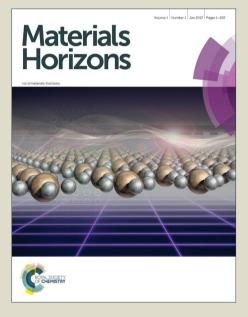
Materials Horizons

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about *Accepted Manuscripts* in the **Information for Authors**.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.



rsc.li/materials-horizons

Conceptual Insights

As the need for discovering new, improved materials continues to increase, it is becoming impractical – and in many cases impossible – to experimentally examine all possible leads. In order to explore deeper into this 'chemical space', many teams are harnessing the power of in silico experiments to perform High Throughput Virtual Screening (HTVS). In this paradigm, computational simulations are used to quickly screen potential materials, filtering away candidates with undesirable properties and allowing experimental searches to concentrate on a few, promising, candidate materials. Key to the success of this technique is the accuracy of the models employed; however, many traditional methods do not have a simple - or indeed known - mapping between the calculated properties, and those observed in the lab and results can be highly biased by the choice of calculation method. This study describes a practical solution to this problem; providing a framework in which current knowledge (e.g. experiments which are already performed) can be embedded within a probabilistic non-parametric mapping allowing a simple 'translation' between calculated and experimental properties. Since this is based upon a defined 'core' of knowledge, the method can also reply 'I don't know' when the translation asked is outside the scope of the knowledge provided. In addition, this study demonstrates how this technique can remove bias introduced by the choice of initial model – as all methods collapse into the same distribution once calculated.

Journal Name

ARTICLE

Received 00th January 20xx,

Accepted 00th January 20xx

DOI: 10.1039/x0xx00000x

www.rsc.org/

A Bayesian Approach to Calibrating High-Throughput Virtual Screening Results and Application to Organic Photovoltaic Materials

Edward O. Pyzer-Knapp^a, Gregor N. Simm^b, and Alán Aspuru Guzik^{a*}

A novel approach for calibrating quantum-chemical properties determined as part of a high-throughput virtual screen to experimental analogs is presented. Information on the moleccular graph is extracted through the use of extended connectivity fingerprints, and exploited using a Gaussian process to calibrate both electronic properties such as frontier orbital energies, and optical gaps and device properties such as short circuit current density, open ciruit voltage and power conversion efficiency. The Bayesian nature of this process affords a value for uncertainty in addition to each calibrated value. This allows the researcher to gain intuition about the model as well as the ability to respect its bounds.

Introduction

High-throughput virtual screening (HTVS) is a popular method for accelerating the discovery of both materials and pharmaceutical leads^{1–17}. In HTVS, simplifications or approximations are often required to make this procedure computationally tractable, resulting in subtle differences between experimental and theoretical property definitions. In these cases, it is likely that a calibration will be performed on the calculated results, in order to facilitate comparison between calibrated and experimental data.

In this study, we investigate the relationship between calculated and experimentally observed values for electronic and device properties of organic photovoltaics. This area has been the subject of many high-throughput screening efforts^{7–9}, such as the Harvard Clean Energy Project (CEP)^{6,18} in which errors in the model due to *in vacuo* calculations and oligomer vs polymer results were accounted for using an empirical linear calibration. Additionally, von Lilienfeld *et al.* have used a similar 'adjustment' method to predict the results of computationally expensive models from much simpler, cheaper, calculations¹⁹.

Here we present an advance upon this calibration technique, which takes into account both quantum chemical information, and information about the molecular graph. In addition, this technique reports an uncertainty alongside each calibration – providing a confidence that the method is being used appropriately.

Electronic Supplementary Information (ESI) available: [details of any supplementary information available should be included here]. See DOI: 10.1039/x0xx00000x

Computational Methods and Theoretical Background

Experimental Results and Theoretical Calculations

We recently reported the Harvard Organic Photovoltaic Dataset (HOPV15),²⁰ which contains experimental results for 266 donor materials from bulk heterojunction devices, alongside corresponding quantum-chemical calculations performed using the BP86^{21,22}, B3LYP^{21,23}, M06-2x^{24,25} and PBE0^{26,27} functionals and the def2-SVP²⁸ basis set on the BP86/def2-SVP optimized geometry.

