) and field cooling (FC) curves showed distinct difference in wide temperature range from 10 up to 300 K, revealing Curie temperature is higher than 300 K. Most importantly, there is no block temperature appearance in ZFC curve, which clearly reveals that there are no ferromagnetic clusters in our sample and provides evidence

for ferromagnetism in the samples.<sup>24</sup> Besides, it can be seen that there is a magnetic moments change at about 120 K in the FC curve. Such magnetic anomaly was found when antiferromagnetic (AFM) correlations competed with ferromagnetic (FM) interactions. This result was similar to that obtained for graphene,<sup>25,26</sup> somewhat like spin glass behavior in frustrated system. Though the origin remained unclear, a coexistence of ferromagnetism and antiferromagnetism was well accepted for a biparticle system.<sup>27</sup> Magnetic measurement results indicate that the raw materials (*h*-MoS<sub>2</sub> powder and NH<sub>4</sub>F) show diamagnetism only (see Fig.S3†). Thus, we conclude that the enhanced ferromagnetism observed in the F-MoS<sub>2</sub> nanosheets is mainly induced by fluorination. Of note, the ferromagnetic impurities such as Fe, Co, Ni could also be excluded by inductively coupled plasma (ICP) results, which contents are all below 10 ppm. Therefore, all above analyses confirmed that the room temperature ferromagnetism in F-MoS<sub>2</sub> nanosheets is inherent nature.<sup>28,29</sup>

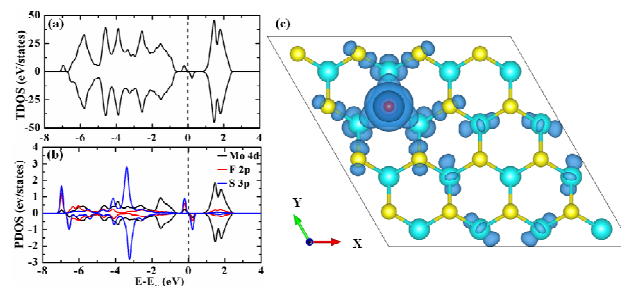
By using the SQUID and vibrating sample magnetometer equipped with a high temperature chamber (VSM, VSM Model EV9, MicroSense, LLC), we could get the information of  $M_s$  dependence on temperature for the F-MoS<sub>2</sub> nanosheets. As shown in Fig. 3c, the dates can be well fitted by the following equation which is common for a ferromagnetic material:<sup>30,31</sup>  $M(T)=A(1-T/T_C)^\beta$ , where  $A$  is a coefficient related to the spontaneous magnetization,  $T_C$  is the ferromagnetic Curie temperature, and  $\beta$  is the critical exponent. The best fitting curve ( $T < T_C$ ) is shown by red curve, where the fitted sample's Curie temperature is about 845 K. Such high  $T_C$  ferromagnetism in few-layer F-MoS<sub>2</sub> nanosheets makes them promising candidates for application in spintronic devices.<sup>32,33</sup> Besides, we also studied the affect of amount of fluoride on the magnetic properties of the F-MoS<sub>2</sub> nanosheets.



**Fig. 3.** (a) Magnetization curves ( $M$ - $H$ ) of the fluorination-MoS<sub>2</sub> nanosheets measured at different temperatures (10–300 K), where the diamagnetic background have been subtracted. (b) Temperature dependence of zero field cooling (ZFC) and field cooling (FC) curves of fluorination-MoS<sub>2</sub> nanosheets under the measuring field of 200 Oe. (c)  $M_s$  dependence on temperature for the fluorination-MoS<sub>2</sub> nanosheets (10–900 K). (d)  $M$ - $H$  curves of fluorination-MoS<sub>2</sub> nanosheets fabricated with different mounts of NH<sub>4</sub>F at 300 K.

We changed the quality of NH<sub>4</sub>F from 2.0 g to 1.5 g and 0.5 g, respectively, in the experiment progress with the other conditions identical, where the concentration of fluorine atoms were 1.3 %, 3.7 % and 7.0 % with the F<sup>-</sup> electron state for the three samples (NH<sub>4</sub>F=0.5 g, 1.0 g and 2.0 g), which were gotten by XPS measurement. As can be seen from Fig. 3d that the F-MoS<sub>2</sub> nanosheets show different  $M_s$ , indicating that ferromagnetism of the samples can be tuned by changing the concentration of the adsorption atoms.

