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DFT-based genetic algorithm search for AuCu nanoalloy electrocatalysts for CO₂ reduction

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Using a DFT-based genetic algorithm (GA) approach, we have determined the most stable structure and stoichiometry for a 309 atom icosahedral AuCu nanoalloy, for potential use as an electrocatalyst for CO₂ reduction. The identified core-shell nano-particle consists of a copper core interspersed with gold atoms having only copper neighbors and a gold surface with a few copper atoms in the terraces. We also present an adsorbate-dependent correction scheme, which enables accurate determination of adsorption energies using a computationally fast, localized LCAO-basis set. These show that it is possible to use the LCAO mode to obtain a realistic estimate of the molecular chemisorption energy for systems where the computation in normal grid mode is not computationally feasible. These corrections are employed when calculating adsorption energies on the Cu, Au and most stable mixed particles. This shows that the mixed Cu₁₃₅@Au₁₇₄ core-shell nanoalloy has a similar adsorption energy, for the most favorable site, to a pure gold nano-particle. Cu, however, has the effect of stabilizing the icosahedral structure that for Au particles is easily distorted when adding adsorbates.

1 Introduction

Efficient electroreduction of CO₂ to fuels or chemicals is a key challenge in artificial photosynthesis from renewable energy, which has received considerable attention recently.¹ Copper and gold are among the most interesting materials for CO₂ reduction, because copper reduces CO₂ to e.g. CO and hydrocarbons, whereas gold reduces CO₂ to CO at a comparatively low overpotential. Alloying the two metals could be a way to break the linear scaling between the binding energies of CO and the precursor COOH.² This is for example seen in CO/CO₂ producing CODH enzymes.³ electrode nano structure have been demonstrated to be important for electrode activity and selectivity, 4,5 variations in the selectivity for CO2 reduction on Au and Cu nano-particles (NP) with particle size have been associated with changes in the surface density of low-coordinated sites.⁶ Au₃Cu bimetallic NP shows improved activity for CO evolution over Au NP.9

The theoretical prediction of structure and composition of gold-copper NP using genetic algorithms (GA) is not a novel idea. However, the size of the particles under investigation has often been smaller than the 2 nm particles considered here, where Au is known to display unique catalytic activity, e.g. for CO oxidation. ^{10,11}

The many-body Gupta potential has been the primary tool for structure prediction ^{12–18} in combination with GAs, basin hopping, ^{19,20} molecular dynamics ¹² or other methods for global minima discovery. The accuracy of Gupta and other empirical potentials on alloy clusters (especially gold rich) can be questioned, since the higher order contributions to the energy are non negligible. ²¹ The need for a computationally

inexpensive potential comes from the fact that the number of homotops (distribution of A and B atoms in an A_aB_b cluster)²² rises combinatorially with particle size for every single composition.

Wilson and Johnston¹⁴ conducted a study of icosahedral Au-Cu particles using the Gupta potential and found Cu in the core and Au in the surface lead to the most stable particles. Darby et al. 13 found that even a single copper atom put in the center could change the lowest energy gold amorphous structure into a high-symmetrical icosahedron. Cheng et al. found Au in the core surrounded by Cu in Au₄₃Cu₁₂ clusters using a tight-binding approach.²³ Metastable particles with Au in the core and Cu in the shell have also been observed both theoretically by simulating particle growth using molecular dynamics²⁴ and experimentally with transmission electron microscopy methods. 25,26 However, above a certain temperature, Au will segregate to the surface due to a lower surface energy. For AuCu clusters DFT studies have so far been limited to smaller particles and has mostly been used for in depth studies of recurring structural motifs.^{27,28}

The trend is thus that AuCu NP are observed in high-symmetrical structures due to Cu that is mostly located in the core with Au, having a lower surface energy, being mainly in the shell. This is in good agreement with experimental observations^{29,30}