Within HOPV15, in order to simplify the conformational landscape, and only sample conformers whose electronic structure contributes to changes in photovoltaic efficiency, calculations are performed on the photovoltaic core – i.e. the molecule with any solubilizing long hydrocarbon chains replaced with a single methyl group. If this reduction resulted in two experimental results referring to the same 'pruned' molecule, the result set containing the larger experimental value for the power conversion efficiency (PCE) was used, since this best represents the potential of the core.

The highest occupied molecular orbital (HOMO), lowest unoccupied molecular orbital (LUMO), power conversion efficiency (PCE), open circuit voltage (V_{OC}), and short circuit current density (J_{SC}), are used in this study, with computational values for PCE, V_{OC} and J_{SC} generated using the Scharber model²⁹, and all computational properties reported as Boltzmann averages over conformers.

Relating Molecular Structure to the Accuracy of Calculated Values

Many models, including density functionals^{23,24}, are constructed by fitting to values contained within a dataset. Since it is possible to reduce this fitting to a set of well-defined parameters, these are known as parametric methods. Parametric methods perform exceptionally well within the bounds of the model, but are



Page 2 of 9

^{a.} Department of Chemistry and Chemical Biology, Harvard University, 12 Oxford Street, Cambridge, MA, 02138, USA

^{b.} Department of Chemistry and Applied Biosciences ETH Zürich Vladimir-Prelog-Weg 2CH-8093 Zürich, Switzerland

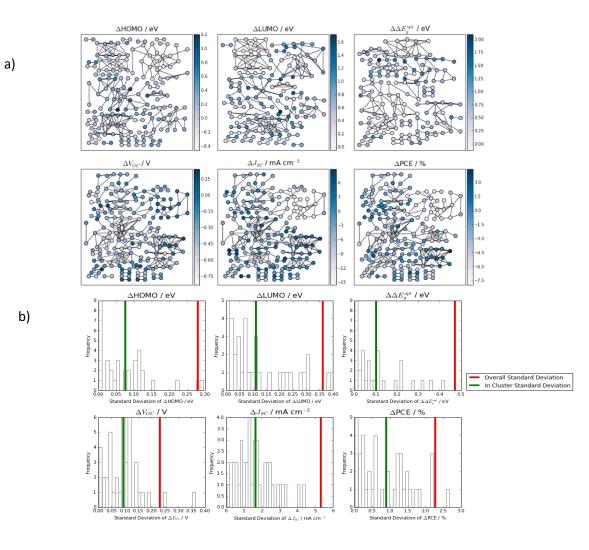


Figure 1 (*a*) Force directed graphs showing how the deviation between calculated and experimental values is cast into molecular space. The error in highest occupied molecular orbital energy, lowest unoccupied molecular orbital energy, optical gap (top row, left to right) and open circuit potential, short circuit current density and power conversion efficiency as calculated by the Scharber model (bottom row, left to right) are displayed via the colour of each node, which is itself representative of a molecule. Nodes are placed using the force-directed placement (FDP) algorithm, and connections are made between these nodes when their Tanimoto similarity between 512 bit, radius 2 Morgan circular fingerprints is > 0.65. (*b*) A histogram shows the distribution of values for the standard deviation of the 'delta' contained within each cluster. If there were no structure-property relationship, we would expect the standard deviation of nodes contained within the average cluster (plotted as a green line) to have a similar value to the standard deviation of the entire set (plotted in red). Since in each case the deviation within the average cluster is lower than this value, we are confident that a mapping between our feature space and the delta value exists.

ill defined when these bounds are broken.^{30,31} A good example of this is the failure of density functionals trained using ground state structures to reproduce transition states accurately³².

Additionally, potentially systematic sources of error are introduced through the difference between some calculated and experimental property definitions²⁹, the use of simplified model systems to reduce computational effort, and the use of an approximate density functional³³. If systematic failings and errors are related to the chemical structures of the molecules in question³⁴, casting this problem into molecular space could afford a method for applying appropriate corrections which take into account the chemical makeup of the molecules in question.