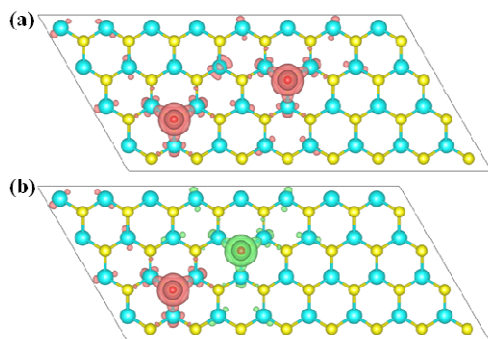
Above results indicate that the observed larger ferromagnetism in F-MoS<sub>2</sub> nanosheets is mainly induced by introducing the F atom absorbed on the MoS<sub>2</sub> nanosheets. To further understand how fluoride influence the ferromagnetism in MoS<sub>2</sub> nanosheets, the geometries, electronic structures, and magnetic behaviour of F-absorbed MoS<sub>2</sub> monolayers were investigated by using first principles calculations. The studied model consists of a 4×4×1 MoS<sub>2</sub> supercell (including 32 S and 16 Mo atoms) and an adatom (F atom). The distance between two adjacent MoS<sub>2</sub> monolayers is larger than 16 Å to avoid the layer-layer interaction. Calculations results indicate that the most energetically favorable absorbed position is at the top of S atom for the adatom. The spin-polarized calculations within density functional theory formalism were performed using the Vienna *ab initio* simulation package (VASP) with the generalized gradient approximation of Perdew–Burke–Ernzerhof functional.<sup>34</sup> The plane wave basis set with kinetic energy cutoff of 500 eV was employed. The sampling of the Brillouin zone was done using a 5×5×1 Monkhorst–Pack grid. All the structures were fully relaxed until the Hellmann–Feynman force was less than 0.01 eV/Å.<sup>35</sup>



**Fig. 4.** (a) The TDOS and (b) PDOS of the F-absorbed MoS<sub>2</sub> monolayers. The black, red, and blue lines represent the nearest Mo atom, the adatom, and the nearest S atom, respectively. (b) The spin resolved charge density isosurface (0.001 Å<sup>-3</sup>) of the F-absorbed MoS<sub>2</sub> monolayers. Mo atoms are indicated by blue balls, S atoms by yellow balls and F atoms by red balls, respectively.

As shown in Fig. 4a, a significant asymmetry between the spin-up state and spin-down state in the density of states (DOS) near the Fermi level suggests the intrinsic ferromagnetism of F-absorbed MoS<sub>2</sub> monolayer. As well known, the value of the Curie temperature depends on the exchange stiffness and the DOS at the Fermi level. The calculated DOS display obvious spin splitting at the conduction band and valence band means the relative high Curie temperature in this configuration, which is identical with our experimental results. The total magnetic moments of this electronic structure is about 1.0  $\mu_B$ . The spin-resolved

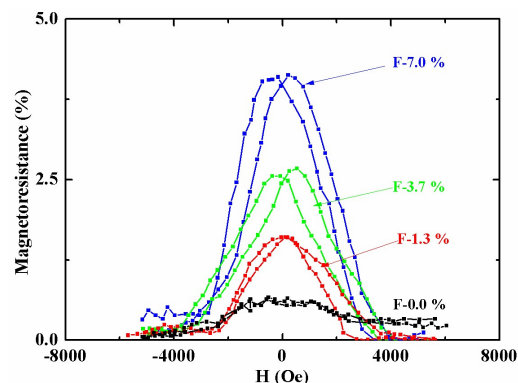
DOS projected on the  $4d$  orbital of the nearest Mo atom,  $p$  orbitals of the adatom and the nearest S atom are also presented in the Fig. 4b. All of them show significant asymmetry between the spin-up state and spin-down state near the Fermi level, illustrating both of them contribute the magnetic moment to the total magnetic moments. Obviously, the PDOS results clearly uncovers that the  $2p$  orbital of the adatom are localized, spin-polarized, and consequently induces polarizations of S  $3p$  and Mo  $4d$  electrons nearby the adatoms. Near the Fermi energy, the spin-up states are full occupied and the spin-down states are unoccupied or partially occupied, which is the origin of magnetic moment for the F-adsorbed MoS<sub>2</sub> monolayer. These results are consistent with the spin resolved charge density distribution of the F-adsorbed MoS<sub>2</sub> monolayers as shown in Fig. 4c. It is worth noting that the overlaps between F  $2p$  states and  $3s$  states of the second-nearest S,  $4d$  states of the third-nearest Mo atoms are still evident, which corresponds to the large spatial extension of spin density of F-MoS<sub>2</sub> monolayer.



