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Widely Available Active Sites on Ni₂P for Electrochemical Hydrogen Evolution – Insights from First Principles Calculations

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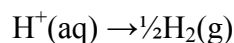
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ABSTRACT: We present insights into the mechanism and the active site for hydrogen evolution on nickel phosphide (Ni₂P). Ni₂P was recently discovered to be a very active non-precious hydrogen evolution catalyst. Current literature attributes the activity of Ni₂P to a particular site on the (0001) facet. In the present study, using Density Functional Theory (DFT) calculations, we show that several widely available low index crystal facets on Ni₂P have better properties for a high catalytic activity. DFT calculations were used to identify moderately bonding nickel bridge sites and nickel hollow sites for hydrogen adsorption and to calculate barriers for the Tafel pathway. The investigated surfaces in this study were the $(10\bar{1}0)$, $(\bar{1}\bar{1}20)$, $(11\bar{2}0)$, $(11\bar{2}1)$ and (0001) facets of the hexagonal Ni₂P crystal. In addition to the DFT results, we present experiments on Ni₂P nanowires growing along the $\langle 0001 \rangle$ direction, which are shown as an efficient hydrogen evolution catalyst. The experimental result adds these nanowires to the variety of different morphologies of Ni₂P, which are all active for HER.

1. Introduction

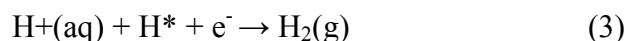
Water splitting by electrolysis or photocatalysis is attracting attention as a prospective sustainable source of hydrogen for energy storage applications^{1,2,3}. The hydrogen evolution reaction (HER) is the cathode half reaction:



The most active and stable catalysts in acid environments for HER are currently based on precious metals⁴. At low over-potentials, HER from Pt(111) is dominated by the Volmer-Tafel mechanism^{5, 6, 7}, which can be written as:



This requires two H atoms to be adsorbed in proximity for fast diffusion and recombination. Alternatively the reaction may take place via a Volmer step (1), followed by the Heyrovski step:



Reaction (3) is considered to be relevant at high over-potentials only⁷

In the last decade, several classes of non-precious materials have been found to be active catalysts for the HER. MoS₂ was proven as a promising non-precious HER catalyst material, which is stable in a wide pH range. However MoS_x catalysts are not as active as platinum and they are only active at undercoordinated edge sites^{8, 9, 10, 11}. Hydrodesulfurization (HDS) catalysts such as Mo₂C and MoB have recently attracted attention as hydrogen evolution catalysts with good stability in both acid and alkaline solution^{12, 13}. Ni₂P has also previously been subject to experimental and theoretical studies for catalysis of hydrodesulfurisation^{13, 14, 15, 16}, water-gas-shift¹⁷, and more recently for hydrogen evolution in acid^{18, 19, 20}. The active sites and the details of the mechanism remain unknown for several of the newly discovered HER catalysts including Ni₂P. The observed Tafel slopes of the Ni₂P catalysts are similar to what is observed for MoS₂ edges, and the specific activity is one of the highest of the non-precious catalysts. Further experimental and theoretical studies can yield new insights for further design of electro-catalysts which is the motivation of this study. In this paper, Density Functional Theory (DFT) calculations are combined with experiments on high aspect ratio crystalline Ni₂P nanowires to understand the mechanism of HER from Ni₂P.

The trends in hydrogen evolution activity over various transition metals^{21, 22}, and various metal- non-metal combinations have been investigated using DFT. The free energy of hydrogen adsorption, ΔG_{H*}, has been established as a descriptor for predicting the exchange current density^{21, 22, 23, 24}, of transition metal catalysts. The best catalysts have free binding energies close to 0.0 eV, which is explained by the Sabatier principle²⁵; Stronger binding results in hydrogen poisoning, leaving no free sites for adsorption. Weaker binding results in a high overpotential needed to adsorb protons. Transition metal catalysts have binding energies that are slightly coverage dependent. On Pt(111) the binding energy calculated with DFT varies from -0.12 eV at low coverage to 0.04 eV at one monolayer coverage²¹. In

addition, it has been shown by DFT calculations that platinum has no significant diffusion barrier between the adsorbing hollow sites on the (111) facet²⁶.

Liu and Rodriguez have published several works on DFT calculations of hydrogen binding energies on the Ni₂P (0001) surface. In 2005, they predicted the hydrogen evolution activity of the Ni₂P to be comparable to that of hydrogenase based on trends in adsorption energies. The ensemble of phosphor atoms available as proton acceptors next to moderately binding metal hollow sites and weak binding Ni-P bridge is proposed to facilitate HER catalysis¹⁸. The binding energies can be compared with the trends in HER exchange current density calculated by Nørskov et al.²¹. In these studies, the binding free energies are presented, which means the calculated potential energy is corrected by +0.24eV accounting entropy and differences in zero point energy. When comparing the binding energies at the relevant coverage^{7,27}, one observes that the metal hollow sites on Ni₂P will be fully occupied and that the Ni-P sites will need an overpotential of at least 0.31V. This does not agree with the very high activity observed in new studies of hollow and multifaceted Ni₂P nanoparticle catalysts^{19,20,28,29}. Another active site may therefore exist, and we investigate this using more detailed DFT calculations. In addition we report experiments showing, that nano wires grown along the (0001) direction are highly active for HER, supporting the results of the present DFT calculations.