To test this hypothesis we construct force-directed graphs for each property for which we have both experimental and simulated

values. In order to cast this problem into chemical space, we extract information on each molecule as the 512-bit extendedconnectivity (Morgan circular) fingerprint³⁵, with the connectivity radius calculated at the two-bond level using the implementation in the RDKit³⁶. Each molecule is a node on this graph and nodes are connected when the Tanimoto similarity³⁷ is > 0.65. The FDP³⁸ algorithm, contained within the package Graphviz³⁹ was then used to structure the graph so that the edge-length was related to the similarity of molecules (i.e. the closer two connected nodes are, the more similar the molecules). Similar approaches have been used to perform 'materials cartography' – mapping chemical structure to properties in order to locate promising new materials⁴⁰. Each node on the graph was coloured to represent the error in the calculation, as defined by the difference between calculated and experimental values.

As can be seen from

Journal Name

Page 4 of 9

Figure 1a, broadly speaking clusters formed which appear on this graph share a similar deviation from experiment – as represented by the fill colour. This suggests that by mapping this problem onto molecular space, we will be able to perform a *per cluster* correction, and improve the accuracy of our calibration between calculated and experimental properties. This is shown in more detail in **Figure 1b** which clearly shows the difference in the standard deviation of 'deltas' within the average cluster (plotted in green) and the standard deviation of the entire set. If there were no correlation between features and deltas, then we would expect these two values to be much closer in magnitude.

Gaussian Processes

Gaussian process regressions (GPs) have been extensively studied by the machine learning community and are known for their sophisticated and consistent theory combined with relative computational tractability^{41,42}. They have been somewhat used within the scientific community to build structure-property relationship rules^{43–46}, although remain remarkably under-utilized, given their potential for strong predictive power.

While a probability distribution describes random variables which are scalars or vectors, a GP describes a distribution over functions. Within the framework of Bayesian inference, we can make predictions on an unknown data based upon input to target mappings described in a prior. This prior includes a covariance function – sometimes known as a kernel – which maps the covariance between function values. We build our covariance function upon the Tanimoto similarity $T(x_{\rho}, x_q)^{37}$, utilizing the popular squared exponential form for the kernel function itself:

$$k(x_{p}, x_{q}) = \sigma_{f}^{2} exp\left(-\frac{1}{2l^{2}}(1 - T(x_{p} - x_{q})^{2})\right) + \sigma_{pq} \delta_{n}^{2} \qquad (1)$$

Here, σ_f^2 is the signal variance, *I* is a length scale, δ_{pq} is the noise variance, and ${\delta_n}^2$ is a Kroneka delta, which takes the value 1 when p=q else 0. The values for these hyperparameters were trained to optimize the log marginal likelihood using the L-BFGS algorithm⁴⁷. In order to make a prediction, the Gaussian process places weights upon the functions of the prior distribution depending on how likely they are to model the target function. Thus the posterior distribution is sampled providing predictions (means), and also uncertainties for these predictions (standard deviations). In contrast to the earlier discussed parametric methods, a Gaussian process is non-parametric and thus very few assumptions need to be made about the target function. Additionally, the accuracy of a non-parametric method will only increase with the size of the prior (i.e. more data).

Results and Discussion

Gaussian processes were used to learn the deviation of computational results from their experimental analogues – i.e. to learn the function which calibrates one to the other.

In order to assess the performance of the calibration in a quantitative manner, we utilize two measures of error; the Pearson R coefficient, and a weighted RMSD. The Pearson R coefficient is a measure of the liner correlation between two variables, and is bounded at 0 (no correlation) and 1 (perfect correlation). In this study, increases in the Pearson R coefficient when the calibration has been applied are strongly indicative of an improvement of the performance. For the RMS Error, we include a weight to each point related to the uncertainty in the prediction returned by the Gaussian process. This weight is derived through normalizing the standard deviations of each prediction returned by the Gaussian process against the most certain point in the prediction. If these were not included, the measures would assume that each point is equally certain, thus removing a key piece of information from the scoring function

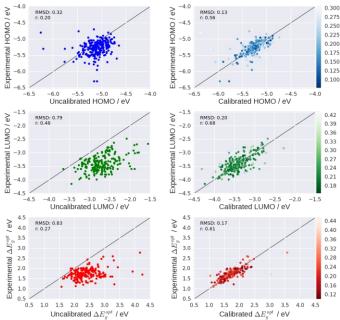


Figure 2 The results of calibrating B3YP/def2-SVP quantum-chemical results for the Highest occupied molecular orbital (HOMO), lowest unoccupied molecular orbital (LUMO) and optical gap to the experimental HOPV15 data set. The uncertainty in the calibrated values is represented in the fill colour; the lighter the colour, the more uncertain the calibration.