2 Methods

2.1 Calculation details

We employ calculations within two different levels of accuracy. The initial GA search for stable stoichiometry and composition of the AuCu icosahedral nanoparticle was performed with an semi-empirical potential based on the Effective Medium Theory (EMT).³¹ The subsequent test of

the obtained particles and calculation of adsorption energies were performed within density functional theory^{32,33} using the GPAW code, 34,35 a real space grid based implementation of the projector-augmented wave (PAW) The electronic wave functions were expanded using a linear combination of atomic orbitals (LCAO),³⁷ GPAW also now has the feature of expanding the wave functions in plane waves. The LCAO mode is faster and less memory intensive than the standard finite difference (FD) mode in codes like GPAW at the cost of energetic accuracy; since we employ an incomplete basis set, we cannot expect the adsorption energies to be consistent with the results from FD grid calculations. We have therefore performed rigorous test of adsorption energies with the LCAO mode and compared with previously published results³⁸ of adsorption energies of all reaction intermediate adsorbates on the (111) and (211) transition metal surfaces and M13 clusters using GPAW in grid mode; these are presented in Section 4. The calculation of adsorption energies on all sites would not have been possible using standard DFT, i.e. the speed-up of CPU hours is on the order of 10^2 - 10^3 compared with a similar adsorption energy calculations performed on Pt₃₀₉. ^{39,40} When reporting adsorption energies calculated with LCAO, we have included the counterpoise correction⁴¹ to avoid the basis set superposition error (BSSE). All DFT calculations were performed with the RPBE exchange-correlation functional⁴² the double zeta polarized basis set and an electronic Fermi smearing of 0.1 eV. All calculations were done spin-paired except reference calculations involving Ni. The 309 icosahedral particles were put in a cube with a side length of 32Å (corresponding to approximately 7Å of vacuum on each side) and a grid sampling of 184 points in each direction. The M13 particles were put in a box with 7Å of vacuum on each side with a 96 point grid sampling in each direction. The (111) and (211) extended surfaces were sampled using a Monkhorst-Pack grid⁴³ with k-points (6,7,1) and (5,6,1) respectively (1in the direction normal to the surface) and a grid spacing of 0.18Å was employed. The presented structures and energies have all undergone geometrical optimization using the BFGS algorithm with a line search mechanism until the force on all individual atoms were less than 0.05 eV/Å.

In addition to the GA, we also perform a screening of all symmetric AuCu icosahedral particles, i.e. particles where all symmetrically equivalent sites under the point group symmetry, also called atom subshells, 44,45 are occupied by the same element. All atoms in the same subshell will have the same distance to the center. In a 309 atom icosahedral particle, there are 11 symmetrically in-equivalent sites, for a bimetallic system this leads to 2¹¹ different particles, that can easily be screened with the EMT potential. The symmetrical NP is represented by a string with 11 elements signifying the occupation of each subshell; e.g. the string {\frac{1}{4}Au^{2}Cu^{3}Cu^{4}Cu^{5}Cu^{6}Cu^{7}Cu^{8}Cu^{9}Cu^{10}Cu^{11}Au}} represents the NP with an Au atom at the center, Au atoms at the corner sites and the rest is Cu. This approach is well known for optimizing the distribution of elements in nanoparticles. \(^{14,45,46}

2.2 Genetic algorithm setup

A GA works by adapting a population of possible solutions to

a problem defined using a fitness function. In this case, the challenge was to find the most stable stoichiometry and internal distribution of copper and gold atoms in an icosahedral nanoalloy containing in total 309 atoms. The fitness metric should not simply be the total energy or energy per atom since one would then not be able to compare different stoichiometries. Instead, we use the mixing energy²² defined as:¹⁴

$$E_{mix} = (E^{AB} - E^{A}N^{A}/N - E^{B}N^{B}/N) / N$$
 (1)

where E^{AB} is the total energy of the mixed cluster, N^A and N^B are the number of A and B atoms respectively in the mixed cluster, E^A and E^B are the energies of the pure icosahedral 309 atom clusters A and B, respectively and N is the total number of atoms in the cluster. A negative mixing energy corresponds to a stable mixed particle. This quantity is sometimes also referred to as excess energy.⁴⁷