2.1. Calculation methods

Calculations were carried out using the GPAW code with projector augmented wave functions on a real space grid and ultra soft pseudopotentials³⁰. The RPBE³¹ functional was used for the exchange-correlation contribution. The Ni₂P crystal structure³² were copied from The Materials Project³³, and imported to the Atomic Simulation Environment (ASE)³⁴. Bulk lattice constants were converged to $a=5.151\text{\AA}$ and $c=3.408\text{\AA}$ using a third order polynomial fit to the energy versus the lattice constant³⁵.

A 1×1 supercell was used for the (0001) and $(11\bar{2}1)$ surfaces, a 3×1 supercell was used for the $(10\bar{1}0)$ surface and a 1×2 supercell was used for the $\pm(11\bar{2}0)$ surfaces. The (0001) surface had five atomic layers, the $(11\bar{2}1)$ and $(10\bar{1}0)$ slabs had four atomic layers, and the $\pm(11\bar{2}0)$ surfaces had three atomic layers. The dipole across the unit cells were all less than 0.10 V

A Monkhorst-Pack³⁶ sampling of $2 \times 2 \times 1$ k-points was used for the large $(11\bar{2}1)$ surfaces, $2 \times 3 \times 1$ k-points was used for the $\pm(11\bar{2}0)$ surfaces, $3 \times 2 \times 1$ k-points for the $(10\bar{1}0)$ slabs and for the (0001) surface a $3 \times 3 \times 1$ k-point sampling was used. The grid spacing was 0.18 Ångström in all calculations.

The slabs were relaxed using the Broyden-Fletcher-Goldfarb-Shanno LineSearch algorithm within ASE until the forces were below 0.01 eV/Ångström. A recalculation was later carried out with double the k-points sampling and a grid spacing of 0.12 Ångström for the most interesting slabs and adsorbate configurations. The resulting differences in adsorption energies did not exceed 0.02 eV, which is below the accuracy usually attributed to DFT.

The choice of facets and surface termination were based on calculations of the minimum energy configuration of all different ways to cut the crystal in the lowest index planes. Several STM and LEED studies^{37, 38, 39, 40} show the Ni₂P surfaces can have a stable phosphor termination, but under the hydro-thermal treatment before testing, it is expected they lose the phosphor layer and expose active metal sites, as on the structures investigated in our calculations. The chosen structures are $(10\bar{1}0)$, $(11\bar{2}0)$, $(\bar{1}\bar{1}20)$, $(11\bar{2}1)$ and the Ni₃P terminated (0001) , as shown in figure 1 along with their (hkjl) indices from a top view.

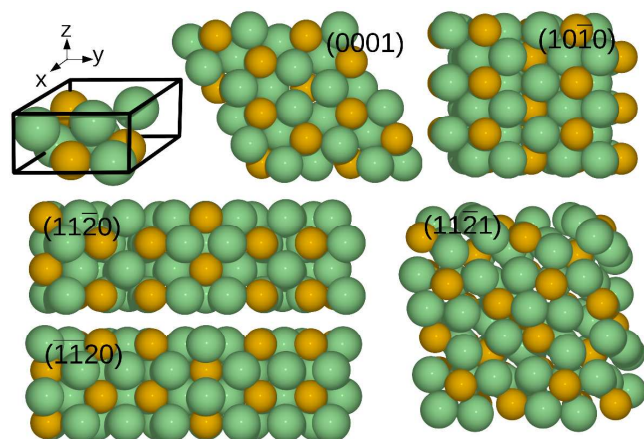


Figure 1: The Ni₂P unit cell and top views of the low index surfaces used in this study. The surface unit cells have been repeated once in the X and Y directions, apart from the $(10\bar{1}0)$, for which a 2×3 supercell is shown. Nickel atoms are represented by green spheres and phosphor by yellow spheres.

The free adsorption energies are found from calculated potential energies by correcting for the gas phase entropy $\Delta S = -S_0(\text{H}_2)$ and the difference in zero point energy ΔZPE using the equation

$$\Delta G = \Delta E - T\Delta S + \Delta ZPE$$

where E denotes the ground state energy at 0 K obtained from DFT. The correction for ZPE and entropy makes an addition of 0.24 eV to the adsorption energy of a hydrogen atom^{21, 41}.

The calculation of adsorption free energies and adsorbate coverages was carried out using the self consistent scheme as in the work by Skulason et al.⁷. The integral adsorption energy $G_{\text{int}}(n)$ is

$$G_{\text{int}} = (G(N, n) - G(N, 0) - n\mu_{\text{H}}), \text{ where } \mu_{\text{H}} = 1/2G_{\text{H}_2} - eU$$

where $G(N, n)$ is the free energy of a surface which includes of N nickel atoms in the top layer and n adsorbed hydrogen atoms. Thus we use a definition of the coverage $\theta = n/N$, where n is the number of hydrogen adsorbates and N is the number of Ni atoms on the surface. The configuration which is relevant at a given chemical potential of protons and electrons is the one with the minimum integral adsorption energy ($G_{\text{int}}(n)$). Using the computational hydrogen electrode approach⁴¹ with self-consistent coverage and adsorption free energies²⁷, the adsorption phase diagram was calculated for every facet. The required over-potential to favor an intermediate step is given through the differential

binding energy, which is given by:

$$\Delta G_{\text{diff}}(n) = (G_{\text{int}}(n) - G_{\text{int}}((n - 1)))$$

We apply $\eta = \Delta G_{\text{diff}}(n)/e$ to find the minimum required over-potential to adsorb the intermediate⁴¹, which enables a surface recombination reaction (the Tafel step). The barriers were calculated using the Nudged Elastic Band (NEB) method with a climbing image⁴².