**Fig. 5.** Spatial distribution of the spin-polarization density of the MoS<sub>2</sub> monolayers with two F atoms absorbed. (a) ferromagnetism coupling, (b) antiferromagnetism coupling. The unit is  $\text{\AA}^{-3}$ , Mo atoms are indicated by blue balls, S atoms by yellow balls and F atoms by red balls, respectively.

To further study the magnetic coupling between the fluorination induced local magnetic moments, we adopted large-size supercells containing 32 primitive cells (denoted as  $8 \times 4 \times 1$ ) and two F ions, which correspond to the fluorination rate of 1/16 in a MoS<sub>2</sub> monolayer. Starting from different initial spin arrangement, self-consistent calculations lead to two types of magnetic orderings. One has the local magnetic moments aligned in a "FM" way. Another has the local magnetic moments aligned in an "AFM" way.<sup>36,37</sup> Calculation results indicate that two stable magnetic structures were obtained: one was FM ( $E_{\text{AFM}} - E_{\text{FM}} = 0.32$  meV), the other was AFM ( $E_{\text{AFM}} - E_{\text{FM}} = -0.48$  meV). It can be seen from Fig. 5a that most of the spin polarized states localized around the F dopants and led to the formation of local magnetic moment while the long range magnetic coupling was mediated between the two moments. However, as the distance between two F dopants decreased, the system would be more stable in the AFM phase, due to the broaden of the dopant band equal to spin split, caused by the larger degree of overlap between the unpaired electrons (Fig. 5b). These results are similar to the reports of C doped BN and defective C<sub>60</sub> systems.<sup>36,38</sup> To

sum up, above calculations results indicate that the spin moments induced by topological fluorine adsorption on MoS<sub>2</sub> nanosheets can be (anti)ferromagnetically coupled to exhibit long rang magnetic ordering, which is corresponding to the experimental results.



**Fig. 6.** MR curves of the fluorination-MoS<sub>2</sub> nanosheets measured at room temperature.

In order to check the intrinsic nature of the ferromagnetism in F-MoS<sub>2</sub> nanosheets and its possible spintronics applications, we measured magnetoresistance (MR) of F-MoS<sub>2</sub> nanosheets by two-terminal method, where the F-MoS<sub>2</sub> nanosheets were pressed into a 5 mm-diameter piece. The MR results for the samples with different F contents are shown in Fig. 6 in the field range of  $-8000 \text{ Oe} < H < 8000 \text{ Oe}$ , where the MR value is defined as  $[R(H) - R(0)]/R(0)$  and  $R(0)$  is the resistance at zero field. It can be seen that all the samples show hysteresis negative MR curves, where the sample with the F content of 7.0% shows the max MR value of 4.1%. Similar negative hysteresis MR has been observed in reduced graphene oxide and VS<sub>2</sub> nanosheets, which is a clear indication of the intrinsic nature of the observed ferromagnetism.<sup>39,40</sup>

In addition, we have quantitatively analyzed the magnetization signal to get the magnetic moment per absorb F atom. Results indicate that the F concentration for sample F-MoS<sub>2</sub>-2.0 is about 7.0%. Therefore, with a F/Mo+S+F ratio of 7%, the number of F dopants per unit mass, was given by  $N_C = (N_A/M_R) \times 7\% = 2.63 \times 10^{20}$  dopants/g,<sup>36</sup> where  $M_R$  was the relative molecular weight of MoS<sub>2</sub> (160.06) and  $N_A$  was the Avogadro's constant of  $6.02 \times 10^{23}$ . The value of the  $M_s$  for the F-MoS<sub>2</sub> nanosheets was 0.07 emu/g as shown in Figure 3a, and thus, the magnetic moment per absorbed F was given by  $\mu_C = M_s / N_C = 1.16 \times 10^{-22} \text{ emu} = 1.25 \times 10^{-2} \mu_B$ . Comparing with our theoretical value of 1.0  $\mu_B$  per F dopant, only 1/8 F dopants seemed to contribute to the ferromagnetism in experiment. Noting that we used a MoS<sub>2</sub> monolayer model in these calculations, for the synthesized F-MoS<sub>2</sub> nanosheets which contain multi-layers, the fluorination rate will be lower than this model. In addition, the reason may be that some absorbed F atoms formed islands on the surfaces of MoS<sub>2</sub> nanosheets. Apart from these, some AFM coupling also make the total magnetic moment from these two F dopants to be zero. Therefore, only those F dopants that had no counterparts on the vicinal site sublattice would contribute to total magnetic moment to F-MoS<sub>2</sub>

nanosheet in our case. However, to what extent each factor will play a role is still a question and our work is on the way.