We employ the GA implemented in the Atomic Simulation Environment⁴⁸ (ASE). Operators previously developed for optimizing the internal distribution of atoms in clusters⁴⁹ have also been employed here, with the removal of the condition that maintains stoichiometry during operations that could change the stoichiometry, e.g. the cut-splice crossover. Note that the core-shell crossover operation has not been used for this work. Two additional substitution operators have been implemented. The first operator substitutes one random atom in the cluster with a different element, directly changing the stoichiometry. In this way all stoichiometries can be examined. Thus we optimize both internal distribution and stoichiometry at the same time. One could suspect that a particle with the composition optimized would never benefit from having a random atom substituted to a different element. However, since the GA maintains a population of particles in different stoichiometries and crosses them, it is possible to approach the optimum of both parameters at the same time. The second additional operator is inspired by the screening of all symmetric AuCu particles. This operator changes all atoms in a subshell to one element type; this effectively incorporates the screening indirectly in the GA.

To sum up the operators used in these GA runs are: Random permutation; permutes two atoms of a different type, center of mass to surface permutation; permutes an atom in the core region with one in the surface, rich to poor (and poor to rich) permutation; permutes two different atoms each from environments rich (poor) in their own type of atoms to environments poor (rich) in their own kind, 49 substitution; see above, symmetric subshell substitution; see above, cut-splice crossover; cuts two particles in half through the center and joins two parts from different parents together.⁵⁰ After each operation the particles undergoes geometrical relaxation. We stress that by using this GA setup, the focus is on the chemical ordering of the elements rather than structural optimization. Changes in structure can only occur during the geometrical relaxation that follows each operation. It is, however, unexpected since moving away from the icosahedral structure requires overcoming an energy barrier. This barrier can be lowered in the case of non-isolated particles.

The GA is initiated with a population of 10 particles in vacuum, each with random Au/Cu stoichiometry. Inclusion of a support material would have been highly relevant, but it has not been taken into account in this study since it would drastically increase both computational time and complexity of particle structure. During each generation, 10 new particles are created using the aforementioned operators, geometrically relaxed and added to the population. The particles are sorted according to the mixing energy (1) and the 10 fittest NP form the population used to create the next generation. When no new NP enters the population for 5 generations, it is assumed that convergence has been reached and the algorithm stops. The algorithm is thus run as a traditional generational GA, where all calculations in a generation must finish before the population is updated and the next generation can commence. It is however also possible to run the algorithm as a pool GA, where the population is continuously updated every time a relaxation is finished and a new structure is created. 51,52

3 Results and discussion

3.1 Finding the optimum

In Fig. 1, the mixing energy as a function of Au/Cu ratio in the particles is shown. The lowest mixing energies as a function of Au/Cu ratio are connected by a red line signifying that no symmetrical particle has a lower mixing energy. In the subshell nomenclature the {\begin{aligned} \cdot Cu \begin{aligned} \cdot Cu \begin{aligned} \cdot Cu \begin{aligned} \cdot Au \begin{aligned} \cd ⁹Au ¹⁰Au ¹¹Au} particle has the lowest mixing energy of the structures in the screening. Two factors are deciding the distribution of atoms in the bimetallic icosahedral NP: 1) Relieve the inherent bulk strain⁵³ by putting the smaller element (Cu) in the core and the larger (Au) on the surface. 2) Maximize the number of stronger bonds; from the cohesive energies one can deduce individual bond strengths¹⁴ Au-Au > Au-Cu > Cu-Cu. Here, it means that Au atoms are dispersed in the core, in subshell 4 and 5, with each Au atom having only Cu neighbors – the gain in bond energy is larger than the loss due to strain. The opposite is true for subshell 1, where a gold atom would also have 12 Cu neighbors; however, here the increase in strain energy is greater than the gain in bond energy. This is in contrast to the findings by Cheng et al., 23 who find an Au atom in the core even for stoichiometries with Cu on the surface albeit with the smaller 55 atom particle. We performed a number of GA runs and in all of them a small