Experimental methods

Materials

The commercially available materials were used as received: nickel acetylacetonate ($\text{Ni}(\text{acac})_2$ for synthesis, VWR), oleic acid (suitable for cell culture, BioReagent, Aldrich), oleylamine (OA approximate C18-content 80-90%, Acros), trioctylphosphine (TOP 90%, technical grade, AcroSeal®, Acros), ethanol (absolute alcohol without additive $\geq 99.8\%$, Aldrich), n-hexane (HiPerSolv CHROMANORM® for HPLC 97%, VWR), sulfuric acid (H_2SO_4 volumetric 1 M, Aldrich), Nafion® (117 solution, Aldrich). Water was first purified using Millipore Milli-Q® Integral water purification system (18.2 M Ω cm resistivity). All syntheses were performed in strict air-free conditions using glove-box and Schlenk line techniques unless stated otherwise.

Synthesis of nickel phosphide nanowires

The synthesis was performed following a previous report⁴³. A stock solution containing 0.75 mmol of $\text{Ni}(\text{acac})_2$, 1.8 mmol of oleic acid and 10 mL of OA was heated at 120 °C. Then the stock solution was very slowly injected (0.05 mL min⁻¹) using a syringe pump into a stirred mixture containing 5 mL of OA and 2.4 mmol of TOP heated at 320 °C under reflux. The reaction mixture was kept at 320 °C under reflux until the stock solution was used up. Over the course of the reaction, the mixture color changes from transparent to dark yellow, orange and finally black. After cooling to room temperature, the products were washed using a mixture of hexane and ethanol, and separated by centrifugation. The supernatant was discarded and the washing step was repeated two more times. After the supernatant was discarded the nanowires were collected on a watch glass and dried in an oven at 50 °C overnight.

Physical characterization

X-ray diffraction (XRD) were carried out on a X'Pert Philips diffractometer in Bragg-Brentano geometry with Cu $K_{\alpha 1}$ radiation ($\lambda = 0.1540$ nm) and a fast Si-PIN multi-strip detector. The tube source was operated at 45 kV and 40 mA. The Scherrer equation was used to calculate the average crystallite size of the powder. The diffraction pattern was analyzed and compared with references in the international center of diffraction data (ICDD). Transmission electron microscope (TEM) images were taken on a Philips (FEI) CM12 with a LaB₆ source operated at 120 kV accelerating voltage. Specimens were prepared by ultrasonic dispersion of the samples in ethanol. The suspension was mixed with a micropipette by several suction-release cycles to ensure representative and reproducible TEM sample. A few drops of the mixed suspension were deposited onto the carbon-coated grid.

Electrochemical measurements

Electrochemical measurements were recorded using a Gamry Instruments Reference 3000™ potentiostat. A traditional three-electrode configuration was used. For polarization and electrolysis measurements, a platinum wire was used as the auxiliary electrode and a double-junction Ag/AgCl (KCl saturated) electrode was used as the reference electrode. Both counter and reference electrodes were rinsed with distilled water and dried with compressed air prior to measurements. Potentials were referenced to a reversible hydrogen electrode (RHE) by adding a value of $(0.2 + 0.059 \times \text{pH})$ V. The current was normalised over the geometric surface area of the electrode. Ohmic drop was corrected using the current interrupt method. A total electrolyte volume of ~ 50 mL was used to fill the glass cell. All potentials were converted and referred to the RHE unless stated otherwise. The electrolyte used throughout all electrochemical experiments was a 1 M H_2SO_4 solution. During electrochemical experiments, the electrolyte was agitated using a magnetic stirrer rotating at 300 rpm. For all electrochemical experiments a glassy carbon electrode ($\sim 0.071 \text{ cm}^2$) was used as the working electrode. The cyclic voltammetry (CV) experiments were conducted in 1 M H_2SO_4 at 25 °C using a scan-rate of 5 mV s^{-1} across a potential window of $-0.3 - 0.1 \text{ V vs. RHE}$.

Pretreatment of the working electrode

Prior to loading of catalyst, the working electrode was pretreated to achieve better performance⁴⁴. First, the electrode was manually polished using alpha alumina powder (CH Instruments, Inc.) with decreasing grain sizes (typically 0.3 and 0.05 μm) on a 73 mm diameter nylon polishing pad (MasterTex, Buehler). Between each polishing step the electrode was rinsed with deionized water and ultrasonicated in distilled water for 10 seconds. Then the electrode was dried using compressed air. This ultrasonication and drying cycle was repeated two more times: once in absolute ethanol and a second time with distilled water. This polishing process resulted in a shiny mirror finish. The bare working electrode was subjected to a constant potential of 2 V vs. RHE in 1 M H_2SO_4 under vigorous stirring at 25 °C over 1 hour. Then, the pretreated electrode was rinsed with absolute ethanol and dried with compressed air prior to catalyst deposition.