## Conclusions

In summary, we have demonstrated that involving of F atoms absorbed on MoS<sub>2</sub> nanosheets can induce intrinsic ferromagnetism, which is verified by both magnetic characterizations and theoretical calculations. The saturation magnetization of F-MoS<sub>2</sub> nanosheets at room temperature was as high as 0.06 emu/g, which is larger than that of the exfoliated MoS<sub>2</sub> nanosheets caused by defects and edge state. What's more, hysteresis negative MR curve can be observed for all the samples, revealing its possible spintronics applications. This finding broaden our horizon and it is anticipate that introduction of non-magnetic absorbed strategy could be an effective way for engineering the intrinsic magnetic properties of 2D nanomaterials.

## Acknowledgements

This work is supported by National Basic Research Program of China (Grant No. 2012CB933101), the National Natural Science Foundation of China (Grant No. 51371093, 51202101, 11474137, 11274146 and 11034004), Program for Changjiang Scholars and Innovative Research Team in University (Grant No. IRT1251) and the Specialized Research Fund for the Doctoral Program of Higher Education (20120211120005).

## Notes and references

*Key Laboratory for Magnetism and Magnetic Materials of MOE, Key Laboratory of Special Function Materials and Structure Design, Ministry of Education, Lanzhou University, Lanzhou 730000, P. R. China. Fax: +86 0931 8914160; Tel: +86 0931 8914160; \*E-mail: [xlweds@lzu.edu.cn](mailto:xlweds@lzu.edu.cn); † Electronic Supplementary Information (ESI) available]. See DOI: 10.1039/b000000x/*