	σ_{f}	δ_n	l
PCE	3.575	1.145	0.327
JSC	19.760	2.167	0.377
VOC	2.000	0.180	2.969
Gap	0.688	0.197	1.133
номо	0.094	0.117	0.458
LUMO	0.233	0.194	0.670

Table 1
 Hyperparameters for each optimized calibrating Gaussian process

 for calibrating the B3LYP functional.
 These values for each additional

ARTICLE

functional are reported in the ESI.

In this way, we do not punish a poor calibration if it is also known to be highly uncertain. This was not done for the Pearson R coefficient, since bounding cases calculated on the distribution of points suggested that the un-weighted metric was representative.

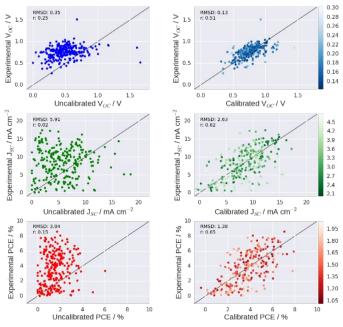
A brief examination of the hyperparameters optimized for each property (**Table 1**) shows that in each case the noise variance is much smaller than the signal variance. This is further evidence of the strength of the feature-to-delta mapping.

We first examine the performance of calibrating the electronic properties of these molecules; the HOMO, LUMO and gap. A plot of predicted and experimental values for these properties is shown in **Figure 2**. This plot shows the results when calibrating the B3LYP^{21,23} functional and def2-SVP basis set²⁸ – the results for the other functionals can be seen in the ESI, but are broadly similar. The hue of each point is related to the certainty of each prediction – the lighter the point, the more uncertain it is. Points were calibrated on a leave-one-out basis, in which the prior was formed using all points except the point being calibrated, with the process being repeated until all points had been calibrated.

It can be seen that whilst for these properties the quantumchemical method performs reasonably, there is significant improvement when the calibration is applied. This is especially true for the calculation of the gap – a key property in the prediction of the performance of photovoltaic materials. This may be due in part to the correction of a systematic error in assuming that the gap can be adequately described by the difference in energy levels of the LUMO and HOMO. For this assumption to hold well, the ionization potential (IP) and electron affinity (EA) would have to be well described by these frontier orbitals, which is not necessarily the case in DFT⁴⁸. The success of this calibration does show, however, that it is not necessary to calculate the EA and IP explicitly to rectify this error. This is of particular importance in the realm of highthroughput virtual screening, where an increase in the necessary number of calculations per molecule can swiftly accumulate, resulting in a significant decrease in the size of the libraries which can be screened.

Macroscopic properties, such as J_{SC} , V_{OC} , and PCE, present additional challenges to successful calibration. Since these properties have additional intermolecular contributions, capturing these in molecular fingerprints may prove challenging. Success in the prediction of lattice energy⁴⁹ and solubility⁵⁰ – properties both strongly related to intermolecular interactions – from the molecular structure does provide hope that these interactions can be somewhat captured, albeit in an implicit manner, in molecular fingerprints. Additionally, experimental measurements of device performance are notoriously noisy, introducing increasing amounts of uncertainty into the model. This method counters this through the application of some controlled noise to the data. The amount of noise was optimized against the log marginal likelihood to provide the function which was most robust to the data it was trained on.

Error! Reference source not found. shows the results of the calibration of $B3LYP^{21,23}/def2-SVP^{28}$ quantum chemical derived



properties to the experimental V_{oc} , J_{sc} and PCE extracted from HOPV15.