Electrode preparation

Before being deposited on the working electrode, the nanowires were annealed at 450 °C for 4 hours under 5% H_2/N_2 gas in order to remove any surfactant present on the nanowires surface^{19, 43}. The catalyst was loaded on a pretreated working electrode *via* drop-casting of 10 μL of catalyst ink, equivalent to a loading of 1.42 mg cm^{-2} . The catalyst suspension was a 500 μL solution consisting of 5 mg of catalyst, 400 μL of distilled water, and 100 μL of absolute ethanol. The slurry was ultrasonicated for 5 hours and mixed with a micropipette by several suction-release cycles prior to deposition to ensure representative and reproducible catalyst sample. The temperature of the ultrasonic bath was kept below 45°C at all time to avoid any undesired heating effect. The nanowires are ultra-sonicated for 5 h to minimize aggregation. Once the catalyst was deposited, the electrode was dried in oven at 50 °C for 10 minutes. 5 μL of 0.2% Nafion[®] solution in absolute ethanol was then drop-casted on the glassy carbon surface and the electrode was dried in air for 5 minutes prior to electrochemical measurement.

3. Results and discussion

XRD measurement was conducted on the annealed Ni₂P nanowires (Fig. 2). According to the Scherrer's equation, the average crystallite size of the nanowires is 11 nm. Ni₂P is the predominant species in the sample, while small amounts of Ni₁₂P₅ and NiO are also present. Recently, Zhang *et al.* reported that pure Ni₁₂P₅ has similar activity to Ni₂P for HER in acid⁴⁵. We showed earlier than NiO is less active than Ni₂P; furthermore, NiO tends to dissolve at a reductive potential in acid. Thus, the HER activity of the nanowires (see below) can be mostly attributed to Ni₂P.

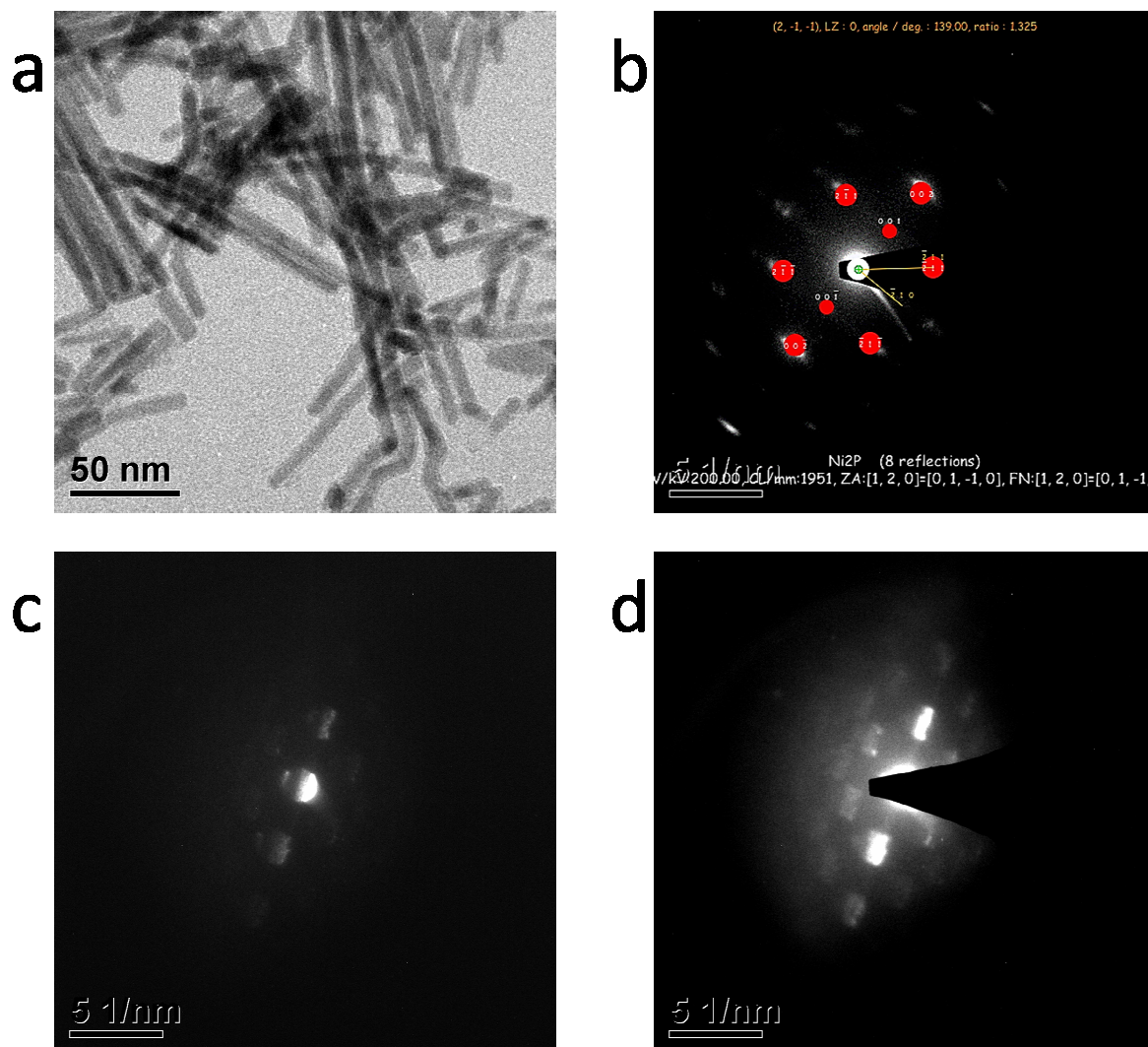


Figure 2: (a) Representative TEM image of the annealed Ni₂P nanowires. (b) Focused SAED of nanowires crystallites indexed using JEMS[®] software. (c) and (d) Unfocused SAED patterns whose complementary information served to identify the growth direction of the nanowires.