number, up to eight, Cu atoms were situated in the terrace sites at the surface. The only difference between the GA particles and the ones from the symmetrical screening is the surface terrace sites (the eighth subshell). Below the surface, the particles are identical to the best from the symmetrical screening. By introducing non-symmetric regions with Cu in the surface, the GA is able to predict particles with lower mixing energy, as is evident in figures 1b and c. We subsequently did DFT calculations on the 50 most stable particles from the screening and the 50 fittest from the GA runs, these are shown in Fig. 1c. Afterwards we tested the optimal amount of Cu in the terrace sites and discovered that EMT predicts 8 and DFT 24 (squares in Fig. 1c). The Cu atoms should be spread out as much as possible with 16 terraces only having one Cu atom and four having two. Fig. 2 is a representation of the most stable particle. The minimum is thus at an Au/Cu ratio of 55/45, which is slightly different from previous work in the literature¹⁴ that found the lowest mixing energy for a 50/50 mixture of Cu and Au.

For calculations of the adsorption energy, we use a particle with one Cu atom in each terrace – that is 20 Cu atoms in the surface in total. This leads to a $Au_{174}Cu_{135}$ particle; a pure Cu core with Au placed in atom subshell 4 and 5 (see Fig. 2), thereby the gold atoms only have copper neighbors in the core. The surface is all gold except for one Cu atom on a terrace site in each terrace.

Pure Cu particles exhibit the icosahedral structure in all sizes, whereas pure Au clusters do not exhibit highly symmetrical ground states but instead more amorphous-like low-symmetry structures. Au clusters are known to fluctuate in structure meaning that there are many local minima close in energy with both low and high symmetry. We found that adding adsorbates can distort metastable icosahedral Au particles into a more energetically favorable amorphous structure. A small amount of Cu can however stabilize high-symmetry structures, even when adding adsorbates. Substituting the two innermost Au subshells with Cu is enough to stabilize the icosahedral 309 Au atom structure against distortion, even though the icosahedral structure is one of the least stable Au structures.

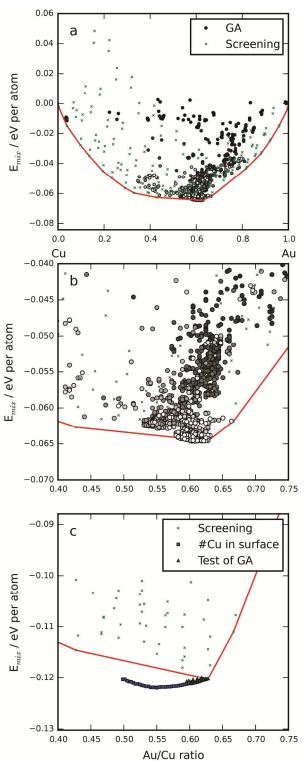


Fig. 1 Mixing energy as a function of Au/Cu ratio in the 309 atom particles. a shows the lowest energy for each possible value of Au/Cu ratio from the screening of 2¹¹ symmetric particles marked with x's. Also shown is all the particles from a single GA run marked by circles. All are calculated with EMT. The circles are shaded corresponding to how far in the GA run the particles were investigated, black is first generation, white is last generation. The red line makes up the convex hull, it connects the lowest points of the screening marking the ground state line as expected from the screening. b shows a zoom of the minimum region of a. c is the minimum region now calculated with DFT.

The findings here underline that it is advantageous to employ calculations at two levels of complexity. The lighter one for fast screening and the heavier one for correct ordering of the best screening results. 57-60

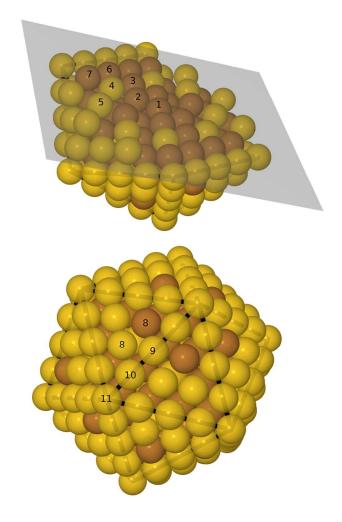


Fig. 2 Top: Cut-through of the optimal particle, the numbers on the atoms indicate the subshell number. Bottom: The whole $Au_{174}Cu_{135}$ particle. Brown: Cu, yellow: Au.