Figure 3 The results of calibrating B3YP/def2-SVP quantum-chemical results for the open-circuit potential (V_{oC}, short circuit current density (J_{sC}), and power conversion efficiency (PCE) to the experimental HOPV15 data set. The uncertainty in the calibrated values is represented in the fill colour; the lighter the colour, the more uncertain the calibration.

These properties were extracted from the quantum-chemical values for the frontier orbitals using the Scharber model²⁹. As with **Figure 2**, the colour of each calibrated point represents the uncertainty of the prediction, with light points being more uncertain.

The calibrated results in Error! Reference source not found. show significant improvements over the raw quantumchemical results, demonstrating the power of this model. It can be seen that the PCE is particularly poorly predicted by the combination of the Scharber model and quantum-chemical results.

While the Pearson R values shown in **Error! Reference source not found.** would suggest that the J_{SC} is predicted the worst, it is the failure of the model to predict PCEs > ~ 5% which is particularly troubling; since these are the very values we are interested in. Since the Scharber model was built using experimental data²⁹, whilst the performance of the raw quantum-chemical inputs is not surprising, it strongly

Journal Name

Materials Horizons Accepted Manuscrip

ARTICLE

illustrated the importance of a model which provides a 'warning' when it is not being used in situations for which it is designed.

Removal of Functional Dependence in High Throughput Screening

The low Pearson R and high weighted RMS error of the uncalibrated values highlight the importance of calibration when applying quantum-chemical results to certain photovoltaic properties. The variability of the raw results when different functionals are used for the calculation indicates another important function of calibration; the removal of functional dependence of DFT. Due to the different approaches employed in the construction of the functionals, properties such as the HOMO, LUMO and gap can have markedly different values when calculated with different functionals. This so-called *functional dependence* can have disastrous effects upon high-throughput virtual screening of organic photovoltaic materials since these differing values produce contrasting values for scoring functions such as the power conversion efficiency. Since these scoring functions are used to suggest molecules for experimental realization, it is important that there is a high degree of confidence in the quantum— chemical result. By removing this functional dependence, we reduce the probability that a particularly high-scoring molecule is simply an artefact of a particular combination of property and functional choice.

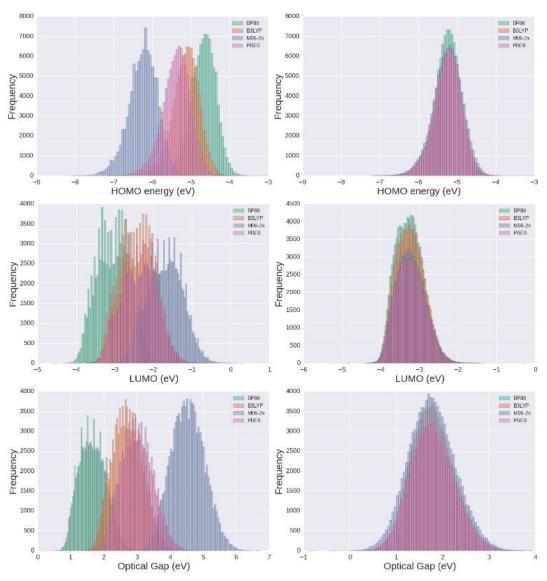


Figure 4 The distributions of the energies for the highest occupied molecular orbital (HOMO), lowest unoccupied molecular orbital (LUMO) and the optical gap before (left hand column) and after (right hand column) calibration calculated using the functionals BP86, B3LYP, M06-2X and PBE0 at the def2-SVP basis. It can be seen that the distributions of the values calculated before calibration have systematic differences, related to the amount of exact exchange present in each functional. After calibration, these distributions are much more closely aligned, hence the values calculated are significantly less dependent upon the initial functional choice; an important property in high throughput screening

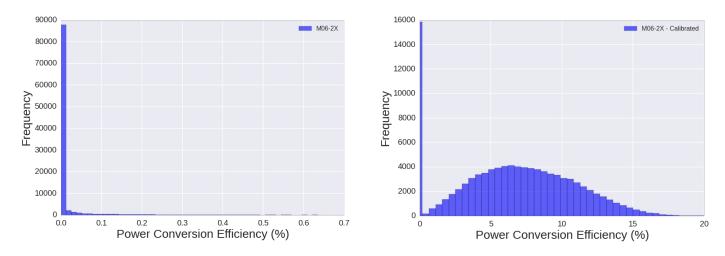


Figure 5 (left) The distribution of values for the Scharber power conversion efficiency calculated using the M06-2x functional at the def2-SVP basis are significantly skewed towards predicting a 0% power conversion efficiency, whilst after calibration (right) the distribution of values for the Scharber power conversion efficiency calculated using the M06-2x functional at the def2-SVP basis follows a much more reasonable distribution.