The transmission electron microscope (TEM) image and selected area electron diffraction (SAED) patterns of the nanowires are shown on Fig. 2a-d. The nanowires are rather uniform and the width of the nanowires observed by TEM is in agreement with the crystallite size calculated from the XRD pattern. The focused SAED pattern (Fig. 2b) was indexed using the JEMS software (P. Stadelmann, JEMS, EPFL). The line linking the (000n) spots is the (0001) direction. The nanowire can be observed on the SAED pattern when the focus is reduced (Fig. 2c and d). Comparison between the indexed (Fig. 2b) and unfocused SAED patterns (Fig. 2c and d) allows the determination of the growth direction. The results indicate that the nanowires grow along the *c*-axis, *i.e.*, the <0001> direction.

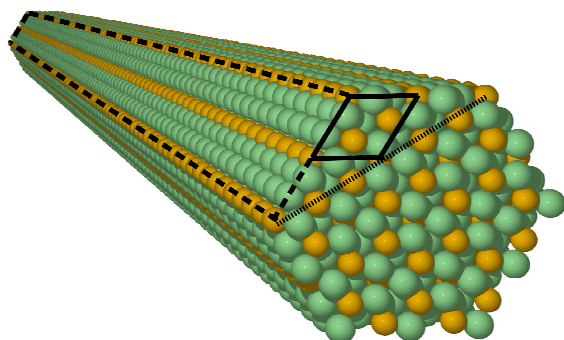


Figure 3: Model of nanowire crystal. This illustrates the exposed facets. The (0001) plane of the unit cell is shown with solid lines, the coarsely dashed lines show a $(10\bar{1}0)$ surface, and the finely dashed line shows the $\pm(11\bar{2}0)$ planes. Ni atoms are shown in green and phosphorus atoms in yellow.

Thus, the nanowires do not expose many (0001) facets, since the (000n) planes exist only in their cross section (see figure 3). The low index facets $\pm(11\bar{2}0)$ and $(10\bar{1}0)$ are assumed to be widely available on the samples. These morphologies would be expected to exhibit a less than optimal performance if the ensemble on the Ni_3P terminated (0001) facet was the only active site.

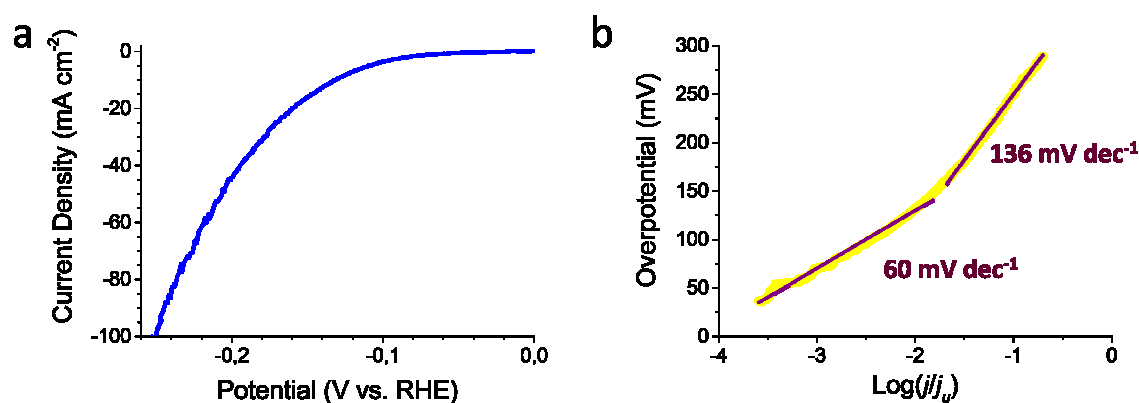


Figure 4: (a) Polarization curve of Ni_2P nanowires in 1 M H_2SO_4 at 25 °C at a scan rate of 5 mV s⁻¹. (b) Tafel analysis of the polarization

Fig. 4 shows the HER activity of the Ni_2P nanowires in 1 M H_2SO_4 . The nanowires are excellent HER catalysts. This result combined with the calculations (see below) does not indicate the (0001) ensemble is the only active site of HER from Ni_2P , although it is not possible to separate contributions from the various facets and sites in the experiments. The overpotential to drive a current density of 10 mA cm⁻² is 133 mV (Fig. 4a). Two Tafel slopes are observed. At $\eta < 125$ mV, the Tafel slope is about 60 mV dec⁻¹, while $\eta \approx 125 - 275$ mV, the Tafel slope is 126 mV dec⁻¹ (Fig. 4b). Despite the difference in sample preparation and morphology, the catalytic activity of the Ni_2P nanowires is very similar to the activity of hollow and multifacet Ni_2P nanoparticles reported by Schaak *et al.*¹⁹, that of polydispersed Ni_2P nanoparticles reported by our group²⁰ and that of high surface area Ni_2P nanoflakes reported by Han *et al.*²⁹. This suggests that widely available facets or sites of Ni_2P are active. This is contrary to MoS_x ,

where only specific morphologies are efficient because the edge sites need to be exposed⁸. The facets with $(a\ b\ -(a+b)\ 0)$ indices are propagating along the growth direction of the nanowires (see figure 3), and are thus expected to be abundant in the sample.