- K. F. Mak, K. He, J. Shan and T. F. Heinz, *Nat. Nanotechnol.* **7**, 494 (2012).
- Y. Li, H. Wang, L. Xie, Y. Lang, G. Hong and H. Dai, *J. Am. Chem. Soc.* **133**, 7296 (2011).
- Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.* **7**, 699 (2012).
- B. Radisavljevic, M. B. Whitwick, A. Kis, *ACS Nano*, **5**, 9934 (2011).
- H. Wang, L. Yu, Y.-H. Lee, Y. Shi, A. Hsu, M. L. Chin, L.-J. Li, M. Dubey, J. Kong and T. Palacios, *Nano Lett.* **12**, 4674 (2012).
- B. Radisavljevic, M. B. Whitwick, A. Kis, *Appl. Phys. Lett.* **101**, 043103 (2012).
- S. Bertolazzi, D. Krasnozhan, A. Kis, *ACS Nano*, **7**, 3246 (2013).
- D. Ovchinnikov, A. Allain, Y. S. Huang, D. Dumcenco, and A. Kis, *ACS Nano*, **8**, 8174 (2014).
- Z. Jin, S. Shin, D. H. Kwon, S. J. Han and Y. S. Min, *Nanoscale*, **6**, 14453 (2014).
- J. Xia, X. Huang, L. Z. Liu, M. Wang, L. Wang, B. Huang, D. D. Zhu, J. J. Li, C. Z. Gu and X. M. Meng, *Nanoscale*, **6**, 8949 (2014).
- C. Ataca, H. S. ahin, E. Aktu'rk, and S. Ciraci, *J. Phys. Chem. C* **115**, 3934 (2011).
- Y. Li, Z. Zhou, S. Zhang, and Z. J. Chen, *J. Am. Chem. Soc.* **130**, 16739 (2008).
- D. Gao, M. Si, J. Li, J. Zhang, Z. Zhang, Z. Yang, and D. Xue, *Nanoscale Res. Lett.* **8**, 129 (2013).
- J. Zhang, J. M. Soon, K. P. Loh, J. Yin, J. Ding, M. B. Sullivan, and P. Wu, *Nano Lett.* **7**, 2370 (2007).
- S. W. Han, Y. H. Hwang, S. H. Kim, W. S. Yun, J. D. Lee, M. G. Park, S. Ryu, J. S. Park, D. H. Yoo, S. P. Yoon, S. C. Hong, K. S. Kim, and Y. S. Park, *Phys. Rev. Lett.* **110**, 247201 (2013).
- Y. D. Ma, Y. Dai, M. Guo, C. W. Niu, Y. T. Zhu, and B. B. Huang, *ACS Nano*, **6**, 1695 (2012).
- O. V. Yazyev and L. Helm, *Phys. Rev. B* **75**, 125408 (2007).
- J. Zeng, K. Chen, C. Sun, *Phys. Chem. Chem. Phys.* **14**, 8032 (2012).
- Q. Feng, N. Tang, F. Liu, Q. Cao, W. Zheng, W. Ren, X. Wan and Y. Du, *ACS Nano*, **7**, 6729 (2013).
- Z. H. Zhang and W. L. Guo, *J. Am. Chem. Soc.* **131**, 6874 (2009).
- J. He, K. Wu, R. Sa, Q. Li, Y. Wei, *Appl. Phys. Lett.* **96**, 082504 (2010).
- Q. Yue, S. L. Chang, S. Q. Qin, J. B. Li, *Phys. Lett. A* **377**, 1362 (2013).
- M. Du, X. L. Li, A. Z. Wang, Y. Z. Wu, X. P. Hao, and M. W. Zhao, *Angew. Chem.* **126**, 3719 (2014).
- D. Q. Gao, G. J. Yang, J. Zhang, Z. H. Zhu, M. S. Si, and D. S. Xue, *Appl. Phys. Lett.* **99**, 052502 (2011).
- H. S. S. R. Matte, K. S. Subrahmanyam, C. N. R. Rao, *J. Phys. Chem. C*, **113**, 9982 (2009).
- L. L. Chen, L. W. Guo, Z. L. Li, H. Zhang, J. J. Lin, J. Huang, S. F. Jin, X. L. Chen, *Sci. Rep.* **3**, 0259911 (2013).
- O. V. Yazyev, *Rep. Prog. Phys.* **73**, 056501 (2010).
- X. Chen, G. Li, Y. Su, X. Qiu, L. Li, and Z. Zou, *Nanotechnology*, **20**(11), 115606 (2009).
- M. Khalid, M. Ziese, A. Setzer, P. Esquinazi, M. Lorenz, H. Hochmuth, M. Grundmann, D. Spemann, T. Butz, G. Brauer, W. Anwand, G. Fischer, W. A. Adeagbo, W. Hergert, and A. Ernst, *Phys. Rev. B*, **80**, 035331 (2009).
- H. E. Stanley, *Introduction to Phase Transitions and Critical Phenomena*, Oxford University Press, Oxford, 1987.
- Z. L. Yang, D. Q. Gao, J. Zhang, Q. Xu, S. P. Shi, K. Tao, and D. S. Xue, *Nanoscale*, **7**, 650 (2015).
- O. V. Yazyev, M. I. Katsnelson, *Phys. Rev. Lett.* **100**, 047209 (2008).
- Y. G. Zhou, Z. G. Wang, P. Yang, X. T. Zu, L. Yang, X. Sun, and F. Gao, *ACS Nano*, **6**(11), 9727 (2012).
- G. Kresse and J. Hafner, *J. Phys.: Condens. Matter* **6**, 8245 (1994).
- J. P. Perdew, K. Burke, and M. Ernzerhof, *Phys. Rev. Lett.* **77**, 3865 (1996).
- C. Zhao, Z. Xu, H. Wang, J. Wei, W. Wang, X. Bai, and E. Wang, *Adv. Funct. Mater.* **24**, 5985 (2014).
- H. H. Qiu, Z. J. Wang, X. L. Sheng, *Physica B*, **421**, 46 (2013).
- Y. H. Kim, J. Choi, K. J. Chang, D. Tomanek, *Phys. Rev. B*, **68**, 125420 (2003).
- S. Qin, X. T. Guo, Y. Q. Cao, Z. H. Ni, Q. Y. Xu, *Carbon*, **78**, 559 (2014).
- D. Gao, Q. Xue, X. Mao, W. Wang, Q. Xu, D. Xue, *J. Mater. Chem. C*, **1**, 5909 (2013).