3.2 Accuracy of LCAO adsorption energies

In order to be able to validate the adsorption energies, we have tested the LCAO mode against values previously reported in the literature. Peterson et al.³⁸ have reported the adsorption energies of O and CO on (111) and (211) surfaces as well as M13 clusters using FD for all stable adsorption sites. We find that the error from using the faster LCAO mode depends mostly on the adsorbate. In Fig. 3, we plot the adsorption energy of CO calculated with LCAO vs. literature data on Ni, Cu, Ag, Pd, Au and Pt. It is evident that the LCAO mode generally yields too low binding energies, but the underestimation is systematic for all elements, except Pd M13, with an acceptable standard deviation of 0.12 eV. For

the further discussions we have subtracted the mean error of 0.28 eV for all calculated CO adsorption energies.

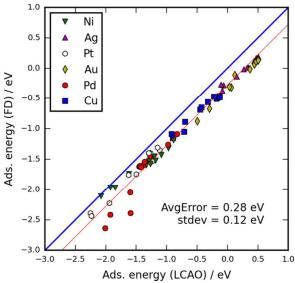


Fig. 3 Adsorption energies of CO with the FD mode vs. LCAO mode in GPAW. The FD mode values are from ref 38. The thick blue line is the no correction line. The thin red line includes the subtraction of the average error of 0.28 eV. The four lowest Pd points (the outliers) are from M13 calculations.

For COOH, we do not have GPAW FD data; instead we used data produced with the planewave code Dacapo (shown in Fig. S1). The correction is very small and we have neglected it in the further discussion.

3.3 Adsorption energies

We calculate adsorption energies of COOH and CO intermediates at the apex site, edge sites and terrace sites of the Au, Cu and AuCu nanoalloys. The most stable adsorption geometries are displayed in Fig 4.

On the Au and Au $_{174}$ Cu $_{135}$ NP, CO and COOH adsorbs strongest on the apex sites, whereas CO and COOH adsorbs stronger on edge sites than apex sites on Cu. On the terraces of the Cu $_{135}$ @Au $_{174}$ core-shell nanoalloy, we find that CO binds to the Cu atom, while COOH binds through the Au atom. Because CO and COOH binds significantly stronger to Cu than Au, it is surprising that COOH binds to Au on the terrace rather than Cu. We observe no indications of the formation of oxygen bonds to the Cu atoms adjacent to the Au-C primary bond, as proposed by Kim et al. 9

It has been suggested the production of CO on Au and Cu goes through a COOH intermediate. ⁶¹⁻⁶³ Fig. **5** show the CO production at 0.35 V overpotential from a kinetic model developed previously. ³ In this model, the production of CO follows the reaction mechanism

$$CO_{2(aq)} + * + H^+ + e^- \rightarrow COOH^*$$
 (R1)

$$COOH^* + H^+ + e^- \rightarrow CO^* + H_2O_{(1)}$$
 (R2)

$$CO^* \to CO_{(aq)} + * \tag{R3}$$

The prefactors for coupled proton-electron transfer have been fitted to experiments on Au. 61 Assuming activation energies to scale with reaction energies; it is possible to describe the activity for CO evolution as function of the COOH and CO binding energies. In Fig. 5 the activity is described as a function of the reaction energy of (R1) and (R2) as calculated from DFT and including 0.25 eV and 0.1 eV stabilizations of COOH and CO* intermediates, due to hydrogen bonding from the solvent. 63

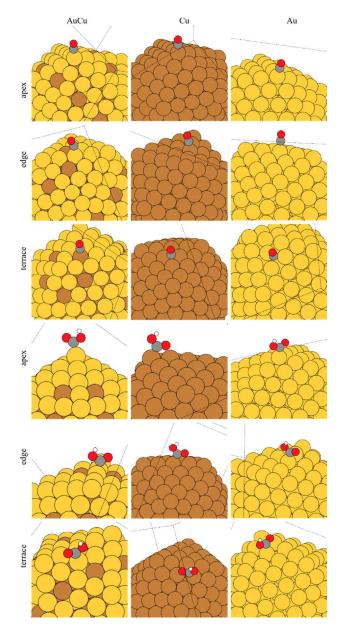


Fig. 4 Geometries of CO and COOH adsorbed at apex, edge and terrace sites on AuCu, Cu and Au nano-particles. Au is yellow, Cu brown, C,

gray, O red, and H is white.

surfaces respectively of the late transition metals.