The DFT calculations show the most strongly binding sites tend to be nickel bridge sites or nickel 3-fold hollow sites. The nickel bridge sites are found in continuous rows on the $(10\ \bar{1}\ 0)$ and $(\bar{1}\bar{1}\ 20)$ facets and adjacent to nickel hollow sites on the $(11\ \bar{2}\ 1)$ facet. (See supplementary materials for geometries and adsorption energies). Adsorption on Ni-P sites and on top P atoms were generally 0.2 eV to 0.4 eV less stable than the metal sites. Thus, adsorption on Ni-P sites and on top P would require a higher over-potential.

It takes a high over-potential to favor a high coverage of atomic hydrogen on all the facets. Thus, Ni₂P is on the weak binding side of the Sabatier volcano. Adsorbates on Ni₂P interact strongly leading to a steeper increase in binding energy with increasing coverage (see supplementary material for the surface phase diagrams), compared to most transition metals including platinum.

The (0001) facet has the hollow site occupied at equilibrium, but as shown in figure 6, it requires a large over-potential of 0.41 V to stabilize the initial Ni-P bridge state for the Tafel step. It is also improbable that adsorbates from separate hollow sites on (0001) can combine since they are not neighbouring. This was confirmed by calculating the diffusion barriers on the (0001) and the $(10\ \bar{1}\ 0)$ facets (See supplementary information). The $(10\ \bar{1}\ 0)$, $(\bar{1}\bar{1}\ 20)$ and the $(11\ \bar{2}\ 1)$ facets only require an over-potential of 0.0 V, 0.06 V and 0.19 V respectively to stabilise neighbouring H* or H*, which are mobile in the row of nickel atoms. These results suggest that H* in Ni hollow sites on the (0001) facet do not contribute to the HER rate at low over-potentials. To investigate further, the barriers for the Tafel step were calculated as described in the following.

Previous calculations on Pt(111) suggest that the Tafel pathway is faster than the Heyrovski pathway at low over-potentials. If proton transfer is faster to H*(Ni₂P) than to H*(Pt(111)), it is possible that isolated metal hollow sites on the (0001) facet play a role. In the following, it is assumed that the surface coverage is in equilibrium with protons in solution, which means that the Tafel step is rate limiting. Calculation of the barrier for proton transfer (Volmer or Heyrovski step) is out of the scope of this paper, since it requires very precise information on the interfacial structure⁴⁶.

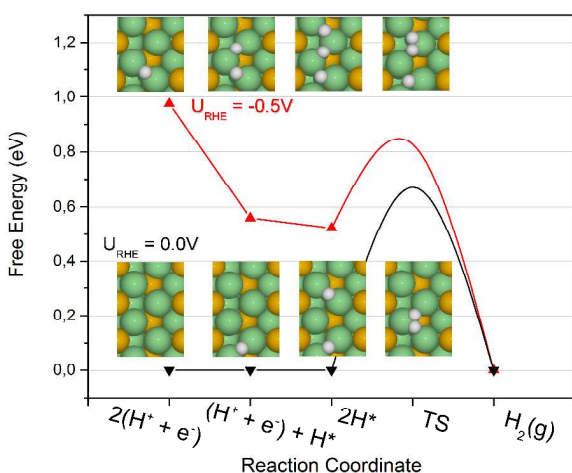


Figure 5: Free Energy pathway of the Volmer-Tafel reaction on the $(10\bar{1}0)$ facet. The configuration of every point are shown by the insets, where Ni is shown in green, P in yellow and H in white. The energy of the transition states are shown by the peak of the splines. These were calculated using the Nudged Elastic Band method with a climbing image⁴¹. This shows how G^{TS} is reduced by applying a higher over-potential.

As shown in figure 5 for, the rate limiting HER barrier is G^{TS} , where G^{TS} is the energy of the transition state relative to the $2(\text{H}^+ + \text{e}^-)$ state. The initial states were the most stable configurations at the lowest possible over-potential according to the surface phase diagrams. As shown from calculations presented in figure 5, G^{TS} can be lowered by further increasing the over-potential until the free energy of protons in solution are at the Tafel transition state level. This agrees well with the exchange current density being a good indication for the activity at higher over-potentials.

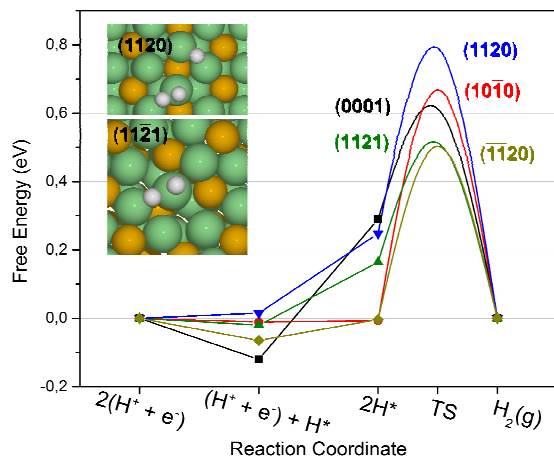


Figure 6: Free Energy diagram assuming a Tafel Mechanism. Calculated barriers for the Tafel step are shown by the peak of the splines for the configurations, which are most stable at 0V vs RHE. The insets show transition states of the two facets with the lowest Tafel barriers. Ni atoms are shown in green, P in yellow H in white.