A good relationship between predicted and experimentally observed properties is paramount in the area of highthroughput screening, where design principles are built from trends in the data, and so confidence in the rankings produced through quantum-chemical calculations is essential.

In order to demonstrate this, 100,000 molecules were randomly selected from the CEPDB⁵¹ to represent the results of a high-throughput virtual screen. The CEPDB contains the properties for each BP86^{21,22}/def2-SVP²⁸ optimized geometry calculated at each of the functionals contained within this study. Since each functional is constructed in a different way, the properties calculated differ with the functional used.

Figure **4** shows the distributions of HOMO and LUMO energies as well as values for the optical gap. These values have been Boltzmann averaged over each conformer in a molecule for the 100,000 molecule set. Values have been calculated using four different functionals: B3LYP^{21,23}, BP86^{21,22}, PBE0^{26,27} and M06-2x^{24,25} at the def2-SVP²⁸ basis. It can be clearly seen in the left hand column of the plot (pre-calibration) that these values are very functional dependent. Since these values are used to calculate the PCE and other photovoltaic properties, **Figure 5, left)** is especially troubling and is most likely due to the overestimation of the gap by this functional which can been seen in the quantum-chemical predictions for this property by the MO6-2x functional contained within the ESI. If **Figure 5, right**), this functional dependence is removed; providing a greater confidence in the values returned.

Conclusions

We show how a Bayesian approach to calibration, here implemented as a Gaussian process with a prior based upon relevant experimental observations, is a robust method for relating the results of quantum chemical calculations to experiment. For each of the properties in the HOPV15 dataset

this is problematic as the properties, and identities, of candidates selected for further study should not depend upon the method used to calculate them. After calibration, however, these distributions strongly overlap, demonstrating a strong reduction in the functional dependence. Thus, whichever functional is chosen for the study, a similar answer is returned - a behaviour much more congruent with experiment, and thus affording a greater confidence that the calculated property is representative of what would be observed if the experiment were to be performed. This argument extends into the PCE, which is the primary fitness function used for many high throughput virtual screening efforts. Since the PCE as calculated by the Scharber model is extremely sensitive to the energy levels of the HOMO and LUMO of the electron donor, and the energy gap between them, small functional dependencies can manifest in vast differences in the distributions and values of PCEs obtained using the Scharber model.

A close examination of the PCEs predicted using a combination of the M06-2X functional and the Scharber model shows that M06-2X predicts essentially all structures to have a PCE of 0 (

this functional alone had been selected, and no calibration performed, the result of the screening would have been incredibly pessimistic, and may have discouraged further research. Once our calibration scheme is applied (studied here - HOMO, LUMO, gap, PCE, V_{oc} , and J_{sc} – values for the weighted RMS error between the calculated and experimental values, and the Pearson r correlation coefficient both improve significantly once the calibration is applied.

Additionally, we show how the application of this calibration method to the results of a high-throughput virtual screen removes the functional dependence of calculated properties, hence providing increased confidence that the result returned is directly comparable to experiment. This also increases the reliability and reproducibility of the candidate rankings, affording increased confidence that any extracted QSPR is an actual trend, and not an artefact of the choice of functional. Finally, the Bayesian nature of our proposed calibration results in a confidence in each calibration point being returned. This is an invaluable tool, since it can inform the user that the scheme is being used for systems for which it is not designed, or for which the prior is not informative.

Acknowledgements

Journal Name

A.A-G. and E.O.P-K acknowledges the Department of Energy through grant DE-SC0008733 for funding. This research would not have been possible without the use of the Harvard FAS Odyssey Cluster and support from FAS Research Computing. The authors also wish to thank IBM for organizing the World Community Grid and the WCG members for their computing time donations.