Within the kinetic model, a good catalyst for CO production should bind COOH sufficiently strong to activate CO₂ and it should not bind CO so strong that the active sites are blocked by CO.³ Activities of (111) terraces and (211) steps on Ag, Au and Cu are indicated based on previous DFT calculations.^{2,64}

The edge sites on the Au NP are predicted to have good activity for CO production, comparable to the activity of Ag(211), but lower than Au(211). The apex sites on the Au₁₇₄Cu₁₃₅ and Au NP have slightly lower activity for CO evolution at the simulated conditions and are comparable to Cu(211). The CO production on these sites is limited by CO poisoning as well as CO₂ activation, so we expect that mass transport conditions are important for CO production, similar to what has been observed on Cu.65 The terrace sites on Au₁₇₄Cu₁₃₅ and Au NP are predicted to be too inactive to activate CO₂. The active sites on the Cu NP are found to bind too strongly for efficient CO production. It should be noted, however, that adsorbate-adsorbate interactions are not included in the model and repulsion from CO* is expected to weaken binding energies resulting in higher activity.⁶⁴ should further be noted, that the model does not include H₂ production, so the selectivity for CO2 reduction is not included by the model.^{8,66} A potentially critical inadequacy if accurate predictions about product selectivity are to be made, however the work is on-going for including H2 production.

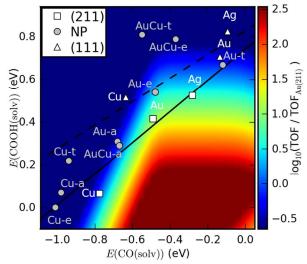


Fig 5 Contour plot of the turn over frequency (TOF) for CO₂ to CO from a kinetic model at 0.35 V overpotential.³ The predicted activity of apex, edge and terrace sites on Au₁₇₄Cu₁₃₅, Au and Cu nanoparticles from our study is shown. Activities of (111) and (211) facets are taken from refs
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4 Conclusions

We have studied catalytically interesting 309 atom AuCu icosahedral nano-particles. In order to determine the most stable stoichiometry and distribution of atoms, we have implemented GA operators that incorporate a screening of cluster symmetric subshells into a GA run. This greatly enhances the efficiency of the GA, since it now has the ability to combine highly symmetrical atomic distributions with slight but important deviations from the symmetry. This feature is present in the most stable mixed AuCu nanoalloy presented here, see Fig. 2. The most stable particle is a Cu₁₃₅@Au₁₇₄ core-shell nanoalloy, it is formed by a Cu core with interspersed Au that all have only Cu neighbors in the core and a surface of Au with a few Cu atoms on the terraces. During the GA runs, the energies were calculated using the EMT potential. Subsequently, the fittest particles were evaluated with DFT to obtain the correct ordering. This showed that almost half the terrace sites should be occupied by Cu (24 out of 60). It is a novel finding that a theoretical method can predict stable nanoalloys with scattered Cu and Au atoms in the shell and core respectively.

Adsorption energies of CO and COOH was determined on the pure Au, Cu and mixed Au₁₇₄Cu₁₃₅ particles. This did not place any of the particles closer to the top of the activity volcano in Fig. 5. We tested the LCAO adsorption energy credibility by comparing with previously published values from the literature; the systematic errors were relatively small probably due to the systematic removal of BSSE.

The fact that the $Au_{174}Cu_{135}$ nanoalloy has approximately the same properties as Au_{309} regarding adsorption of molecular species on the most energetically favorable site is interesting from an economical point of view; it could be possible to reduce the amount of precious metal in a catalyst since the active metals would only be situated on the surface. This has previously also been observed for Pt-Ni particles.⁶⁷

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6 Notes and references

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