The results for a Tafel pathway are summarized in figure 6, with comparison between the studied facets. The adsorption energies which are easy to calculate are usually a good descriptor, since they are

expected to scale with the transition state energies. In the case of a Tafel pathway, it is more accurate to use calculated G^{TS} to compare the activity from the different facets, since it is the highest barrier in the reaction pathway, which limits the rate.

Observing the calculated G^{TS} in figure 6 for the Tafel steps, it is clear that the $(\overline{11}20)$ and the $(11\overline{2}1)$ facets should have the highest exchange current density and thus be the most active facets. The reason why these sites have the lowest combination barrier may be found in the geometry of the Ni atom binding the transition state. The transition states for $(\overline{11}20)$ and $(11\overline{2}1)$ are presented in the insets in figure 6. The atomic hydrogen move from stable bridge sites on top of the nickel atom, which on these two facets is coordinated to six nickel atoms and two or three phosphor atoms. It is possible to imagine that a few defect sites, not investigated here, could be participating in the catalysis. However, the most active sites investigated in this study have binding energies very close to the optimal value and they are calculated to be very active. Other sites such as corners or kinks, would only be relevant if they have a much higher activity than the low index facets. Furthermore, we expect that sites with lower coordination will bind too strongly, thus being less active according to the Sabatier principle.

Conclusion

In summary, the DFT calculations have shown that phosphor stabilized Ni-bridge sites found on several available facets of Ni_2P nanowires provide moderate binding to hydrogen atoms. The sites with highest exchange current density estimated from the Tafel transition state energy appear to be the Nickel bridge sites on either the $(\overline{11}20)$ or the $(11\overline{2}1)$ facet. The high activity measured from nanowires grown along the $\langle 0001 \rangle$ direction suggests that the active site of Ni_2P is unlikely the (0001) facet as previously speculated¹⁸. Thus the experiments support the calculations, although contributions from various sites or facets cannot be completely isolated in the electrochemical measurements.

This work shows that Ni_2P has many very active sites for HER, which explains the good performance. An interesting outlook is the calculation of other metal phosphor compounds, to look for sites with slightly stronger hydrogen binding energies or less adsorbate repulsion, which could be promising candidates for new catalyst materials.

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References

- 1 A. Züttel, in *Hydrogen as a Future Energy Carrier*, Wiley-VCH Verlag GmbH & Co. KGaA, 2008, pp. 1–6.
- 2 N. S. Lewis and D. G. Nocera, *Proceedings of the National Academy of Sciences*, 2006, **103**, 15729–15735.

- 3 M. G. Walter, E. L. Warren, J. R. McKone, S. W. Boettcher, Q. Mi, E. A. Santori and N. S. Lewis, *Chemical reviews*, 2010, **110**, 6446–6473.
- 4 S. Trasatti, *Journal of Electroanalytical Chemistry and Interfacial Electrochemistry*, 1972, **39**, 163 – 184.
- 5 J. X. Wang, T. E. Springer and R. R. Adzic, *Journal of the Electrochemical Society*, 2006, **153**, A1732–A1740.
- 6 K. Kunimatsu, T. Senzaki, M. Tsushima and M. Osawa, *Chemical Physics Letters*, 2005, **401**, 451 – 454.
- 7 E. Skulason, G. S. Karlberg, J. Rossmeisl, T. Bligaard, J. Greeley, H. Jonsson and J. K. Nørskov, *Phys. Chem. Chem. Phys.*, 2007, **9**, 3241–3250.
- 8 B. Hinnemann, P. G. Moses, J. Bonde, K. P. Jørgensen, J. H. Nielsen, S. Horch, I. Chorkendorff and J. K. Nørskov, *Journal of the American Chemical Society*, 2005, **127**, 5308–5309.
- 9 T. F. Jaramillo, K. P. Jørgensen, J. Bonde, J. H. Nielsen, S. Horch and I. Chorkendorff, *Science*, 2007, **317**, 100–102.
- 10 J. Kibsgaard, Z. Chen, B. N. Reinecke and T. F. Jaramillo, *Nat Mater*, 2012, **11**, 963–969.
- 11 J. R. McKone, B. F. Sadtler, C. A. Werlang, N. S. Lewis and H. B. Gray, *ACS Catalysis*, 2013, **3**, 166–169.
- 12 H. Vrubel and X. Hu, *Angewandte Chemie International Edition*, 2012, **51**, 12703–12706.
- 13 C. G. Morales-Guio, L.-A. Stern and X. Hu, *Chemical Society Reviews*, 2014, **43**, 6555–6569.
- 14 P. Liu, J. A. Rodriguez, T. Asakura, J. Gomes and K. Nakamura, *The Journal of Physical Chemistry B*, 2005, **109**, 4575–4583.
- 15 G. H. Layan Savithra, E. Muthuswamy, R. H. Bowker, B. A. Carrillo, M. E. Bussell and S. L. Brock, *Chemistry of Materials*, 2013, **25**, 825–833.
- 16 S. T. Oyama, T. Gott, H. Zhao and Y.-K. Lee, *Catalysis Today*, 2009, **143**, 94 – 107.
- 17 P. Liu, J. A. Rodriguez, Y. Takahashi and K. Nakamura, *Journal of Catalysis*, 2009, **262**, 294 – 303.
- 18 P. Liu and J. A. Rodriguez, *Journal of the American Chemical Society*, 2005, **127**, 14871–14878.
- 19 E. J. Popczun, J. R. McKone, C. G. Read, A. J. Biacchi, A. M. Wiltrout, N. S. Lewis and R. E. Schaak, *Journal of the American Chemical Society*, 2013, **135**, 9267–9270.