Notes and references

- 1 E. O. Pyzer-Knapp, C. Suh, R. Gomez-Bombarelli, J. Aguilera-Iparraguirre and A. Aspuru-Guzik, *Annu. Rev. Mater. Res.*, 2015, **45**, null.
- B. Huskinson, M. P. Marshak, C. Suh, S. Er, M. R. Gerhardt, C. J. Galvin, X. Chen, A. Aspuru-Guzik, R. G. Gordon and M. J. Aziz, *Nature*, 2014, 505, 195–198.
- 3 B. K. Shoichet, *Nature*, 2004, **432**, 862–865.
- 4 J. Bajorath, Nat. Rev. Drug Discov., 2002, 1, 882–894.
- 5 M. D. Halls and K. Tasaki, *J. Power Sources*, 2010, **195**, 1472–1478.
- 6 J. Hachmann, R. Olivares-Amaya, S. Atahan-Evrenk, C. Amador-Bedolla, R. S. Sanchez-Carrera, A. Gold-Parker, L. Vogt, A. M. Brockway and A. Aspuru-Guzik, *J Phys Chem Lett*, 2011, 2, 2241–2251.
- 7 I. Y. Kanal, S. G. Owens, J. S. Bechtel and G. R. Hutchison, *J. Phys. Chem. Lett.*, 2013, **4**, 1613–1623.
- 8 N. M. O'Boyle, C. M. Campbell and G. R. Hutchison, *J. Phys. Chem. C*, 2011, **115**, 16200–16210.
- 9 Y. Shu and B. G. Levine, J. Chem. Phys., 2015, 142, 104104.
- 10 Y. J. Colón, D. Fairen-Jimenez, C. E. Wilmer and R. Q. Snurr, *J. Phys. Chem. C*, 2014, **118**, 5383–5389.
- M. D. Halls, P. J. Djurovich, D. J. Giesen, A. Goldberg, J. Sommer, Eric McAnally and M. E. Thompson, *New J. Phys.*, 2013, **15**, 105029.
- 12 M. D. Halls, D. J. Giesen, T. F. Hughes, A. Goldberg and Y. Cao, 2013, vol. 8829, pp. 882926–882926–6.
- 13 E. S. Kadantsev, P. G. Boyd, T. D. Daff and T. K. Woo, *J. Phys. Chem. Lett.*, 2013, **4**, 3056–3061.
- 14 S. Curtarolo, G. L. W. Hart, M. B. Nardelli, N. Mingo, S. Sanvito and O. Levy, *Nat. Mater.*, 2013, **12**, 191–201.
- 15 M. Korth, Phys. Chem. Chem. Phys., 2014, 16, 7919–7926.
- 16 C. E. Wilmer, M. Leaf, C. Y. Lee, O. K. Farha, B. G. Hauser, J. T. Hupp and R. Q. Snurr, *Nat. Chem.*, 2012, **4**, 83–89.
- 17 A. Jain, S. P. Ong, G. Hautier, W. Chen, W. D. Richards, S. Dacek, S. Cholia, D. Gunter, D. Skinner, G. Ceder and K. A. Persson, *APL Mater.*, 2013, **1**, 011002.
- J. Hachmann, R. Olivares-Amaya, A. Jinich, A. L. Appleton, M.
 A. Blood-Forsythe, L. R. Seress, C. Roman-Salgado, K. Trepte, S.

Atahan-Evrenk, S. Er, S. Shrestha, R. Mondal, A. Sokolov, Z. Bao and A. Aspuru-Guzik, *Energy Env. Sci*, 2014, **7**, 698.

- 19 R. Ramakrishnan, P. O. Dral, M. Rupp and O. A. von Lilienfeld, J. Chem. Theory Comput., 2015, **11**, 2087–2096.
- 20 Edward O. Pyzer-Knapp, Gregor Simm, Trevor Lutzow, Kewei Li, Laszlo Serres, Johannes Hachmann and Alan Aspuru-Guzik, Submitted.
- 21 A. D. Becke, *Phys Rev A*, 1988, **38**, 3098–3100.
- 22 J. P. Perdew, Phys Rev B, 1986, 33, 8822–8824.
- 23 A. D. Becke, J. Chem. Phys., 1993, 98, 5648–5652.
- 24 Y. Zhao and D. Truhlar, Theor. Chem. Acc., 2008, 120, 215–241.
- 25 Y. Zhao and D. G. Truhlar, *J. Chem. Theory Comput.*, 2007, **3**, 289–300.
- 26 J. P. Perdew, M. Ernzerhof and K. Burke, J. Chem. Phys., 1996, 105, 9982–9985.
- J. P. Perdew, K. Burke and M. Ernzerhof, *Phys Rev Lett*, 1996, 77, 3865–3868.
- 28 F. Weigend and R. Ahlrichs, *Phys Chem Chem Phys*, 2005, **7**, 3297–3305.
- M. C. Scharber, D. Mühlbacher, M. Koppe, P. Denk, C.
 Waldauf, A. J. Heeger and C. J. Brabec, *Adv. Mater.*, 2006, 18, 789–794.
- 30 P. J. Silva and M. J. Ramos, *Comput Theor Chem*, 2011, **966**, 120–126.
- 31 J.-W. Song, T. Tsuneda, T. Sato and K. Hirao, *Theor Chem Acc*, 2011, **130**, 851–857.
- 32 X. Xu, I. M. Alecu and D. G. Truhlar, J. Chem. Theory Comput., 2011, 7, 1667–1676.
- A. J. Cohen, P. Mori-Sánchez and W. Yang, *Chem Rev*, 2012, 112, 289–320.
- 34 D. E. Edwards, D. Y. Zubarev, A. Packard, W. A. Lester and M. Frenklach, *Phys. Rev. Lett.*, 2014, **112**, 253003.
- 35 D. Rogers and M. Hahn, J. Chem. Inf. Model., 2010, 50, 742– 754.
- 36 G. Landrum, RDKit: Open-source cheminformatics, rdkit.org
- 37 T. T. Tanimoto, An elementary mathematical theory of classification and prediction, IBM Report (November, 1958), cited in: G. Salton, Automatic Information Organization and Retrieval, McGraw-Hill New York, 1968.
- 38 T. M. J. Fruchterman and E. M. Reingold, Softw Pr. Exper, 1991, 21, 1129–1164.
- 39 E. R. Gansner and S. C. North, Softw Pr. Exp, 2000, 30, 1203– 1233.
- 40 O. Isayev, D. Fourches, E. N. Muratov, C. Oses, K. Rasch, A. Tropsha and S. Curtarolo, *Chem. Mater.*, 2015, **27**, 735–743.
- 41 M. Seeger, Int J Neural Syst, 2004, **14**, 69–106.
- 42 C. E. Rasmussen, Gaussian processes for machine learning, MIT Press, 2006.
- 43 O. Obrezanova, G. Csányi, J. M. R. Gola and M. D. Segall, *J. Chem. Inf. Model.*, 2007, **47**, 1847–1857.
- 44 F. R. Burden, J. Chem. Inf. Comput. Sci., 2001, 41, 830–835.
- 45 A. Schwaighofer, T. Schroeter, S. Mika, J. Laub, A. ter Laak, D. Sülzle, U. Ganzer, N. Heinrich and K.-R. Müller, J. Chem. Inf. Model., 2007, 47, 407–424.
- 46 P. Gao, A. Honkela, M. Rattray and N. D. Lawrence, *Bioinformatics*, 2008, **24**, i70–i75.
- R. Byrd, P. Lu, J. Nocedal and C. Zhu, *SIAM J Sci Comput*, 1995, 16, 1190–1208.
- 48 S. Kümmel and L. Kronik, *Rev. Mod. Phys.*, 2008, **80**, 3–60.
- 49 C. Ouvrard and J. B. O. Mitchell, *Acta Cryst Sect B*, 2003, **59**, 676–685.

ARTICLE

- 50 J. L. McDonagh, N. Nath, L. D. Ferrari, T. van Mourik and J. B. O. Mitchell, *J. Chem. Inf. Model.*, 2014, **54**, 844–856.
- 51 *Harvard. Clean Energy Project. Database* cepdb.molecularspace.org