- 20 L. Feng, H. Vrubel, M. Bensimon and X. Hu, *Physical Chemistry Chemical Physics*, 2014, **16**, 5917-5921.
- 21 J. K. Nørskov, T. Bligaard, A. Logadottir, J. Kitchin, J. Chen, S. Pandalov and U. Stimming, *Journal of The Electrochemical Society*, 2005, **152**, J23–J26.
- 22 M. E. Bjorketun, A. S. Bondarenko, B. L. Abrams, I. Chorkendorff and J. Rossmeisl, *Phys. Chem. Chem. Phys.*, 2010, **12**, 10536–10541.
- 23 J. O. Bockris, A. K. Reddy and M. Electrochemistry, *Electrodics in Chemistry, Engineering, Biology, and Environmental Science*, Springer, 2000.
- 24 J. Greeley, J. K. Nørskov, L. A. Kibler, A. M. El-Aziz and D. M. Kolb, *ChemPhysChem*, 2006, **7**, 1032–1035.
- 25 R. Parsons, *Trans. Faraday Soc.*, 1958, **54**, 1053–1063.
- 26 G. W. Watson, R. P. K. Wells, D. J. Willock and G. J. Hutchings, *The Journal of Physical Chemistry B*, 2001, **105**, 4889–4894.
- 27 J. Rossmeisl, J. K. Nørskov, C. D. Taylor, M. J. Janik and M. Neurock, *The Journal of Physical Chemistry B*, 2006, **110**, 21833–21839.
- 28 Y. Pan, Y. Liu, J. Zhao, K. Yang, J. Liang, D. Liu, W. Hu, D. Liu, Y. Liu and C. Liu, *Journal of Materials Chemistry A*, 2015, **3**, 1656–1665.
- 29 A. Han, S. Jin, H. Chen, H. Ji, Z. Sun and P. Du, *J. Materials Chemistry A*, 2015, **3**, 1941-1946.
- 30 J. J. Mortensen, L. B. Hansen and K. W. Jacobsen, *Phys. Rev. B*, 2005, **71**, 035109.
- 31 B. Hammer, L. B. Hansen and J. K. Nørskov, *Phys. Rev. B*, 1999, **59**, 7413–7421.
- 32 E. Larsson, *Arkiv foer Kemi*, 1965, **23**, 335–365.
- 33 A. Jain, G. Hautier, C. J. Moore, S. P. Ong, C. C. Fischer, T. Mueller, K. A. Persson and G. Ceder, *Computational Materials Science*, 2011, **50**, 2295–2310.
- 34 S. R. Bahn and K. W. Jacobsen, *Comput. Sci. Eng.*, 2002, **4**, 56–66.
- 35 A. B. Alchagirov, J. P. Perdew, J. C. Boettger, R. C. Albers and C. Fiolhais, *Phys. Rev. B*, 2003, **67**, 026103.
- 36 H. J. Monkhorst and J. D. Pack, *Physical Review B*, 1976, **13**, 5188–5192.
- 37 D. Kanama, S. T. Oyama, S. Otani and D. F. Cox, *Surface Science*, 2004, **552**, 8 – 16.

- 38 M. G. Moula, S. Suzuki, W.-J. Chun, S. Otani, S. T. Oyama and K. Asakura, *Surface and Interface Analysis*, 2006, **38**, 1611–1614.
- 39 Q. Li and X. Hu, *Phys. Rev. B*, 2006, **74**, 035414.
- 40 D. Guo, Y. Nakagawa, H. Ariga, S. Suzuki, K. Kinoshita, T. Miyamoto, S. Takakusagi, K. Asakura, S. Otani and S. T. Oyama, *Surface Science*, 2010, **604**, 1347 – 1352.
- 41 J. K. Nørskov, J. Rossmeisl, A. Logadottir, L. Lindqvist, J. R. Kitchin, T. Bligaard and H. Jónsson, *The Journal of Physical Chemistry B*, 2004, **108**, 17886–17892.
- 42 G. Henkelman and H. Jónsson, *The Journal of Chemical Physics*, 2000, **113**, 9978.
- 43 Y. Chen, H. She, X. Luo, G.-H. Yue and D.-L. Peng, *Journal of Crystal Growth*, 2009, **311**, 1229 – 1233.
- 44 L.-A. Stern and X. Hu, *Faraday Discuss.*, 2014, doi: 10.1039/C4FD00120F -.
- 45 Z. Huang, Z. Chen, Z. Chen, C. Lv, H. Meng and C. Zhang, *ACS Nano*, 2014, **8**, 8121–8129.
- 46 J. Rossmeisl, K. Chan, R. Ahmed, V. Tripković and Mårten E. Björketun, *Physical Chemistry Chemical Physics*, 2013, **15**, 10321–10325.